

Search for Cosmological μ -Variation from High-Redshift H_2 Absorption; A Status Report

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Abstract Observations of H_2 spectra in the line-of-sight of distant quasars may reveal a variation of the proton-electron mass ratio $\mu = m_p/m_e$ at high redshift, typically for $z > 2$. Currently four high-quality systems (Q0347-383, Q0405-443, Q0528-250 and J2123-050) have been analyzed returning a constraint $\Delta\mu/\mu < 1 \times 10^{-5}$. We present data and a μ -variation analysis of another system, Q2348-011 at redshift $z_{abs} = 2.42$, delivering $\Delta\mu/\mu = (-1.5 \pm 1.6) \times 10^{-5}$. In addition to observational data the status of the laboratory measurements is reviewed. The future possibilities of deriving a competitive constraint on $\Delta\mu/\mu$ from the known high-redshift H_2 absorbers is investigated, resulting in the identification of a number of potentially useful systems for detecting μ -variation.

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1 Cosmological μ Variation and the Laboratory Database

The proton-electron mass ratio $\mu = m_p/m_e$ is one of the constants suitable for targeting a possible variation of a physical constant on a cosmological time scale. Although μ is a dimensionless constant, lending itself for an observational approach, it is not a *fundamental* constant in a strict sense, like the fine-structure constant α , because μ involves all the binding physics of the quarks inside the proton. However, its close connection to m_e/Λ_{QCD} makes it the parameter determining the measure of the strong force with respect to the electroweak sector, and is therefore of fundamental importance. Variation of μ can be straightforwardly derived from a comparison between spectral lines observed at vastly differing redshifts. Values of wavelengths at high redshift λ_z^i are compared to laboratory wavelengths λ_0^i for as many as possible H₂ lines [40] using sensitivity coefficients for each individual spectral line $K_i = d \ln \lambda_i / d \ln \mu$. These K_i are calculable to the 1% accuracy level [43].

The laboratory wavelengths of the spectral lines in the Lyman and Werner bands of H₂ have been investigated at ever improved accuracy using tunable vacuum ultraviolet (VUV) lasers since over a decade [11, 13, 17, 35, 37, 44], also including spectra of HD [12, 16]. In addition the novel Fourier-transform spectrometer in the VUV range at the Soleil synchrotron was used to perform direct absorption spectra focusing on HD to cover the entire range of Lyman and Werner bands for this deuterated molecule [15]. The most accurate wavelength positions are obtained from a third indirect spectroscopic method based on the highly accurate determination of level energies for certain anchor levels [9], where the determination of Lyman and Werner wavelengths results from an independent measurement of spacings between excited states [38]. The full set of recommended wavelengths is reported in Bailly *et al.* [4]. Another compilation of these data, both for H₂ and HD is given in the appendix of Malec *et al.* [24], also including a listing of K_i sensitivity coefficients [15, 43], oscillator strengths of the absorption lines [2], and of damping coefficients [1].

A quasar absorption spectrum can be treated in different ways to detect a possible $\Delta\mu/\mu$. A straightforward method is to derive line positions by fitting H₂ resonances, and thereupon compare the results with the laboratory wavelengths; by deriving reduced redshifts $\zeta_i = (z_i - z_A)/(1 + z_A)$ with z_i the redshift determined for each line ($z_i = (\lambda_z^i/\lambda_z^0) - 1$) and z_A the redshift of the absorber, such a method gives a graphical insight into a possible variation of μ [37, 43]. In contrast to this line-by-line fitting method a *comprehensive* fitting method has advantages. There, an H₂ spectrum is generated, based on the available molecular parameters (wavelengths [4, 24], oscillator strengths [2], and damping parameters [1]). The value for $\Delta\mu/\mu$ is then a single fitting parameter comparing the observed spectrum with a model spectrum; this is accomplished via least squares fitting with the VPFIT software [47]. Each line is convolved with the instrumental profile (pertaining to the spectral resolving power $R = \lambda/\Delta\lambda$ of the spectrometer), a natural linewidth Γ , and a Doppler parameter b . The main advantage of the method is that a composite velocity structure can be addressed. The column densities $N(J)$ for each rotational state can be tied between velocity components or left independent. Further details

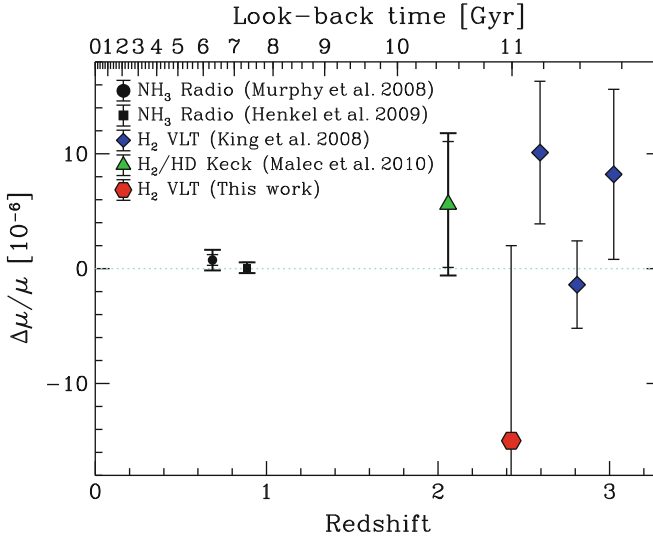


Fig. 1 Current result on cosmological μ -variation based on the five available high-quality H_2 absorption systems as a function of redshift: *diamonds* relate to Q0347–383, Q0405–443 and Q0528–250 [19], *triangle* is the result from J2123–005 [24], while the black points relate to the ammonia data [10, 25] for redshifts $z < 1$. The hexagonal point represents the result of the present analysis of Q2348–011

and advantages of the comprehensive fitting method are discussed in Refs. [19, 24]. By this means the spectral information on overlapping lines and of saturated lines can be included. In the example of the analysis of the Keck-HIRES spectrum of the quasar J2123–005 it is shown how the velocity structure can be investigated in much detail in terms of computed *composite residuals* which demonstrate a complex velocity structure that should be accounted for [24]. These optimized comprehensive fitting methods were also implemented in the analyses of the three systems Q0347–383, Q0405–443 and Q0528–250 by King *et al.* [19].

The resulting constraints on $\Delta\mu/\mu$ for the four systems mentioned are shown in Fig. 1. The figure also includes the result of the Q2348–011 analysis, details of which will be presented in the next section. The tighter constraints stemming from radio-astronomical observations of NH_3 at redshifts $z < 1$ are shown as well [10, 25]. The data show the present status of a possible variation of the proton-electron mass ratio to be $\Delta\mu/\mu < 1 \times 10^{-5}$ at redshifts $z = 2 - 3.5$.

2 VLT-UVES Observation of Q2348–011

H_2 absorption features associated with the Damped Lyman- α (DLA) system at $z = 2.42$ along the sightline of Q2348–011 (SDSS J235057.88–005209.8, $z_{\text{em}} \simeq 3.02$) were first reported by Petitjean *et al.* [32]. Since then this system has been the subject

of several investigations but none of them focused on μ -variation (e.g. Ledoux *et al.* [20], Noterdaeme *et al.* [31]), but rather on the molecular column densities, the metallicity and the physical conditions of the gas contained in this DLA-system. For an investigation of μ -variation new observations of well-calibrated spectra of Q2348–011 were recorded, and will be presented here.

2.1 Data

The absorption spectrum of the quasar Q2348–011 was obtained with the Ultraviolet and Visual Echelle Spectrograph (UVES) at the Very Large Telescope (VLT) of the European Southern Observatory (ESO) in Paranal, Chile. Observations were carried out in visitor mode on four consecutive nights (2007 August 18–21). Q2348–011 is not a particularly bright object ($R_{mag} = 18.31$) so long exposures were needed. An overall of 15 spectra were recorded with an exposure time between 3600 and 4800 seconds each; they make up for a combined total of 19.25 hours of observation. The slit widths were set to $0.''8 - 1.''0$ and $0.''7 - 0.''8$ in, respectively, blue and red ranges. They were kept constant for all exposures irrespective of the noticeably varying seeing conditions, which were on average $1.''2$. These settings yield a resolution of $R \sim 51700$ (in the blue-UV range where all H_2 lines are observed) and 55300 (in the red). This corresponds to a resolution element of 2.46 and 2.31 $km\ s^{-1}$ in terms of velocity units, respectively. For optimum wavelength calibration the 15 science exposures were each attached with a ThAr calibration frame, obtained immediately after each science exposure. Special care was taken to use pre-selected Th-Ar reference lines following optimization procedures [26]. The final spectrum of Q2348–011 covers (vacuum-heliocentric) wavelengths 3572–9467 Å, with gaps at 4520–4621 and 7505–7665 Å. Parts of the spectrum are displayed in Fig. 2.

2.2 Analysis

The wavelength range covered by the VLT-UVES spectrum provides 58 H_2 transitions for rotational levels $J = 0 - 5$. The H_2 features fall in the wavelength range between 3580 and 3860 Å for this absorber at $z = 2.42$. This part of the combined spectrum has a signal-to-noise ratio (S/N) of 25. Note that no associated HD spectral features have been detected at a significant level. All the transitions observed arise from the Lyman band. Molecular hydrogen is present in 7 velocity features (the two left-most are blended together but they count at least as two), so in principle this spectrum would provide a sample of 58×7 absorption lines. However, besides the DLA at $z = 2.42$ (neutral content $\log N(H\ I) = 20.50 \pm 0.10$, [20]), which contains molecular hydrogen, there is an additional DLA at $z = 2.62$, which has a larger column density of neutral hydrogen ($\log N(H\ I) = 21.30 \pm 0.08$ [31]), but no signs of molecular hydrogen. The strong Lyman- β feature of the second DLA ($z = 2.62$) at

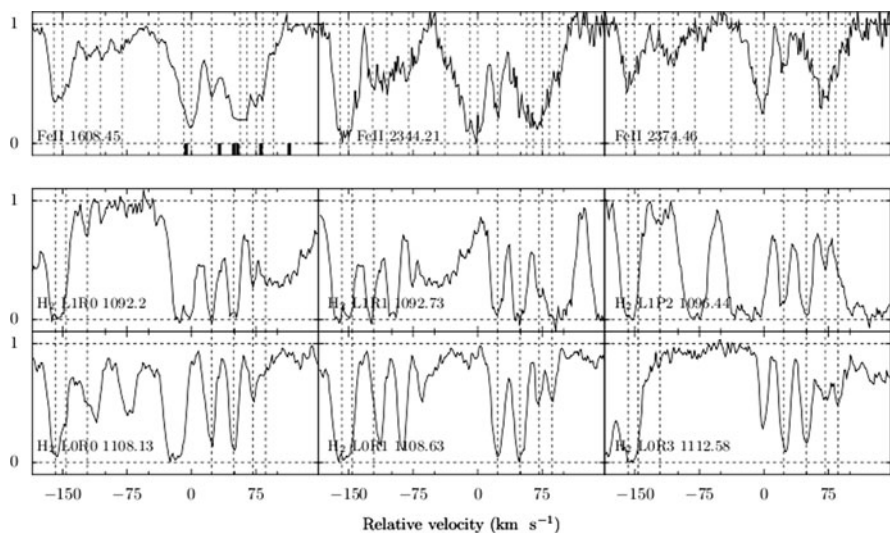


Fig. 2 Lower panel: H_2 velocity structure in the Q2348–011 DLA system at $z = 2.42$; upper panel: Fe II velocity structure of the same absorber. The positions of 7 different velocity features of H_2 are indicated by the blue dashed lines, and the positions of 14 velocity components of Fe II are indicated by the dashed lines. Note that the strongest velocity component in Fe II absorption profile (at 0 km s^{-1}) does not have an associated H_2 component. The cleanest Fe II transition (1608.45 \AA) is blended with Si IV (1402.77 \AA) of an additional absorber at $z = 2.93$. The ticks in the Fe II panel indicate positions of the overlapping Si IV lines

3710 \AA obscures 4 H_2 transitions (24 velocity components) and damps some others (see Fig. 3), while also the Lyman- γ blends 4 H_2 lines (28 velocity components). This DLA gives rise to a cutoff in the absorption spectrum at 3580 \AA , which makes this DLA an unfortunate coincidence along the sightline of Q2348–011. In addition there exist strong H I absorbers, almost at the column densities of sub-DLA's, at $z = 2.73$ and $z = 2.93$ obscuring several H_2 lines.

Another distinctive feature of this spectrum is self-blending of H_2 : since the velocity structure of H_2 consists of numerous components spread over $\sim 300 \text{ km s}^{-1}$, nearby lying lines exhibit mutual overlap of velocity components. For example, velocity components 4-7 of H_2 L0R1 lie in the same region as velocity components 1-3 of H_2 L1P5. This kind of overlapping can be disentangled using the comprehensive fitting method, if a sufficient number of non-overlapping lines is available. However, the non-overlapping lines are blended in addition with the absorbers of the Lyman- α forest, and the overall amount of clear features is strongly reduced.

Some of the molecular hydrogen features are contaminated by metal lines associated with the two DLAs, namely, L2R0 with Si II (1020.69 \AA) at $z = 2.62$, and L3R1, L3P1, L1P2, L1R3, L0P2, L0R3 are blended with various Fe II lines associated with the $z = 2.42$ DLA. Further, L1R5 and L1P4 may be blended with an additional C IV system at $z = 1.44$ but this is difficult to reconstruct unambiguously.

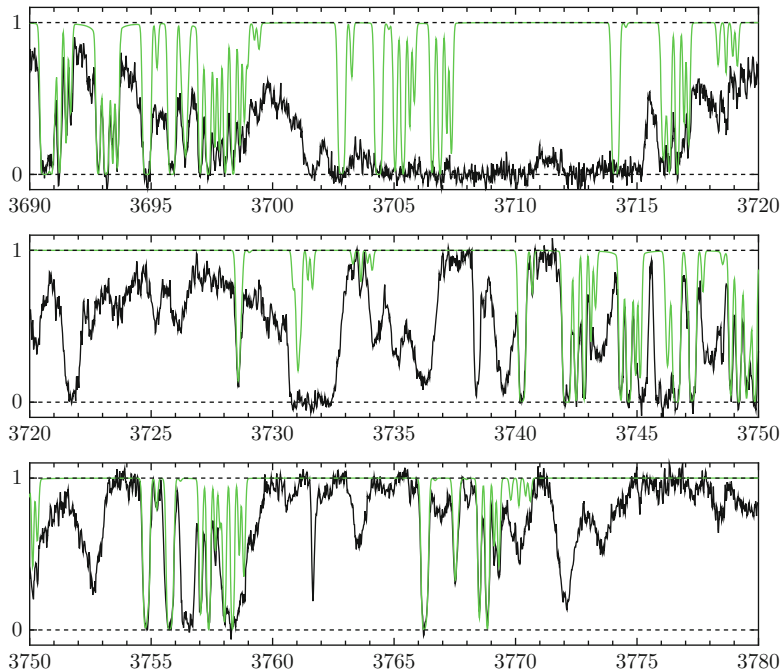


Fig. 3 H_2 absorption model superimposed on a portion of the Q2348–011 spectrum. The strong feature centered at 3710 Å is Lyman- β of the DLA at $z = 2.62$

Generally, the sightline towards the quasar Q2348–011 is rich in interesting absorbers: besides the two DLAs there are other strong metal absorbers at $z = 2.97$, 2.93, 2.72, 2.58 (C IV, Si IV, N V), and at $z = 0.86$, 0.77 (Mg II). The latter is remarkable in having numerous components spread over 750 km s^{-1} . Noteworthy coincidence is that the cleanest Fe II transition (1608.45 Å) associated with $z = 2.42$ is blended with Si IV (1402.77 Å) of the absorber at $z = 2.93$ (see Fig. 2). Also note that the most prominent velocity component in metal transitions does not have an associated H_2 component.

For the final fit deriving $\Delta\mu/\mu$ 35 regions with 36 different molecular hydrogen transitions have been selected. Each line is modeled as a Voigt profile. In the comprehensive fitting method the parameters (column density N , Doppler parameter b , and position z) are connected for different lines. Each J -level for H_2 has a different ground state population, so all transitions from the same J -level have the same value for N , b , and z . It is also assumed that a given velocity component has the same z and b value in all J -levels; physically this can be interpreted that the molecular hydrogen of all states resides in the same absorbing cloud. The value for $\Delta\mu/\mu$ is a single fitting parameter for all molecular hydrogen transitions. The absorption model is optimized to fit data by using the VPFIT software.

As analyzed by Malec *et al.* in the example of the J2123–005 quasar system, more than one velocity component can underlie a molecular absorption feature as

distinguishable by inspection. Their presence can be confirmed or rejected relying on goodness-of-fit and by inspecting *composite residuals* [24]. As the two left-most features of H₂ absorption are blended, more components could in principle be present close to or between the two blended features. A fit with 7+1 velocity components (VC) instead of 7, where the additional VC was included between the two left-most VCs, has improved the χ^2_{ν} per degree-of-freedom by several hundredths, but the additional velocity component was rejected as unnecessary in transitions from $J = 4$ and $J = 5$ rotational levels. The 7+1 VC model is regarded as the fiducial result, because it gives the best match with the observed spectrum, i.e. it is statistically preferred as it returns the smallest χ^2_{ν} . A more detailed analysis will be presented in an upcoming paper [3].

3 Results

The fiducial absorption model led to the following result: $(\Delta\mu/\mu)_{7+1VC} = (-1.5 \pm 1.6) \times 10^{-5}$. Table 1 provides molecular cloud properties with their statistical uncertainties as derived from the final fitting attempt when $\Delta\mu/\mu$ was set as a free parameter.

The analysis of H₂ absorption in the Q2348–011 spectrum yields a competitive result on $\Delta\mu/\mu$, but it is three times less tight compared to the constraints from previous analyses (see Fig. 1). One of the major causes that led to the relatively large uncertainty is the low S/N of the spectrum. The second cause is that the spectrum is obscured by the neutral hydrogen features of the additional strong DLA at $z = 2.62$. The H₂ transitions falling at the shortest wavelengths covered by the spectrum are especially useful in μ -variation analysis, since they exhibit the larger K -coefficients (transitions in the Lyman band with higher vibration excitation). In the spectrum of Q2348–011 they are not detected due to the H I absorption produced by the additional DLA. Although the total number of H₂ transitions available in this spectrum is reduced, it is at least partially compensated by the presence of numerous velocity components associated with each transition. Altogether, each transition contributes a higher information content when it is imprinted in several velocity components. The complex absorption structure was successfully modelled by means of the comprehensive fitting method.

4 Potential High-Redshift H₂ Absorbers for Future Detection of μ -Variation

Detection of cosmological μ -variation, or putting strong constraints on the variation, should preferably be based on large numbers of high redshift H₂ absorbing systems. For comparison the recent study on a possible spatial dipole for the fine structure constant α relies on the analysis of almost 300 absorbing systems [48].

Table 1 Parameters and $1\text{-}\sigma$ statistical uncertainties pertaining to the $7 + 1$ VC fiducial fit model for the molecular hydrogen absorption in the Q2348-011 spectrum. The 8 velocity components are at redshifts z and have Doppler widths b . The second row gives the velocity relative to $z = 2.426318$ in units of km s^{-1}

VC	z Δv [km s^{-1}]	b [km s^{-1}]	J = 0	J = 1	J = 2	J = 3	J = 4	J = 5
			$\log N(\text{H}_2)$ [cm^{-2}]					
1a	2.424509(1)	6.73 ± 0.15	15.27 ± 0.15	16.45 ± 0.05	15.50 ± 0.14	15.92 ± 0.07	15.30 ± 0.03	15.12 ± 0.02
1b	2.424546(4)	9.94 ± 0.29	15.45 ± 0.08	15.53 ± 0.40	15.79 ± 0.07	15.67 ± 0.10	–	–
2	2.424652(2)	1.25 ± 0.24	14.97 ± 0.12	16.98 ± 0.50	14.37 ± 0.27	14.77 ± 0.21	14.20 ± 0.06	13.62 ± 0.13
3	2.424929(2)	4.15 ± 0.33	14.29 ± 0.03	14.47 ± 0.05	14.51 ± 0.03	13.67 ± 0.34	12.95 ± 1.00	–
4	2.426582(1)	4.56 ± 0.13	15.66 ± 0.02	15.84 ± 0.10	15.60 ± 0.03	15.62 ± 0.04	14.29 ± 0.04	13.81 ± 0.11
5	2.426883(1)	2.97 ± 0.08	16.24 ± 0.08	17.98 ± 0.04	17.27 ± 0.10	16.99 ± 0.13	14.31 ± 0.08	14.08 ± 0.07
6	2.427136(1)	3.10 ± 0.12	14.72 ± 0.05	15.22 ± 0.03	14.93 ± 0.04	14.86 ± 0.06	14.09 ± 0.06	13.67 ± 0.16
7	2.427310(1)	4.36 ± 0.22	14.45 ± 0.05	15.27 ± 0.03	14.75 ± 0.05	15.09 ± 0.05	14.29 ± 0.04	13.89 ± 0.16

However, for H_2 , up to now only five systems (Q0347–383, Q0405–443, Q0528–250, J2123–050 and Q2348–011) of sufficient quality are found, and the results obtained in comprehensive fitting procedures are summarized in Fig. 1. We note that alternative line-by-line fitting procedures on the Q0347–383 and Q0405–443 systems have recently been reported [41, 49, 50]. Taking the weighted mean of the H_2 results of Fig. 1 yields $\Delta\mu/\mu = (2.9 \pm 2.6) \times 10^{-6}$ for redshift ranges between $z = 2.0 - 3.1$, which may be considered as the status by September 2010. The data provide firm evidence that μ -variation is below the 10^{-5} level for look-back times in the range of 10–12 Gyrs.

A recent analysis of the J1337+315 spectrum by Srianand *et al.* [31] yielded a constraint only at the 10^{-4} level. This is due to the combination of low brightness of the quasar ($R_{mag} = 18.08$) and low H_2 content at the absorbing cloud ($\log_{10}N(H_2) = 14.1$). The present analysis of the Q2348–011 H_2 absorption system demonstrate, that even when the S/N of a spectrum is low, a competitive constraint on $\Delta\mu/\mu$ can still be delivered if other circumstances are favorable. In case of Q2348–011 the low S/N is compensated by the presence of at least 7 sufficiently strong velocity components. However, in this regard the Q2348–011 absorbing system is rather an exception than a rule. In other known DLAs H_2 absorption is usually detected in no more than three visually distinguishable spectral features.

This raises the question how many of the thousands of quasar systems on the sky (see Fig. 4 and [46]), of which there are some 1000 identified as DLAs could produce a competitive result, i.e. at the level $\delta(\Delta\mu/\mu) \leq 5 \times 10^{-6}$.

Based on our experience in analyzing spectra, we postulate that high quality H_2 absorption systems should obey the following conditions:

- An apparent magnitude of $V_{mag} < 18$ of the quasar background source for providing a SNR of about 50 after a reasonable amount of data collection (15–20 h).
- A column density of $N(H_2)$ preferably in the range between 10^{14} and 10^{18} cm^{-2} yielding sufficient but not too saturated H_2 absorption.
- An H_2 absorption system at redshift $z > 2$; only then a sufficient number of lines will shift into the atmospherically transparent window and the throughput

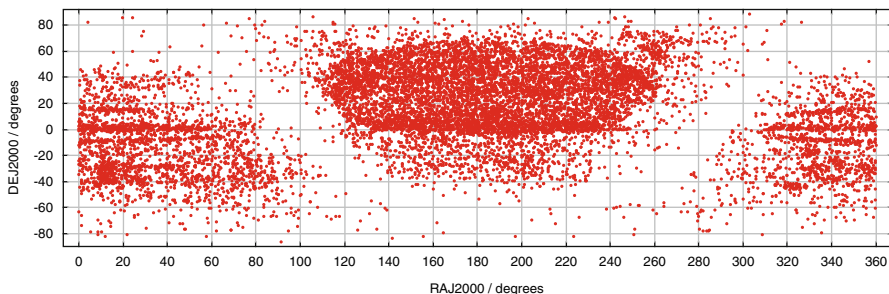


Fig. 4 Known quasars with magnitude $V < 18$ as obtained from [46]

window of Keck-HIRES and VLT-UVES; preferably some Werner lines should be observed for dealing with systematics [24].

- No occurrence of additional DLA's and only a limited amount of metal absorption systems along the same line of sight; the Q2348-011 absorber is an example of an unfavorable situation with an additional strong DLA and a few very strong HI absorbers, as well as several metal bearing systems.
- In principle, a firm constraint on the velocity structure of the H₂ absorber is only possible with the comprehensive fitting method. The present example of Q2348-011 demonstrates the applicability of this method in case of a complex structure with 7 velocity components. For line-by-line fitting methods the H₂ spectrum should preferably exhibit a single absorber; that a, e.g. two-component model, is statistically worse than a single component model should be demonstrated.

We have inspected the existing literature on high-redshift H₂ absorbers in order to identify possibly useful systems that may be targeted in future for μ -variation. In Table 2 all systems in which H₂ has been detected so far are listed. Relevant details for all systems, such as absorption and emission redshifts, their position on the sky,

Table 2 Listing of known high redshift H₂ absorption systems with some relevant parameters. Bessel R photographic magnitude taken from the SuperCOSMOS Sky Survey [8]. The five high-quality systems analyzed so far are offset. The column densities $N(\text{H}_2)$, $N(\text{HD})$ and $N(\text{H I})$ are given on a \log_{10} scale in $[\text{cm}^{-2}]$

Quasars	z_{abs}	z_{em}	RA(J2000)	Decl.(J2000)	$N(\text{H}_2)$	$N(\text{HD})$	$N(\text{H I})$	R_{mag}	Ref.
Q0347-383	3.02	3.21	03:49:43.64	-38:10:30.6	14.5		20.6	17.48	[19, 37]
Q0405-443	2.59	3.00	04:07:18.08	-44:10:13.9	18.2		20.9	17.34	[19, 37]
Q0528-250	2.81	2.81	05:30:07.95	-25:03:29.7	18.2	^a	21.1	17.37	[19]
J2123-005	2.06	2.26	21:23:29.46	-00:50:52.9	17.6	13.8	19.2	15.83	[24]
Q2348-011^b	2.42	3.02	23:50:57.87	-00:52:09.9	18.4		20.5	18.31	[20, 31] ^c
Q0013-004	1.97	2.09	00:16:02.40	-00:12:25.0	18.9		20.8	17.89	[34]
HE0027-184	2.42	2.55	00:30:23.62	-18:19:56.0	17.3		21.7	17.37	[27]
Q0551-366	1.96	2.32	05:52:46.18	-36:37:27.5	17.4		20.5	17.79	[23]
Q0642-506	2.66	3.09	06:43:26.99	-50:41:12.7	18.4		21.0	18.06	[28]
FJ0812+320	2.63	2.70	08:12:40.68	+32:08:08.6	19.9	15.4	21.4	17.88	[5, 42]
Q0841+129	2.37	2.48	08:44:24.24	+12:45:46.5	14.5		20.6	17.64	[33]
Q1232+082	2.34	2.57	12:34:37.58	+07:58:43.6	19.7	15.5	20.9	18.40	[14, 45]
J1237+064	2.69	2.78	12:37:14.60	+06:47:59.5	19.2 ^d	14.5	20.0	18.21	[29]
Q1331+170^e	1.78	1.78	13:33:35.81	+16:49:03.7	19.7	14.8	21.2	16.26	[5, 6]
Q1337+315	3.17	3.17	13:37:24.69	+31:52:54.6	14.1		21.4	18.08	[39]
Q1439+113	2.42	2.58	14:39:12.04	+11:17:40.5	19.4	14.9	20.1	18.07	[30]
Q1443+272	4.22	4.42	14:43:31.18	+27:24:36.4	18.3		21.0	18.81	[22]
Q1444+014	2.08	2.21	14:46:53.04	+01:13:56.0	18.3		20.1	18.10	[21]
Q2318-111	1.99	2.56	23:21:28.69	-10:51:22.5	15.5		20.7	17.67	[27]
Q2343+125	2.43	2.52	23:46:25.42	+12:47:43.9	13.7		20.4	20.18	[7, 32]

^a A recent re-analysis has shown that this system contains HD [18].

^b Multiple DLAs in this system.

^c Studied in this work for μ -variation.

the obtained column densities for H_2 , HD and H I, and the magnitude are listed as well. Inspection of these systems leads to a quality assessment of the known H_2 absorber systems. Some of the systems can be discarded based on their low redshift (Q0551 – 366, Q2318 – 111, Q0013 – 004 and Q1331 + 170), on their low H_2 column density (Q2343 + 125 and Q1337 + 315), too low brightness, or combinations thereof. It should be noted that this assessment is linked to the status of present day technology. Future larger telescopes, or space-based telescopes equipped with high resolution spectrometers ($R > 50\,000$) may alter this view.

Based on the considerations given we have identified five quasar systems that have the potential to achieve a constraint at the $\delta(\Delta\mu/\mu) \leq 5 \times 10^{-6}$ level. These systems have been investigated to some extent, but larger numbers of exposures yielding better S/N, and ThAr attached calibration frames will be needed to reach the desired accuracy. Of the five systems, four have been observed with VLT-UVES, while one (FJ0812+320) has been studied with Keck-HIRES.

- **Q0642–506** with redshift $z_{abs} = 2.66$ provides a good number of H_2 lines with a high column density giving rise to saturation of only the strongest lines. The occurrence of only a single velocity components renders the spectrum relatively simple, also accessible for a line-by-line treatment. Spectra are shown in [28].
- **HE0027–184** is of similar good quality. It has a favorable magnitude and a molecular column density which allows for observation of a large number of lines to be observed. Indeed in the archived spectra with data observed in 2004 many lines are found, possible even up to $J = 6$ [27].
- **FJ0812+320** has a strong saturated H_2 component, which makes it less ideal. The system has been studied in particular for its abundant presence of the widest variety of metal lines [36] and of HD [5,42]. The system is nevertheless of importance because it is the only system of reasonable quality at northern declination, except for the special $z = 4.22$ system discussed below.
- **J1237+064** at $z_{abs} = 2.69$ shares some of the disadvantages of the previous system, in particular the very high column density giving rise to saturation. However, the velocity structure exhibits two weaker components, such that high- J lines can be obtained in the strong components, and the low- J lines in the weak components; HD is abundantly present and can be included in the analysis. This makes J1237+064 a target with a high potential. Further details on the system in [29].
- **Q1443+272** is a special case, and its unusually high redshift ($z_{abs} = 4.22$) alters some of the arguments. The system has a very low brightness, but that is partly compensated because the H_2 lines shift further into the visible, where scattering is less and CCD's are more sensitive. In addition the high redshift shifts many more H_2 lines into the observable range. However, for these high redshift objects the Lyman- α forest becomes denser, which is a disadvantage. In summary this system has a good potential, if sufficient observation time is made available. Further details on the system in [22].

5 Conclusion

We have reviewed the status of investigations into variation of the proton-electron mass ratio from analyses of highly redshifted H₂ absorption systems. As of now five results have been obtained putting a tight constraint at the $\Delta\mu/\mu < 1 \times 10^{-5}$ level. We present spectroscopic data on the fifth H₂ absorption system, Q2348–011, and have show how a complex velocity structure can be unraveled using the comprehensive fitting method. An attempt is made to list in full all relevant quasar systems where H₂ has been observed so far, followed by an assessment as to which systems could produce a competitive result on $\Delta\mu/\mu$, provided that high-quality well-calibrated observations are performed.

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