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2 **Diurnal Variation of Tropical Ice Cloud Microphysics:**
3 **Evidence from Global Precipitation Measurement Microwave Imager**
4 **(GPM-GMI) Polarimetric Measurements**

5
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15 **Key points:**

16
17 1. The first study of the diurnal cycle of tropical ice cloud microphysics using spaceborne high-

18 frequency microwave polarimetric measurements.

19 2. The polarimetric difference's (PD's) diurnal cycle precedes that of the ice cloud occurring

20 frequency and total mass by about 2 hours.

21 3. The diurnal cycle of PD can be explained by changes of the ice crystal axial ratio using

22 radiative transfer calculations.

23

24 **Abstract**

25 The diurnal variation of tropical ice clouds has been well observed and examined in terms of the
26 occurring frequency and total mass but rarely from the viewpoint of ice microphysical
27 parameters. It accounts for a large portion of uncertainties in evaluating ice clouds' role on
28 global radiation and hydrological budgets. Owing to the advantage of precession orbit design and
29 paired polarized observations at a high-frequency microwave band that is particularly sensitive to
30 ice particle microphysical properties, three years of polarimetric difference (PD) measurements
31 using the 166 GHz channel of Global Precipitation Measurement Microwave Imager (GPM-GMI)
32 are compiled to reveal a strong diurnal cycle over tropical land (30°S–30°N) with peak
33 amplitude varying up to 38%. Since the PD signal is dominantly determined by ice crystal size,
34 shape, and orientation, the diurnal cycle observed by GMI can be used to infer changes in ice
35 crystal properties. Moreover, PD change is found to lead the diurnal changes of ice cloud
36 occurring frequency and total ice mass by about 2 hours, which strongly implies that
37 understanding ice microphysics is critical to predict, infer, and model ice cloud evolution and
38 precipitation processes.

39

40 **Plain Language Summary:**

41 Along with the rising and setting of the Sun, our Earth's clouds vary between day and night. For
42 tropical ice clouds, there are many evidences showing that the cloud cover and thickness change
43 dramatically during a day (referred to as the "diurnal cycle") but few hints of what happens to
44 individual ice crystals. Using a new spaceborne satellite instrument called GPM-GMI, we found
45 that ice crystal size or shape also exhibits a strong diurnal cycle over tropical land but not over

46 tropical ocean. We can further infer that the microscopic variation of ice crystals may be an
47 important reason for the change of ice cloud cover and thickness later on.

48

49 **1. Introduction**

50 Upper tropospheric ice clouds, namely anvils and cirrus, spread to cover vast areas and persist
51 for hours to days. They are a major modulator of Earth's radiation and thus play an important
52 role in weather and climate changes (e.g., Hartmann et al., 1984; Raymond and Zeng, 2000;
53 Waliser et al., 2009). Studies have shown that the cloud radiative effect (CRE) of the clouds
54 strongly depends on physical details, such as cloud top height (Kiehl, 1994), thickness (Hong et
55 al., 2016; Hartmann and Berry, 2017), overlaying situation (Hartmann et al., 2001), and
56 microphysical properties (Liou et al., 2002; Zeng et al., 2009a, b; Tang et al., 2017). Fu and Liou
57 (1993), for example, showed that the radiative heating rate for a layer of ice cloud with a fixed
58 ice water path could differ by a factor of 10 when the mean effective radius of the ice particles
59 varies by a factor of 5. Reducing the uncertainty of CRE requires a better understanding of ice
60 cloud microphysical properties, which is essential not only to remote sensing of the bulk optical
61 properties but also to the simulation of CRE (Tang et al., 2017). In addition, ice clouds,
62 especially the anvils and mixed-phase clouds in deep convections, contribute much to tropical
63 precipitation and vertical circulations (e.g., Raymond and Zeng, 2000; Liu and Zipser, 2008).
64 Since ice clouds play a critical role in the global water and energy cycles, it is imperative to
65 determine their microphysical processes as well as their macro-characteristics.

66 One of the large variabilities of tropical ice clouds is their strong response to the diurnal
67 variation of solar heating at the Earth's surface (Yang and Slingo, 2001; Tian et al., 2004; Hong
68 et al., 2006). The diurnal variations of ice clouds have been observed using various cloud

69 properties such as occurring frequency, total mass, and cloud cover with a variety of measuring
70 techniques (e.g., Tian et al., 2004; Liu and Zipser, 2008; Eriksson et al., 2010; Millan et al., 2013;
71 Jiang et al., 2014) in conjunction with some other key atmospheric variables such as water vapor
72 (Soden, 2000), surface temperature (Dai et al., 2004) and convective activity (Chung et al., 2007).
73 These metrics have been extensively studied over the past two decades especially owing to the
74 rapid advance in spaceborne instrument technology, as well as the increase in a satellite's spatial
75 and temporal coverage. However, microphysical properties of ice crystals, such as crystal shape,
76 orientation, and size distribution, have rarely been investigated, since it is a challenge to observe
77 ice crystal microphysical properties accurately from space. Some of the integrated properties
78 (e.g., effective diameter) are derived products from primary properties, such as mass, coverage,
79 and type, in most satellite observations (e.g., CloudSat) with arbitrary assumptions. Although
80 new retrieval techniques such as one-dimensional variation method (1DVAR) and neural
81 network can retrieve all desired variables together at least theoretically, oversimplified
82 microphysical assumptions are still required in the retrieval model or training database, which
83 directly impact their outputs (e.g., Boukabara et al., 2011; Duncan and Kummerow, 2016).
84 Hence, how to observe and analyze the diurnal variation of cloud ice microphysics from space is
85 inherently a difficult yet interesting topic.

86 Polarimetric difference (PD) between paired vertically and horizontally polarized (V-pol and
87 H-pol hereafter for abbreviation) channels at a fixed incidence angle is a relatively matured
88 product and can be measured precisely from space. Recently, the polarimetric data have been
89 used to study ice microphysical properties (e.g., Skofronick-Jackson et al. 2015; Homeyer and
90 Kumjian, 2015; Hioki et al., 2016). In particular, studies of PD from passive and active
91 spaceborne measurements at visible and microwave wavelengths show that nonspherical ice

92 crystals are not randomly oriented but usually horizontally oriented (Prigent et al., 2001, 2005;
93 Davis et al., 2005, 2007; Zhou et al., 2012; Defer et al., 2014). Gong and Wu (2017) (GW17
94 hereafter) first identified the similar features from the GPM-GMI 166 and 89 GHz observations.
95 Taking advantage of GPM's precession orbit, this paper aims to better understand the diurnal
96 variation of PDs in the tropics and its implication for ice crystal's microphysical properties and
97 how they vary together with macroscopic quantities such as cloud occurring frequency and mass.

98

99 **2. Diurnal Variations of 166 GHz Polarimetric Radiance Difference**

100 The GPM core satellite has a dual-frequency precipitation radar (DPR) and a passive GMI
101 instrument with 13 channels between 10 and 190 GHz (Hou et al. 2014; Skofronick-Jackson et al.
102 2015). Among the 13 channels, the 10.65, 18.7, 36.5, 89, and 166 GHz channels have V-pol and
103 H-pol. In this study, only the 166 GHz polarization data are used because this channel is
104 sensitive to cloud ice scattering and its PD is not sensitive to water cloud emission. Furthermore,
105 the PD at this channel is least contaminated by surface polarizations in the tropics (GW17). The
106 GPM core satellite flies at the altitude of 407 km in a non-Sun-synchronized orbit that covers a
107 latitude range between 68°S and 68°N. It slowly progresses over the local solar time (LST),
108 making it ideal to study the diurnal variability of clouds and precipitation. The GMI has a
109 forward conical scan viewing geometry off-nadir at an angle of 48.5° (see Skofronick-Jackson et
110 al. 2015 for details).

111 The current study uses Version 1.4 of Level-1C inter-calibrated GMI brightness temperature
112 (TB) observations at 166 GHz during February 2014 – December 2016 (GPM X-Cal Working
113 Group, 2015). The diurnal cycle is not separated further by seasons, as the seasonality in the
114 tropics between 30°S and 30°N is small. Furthermore, since it takes approximately 45 days for

115 the GPM satellite to preprocess through all LSTs, integration over three years on a time interval of 2
116 hours can generate robust and useful cloud statistics. The GMI 183±3 GHz radiance is used to
117 screen the “confident cloudy-sky” scenes for estimating the cloud occurring frequency (OF) and
118 the mean cloudy-sky TB (TB_{cld}). The colder TB_{cld} is, the thicker the ice cloud is. As at 166 GHz,
119 $-TB_{\text{cld}}$ is roughly linearly correlated with the column-integrated ice water path (IWP) as long as
120 the cloud is not completely opaque (Gong and Wu, 2014). Therefore, $-TB_{\text{cld}}$ can be used as a
121 proxy for IWP along the GMI’s line of sight. Both 166 and 183±3 GHz channels belong to the
122 high-frequency group that scans simultaneously with the same footprint size of 4.2 X 7.0 km².
123 The “3 σ cloud detection” threshold is employed to distinguish between clear- and cloudy-sky
124 scenes, which is identical to Gong and Wu [2014; 2017]. This method first identifies the peak
125 (TB_{peak}) and standard deviation (σ) of the probability density function (PDF) of the TB
126 observations in a given region or a latitude belt over the period of interest, and then the cloudy-
127 sky scene is defined at places where $TB < TB_{\text{peak}} - 3\sigma$. In practice, this procedure is iterated twice
128 to remove the broadening effect introduced by clouds (i.e., bias in σ). Since a large amount of
129 “possible ice cloud present” scenes with relatively warmer TB is excluded using this method, OF
130 is much smaller and TB_{cld} is cold-biased. However, to fully describe the ice cloud polarization
131 feature, it is necessary to account for the polarization signal corresponding to these “warmer ice
132 cloud” scenes that are in general populated with cirrus and thin anvil clouds. Hence, another
133 threshold is proposed in the next paragraph to separate the clear-sky and cloudy-sky polarimetric
134 signatures.

135 Three years of the 166 GHz observations in the tropics (30°S-30°N) are sorted to the nearest
136 LST slot that is separated by a 2-hour interval. Then, the 2D PDF integration is performed for
137 each of the 12 LSTs according to its V-pol TB (TB_V) and PD values. PD is defined as

138
$$PD \equiv TB_V - TB_H \quad (\text{Eqn 1})$$

139 A positive PD implies that the upwelling microwave radiation must have a shorter cloud optical
140 length in V-pol (τ_V) than in H-pol (τ_H) due to either scattering or absorption along the path and
141 vice versa. Once the PDF is compiled, the mean and standard deviation of the PD value within
142 each TB_V bin is then calculated to draw the final curves shown in Fig. 1a for the oceanic and Fig.
143 1b for the landmasses, respectively. The cloudy-sky thresholds are drawn in the black dash lines.
144 Following GW17, the bottom-up bell curve is referred to as the “PD-TB” relationship. The same
145 as found in GW17, the TB_V value corresponded to the peak PD robustly located at approximately
146 200 K, and the broadness of the spread is strikingly similar across different LSTs. These features
147 again reveal an apparently nearly universal bell-curve in the PD-TB relationship. These curves
148 spread out at the peak PD but converge to roughly the same TB values as PD values approach
149 zero. At the warm TB end when the curves converge, another threshold (dash-dot lines) is set, to
150 the left of which is considered a full cloudy-sky PD-TB relationship. Since the dynamic range of
151 TB_V for the PD-TB relationship holds the same during a day, the curvature of the curve is solely
152 determined by its peak value of PD (PD_{peak}).

153 The PD_{peak} from tropical clouds exhibits a clear diurnal cycle over land but not over ocean.
154 The amplitude of PD_{peak} variation over land is $\sim 38\%$ and that over ocean is only $\sim 6\%$.
155 Furthermore, the maximum and minimum of PD_{peak} occur at local morning (9:00 LST) and
156 afternoon (16:00 LST) over land, while the difference is not statistically significant over ocean.
157 The general conclusion of local morning maximum and afternoon minimum of PD_{peak} over land
158 is robust against using different time intervals (1 hr and 3 hrs), although the exact LST may vary
159 by ± 1 hr. As will be shown in the next section, PD_{peak} , or the curvature of the curve, is the key
160 variable to correspond to the column-integrated optical axial ratio.

161 The diurnal variation of PD_{peak} is better exhibited in Fig. 2a (ocean) and 2b (land) where
162 PD_{peak} clearly shows an anticorrelation with cloud occurrence frequency and mean cloud TB.
163 Compared to the diurnal cycle of upper-tropospheric clouds previously identified using Tropical
164 Rainfall Measurement Mission (TRMM)'s precipitation radar and visible and infrared scanner,
165 the diurnal cycle of PD_{peak} also anticorrelates with that of the area and cloud fraction for both
166 ocean and land surface conditions. In those analyses, over ocean, the diurnal cycle of PD_{peak}
167 precedes that of the precipitation rate (PR) by ~ 2 hrs, but for land cases, it lags the PR diurnal
168 cycle by ~ 2 hrs (reference: Fig. 7 of Hong et al., 2006; Fig. 3 of Liu and Zipser, 2008).

169 To showcase the direct correlation of the diurnal cycles between PD_{peak} and the macro-
170 quantities of the same group of clouds, Fig. 2 also displays the time series of PD_{peak} , OF, and
171 TB_{cloud} . As mentioned previously, to compute the latter two variables, all cloudy-sky scenes with
172 TB below the " 3σ thresholds," shown as the dashed lines in Fig. 1, are used here. In addition, to
173 remove some casual noise produced by only taking PD_{peak} values, PD from a relatively small TB
174 window ($180K < TB_V < 210K$) is averaged (referred to as $\overline{PD_{peak}}$), showing that the major
175 conclusions hold robustly against using 1 hr time interval, as well as against variations of the
176 thresholds (e.g., using dash or dash-dot lines in Fig. 1 to compute OF and TB_{cloud} or using slightly
177 different windows to compute $\overline{PD_{peak}}$).

178 Similar to the findings from TRMM data, the diurnal cycle of $\overline{PD_{peak}}$ is the opposite of
179 OF and total mass ($-TB_{cloud}$), while the latter two go together tightly. However, the largest
180 correlation coefficient is found when $\overline{PD_{peak}}$ leads the other two quantities by ~ 2 hours, when
181 the correlation coefficient exceeds -0.9. That is to say, when $\overline{PD_{peak}}$ reaches its minimum
182 (maximum) during the day, ice cloud coverage and total ice cloud mass reach their maximum
183 (minimum) within ~ 2 hrs. This is the first study implying that ice crystal microphysics may play

184 a key role on the formation/dissipation processes of ice clouds. Furthermore, this leading time of
185 2 hrs is robust against using a time interval of 1 hr, but it would be “distorted” to 3 hrs if we use
186 3 hrs as the time interval for the analysis. Such a sensitivity exercise emphasizes the importance
187 of having enough samples and small enough time steps in order to investigate the diurnal
188 variation.

189 The only exception is between $\overline{PD_{peak}}$ and OF over ocean where the correlation is the
190 largest without 5 hour advancement. This result changes with choosing different time interval so
191 it’s not a robust result. Since the diurnal cycle of $\overline{PD_{peak}}$ is too small over ocean to beat its own
192 variability (error bars on Fig. 2a), the correlation among these three cloud parameters are not
193 statistically significant over ocean anyway.

194 In summary, a strong diurnal cycle of GMI’s 166 GHz PD is found for tropical ice clouds
195 over land and a much weaker diurnal cycle over the ocean. The diurnal variations of PD lead
196 those of the total ice cloud coverage and mass by ~ 2 hrs.

197

198 **3. Interpretation and Radiative Transfer Model Simulations**

199 As the observed PD is dominantly positive, by definition, ice particles in the ice cloud layer
200 must possess anisotropic characteristics in terms of the bulk microphysical properties along the
201 line of sight. In other words, the ice particles within this ice cloud layer must be aspherical
202 microphysically and, more importantly, also aspherical on an integrated sense. Previous studies
203 hence proposed that the majority of aspherical ice crystals should have oriented preferably along
204 the horizontal direction rather than randomly oriented (Prigent et al., 2001, 2005; Davis et al.,
205 2005, 2007; Zhou et al., 2012, 2013; Defer et al., 2014). GW17 too one step further that
206 explained the entire bottom-up bell curve feature of the PD-TB relationship with a conceptual

207 model. In this model, when a cloud is optically thick, both V- and H-pol radiations are saturated
208 at the same level, causing PD to diminish. By applying this hypothesis to a radiative transfer
209 model (RTM), GW17 simulated the observed PD-TB relationship for 89, 166, and 640 GHz at
210 the same time by controlling one parameter, called the “optical aspect ratio” (AR), defined as
211 $AR \equiv \tau_V / \tau_H$. The ratio is not the typical “axial ratio” parameter often used to describe the
212 eccentricity of an aspherical ice particle; it is the column integrated optical property difference
213 between V-pol and H-pol at a given frequency, in the present case, 166 GHz. Further, by
214 adjusting the value of AR, GW17 found the curvature of the PD-TB relationship changes, but the
215 warm and cold ends remain converging to 0 (as expected theoretically) at roughly the same TB_V
216 values. In other words, PD_{peak} is directly tied to the value of AR.

217 Following on the GW17 approach, the observed diurnal variation of the PD-TB relationship
218 is simulated using different AR values in the same RTM as detailed in GW17. Briefly speaking
219 about the setup, two parallel simulations are run at the same time; one with a regular extinction
220 profile and the other profile is arbitrarily multiplied by a fixed AR value. Gamma distribution is
221 employed for the particle size distribution, with the effective diameter fixed at 160 μm . By
222 varying the cloud optical depth to create the dynamic range of TB, the difference between the
223 two parallel runs creates the “PD-TB” bell curve observed (c.f., Section 4.2 of GW17). In this
224 study, AR is changed from 1.0 to 2.0 with increments of 0.02 for each set of simulation. The
225 “best-fit” AR value is then defined such that the standard deviation between the simulated PD
226 and the observed counterpart reaches minimum (see Fig. 3a). In this way, the AR value is
227 “retrieved” for different LSTs and its time series are displayed in Fig. 3b and 3c over the ocean
228 and land surfaces, respectively. This AR “retrieval” is technically similar to a standard retrieval
229 process except for no actually retrieving a physical quantity with a set of rigorous algorithms.

230 As expected, the variation of the “best-fit” AR is tightly correlated with the diurnal variation
231 of $\overline{PD_{peak}}$. When the diurnal variation amplitude is small compared to the internal variability at a
232 given LST, the “best-fit” AR is insensitive to the small change of $\overline{PD_{peak}}$. That’s the case over
233 tropical ocean. On the other hand, when the tropical land convection and the associated anvil and
234 cirrus clouds start to experience the most dramatic diurnal change from noon to late night, the
235 “best-fit” AR value also changes accordingly.

236 The value of AR physically corresponds to an integrated result from a few ice microphysical
237 parameters that impact τ_V/τ_H by definition. $\tau = \int_{z_{base}}^{z_{top}} \int_{D_{min}}^{D_{max}} N(D)A(D)dD dz$, where N is the
238 number density, A is the cross section of the particle size perpendicular to the incoming radiation,
239 D is the diameter, and z_{base} and z_{top} are the cloud base and top heights, respectively (Heymsfield
240 et al., 2003). With V-pol and H-pol detecting radiation from the surface that penetrates up
241 through the same ice cloud layer, only the aspherical particle shape and diameter at two specific
242 directions that are perpendicular to V-pol and H-pol, assuming that the 100% horizontal-
243 alignment is valid. Under this assumption, large (small) AR corresponds to high (low) axial ratio
244 of the ice particles. That means over land, the majority of ice particles tend to become more and
245 more sphere-like from morning (9 am) to afternoon (4 pm), and from then on horizontally
246 oriented non-spherical ice crystals start to dominate the PD signal again.

247 This proposed diurnal variation of ice crystal orientation is consistent with the physics model
248 of radiation-induced ice crystal growth (or the radiative effect on microphysics [(e.g., Zeng
249 2008])). In the physics model, horizontally oriented ice crystals near the cloud top, in contrast to
250 the vertically oriented ones, emit longwave radiation effectively to space, subsequently have low
251 surface temperature, and thus can grow to precipitation-sized particles, which explains the
252 positive value of the PD observed. Meanwhile, ice crystals near the cloud top absorb solar

253 radiation during daytime and consequently the radiation-induced crystal growth is modulated by
254 solar radiation, which explains the diurnal variation of the PD shown in Figs. 1 and 2 (Zeng et al.,
255 2016). The model also explains the delay of peak of ice cloud mass and coverage, since
256 horizontally oriented ice crystals with larger PD (i.e., flatter) tend to quickly grow into
257 precipitation-sized ones and fall out as precipitation, while ice crystals that are more sphere-like
258 grow slower and stay as ice cloud for a longer time.

259 As for the relationship of AR to particle size, GW17 found its relationship to AR is quite
260 complicated and nonlinear using a fully polarized RTM simulation (see Fig. 9 in GW17). GW17
261 found that the ice crystal with “column” shape overall matched the best with the observations,
262 and the curvature of the PD-TB curve rises with the effective radius, But for certain other shapes,
263 such as rosette with 7 branches, simulation results do not show the clear linear relationship
264 between PD_{peak} and effective radius. Therefore, one cannot rush to the conclusion that the
265 particle size growth exhibits a diurnal cycle. Other factors might also contribute to the diurnal
266 variation of PD but are likely secondary (e.g., liquid cloud that can damp the PD and TB signals
267 also exhibit its own diurnal cycle).

268 The 100% horizontal alignment is unlikely to occur in reality, and the mixing of different ice
269 particle shapes may also vary diurnally. The better-mixed the shapes are (i.e., more percentage of
270 random-orientation), the smaller the AR value is because AR is expected to vanish for the
271 extreme situation of 100% random orientation. If the percentage of randomly oriented particles is
272 the only or major contributor to the diurnal variation of PD_{peak} , then it indicates that along with
273 the growth of deep convective systems from late morning to the early evening, ice particles also
274 tend to become more and more randomly oriented. This makes sense as the vertical motion in
275 deep convective cores tends to generate and maintain a more turbulent environment than that in

276 the ambient environment of the anvils and cirrus. However, this could not explain the ~ 2 hrs
277 delay of the peak/trough of the diurnal cycles of ice cloud mass and coverage. Also, as indicated
278 in the Appendix A, the PD_{peak} information indeed mostly comes from the differentiation of τ_V
279 and τ_H for anvils, instead of from deep convective clouds as 166 GHz quickly saturates at the top
280 several kilometers of the deep convective core. Therefore, the degree of turbulence, or
281 randomness, inside the deep convective cloud is not a major contributor to the value of AR or
282 PD_{peak} .

283

284 **4. Conclusions**

285 High-frequency passive microwave measurements, such as GMI 166 GHz channel, are
286 sensitive to ice cloud scattering. Vertically and horizontally polarized microwave observations
287 are found to differ in cloudy-sky because of the disparity of their optical thicknesses through an
288 ice cloud layer with anisotropic bulk microphysical properties. As a result, the polarimetric
289 radiance difference, or PD, can be used to infer cloud microphysical properties (i.e., ice crystal
290 shape, size, and orientation).

291 The GPM precession orbit provides a valuable chance to study the diurnal variations of PD at
292 166 GHz in the tropics for the first time. The PD-TB relationship bears with a bottom-up bell
293 curve that was investigated extensively in an earlier study of Gong and Wu [2017]. All bell
294 curves can be represented by the peak PD amplitude, called PD_{peak} , which serves as a good
295 measure to characterize the diurnal variation of the PD-TB curve. Similar to other macrophysical
296 quantities like the area of coverage, occurring frequency, and total mass of ice clouds, PD_{peak} is
297 found to vary by only 6% over tropical ocean, but vary dramatically ($\sim 38\%$) over tropical land.

298 The maximum and minimum of PD_{peak} occur at local morning and afternoon over tropical land,
299 which lead the diurnal cycle of ice cloud mass and coverage by ~ 2 hrs with an opposite sign.

300 The observed diurnal variation of the PD-TB relationship can be explained quantitatively
301 using a set of different AR values, because AR reflects a bulk effect of ice crystal shape and size
302 in radiative transfer, and the diurnal variation in PD is linked directly to a variation in AR. Other
303 microphysical properties, however, may also play a role in altering the PD-TB relationship as
304 observed and need to be explored with more observations and accurate model representations of
305 ice microphysical properties. Since radiation impacts ice crystal growth via crystal surface
306 temperature (e.g., Zeng, 2008; Zeng et al., 2016), it can subsequently change AR, which in turn
307 changes PD. Although we cannot fully understand the delay and anticorrelation of cloud mass
308 and coverage with respect to PD, the delayed occurrence of ice cloud macro-physical properties
309 indicates the importance of ice microphysics in formation/decay of ice clouds. It is imperative to
310 include these processes in models to replicate the realistic cloud diurnal cycle and its radiative
311 and hydrological effects.

312 This study exemplifies the power of passive microwave polarimetry in remotely sensing
313 ice cloud microphysical properties. The 166 GHz is the highest frequency of polarized channels
314 in GMI. Because ice cloud scattering is more significant at high frequencies, careful fusion of
315 observations from multi-frequency polarized microwave/IR channels makes it feasible to infer
316 both bulk ice particle size and shape simultaneously and allows global studies of ice cloud
317 microphysics on diurnal, seasonal, and interannual time scales. Although it has been well known
318 that ice microphysics are strongly tied together with the evolution of ice cloud macro-properties,
319 this work identifies that understanding the variations in ice microphysics are necessary to predict,
320 infer, and model the bulk properties of the ice clouds and their broader evolution (e.g.,

321 precipitation process). This study is a good exercise to manifest that new satellite
322 instrumentation (e.g., GMI) with novel observational techniques (e.g., paired high-frequency
323 microwave channels) may offer a valuable means to observe and understand the entire cloud-
324 precipitation processes.

325

326 **Appendix A: Where does the PD_{peak} come from?**

327 In Gong and Wu (2017), the bottom-up bell curves of the PD-TB relationship from 89, 166,
328 and 640 GHz observations are explained by the different penetration depth of V-pol and H-pol
329 measurements through an ice cloud layer consisting of horizontally oriented nonspherical ice
330 crystals. In this appendix, we present additional evidence to support this explanation. More
331 importantly, this exercise answers directly to the types of ice cloud, their vertical levels, and how
332 they contribute to the largest PD signal (PD_{peak}).

333 Two years (2014–2015) of collocated CloudSat-GMI data are compiled together to composite
334 the Contoured Frequency by Altitude Diagram (CFAD) in Fig. A1b for the clouds that satisfy the
335 criteria such that $7 \text{ K} < PD < 20 \text{ K}$ and $160 \text{ K} < TB_v < 220 \text{ K}$ (enclosed by the red rectangle box
336 in Fig. A1a). This dataset is produced by Dr. Joe Turk at the Jet Propulsion Lab and details about
337 this dataset can be found in Turk [2017]. To generate robust statistics, we employed two years of
338 data since collocated GPM-CloudSat scenes are sparse. For CloudSat, 0 dBZ roughly separates
339 the thin ice clouds (anvils, cirrus) with low reflectivity from thick ice clouds (thick anvils, deep
340 convective cores) with high reflectivity.

341 As shown in CFADs for V-pol and H-pol, the largest disparity comes from clouds with
342 reflectivity below 0 dBZ and height above 10 km. These clouds are apparently thin anvils (that at
343 least 166 GHz V-pol can partially penetrate through) and cirrus clouds. When the ice cloud

344 becomes thicker and lower, 166 GHz passive microwave observations quickly saturates, and
345 although 166 GHz are still sensitive to these clouds (down to cloud height as low as 2 km and as
346 strong as 25 dBZ), the PD of 166 GHz no longer shows the disparity.

347 One possible explanation of the curvature of the PD-TB relationship that is proposed in
348 GW17 is that turbulence in deep convective cores likely forces the nonspherical ice crystals to be
349 well mixed (i.e., randomly oriented), in which way the PD eventually vanishes when TB gets
350 cold. This exercise demonstrates that although this mechanism is quite plausible, 166 GHz PD
351 measurements cannot identify the turbulence condition in thick ice clouds and therefore cannot
352 be used to verify such a hypothesized mechanism.

353 **Data Access**

354 GMI Level-1 calibrated radiance data and CloudSat-GPM collocation dataset can be acquired
355 from the Precipitation Processing System's (PPS) FTP server at [arthurhou.pps.eosdis.nasa.](http://arthurhou.pps.eosdis.nasa.gov)
356 [gov](http://arthurhou.pps.eosdis.nasa.gov). The radiative transfer model is coded and modified by the authors for this specific study.
357 Interested readers are encouraged to contact the lead author directly for a copy of the model. All
358 data used to make the figures in this manuscript are included in the supplementary materials. A
359 README file illustrates the file format, content, and variable meanings.

360

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373

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497

498 **Figure Titles:**

499 **Figure 1:** PD-TB_v curves at different local solar times (colors) over tropical ocean (a) and land
500 (b). A curve is a simple connection of the mean values at each TB bin of the 2D PDF integrated
501 from year 2015 of GMI observations at 166 GHz between 30°S and 30°N. The black dash lines
502 indicate the 3σ threshold for computing the cloud occurring frequency and mean TB, and dash-
503 dot lines exhibit the threshold for describing the PD-TB relationship.

504 **Figure 2:** Top: time series of the mean PD value for clouds with TB in the range of [180, 210] K
505 (black), mean cloud occurring frequency (OF; red), and mean TB (TB_{cl*d*}; blue; axis value is
506 upside-down to match the physical meaning of total column ice mass) for clouds below the 3σ
507 threshold (red) at tropical ocean (a) and land (b) during 2015. Bottom: lag-correlation between

508 *the time series of PD and OF (red) and PD and TB_{cld} (blue) over ocean (c) and land (d).*

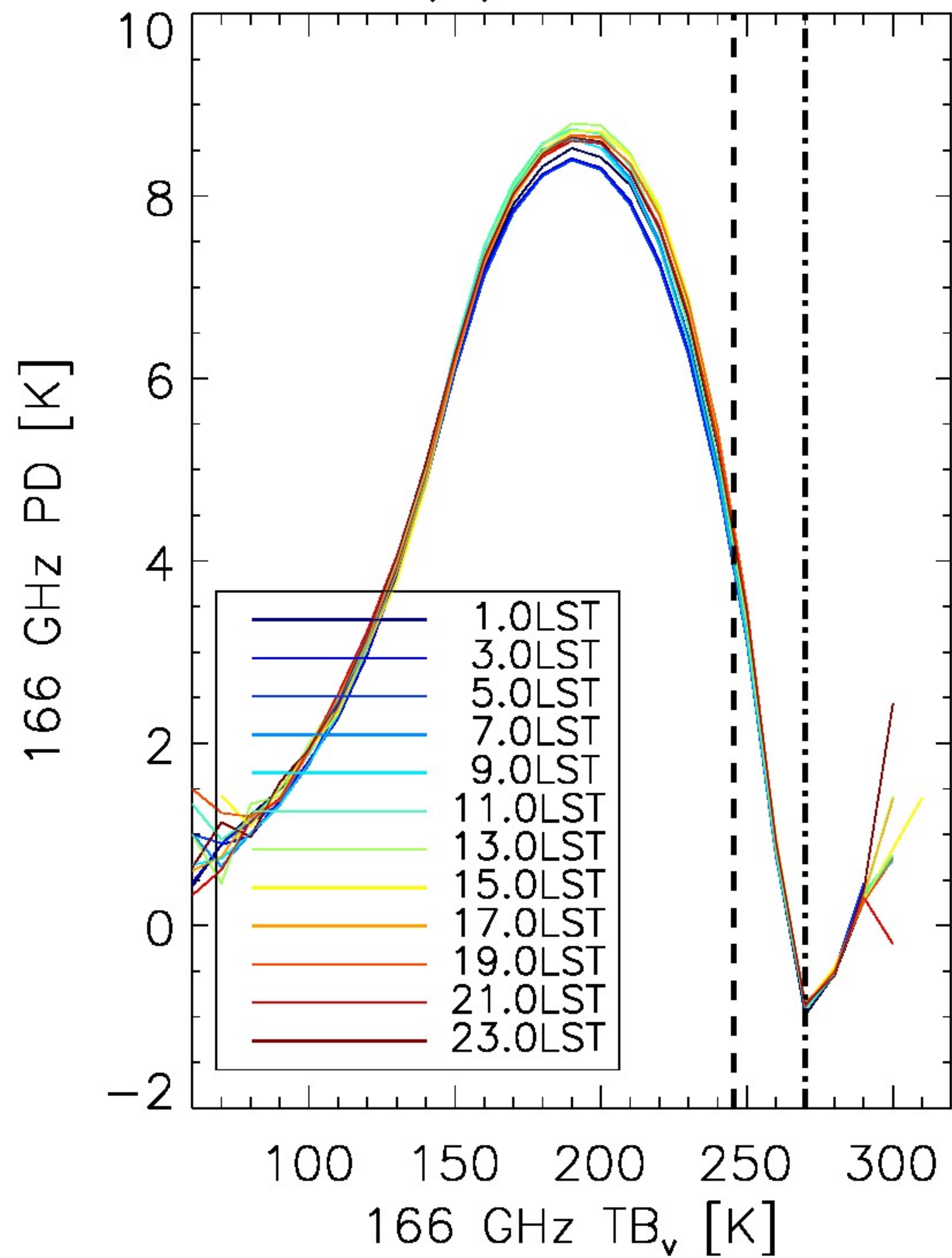
509 *Positive/negative lag time corresponds to a lead/lag of the former time series to the latter ones.*

510 **Figure 3:** *(a) Example of how simulation compares with observations at 166 GHz. The simulated*
511 *166 GHz PD- TB_v relationships are shown in black lines with different AR values (AR varies*
512 *from thin to bold between 1.0 and 1.5 with an increment of 0.04); the top of the simulated results*
513 *shows the mean of the observed ensemble for the entire year of 2015 between the equator and*
514 *10°N with the variability shown in pink envelope (adapted from Fig. 8b of Gong and Wu, 2017).*
515 *The best-fit AR values derived from (a) for each time frame are plotted as stars in (b) and (c) for*
516 *ocean and land situations, respectively. Observations are taken from the curves shown in Fig. 1*
517 *and color legend is the same as Fig. 1. See the context for the definition of “best-fit”.*

518 **Figure A1:** *(a) Two-dimensional Probability Density Function (PDF) of PD (i.e., ΔTB) and TB_v*
519 *derived from GMI 166 GHz oceanic measurements during July 2015 collected between the*
520 *equator and 10°N. This figure is adapted from Fig. 3c of Gong and Wu (2017); (b) CFADs of*
521 *clouds selected from the red rectangle box enclosed area in (a). The black and red contours*
522 *correspond to 166 GHz V-pol and H-pol CFAD, respectively.*

Figure 1.

(a) Ocean



(b) Land

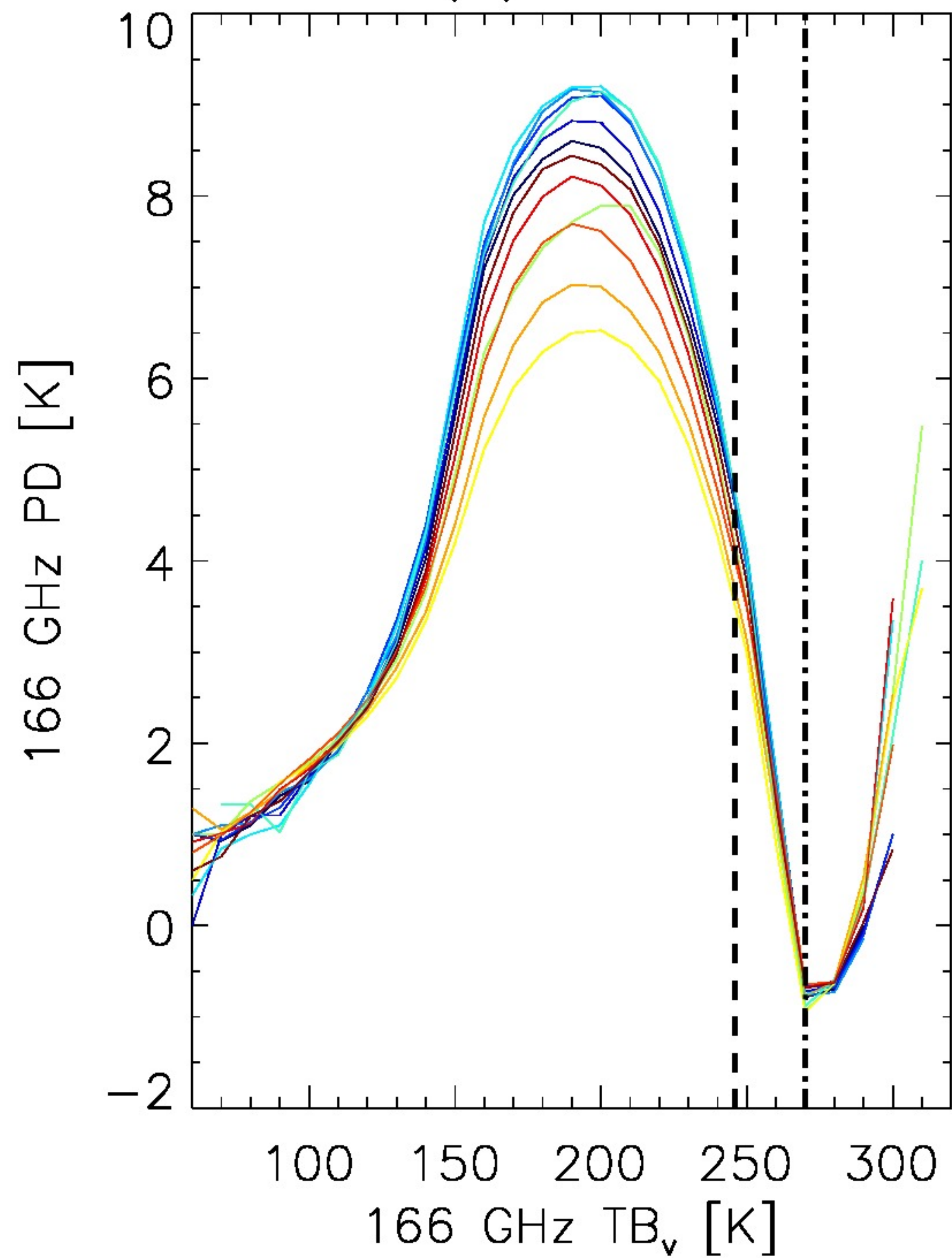
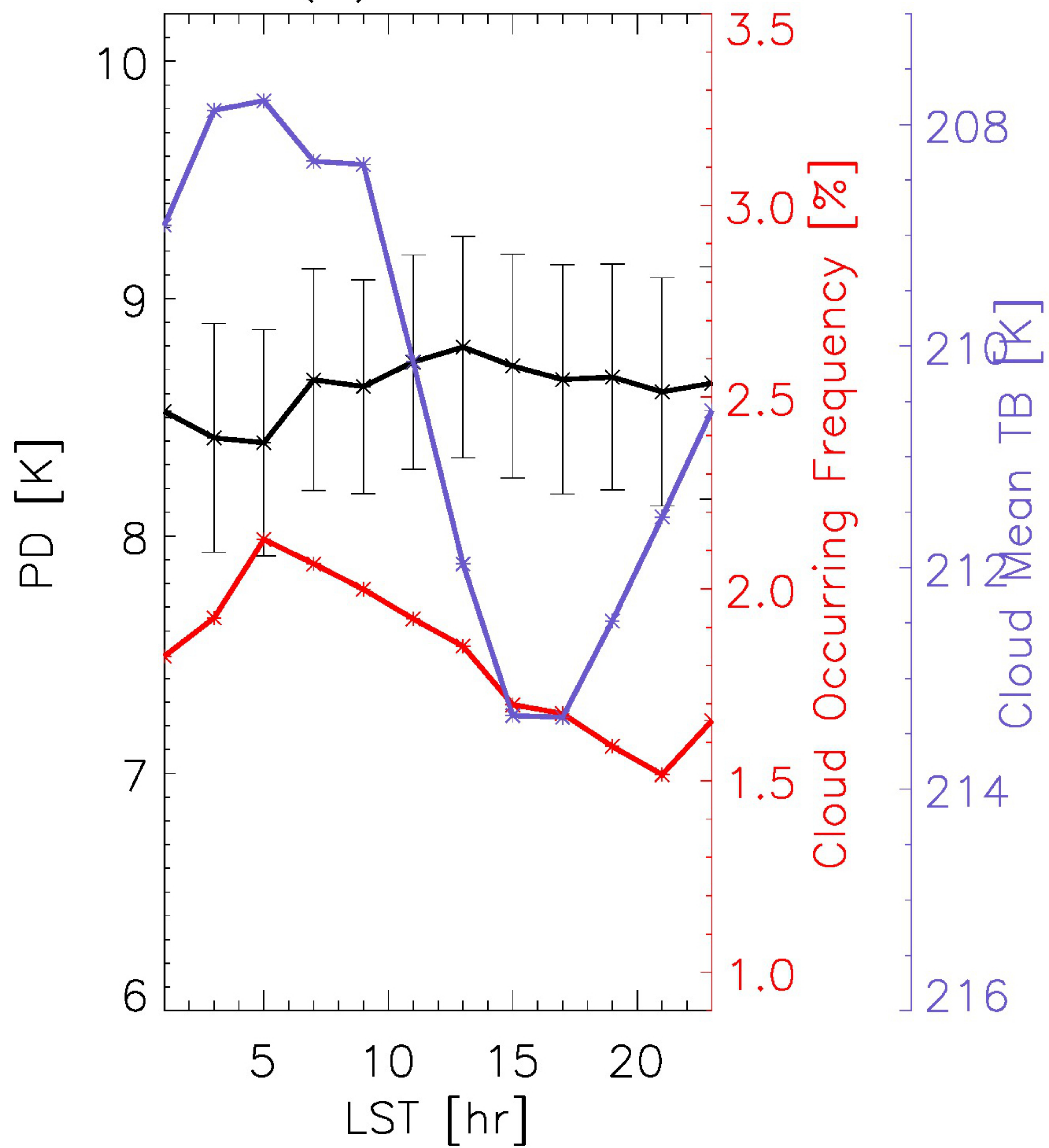
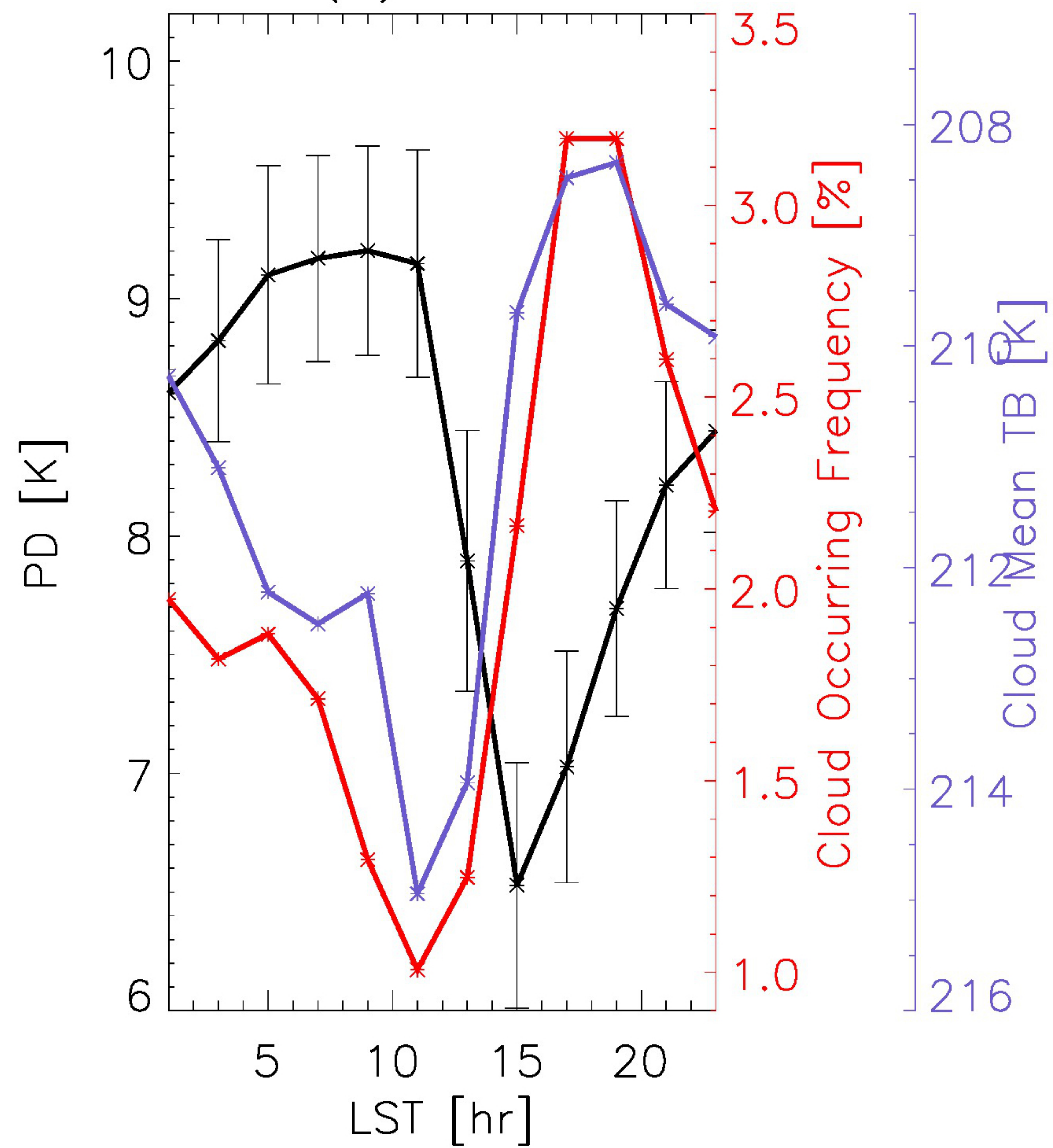


Figure 2.

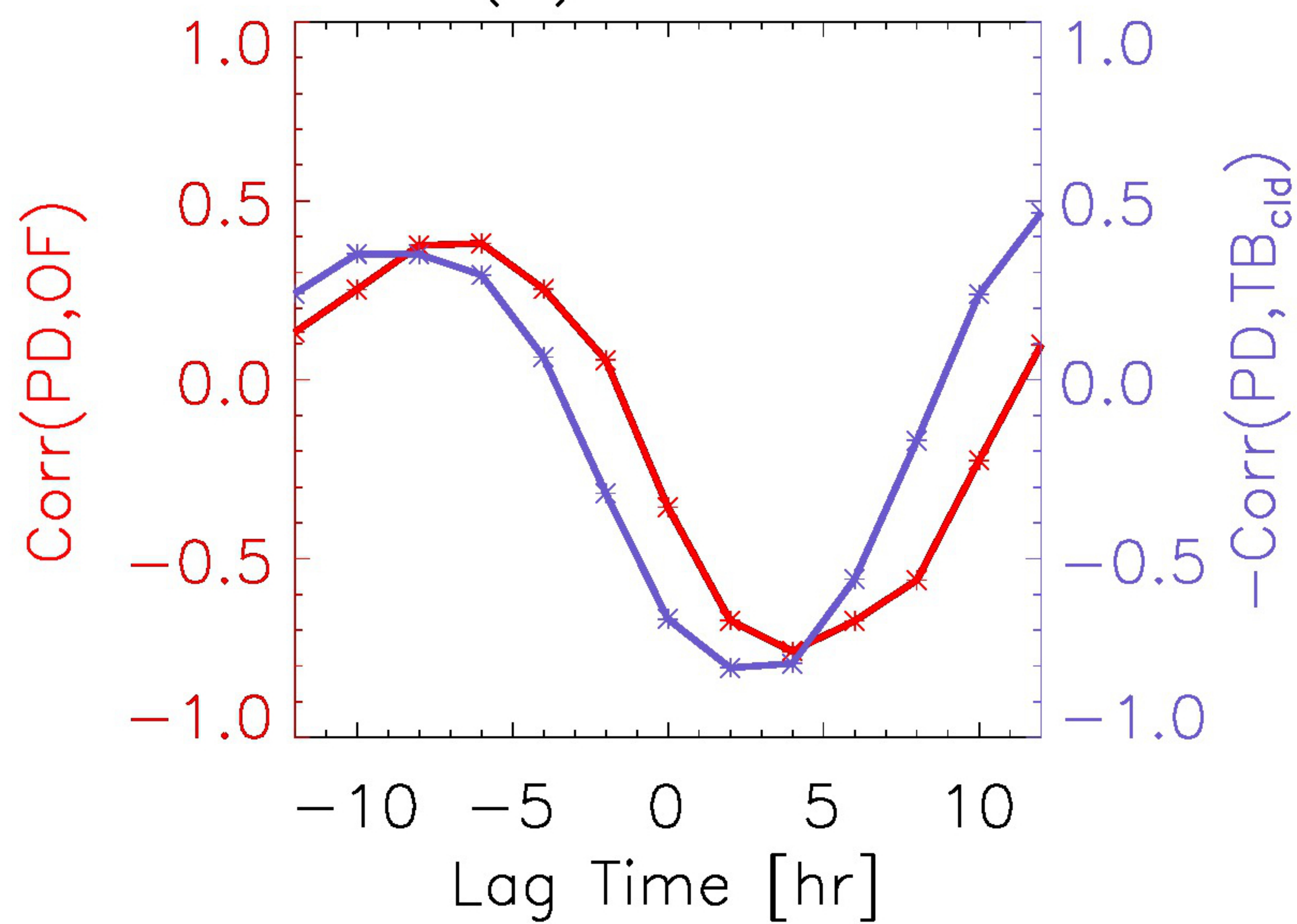
(a) Ocean



(b) Land



(c) Ocean



(d) Land

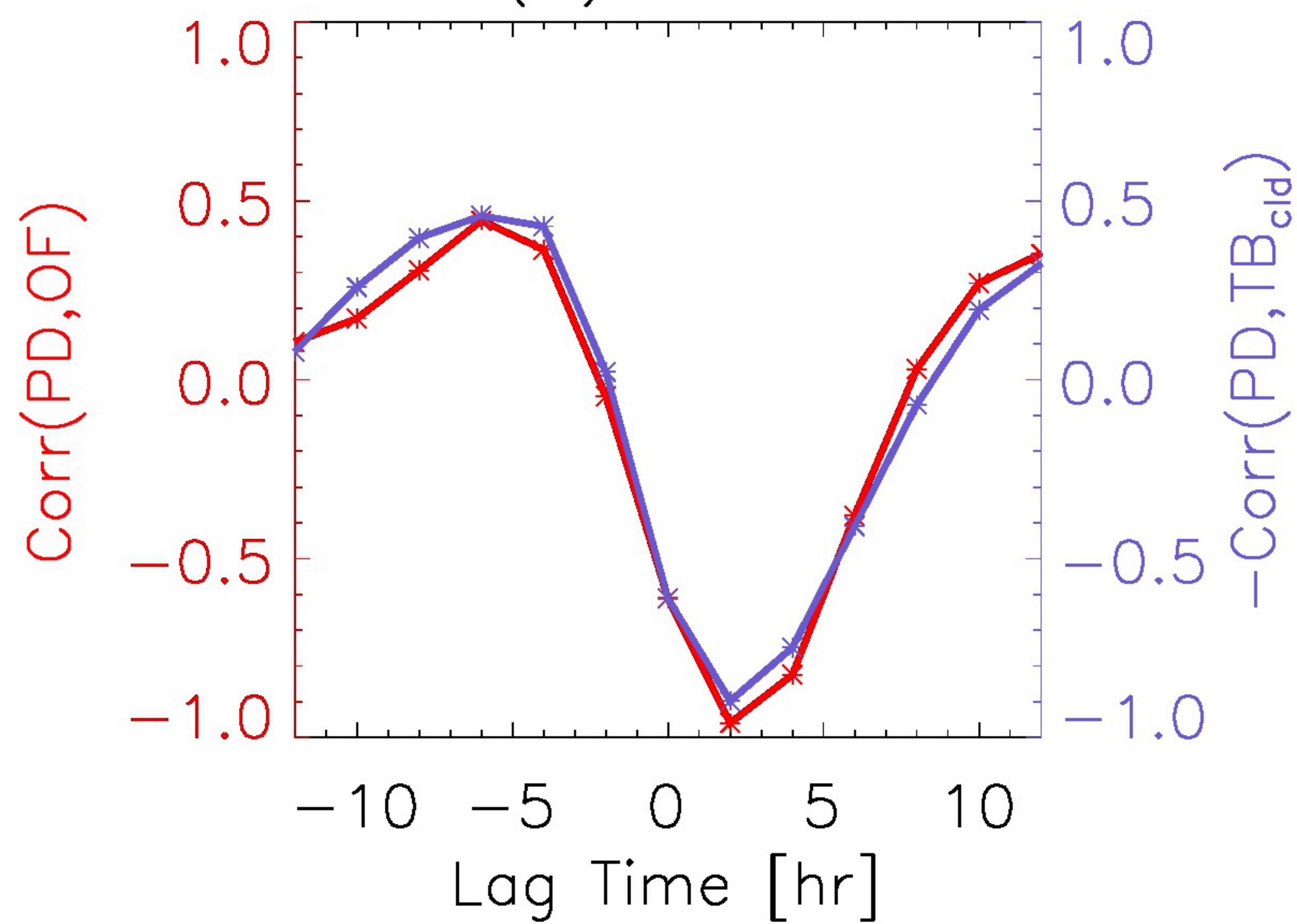
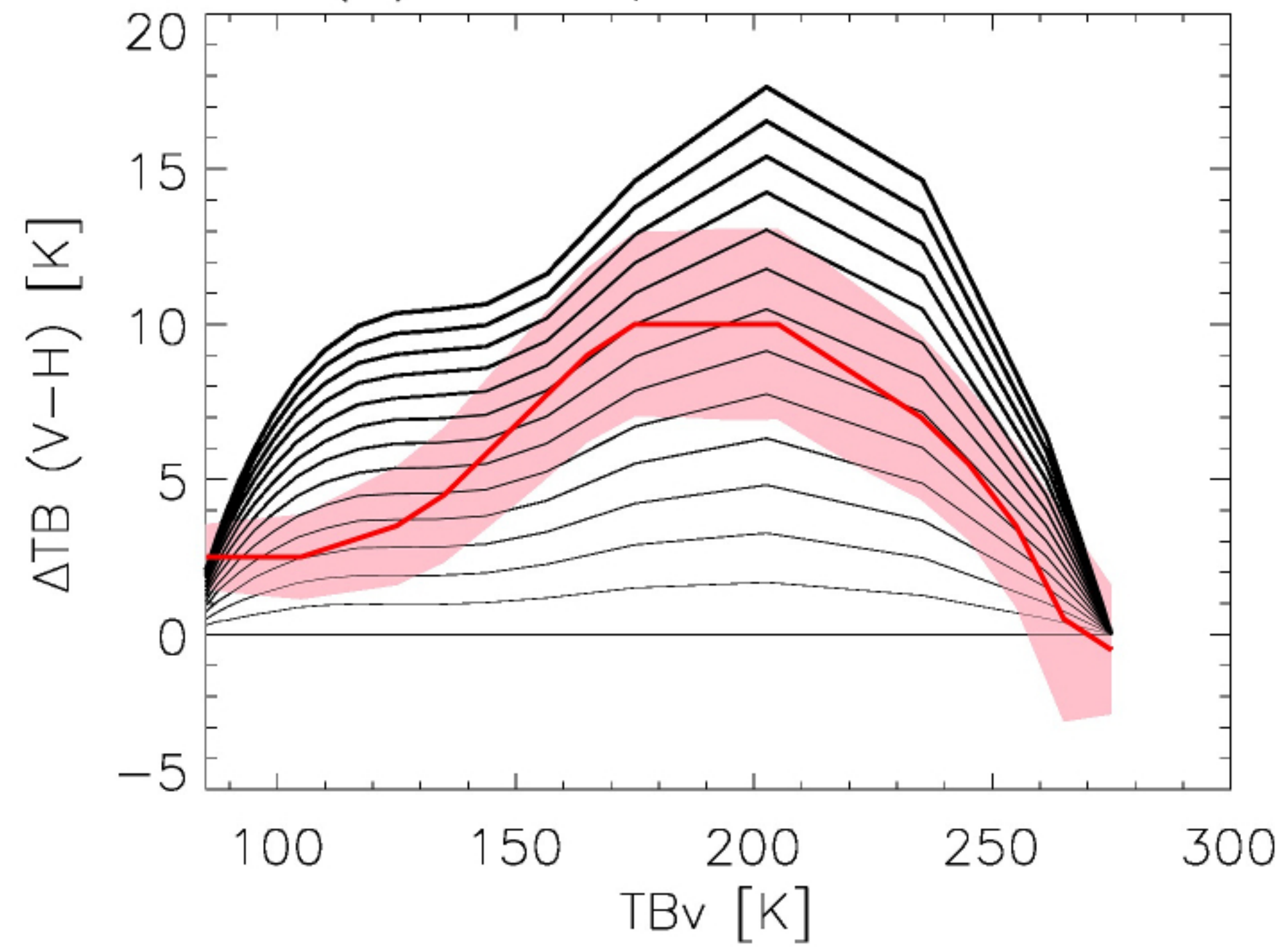
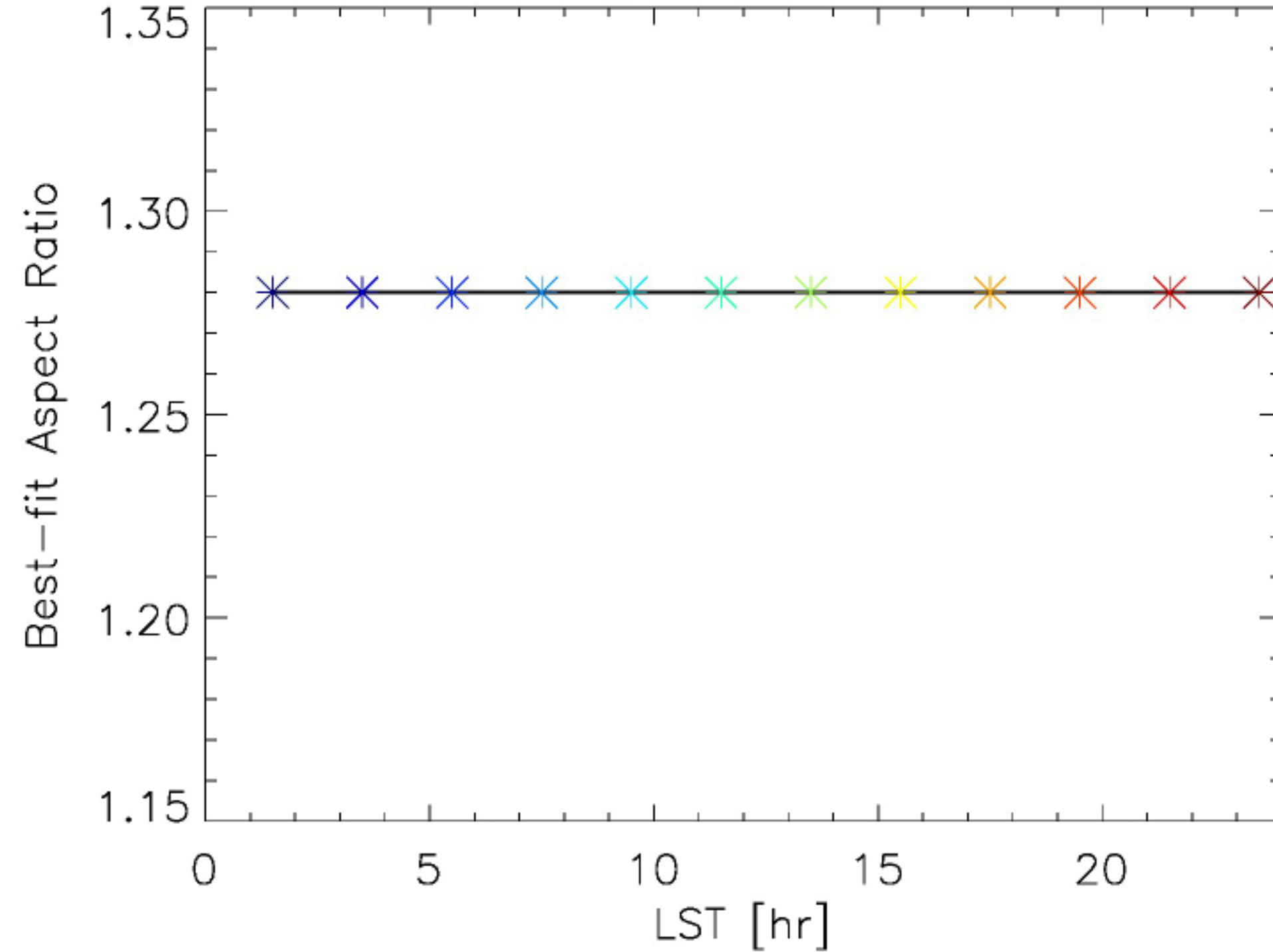


Figure 3.

(a) Example of Best-fit



(b) Ocean



(c) Land

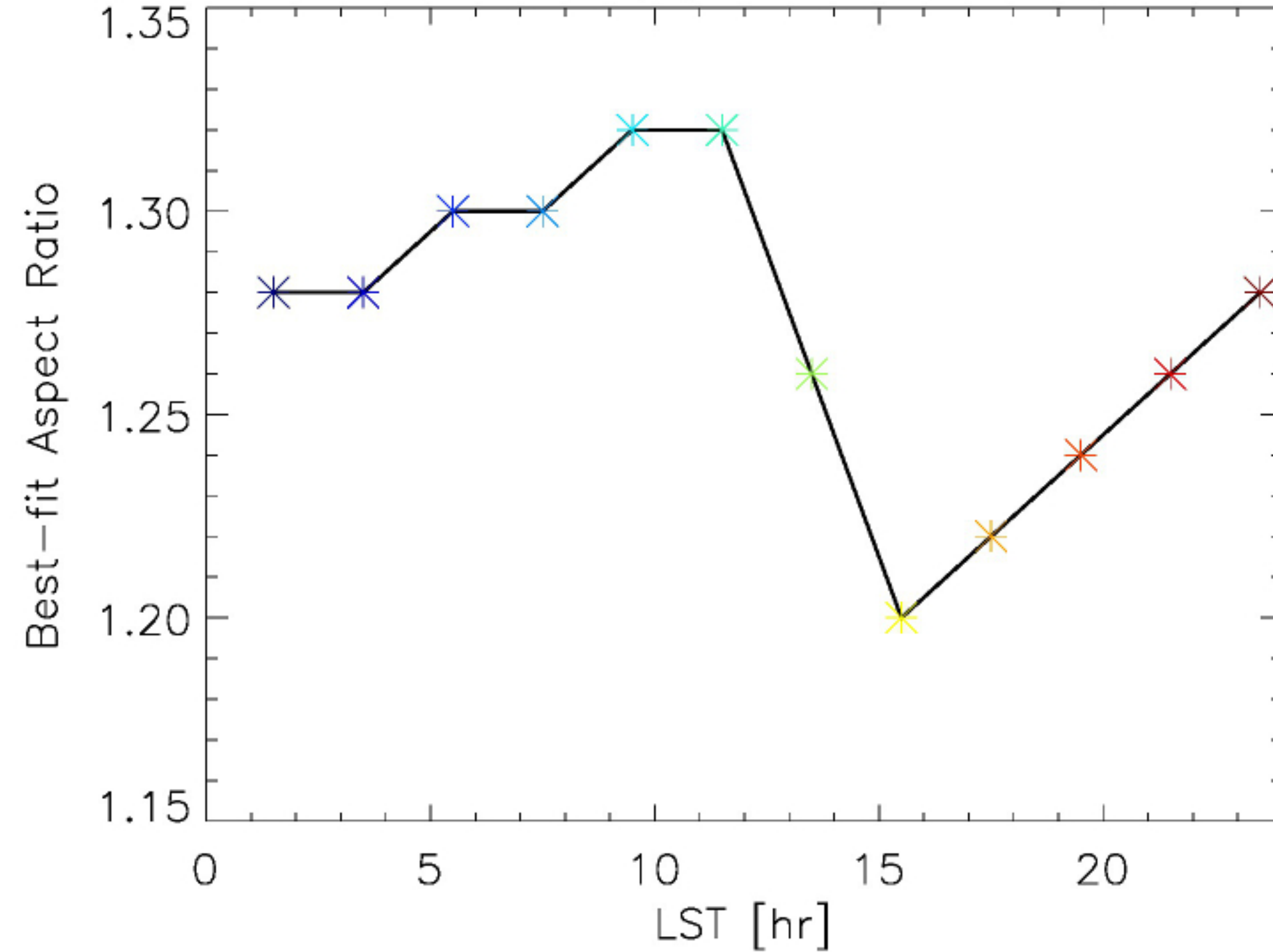
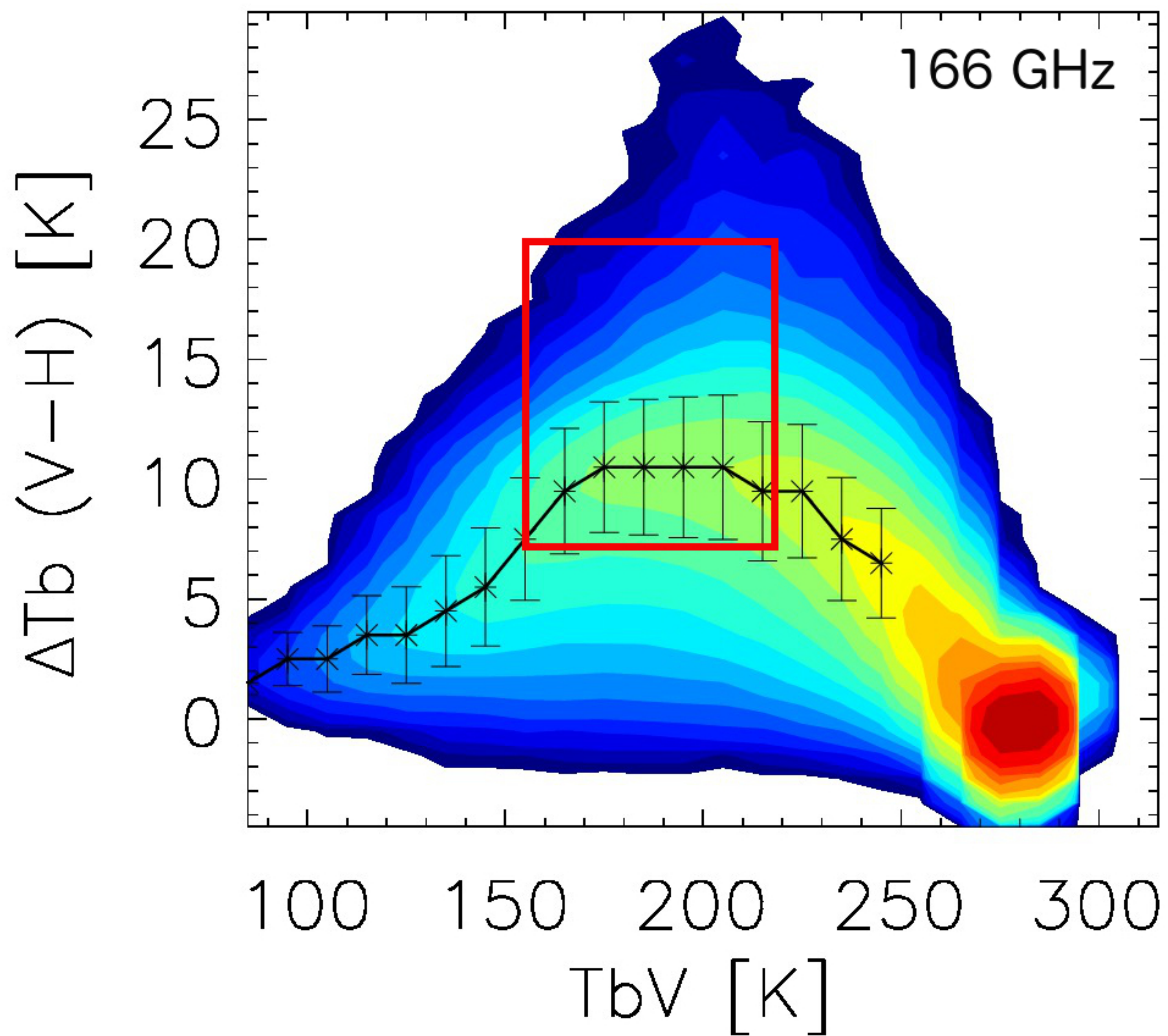


Figure A1.

(a) PDF, Ocean, [0, 10N], July 2015



(b) CFAD, [30S, 30N], 2014-2015

