

1 **Kinematic and Microphysical Significance of Lightning Jumps**  
2 **versus Non-Jump Increases in Total Flash Rate**

3 CHRISTOPHER J. SCHULTZ \*

*Department of Atmospheric Science, The University of Alabama in Huntsville and NASA MSFC, Huntsville, Alabama*

4 LAWRENCE D. CAREY

*Department of Atmospheric Science, The University of Alabama in Huntsville, Huntsville, Alabama*

5 ELISE V. SCHULTZ

*Earth System Science Center, The University of Alabama in Huntsville, Huntsville, Alabama*

6 RICHARD J. BLAKESLEE

*NASA Marshall Space Flight Center, Huntsville, Alabama*

---

\* *Corresponding author address:* Christopher J. Schultz, UAH/NASA Marshall Space Flight Center, 320 Sparkman Drive, Huntsville, AL 35805.  
E-mail: schultz@nsstc.uah.edu

## ABSTRACT

7  
8 Thirty-nine thunderstorms are examined using multiple-Doppler, polarimetric and total  
9 lightning observations to understand the role of mixed phase kinematics and microphysics in  
10 the development of lightning jumps. This sample size is larger than those of previous studies  
11 on this topic. The principal result of this study is that lightning jumps are a result of mixed  
12 phase updraft intensification. Larger increases in intense updraft volume ( $\geq 10 \text{ m s}^{-1}$ ) and  
13 larger changes in peak updraft speed are observed prior to lightning jump occurrence when  
14 compared to other non-jump increases in total flash rate. Wilcoxon-Mann-Whitney Rank  
15 Sum testing yields p-values  $\leq 0.05$ , indicating statistical independence between lightning  
16 jump and non-jump distributions for these two parameters. Similar changes in mixed phase  
17 graupel mass magnitude are observed prior to lightning jumps and non-jump increases in  
18 total flash rate. The p-value for graupel mass change is  $p=0.096$ , so jump and non-jump dis-  
19 tributions for graupel mass change are not found statistically independent using the  $p=0.05$   
20 significance level. Timing of updraft volume, speed and graupel mass increases are found  
21 to be 4 to 13 minutes in advance of lightning jump occurrence. Also, severe storms without  
22 lightning jumps lack robust mixed phase updrafts, demonstrating that mixed phase updrafts  
23 are not always a requirement for severe weather occurrence. Therefore, the results of this  
24 study show that lightning jump occurrences are coincident with larger increases in intense  
25 mixed phase updraft volume and peak updraft speed than smaller non-jump increases in  
26 total flash rate.

27

# 28 1. Introduction

29 Sudden increases in total flash rates are denoted as lightning jumps. Research about  
30 lightning jumps has primarily focused on the correlation between lightning jumps and severe  
31 weather<sup>1</sup> occurrence (e.g., Williams et al. 1999, Schultz et al. 2009, Gatlin and Goodman  
32 2010, Schultz et al. 2011, Rudlosky and Fuelberg 2013). However, these studies lack analysis  
33 of the microphysical and dynamical mechanisms which lead to a rapid increase in total flash  
34 rate.

35 Several studies observed good correlation between total lightning trends and mixed phase  
36 ice mass or updraft volume, but poorer correlation between total lightning and maximum  
37 updraft speed over the entire lifecycle of thunderstorms (e.g., Workman and Reynolds 1949,  
38 Goodman et al. 1988, Tuttle et al. 1989, Dye et al. 1989, Carey and Rutledge 1996, Lang  
39 and Rutledge 2002, Wiens et al. 2005, Tessendorf et al. 2005, Kuhlman et al. 2006, Deier-  
40 ling et al. 2008, Deierling and Petersen 2008). These studies relied on the observed con-  
41 nection between kinematics, microphysics and electrification within thunderstorms via the  
42 non-inductive charging mechanism (e.g., Takahashi 1978, Saunders et al. 2006).

43 Electrification within thunderstorms is found to occur on the order of the quarter to half  
44 life of a ordinary thunderstorm<sup>2</sup>. Research shows that initial electrification in the primary  
45 development of thunderstorms is approximately 10-15 minutes (e.g., Dye et al. 1986, Bringi  
46 et al. 1997). Lightning jumps themselves also occur on timescales which are on the order of  
47 several minutes (Goodman et al. 1988, Williams et al. 1989, Williams et al. 1999). There-  
48 fore, storm properties that are well-correlated to total flash rate on longer timescales may

---

<sup>1</sup>Defined as the presence of hail  $\geq 2.54$  cm, winds  $\geq 26$  m s<sup>-1</sup> or a tornado.

<sup>2</sup>The duration of an ordinary thunderstorm is 30-60 minutes (Byers and Braham 1949).

49 not represent the same mechanisms which result in lightning jumps.

50 Schultz et al. (2015) examined the correspondence between lightning jumps and trends in  
51 mixed phase graupel mass, maximum updraft speed and updraft volume on 15 minute time  
52 scales for 4 thunderstorms of varying morphology. These specific parameters were chosen  
53 because of their strong correlations to total flash rate from studies mentioned previously  
54 in this paper. Schultz et al. (2015) showed that lightning jumps occur when the  $10 \text{ m s}^{-1}$   
55 updraft volume and mixed phase graupel mass increase prior to jump occurrence. They  
56 also determined that maximum updraft speed increases in 8 of the 12 flash rate periods  
57 examined. However, Schultz et al. (2015) did not robustly demonstrate how the kinematic  
58 and microphysical mechanisms examined differ between lightning jumps and other non-jump  
59 increases in total flash rate because of a small sample size.

60 Therefore, the goal of this research is to determine whether there are statistically signif-  
61 icant differences between lightning jumps and non-jump increases in total flash rate using  
62 a large sample of thunderstorm observations. Analysis of the kinematic and microphys-  
63 ical characteristics will assess if larger changes in the magnitude of mixed phase graupel  
64 mass, updraft volume or maximum updraft speed occur prior to lightning jumps versus  
65 other non-jump increases in total flash rate. This analysis will also evaluate the tempo-  
66 ral correspondence between the  $2\sigma$  lightning jump algorithm (Schultz et al. 2009, Schultz  
67 et al. 2011, Schultz et al. 2015) and the underlying kinematic and microphysical thunder-  
68 storm characteristics needed for rapid electrification. The  $2\sigma$  algorithm is currently being  
69 used experimentally at the National Oceanic and Atmospheric Administration's Hazardous  
70 Weather Testbed (NOAA HWT; Calhoun 2015) in preparation for the launch of GOES-R's  
71 Geostationary Lightning Mapper (GLM; Goodman et al. 2013).

## 73 2. Data and Methods

74 The data, study domain<sup>3</sup> and analysis methods are similar to those of Schultz et al.  
75 (2015) for continuity between the results of that study and the present study. The focus  
76 is using total lightning, polarimetric and multi-Doppler data and analysis to characterize  
77 kinematic and microphysical changes within a thunderstorm prior to any increase in total  
78 flash rate. This research provides more comprehensive statistical metrics related to the phys-  
79 ical mechanisms hypothesized to modulate electrification and lightning production within  
80 thunderstorms.

81 A total of 39 thunderstorms are used in this analysis. Convective intensity of the thun-  
82 derstorms examined ranges from weak ordinary multicellular convection and low-topped  
83 winter convection to bowing segments within quasi-linear convective systems (QLCS) and  
84 supercells (Table 1). Of the 39 thunderstorms, 20 thunderstorms contain at least 1 lightning  
85 jump, and 19 possess zero lightning jumps while they are within the multi-Doppler domain.  
86 Twenty-three of the 39 thunderstorms are multicellular thunderstorms, 10 thunderstorms  
87 are supercells, 3 are low topped supercell storms, 2 are bowing segments within QLCSs and  
88 1 storm is in the outer bands of a remnant tropical cyclone. In total, 214-fifteen minute  
89 analysis periods prior to an increase in total flash rate (both jumps and non-jump increases)  
90 are analyzed from these 39 thunderstorms. Properties which are examined in this analysis  
91 include: mixed phase graupel mass (-10° to -40°C), 5 and 10 m s<sup>-1</sup> mixed phase updraft

---

<sup>3</sup>See Fig. 1 of Schultz et al. (2015).

92 volume, and maximum updraft speed.

93

94 *a. Radar Data*

95 The same radar data and methodologies used in Schultz et al. (2015) are employed in this  
96 study. The University of Alabama in Huntsville’s (UAH) Advanced Radar for Operational  
97 Research (ARMOR; Schultz et al. 2012, Knupp et al. 2014) and the National Weather  
98 Service’s (NWS) radar located at Hytop, AL (KHTX; Crum and Albery 1993) are used  
99 for three-dimensional retrieval of velocity and bulk characterization of hydrometeor types  
100 within thunderstorms. ARMOR can be taken out of its default 5 tilt operational mode to  
101 collect higher temporal resolution and larger volumetric data for research purposes. All  
102 radar data are corrected for attenuation and differential attenuation (Bringi et al. 2001).  
103 Aliased velocities are unfolded using NCAR’s SOLO software (Oye et al. 1995) and ground  
104 clutter, side lobe and second trip echoes are also removed from all radar data. Data are  
105 gridded to a Cartesian coordinate system using a grid spacing of 1 km x 1 km x 1 km on a  
106 grid of 300 km x 300 km x 19 km. This spacing is chosen because of the resolution limita-  
107 tions of the longer baseline used for the ARMOR-KHTX domain (e.g., Davies-Jones 1979,  
108 Deierling and Petersen 2008). A Cressman weighting scheme is implemented using 1 km  
109 radius of influence centered at each grid point with NCAR’s REORDER software (Oye and  
110 Case 1995). Individual thunderstorms are identified and semi-objectively tracked using the  
111 Thunderstorm Identification Tracking Analysis and Nowcasting (TITAN; Dixon and Wiener  
112 1993) algorithm to assign radar and lightning characteristics to individual storms. This

113 tracking method is the same used in previous lightning jump studies (Schultz et al. 2009,  
114 Schultz et al. 2011, Schultz et al. 2015).

115 The National Center for Atmospheric Research’s (NCAR) Custom Editing and Display  
116 of Reduced Information in Cartesian Space (CEDRIC; Mohr et al. 1986) is used to per-  
117 form multi-Doppler synthesis. Vertical velocity retrievals are calculated using radial velocity  
118 measurements from two or more radars and a reflectivity based hydrometeor fall speed re-  
119 lationship to solve a set of linear equations (e.g., Armijo 1969, O’Brien 1970, Brandes 1977,  
120 Ray et al. 1980, Deierling and Petersen 2008, Schultz et al. 2015). Horizontal velocity com-  
121 ponents  $u$  and  $v$  derived from radial velocity measurements from both radars and are used  
122 to solve for the vertical velocity component ( $w$ ) by integrating the anelastic continuity equa-  
123 tion.

124 Similar to Schultz et al. (2015), the variational integration technique is utilized in this  
125 study to evaluate trends in updraft within thunderstorms (e.g., O’Brien 1970, Matejka and  
126 Bartels 1998, their Section 2e). The variational technique is chosen for this analysis for con-  
127 tinuity between the methods used in this study and other studies using the ARMOR-KHTX  
128 baseline (e.g., Deierling and Petersen 2008, Johnson 2009, Mecikalski et al. 2015, Carey et al.  
129 2016). The advantage of the variational integration technique is that it redistributes errors  
130 from both boundary conditions to produce profiles of vertical air motion and divergence that  
131 converge to a solution (O’Brien 1970, Matejka and Bartels 1998). The downward integration  
132 scheme could also be utilized for similar analysis of updrafts.

133 Vertical velocity is set at  $0 \text{ m s}^{-1}$  at the upper and lower bounds of integration (0 and  
134 17 km). Integration of the anelastic mass continuity equation is performed from the upper  
135 and lower bounds of integration for all points within the multi-Doppler domain. Upward

136 integration is performed from 0 km up to 3 km and downward integration is performed from  
137 the upper boundary from 17 km down to (and including) 3 km. Upward integration is only  
138 used below 3 km because of potential errors introduced into the calculation of divergence  
139 and vertical velocity at low-levels from radar beam height limitations (i.e., radar data do  
140 not extend all of the way to the surface). However, updraft information used for analysis  
141 is limited to the  $-10^{\circ}$  to  $-40^{\circ}\text{C}$ , which is  $\geq 4$  km in the 39 cases examined and these levels  
142 utilize calculations from the downward integration of continuity. Integration of the anelastic  
143 continuity equation results in an estimate of vertical velocity for each  $1 \text{ km}^3$  volume where  
144 u, v and divergence are calculated in the vertical column.

145 Analysis of updraft speed and volume are limited to the mixed phase region of the thun-  
146 derstorm (i.e., between  $-10^{\circ}\text{C}$  and  $-40^{\circ}\text{C}$  isotherms) because the mixed phase region is where  
147 charge development and separation take place to ultimately lead to electrical breakdown  
148 (e.g., Dye et al. 1986, Carey and Rutledge 1996, Bringi et al. 1997, Deierling and Petersen  
149 2008, Calhoun et al. 2013). Maximum updraft speed and a sum of  $1 \text{ km}^3$  updraft volumes  
150 with speeds  $\geq 5 \text{ m s}^{-1}$  and  $10 \text{ m s}^{-1}$  are computed from the multi-Doppler Cartesian grids  
151 for all multi-Doppler syntheses in which a thunderstorm is identified and tracked by TITAN.  
152 The longer baseline between the radars means that updraft values calculated in the study are  
153 smaller in magnitude than the true updraft observed if higher resolution observations were  
154 available. However, the trend in the updraft can still be characterized, especially with hori-  
155 zontal resolution  $\leq 1.5 \text{ km}$  in the domain used in this study. This longer baseline approach  
156 is used in similar lightning/updraft studies like Deierling and Petersen (2008), Mecikalski  
157 et al. (2015) and Carey et al. (2016).

158 Particle identification is performed using ARMOR radar data and the NCAR Particle



159 Identification Algorithm (PID; Vivekanandan et al. 1999) modified for C-band observations  
 160 (Deierling et al. 2008, Johnson 2009, Schultz et al. 2015) to identify the dominant scatterer  
 161 observed in each ARMOR radar volume. Graupel/small hail category is the primary hydrom-  
 162 eteor of interest in this study because of graupel’s strong tie to electrification and lightning  
 163 production in thunderstorms through NIC processes (e.g., Carey and Rutledge 1996, Saun-  
 164 ders et al. 2006, Deierling et al. 2008). Graupel mass is calculated using a z-M relationship of,  
 165

$$mass \ (g \ m^{-3}) = 0.0052 \times z^{0.5}, \quad (1)$$

166 from Heymsfield and Miller (1988). The letter z represents the reflectivity factor in linear  
 167 units of  $mm^6 \ m^{-3}$ , and this calculation is made for each volume where graupel/small hail  
 168 is identified to be the dominant particle in a  $1 \ km^3$  volume. All  $1 \ km^3$  between  $-10^\circ C$  and  
 169  $-40^\circ C$  that also fell inside the thunderstorm’s TITAN footprint are then used to calculate a  
 170 total mass for the storm at each ARMOR radar volume time.

171

172 *b. Lightning Data*

173 The same lightning data and methods used in Schultz et al. (2015) are employed in this  
 174 study. Total lightning information is collected by the North Alabama Lightning Mapping  
 175 Array (NALMA, Koshak et al. 2004, Goodman et al. 2005). Very high frequency (VHF)  
 176 source points are combined into corresponding flashes using a flash clustering algorithm de-  
 177 veloped by McCaul et al. (2009). This cluster algorithm requires that all VHF source points  
 178 0.3 s apart in time and that satisfy an azimuth and range dependent spatial separation re-

179 striction are grouped into a single lightning flash<sup>4</sup>. A flash must have a minimum of 10 VHF  
180 source points to be considered in this analysis.

181

## 182 1) THE LIGHTNING JUMP

183 The sigma-level configuration of the  $2\sigma$  lightning jump algorithm is used to categorize  
184 jump and non-jump increases in total flash rate within this study (Schultz et al. 2009, Schultz  
185 et al. 2011). Sigma-level is represented by,

186

$$\text{sigma-level} = \frac{DFRDT_{t_o}}{\sigma(DFRDT_{t-2,t-4,t-6,t-8,t-10})}, \quad (2)$$

187 where  $DFRDT_{t_o}$  represents the time rate of change of the total flash rate at the current time,  
188 and  $\sigma(DFRDT_{t-2,t-4,t-6,t-8,t-10})$  represents the standard deviation of time rate of change for  
189 the previous 12 minutes of lightning data starting at  $t - 2$ . Please see Appendix A or Schultz  
190 et al. (2011), Chronis et al. (2015) and Schultz et al. (2015) for more detail on the calculation  
191 of the  $2\sigma$  lightning jump algorithm. All increases in total flash rate (i.e., positive sigma-level)  
192 are examined, and flash rate increases are binned into two groups by their sigma-level. A  
193 non-jump increase in flash rate has a sigma-level  $<2$  (hereafter defined as the 0-2 category)  
194 and a lightning jump has a sigma-level  $\geq 2$  (hereafter defined as the 2+ category).

195

---

<sup>4</sup>For more information on the spatial requirements, see McCaul et al. (2009) and references within the article.

196 *c. Analysis Windows*

197 Trends in updraft speed, updraft volume and graupel mass are determined in the follow-  
198 ing manner. First, the time of the flash rate increase ( $t_o$ ) is used to identify radar volumes  
199 within  $\pm 2$  minutes of occurrence. If two radar volumes are available, the one closest to  
200 the time of the flash rate increase is used. Similarly, the closest radar volume to the time  
201 15 minutes prior to the flash rate increase (i.e.,  $t-15$ ) is also identified. This radar volume  
202 also must occur within  $\pm 2$  minutes of the  $t-15$  time. Next the local trend in each radar  
203 derived parameter is determined by subtracting the value at time  $t-15$  from time  $t_o$ . The  
204 magnitude of the change is placed into the corresponding sigma-level category. A 15 minute  
205 analysis window is chosen because it is on the order of the quarter to half life of an ordinary  
206 thunderstorm (Byers and Braham 1949), its the approximate amount of time for the onset  
207 of electrification (Dye et al. 1986, Bringi et al. 1997) and this period allows for 2-3 radar  
208 updates from the WSR-88D radars to obtain trends in other intensity metrics like maximum  
209 expected size of hail (MESH; Witt et al. 1998) or azimuthal shear. Lengthening the analysis  
210 window could also incorporate data which is less likely to be attributed to the development  
211 of a lightning jump (Schultz et al. 2015).

212

213 *d. Statistical Significance*

214 Assessment of statistical independence between the jump and non-jump distributions is  
215 made for each kinematic or microphysical quantity in this study (i.e., mixed phase updraft  
216 volume, updraft speed or graupel mass). The Wilcoxon-Mann-Whitney Rank-Sum Test is

217 used to determine the degree of independence between the 0-2 and 2+ sigma-level data dis-  
218 tributions for updraft volume, updraft speed and graupel mass (Wilks 1995, pp. 159-163).  
219 The use of the rank sum test is ideal for this dataset because the sampling distribution of the  
220 data is unknown and this test is resistant to any potential outliers. Z-scores and p-values for  
221 each of the comparisons are presented to illustrate the level of significance between the 0-2  
222 and 2+ sigma-level categories. The null hypothesis is that the 0-2 and 2+ sigma-level are  
223 drawn from the same distribution for each parameter examined in this study. Thus, if the  
224 p-value is  $p \leq 0.05$ , the null hypothesis is rejected and the property is more likely observed  
225 with lightning jump occurrence than a general increase in flash rate. If the p-value is  $p > 0.05$ ,  
226 the null hypothesis is supported and the kinematic/microphysical property is observed for  
227 any increase in total flash rate and not solely for lightning jumps.

228

### 229 **3. Results**

230 Parameters of mixed phase graupel mass, updraft volume and updraft speed are examined  
231 to determine differences in the kinematic and microphysical growth within a thunderstorm  
232 prior to lightning jumps and non-jump increases in total flash rate. Changes in these quan-  
233 tities in a 15 minute analysis window will help determine the degree to which well-correlated  
234 parameters observed in previous studies can differentiate between lightning jumps and non-  
235 jump increases in total flash rate. Ultimately, lightning jumps could be used to infer a higher  
236 likelihood that a specific physical process is present if a lightning jump is observed, especially  
237 for physical parameters which are not readily available (e.g., updraft speed, updraft volume).

239 *a. Mixed Phase Graupel Mass*

240 Growth of mixed phase graupel mass in this sample of thunderstorms is observed prior  
 241 to the majority of total flash rate increases (Fig. 1). The median changes for the 0-2 and 2+  
 242 sigma-level categories are  $5.70 \times 10^7$  kg and  $7.15 \times 10^7$  kg, respectively. There is considerable  
 243 overlap of the inner quartile ranges (IQR) within each sigma-level category. Wilcoxon-Mann-  
 244 Whitney Rank Sum testing illustrates that the two distributions are statistically similar  
 245 (Table 2). The 0-2 and 2+ graupel mass distributions result in a Z-score of 1.065, with a  
 246 one tailed p-value of 0.096. This p-value is larger than the  $p=0.05$  value used to determine  
 247 statistical independence, so the null hypothesis of similar distributions is supported, and  
 248 larger increases in mixed phases graupel mass are not observed to statistically discriminate  
 249 between lightning jump and non-jump increases in total flash rate<sup>5</sup>.

250

251 *b. Updraft Volume*

252 The change in mixed phase  $5 \text{ m s}^{-1}$  updraft volume also does not discriminate between  
 253 lightning jump and non-jump increases in total flash rate. Small differences are observed  
 254 in the distributions between the two sigma-level categories (Fig. 2A). Medians of the 0-2  
 255 and 2+ sigma-level categories are 66 and 125  $\text{km}^3$ , respectively. Wilcoxon-Mann-Whitney

---

<sup>5</sup>These calculations only include volumes of the storm where graupel is identified as the dominant type of hydrometeor.

256 Rank Sum Testing shows that these two distributions are statistically similar Fig. 2A. The  
257 Z-score and p-value of 1.323 and 0.093 for  $5 \text{ m s}^{-1}$  updraft volume change supports the null  
258 hypothesis since the p-value is larger than the  $p=0.05$  independence threshold. This result  
259 demonstrates that larger increases in  $5 \text{ m s}^{-1}$  updraft volume are not observed to statistically  
260 discriminate between lightning jumps and non-jump increases in total flash rate.

261 The change in mixed phase  $10 \text{ m s}^{-1}$  updraft volume prior to flash rate increases shows  
262 larger differences between the jump and non-jump categories. Median growth of the  $10 \text{ m s}^{-1}$   
263 updraft volume in the 0-2 and 2+ sigma-level category is 16 and  $62 \text{ km}^3$ , respectively  
264 (Fig. 2B). Wilcoxon-Mann-Whitney Rank Sum Testing demonstrates that the two distri-  
265 butions are different. The Z-score and p-value for  $10 \text{ m s}^{-1}$  updraft volume change are 1.987  
266 and 0.0234 (Table 2). Thus, the null hypothesis is rejected for  $10 \text{ m s}^{-1}$  updraft volume,  
267 and larger increases in  $10 \text{ m s}^{-1}$  updraft volume are observed to statistically discriminate  
268 between lightning jumps and non-jump increases in total flash rate.

269

### 270 *c. Peak Updraft Speed*

271 Change in peak mixed phase updraft speed reveals a major difference in the distributions  
272 of the two sigma-level categories (Fig. 3). Medians of the 0-2 and 2+ sigma-level categories  
273 are 1 and  $5 \text{ m s}^{-1}$  from the  $1 \text{ km} \times 1 \text{ km} \times 1 \text{ km}$  resolution data used in this analysis.  
274 Wilcoxon-Mann-Whitney Rank Sum Testing shows that the two populations are different in  
275 Fig. 3. The Z-score for the change in peak updraft speed is 3.286, with a p-value of  $5.0 \times 10^{-4}$ .  
276 This indicates that the null hypothesis of similar distributions for jump and non-jump in-

277 creases in total flash rate is rejected at the  $p=0.05$  significance level (Table 2). Thus, a larger  
278 magnitude change in the mixed phase maximum updraft speed in a thunderstorm is more  
279 likely associated with the development of a  $2\sigma$  lightning jump than non-jump increases in  
280 total flash rate.

281

#### 282 *d. Timing of Increases*

283 Figure 4 shows the difference in time between the time of 0-2 and 2+ sigma-level increases  
284 in total flash rate and the maximum increase in each of the 3 parameters (graupel mass,  
285  $10 \text{ m s}^{-1}$  updraft volume, maximum updraft speed). The time of the lightning increase is  
286 subtracted from the time of the peak increase in the 3 parameters to maintain a reference  
287 frame centered on the time of the lightning increase. In general, the largest increase in  
288 graupel mass,  $10 \text{ m s}^{-1}$  updraft volume and maximum updraft speed during each 15 minute  
289 analysis window is occurring on the order of 4 to 13 minutes prior to all increases in the  
290 total flash rate.

291

## 292 **4. Discussion**

### 293 *a. The Importance of Peak Updraft Speed and $10 \text{ m s}^{-1}$ Updraft Volume*

294 Table 2 shows that the peak updraft speed is one of two parameters examined that  
295 demonstrates statistical independence between lightning jumps and non-jump increases in

296 flash rate (the other being  $10 \text{ m s}^{-1}$  updraft volume). Therefore, the maximum updraft is not  
297 necessarily well-correlated to the total flash rate over the entire lifetime of a thunderstorm,  
298 but the observations in this study indicate the increased likelihood that larger increases in  
299 maximum updraft speed are observed prior to the development of lightning jumps on shorter  
300 timescales (i.e.,  $< 15$  minutes).

301 However, this discussion goes beyond the timescale at which correlations are made in  
302 these analyses. The peak updraft speed and  $10 \text{ m s}^{-1}$  updraft volume are found to be higher  
303 than the fall speeds of ice hydrometeors responsible for electrification in thunderstorms. Ice  
304 crystals and graupel/small hail contribute to electrification of thunderstorms and their typ-  
305 ical fall speeds have been found to be  $\leq 10 \text{ m s}^{-1}$  (e.g., Dye et al. 1983, Dye et al. 1986,  
306 Musil et al. 1986, Musil and Smith 1989). The literature also shows that lightning propaga-  
307 tion typically avoids regions of peak updraft speed and intense updraft volume due to lower  
308 concentrations of precipitation size ice and a lack of available charge (e.g., Wiens et al. 2005,  
309 Payne et al. 2010, Emersic et al. 2011, Calhoun et al. 2013, Kozlowski and Carey 2014).  
310 These regions are referred to as “lightning holes.” Therefore, the outstanding question re-  
311 mains; *why* do these intense updraft characteristics matter to rapid lightning production?

312 Data from the 10 April 2009 case in Schultz et al. (2015) provides the best observational  
313 evidence of the importance of  $10 \text{ m s}^{-1}$  updraft volume and peak updraft speed working in  
314 combination to influence the total flash rate. Figure 5 shows constant altitude plan posi-  
315 tion indicator (CAPPI) at 6 km and a north-south oriented cross section through the most  
316 intense part of this developing supercell 8 minutes prior to lightning jump occurrence at  
317 1720 UTC. Flashes during this period of time are primarily initiating in regions of weaker  
318 updraft (e.g.,  $< 10 \text{ m s}^{-1}$ ). Much of the lightning activity is to the north or south of the



319 main updraft and contain convex hull-derived flash footprints<sup>6</sup>  $\geq 50 \text{ km}^2$ . Figure 6 shows a  
320 north-south oriented cross section through the same supercell at 1739 UTC 9 minutes after  
321 two consecutive lightning jumps at 1728 and 1730 UTC. The highest density of flashes is  
322 now occurring above and along the sides of the core updraft region. The location observed  
323 to have the largest number of flashes also corresponds to the region where the smallest flash  
324 footprints are found. During this period of time, the  $10 \text{ m s}^{-1}$  updraft volume and peak  
325 updraft speed increase by over  $100 \text{ km}^3$  and  $20 \text{ m s}^{-1}$ , respectively. Thus, it appears that  
326 the expansion of the  $10 \text{ m s}^{-1}$  updraft volume results in a larger three dimensional volume  
327 of weaker updraft and a larger interface between the updraft and downdraft regions. This  
328 leads to more frequent lightning flashes with smaller flash footprints in regions around the  
329 thunderstorm updraft. These regions near the updraft are known for turbulent motion (e.g.  
330 Knupp and Cotton 1982, Pantley and Lester 1990, Lane et al. 2003, Bedka et al. 2015,  
331 Behnke and Bruning 2015).

332 The measurements within this study are not at sufficient spatial and temporal resolution  
333 to examine this hypothesis beyond this inference. It is likely that the lightning jump is  
334 due to a combination of the increase in  $10 \text{ m s}^{-1}$  updraft volume (i.e., more cloud water,  
335 particle charging) and turbulence (i.e., smaller, more numerous charge regions; Bruning and  
336 MacGorman 2013); however, this hypothesis also relies on the ability of opposite charges  
337 to separate from each other in regions of higher turbulence (e.g., Bruning and MacGorman  
338 2013).

---

<sup>6</sup>Flash footprint (i.e., approximate area the flash occupies in space) calculations are made in the same manner as Schultz et al. (2015) using the convex hull methodology outlined in Bruning and MacGorman (2013).

340 *b. The Less Definitive Role of Graupel Mass for Lightning Jumps*

341 Another outcome of this study is that changes in graupel mass are not shown to be sta-  
342 tistically robust indicators that separate jumps and non-jump increases in total flash rate.  
343 Figure 1 shows an increase in graupel mass during the 15 minutes prior to most increases  
344 in total flash rate. This indicates that graupel mass changes play a similar role for both  
345 jump and non-jump increases in total flash rate. Previous studies which show ice mass and  
346 total flash rates are well-correlated over longer periods of time (i.e., entire lifecycle of the  
347 storm) also provide plausibility to this hypothesis. Deierling et al. (2008)'s Figs. 11 and 12  
348 specifically illustrate that the same ice/graupel mass magnitude results in total flash rates  
349 which differ by as much as a factor of 10. This means that the relationship is not linear and  
350 one specific graupel mass does not result in one specific flash rate. Similarly, Schultz et al.  
351 (2015) shows that similar changes in graupel mass result in different flash rates and DFRDT  
352 values (e.g., their Table 1). Therefore, the rate of change of the graupel mass also is not  
353 directly related to the rate of change of the flash rate.

354

355 *c. Kinematic and Microphysical Characteristics of Severe Storms without Jumps*

356 The lightning jump algorithm in its current form will not be a stand alone warning  
357 algorithm. There are several scenarios where severe weather is produced and lightning pro-  
358 duction is small or non-existent (e.g., Butts 2006, Schultz et al. 2009, Schultz et al. 2011).

359 In fact, nearly 40% (64/161) of the missed severe weather events by the lightning jump in  
360 Schultz et al. (2011) were due to cold season and tropical cyclone storms that produce very  
361 little lightning. These environments mainly consisted of very little thermally buoyant energy  
362 (e.g., CAPE  $\leq 500 \text{ J kg}^{-1}$ ) and strong 0-3 km wind shear (not shown).

363 The 39 thunderstorm dataset contains 6 thunderstorms that fit the low topped, cold sea-  
364 son or tropical classification which also lack lightning jumps. All 6 of these storms are severe  
365 and produce hail, high winds or tornadoes. The median (mean) increase in graupel mass for  
366 these types of severe storms is  $1.96 \times 10^7 \text{ kg}$  ( $3.54 \times 10^7 \text{ kg}$ ), and the median trend in graupel  
367 mass for these 5 storms falls below the 25<sup>th</sup> percentile for trends in graupel mass prior to  
368 lightning jumps of  $2.31 \times 10^7 \text{ kg}$  (Fig. 1). Mixed phase  $10 \text{ m s}^{-1}$  updraft volume growth  
369 and peak updraft speed intensification are also weak. The median (mean)  $10 \text{ m s}^{-1}$  updraft  
370 volume increase is  $0 \text{ km}^3$  ( $15 \text{ km}^3$ ) for these types of severe storms. Furthermore, median  
371 (mean) increases in the peak mixed phase updraft speed are only on the order of  $0.4 \text{ m s}^{-1}$   
372 ( $1.5 \text{ m s}^{-1}$ ) prior to their peak increase in total flash rate. Thus, there is a lack of mixed  
373 phase updraft growth or a total absence of  $10 \text{ m s}^{-1}$  updraft volume within this set of storms  
374 (3 of the 6 cases have a  $10 \text{ m s}^{-1}$  updraft volume of  $0 \text{ km}^3$ ). Weaker magnitude changes in  
375 peak vertical velocity are also observed in these storms (Fig. 3). These weaker mixed phase  
376 kinematic properties limit the storm's potential to produce lightning and lightning jumps  
377 prior to severe weather occurrence.

378 The weak mixed phase updraft magnitudes and changes in magnitude observed in this  
379 study are similar to those found in other shallow severe storms in previous observational and  
380 modeling studies (e.g., McCaul and Weisman 1996, Cantrell 1995, Knupp et al. 1998, Eastin  
381 and Link 2009). This indicates that severe weather production does not always require ro-

382 bust mixed phase updrafts. This is why lightning trends in these types of storm may not  
383 always be useful for providing lead time on severe weather occurrence because of the limited  
384 size of the updraft or lack of strong mixed updraft speeds in cold season, low topped, or  
385 tropical cyclone severe thunderstorm environments.

386

## 387 5. Conclusions

388 The results of this work provide a comprehensive statistical evaluation for physical pa-  
389 rameters which are hypothesized to modulate electrification and lightning in thunderstorms.  
390 A large dataset of 39 thunderstorms with 214-fifteen minute analysis windows are used to  
391 assess trends in mixed phase graupel mass, updraft volume and updraft speed prior to light-  
392 ning jumps and other non-jump increases in total flash rate. The following conclusions were  
393 made from this analysis:

394

395 • Graupel mass is not observed to be a statistically significant discriminator between  
396 lightning jumps (i.e., 2+ sigma-level) and non-jump (i.e., 0-2 sigma-level) increases  
397 in total flash rate. The one-tailed p-value for the independence test of the jump and  
398 non-jump distributions is  $p=0.096$ .

399

400 • The change in  $5 \text{ m s}^{-1}$  updraft volume is also not observed to be statistically significant  
401 discriminator between lightning jumps and non-jump increases in total flash rate. The

402 one-tailed p-value for the independence test of the jump and non-jump distributions  
403 is  $p=0.093$ .

404

405 • Larger increases in  $10 \text{ m s}^{-1}$  updraft volume are observed for lightning jumps versus  
406 those observed with non-jump increases in total flash rate ( $p=0.0234$ ). The median  
407 change in  $10 \text{ m s}^{-1}$  updraft volume for jump and non-jump categories is  $62 \text{ km}^3$  and  
408  $16 \text{ km}^3$ , respectively.

409

410 • Larger magnitude increases in peak updraft speed are observed for lightning jumps  
411 versus those observed with non-jump increases in total flash rate ( $p=5.0 \times 10^{-4}$ ). The  
412 median change in maximum updraft speed is  $5 \text{ m s}^{-1}$  and  $1 \text{ m s}^{-1}$  for jump and non-  
413 jump increases, respectively (Fig. 3).

414

415 • Very little difference is found in the timing between peak increase in each of the three  
416 kinematic/microphysical parameters (mixed phase graupel mass,  $10 \text{ m s}^{-1}$  updraft  
417 volume and peak maximum updraft speed) relative to the time of the total flash rate  
418 increase. In general, growth occurs between 4 and 13 minutes in advance of most flash  
419 rate increases.

420

421 • A sample of 6 severe thunderstorms that did not produce lightning jumps demonstrate  
422 that the main characteristic lacking in these storms is mixed phase updraft. These

423 storms lack significant changes in  $10 \text{ m s}^{-1}$  updraft volume and the magnitude of the  
424 peak updraft speed in the mixed phase region during their largest increases in total  
425 flash rate.

426

427 These strong statistical results support the use of lightning jumps to infer changes in stronger  
428 updraft characteristics in thunderstorms. Often these physical parameters are not readily  
429 available in operational datasets, and thus the lightning data can provide some indication  
430 on the trend of the mixed phase updraft (growing vs weakening). Future work will need  
431 to demonstrate the physical connections between mixed phase updraft growth and severe  
432 weather production in thunderstorms.

433

#### 434 *Acknowledgments.*

435 The authors would like to acknowledge Dr. Steven J. Goodman and GOES-R Risk Re-  
436 duction Research funding for support of this research. C. Schultz would like to acknowledge  
437 support from the NASA Pathways Intern Program at Marshall Space Flight Center, namely  
438 Julie Clift and Christopher Randall. The authors are; thankful for technical support with  
439 radar data processing from Lamont Bain, Retha Mecikalski, and Danielle Kozlowski for parts  
440 of the June 11 and April 10 events. Furthermore, productive conversations with Drs. Themis  
441 Chronis, Phil Bitzer, Hugh Christian, Walt Petersen, Kristin Calhoun and Eric Bruning ben-  
442 efited the outcomes of this research. We also gratefully acknowledge the technical support  
443 for maintenance of the operational instrumentation, namely Dustin Phillips, Patrick Gatlin,

444 Chris Phillips for the maintenance of the ARMOR radar. The authors would also like to  
445 recognize Jeff Bailey and Blair Breitreiter for the continued maintenance of the North Al-  
446 abama Lightning Mapping Array. Finally, the authors would like to thank Editor Dr. Paul  
447 Markowski and two anonymous reviewers for their helpful comments which improved the  
448 content and technical writing of this manuscript.

449

450

## APPENDIX

451

452

### Appendix A

#### 453 *a. $2\sigma$ Lightning Jump Algorithm*

454 Although the lightning jump algorithm has been described in previous work (e.g., Schultz  
455 et al. 2009, Schultz et al. 2011, Chronis et al. 2015) it is good to review the formulation of the  
456 algorithm for reference to this work. The primary source of lightning data for this algorithm  
457 has been lightning mapping arrays with the goal of ultimately utilizing GLM once GOES-R  
458 data are operationally available.

459 The algorithm starts with 14 minutes of total lightning data which has been assigned to  
460 a specific thunderstorm. For this example,  $t_0$  is the most recent minute of data, and  $t-13$  is  
461 the oldest minute of data. First, 1 minute flash rates are combined to produce an average

462 flash rate every 2 minutes. For example, the average flash rate for time  $t_0$  and time  $t-1$  is,

$$FR_{avg}(t_0)(flashes\ min^{-1}) = \frac{FR_{t_0} + FR_{t-1}}{2\ minutes}, \quad (A1)$$

463 while the average flash rate for times  $t-12$  and  $t-13$  would be,

$$FR_{avg}(t-12)(flashes\ min^{-1}) = \frac{FR_{t-12} + FR_{t-13}}{2\ minutes}. \quad (A2)$$

464

465 Now there are a total of 7 1-minute average flash rates:  $FR_{avg}(t_0)$ ,  $FR_{avg}(t-2)$ ,  $FR_{avg}(t-4)$ ,  $FR_{avg}(t-6)$ ,  $FR_{avg}(t-8)$ ,  $FR_{avg}(t-10)$ ,  $FR_{avg}(t-12)$ . Next, subsequent  $FR_{avg}$  times  
 466 4),  $FR_{avg}(t-6)$ ,  $FR_{avg}(t-8)$ ,  $FR_{avg}(t-10)$ ,  $FR_{avg}(t-12)$ . Next, subsequent  $FR_{avg}$  times  
 467 are subtracted from each other to obtain the rate of change of the total flash rate, or more  
 468 commonly known as DFRDT. For the rate of change in the flash rate between  $FR_{avg}(t_0)$  and  
 469  $FR_{avg}(t-2)$  the equation would be,

$$DFRDT_{t_0} = \frac{FR_{avg}(t_0) - FR_{avg}(t-2)}{2\ minutes} = DFRDT(flashes\ min^{-2}), \quad (A3)$$

470 while for  $FR_{avg}(t-10)$  and  $FR_{avg}(t-12)$  the equation would be,

$$DFRDT_{t-10} = \frac{FR_{avg}(t-10) - FR_{avg}(t-12)}{2\ minutes} = DFRDT(flashes\ min^{-2}). \quad (A4)$$

471

472 Now there are a total of 6 DFRDT values for the algorithm to use to identify a lightning  
 473 jump ( $DFRDT_{t_0}$ ,  $DFRDT_{t-2}$ ,  $DFRDT_{t-4}$ ,  $DFRDT_{t-6}$ ,  $DFRDT_{t-8}$ ,  $DFRDT_{t-10}$ ).  $DFRDT_{t_0}$   
 474 is the current rate of change of the total flash rate in the storm, while  $DFRDT_{t-2}$ ,  $DFRDT_{t-4}$ ,  
 475  $DFRDT_{t-6}$ ,  $DFRDT_{t-8}$  and  $DFRDT_{t-10}$  are used to calculate the standard deviation of the  
 476 rate of change of the total flash rate in the storm between time  $t-2$  up to (and not in-  
 477 cluding)  $t-14$ . The result is the sigma-level calculation found in Eqn. 2. A sigma-level



478 value  $\geq 2$  identifies a lightning jump, while a sigma-level value  $< 2$  is identified as a non-  
479 jump increase in the total flash rate. This representation of the  $2\sigma$  lightning jump algorithm  
480 provides users with more information than the previous algorithm (i.e., a yes/no answer  
481 that the  $2\sigma$  lightning jump threshold has been exceeded) by allowing the user to determine  
482 how far above or below any increase in total flash rate is relative to the dynamic  $2\sigma$  threshold.

483

## REFERENCES

- 486 Armijo, L., 1969: A theory for the determination of wind and precipitation velocities with  
487 Doppler radars. *J. Atmos. Sci.*, **26**, 570–573.
- 488 Bedka, K. M., C. Wang, R. Rogers, L. D. Carey, W. Feltz, and J. Kanak, 2015: Examining  
489 deep convective cloud evolution using total lightning, WSR-88D, and GOES-14 Super  
490 Rapid Scan datasets. *Wea. Forecasting*, doi:<http://dx.doi.org/10.1175/WAF-D-14-00062>.  
491 1.
- 492 Behnke, S. A. and E. C. Bruning, 2015: Changes to the turbulent kinematics of a volcanic  
493 plume inferred from lightning data. *Geophys. Res. Lett.*, **42**, 4232–4239.
- 494 Brandes, E. A., 1977: Flow in severe thunderstorms observed by dual-Doppler radar. *Mon.*  
495 *Wea. Rev.*, **105**, 113–120.
- 496 Bringi, V. N., T. D. Keenan, and V. Chandrasekar, 2001: Correcting C-band radar reflec-  
497 tivity and differential reflectivity data for rain attenuation: A self-consistent method with  
498 constraints. *IEEE Trans. On Geo. and Rem. Sens.*, **39**, 1906–1915.
- 499 Bringi, V. N., K. R. Knupp, A. Detwiler, L. Liu, I. J. Caylor, and R. A. Black, 1997: Evo-  
500 lution of a Florida thunderstorm during the Convection and Precipitation/Electrification  
501 Experiment: The case of 9 August 1991. *Mon. Wea. Rev.*, **125**, 2131–2160.
- 502 Bruning, E. C. and D. M. MacGorman, 2013: Theory and observations of controls on light-  
503 ning flash spectra. *J. Atmos. Sci.*, **70**, 4012–4029.

504 Butts, D., 2006: *An examination of the relationship between cool season tornadoes and*  
505 *cloud-to-ground lightning flashes*. MS Thesis, Texas A and M University, 109 pp.

506 Byers, H. R. and R. R. Braham, 1949: *The Thunderstorm*. U. S. Government Printing Office,  
507 287 pp.

508 Calhoun, K. M., 2015: Forecaster use of total lightning data for short-term forecasts and  
509 warnings in the Hazardous Weather Testbed. *7th Conf. on Meteorological Applications of*  
510 *Lightning Data*, Phoenix, AZ.

511 Calhoun, K. M., D. R. MacGorman, C. L. Ziegler, and M. I. Biggerstaff, 2013: Evolution of  
512 lightning activity and storm charge relative to dual-Doppler analysis of a high-precipitation  
513 supercell storm. *Mon. Wea. Rev.*, **141**, 2199–2223.

514 Cantrell, L. E., Jr., 1995: *The Role of Vertical Buoyancy Distribution in Simulated Low-*  
515 *Topped Supercells*. M.S. Thesis, Texas A & M University, 135 pp.

516 Carey, L. D., W. J. Koshak, H. Peterson, and R. M. Mecikalski, 2016: The kinematic and  
517 microphysical control of lightning rate, extent, and NO<sub>x</sub> production. *J. Geophys. Res.*,  
518 **121**, 7975–7989.

519 Carey, L. D. and S. A. Rutledge, 1996: A multiparameter radar case study of the micro-  
520 physical and kinematic evolution of a lightning producing storm. *Meteorol. Atmos. Phys.*,  
521 **59**, 33–64.

522 Chronis, T., L. D. Carey, C. J. Schultz, E. V. Schultz, K. M. Calhoun, and S. J. Goodman,  
523 2015: Exploring lightning jump characteristics. *Wea. Forecasting*, Accepted, in publica-  
524 tion.

- 525 Crum, T. D. and R. L. Alberty, 1993: The WSR-88D and the WSR-88D Operational Support  
526 Facility. *Bull. Amer. Meteor. Soc.*, **74**, 1669–1687.
- 527 Davies-Jones, R. P., 1979: Dual-Doppler radar coverage area as a function of measurement  
528 accuracy and spatial resolution. *J. Appl. Meteor.*, **18**, 1229–1233.
- 529 Deierling, W. and W. A. Petersen, 2008: Total lightning activity as an indicator of updraft  
530 characteristics. *J. Geophys. Res.*, **113**, doi:10.1029/2007JD009598.
- 531 Deierling, W., W. A. Petersen, J. Latham, S. Ellis, and H. J. Christian, 2008: The relation-  
532 ship between lightning activity and ice fluxes in thunderstorms. *J. Geophys. Res.*, **113**,  
533 doi:10.1029/2007JD009700.
- 534 Dixon, M. and G. Wiener, 1993: TITAN: Thunderstorm Identification, Tracking Analysis  
535 and Nowcasting-A radar-based methodology. *J. Atmos. Ocean Tech.*, **10**, 785–797.
- 536 Dye, J. E. and Coauthors, 1986: Early electrification and precipitation development in a  
537 small, isolated Montana cumulonimbus. *J. Geophys. Res.*, **91**, 1231–1247.
- 538 Dye, J. E., B. E. Martner, and L. J. Miller, 1983: Dynamical-microphysical evolution of a  
539 convective storm in a weakly-sheared environment. Part I: Microphysical observations and  
540 interpretation. *J. Atmos. Sci.*, **40**, 12083–2096.
- 541 Dye, J. E., W. P. Winn, J. J. Jones, and D. W. Breed, 1989: The electrification of New  
542 Mexico thunderstorms. 1. Relationship between precipitation development and the onset  
543 of electrification. *J. Geophys. Res.*, **94**, 8643–8656.

544 Eastin, M. D. and M. C. Link, 2009: Minature supercells in an offshort outer rainband of  
545 Hurricane Ivan (2004). *Mon. Wea. Rev.*, **137**, 2081–2104.

546 Emersic, C., P. L. Heinselman, D. R. MacGorman, and E. C. Bruning, 2011: Lightning  
547 activity in a hail-producing storm observed with phased array radar. *Mon. Wea. Rev.*,  
548 **139**, 1809–1824, doi:10.1175/2010MWR3574.1.

549 Gatlin, P. N. and S. J. Goodman, 2010: A total lightning trending algorithm to identify  
550 severe thunderstorms. *J. Atmos. Oceanic Technol.*, **27**, 3–22.

551 Goodman, S. J., D. E. Buechler, P. D. Wright, and W. D. Rust, 1988: Lightning and  
552 precipitation history of a microburst-producing storm. *Geophys. Res. Lett.*, **15**, 1185–1188.

553 Goodman, S. J. and Coauthors, 2005: The North Alabama Lightning Mapping Array: Re-  
554 cent severe storm observations and future prospects. *Atmos. Res.*, **76**, 423–437.

555 Goodman, S. J. and Coauthors, 2013: The GOES-R Geostationary Lightning Mapper  
556 (GLM). *Atmos. Res.*, **125-126**, 34–49.

557 Heymsfield, A. J. and K. M. Miller, 1988: Water vapor and ice mass transported into the  
558 anvils of CCOPE thunderstorms: Comparison with storm influx and rainout. *J. Atmos.*  
559 *Sci.*, **45**, 3501–3514.

560 Johnson, E. V., 2009: *Behavior of Lightning and Updrafts for Severe and Non Severe Thun-*  
561 *derstorms in Northern Alabama*. M.S. Thesis, University of Alabama-Huntsville, 70 pp.

562 Knupp, K. R. and Coauthors, 2014: Meteorological overview of the devastating 27 April

563 2011 tornado outbreak. *Bull. Amer. Meteor. Soc.*, **95**, 1041–1062, doi:[http://dx.doi.org/](http://dx.doi.org/10.1175/BAMS-D-11-00229.1)  
564 10.1175/BAMS-D-11-00229.1.

565 Knupp, K. R. and W. R. Cotton, 1982: An intense, quasi-steady thunderstorm over moun-  
566 tainous terrain: Part III: Doppler radar observations of the turbulent structure. *J. Atmos.*  
567 *Sci.*, **39**, 359–368.

568 Knupp, K. R., J. R. Stalker, and E. W. McCaul, 1998: An observational and numerical  
569 study of a mini-supercell storm. *Atmos. Res.*, **49** (1), 35–63.

570 Koshak, W. J. and Coauthors, 2004: North Alabama Lightning Mapping Array (LMA):  
571 VHF source retrieval algorithm and error analysis. *J. Atmos. Ocean. Tech.*, **21**, 543–558.

572 Kozlowski, D. M. and L. D. Carey, 2014: An analysis of lightning holes in  
573 Northern Alabama severe storms using a lightning mapping array and dual-  
574 polarization radar. *Preprints, 5th Inter. Lightning Meteorology Conf.*, Tuscon, AZ,  
575 <http://www.vaisala.com/en/events/ildcilmc/Pages/ILDC-2014-archive.aspx>.

576 Kuhlman, K. M., C. L. Zeigler, E. R. Mansell, D. R. MacGorman, and J. M. Straka, 2006:  
577 Numerically simulated electrification and lightning of the 29 June 2000 STEPS supercell  
578 storm. *Mon. Wea. Rev.*, **134**, 2734–2757.

579 Lane, T. P., R. D. Sharman, T. L. Clark, and H. M. Hsu, 2003: An investigation of turbulence  
580 generation mechanisms above deep convection. *J. Atmos. Sci.*, **60**, 1297–1321.

581 Lang, T. J. and S. A. Rutledge, 2002: Relationships between convective storm kinematics,  
582 precipitation, and lightning. *Mon. Wea. Rev.*, **130**, 2492–2506.

583 Matejka, T. and D. L. Bartels, 1998: The accuracy of vertical air velocities from Doppler  
584 radar data. *Mon. Wea. Rev.*, **92**, 92–117.

585 McCaul, E. W., S. J. Goodman, K. M. LaCasse, and D. J. Cecil, 2009: Forecasting lightning  
586 threat using cloud-resolving model simulations. *Weather and Forecasting*, **24 (3)**, 709–729,  
587 doi:10.1175/2008WAF2222152.1.

588 McCaul, E. W., Jr. and M. L. Weisman, 1996: Simulations of shallow supercell storms in  
589 landfalling hurricane environments. *Mon. Wea. Rev.*, **124**, 408–429.

590 Mecikalski, R. M., A. L. Bain, and L. D. Carey, 2015: Radar and lightning observations of  
591 deep moist convection across Northern Alabama during DC3: 21 May 2012. *Mon. Wea.*  
592 *Rev.*, **143**, 2774–2794.

593 Mohr, C. G., L. J. Miller, R. L. Vaughn, and H. W. Frank, 1986: On the merger of mesoscale  
594 datasets into a common Cartesian format for efficient and synthetic analysis. *J. Atmos.*  
595 *Oceanic. Technol.*, **3**, 141–161.

596 Musil, D. J., A. J. Heymsfield, and P. L. Smith, 1986: Microphysical characteristics of a  
597 well-developed weak echo region in a High Plains supercell thunderstorm. *J. Clim. and*  
598 *Appl. Meteor.*, **25**, 1037–1051.

599 Musil, D. J. and P. L. Smith, 1989: Interior characteristics at mid-levels of thunderstorms  
600 in the Southeastern United States. *Atmos. Res.*, **24**, 149–167.

601 O’Brien, J. J., 1970: Alternative solutions to the classical vertical velocity problem. *J. Appl.*  
602 *Meteor.*, **9**, 197–203.

603 Oye, D. and M. Case, 1995: REORDER: A Program for Gridding Radar Data. Installa-  
604 tion and User Manual for the UNIX Version. NCAR Atmospheric Technology Division,  
605 Boulder, CO, 19 pp.

606 Oye, D., C. Mueller, and S. Smith, 1995: Software for radar translation, visualization, editing  
607 and interpolation. *Preprints, 27th Conf. on Radar Meteorology*, 359–361, Vail, CO, Amer.  
608 Met. Soc.

609 Pantley, K. C. and P. F. Lester, 1990: Observations of severe turbulence near thunderstorm  
610 tops. *J. Appl. Meteor.*, **60**, 1171–1179.

611 Payne, C. D., T. J. Schuur, D. R. MacGorman, M. I. Biggerstaff, K. M. Kuhlman, and W. D.  
612 Rust, 2010: Polarimetric and electrical characteristics of a lightning ring in a supercell  
613 storm. *Mon. Wea. Rev.*, **138**, 2405–2425.

614 Ray, P. S., C. L. Ziegler, W. Bumgarner, and R. J. Serafin, 1980: Single and multiple-Doppler  
615 radar observations of tornadic storms. *Mon. Wea. Rev.*, **108**, 1607–1625.

616 Rudlosky, S. D. and H. E. Fuelberg, 2013: Documenting storm severity in the Mid-Atlantic  
617 region using lightning and radar information. *Mon. Wea. Rev.*, **141**, 3186–3202, doi:10.  
618 1175/MWR-D-12-00287.1.

619 Saunders, C. P. R., H. Bax-Norman, C. Emersic, E. E. Avila, and N. E. Castellano, 2006:  
620 Laboratory studies of the effect of cloud conditions on graupel/crystal charge transfer in  
621 thunderstorm electrification. *Quart. J. Roy. Meteor. Soc.*, **132**, 2653–2673.

622 Schultz, C. J., L. D. Carey, E. V. Schultz, and R. L. Blakeslee, 2015: Insight into the



623 physical and dynamical processes that control rapid increases in total flash rate. *Mon.*  
624 *Wea. Forecasting*, **30**, 1591–1621.

625 Schultz, C. J. and Coauthors, 2012: Dual-polarization tornadic debris signatures Part I:  
626 Examples and utility in an operational setting. *Electronic J. Operational Meteor.*, **13**,  
627 120–137.

628 Schultz, C. J., W. A. Petersen, and L. D. Carey, 2009: Preliminary development and evalu-  
629 ation of lightning jump algorithms for the real-time detection of severe weather. *J. Appl.*  
630 *Meteor.*, **48**, doi:10.1175/2009JAMC2237.1.

631 Schultz, C. J., W. A. Petersen, and L. D. Carey, 2011: Lightning and severe weather: A  
632 comparison between total and cloud-to-ground lightning trends. *Wea. Forecasting*, **26**,  
633 744–755, doi:10.1175/WAF-D-10-05026.1.

634 Takahashi, T., 1978: Riming electrification as a charge generation mechanism in thunder-  
635 storms. *J. Atmos. Sci.*, **35**, 1536–1548.

636 Tessendorf, S. A., L. J. Miller, K. C. Wiens, and S. A. Rutledge, 2005: The 29 June 2000  
637 supercell observed during STEPS. Part I: Kinematics and microphysics. *JAS*, **62**, 4127–  
638 4150.

639 Tuttle, J. D., V. N. Bringi, H. D. Orville, and F. J. Kopp, 1989: Multiparameter radar study  
640 of a microburst - comparison with model results. *J. Atmos. Sci.*, **46**, 601–620.

641 Vivekanandan, J., D. S. Zrnic, S. M. Ellis, R. Oye, A. V. Ryzhkov, and J. Straka, 1999:  
642 Cloud microphysics retrieval using S-band dual-polarization radar measurements. *Bull.*  
643 *Amer. Met. Soc.*, **80**, 381–388.

- 644 Wiens, K. C., S. A. Rutledge, and S. A. Tessendorf, 2005: The 29 June 2000 supercell  
645 observed during steps. Part II: Lightning and charge structure. *J. Atmos. Sci.*, **62**, 4151–  
646 4177.
- 647 Wilks, D. S., 1995: *Statistical Methods in the Atmospheric Sciences*. Academic Press, 467  
648 pp.
- 649 Williams, E. R. and Coauthors, 1999: The behavior of total lightning activity in severe  
650 Florida thunderstorms. *Atmos. Res.*, **51**, 245–265.
- 651 Williams, E. R., M. E. Weber, and R. E. Orville, 1989: The relationship between lightning  
652 type and convective state of thunderclouds. *J. Geophys. Res.*, **94**, 13 213–13 220.
- 653 Witt, A., M. D. Eilts, G. J. Stumpf, J. T. Johnson, E. D. Mitchell, and K. W. Thomas, 1998:  
654 An enhanced hail detection algorithm for the WSR-88D. *Wea. Forecasting*, **13**, 286–303.
- 655 Workman, E. J. and S. E. Reynolds, 1949: Electrical activity as related to thunderstorm cell  
656 growth. *Bull. Amer. Meteor. Soc.*, **30**, 142–144.

## 657 List of Tables

658	1	Dates, number of storms, storm type and the number of storms from each day	
659		used in this analysis.	35
660	2	Z-scores and p-values using Wilcoxon-Mann-Whitney Rank Sum Testing be-	
661		tween the 0-2 and 2+ sigma-level categories for Graupel Mass Change (kg),	
662		5 and 10 m s <sup>-1</sup> updraft volume change (km <sup>3</sup> ) and maximum vertical velocity	
663		change (m s <sup>-1</sup> )	36

TABLE 1. Dates, number of storms, storm type and the number of storms from each day used in this analysis.

Location	Number	Type	Jump	No Jump
3 May 2006	1	multicell	1	0
19 July 2006	2	multicell	2	0
3 April 2007	3	supercell	2	1
4 April 2007	1	QLCS	1	0
1 June 2007	4	multicell	0	4
7 July 2007	2	multicell	2	0
17 August 2007	6	multicell	5	1
14 September 2007	1	tropical	0	1
10 April 2009	3	supercell	3	0
13 April 2009	1	low topped	0	1
21 January 2010	2	low topped	1	1
12 March 2010	1	QLCS	1	0
26 October 2010	3	supercell	0	3
27 April 2011	1	supercell	1	0
18 May 2012	1	multicell	1	0
21 May 2012	1	multicell	0	1
11 June 2012	4	multicell	0	4
14 June 2012	2	multicell	0	2

TABLE 2. Z-scores and p-values using Wilcoxon-Mann-Whitney Rank Sum Testing between the 0-2 and 2+ sigma-level categories for Graupel Mass Change (kg), 5 and 10 m s<sup>-1</sup> updraft volume change (km<sup>3</sup>) and maximum vertical velocity change (m s<sup>-1</sup>)

	Graupel Mass	5 m s <sup>-1</sup>	10 m s <sup>-1</sup>	MaxVV
Z-Score	1.065	1.323	1.987	3.286
p-value (one tailed)	0.096	0.093	0.0234	5.0×10 <sup>-4</sup>

## 664 List of Figures

- 665 1 Box plots of storm graupel mass change (kg) versus the sigma-level of the  
666 subsequent increase in total flash rate. Data from non-jump increases in total  
667 flash rates are in blue (i.e., 0-2 sigma-level), while data from lightning jump  
668 events are in orange. Median, 25<sup>th</sup>, and 75<sup>th</sup> changes are to the right of each  
669 box, and the population size of each bin is on the left. “X” marks indicate  
670 the individual data points within each sigma-level category. 39
- 671 2 Same as Fig. 1, but for storm 5 m s<sup>-1</sup> (Panel A) and 10 m s<sup>-1</sup> (Panel B) updraft  
672 volume change (km<sup>3</sup>) versus the sigma-level of the subsequent increase in total  
673 flash rate. 40
- 674 3 Same as Fig. 1, but for the change in peak updraft speed (m s<sup>-1</sup>) versus the  
675 sigma-level of the subsequent increase in total flash rate. 41
- 676 4 Same as Fig. 1, but for the timing of the peak increase in graupel mass  
677 (Panel A), 10 m s<sup>-1</sup> updraft volume (Panel B) and maximum updraft up-  
678 draft speed (Panel C) minus the time of flash rate increase distributed versus  
679 the sigma-level of the subsequent increase in total flash rate. 42

- 680 5 Reflectivity, updraft velocity, flash extent density and mean flash footprint  
681 prior to lightning jump occurrence in a storm transitioning from multicell to  
682 supercell on 10 April 2009. Panel A is a CAPPI of reflectivity at 6 km with  
683 reflectivity (shaded every 5-dB starting at 5-dBZ, velocity (black contours in  
684  $10 \text{ m s}^{-1}$  increments starting at  $10 \text{ m s}^{-1}$ ) and lightning flash origin (black  
685 dots within 2 minutes of radar volume time) are overlaid. The gray dashed  
686 rectangle represents the region which lightning data for flash extent density  
687 and mean flash footprint are calculated from in Panels B and C. Panel B is  
688 flash extent density (flashes  $\text{km}^{-2}$ ) in  $1 \text{ km} \times 1 \text{ km}$  bins within 2 minutes of  
689 radar volume start time with reflectivity from ARMOR (solid black contours  
690 every 10-dB, starting at 10-dBZ) and vertical velocity (blue contours starting  
691 at  $5 \text{ m s}^{-1}$ , then in  $10 \text{ m s}^{-1}$  increments after  $10 \text{ m s}^{-1}$ ) overlaid. Panel C is  
692 mean flash footprint ( $\text{km}^2$ ) within 2 minutes of radar volume start time with  
693 with reflectivity from ARMOR (solid black contours every 10-dB, starting at  
694 10-dBZ) and vertical velocity (blue dashed contours starting at  $5 \text{ m s}^{-1}$ , then  
695 in  $10 \text{ m s}^{-1}$  increments after  $10 \text{ m s}^{-1}$ ) overlaid. 43
- 696 6 Same as Fig. 5 but for 1739 UTC on 10 April 2009. 44

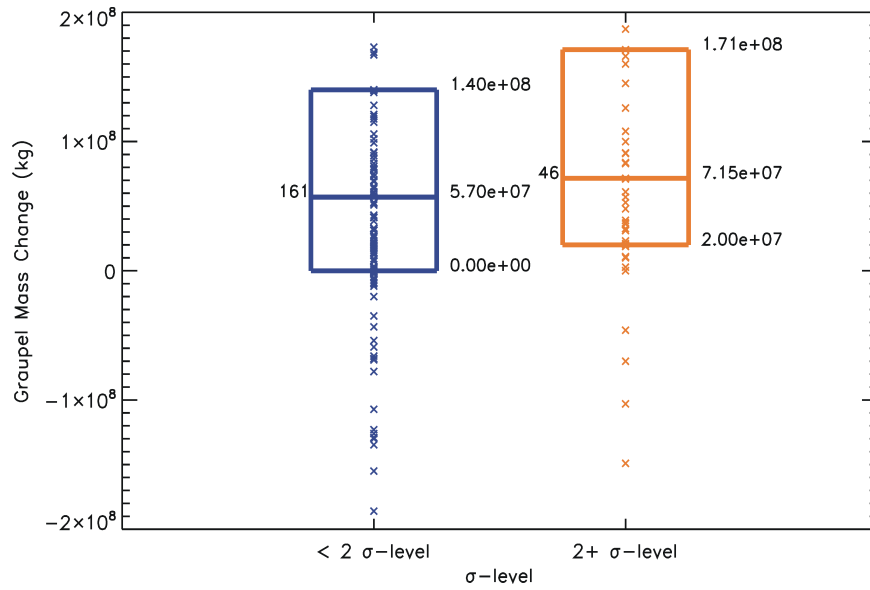


FIG. 1. Box plots of storm graupel mass change (kg) versus the sigma-level of the subsequent increase in total flash rate. Data from non-jump increases in total flash rates are in blue (i.e., 0-2 sigma-level), while data from lightning jump events are in orange. Median, 25<sup>th</sup>, and 75<sup>th</sup> changes are to the right of each box, and the population size of each bin is on the left. “X” marks indicate the individual data points within each sigma-level category.



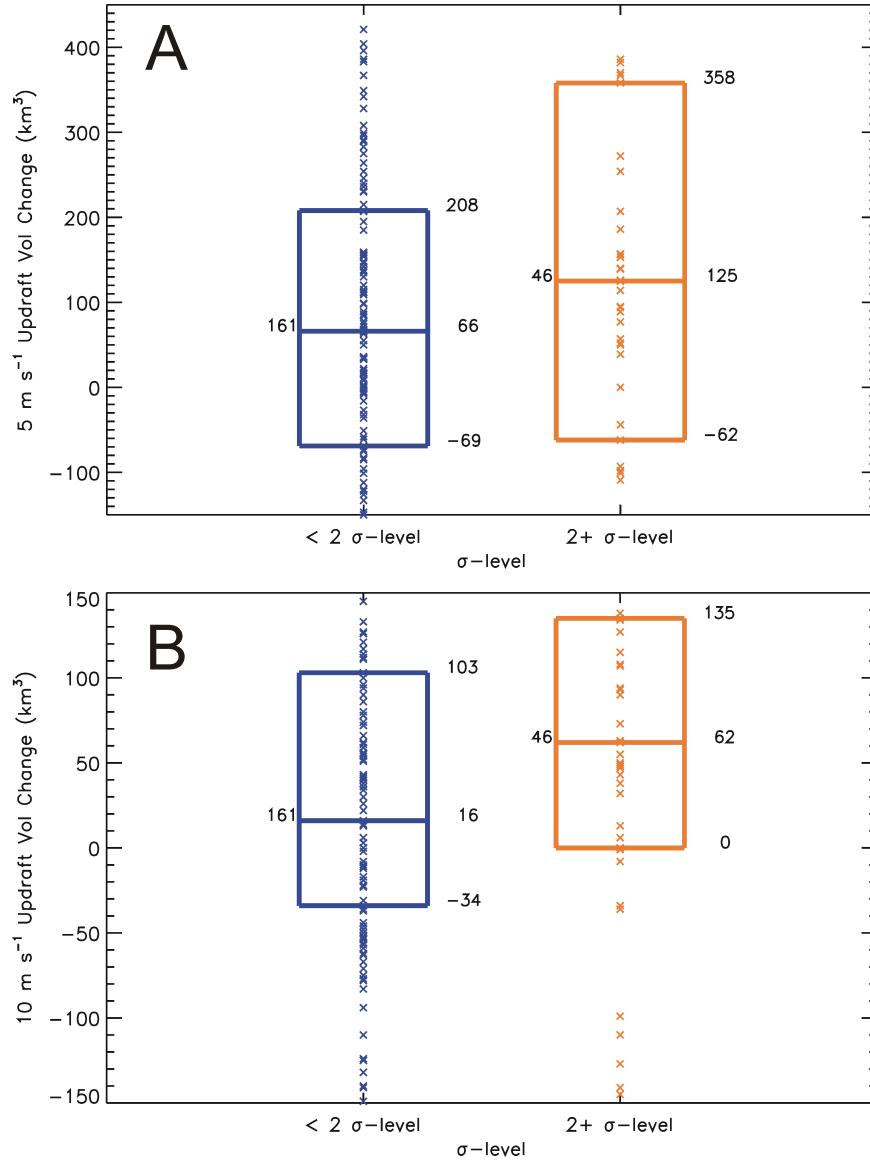


FIG. 2. Same as Fig. 1, but for storm 5 m s<sup>-1</sup> (Panel A) and 10 m s<sup>-1</sup> (Panel B) updraft volume change (km<sup>3</sup>) versus the sigma-level of the subsequent increase in total flash rate.

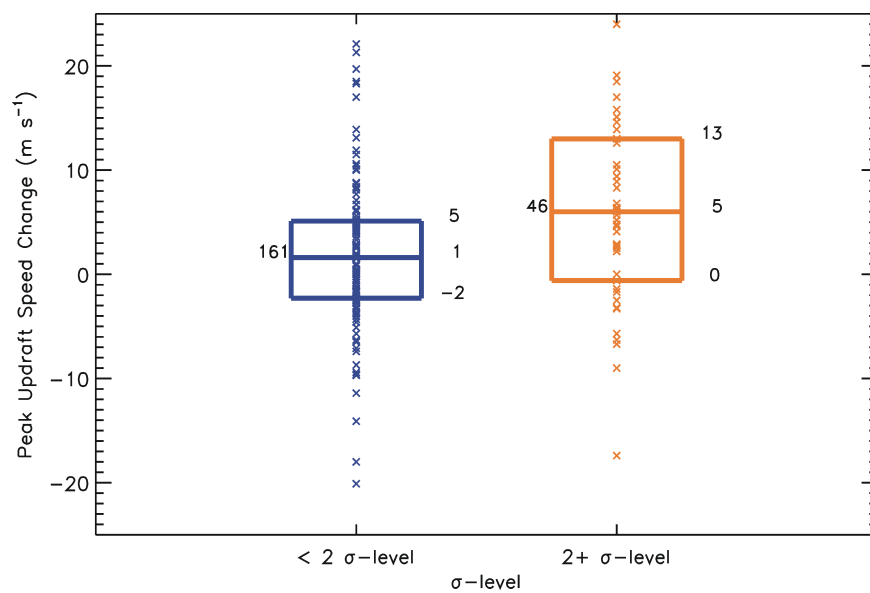


FIG. 3. Same as Fig. 1, but for the change in peak updraft speed ( $\text{m s}^{-1}$ ) versus the sigma-level of the subsequent increase in total flash rate.

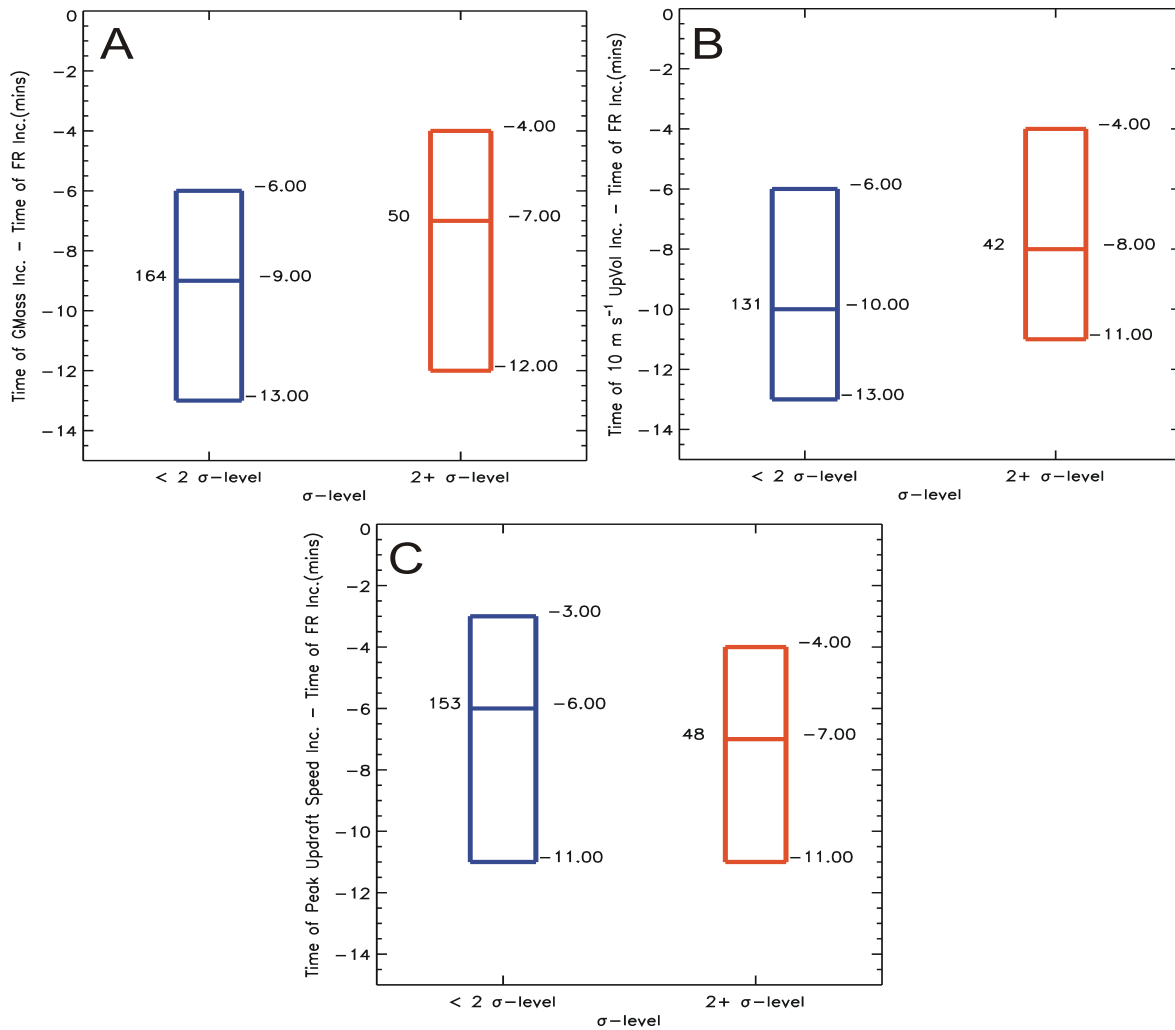


FIG. 4. Same as Fig. 1, but for the timing of the peak increase in graupel mass (Panel A),  $10 \text{ m s}^{-1}$  updraft volume (Panel B) and maximum updraft updraft speed (Panel C) minus the time of flash rate increase distributed versus the sigma-level of the subsequent increase in total flash rate.

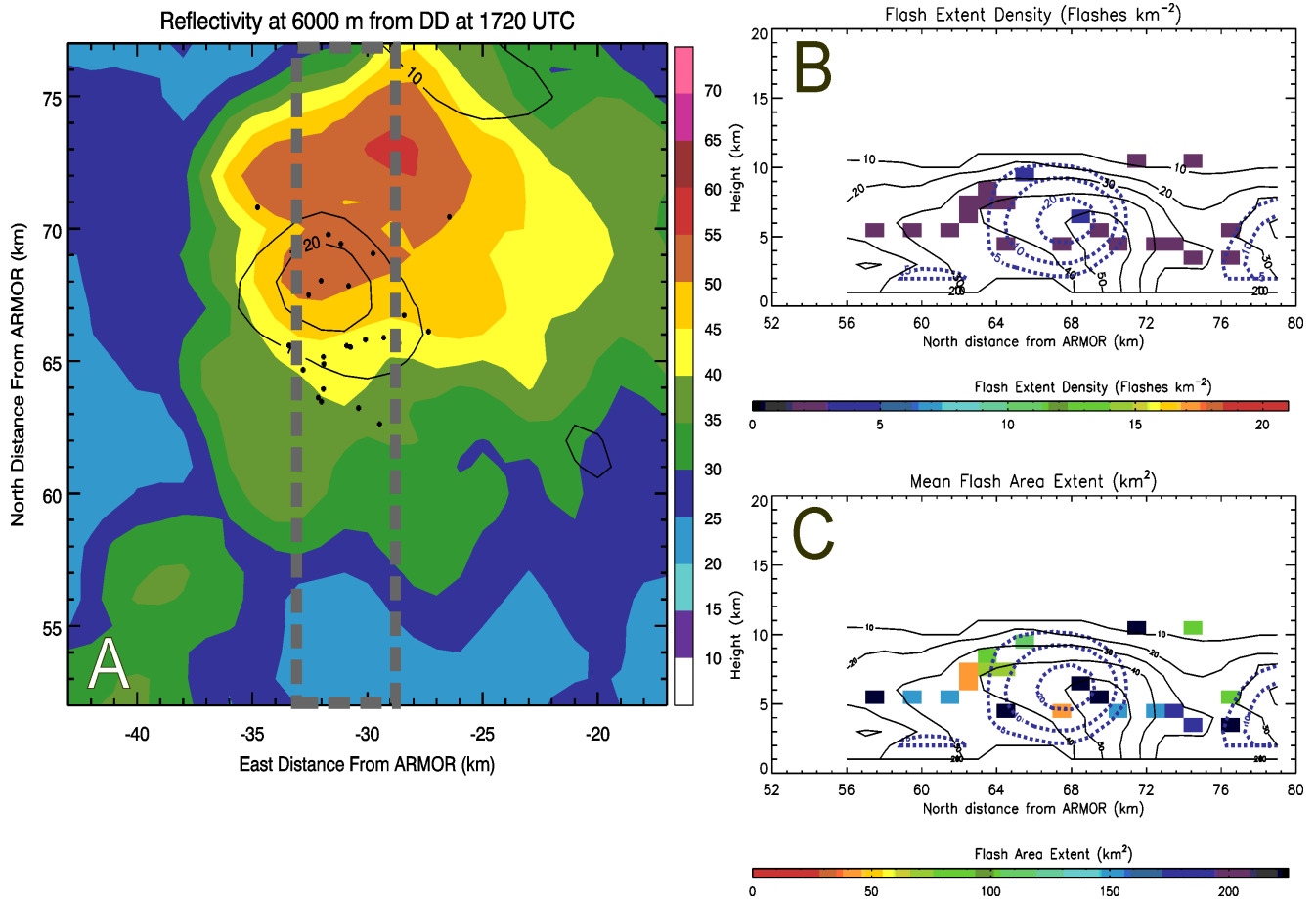


FIG. 5. Reflectivity, updraft velocity, flash extent density and mean flash footprint prior to lightning jump occurrence in a storm transitioning from multicell to supercell on 10 April 2009. Panel A is a CAPPI of reflectivity at 6 km with reflectivity (shaded every 5-dB starting at 5-dBZ), velocity (black contours in  $10 \text{ m s}^{-1}$  increments starting at  $10 \text{ m s}^{-1}$ ) and lightning flash origin (black dots within 2 minutes of radar volume time) are overlaid. The gray dashed rectangle represents the region which lightning data for flash extent density and mean flash footprint are calculated from in Panels B and C. Panel B is flash extent density ( $\text{flashes km}^{-2}$ ) in  $1 \text{ km} \times 1 \text{ km}$  bins within 2 minutes of radar volume start time with reflectivity from ARMOR (solid black contours every 10-dB, starting at 10-dBZ) and vertical velocity (blue contours starting at  $5 \text{ m s}^{-1}$ , then in  $10 \text{ m s}^{-1}$  increments after  $10 \text{ m s}^{-1}$ ) overlaid. Panel C is mean flash footprint ( $\text{km}^2$ ) within 2 minutes of radar volume start time with with reflectivity from ARMOR (solid black contours every 10-dB, starting at 10-dBZ) and vertical velocity (blue dashed contours starting at  $5 \text{ m s}^{-1}$ , then in  $10 \text{ m s}^{-1}$  increments after  $10 \text{ m s}^{-1}$ ) overlaid.

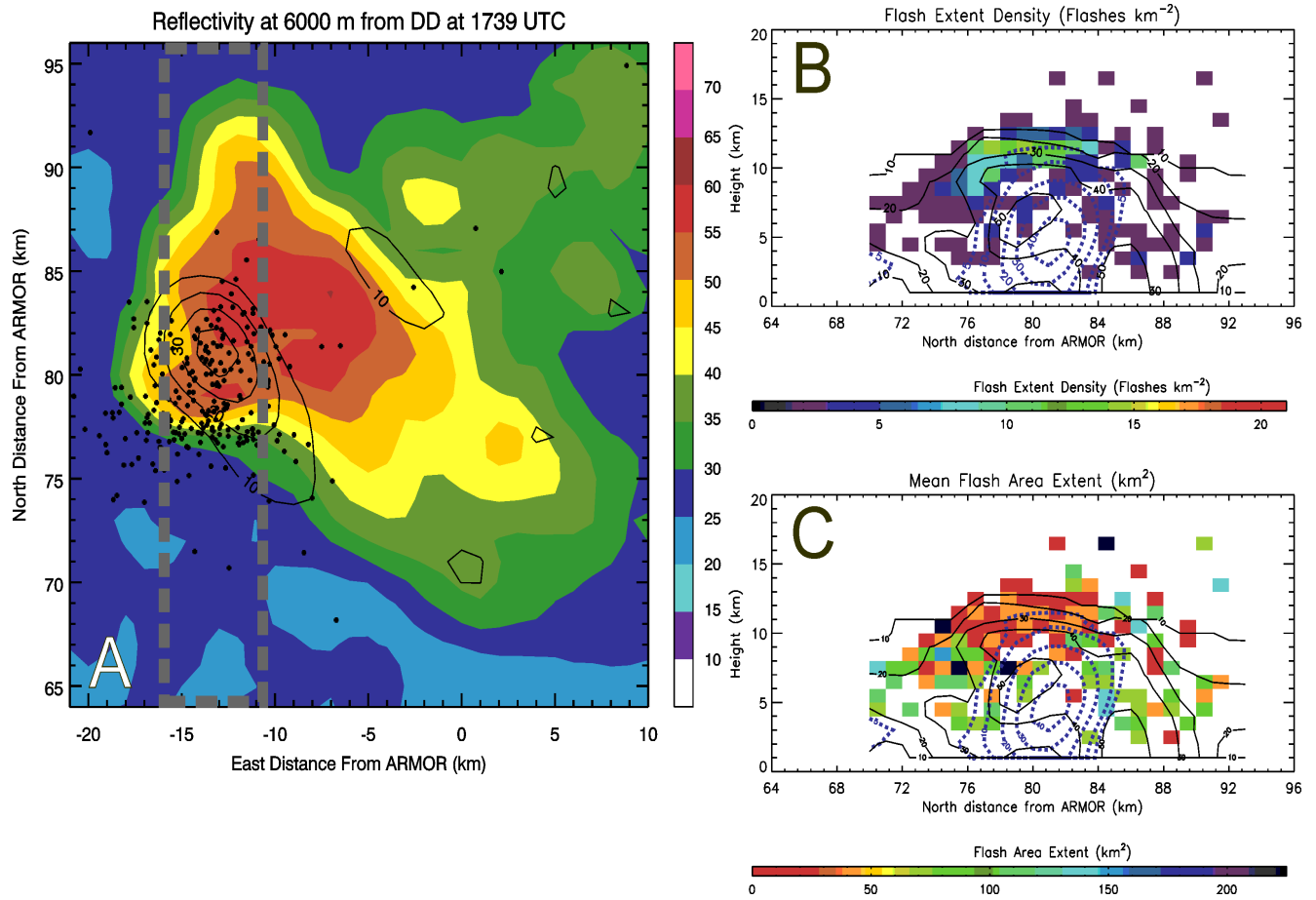


FIG. 6. Same as Fig. 5 but for 1739 UTC on 10 April 2009.