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Published in: Monthly Notice of the Royal Astronomical Society

DOI: 10.1093/mnras/stv992

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Document Version Publisher's PDF, also known as Version of record

Publication date: 2015

Link to publication in University of Groningen/UMCG research database

Citation for published version (APA): Yildiz, M. K., Serra, P., Oosterloo, T. A., Peletier, R. F., Morganti, R., Duc, P-A., ... Karabal, E. (2015). Star formation in the outer regions of the early-type galaxy NGC 4203. Monthly Notice of the Royal Astronomical Society, 451(1), 103-113. https://doi.org/10.1093/mnras/stv992

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Star formation in the outer regions of the early-type galaxy NGC 4203

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Accepted 2015 May 1. Received 2015 April 28; in original form 2015 February 24

ABSTRACT

NGC 4203 is a nearby early-type galaxy surrounded by a very large, low-column-density H I disc. In this paper, we study the star formation efficiency in the gas disc of NGC 4203 by using the UV, deep optical imaging and infrared data. We confirm that the H I disc consists of two distinct components: an inner star-forming ring with radius from ~ 1 to $\sim 3 R_{eff}$ and an outer disc. The outer H I disc is nine times more massive than the inner H I ring. At the location of the inner H I ring, we detect spiral-like structure both in the deep g' - r' image and in the 8 μ m *Spitzer*-Infrared Array Camera image, extending in radius up to $\sim 3 R_{eff}$. These two gas components have a different star formation efficiency likely due to the different metallicity and dust content. The inner component has a star formation efficiency very similar to the inner regions of late-type galaxies. Although the outer component has a very low star formation efficiency, it is similar to that of the outer regions of spiral galaxies and dwarfs. We suggest that these differences can be explained with different gas origins for the two components such as stellar mass loss for the inner H I ring and accretion from the inter galactic medium for the outer H I disc. The low-level star formation efficiency in the outer H I disc is not enough to change the morphology of NGC 4203, making the depletion time of the H I gas much too long.

Key words: galaxies: elliptical and lenticular, cD – galaxies: evolution – galaxies: individual: NGC 4203 – galaxies: ISM – galaxies: structure.

1 INTRODUCTION

Although we have long known that early-type galaxies (ETGs) can harbour a stellar disc, their actual distribution of bulge-to-disc ratios B/D has been debated for many decades (e.g. Spitzer & Baade 1951; Sandage, Freeman & Stokes 1970; van den Bergh 1976). Recently, a number of studies based on broad-band imaging and integral-field spectroscopy at optical wavelengths have concluded that most ETGs host a stellar disc, and that ETGs as a family cover a similar range of B/D ratio as spirals (Cappellari et al. 2011; Laurikainen et al. 2011; Kormendy & Bender 2012; Krajnović et al. 2013; Weijmans et al. 2014). Discs are detected also with a number of additional techniques such as narrow-band imaging of emission lines (e.g. H β), imaging of dust absorption, or the observation of the kinematics of planetary nebulae and globular clusters (Romanowsky et al. 2003; Sarzi et al. 2006; Pota et al. 2013). Taken together, these results suggest that galaxy evolution is dominated by processes which allow the growth and survival of stellar discs in galaxies of both early and late types.

Discs have also been found in the cold gas phase of ETGs by means of both neutral hydrogen (H1) and CO observations (e.g. van Driel & van Woerden 1991; Young 2002; Morganti et al. 2006; Oosterloo et al. 2007, 2010; Crocker et al. 2011). More recently, ATLAS^{3D} observations of H1 and CO have shown that \sim 50 per cent of all ETGs contain some cold gas (Young et al. 2011; Serra et al. 2012, hereafter S12). Molecular gas is typically found in small gas discs in the central regions and is linked to small amounts of star formation (SF; Alatalo et al. 2013; Davis et al. 2013). In contrast, in \sim 20 per cent of ETGs outside clusters, H1 has been found distributed in low-column-density discs or rings with typical

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Figure 1. SDSS g-band (left) and new deep WSRT H I (right: see Section 2.1) images of NGC 4203. The size of the images is \sim 74 × 74 kpc² (17 arcmin × 17 arcmin). The colour bar refers to the H I image. The H I beam is shown in the bottom-left corner of the H I image. The blue boxes represent the size of the images shown in Fig. 2.

sizes of many tens of kpc, much larger than the stellar body (S12). If the column density of the gas is high enough, these H_I discs are potential fuel for SF. Contrary to the general idea that galaxies evolve from late to early type, these H_I reservoirs could generate the growth of stellar discs and therefore trigger the transition of the host from an early- to a late-type morphology (Cortese & Hughes 2009; Fang et al. 2012; Salim et al. 2012).

In this context, many recent studies have focused on the presence of SF in the outer discs of ETGs (Salim & Rich 2010; Lemonias et al. 2011; Moffett et al. 2012; Salim et al. 2012). Such SF is typically traced by UV light emitted by massive stars of age up to several hundred Myr (Kauffmann et al. 2007; Bianchi 2011; Salim et al. 2012). Moffett et al. (2012) show that ~40 per cent of 38 red sequence and blue cloud ETGs with stellar mass below ~4 $\times 10^{10}$ M_{\odot} have extended-UV discs, which is evidence for disc growth in these objects.

In this paper, we aim to bring these two lines of investigation (H₁ and UV) together and study SF in the large H₁ disc of the ETG NGC 4203. This is a nearby, nearly face-on lenticular. Fig. 1 shows an SDSS *g*-band image of the galaxy (left-hand panel), and Table 1 provides some basic information. NGC 4203 has been observed at many wavelengths. Optical studies show that NGC 4203 has a relatively massive stellar disc (B/D = 0.4–0.5; Burstein 1979; Krajnović et al. 2013). Early H₁ observations revealed a low-columndensity H₁ disc with a complex morphology (Burstein & Krumm 1981; van Driel et al. 1988). More recent observations show that the H₁ disc extends to very large radius and is characterized by spiral-like H₁ arms (S12). This gas rotates regularly around the galaxy and is kinematically aligned with the stellar disc (Serra et al. 2014).

Despite the large H_I reservoir ($\sim 10^9 M_{\odot}$), the stellar population is uniformly old ($\sim 10 \text{ Gyr}$) inside 1 R_{eff} (McDermid et al. 2015) suggesting that the H_I is not fuelling a high level of SF in this region. Yet, there are signs of possible low-level SF within the stellar body. For instance, NGC 4203 contains 2.5 × 10⁷ M_☉ of H₂ in the centre (Welch & Sage 2003; Young et al. 2011). *Hubble Space Telescope* observations show nuclear dust structures in the very centre of the

Table 1. General properties of NGC 4203.

Quantity	Value	Unit	Reference
Hubble type	SAB0	_	(1)
RA (J2000)	12h15m05s06	HMS	(2)
Dec. (J2000)	+33 ^d 11 ^m 50 ^s 382	DMS	(2)
Helio. rad. vel.	1086 ± 4	$\rm km~s^{-1}$	(1)
Distance	14.7	Mpc	(3)
M_K	-23.4	mag	(3)
A _B	0.05	mag	(3)
Effective radius	29.5	arcsec	(3)
$\log(M/L)_{\rm star}$	0.5	$M_{\odot}/L_{\odot r}$	(4)
$\log M_{\rm star}$	10.6	M	(4)
Galactic $E(B - V)$	0.012	mag	(5)

References: (1) de Vaucouleurs et al. (1991); (2) Anderson & Ulvestad (2005); (3) Cappellari et al. (2011); (4) Cappellari et al. (2013); (5) Schlegel, Finkbeiner & Davis (1998).

galaxy (300 pc radius; Erwin & Sparke 2003). The dust is visible also on a large scale in an 8 μ m image, where the flux, mostly from polycyclic aromatic hydrocarbons (PAHs), appears distributed in a spiral-like pattern (Pahre et al. 2004; we will make use of this image as well as of a 24 μ m image in support of our analysis of SF in NGC 4203).

Here we use new, deep H₁ observations and UV imaging from the *Galaxy Evolution Explorer* (*GALEX*) to study the relation between H₁ and SF as well as the impact of this large gas reservoir on NGC 4203. In Section 2, we describe the data (H₁, FUV, infrared as well as ancillary optical imaging) and the methodology used for our work. In Section 3, we discuss global properties of the H₁, FUV, infrared and optical images, present the derivation of star formation rate (SFR) from the data and compare the SF efficiency in the H₁ disc with that of late-type and dwarf galaxies. In Section 4, we discuss the origin of the H₁ gas and its impact on the host galaxy. In Section 5, we give a summary of the work.

2 DATA AND ANALYSIS

2.1 WSRT data reduction

We select NGC 4203 from the sample of ETGs with extended H I discs presented by S12. We observe this galaxy for 9×12 h using the Westerbork Synthesis Radio Telescope (WSRT). Compared to the original 12 h observation by S12, this longer integration allows us to image the H I down to a similar column density but at a significantly higher angular resolution. This is desirable to compare the H I image to the UV image used to study SF (see below).

We reduce separately the nine H_I data sets (12 h integration each) following standard techniques with the MIRIAD package (Sault, Teuben & Wright 1995). Having calibrated and subtracted the continuum from the data sets, we combine them in the UV plane and image with robust = 0.4 weighting, obtaining an H_I cube with an angular resolution of 29.3×16.7 arcsec² (PA = 0.5 deg). The velocity resolution is 16 km s^{-1} after Hanning smoothing. The noise level of the data cube is $0.12 \text{ mJy beam}^{-1}$.

We create an H_I image by summing flux at all velocities using only pixels included in a mask. In order to include diffuse, faint emission in the mask, we select pixels above a certain threshold in (i) the original data cube; (ii) four data cubes smoothed spatially with a 2D Gaussian of full width at half-maximum (FWHM) 30, 45, 60 and 90 arcsec, respectively; (iii) three data cubes smoothed in velocity with a Hanning filter of width 32, 64 and 112 km s⁻¹; (iv) six data cubes obtained by smoothing in velocity the 30 and 60 arcsec resolution cubes as in point (iii). We create the mask using a 3 σ threshold except for the 90 arcsec resolution cube, for which we adopt a 4 σ threshold.

We determine a detection limit from the original data cube as a 5σ signal in a single velocity resolution element, obtaining 2.3 × 10^{19} cm⁻². We obtain the final H_I image by applying the WSRT primary beam correction to all pixels above this detection limit in the moment zero image derived from the masked H_I cube. We show the H_I image in Fig. 1 (right-hand panel) and a zoom-in on the central 300 arcsec in Fig. 2 (top-left panel).

2.2 GALEX data reduction

We obtain *GALEX* UV images of NGC 4203 from the Mikulski Archive for Space Telescopes. We use data from the Calibration Imaging, whose goals and specifications are described in Martin & GALEX Team (2005). The *GALEX* field of view (FOV) is \sim 1.2 deg in diameter and the instrument has two main bands: the far-UV (FUV) and the near-UV (NUV). The effective wavelengths for the FUV and NUV channels are 1516 and 2373 Å, with 4.5 and 6 arcsec resolution (FWHM), respectively. Details can be found in Morrissey et al. (2005, 2007). Our study is mostly based on the FUV data because the FUV emission is more sensitive to recent SF than the NUV emission (Bianchi 2011).

Since NGC 4203 has not been observed as a target galaxy, the position of the galaxy is different in all the 10 *GALEX* images. Therefore, we have selected five *GALEX* images where NGC 4203 is sufficiently far from the edge of the *GALEX* FOV to grant good image quality and reasonable background uniformity. The total exposure time of these five UV images is 1186 s. For the purpose of comparing the distribution of UV emission and H₁ gas, we analyse a 1500×1500 arcsec² *GALEX* image centred on NGC 4203. Fig. 2 shows the central 300 arcsec of the FUV image, which is the only region where emission from SF in the H₁ disc is clearly visible by eye.

We make use of the SDSS photometric catalogue (Ahn et al. 2012) to mask background and foreground objects based on their type (unresolved or extended), brightness, colour and Petrosian radius. For the unresolved sources, we adopt a radius equal to the GALEX point spread function (PSF), otherwise we use the Petrosian radius of the object itself. Following visual inspection of the UV images, we choose to mask the objects with $g' - r' \leq 1$ and 1.5, and with g < 19.5 and 21.1 mag for the unresolved and resolved sources, respectively. We estimate the local background from the masked images by applying a moving mean filter of size $300 \times$ 300 pixel^2 ($450 \times 450 \text{ arcsec}^2$). The filtering is performed iteratively by clipping three times at a 4.5σ level. To avoid overestimating the background, we mask a large area (300 arcsec diameter) around the centre of the galaxy. We subtract this background from the original, masked image. Additionally, a flat residual background is subtracted from the UV image before proceeding with our analysis (see Section 3.2).

Finally, we correct the FUV image for the effect of Galactic extinction. We use E(B - V) = 0.012 based on the dust images by Schlegel et al. (1998) and assume that $A_{FUV} = 8.24 \times E(B - V)$ and $A_{NUV} = 8.2 \times E(B - V)$ (Wyder et al. 2007). Thus, we use the following equations for the Galactic correction:

$$I_{\rm FUV,corr} = I_{\rm FUV} \times 1.1 \tag{1}$$

$$I_{\rm NUV, corr} = I_{\rm NUV} \times 1.08.$$
⁽²⁾

2.3 Ancillary optical and infrared images

We support our study of SF, which is primarily based on the UV images described above, with ancillary images at 8, 24 μ m and optical wavelengths. The 8 μ m mid-infrared emission is generally considered to trace the PAHs heated by UV photons (Leger & Puget 1984; Sellgren 1984). It is thought that 2/3 of the 8 μ m emission from a galaxy originates from heating by stellar populations 100 Myr or younger (Calzetti et al. 2005; Calzetti 2013), and that PAH emission is stronger for increasing metallicity and dust content of the interstellar medium (e.g. Engelbracht et al. 2005; Marble et al. 2010; Calzetti 2013). The 24 μ m emission is generally produced by dust following absorption of UV photons emitted by young stars, and traces SF on a time-scale of 10 Myr (Calzetti et al. 2005; Leroy et al. 2008).

In this paper, we use archival 3.6, 8 and 24 μ m images from Fazio et al. (2004) obtained with the Infrared Array Camera and the Multiband Imaging Photometer on the *Spitzer Space Telescope* (Rieke et al. 2004; Werner et al. 2004). We remove the background from each image by subtracting the average pixel value calculated in a number of empty regions.

The 8 μ m image includes a contribution from stars in the galaxy, which we remove by subtracting the 3.6 μ m image scaled by a factor $X_{3,6}$ as in Shapiro et al. (2010). These authors find that $X_{3,6}$ has very low scatter about a median value of 0.264 for ETGs, and we adopt this value for NGC 4203. The bottom-left panel of Fig. 2 shows the resulting 8 μ m non-stellar image. Similarly, the 24 μ m flux is also contaminated by emission from old stellar populations and hot, circumstellar dust (e.g. Salim et al. 2009; Kelson & Holden 2010; Martini, Dicken & Storchi-Bergmann 2013). To remove this emission, we apply the same scaling method as used for the 8 μ m. We have adopted a scaling factor from Martini et al. (2013) that study the empirical 24 μ m scaling factor for dust-free ETGs.

We also use deep optical images of NGC 4203 that are obtained in the g' and r' bands using the MegaCam camera on the



Figure 2. Zoomed-in images of the deep H_I (top-left), *GALEX*-FUV (top-right), *Spitzer* 8 µm infrared (bottom-left) and deep MegaCam g' - r' colour (bottom-right) images. All images show a field of size 300 arcsec × 300 arcsec (~22 kpc × 22 kpc). Dark areas correspond to redder regions in the deep MegaCam colour image. The red contours on the H_I and FUV images represent an H_I column density of 1.8×10^{20} cm⁻². The location of the inner H_I ring is shown by the green dashed lines.

Canada–France–Hawaii Telescope. These observations are obtained as part of the MATLAS/ATLAS^{3D} Large Programme. The observing strategy and data reduction procedures with the LSB-Elixir pipeline optimized for the detection of extended low surface brightness features are described in Duc et al. (2015). Further processing is carried out to remove the extended artificial haloes surrounding two nearby bright stars. We determine the stellar PSFs, including their extended wings, with a manual procedure (as explained in Duc et al. 2015). In the resulting image (see Fig. 3, top), the ETG exhibits a rather regular extended outer stellar halo, and a dwarf galaxy can be seen east of NGC 4203.

We compute a model of the NGC 4203 with the ellipse-fitting procedure of IRAF for both the g'- and r'-band images and then subtract these models from the original images. The resulting residual image together with the g'-band image of the dwarf galaxy is shown at the bottom panel of Fig. 3, and it is obvious that some extended spiral/ring structures show up in the main body of the lenticular galaxy.



Figure 3. Top: composite g' + r' MegaCam image of NGC 4203 with the dust lanes superimposed in red. The dust lanes are obtained by summing the residual g' + r' images after a galaxy model subtraction (see Section 2.3). This image also shows the disrupted companion dwarf galaxy at the east of the galaxy. Bottom: MegaCam g'-band image of the disrupted companion dwarf galaxy and residual map of NGC 4203. It exhibits a possible faint tail east of our galaxy which is the same direction as the two tidal tails of the companion dwarf (see Section 3.1). The white bar shows a length of 1 arcmin.

2.4 SFR calculation

We estimate the SFR using the various tracers described above (FUV, non-stellar 8 μ m and 24 μ m). In order to convert FUV emission into SFR, we use a conversion from Bigiel et al. (2010),

$$\Sigma_{\rm SFR}[M_{\odot} \rm yr^{-1} \, \rm kpc^{-2}]$$

= 0.68 × 10⁻²⁸ × I_{FUV}[erg s⁻¹Hz⁻¹ kpc⁻²], (3)

where $I_{\rm FUV}$ is the FUV intensity per unit area and the initial mass function (IMF) is assumed to be of a Kroupa type. This estimate could in principle be corrected for internal extinction on the basis of the H_I image using the relation $N({\rm H_{I}})/E(B - V) = 5 \times$ 10^{21} cm⁻² mag⁻¹ presented by Bohlin, Savage & Drake (1978). This relation is derived for the Milky Way and therefore is likely to provide an upper limit on the extinction in metal-poor environments. In the case of NGC 4203, the peak H_I column density is ~4.5 × 10^{20} cm⁻² and, thus, the maximum extinction should be around ~25 per cent. In the outer regions, where the H_I column density is much lower (~3 × 10^{19} cm⁻²), the extinction decreases to a negligible ~2 per cent. In this paper, we do not apply this correction but instead trace SF in dusty environments by using infrared imaging.

Although we do not expect a large amount of dust in the outer regions of NGC 4203, Fig. 2 shows that in the inner regions dust is present and this will affect the calculation of SFR. As mentioned

in Section 2.3, the FUV emission absorbed by dust is re-emitted at infrared wavelengths. We use the conversion from Wu et al. (2005) to estimate SFR from the 8 μ m non-stellar image:

$$\Sigma_{\rm SFR}[\rm M_{\odot} yr^{-1} kpc^{-2}] = 8.43 \times 10^{-3} I_{8\,\mu\rm m} (\rm non-stellar), \qquad (4)$$

where $I_{8 \mu m}$ is the non-stellar 8 μm intensity in MJy sr⁻¹.

We also combine the FUV and non-stellar 24 μ m images to estimate the SFR as in Leroy et al. (2008):

$$\Sigma_{\text{SFR}}[M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}] = 8.1 \times 10^{-2} I_{\text{FUV}} + 3.2 \times 10^{-3} I_{24 \text{ µm}} (\text{non-stellar}).$$
(5)

where I_{FUV} and I_{24} are in MJy sr⁻¹, and similar to equation (3) a Kroupa-type IMF is used.

Below we study the SF at the original resolution of the FUV and infrared images as well as at the resolution of the H I image in order to compare the distribution of neutral gas and SFR.

3 GLOBAL PROPERTIES OF FUV, HIAND ANCILLARY IMAGES

3.1 Flux distribution of the images

Early studies emphasize that the H_I properties of NGC 4203 are similar to those of spiral galaxies. For example, Burstein & Krumm (1981) and van Driel et al. (1988) find a distorted rotation curve from the H_I velocity field, possibly suggesting an oval distortion or a kinematical warp (Bosma 1978). They find that the H_I disc is three times larger than the optical body and this is confirmed later by S12 using deeper data.

As described in Section 2, our new H_I image has similar sensitivity to that of S12 but a much better resolution. We detect H_I down to a column density of 2.3×10^{19} cm⁻² at a resolution of 29.3×16.7 arcsec². Using the new H_I observations, one can easily see a dense gas ring whose size is similar to that of the stellar body (Fig. 1). This ring is embedded in a much larger H_I disc characterized by long spiral-like arms. The H_I disc extends out to ~40 kpc from the galaxy centre and contains ~9 times more H_I gas than the inner ring. Table 2 lists some basic galaxy properties derived from the H_I image.

Fig. 2 shows that NGC 4203 hosts some diffuse FUV emission at the location of the inner H I ring. Additionally, at the same location, we see spiral-like structure in both the 8 μ m and g' - r' images, which is clumpy and almost identical in these images. Moreover, the extended spiral/ring structures in the model-subtracted deep optical image (bottom panel of Fig. 3) correspond to the dust lanes visible in the MegaCam g' - r' colour image.

In contrast, the bright H₁ arm ~ 2 arcmin west of the galaxy (visible in the same figure and referred to as western H₁ arm in the rest of this paper) does not exhibit any clear FUV emission. A similar picture is evident from the 8 µm non-stellar image and the g' - r' image. While there is clear emission from PAHs and dust absorption in the H₁ ring, these features are absent from the region occupied by the western H₁ arm. The question is why these two regions have such different levels of the FUV and 8 µm emission despite the very similar H₁ column density. We will come back to this point in Section 4.1.

As indicated in Section 2.3, the deep optical images show a dwarf satellite \sim 7.2 arcmin (32 kpc) east of the NGC 4203. This dwarf is red and at a similar redshift with $V_{\text{helio}} = 1067 \pm 24 \,\text{km s}^{-1}$ (Adelman-McCarthy et al. 2007). The stellar mass of the dwarf is roughly \sim 10⁸ M_☉ from the SDSS *i*-band photometry.

Table 2.	Observational	results	for	NGC	4203
I able 2.	Observational	results	101	1100	7205.

Quantity	Value	Unit	$\pm^{b, c, d}$
og <i>M</i> (H I)	9.12	M _O	2 per cent
ogM(H I) _{Ring}	8.12	M _O	5 per cent
ogM(H I) _{Disc}	9.07	Mo	2 per cent
$og[M(H_I)/M_{stellar}]$	-1.442	_	_
FUV ^a	16.10	AB mag	0.05
FUV(R < 100 arcsec)	16.92	AB mag	0.06
SFR _{Ring,8µm}	8.4×10^{-3}	$M_{\odot} yr^{-1}$	3.0×10^{-3}
SFR _{Ring,FUV+24} µm	4.3×10^{-3}	$M_{\odot} yr^{-1}$	2.0×10^{-3}
SFR _{Ring, FUV}	3.8×10^{-3}	M_{\odot} yr ⁻¹	1.0×10^{-3}
SFR _{Disc}	9.0×10^{-4}	$M_{\odot} yr^{-1}$	4.0×10^{-4}
SFR _{Ring + Disc, FUV}	4.7×10^{-3}	M_{\odot} yr ⁻¹	1.0×10^{-3}
$\overline{\Sigma}$ SFR _{Ring+Disc,FUV}	$5.9 imes 10^{-5}$	$M_{\odot} yr^{-1} kpc^{-2}$	1.3×10^{-5}

Note. ^{*a*}We take into account the Galactic extinction correction for all the values below. Look at the text for the details.

^bThe errors are in the same unit as the values.

^{*c*}The errors in H_I mass values are derived from the noise level (1σ) of the mask file that is used to create the H_I image. We use an annulus that has the same size of the H_I beam.

 d The errors in the SFR and FUV magnitude values are calculated from the uncertainties in the photometric measurements of the FUV, 8 µm and 24 µm images (Calzetti et al. 2007; Morrissey et al. 2007; Leroy et al. 2008).

The bottom panel of Fig. 3 reveals prominent tidal tails associated with the dwarf and a faint stream, which has the same direction as these tidal tails. From this point, it is clear that the satellite dwarf galaxy is interacting with NGC 4203, for example, these features may be associated with a single on-going minor merger event. We will discuss this point in Section 4.2.

3.2 HI, FUV and infrared radial profiles

We derive the radial profile of H₁, UV and infrared emission along circular annuli whose width increases logarithmically from inside out. The inner region of the radial profiles (1 $R_{\rm eff}$) is contaminated by old stars in the bulge and, therefore, we do not include it in our SFR calculation.

Fig. 4 shows the resulting H_I and UV radial profiles together. The first peak of the H_I profile indicates the location of the inner H_I ring \sim 30–80 arcsec from the centre. We find that the slope of the declining FUV profile becomes flatter in this region consistent with the presence of diffuse FUV emission associated with the SF (see Fig. 2). Additionally, the FUV–NUV colour becomes significantly bluer at the same location. This might be because the young stellar populations are dominant in this region. The second peak in the H_I profile of Fig. 4 (\sim 120–150 arcsec) is caused by the high-columndensity gas in the western H_I arm (see Fig. 2). Unlike for the first peak of the H_I profile, here the FUV profile does not seem to depart from the general declining trend. However, the FUV–NUV colour seems to become bluer at this location.

The most significant diffuse FUV emission in NGC 4203 comes from a region close to the UV-bright galaxy bulge. The emission from the bulge may contaminate a large region around it because the *GALEX* PSF is characterized by broad wings out to at least 60 arcsec. To quantify this effect, we create a bulge model with a Sérsic surface brightness profile and convolve it with the FUV PSF. We find that the average contamination from the bulge is ~29.2 AB mag arcsec⁻² between 50 and 60 arcsec. This contribution to the total emission is very low and we ignore this effect in our SFR calculation.

Figure 4. Top panel: radial profiles of the H_I (blue) FUV (red) and NUV (yellow) emission. The left *y*-axis shows the UV magnitude and the right *y*-axis shows the H_I surface density. The green vertical line shows the effective radius of the galaxy. The grey shaded region indicates the inner H_I ring. Bottom: radial FUV–NUV colour profile. The error bars represent the uncertainty in determining the background for both images.

In order to take into account the significant 8 µm non-stellar emission in the inner HI ring (see Fig. 2), we derive the SFR radial profiles based on this 8 µm emission as well as the 24 µm emission (see Section 2.3 and equations 4 and 5). We show the resulting SFR radial profiles together with the FUV one in Fig. 5. All profiles show a similar trend as their slope changes at the location of the HI ring (in fact, at 8 µm the profile increases significantly between 30 and 60 arcsec). The infrared-based SFR profiles are up to a factor of \sim 3 higher than the profile based on the FUV emission alone. This difference is larger than the errors in the photometric measurements and is likely due to infrared emission from obscured SF. The photometric uncertainty for the Spitzer 8 and 24 µm images is \sim a factor of 2, and this uncertainty is also similar for the FUV measurements (Calzetti et al. 2007; Salim et al. 2007). However, at present, the depth of the Spitzer images and the small FOV imply that the UV images are the only reliable tracer of SF at larger radius. Table 2 shows the SFR efficiencies at the location of the inner H I ring for all three tracers as well as their approximate uncertainties.

3.3 SFR in terms of HI column density

In Fig. 6, we compare H I and SFR surface density by plotting the H I and FUV radial profiles of Fig. 5 against one another. In this figure, points are colour-coded according to the distance from the galaxy

centre. The reason that the track described by our measurements is not monotonic is the result of the two HI peaks (one corresponding to the inner HI ring and the other to the outer HI arm; see Figs 2 and 4) while the UV radial profile decreases monotonically with radius. For comparison, the yellow area shows the spatially resolved relation between H I and SFR surface density inside r_{25} for a sample of nearby (D < 12 Mpc) spirals covering a range of morphologies (Bigiel et al. 2008). In this inner region, interstellar medium of these galaxies is dominated by molecular hydrogen. The green vertically and horizontally shaded areas show the same relation for the outer parts (i.e. $1-2 \times r_{25}$) of dwarf and spiral galaxies, respectively (Bigiel et al. 2010). Details about their sample, e.g. the THINGS sample, are given in Walter et al. (2008). This figure shows that the SF efficiency in the bright H1 ring around NGC 4203 is consistent with that of late-type galaxies. Moreover, the SF efficiency in the outer regions of NGC 4203 is also consistent with that of outer regions of spiral and dwarf galaxies.

So far we have compared H_I and SFR surface density at fixed radius. Below, we perform a pixel-by-pixel comparison as implemented by previous authors (e.g. Bigiel et al. 2008, 2010; Boissier et al. 2012). To do so, we smooth and re-grid the FUV and infrared images to the resolution and coordinate grid of the H_I image. We find this comparison to be characterized by a very large scatter. For this reason, we bin the points along the H_I column density axis, and

Figure 5. Comparison of different SFR radial profiles together with the H I profile. We calculate the SFR with equations (3)–(5). Although we show the central emission from all the bands, we do not include these emissions in our SFR calculation. The dashed part of the profiles traces emission from the bulge and this is not included in the calculation of SFR in the rest of paper. The grey shaded area and H I profile are the same as in Fig. 4.

Figure 6. Comparison between H I and FUV surface brightness obtained by plotting the radial profiles of Fig. 4 against one another. The triangles indicate H I detected regions with no detectable FUV emission; we estimate the upper limit as equal to the rms of the FUV background image. The colour bar shows the distance from the centre of the galaxy. The yellow area shows the relation between H I and SFR in the sample of late-type galaxies of Bigiel et al. (2008). The vertical and horizontal green shaded areas show the location of dwarf galaxies and outer regions of spirals, respectively (Bigiel et al. 2010). In this figure, the left axis shows the FUV surface brightness and the right axis shows the SFR calculated with equation (3).

Figure 7. SFR surface density as a function of H 1 column density based on the pixel-by-pixel method. Here, the FUV pixels are binned as a function of H 1 column density. The light blue shaded area shows the scatter of FUV pixels in the bins. The error bars are similar to those in Fig. 4 and represent the uncertainty in determining the background. The grey error bar on the bottom-right shows the result of the same analysis for all the pixels without H 1 emission. The green asterisk and red pentagon show the inner H 1 ring and the western H 1 arm, respectively. The two green lines above the green asterisk represent the SFRs in the inner H 1 ring that are calculated from equation (5) (for FUV+24 μ m) and 4 (for 8 μ m), respectively. The blue points represent the upper limits, calculated by using the rms of the FUV background image. The yellow area and green lines are the same as in Fig. 6.

show the average SFR surface density of each bin as well as the scatter in Fig. 7.

 $\log_{10} \overline{\Sigma} (SFR) [M_{\odot} yr^{-1} kpc^{-2}]$

-5

-6

This plot shows a weak trend between H1 and SFR surface density as well as a large scatter, represented by the blue bars. This large variation of SF efficiency at fixed H1 column density is also clear when one compares the HI ring and the western HI arm mentioned above. The green asterisks and red pentagon indicate the average values of SFR surface density and H1 column density in these two regions, respectively, and are calculated on the basis of the FUV image. Combining the 24 µm image with the FUV image, the SFR estimate increases further, as illustrated by the first green line. If we calculate the SFR with the 8 µm non-stellar emission, we obtain the value that is indicated by the second green bar. As argued above, this strengthens our claim that SF efficiency in the inner HI ring is as high as that of spiral galaxies, while it is much lower in the western HI arm despite the similar gas density. However, we note that even if the SFR is low in the western H1 arm, this SFR is still consistent with the outer regions of spirals and dwarf galaxies.

4 DISCUSSION

4.1 SF efficiency

The previous section demonstrates that at any given H₁ column density there is a large range of SFR surface densities and, therefore, SF efficiencies in the H₁ disc of NGC 4203. This is highlighted by the comparison between the H₁ ring and the western H₁ arm presented in Fig. 7. There are a number of possible explanations for the difference between these two regions: (i) there may be no direct relation between H₁ and SF; (ii) the H₁ column density of the arm might appear to be higher because of projection effects (e.g. the H₁ arm stretches away from the plane of the disc), and the actual density of the gas could be too low to form stars; (iii) the temperature of the H_I in the arm may be higher than that in the ring, making it more difficult to form stars; (iv) the metallicity and dust content of the inner and outer regions might be different due to a difference in formation processes; (v) the western H_I arm may be stable against gravitational instabilities, preventing the formation of stars.

Both the second and third points would result in a broadening of the H_I spectral profile at the location of the H_I arm. This is, however, not observed, leaving us with no clear indication in favour of these options. On the other hand, the fourth point seems an attractive explanation. Miller, Knezek & Bregman (1999) observe a quasar ~ 2 arcmin south-east of the galaxy. Interestingly, they detect one of the closest known damped Ly α (DLA) system. This DLA system belongs to NGC 4203, and the metallicity of the gas is found to be $[Fe/H] = -2.29 \pm 0.1$. They also show a number of heavyelement abundances such as [Mn/H] < -0.68 and [Mg/H] > -2.4, respectively. The velocity of the lines is identical to that of the HI emission from the disc of NGC 4203 at the position of the quasar within $(\pm 10 \text{ km s}^{-1})$. Although the metal abundance values are very low, the depletion of the heavy elements on to small dust grains can affect the observed measurements. Vladilo (2002), for instance, shows that the corrected [Fe/H] abundance of a DLA system can be on average 0.5 dex (in some extreme cases 1 dex) higher than that of observed. Even if we consider 1 dex depletion effect, the metallicity of the outer H I disc remains very low (i.e. [Fe/H] ~ -1.29).

This is likely quite different from the inner H_I ring, which is known to contain a considerable amount of dust (Fig. 2). In particular, we know that the PAH molecules traced by the 8 μ m emission are vulnerable to highly energetic photons in dust-poor environments due to low dust shielding (e.g. Boselli, Lequeux & Gavazzi 2004; Calzetti et al. 2005), and the PAH emission generally vanishes in regions with metallicity below solar (Engelbracht et al. 2005). Thus, the metallicity in the inner region should be significantly higher than that in the outer regions. Consequently, different formation processes for the inner and outer H₁ regions can explain these different metallicities.

The last possibility is the differences of the gravitational instabilities of the H_I gas in the inner and outer regions. For example, SF can occur in regions where the gas disc is unstable against large-scale collapse. If the gas is stable against it, Coriolis forces counteract the self-gravity of the gas and suppress SF (Leroy et al. 2008, and references therein). This might explain the different SF efficiencies. However, in order to understand this effect, we need further investigation.

4.2 Origin of the HI gas

There are different scenarios that could explain the origin of the H_I disc of NGC 4203. One possibility is a high-angular-momentum, gas-rich major merger. Such a merger could result in the formation of a large gas disc (e.g. Barnes 2002). At the same time, it would not necessarily trigger the central burst of SF usually thought to occur during mergers (e.g. Di Matteo et al. 2007; Serra et al. 2008), as evidenced by the old stellar populations in the centre of NGC 4203. If this scenario is correct, we should be able to see a residual from the major merger close to the centre, or a distortion in the outer stellar halo of NGC 4203 in the deep optical image. However, we do not detect any of these features in Fig. 3. Another problem with the major merger scenario is the extremely low metallicity of the outer disc which may rule out that the H_I once belonged to a large galaxy.

The other possibility is that the gas has been accreted from the intergalactic medium (IGM). The disturbed appearance of the H I disc may be explained by some minor interaction with a satellite. Indeed, as explained in Section 3.1, there is a nearby dwarf satellite, which is most likely interacting with our galaxy. Although this dwarf is positioned at the end of one of the H I arms of NGC 4203 (Fig. 1), we do not detect any H I emission at that position and, therefore, rule out that it is the donor of the H I gas.

However, there is another possibility that this gas can be captured from another gas-rich dwarf galaxy in a single encounter. Since the H₁ disc is almost regular, this event would happen at least a couple Gyr ago, making it difficult to detect. This is consistent with the conclusion of van Driel et al. (1988) based on a study of the environment around the galaxy.

In conclusion, accretion from the IGM would explain the low metallicity of the outer disc, while the inner H I ring may have been enriched of metals by stellar mass loss in the galaxy. Moreover, the disturbed appearance of the outer disc can be explained by an interaction with the satellite visible in Figs 1 and 3.

4.3 Impact of the HI on the host galaxy

We have long known that a large fraction of all ETGs host low levels of on-going or recent SF. This SF activity is detected in their centre using optical spectroscopy (e.g. Trager et al. 2000) or UV photometry (e.g. Yi et al. 2005). Similar to extended-UV discs in late-type galaxies, SF can also be detected at large radius in ETGs using UV and deep optical imaging (e.g. Salim & Rich 2010; Fang et al. 2012; Salim et al. 2012; Duc et al. 2015). In a few cases, it has been possible to establish a link between the presence of H_I and SF at the outskirts of ETGs (e.g. Cortese & Hughes 2009; Donovan et al. 2009; Thilker et al. 2010; Duc et al. 2014). Here we explore

such link for NGC 4203, and in this section we quantify the impact that the SF hosted within the H₁ disc has on the morphology of the host.

Table 2 shows the result of our SFR calculation based on the various tracers discussed above, and performed separately for the inner H_I ring and the remaining, outer H_I disc. In particular, the inner H_I ring hosts a significant amount of dust-obscured SF, and the SFR is best estimated using a combination of 24 μ m and FUV emission.

The inner H₁ ring and outer H₁ disc are very different from each other. The former contains only ~10 per cent of the total H₁ mass but hosts ~97 per cent of the total SF. The latter hosts the remaining ~90 per cent of the H₁ mass but only ~3 per cent of the total SF. As a result, the H₁ depletion time is much longer in the outer disc (~10³ Gyr) than in the inner H₁ ring (~5 Gyr). If nothing happens to the galaxy, the H₁ disc could survive for hundreds of a Hubble time at the current SFR.

Overall, the detected amount of SF has a negligible effect on the stellar disc growth and, therefore, the morphology of NGC 4203. Based on the total stellar mass reported by Cappellari et al. (2013) and on the B/D ratio given in Krajnović et al. (2013), the mass of the stellar disc is $\sim 2.7 \times 10^{10} M_{\odot}$. At the current SFR, this mass would grow by a mere 0.1 per cent in 1 Gyr. The SFR per unit stellar mass in the stellar disc is of just $\sim 10^{-12} \text{ yr}^{-1}$.

If we suddenly convert all the gas in the outer H₁ disc of NGC 4203 to stars, the *R*-band surface brightness of this region would be 26.8 mag arcsec⁻². This is ~2 mag lower than Malin 1, a giant low-surface-brightness galaxy that has a low surface brightness and very large H₁ disc (Lelli, Fraternali & Sancisi 2010). In conclusion, it is unlikely that NGC 4203 will become a galaxy similar to Malin 1.

5 CONCLUSIONS

NGC 4203 has an H I disc that consists of two separate components. The first component is the inner star-forming H I ring. This ring contains metal-rich gas and a large amount of dust. The SF efficiency in the H I ring is comparable to that of inner regions of spiral galaxies and the H I depletion time is of the order of a few Gyr. The second component is the outer H I disc, where the metallicity is likely to be low. The SF efficiency in this disc is much lower than that of centre of spirals even when the H I column density is high (e.g. in the western H I arm) and the H I depletion time is of hundreds of the Hubble time. However, the SF efficiency in the outer H I disc is still consistent with that of the outer regions of spiral galaxies and dwarfs.

These two components of the H I structure may be formed in different ways. On the one hand, the H I ring is most likely enriched by stellar processes within the galaxy. On the other hand, the outer H I disc might be accreted from the low-metallicity IGM because the outer disc is formed with metal-poor, high-column-density gas and we cannot find a suitable donor. Subsequent interaction with satellites, which is possibly on-going, may explain the current distorted morphology of the H I disc.

We also detect spiral-like structure in the 8 μ m and deep g' - r' images. This spiral-like structure is extending to 3 effective radii (~6 kpc), and, because of its clumpiness, is likely organized by the SF process.

Finally, despite the large H₁ reservoir and the detection of some SF, the stellar disc of NGC 4203 is growing at a very low rate and the morphology of the galaxy is unlikely to change significantly in the foreseeable future.

ACKNOWLEDGEMENTS

MKY acknowledges a PhD scholarship from the Council of Higher Education of Turkey and is supported by the Erciyes University.

We would like to thank the GALEX team for their great work to make the data public and available. This work uses observations made with the NASA Galaxy Evolution Explorer. *GALEX* is operated for NASA by Caltech under NASA contract NAS5-98034. Part of this work is based on observations made with the *Spitzer Space Telescope*, which is operated by the Jet Propulsion Laboratory(JPL), California Institute of Technology under a contract with NASA. This research has made use of the NASA/IPAC Extragalactic Database (NED) which is operated by JPL, Caltech, under contract with NASA.

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