



## Philosophical Magazine Series 6

ISSN: 1941-5982 (Print) 1941-5990 (Online) Journal homepage: <http://www.tandfonline.com/loi/tphm17>

# XLVI. A reciprocal relation in diffraction

A.A. Michelson

To cite this article: A.A. Michelson (1905) XLVI. A reciprocal relation in diffraction , Philosophical Magazine Series 6, 9:52, 506-507, DOI: [10.1080/14786440509463300](https://doi.org/10.1080/14786440509463300)

To link to this article: <http://dx.doi.org/10.1080/14786440509463300>



Published online: 08 Jun 2010.



Submit your article to this journal [↗](#)



Article views: 1



View related articles [↗](#)



Citing articles: 4 View citing articles [↗](#)

Full Terms & Conditions of access and use can be found at  
<http://www.tandfonline.com/action/journalInformation?journalCode=6phm20>

XLVI. *A Reciprocal Relation in Diffraction.*

By A. A. MICHELSON\*.

SUPPOSE the vibration at the surface  $A$  of a sphere whose centre is at  $O$  to be a known function of  $x$  and  $y$ , the origin being a point  $P$  on the sphere,

$$V = \phi \cos nt + \psi \sin nt. \quad \dots \quad (1)$$

Then the vibration † on a sphere  $B$  passing through  $O$  whose centre is at  $P$  will be

$$W = -\frac{1}{\lambda f} \iint \phi dx dy \sin n(t - \tau) - \frac{1}{\lambda f} \iint \psi dx dy \cos n(t - \tau) \quad (2)$$

Putting

$$\tau = \frac{f}{a} \left( 1 - \frac{x\xi + y\eta}{f^2} \right), \quad t - \frac{f}{a} = t_1, \quad \frac{n\xi}{fa} = u, \quad \frac{n\eta}{fa} = v,$$

$$W = -\frac{1}{\lambda f} \left[ \iint \phi dx dy \cos(ux + vy) - \iint \psi dx dy \sin(ux + vy) \right] \sin nt \left. \vphantom{W} \right\} (3)$$

$$- \frac{1}{\lambda f} \left[ \iint \phi dx dy \sin(ux + vy) + \iint \psi dx dy \cos(ux + vy) \right] \cos nt$$

or

$$W = P \sin nt_1 + Q \cos nt_1. \quad \dots \quad (4)$$

If now a spherical mirror be made to coincide with the sphere  $B$  an image of the source will be formed at  $A$ .

This image may also be considered as the resultant of the vibrations at  $B$ . Hence, if we designate by  $DW$  the

\* Communicated by the Author.

† Scientific Papers of Lord Rayleigh, vol. iii. p. 80. The results given by Lord Rayleigh for the *intensity* of the diffraction figure in the focal plane do not apply to the phase of the vibration. This restriction is removed if the surface considered be the sphere  $B$ ; for the distance between two points on the spheres is

$$\rho^2 = (x - \xi)^2 + (y - \eta)^2 + (z - \zeta)^2,$$

$$\text{or if } f^2 = x^2 + y^2 + z^2 = \xi^2 + \eta^2 + (\zeta - f)^2$$

$$\rho^2 = f^2 - 2x\xi - 2y\eta - 2z\zeta + 2f\zeta.$$

If  $\xi$  and  $\eta$  are small,  $\zeta$  will be of the second order, and so is  $f - z$ , so that  $(z - f)\zeta$  is of the fourth order and may be neglected. Hence

$$\rho^2 = f^2 - 2x\xi - 2y\eta$$

or approximately

$$\rho = f - \frac{x\xi + y\eta}{f}.$$

operation by which  $W$  is obtained from  $V$ ,

$$\left. \begin{aligned} W &= DV \\ V &= DW \end{aligned} \right\} \dots \dots \dots (5)$$

or  $V = DDV. \dots \dots \dots (6)$

These equations hold for light of any colour, and—with corresponding extension of the definition of  $D$ —to any combination of colours.

The analogy with Fourier's formula is apparent. In fact, the second of equations (5) is, putting  $n/fa = p$ ,

$$V = -\frac{1}{\lambda f} \left[ \iint Q d\xi d\eta \cos p(x_1\xi + y_1\eta) - \iint P d\xi d\eta \sin p(x_1\xi + y_1\eta) \right] \cos nt_2$$

$$- \frac{1}{\lambda f} \left[ \iint Q d\xi d\eta \sin p(x_1\xi + y_1\eta) + \iint P d\xi d\eta \cos p(x_1\xi + y_1\eta) \right] \sin nt_2$$

Substituting the values of  $P$  and  $Q$  from (4),

$$V = \frac{1}{\lambda^2 f^2} \left\{ \begin{aligned} &\iiint dx dy d\xi d\eta \psi(x,y) \sin [(x_1-x)p\xi + (y_1-y)p\eta] \\ &+ \frac{1}{\lambda^2 f^2} \iiint dx dy d\xi d\eta \phi(x,y) \cos [(x_1-x)p\xi + (y_1-y)p\eta] \end{aligned} \right\} \cos nt_2$$

$$+ \frac{1}{\lambda^2 f^2} \left\{ \begin{aligned} &\iiint dx dy d\xi d\eta \phi(x,y) \sin [(x_1-x)p\xi + (y_1-y)p\eta] \\ &+ \frac{1}{\lambda^2 f^2} \iiint dx dy d\xi d\eta \psi(x,y) \cos [(x_1-x)p\xi + (y_1-y)p\eta] \end{aligned} \right\} \sin nt_2$$

The first and third integrals are identically zero, so that putting

$$u = \frac{n\xi}{fa} \quad \text{and} \quad v = \frac{n\eta}{fa}$$

and disregarding the phase difference between  $t$  and  $t_2$ , we have by (1)

$$\frac{4\pi^2}{f^2} \phi(x_1, y_1) = \iiint dudv dx dy \phi(x,y) \cos [(x_1-x)u + (y_1-y)v]$$

This, disregarding the intensity factor  $f^2$ , is the Fourier formula extended to two dimensions.

Formulæ (5) express the fact that if

$W$  is the diffraction image of  $V$ , then

$V$  is the diffraction image of  $W$ .

In applying the formulæ it must be remembered that  $V$  and  $W$  represent the vibration—not merely the intensity.