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Erhard G. Schultchen

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USER'S MANUAL FOR AXCST
FINITE ELEMENT PROGRAM

by

Erhard G. Schultchen

Celal N. Kostem

Department of Civil Engineering
Fritz Engineering Laboratory
Lehigh University
Bethlehem, Pennsylvania

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1. INTRODUCTION

AXCST is a program for the stress and displacement analysis of axi-symmetric elasticity problems using the Finite Element method. In the program triangular elements with a linear displacement field are used.

The major blocks of this program are similar to those of the CSTES program (Ref. 1). The AXCST program can be used for the linear elastic analysis of axi-symmetric solids with axi-symmetric boundary conditions of any type. In addition, the program can solve thermal stress problems with arbitrary axi-symmetric temperature distribution.

2. ANALYTIC FORMULATION

2.1 Displacement Field

Stresses and displacements of a solid with axi-symmetric geometry, boundary conditions, and loading are also axi-symmetric. If the problem is formulated in cylindrical coordinates r , θ , and z , this implies that stresses and displacements are independent of θ . Thus, the three-dimensional elasticity problem will be reduced to a two-dimensional case, where the coordinates r and z completely define the stresses and displacements of the body in all directions.

The Finite Element approach for solving elasticity problems requires the definition of a displacement field within the element. The simplest form of this function is based on a complete first order polynomial in the following form:

$$u(r,z) = e_1 + e_2 r + e_3 z \quad (1)$$

$$v(r,z) = e_4 + e_5 r + e_6 z$$

This assumption immediately defines the strain variation within each element, which is found to be constant except for ϵ_θ . Since the tangential strain varies as an inverse function of r , the normal stresses throughout the element will vary in radial direction as well.

The linear displacement function defined in Eq. (1) satisfies the compatibility conditions along the edges of the elements as long as nodal point continuity is maintained.

Expressing the unknown coefficients e_1, \dots, e_6 in terms of nodal point displacements, the displacement function can be rewritten:

$$\begin{bmatrix} u(r,z) \\ v(r,z) \end{bmatrix} = \frac{1}{2\Delta} \begin{bmatrix} 1 \\ r \\ z \end{bmatrix}^T \begin{bmatrix} a_i & a_j & a_k \\ b_i & b_j & b_k \\ c_i & c_j & c_k \end{bmatrix} \begin{bmatrix} u_i & v_i \\ u_j & v_j \\ u_k & v_k \end{bmatrix} \quad (2)$$

where

$$\begin{aligned} a_i &= r_j z_k - r_k z_j & b_i &= z_j - z_k & c_i &= r_k - r_j \\ a_j &= r_k z_i - r_i z_k & b_j &= z_k - z_i & c_j &= r_i - r_k \\ a_k &= r_i z_j - r_j z_i & b_k &= z_i - z_j & c_k &= r_j - r_i \end{aligned} \quad (3)$$

$$2\Delta = b_j c_k - b_k c_j \quad |\Delta| = \text{area of triangle.}$$

It should be noted that the axi-symmetric CST-elements are actually complete rings with triangular cross-section, while the nodal "points" at which they are connected are circular lines in plane view (Fig. 1).

2.2 Fundamental Matrices

In the Finite Element formulation of elasticity problems four fundamental matrices are used:

(1) B - matrix,

relates the elemental strains to the nodal point displacements.

(2) D - matrix,

relates the elemental stresses to the elemental strains, known as the elasticity matrix.

(3) k^e - matrix,

relates the nodal point reactions to the nodal point displacements of each element, known as the element stiffness matrix.

(4) K - matrix,

relates the nodal point reactions to the nodal point displacements of the whole structure, known as the global stiffness matrix.

2.3 B - Matrix

From the theory of elasticity the following relationships can be found for the axi-symmetric state of strain (Ref. 2):

$$\begin{aligned}
 \epsilon_r &= \frac{\partial u}{\partial r} & \gamma_{rz} &= \frac{\partial u}{\partial z} + \frac{\partial v}{\partial r} \\
 \epsilon_z &= \frac{\partial v}{\partial z} & \gamma_{\theta r} &= 0 \\
 \epsilon_\theta &= \frac{u}{r} & \gamma_{\theta z} &= 0
 \end{aligned}
 \tag{4}$$

Substituting the displacement function (Eq. (2)) into the strain-displacement relations (Eq. (4)), the elemental strains can be expressed in terms of nodal point displacements:

$$\underline{\epsilon} = \underline{B} \underline{\delta}^e$$

$$\begin{bmatrix} \epsilon_r \\ \epsilon_z \\ \epsilon_\theta \\ \gamma_{rz} \end{bmatrix} = \frac{1}{2\Delta} \begin{bmatrix} b_i & - & b_j & - & b_k & - \\ - & c_i & - & c_j & - & c_k \\ p_i & - & p_j & - & p_k & - \\ c_i & b_i & c_j & b_j & c_k & b_k \end{bmatrix} \begin{bmatrix} u_i \\ v_i \\ u_j \\ v_j \\ u_k \\ v_k \end{bmatrix} \quad (5)$$

where

$$p_i = \frac{a_i}{r} + b_i + \frac{c_i z}{r}$$

$$p_j = \frac{a_j}{r} + b_j + \frac{c_j z}{r}$$

$$p_k = \frac{a_k}{r} + b_k + \frac{c_k z}{r}$$

It is convenient to split the \underline{B} - matrix into a constant and a variable part:

$$\underline{B} = \underline{B}_1 + \underline{B}_2$$

\underline{B}_1 defines here the element strains at the center of the triangle, while \underline{B}_2 takes into account the variation of ϵ_θ throughout the element.

$$\underline{B}_2 = \frac{1}{2\Delta} \begin{bmatrix} b_i & 0 & b_j & 0 & b_k & 0 \\ 0 & c_i & 0 & c_j & 0 & c_k \\ p_{ci} & 0 & p_{cj} & 0 & p_{ck} & 0 \\ c_i & b_i & c_j & b_j & c_k & b_k \end{bmatrix} \quad (6)$$

$$\underline{B}_2 = \frac{1}{2\Delta} \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ q_i & 0 & q_j & 0 & q_k & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

where

$$p_{cm} = \frac{a_m}{r_c} + b_m + \frac{c_m z_c}{r_c} \quad m = i, j, k$$

$$q_m = \frac{a_m}{r} + \frac{c_m z}{r} - \frac{a_m}{r_c} - \frac{c_m z_c}{r_c}$$

$$r_c = \frac{1}{3} (r_i + r_j + r_k)$$

coordinates of
center of triangle

$$z_c = \frac{1}{3} (z_i + z_j + z_k)$$

The expressions for p_{ci} , p_{cj} , and p_{ck} can be rewritten in the following form:

$$p_{ci} = p_{cj} = p_{ck} = \frac{2\Delta}{3r_c} = p_c$$

2.4 Elasticity Matrix

The elasticity matrix \underline{D} defines the stress-strain relationship within each element. For an isotropic material the \underline{D} matrix is of the following form:

$$\underline{\sigma} = \underline{D} \underline{\epsilon}$$

$$\begin{bmatrix} \sigma_r \\ \sigma_z \\ \sigma_\theta \\ \tau_{rz} \end{bmatrix} = E \frac{(1-\nu)}{(1+\nu)(1-2\nu)} \begin{bmatrix} 1 & \frac{\nu}{1-\nu} & \frac{\nu}{1-\nu} & - \\ \frac{\nu}{1-\nu} & 1 & \frac{\nu}{1-\nu} & - \\ \frac{\nu}{1-\nu} & \frac{\nu}{1-\nu} & 1 & - \\ - & - & - & \frac{1-2\nu}{2(1-\nu)} \end{bmatrix} \begin{bmatrix} \epsilon_r \\ \epsilon_z \\ \epsilon_\theta \\ \gamma_{rz} \end{bmatrix} \quad (7)$$

2.5 Element Stiffness Matrix

A general expression for the element stiffness matrix, which relates nodal point reactions and nodal point displacements of each element, can be derived by applying the principle of virtual work (Ref. 3).

$$\underline{k}^e = \int_V \underline{B}^T \underline{D} \underline{B} dV \quad (8)$$

This integration has to be carried out over the total volume of the ring-element, however, due to axi-symmetry this volume integral can be simplified to an integral over the triangular cross section of the element.

$$\underline{k}^e = 2\pi \int_{\Delta} \underline{B}^T \underline{D} \underline{B} r dr dz \quad (9)$$

Separating the constant and variable parts of the \underline{B} - matrix results in the following summation:

$$\begin{aligned} \underline{k}^e = & 2\pi \int_{\Delta} \underline{B}_1^T \underline{D} \underline{B}_1 r dr dz + 2\pi \int_{\Delta} \underline{B}_1^T \underline{D} \underline{B}_2 r dr dz \\ & + 2\pi \int_{\Delta} \underline{B}_2^T \underline{D} \underline{B}_1 r dr dz + 2\pi \int_{\Delta} \underline{B}_2^T \underline{D} \underline{B}_2 r dr dz \end{aligned} \quad (10)$$

Noting that $\int_{\Delta} \underline{B}_2 r dr dz = \int_{\Delta} \underline{B}_2^T r dr dz = 0$ and $\int_{\Delta} r dr dz = r_c |\Delta|$, the element stiffness matrix Δ can be expressed as Δ

$$\underline{k}^e = \underline{k}_1 + \underline{k}_2$$

or

(11)

$$\underline{k}^e = 2\pi r_c |\Delta| \underline{B}_1^T \underline{D} \underline{B}_1 + 2\pi \int_{\Delta} \underline{B}_2^T \underline{D} \underline{B}_2 r dr dz$$

The numerical evaluation of \underline{k}_1 presents no difficulties, however, problems arise in the integration of certain terms of \underline{k}_2 . From Eqs. (6) and (7)

$$\underline{B}_2^T \underline{D} \underline{B}_2 = \frac{E}{4\Delta^2} \frac{(1-\nu)}{(1+\nu)(1-2\nu)} \begin{bmatrix} q_i q_i & - & q_j q_i & - & q_k q_i & - \\ - & - & - & - & - & - \\ q_i q_j & - & q_j q_j & - & q_k q_j & - \\ - & - & - & - & - & - \\ q_i q_k & - & q_j q_k & - & q_k q_k & - \\ - & - & - & - & - & - \end{bmatrix} \quad (12)$$

where

$$q_m q_n = q_n q_m = \left(\frac{a_m}{r} + \frac{c_m z}{r} - \frac{a_m}{r_c} - \frac{c_m z}{r_c} \right) \left(\frac{a_n}{r} + \frac{c_n z}{r} - \frac{a_n}{r_c} - \frac{c_n z}{r_c} \right)$$

$$m = i, j, k$$

$$n = i, j, k$$

Integration of these terms leads to

$$\int_{\Delta} q_m q_n r dr dz = |\Delta| \left[a_m a_n \left(\frac{F_1}{\Delta} - \frac{1}{r_c} \right) + (a_m c_n + a_n c_m) \left(\frac{F_2}{\Delta} - \frac{z_c}{r_c} \right) + c_m c_n \left(\frac{F_3}{\Delta} - \frac{z_c^2}{r_c} \right) \right] \quad (13)$$

In the derivation of this expression the following integrals have to be evaluated:

$$\begin{aligned} \int_{\Delta} dr dz &= |\Delta| & \int_{\Delta} \frac{1}{r} dr dz &= F_1 \\ \int_{\Delta} r dr dz &= |\Delta| r_c & \int_{\Delta} \frac{z}{r} dr dz &= F_2 \\ \int_{\Delta} z dr dz &= |\Delta| z_c & \int_{\Delta} \frac{z^2}{r} dr dz &= F_3 \end{aligned} \quad (14)$$

For the general case $r_i \neq r_j \neq r_k$ and $r_i \neq 0$, $r_j \neq 0$, $r_k \neq 0$ (Fig. 2), the integrals F_1 , F_2 , and F_3 can be written as

$$F_1 = \int_{\Delta} \frac{1}{r} dr dz = I_{11} + I_{12} + I_{13}$$

$$I_{11} = \frac{a_i}{c_i} (\text{Lnr}_k - \text{Lnr}_j) + b_i \quad \text{along } k-j$$

$$I_{12} = \frac{a_j}{c_j} (\text{Lnr}_i - \text{Lnr}_k) + b_j \quad \text{along } i-k$$

$$I_{13} = \frac{a_k}{c_k} (\text{Lnr}_j - \text{Lnr}_i) + b_k \quad \text{along } j-i$$

$$F_2 = \int_{\Delta} \frac{z}{r} dz dr = I_{21} + I_{22} + I_{23}$$

$$I_{21} = -\frac{1}{2} \frac{a_i^2}{c_i^2} (\text{Lnr}_k - \text{Lnr}_j) - \frac{a_i}{c_i} b_i - \frac{1}{4} \frac{b_i}{c_i} b_i (r_k + r_j)$$

$$I_{22} = -\frac{1}{2} \frac{a_j^2}{c_j^2} (\text{Lnr}_i - \text{Lnr}_k) - \frac{a_j}{c_j} b_j - \frac{1}{4} \frac{b_j}{c_j} b_j (r_i + r_k)$$

$$I_{23} = -\frac{1}{2} \frac{a_k^2}{c_k^2} (\text{Lnr}_j - \text{Lnr}_i) - \frac{a_k}{c_k} b_k - \frac{1}{4} \frac{b_k}{c_k} b_k (r_j + r_i)$$

$$F_3 = \int_{\Delta} \frac{z^2}{r} dz dr = I_{31} + I_{32} + I_{33}$$

$$I_{31} = +\frac{1}{3} \frac{a_i^3}{c_i^3} (\text{Lnr}_k - \text{Lnr}_j) + \frac{a_i^2}{c_i^2} b_i + \frac{1}{2} \frac{a_i}{c_i} \frac{b_i}{c_i} b_i (r_k + r_j) \\ + \frac{1}{9} \frac{b_i^3}{c_i^3} (r_k^3 - r_j^3)$$

$$I_{32} = +\frac{1}{3} \frac{a_j^3}{c_j^3} (\text{Lnr}_i - \text{Lnr}_k) + \frac{a_j^2}{c_j^2} b_j + \frac{1}{2} \frac{a_j}{c_j} \frac{b_j}{c_j} b_j (r_i + r_k) \\ + \frac{1}{9} \frac{b_j^3}{c_j^3} (r_i^3 - r_k^3)$$

$$I_{33} = + \frac{1}{3} \frac{a_k^3}{c_j^3} (\text{Lnr}_j - \text{Lnr}_i) + \frac{a_k^2}{c_k^2} b_k + \frac{1}{2} \frac{a_k}{c_k} \frac{b_k}{c_k} b_k (r_j + r_i) \\ + \frac{1}{9} \frac{b_k^3}{c_k^3} (r_j^3 - r_i^3)$$

It can be shown that certain simplifications of these expressions are possible; for instance, in the evaluation of F_1 the sum of $b_i + b_j + b_k$ is indentically zero. Whereas this is true for the general case of $r_m \neq r_n$ and $r_m \neq 0$, algebraic operations of this type are not permissible in certain special cases where terms of I_{pq} become improper. These special cases can be classified into two groups:

$$\text{Type 1: } r_m = r_n$$

$$\text{Type 2: } r_m = 0$$

A Type 1 singularity arises when two nodal points, m and n , have the same value of r (for instance, $r_k = r_j$). This means that the corresponding value of c_i (Eq. (3)) becomes zero. Since c_i appears in the denominator of several terms of I_{pq} , these expressions will become infinite or at least improper. However, from Fig. 2 it can be seen that in this case the sub-integral over the area $k j j' k'$ does not exist, which means that $I_{11} = I_{21} = I_{31} = 0$.

A singularity of Type 2 is encountered in the case of $r_m = 0$. Those terms of I_{pq} which become improper for $r_m = r_i = 0$ can be reduced to the following form:

$$T_{pi} = \left(\frac{a_j^p}{c_j^p} - \frac{a_k^p}{c_j^p} \right) \text{Lnr}_i \quad p = 1, 2, 3 \quad (16)$$

The evaluation of T_{pi} for $r_i = 0$ leads to an indefinite expression of the form $0 \cdot \infty$; however, by applying L'Hospital's Rule it can be shown that the limit of all T_{pm} terms is zero.

Although it is possible that both types of singularities occur simultaneously ($r_m = r_n = 0$), this condition requires no special attention since the conclusions drawn previously are still valid in this case.

The foregoing paragraphs show that the evaluation of the element stiffness matrix for axi-symmetric triangular elements is relatively complicated compared to the CST (constant strain triangle) element in plane stress problems. The reason for this lies in the difficulties encountered in generating the \underline{k}_2 - matrix.

Therefore, a number of developers of Finite Element programs tried to circumvent these difficulties by using certain approximations (Refs. 4,5,6,7).

The simplest way of doing this is by assuming $\underline{k}_2 = 0$ (Ref. 4). Recalling that \underline{k}_2 takes into account the variation of ϵ_θ inside the element, this approximation is certainly justified for large values of r_m . On the other hand, it can be shown that the elements of \underline{k}_2 are of the same magnitude as those of \underline{k}_1 for elements in the vicinity of the axis of rotation. Therefore, the use of this approximation may cause large inaccuracies.

It was shown that the existence of certain special conditions, for which a number of sub-integrals become singular or do not exist,

eliminates the possibility of simplifying the expressions for the integrals F_1, \dots, F_3 .

To circumvent this difficulty it was suggested (Ref. 5) to assign a small quantity to those values that cause the integrals to become singular. Test runs show that the approximation $r_m = \epsilon$ instead of $r_m = 0$ gives results which seem to converge to the true value. On the contrary, the approximation $r_m - r_n = \epsilon$ results in an excessive error build-up even for computers with a relatively long word length. For this reason, it seems to be questionable whether an approximation of this type should be used.

Another very common method of evaluating the element stiffness matrix for axi-symmetric triangular elements is the use of numerical integration techniques (Refs. 6,7). Obviously, this procedure gives only approximate results, regardless of whether those special conditions for the nodal point coordinates are met or not. In general, it appears to be doubtful whether numerical integration has a distinct advantage as far as execution time and storage space is concerned, when compared with the exact evaluation procedure.

2.6 Global Stiffness Matrix

One of the major problems in a Finite Element program is the assemblage of the global stiffness matrix K of the structure. A detailed description of this procedure is given in Ref. 1.

2.7 Auxiliary Matrices

In order to make the analytic formulation of the Finite Element method complete, two additional relationships have to be established.

Since the AX CST program includes the analysis of thermal strain problems, expressions have to be derived to transform the thermal strain loading condition $\underline{\epsilon}^t$ to equivalent nodal point loads \underline{F}^t .

$$\underline{\epsilon}^t = \begin{bmatrix} \epsilon_r^t \\ \epsilon_z^t \\ \epsilon_\theta^t \\ \gamma_{r\theta}^t \end{bmatrix} = \begin{bmatrix} \alpha t \\ \alpha t \\ \alpha t \\ 0 \end{bmatrix} \quad (17)$$

By applying the principle of virtual work to this problem the following relationship can be obtained:

$$\begin{aligned} \underline{F}^t &= \int_v \underline{B}^T \underline{D} \underline{\epsilon}^t dV \\ &= 2\pi \int_{\Delta} (\underline{B}_1^T + \underline{B}_2^T) \underline{D} \underline{\epsilon}^t r dr dz \end{aligned} \quad (18)$$

Recalling that $\int_{\Delta} \underline{B}_2^T r dr dz = 0$, this expression can be simplified:

$$\begin{aligned} \underline{F}^t &= 2\pi \underline{B}_1^T \underline{D} \underline{\epsilon}^t \int r dr dz \\ &= 2\pi \underline{B}_1^T \underline{D} \underline{\epsilon}^t r_c |\Delta| \end{aligned} \quad (19)$$

After evaluating $\underline{B}_1^T \underline{D} \underline{\epsilon}^t$ explicitly the nodal point loads \underline{F}^t finally can be written as

$$\underline{F}^t = \frac{E}{(1-2\nu)} \cdot \alpha \cdot t \pi r_c \cdot \frac{|\Delta|}{\Delta} \begin{bmatrix} b_i + p_c \\ c_i \\ b_j + p_c \\ c_j \\ b_k + p_c \\ c_k \end{bmatrix} \quad (20)$$

As a last step in a Finite Element analysis, the elemental stresses have to be expressed in terms of nodal point displacements. For this purpose the following relationship has to be used:

$$\begin{aligned} \underline{\sigma} &= \underline{D} \underline{B} \underline{\delta}^e - \underline{D} \underline{\epsilon}^t \\ &= \underline{D} (\underline{B} \underline{\delta}^e - \underline{\epsilon}^t) \end{aligned} \quad (21)$$

3. GENERAL PROGRAM DESCRIPTION

3.1 Application to Boundary Value Problems

The AX CST program can be used for solving the following type of elasticity problems:

- (a) Surface Forces, (S)
- (b) Body Surfaces, (F)
- (c) Displacement Boundary Conditions
- (d) Thermal Strain

It is understood, of course, that these loading conditions are axisymmetrical.

As a characteristic of the Finite Element Method, all force and displacement boundary conditions have to be prescribed at the nodal points, as nodal point loads or as nodal point displacements. This makes it possible to prescribe displacements not only at the surface, but also in the interior of the body.

To keep the program as general as possible no attempt was made to incorporate the evaluation of nodal point loads due to distributed surface or body forces. The procedure for obtaining the nodal point loads due to arbitrary distributed forces will be outlined in Section 3.2.

Only in the case of thermal strain problems the program will generate the nodal point loads using data for the coefficient of thermal expansion and the temperature change, which have to be prescribed for every element. Again, more specific features like temperature

interpolation or temperature variation within the element was not included, but can be handled by assigning nodal point loads directly.

3.2 Nodal Point Loads due to Distributed Forces

The evaluation of nodal point loads due to distributed surface forces is commonly done by applying the principle of virtual work (Refs. 5,8).

Using the following notation

S_n = nodal point load due to distributed surface forces

$\bar{\Delta}_n$ = virtual displacement at nodal point n in the direction of S_n

$\underline{f}^T = \begin{bmatrix} f_r(r,z) \\ f_z(r,z) \end{bmatrix}$ = distributed surface forces

$\underline{\Delta}_s = \begin{bmatrix} \bar{u}_s(r,z) \\ \bar{v}_s(r,z) \end{bmatrix}$ = virtual displacement of the surface s

$\int_s ds$ = integral over the element surface

the nodal point load can be written as

$$S_n \cdot \bar{\Delta}_n = \int_s \underline{f}^T \underline{\Delta}_s ds \quad (22)$$

By assuming $\bar{\Delta}_n = 1$ it follows

$$S_n = \int_s \left[f_r \cdot \bar{u}_s + f_z \cdot \bar{v}_s \right] ds \quad (23)$$

Here \bar{u}_s and \bar{v}_s represent a compatible displacement field, having a unit displacement at nodal point n in the direction of S_n .

To illustrate this procedure the nodal point loads due to uniform pressure acting on a plane $z_s = \text{constant}$ will be evaluated (Fig. 3).

From Eq. (23) follows

$$S_{iz} = \int_s \left[f_r \cdot \bar{u}_s + f_z \bar{v}_s \right] ds \quad (24)$$

Since $f_r(r, z) = 0$ and $f_z(r, z) = -p_z$ on surface s , Eq. (25) can be rewritten as

$$S_{iz} = -p_z \int_s \bar{v}_s ds \quad (25)$$

A compatible displacement function \bar{v}_s with $\bar{v}_i = 1$ can be obtained from Eq. (2)

$$\begin{aligned} \bar{v}_s(r, z_s) &= \frac{1}{2\Delta} \begin{bmatrix} 1 \\ r \\ z_s \end{bmatrix}^T \begin{bmatrix} a_i & a_j & a_k \\ b_i & b_j & b_k \\ c_i & c_j & c_k \end{bmatrix} \begin{bmatrix} \bar{v}_i = 1 \\ 0 \\ 0 \end{bmatrix} \\ &= \frac{1}{2\Delta} (a_i + b_i r + c_i z_s) \end{aligned} \quad (26)$$

Since the surface integral $\int_s ds$ is independent of z in this particular case, the following simplification can be made:

$$\int_s ds = 2\pi \int r dr \quad (27)$$

Using Eqs. (25), (26), and (27) the nodal point load S_{iz} can be written as

$$S_{iz} = -\frac{2\pi p_z}{2\Delta} \int_{r_i}^{r_k} [(a_i + c_i z_s) r + b_i r^2] dr \quad (28)$$

After executing this integration and replacing a_i , b_i , c_i , and 2Δ , according to Eq. (3), the following expression for S_{iz} will be obtained:

$$S_{iz} = -\frac{\pi p_z}{3} (r_k - r_i) (r_k + 2r_i) \quad (29a)$$

Using the same procedure to evaluate S_{jz} and S_{kz} yields

$$S_{kz} = -\frac{\pi p_z}{3} (r_k - r_i) (r_i + 2r_k) \quad (29b)$$

$$S_{jz} = 0 \quad (29c)$$

Another very common type of loading, uniform pressure on a surface with $r_s = \text{constant}$ (Fig. 4), will result in the following expressions for the nodal point loads:

$$S_{ir} = S_{jr} = -\pi p_r r_s (z_j - z_i) \quad (30a,b)$$

$$S_{kr} = 0 \quad (30c)$$

The generality of this procedure makes it possible to evaluate nodal point loads for arbitrary surface forces, although the expressions become rather involved for more complicated cases. This procedure is also applicable to body forces if the integration is executed over the volume of the element instead of its surface.

3.3 Solution of Simultaneous Equations

The equilibrium equation for the total system of elements can be written as

$$\underline{F} = \underline{K} \underline{\delta} \quad (31)$$

where \underline{K} = global stiffness matrix
 \underline{F} = nodal force vector
 $\underline{\delta}$ = nodal point displacement vector
 N = total number of nodal points

The vectors \underline{F} and $\underline{\delta}$ consist of N sub-vectors \underline{F}_n and $\underline{\delta}_n$, which contain the force and displacement components for each nodal point.

$$\underline{F}_n = \begin{bmatrix} F_{r_n} \\ F_{z_n} \end{bmatrix} \quad \underline{\delta}_n = \begin{bmatrix} u_n \\ v_n \end{bmatrix} \quad (32)$$

In the AX CST program an over-relaxation type of approach is used for solving the large system of simultaneous equations (Refs. 1,9). The basic formula for this iterative procedure can be written as:

$$\delta_n^{(j+1)} = \delta_n^{(j)} + \beta \underline{k}_{nn}^{-1} \left[\underline{F}_n - \sum_{i=1}^{n-1} \underline{k}_{ni} \delta_i^{(j+1)} - \sum_{i=n+1}^N \underline{k}_{ni} \delta_i^{(j)} \right] \quad (33)$$

where β = over-relaxation factor (1.85)
 j = number of iteration cycle
 \underline{k}_{ni} = (2x2) sub-matrix of global stiffness matrix
 \underline{k}_{nn}^{-1} = nodal point flexibility matrix

The use of this over-relaxation method offers a number of distinct advantages:

- (a) The large but sparsely populated global stiffness matrix can be stored very effectively, that is, for a given storage space a maximum number of nodal points can be handled.
- (b) The error build-up and propagation, a very serious problem for computers with short word length, is reduced to a minimum.
- (c) No restrictions concerning the "band-width" have to be imposed on the enumeration of the nodal points.

On the other hand, iterative methods of this type tend to require more execution time as compared to direct methods of solving simultaneous equations. This is particularly true for certain types of problems, where the convergence of the solution presents difficulties. However, the same reasons that cause a slow convergence of the iterative solution

might cause an increased error build-up for direct methods. This error remains undetectable unless elaborate error analysis is performed.

3.4 Displacement Boundary Conditions

From Eq. (33) the displacement increment per iteration cycle can be written as

$$\Delta \delta_n^{(j)} = \underline{k}_{nn}^{-1} \left[\underline{F}_n - \sum_{i=1}^{n-1} \underline{k}_{ni} \delta_i^{(j+1)} - \sum_{j=n+1}^N \underline{k}_{ni} \delta_i^{(j)} \right] \quad (34)$$

or

$$\Delta \delta_n^{(j)} = \underline{k}_{nn}^{-1} \cdot \underline{P}_n^{(j)} \quad (35a)$$

$$\begin{bmatrix} \Delta u_n \\ \Delta v_n \end{bmatrix}^{(j)} = \begin{bmatrix} f_{rr} & f_{rz} \\ f_{zr} & f_{zz} \end{bmatrix} \begin{bmatrix} P_r \\ P_z \end{bmatrix}^{(j)} \quad (35b)$$

\underline{P}_n denotes here the difference between the nodal point load \underline{F}_n and those forces caused by the displacement of adjacent nodal points. \underline{P}_n is commonly referred to as "unbalanced force".

Whereas Eq. (35) is valid for all nodal points without restraints, certain modifications of the nodal point flexibility matrix \underline{k}_{nn}^{-1} are necessary for those points where displacement boundary conditions are prescribed (Ref. 9).

In general, two types of restraints have to be distinguished:

- (1) the nodal point is fixed in all directions
- (2) the nodal point is free to move along a line with an angle φ

Denoting the support reactions with \underline{S}_n , Eq. (35) can be re-written for a restrained nodal point as follows:

$$\begin{bmatrix} \Delta u_n \\ \Delta v_n \end{bmatrix} = \begin{bmatrix} f_{rr}^* & f_{rz}^* \\ f_{zr}^* & f_{zz}^* \end{bmatrix} \begin{bmatrix} P_r + S_r \\ P_z + S_z \end{bmatrix} \quad (36)$$

For a Type 1 restraint the displacement increments $\Delta \delta_n$ are identical zero. It is apparent that this condition can only be met by setting the modified stiffness matrix \underline{k}_{nn}^{*-1} equal to zero.

$$f_{rr}^* = f_{rz}^* = f_{zr}^* = f_{zz}^* = 0 \quad (37)$$

For a Type 2 restraint (Fig. 5) Eq. (36) reads:

$$\begin{bmatrix} \Delta u_n \\ \Delta u_n \cdot \tan \varphi \end{bmatrix} = \begin{bmatrix} f_{rr}^* & f_{rz}^* \\ f_{zr}^* & f_{zz}^* \end{bmatrix} \begin{bmatrix} P_r - S \cdot \sin \varphi \\ P_z + S \cdot \cos \varphi \end{bmatrix} \quad (38)$$

Eliminating the unknown support reaction S results in the following expressions for \underline{k}_{nn}^{*-1}

$$\begin{aligned} f_{rr}^* &= \frac{f_{rr} - \alpha f_{rz}}{1 - \alpha \tan \varphi} & f_{rz}^* &= \frac{f_{rz} - \alpha f_{yy}}{1 - \alpha \tan \varphi} \\ f_{zr}^* &= f_{rr}^* \tan \varphi & f_{zz}^* &= f_{rz}^* \cdot \tan \varphi \end{aligned} \quad (39)$$

where

$$\alpha = \frac{f_{rr} \tan \varphi - f_{rz}}{f_{rz} \tan \varphi - f_{zz}}$$

Special attention has to be paid to the case where φ equals $\pi/2$. This corresponds to the case where nodal point is free to move in z - direction. Since $\tan\varphi$ is infinite in this case, Eq. (38) has to be rewritten

$$\begin{bmatrix} 0 \\ \Delta v_n \end{bmatrix} = \begin{bmatrix} f_{rr}^* & f_{rz}^* \\ f_{zr}^* & f_{zz}^* \end{bmatrix} \begin{bmatrix} P_r + S_r \\ P_z \end{bmatrix} \quad (40)$$

Expressing S_r in terms of P_r and P_z yields

$$f_{rr}^* = f_{rz}^* = f_{zr}^* = 0$$

$$f_{zz}^* = f_{zz} - \frac{f_{rz}^2}{f_{rr}}$$

3.5 General Description of Input

The minimum amount of information that has to be provided for a Finite Element analysis includes the following items:

1) Control Variables

Number of nodal points, number of elements, number of restrained nodal points, control variables for output. In case an iterative method for solving the system of simultaneous equations is used, a convergence limit and/or a maximum number of iteration cycles has to be provided.

2) Element Data

Nodal point numbers i, j, k
Material properties (Young's modulus, Poisson's ratio)

3) Nodal Point Data

Coordinates of nodal point.
Nodal point loads.

4) Restrained Nodal Point Data

Number of restrained nodal point.
Type of restraint and direction in which the nodal point is free to move.

In the AXCST program two features were incorporated which require additional input information. These features are:

- 1) The program will generate nodal point loads due to thermal strain. For this purpose the coefficient of thermal expansion and the temperature change has to be provided for each element.
- 2) The program will read in initial values for the nodal point displacements. These values have the character of initial guesses for unrestrained nodal points, whereas they are treated as specified displacements for restrained nodal points.

It is possible to incorporate numerous other features in this program. Of particular interest would be program parts for generating nodal point loads due to various types of surface and body forces. However, to keep the input as general and at the same time as simple as possible, no provisions of this type were made. For the same reason all input data have to be punched on separate cards for each element or nodal point respectively. It is apparent that for specific types of problems the amount of input data could be greatly reduced. In those cases it is left to the users to modify the input of this program according to the specific purpose.

3.6 General Description of Output

The Finite Element analysis of elasticity problems provides the following numerical results:

- 1) Nodal point displacements
- 2) Element strains and stresses

These nodal point displacements define the displacements throughout the body uniquely since the linear displacement function used in the AXCST program satisfies compatibility along the edges of the elements (Section 2.1). However, the element stresses and strains will generally not match along the interface of two adjacent elements.

From comparisons with elasticity problems, of which an exact solution is known, it was found that the best numerical results are obtained by assigning the element stresses and strains to the center of the element. To obtain stress and strain values also for the nodal points a number of authors applied certain weighted average procedures. Unfortunately, the criteria used by different authors to determine the weighting coefficients contradict each other in a certain way.

A second problem arises in connection with nodal point stresses. In many elasticity problems the maximum stresses occur at the boundary of the body rather than in its interior. In this case the stresses at the boundary points will be biased noticeably, since their value is based on the average of the interior element stresses, not on their extrapolated values. Therefore, it seems to be doubtful whether the nodal point stresses should be used at all as a result of the Finite Element analysis.

The AXCST program provides the following output:

1) Input Data

Control variables

Element data (including center coordinates)

Nodal point data

Restrained nodal point data

2) Force Unbalance (printed in specified cycle intervals)3) Displacement and Stresses (printed in specified cycle intervals)

Number of cycle, corresponding force unbalance

Nodal point displacements

Element stresses (σ_r , σ_z , σ_θ , τ_{rz} , σ_1 , σ_2 , direction of σ_1)

In case there are distinct reasons to use nodal point stresses as a result of the analysis, their values can be calculated and printed out by specifying the proper control variables (see Section 4.1).

4. USER'S GUIDE4.1 Input Information

A. Header Card - FORMAT (1X,52H.....)

B. Control Card - FORMAT (6I4,6X,E10.3)

<u>Cols.</u>	<u>Input</u>
1 - 4	Number of elements
5 - 8	Number of nodal points
9 - 12	Number of restrained boundary points
13 - 16	Print interval of force unbalance
17 - 20	Print interval of displacements and stresses
21 - 24	Maximum number of cycles
31 - 40	Convergence limit for unbalanced force
50	0 - output of nodal point stresses suppressed 1 - output of nodal point stresses executed

C. Element Data - 1 Card per element FORMAT (3I4,8X,4F10.0)

<u>Cols.</u>	<u>Input</u>
1 - 4	Nodal point number i
5 - 8	Nodal point number j
9 - 12	Nodal point number k
21 - 30	Modulus of elasticity
31 - 40	Poisson's ratio
41 - 50	Coefficient of thermal expansion
51 - 60	Temperature change

D. Nodal Point Data - 1 Card per nodal point FORMAT (6F10.0)

<u>Cols.</u>	<u>Input</u>
1 - 10	R - Ordinate
11 - 20	Z - Ordinate
21 - 30	R - Load
31 - 40	Z - Load
41 - 50	R - Displacement (*)
51 - 60	Z - Displacement (*)

Note: (*) On free nodal points these are initial guesses; on restrained nodal points these are specified displacements.

E. Boundary Point Data - 1 Card per boundary point FORMAT (2I4,2X,F10.0)

<u>Cols.</u>	<u>Input</u>
1 - 4	Nodal point number
8	0 - nodal point is fixed in both directions 1 - nodal point is fixed in R direction 2 - nodal point is free to move along a line of slope S

<u>Cols.</u>	<u>Input</u>
11 - 20	Slope S (type 2 boundary point only)

Note: Cards of type C and type D have to be placed in their natural sequence. To avoid errors the element or nodal point number should be punched in columns 71 - 80.

4.2 Assumptions, Limitations, Remarks

A. Limitations

- 1) Maximum number of elements: 600
- 2) Maximum number of nodal points: 310
- 3) Maximum number of nodal points adjacent to a certain point: 8
- 4) Elements and nodal points have to be numbered in natural sequence (data cards have to be placed in the same order).

B. Assumption

Over-relaxation factor: 1.85

C. Initial guesses for displacements

For boundary condition type 2 (nodal point is forced to move along a line of slope S): Initial guesses of displacements perpendicular to the prescribed slope will be treated as imposed displacements.

D. Direction of principal stresses

The calculated value is the angle between the positive R axis and the direction of the maximum positive (or minimum negative) principal stress. In case that both values ($\sigma_r - \sigma_z$) and τ_{rz} are very small this angle becomes meaningless.

E. Required field length

CM - 100,000_g

By using CM - 140,000_g the program can handle up to 530 nodal points and 1000 elements.

5. REFERENCES

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6. PROGRAM LISTING

```

PROGRAM AXCST (INPUT,TAPE1=INPUT,OUTPUT,TAPE2=CUTPUT)
AXCST 1
*
AXCST 2
*
FINITE ELEMENT ANALYSIS OF AXI-SYMMETRIC SOLIDS
AXCST 3
*
AXCST 4
*
DIMENSION NPI(600),NPJ(600),NPK(600),E(600),P(600),AL(600),DT(600)
AXCST 5
DIMENSION TH(600),SR(600),SZ(600),ST(600),TA(600)
AXCST 6
DIMENSION R(310),Z(310),RLCAD(310),ZLOAD(310),DR(310),DZ(310)
AXCST 7
DIMENSION SRR(310,9),SRZ(310,9),SZR(310,9),SZZ(310,9),NN(310,9)
AXCST 8
DIMENSION FRR(310),FRZ(310),FZR(310),FZZ(310)
AXCST 9
DIMENSION B(6,6),D(6,6),S(6,6)
AXCST 10
DIMENSION AA(3),BB(3),CC(3),AO(3),RA(3),RB(3),LL(3)
AXCST 11
EQUIVALENCE (AL(600),TH(600)),(DT(600),SR(600))
AXCST 12
IN=1
AXCST 13
IO=2
AXCST 14
DATA RELA/1.85/
AXCST 15
*
AXCST 16
*
INPUT
AXCST 17
*
AXCST 18
*
WRITE(IO,300)
AXCST 19
READ(IN,400)
AXCST 20
WRITE(IO,400)
AXCST 21
WRITE(IO,301)
AXCST 22
READ(IN,401) NELE,NNOP,NBOP,INUN,INST,MAXC,TOLE,NPST
AXCST 23
WRITE(IO,302) NELE
AXCST 24
WRITE(IO,303) NNOP
AXCST 25
WRITE(IO,304) NBOP
AXCST 26
WRITE(IO,305) INUN
AXCST 27
WRITE(IO,306) INST
AXCST 28
WRITE(IO,307) MAXC
AXCST 29
WRITE(IO,308) TOLE
AXCST 30
WRITE(IO,309) RELA
AXCST 31
READ(IN,402) (NPI(M),NPJ(M),NPK(M),E(M),P(M),AL(M),DT(M),M=1,NELE)
AXCST 32
READ(IN,403) (R(N),Z(N),RLOAD(N),ZLOAD(N),DR(N),DZ(N),N=1,NNOP)
AXCST 33
WRITE(IO,310)
AXCST 34
WRITE(IO,311) (N,R(N),Z(N),RLOAD(N),ZLOAD(N),DR(N),DZ(N),N=1,NNOP)
AXCST 35
WRITE(IO,312)
AXCST 36
DO 101 N=1,NNOP
AXCST 37
DO 100 L=1,9
AXCST 38
SRR(N,L)=SRZ(N,L)=SZR(N,L)=SZZ(N,L)=0.0
AXCST 39
100 NN(N,L)=0
AXCST 40
101 NN(N,1)=N
AXCST 41
*
AXCST 42
*
FORMATION OF ELEMENT STIFFNESS MATRIX
AXCST 43
*
AXCST 44
*
PI=4.0*ATAN(1.0)
AXCST 45
CD=90.0/PI
AXCST 46
DO 112 M=1,NELE
AXCST 47
DO 102 II=1,6
AXCST 48
DO 102 JJ=1,6
AXCST 49
102 B(II,JJ)=D(II,JJ)=S(II,JJ)=0.0
AXCST 50
I=LL(1)=NPI(M)
AXCST 51
J=LL(2)=NPJ(M)
AXCST 52
K=LL(3)=NPK(M)
AXCST 53
AA(1)=R(J)*Z(K)-R(K)*Z(J)
AXCST 54
AA(2)=R(K)*Z(I)-R(I)*Z(K)
AXCST 55
AA(3)=R(I)*Z(J)-R(J)*Z(I)
AXCST 56
BB(1)=Z(J)-Z(K)
AXCST 57

```

BB(2)=Z(K)-Z(I)	AXCST 58
BB(3)=Z(I)-Z(J)	AXCST 59
CC(1)=R(K)-R(J)	AXCST 60
CC(2)=R(I)-R(K)	AXCST 61
CC(3)=R(J)-R(I)	AXCST 62
RC=(R(I)+R(J)+R(K))/3.0	AXCST 63
ZC=(Z(I)+Z(J)+Z(K))/3.0	AXCST 64
WRITE(IO,313) M,I,J,K,E(M),P(M),AL(M),DT(M),RC,ZC	AXCST 65
ZR=ZC/RC	AXCST 66
AREA=(BB(2)*CC(3)-BB(3)*CC(2))/2.0	AXCST 67
IF(AREA.EQ.0.0) GO TO 135	AXCST 68
ARE=ABS(AREA)	AXCST 69
TT=2.0*AREA/(3.0*RC)	AXCST 70
PP=1.0-2.0*P(M)	AXCST 71
TH(M)=E(M)*AL(M)*DT(M)/PP	AXCST 72
THE=PI*RC*TH(M)*SIGN(1.0,AREA)	AXCST 73
DO 103 L=1,3	AXCST 74
LB=2*L	AXCST 75
LA=LB-1	AXCST 76
LC=LL(L)	AXCST 77
RLOAD(LC)=RLOAD(LC)+THE*(BB(L)+TT)	AXCST 78
ZLOAD(LC)=ZLOAD(LC)+THE*CC(L)	AXCST 79
B(1,LA)=B(4,LB)=BB(L)	AXCST 80
B(2,LB)=B(4,LA)=CC(L)	AXCST 81
103 B(3,LA)=TT	AXCST 82
DC=E(M)*PI*RC/(4.0*ARE*(1.0+P(M)))	AXCST 83
DB=DC*2.0/PP	AXCST 84
DA=DB*(1.0-P(M))	AXCST 85
DB=DB*P(M)	AXCST 86
DD=DA/RC	AXCST 87
D(1,1)=D(2,2)=D(3,3)=DA	AXCST 88
D(1,2)=D(1,3)=D(2,1)=D(2,3)=D(3,1)=D(3,2)=DB	AXCST 89
D(4,4)=DC	AXCST 90
DO 104 II=1,4	AXCST 91
DO 104 JJ=1,6	AXCST 92
S(II,JJ)=0.0	AXCST 93
DO 104 KK=1,4	AXCST 94
104 S(II,JJ)=S(II,JJ)+D(II,KK)*B(KK,JJ)	AXCST 95
DO 105 II=1,4	AXCST 96
DO 105 JJ=1,6	AXCST 97
105 D(II,JJ)=S(II,JJ)	AXCST 98
DO 106 II=1,6	AXCST 99
DO 106 JJ=1,6	AXCST100
S(II,JJ)=0.0	AXCST101
DO 106 KK=1,4	AXCST102
106 S(II,JJ)=S(II,JJ)+B(KK,II)*D(KK,JJ)	AXCST103
BLOI=BLOJ=BLOK=0.0	AXCST104
IF(R(I).GT.0.0) BLOI=ALOG(R(I))	AXCST105
IF(R(J).GT.0.0) BLOJ=ALOG(R(J))	AXCST106
IF(R(K).GT.0.0) BLOK=ALOG(R(K))	AXCST107
A0(1)=BLOK-BLOJ	AXCST108
A0(2)=BLOI-BLOK	AXCST109
A0(3)=BLOJ-BLOI	AXCST110
RA(1)=R(K)+R(J)	AXCST111
RA(2)=R(I)+R(K)	AXCST112
RA(3)=R(J)+R(I)	AXCST113
RB(1)=RA(1)*RA(1)-R(K)*R(J)	AXCST114
RB(2)=RA(2)*RA(2)-R(I)*R(K)	AXCST115

RB(3)=RA(3)*RA(3)-R(J)*R(I)	AXCST116
FA=FB=FC=0.0	AXCST117
DO 107 L=1,3	AXCST118
IF(CC(L).EQ.0.0) GO TO 107	AXCST119
AC=AA(L)/CC(L)	AXCST120
BC=BE(L)/CC(L)	AXCST121
CO=AC*AC(L)	AXCST122
BO=BC*BB(L)	AXCST123
FA=FA+CO+BB(L)	AXCST124
FB=FB-AC*(CO/2.0+BB(L))-BO*RA(L)/4.0	AXCST125
FC=FC+AC*AC*(CO/3.0+BB(L))+BO*(AC*RA(L)/2.0+BC*RE(L)/9.0)	AXCST126
107 CONTINUE	AXCST127
FA=FA/AREA-1.0/RC	AXCST128
FB=FB/AREA-ZR	AXCST129
FC=FC/AREA-ZR*ZC	AXCST130
DO 108 II=1,3	AXCST131
DO 108 JJ=1,3	AXCST132
IS=2*II-1	AXCST133
JS=2*JJ-1	AXCST134
108 S(IS,JS)=S(IS,JS)+DC*(AA(II)*AA(JJ)*FA+(AA(II)*CC(JJ)+AA(JJ)*CC(II	AXCST135
1))*FB+CC(II)*CC(JJ)*FC)	AXCST136
*	AXCST137
* FORMATION OF GLOBAL STIFFNESS MATRIX	AXCST138
*	AXCST139
DO 111 II=1,3	AXCST140
DO 111 JJ=1,3	AXCST141
LM=LL(II)	AXCST142
MM=0	AXCST143
109 MM=MM+1	AXCST144
IF(MM.GE.10) GO TO 136	AXCST145
IF(NN(LM,MM).EQ.LL(JJ)) GO TO 110	AXCST146
IF(NN(LM,MM).NE.0) GO TO 109	AXCST147
110 NN(LM,MM)=LL(JJ)	AXCST148
IB=2*II	AXCST149
IA=IB-1	AXCST150
JB=2*JJ	AXCST151
JA=JB-1	AXCST152
SRR(LM,MM)=SRR(LM,MM)+S(IA,JA)	AXCST153
SRZ(LM,MM)=SRZ(LM,MM)+S(IA,JB)	AXCST154
SZR(LM,MM)=SZR(LM,MM)+S(IB,JA)	AXCST155
111 SZZ(LM,MM)=SZZ(LM,MM)+S(IB,JB)	AXCST156
112 CONTINUE	AXCST157
DO 114 N=1,NNOP	AXCST158
NO=10	AXCST159
113 NO=NO-1	AXCST160
IF(NN(N,NO).EQ.0) GO TO 113	AXCST161
114 NN(N,1)=NO	AXCST162
DO 115 N=1,NNOP	AXCST163
DET=SRR(N,1)*SZZ(N,1)-SRZ(N,1)*SZR(N,1)	AXCST164
FRR(N)=SZZ(N,1)/DET	AXCST165
FRZ(N)=-SRZ(N,1)/DET	AXCST166
FZR(N)=-SZR(N,1)/DET	AXCST167
115 FZZ(N)=SRR(N,1)/DET	AXCST168
*	AXCST169
* DISPLACEMENT BOUNDARY CONDITIONS	AXCST170
*	AXCST171
IF(NBOP.EQ.0) GO TO 121	AXCST172
WRITE(IO,314)	AXCST173

WRITE(IO,315)	AXCST174
DO 120 L=1,NBOP	AXCST175
READ(IN,404) NP,NT,SLO	AXCST176
WRITE(IO,316) NP,NT,SLO	AXCST177
IF(NT-1) 118,117,116	AXCST178
116 DET=(FRR(NP)*SLO-FRZ(NP))/(FZR(NP)*SLO-FZZ(NP))	AXCST179
COF=1.0-DET*SLO	AXCST180
FRR(NP)=(FRR(NP)-DET*FZR(NP))/COF	AXCST181
FRZ(NP)=(FRZ(NP)-DET*FZZ(NP))/COF	AXCST182
FZR(NP)=FRR(NP)*SLO	AXCST183
FZZ(NP)=FRZ(NP)*SLO	AXCST184
GO TO 120	AXCST185
117 FZZ(NP)=FZZ(NP)-FZR(NP)*FRZ(NP)/FRR(NP)	AXCST186
GO TO 119	AXCST187
118 FZZ(NP)=0.0	AXCST188
119 FRR(NP)=FRZ(NP)=FZR(NP)=0.0	AXCST189
*	AXCST190
* ITERATIVE SOLUTION OF SIMULTANEOUS EQUATIONS.	AXCST191
*	AXCST192
120 CONTINUE	AXCST193
121 NCY=0	AXCST194
IUN=INUN	AXCST195
IST=INST	AXCST196
IND=1	AXCST197
122 WRITE(IO,317)	AXCST198
123 SUM=0.0	AXCST199
DO 125 N=1,NNOP	AXCST200
NUM=NN(N,1)	AXCST201
FR=RLOAD(N)	AXCST202
FZ=ZLOAD(N)	AXCST203
DO 124 L=2,NUM	AXCST204
M=NN(N,L)	AXCST205
FR=FR-SRR(N,L)*DR(M)-SRZ(N,L)*DZ(M)	AXCST206
124 FZ=FZ-SZR(N,L)*DR(M)-SZZ(N,L)*DZ(M)	AXCST207
DDR=FRR(N)*FR+FRZ(N)*FZ-DR(N)	AXCST208
DDZ=FZR(N)*FR+FZZ(N)*FZ-DZ(N)	AXCST209
DR(N)=DR(N)+RELA*DDR	AXCST210
DZ(N)=DZ(N)+RELA*DDZ	AXCST211
125 SUM=SUM+ABS(SRR(N,1)*DDR+SRZ(N,1)*DDZ)+ABS(SZR(N,1)*EDR+SZZ(N,1)*D	AXCST212
DDZ)	AXCST213
NCY=NCY+1	AXCST214
IF(NCY.LT.IUN) GO TO 126	AXCST215
IUN=IUN+INUN	AXCST216
WRITE(IO,318) NCY,SUM	AXCST217
126 IF(SUM.LE.TOLE.OR.NCY.GE.MAXC) IND=2	AXCST218
IF(NCY.GE.IST) GO TO 127	AXCST219
GO TO (123,127) IND	AXCST220
*	AXCST221
* NODAL POINT DISPLACEMENTS	AXCST222
*	AXCST223
127 WRITE(IO,319) NCY,SUM	AXCST224
WRITE(IO,320)	AXCST225
DO 128 N=1,NNOP,2	AXCST226
L=N+1	AXCST227
IF(L.GT.NNOP) L=NNOP	AXCST228
128 WRITE(IO,321) N,DR(N),DZ(N),L,DR(L),DZ(L)	AXCST229
*	AXCST230
* ELEMENT STRESSES	AXCST231

```

*
WRITE(I0,322)
WRITE(I0,323)
IST=IST+INST
DO 129 M=1,NELE
I=NPI(M)
J=NPJ(M)
K=NPK(M)
BI=Z(J)-Z(K)
BJ=Z(K)-Z(I)
BK=Z(I)-Z(J)
CI=R(K)-R(J)
CJ=R(I)-R(K)
CK=R(J)-R(I)
AREA=(BJ*CK-BK*CJ)/2.0
TT=2.0*AREA/(R(I)+R(J)+R(K))
ER=BI*DR(I)+BJ*DR(J)+BK*DR(K)
EZ=CI*DZ(I)+CJ*DZ(J)+CK*DZ(K)
ET=TT*(DR(I)+DR(J)+DR(K))
GA=CI*DR(I)+CJ*DR(J)+CK*DR(K)+BI*DZ(I)+BJ*DZ(J)+BK*DZ(K)
DC=E(M)/(4.0*AREA*(1.0+P(M)))
DB=DC*2.0/(1.0-2.0*P(M))
DA=DB*(1.0-P(M))
DB=DB*P(M)
SR(M)=DA*ER+DB*(EZ+ET)-TH(M)
SZ(M)=DA*EZ+DB*(ER+ET)-TH(M)
ST(M)=DA*ET+DB*(ER+EZ)-TH(M)
TA(M)=DC*GA
129 CALL STRESS (SR(M),SZ(M),ST(M),TA(M),M,CD)
GO TO (122,130) IND
*
*
NODAL POINT STRESSES
*
130 IF(NPST.NE.1) GO TO 500
WRITE(I0,324)
WRITE(I0,325)
DO 134 N=1,NNOP
SSR=SSZ=SST=TTA=FR=FZ=FT=0.0
DO 133 M=1,NELE
I=NPI(M)
J=NPJ(M)
K=NPK(M)
IF(N.EQ.I) GO TO 132
IF(N.EQ.J) GO TO 131
IF(N.NE.K) GO TO 133
I=NPK(M)
K=NPI(M)
GO TO 132
131 I=NPJ(M)
J=NPI(M)
132 QR=ABS(R(J)+R(K)-2.0*R(I))
QZ=ABS(Z(J)-Z(I))+ABS(Z(K)-Z(I))
Q=QR+QZ
FFR=QR/Q
FFZ=QZ/Q
FR=FR+FFR
FZ=FZ+FFZ
FT=FT+1.0

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AXCST232
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AXCST286
AXCST287
AXCST288
AXCST289

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```

SSR=SSR+SR(M)*FFR
SSZ=SSZ+SZ(M)*FFZ
SST=SST+ST(M)
TTA=TTA+TA(M)
133 CONTINUE
PSR=SSR/FR
PSZ=SSZ/FZ
PST=SST/FT
PTA=TTA/FT
134 CALL STRESS (PSR,PSZ,PST,PTA,N,CD)
GO TO 500
135 WRITE(IC,326) M
GO TO 500
136 WRITE(IO,327) LM
300 FORMAT(1H1,/)
301 FORMAT(///,* INPUT DATA*,///)
302 FORMAT(* NUMBER OF ELEMENTS =*,I5,/)
303 FORMAT(* NUMBER OF NODAL POINTS =*,I5,/)
304 FORMAT(* NUMBER OF BOUNDARY POINTS =*,I5,/)
305 FORMAT(* PRINT INTERVAL UNBALANCE =*,I5,/)
306 FORMAT(* PRINT INTERVAL STRESSES =*,I5,/)
307 FORMAT(* CYCLE LIMIT =*,I5,/)
308 FORMAT(* TOLERANCE LIMIT =*,E13.3,/)
309 FORMAT(* OVER-RELAXATION FACTOR =*,F13.3,/)
310 FORMAT(1H1,/,/* POINT*,3X,*R-ORD*,10X,*Z-ORD*,9X,*R-LOAD*,9X,*Z-LOAD*
1AD*,10X,*R-DISP*,11X,*Z-DISP*,//)
311 FORMAT(I4,1X,4F15.5,2F17.8)
312 FORMAT(1H1,/,/* ELEMENT*,4X,*I*,7X,*J*,7X,*K*,6X,*E-MODULUS*,6X,*PAXISSON*,7X,*ALPHA*,8X,*DELTA T*,11X,*R-CENT*,7X,*Z-CENT*,//)
313 FORMAT(1X,4(I4,4X),E11.3,F12.4,F15.8,F13.3,5X,2F13.5)
314 FORMAT(1H1,/,/* DISPLACEMENT BOUNDARY CONDITIONS*,//)
315 FORMAT(* POINT*,9X,*TYPE*,13X,*SLOPE*,//)
316 FORMAT(I4,13X,I1,F20.6)
317 FORMAT(1H1,/,/* CYCLE*,6X,*FORCE UNBALANCE*,/)
318 FORMAT(I5,8X,E12.5)
319 FORMAT(1H1,*CYCLE !*,I6,15X,*FORCE UNBALANCE !*,E15.5,///)
320 FORMAT(//,2(* POINT*,9X,*R-DISPLACEMENT*,6X,*Z-DISPLACEMENT*,18X),
1//)
321 FORMAT(2(I4,5X,2E20.7,18X))
322 FORMAT(1H1,/,/* ELEMENT STRESSES*)
323 FORMAT(////,* ELEMENT*,11X,*SIGMA R*,11X,*SIGMA Z*,11X,*SIGMA T*,11X,*TAU R-Z*,11X,*SIG-MAX*,11X,*SIG-MIN*,11X,*DIRECTION*,//)
324 FORMAT(1H1,/,/* NODAL POINT STRESSES (WEIGHTED AVERAGE OF ELEMENT STRESSES)*
)
325 FORMAT(////,* POINT *,11X,*SIGMA R*,11X,*SIGMA Z*,11X,*SIGMA T*,11X,*TAU R-Z*,11X,*SIG-MAX*,11X,*SIG-MIN*,11X,*DIRECTION*,//)
326 FORMAT(1H1,10X,*ZERO AREA ELEMENT NUMBER !*,I5)
327 FORMAT(1H1,10X,*MORE THAN 8 NODAL POINTS ADJACENT TO POINT !*,I5)
400 FORMAT(1X,52H)
401 FORMAT(6I4,E16.3,I10)
402 FORMAT(3I4,8X,4F10.0)
403 FORMAT(6F10.0)
404 FORMAT(2I4,F12.0)
500 CALL EXIT
END

```

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AXCST290
AXCST291
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AXCST301
AXCST302
AXCST303
AXCST304
AXCST305
AXCST306
AXCST307
AXCST308
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AXCST343
AXCST344

```



```
* SUBROUTINE STRESS (R,Z,T,G,N,CD) AXCST345
* DIRECTION OF PRINCIPAL STRESSES AXCST346
* AXCST347
IO=2 AXCST348
SS=(R+Z)/2.0 AXCST349
SD=R-Z AXCST350
SA=SQRT(SD*SD/4.0+G*G) AXCST351
SMAX=SS+SA AXCST352
SMIN=SS-SA AXCST353
TT=SIGN(1.0,G) AXCST354
IF(SD.NE.0.0) GO TO 100 AXCST355
FI=TT*45.0 AXCST356
GO TO 101 AXCST357
100 FI=CD*ATAN(2.0*G/SD) AXCST358
IF(SD.LT.0.0) FI=TT*90.0+FI AXCST359
101 WRITE(IO,300) N,R,Z,T,G,SMAX,SMIN,FI AXCST360
300 FORMAT(I5,3X,6F18.4,F18.2) AXCST361
RETURN AXCST362
END AXCST363
AXCST364
```

7. ACKNOWLEDGEMENTS

Thanks and appreciation are due to Mrs. K. Michele Kostem for her help in the preparation of the report and to Mrs. Jane Lenner for typing and proofreading the manuscript.

8. FIGURES

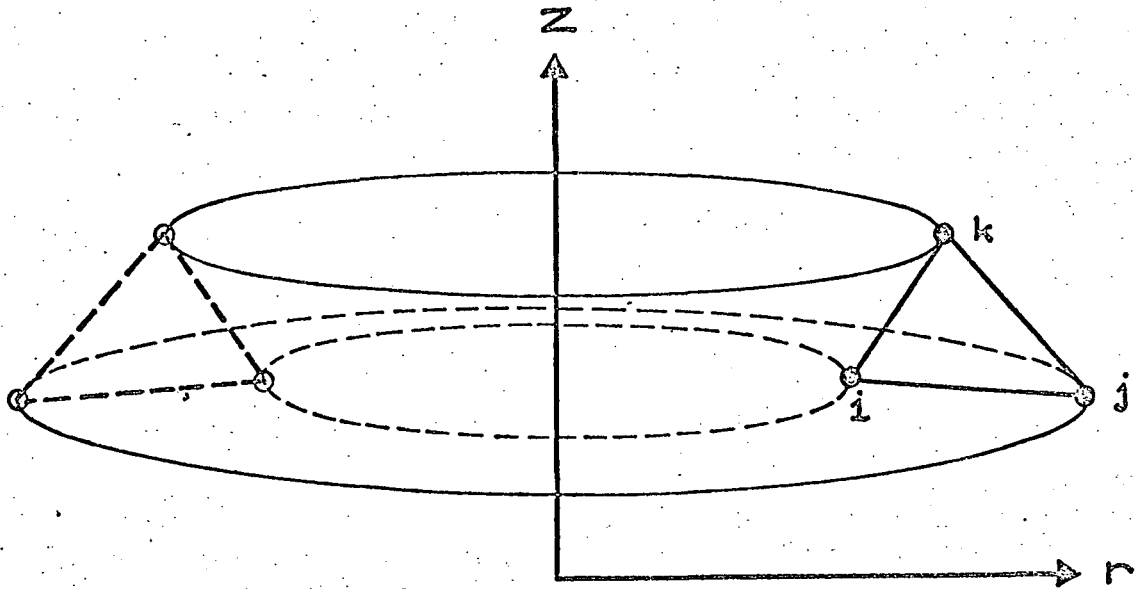


Fig. 1: Triangular Ring Element

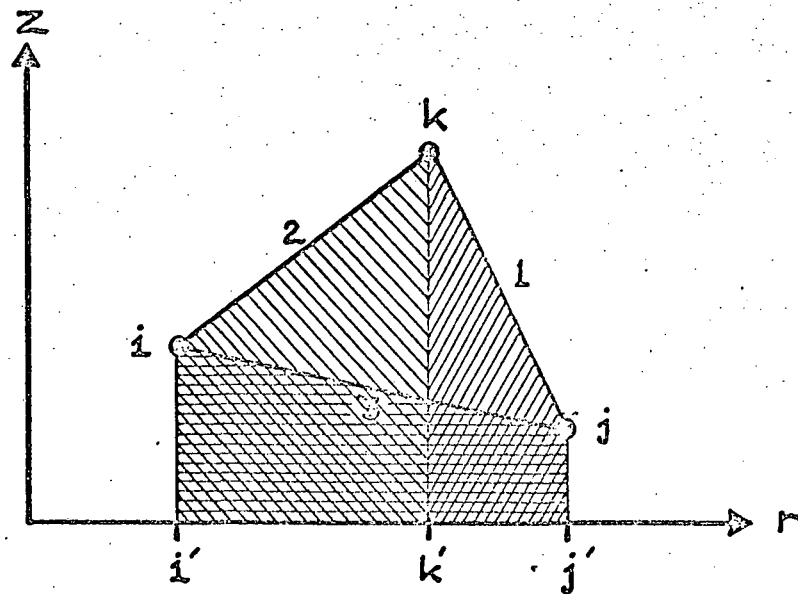


Fig. 2: Integration Scheme

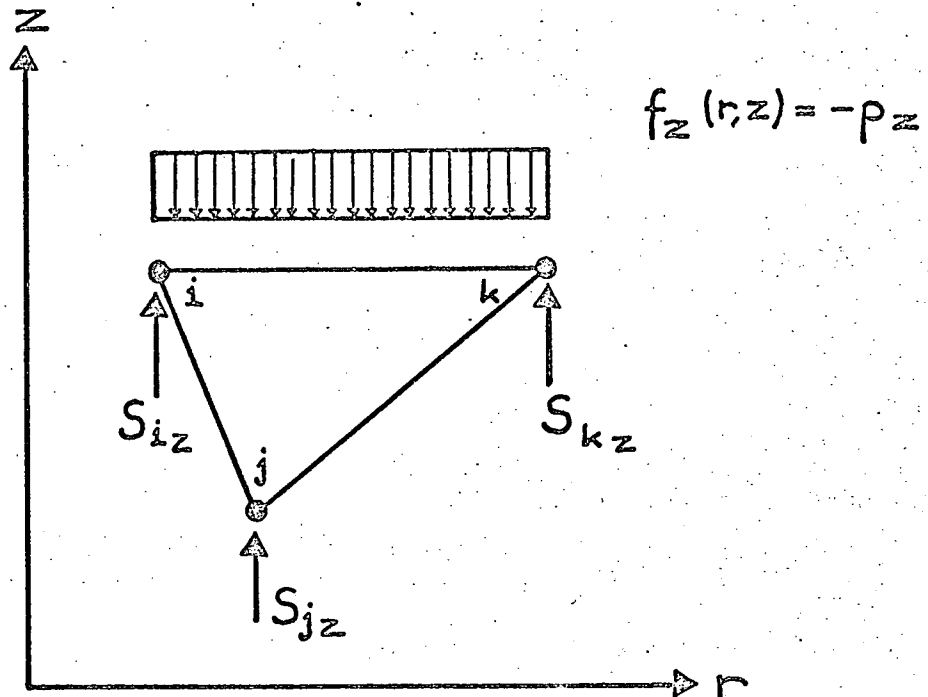
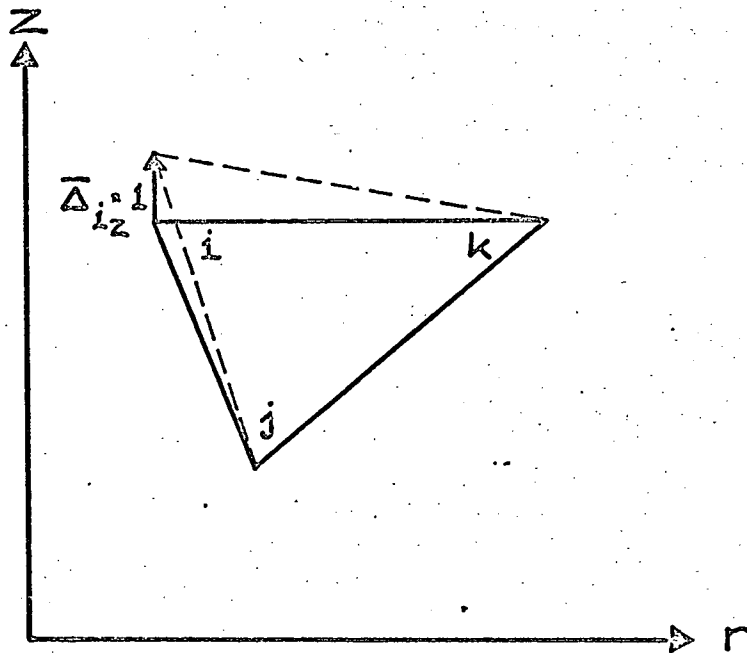


Fig. 3a: Element Geometry and Loading

Fig. 3b: Virtual Displacement $\bar{\Delta}_{iz}$

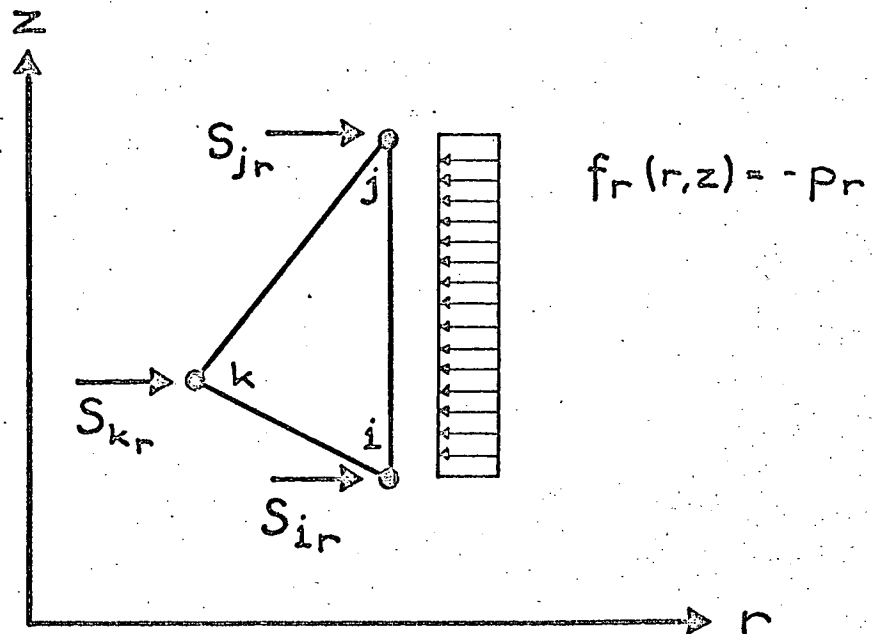


Fig. 4: Element Geometry and Loading

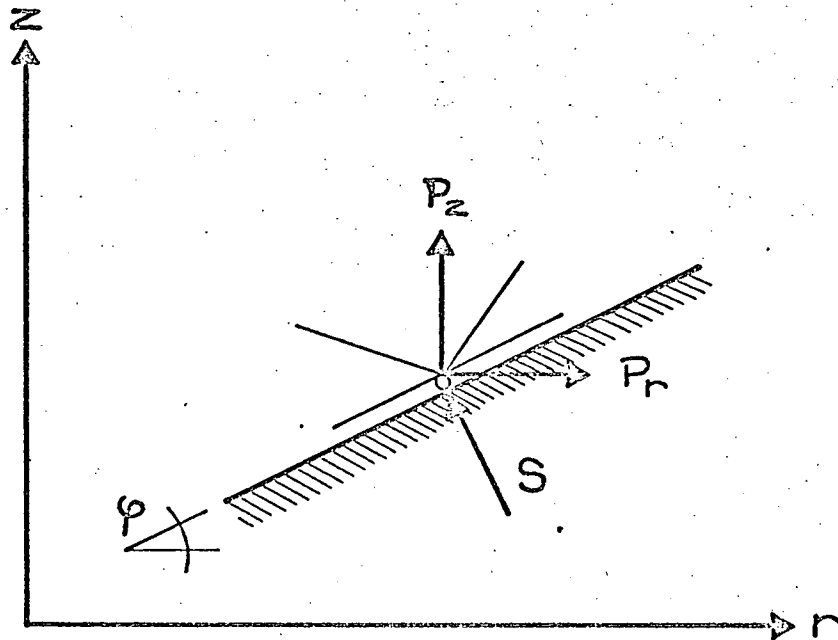
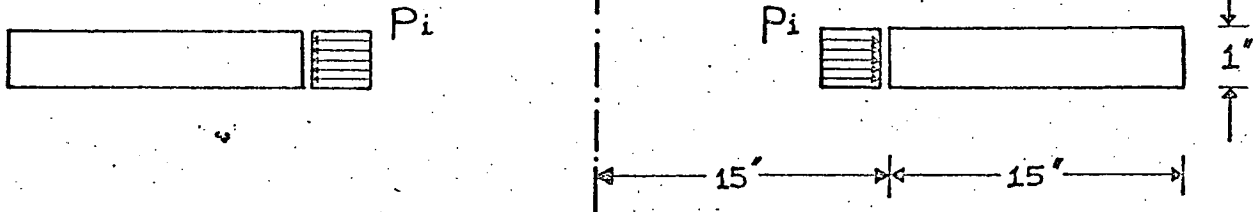


Fig. 5: Restrained Boundary Point

9. APPENDIX

EXAMPLE: Thick-walled cylinder under internal pressure (plane stress solution)

Geometry



Elastic Constants and Loading

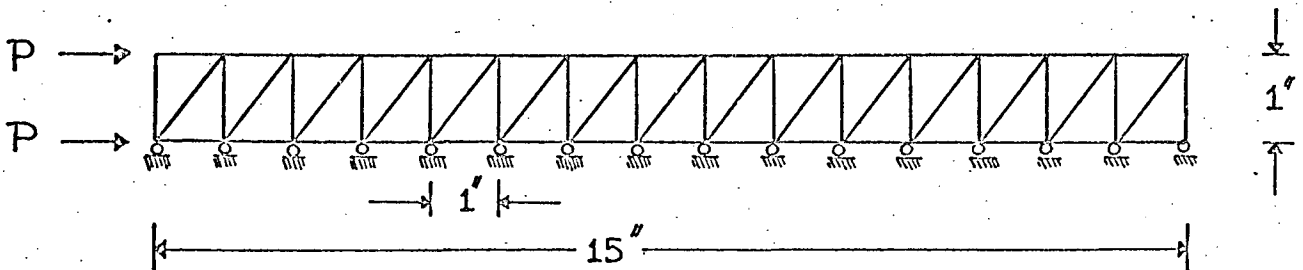
$$E = 30,000 \text{ ksi}$$

$$p_i = 100 \text{ ksi}$$

$$\nu = 0.30$$

$$P = 4712.39 \text{ kip}$$

Element Layout



Number of Elements: 30

Number of Nodal Points: 32

Units

Stresses: ksi

Length, displacements: in.

Note: Solid lines denote exact solution; dots and circles denote Finite Element solution

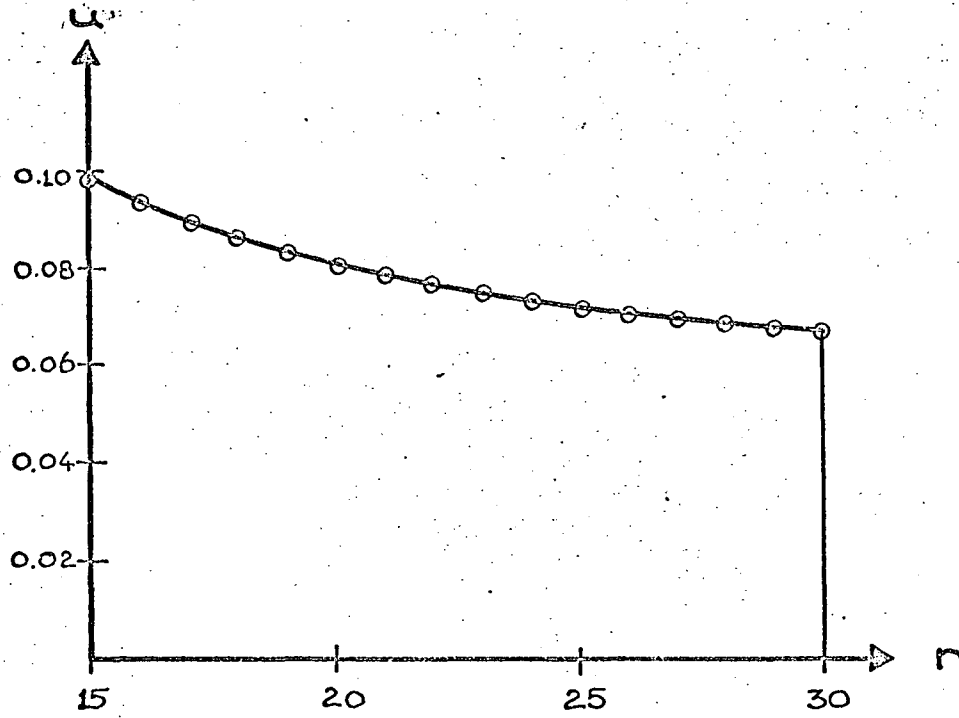
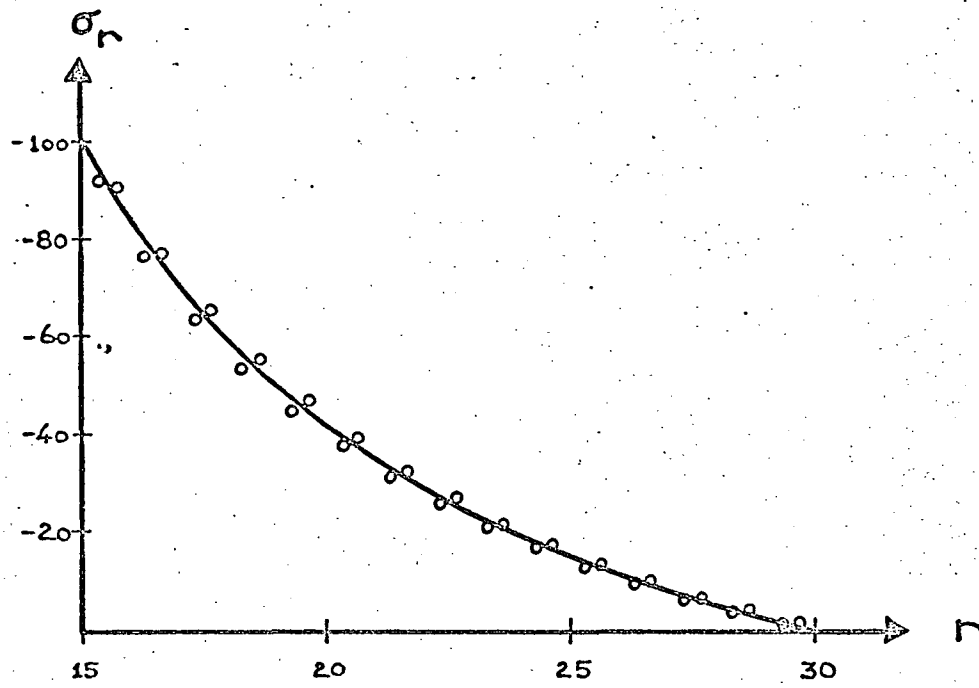
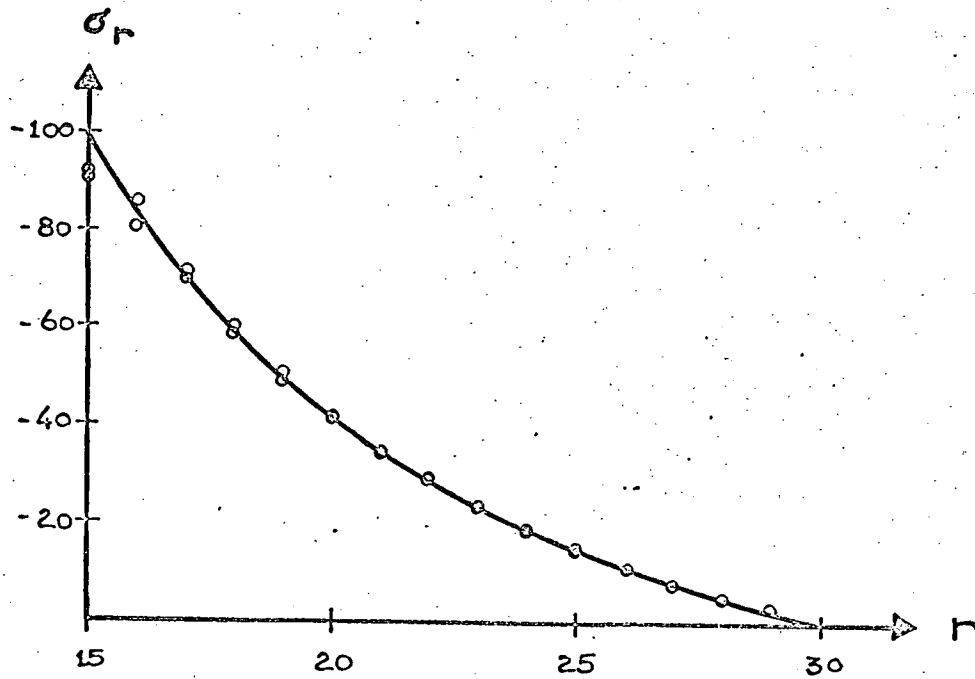
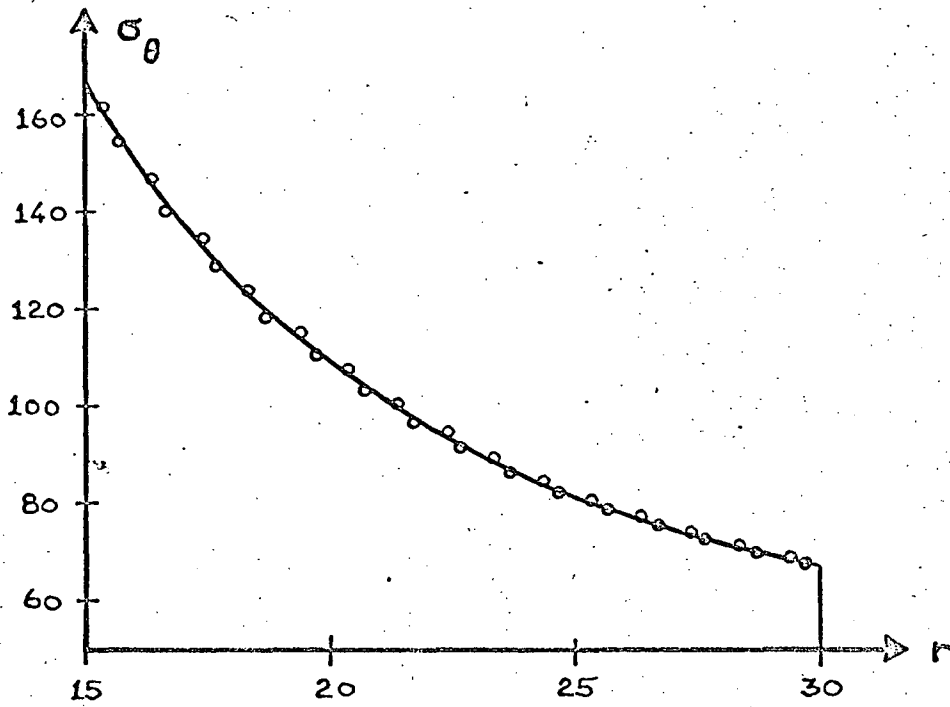
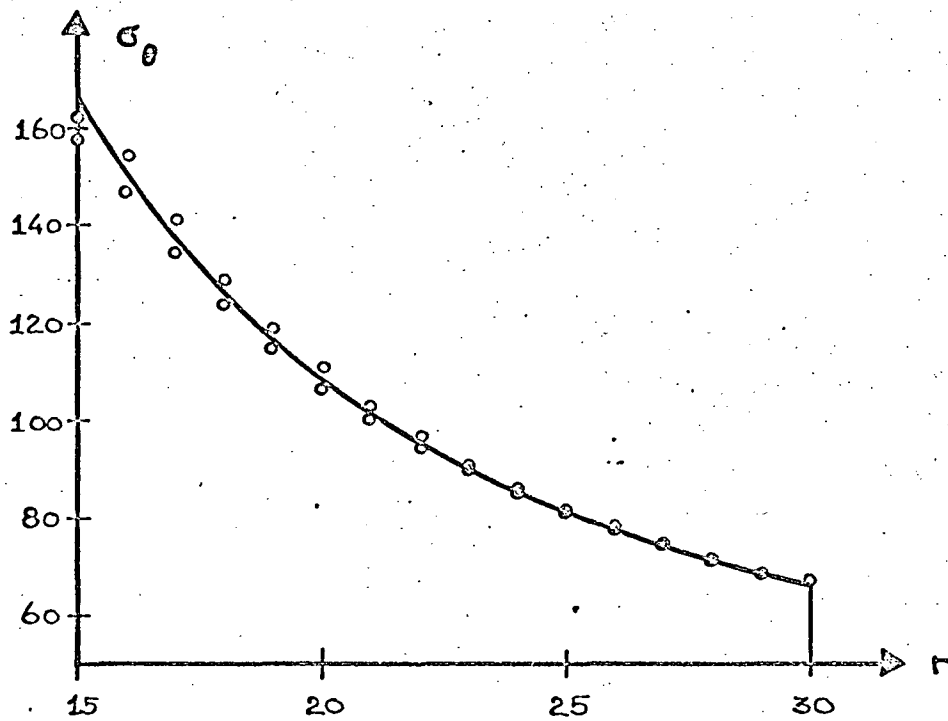


Fig. 1: Nodal Point Displacements

Fig. 2a: Element Stresses σ_r Fig. 2b: Nodal Point Stresses σ_r

Fig. 3a: Element Stresses σ_θ Fig. 3b: Nodal Point Stresses σ_θ