



## Numerical and experimental analysis on green laser crystallization of amorphous silicon thin films

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### ABSTRACT

The effect of laser fluence on the crystallization of amorphous silicon irradiated by a frequency-doubled Nd:YAG laser is studied both theoretically and experimentally. An effective numerical model is set up to predict the melting threshold and the optimized laser fluence for the crystallization of 200-nm-thick amorphous silicon. The variation of the temperature distribution with time and the melt depth is analyzed. Besides the model, the Raman spectra of thin films treated with different fluences are measured to confirm the phase transition and to determine the optimized fluence. The calculating results accord well with those obtained from the experimental data in this research.

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## 1. Introduction

Polycrystalline silicon (poly-Si) holds great potential for a number of applications, including thin-film transistors for active-matrix liquid-crystal displays (AMLCD) [1], sensors [2] and solar cells [3]. Laser crystallization of amorphous silicon (a-Si) is regarded as the most efficient method for fabricating polycrystalline silicon on glass [4]. However, the traditionally used excimer laser crystallization (ELC) has many disadvantages, such as the narrow process window of laser energy density, very high absorption of ultraviolet wavelength in a-Si and the costly maintenance. So in many recent works, green lasers, especially frequency-doubled Nd:YAG ( $\lambda = 532$  nm), have been used instead of excimer lasers as the crystallization laser source [5–8].

Early experimental results revealed that the phase-change from a-Si to poly-Si depends strongly on the absorbed energy from the incident laser pulse [4]. However, the energy absorption

and the temperature distribution in Si thin films are difficult to measure directly. So it is necessary and significant to make a theoretical analysis of the temperature field of the Si film when irradiated by pulsed laser. The heat transfer models for laser material processing have previously been reported and are used by some researchers for the simulation of ELC [9,10] or the laser interference crystallization [11]. But no systematic numerical study including the laser fluence effect on green laser crystallization has been reported.

According to Ref. [4], three main crystal-growth regimes (partial, near-complete and complete melting regimes) exist, depending on the laser fluence. Large grains can only be produced at the near-complete regime, which indicates a very narrow energy density window. The so-called “critical fluence” ( $F_c$ ) [9] distinguishes between the partial melting regime and near-complete melting regime, indicating the optimized fluence for the crystallization of a-Si films. Hence, it is of great importance to determine the value of  $F_c$ .

In this work, the effect of laser fluence on green laser-induced crystallization of a-Si is investigated numerically and experimentally. The finite difference method is used to solve the heat conduction equation and to predict the value of the melting

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threshold and of  $F_c$ . A laser crystallization experiment was also developed for fabricating poly-Si films. Results of the calculation are compared with experimental data of other researchers and our own in this work.

**2. Mathematical model**

When the pulsed laser irradiates the surface of silicon, only a fraction of the incident energy can be absorbed and taken as the heat source because of the effects of surface reflection. The temperature of the thin film begins to increase with the absorbed energy. The film melts to a liquid when the melt point ( $T_m$ ) of a-Si is reached. Then it cools down and solidifies to polycrystalline phase after the irradiation.

The crystallization process of a-Si can be simulated with a one-dimensional non-linear heat conduction equation:

$$c(T)\rho \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left[ K(T) \frac{\partial T}{\partial x} \right] + S(x, t) \tag{1}$$

where  $c$  is the specific heat,  $\rho$  is the density, the temperature,  $T(x, t)$ , varies with depth  $x$  and time  $t$ ,  $K$  is the thermal conductivity and  $S(x, t) = S(x)S(t)$  is the laser source term. The laser beam is assumed to be Gaussian shape, so the time-dependent part of the source term is given by

$$S(t) = \frac{1}{\Gamma\sqrt{\pi}} \exp\left(-\frac{(t-t')^2}{\Gamma^2}\right) \tag{2}$$

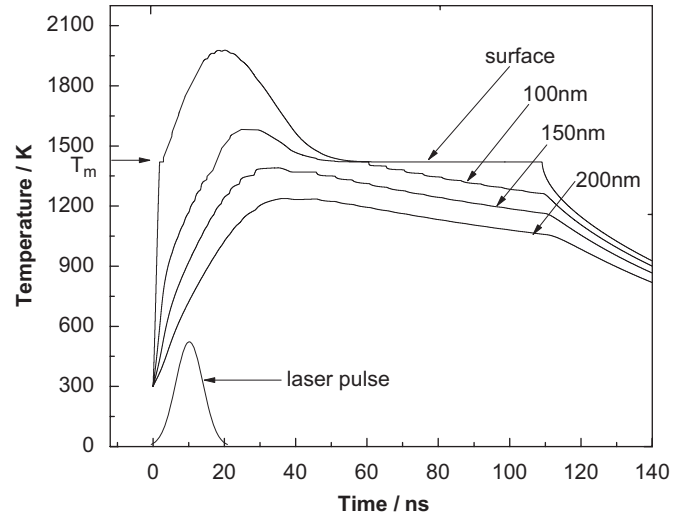
where  $\Gamma$  denotes the full-width duration of the pulse at 1/e intensity and  $t'$  is the time when the pulse reaches its maximum value.

$$S(x) = E_0(1 - R_i)\alpha_i \exp(-\alpha_i x) \quad i = s \text{ or } l \tag{3}$$

where  $E_0$  is the power density of the incident laser;  $R_i$  and  $\alpha_i$  represent the reflectivity at the interface and the optical absorption coefficient of a-Si, respectively. Their value can be very different when the interface is solid phase ( $R_s, \alpha_s$ ) or liquid phase ( $R_l, \alpha_l$ ).

The specific heat  $c$  and the thermal conductivity  $K$  in Eq. (1) are generally temperature dependent, so it is difficult to solve the equation by the analytical approach. The finite difference method is used in this work to solve for  $T(x, t)$ , by dividing the irradiated region into equal sections of length  $\Delta x$  and time into equal sections of  $\Delta t$ . The expressions and values of the physical properties used in this simulation are shown in Table 1.

Initially, the temperature of a-Si is equal to the ambient temperature ( $T_0 = 300$  K). So the boundary condition and initial



**Fig. 1.** Results of the simulation showing the temperature versus time profiles at the surface, 100, 150 and 200 nm layers of the 200 nm a-Si, which is irradiated by laser fluence of 600 mJ/cm<sup>2</sup>.

condition are listed as follows:

$$T(x, t)|_{t=0} = T_0 \tag{4}$$

$$T(x, t)|_{x \rightarrow \infty} = T_0 \tag{5}$$

$$\frac{\partial T(x, t)}{\partial x} \Big|_{x=0} = 0 \tag{6}$$

The results of the modeling of the time-dependent temperature profile for a-Si:H thin film irradiated with a pulsed 2ωNd:YAG laser are shown in Fig. 1. The laser fluence is 600 mJ/cm<sup>2</sup>. The shape of the laser pulse is also shown in the bottom left corner. The temperature at the surface rises rapidly until the melting point is reached. The surface begins to melt shortly after the pulse reaches its peak. When the surface layer is heated to the melt point  $T_m$ , the latent heat of melting must be considered. The temperature of the melting layer will remain at  $T_m$  until the energy of  $\Delta Q = \rho L \Delta x$  is absorbed subsequently. Also, the latent heat should be taken into account when the liquid silicon release heat at the crystallization process.

The temperature profiles at depths 100, 150 and 200 nm are also plotted. It is noted that the temperature at the layer 200 nm below the surface is still lower than  $T_m$ . Therefore, a laser fluence of 600 mJ/cm<sup>2</sup> is not sufficient to melt the entire a-Si thin film of 200 nm. This is confirmed by our experiment, which will be discussed later.

There is an interest on how much laser fluence is required for a-Si crystallization. So the temperatures of surface and bottom layers irradiated by different laser fluences are calculated in Fig. 2. It can be seen that the maximum temperatures for both layers increase approximately linearly with laser fluence. As for an a-Si:H thin film of thickness 200 nm, the temperature of the surface layer reaches  $T_m$  when the laser fluence is about 210 mJ/cm<sup>2</sup>, and the bottom layer begins melting only when laser fluence is over 800 mJ/cm<sup>2</sup>. Then near-complete melting can be achieved.

The calculated threshold energy density for green laser crystallization, 210 mJ/cm<sup>2</sup>, is about two times as much as that of ELC. This is mainly due to the reduced absorption of green light ( $1.25 \times 10^5 \text{ cm}^{-1}$ ) in a-Si compared with that of ultraviolet light ( $1.43 \times 10^6 \text{ cm}^{-1}$ ), so more energy is needed to reach the melting point.

**Table 1**  
Physical properties used in this simulation

Physical properties	Value or expression	Reference
Absorption coefficient (cm <sup>-1</sup> )	$\alpha_s = 1.25 \times 10^5$ ; $\alpha_l = 10^6$	[11]
Density (g/cm <sup>3</sup> )	$\rho_s = 2.20$ ; $\rho_l = 2.53$	
Specific heat (J/g K)	$C(T) = 0.952 + 0.171T/1685$	[10]
Reflectivity	$R_s = 0.4$ ; $R_l = 0.73$	[10,11]
Thermal conductivity (W/cm K)	$K(T) = 4.828 \times 10^{-11} (T-900)^3 + 4.828 \times 10^{-9} (T-900)^2 + 3.714 \times 10^{-6} (T-900) + 3.714 \times 10^{-2}$	[10]
Latent heat (J/g)	$L = 1320$	[11]
Melting point of a-Si (K)	$T_m = 1420$	[12,13]

Subscripts s and l denote the solid and liquid phase, respectively.

### 3. Experiment procedure

A 500-nm-thick SiO<sub>2</sub> layer was deposited by Plasma-Enhanced Chemical Vapor Deposition (PECVD) on corning glass substrate, which acted as a thermal barrier. Then, a 200-nm-thick hydrogenated amorphous silicon (a-Si:H) film was deposited on the SiO<sub>2</sub> layer by PECVD of SiH<sub>4</sub>/SiCl<sub>2</sub>H<sub>2</sub>. The gas pressure of the deposition was 67–133 Pa. Dehydrogenation of a-Si was done by a furnace annealing at 500 °C for 2 h.

A laser crystallization apparatus used in this work consisted of a 2ωYAG laser (532 nm) with 10-ns pulse width, a beam shaping system and a translation stage. The schematic diagram of our experiment is given in Fig. 3. The beam shaping system, including a fly-eye lens array, transformed the incoming Gaussian laser beam into a flattop one. A wave-front analyzer was used to examine the intensity uniformity on the image plane. The samples placed on the stage were scanned at a speed of 1.0 cm/s. After the crystallization, Raman spectra of thin films were measured by a Jobin Yvon LabRam-1B micro Raman spectrometer.

### 4. Results and discussion

Raman scattering measurements are performed to investigate the microstructure. As can be seen in Fig. 4, a sharp peak appears

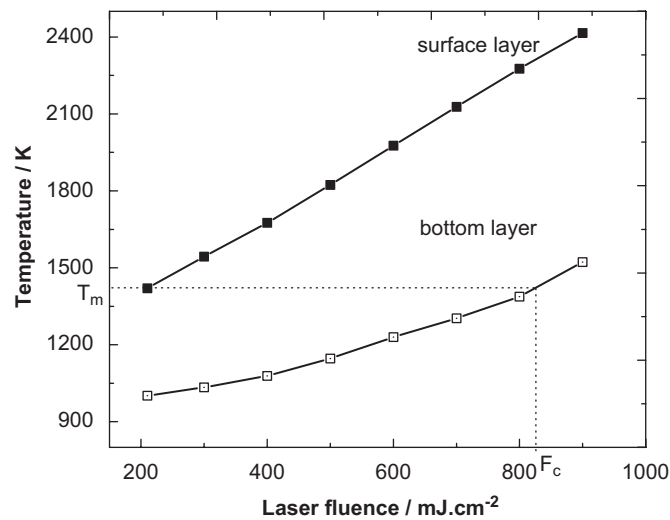


Fig. 2. Calculated maximum temperatures of the surface and bottom layers (200 nm below surface layer) at different laser fluences.

between 510 and 520 cm<sup>-1</sup> after the laser treatment. That peak is assigned to the TO phonon scattering of poly-Si and resulted mainly from the phonon confinement effects [14], which indicates the phase transition from a-Si to poly-Si. The same TO peaks can be obtained by some other different crystallization techniques, such as ELC [15], low-pressure chemical vapor deposition [14] and so on.

Further investigation of the Raman spectra shows that laser fluence used in the irradiation of the thin films may exert a major influence on the shift and full-width at half-maximum (FWHM) of the TO peak. Fig. 5 shows the intensity-normalized Raman spectra of poly-Si crystallized using different energy densities. It can be seen that the asymmetric profile of the Raman spectra tends to be narrower and more symmetric when the laser fluence increased from 400 to 850 mJ/cm<sup>2</sup>. Also, the TO peaks shift towards the frequency of 520 cm<sup>-1</sup>, which is characteristic of crystalline silicon. This phenomenon indicates the enlargement of the grain size and crystalline percentage in the poly-Si thin films [15]. However, when the laser energy density is over 900 mJ/cm<sup>2</sup>, the position of the TO peak begins to shift back.

In detail, the shift of the TO peaks at different laser fluences is shown in Fig. 6. The calculated melt depths of 200 nm a-Si from the heat conduction equations are also plotted.

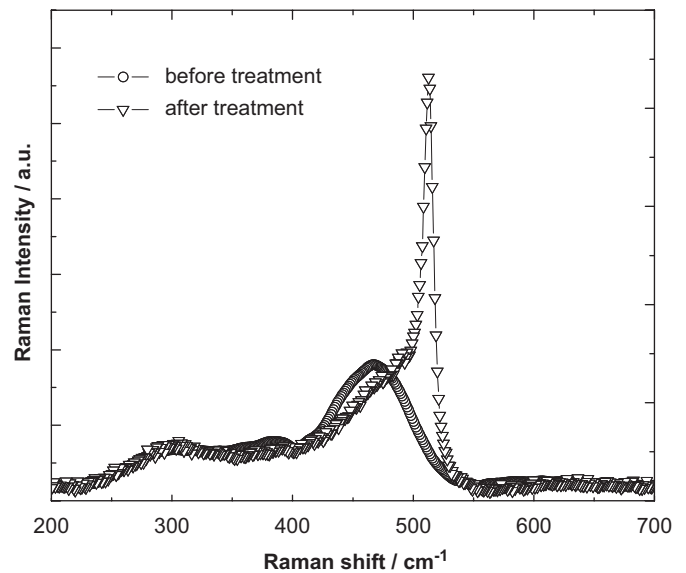


Fig. 4. Raman scattering spectra of a-Si before and after laser treatment.

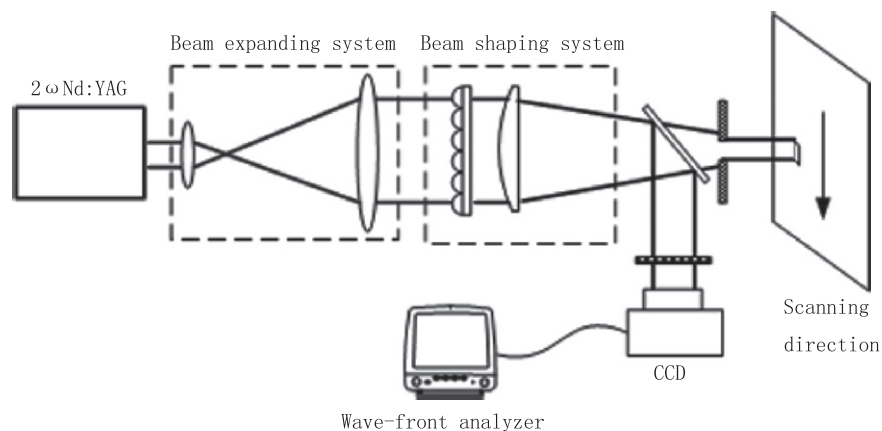


Fig. 3. Schematic diagram of laser crystallization.

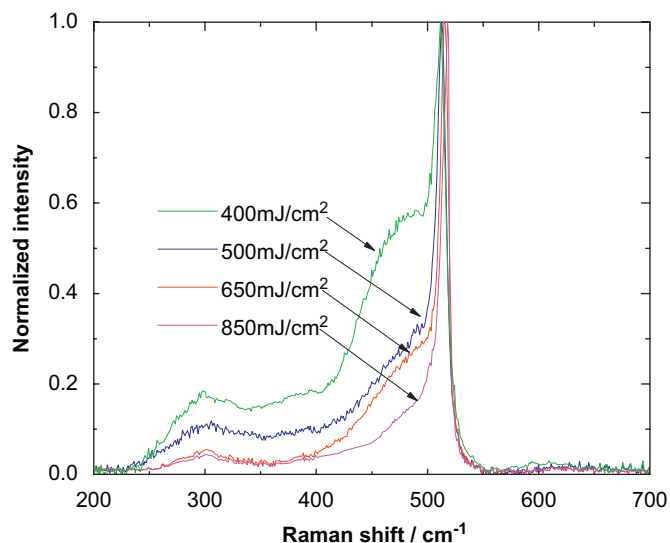


Fig. 5. Normalized Raman spectra of polycrystalline silicon crystallized by different laser fluences.

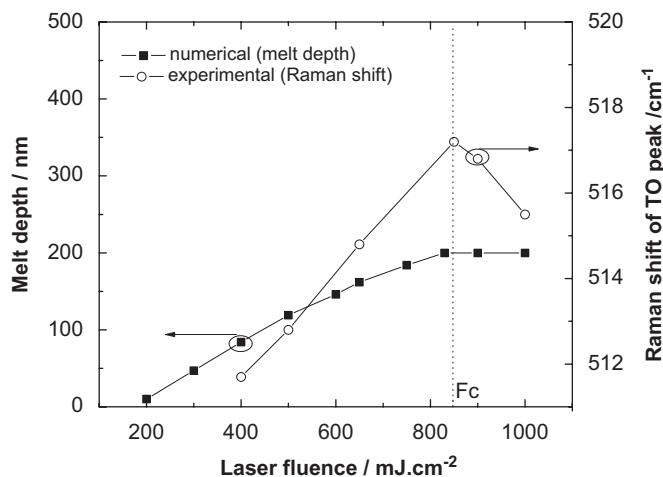


Fig. 6. Calculated melt depths and Raman shift of TO peaks versus laser fluence.

The critical  $F_c$  can be used to identify the partial and near-complete melting regimes. In Fig. 6, the fluence lower than  $F_c$  is in the partial melting regime. The calculated maximum melt depth is lower than the thickness of the Si film (200 nm) and it increases almost linearly with laser fluence. If the fluence is larger than  $F_c$ , the regime of complete melting is obtained and the melt depth remains constant with the thickness of the Si film. The near-complete melting occurs at the fluence lower than and near  $F_c$ ; in this work, it is about 830 mJ/cm<sup>2</sup>. On the other hand, in our experiment, the TO peaks in Raman spectra have the maximum shift when the fluence is about 850 mJ/cm<sup>2</sup>. So from Fig. 6, the

value of  $F_c$  obtained from the simulation results agrees fairly well with the experimental one.

## 5. Conclusions

In this paper, the green laser crystallization of a-Si is investigated numerically and experimentally. The model is used to predict the optimum laser parameters, including the threshold and critical fluence for crystallization. It also provides better understanding on the dynamics during the laser-material interaction process. The optimized critical fluence calculated by this model agrees well with that determined from Raman spectra. Near-complete melting is achieved by laser fluence of about 850 mJ/cm<sup>2</sup> in this work. In this region, the melt depth just approaches the thickness of the Si film and the largest grain size and crystalline fraction can be found.

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