# Model of semidiurnal pseudo tide in the high-latitude upper mesosphere

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Journal of Atmospheric and Solar Terrestrial Physics August 24, 2011 Abstract: We present numerical results for the m = 1 meridional winds of semidiurnal oscillations in the high-latitude upper mesosphere, which are generated in the Numerical Spectral Model (NSM) without solar excitations of the tides. Identified with heuristic computer runs, the pseudo tides attain amplitudes that are, at times, as large as the non-migrating tides produced with standard solar forcing. Under the influence of parameterized gravity waves, the nonlinear NSM generates internal oscillations like the quasi-biennial oscillation, that are produced with periods favored by the dynamical properties of the system. The Coriolis force would favor at polar latitudes the excitation of the 12-hour periodicity. This oscillation may help explain the large non-migrating semidiurnal tides that are observed in the region with ground-based and satellite measurements.

## 1. Introduction

The semidiurnal tide is one of the largest dynamical features of the mesosphere and lower thermosphere. At mid- and high-latitudes, the semidiurnal amplitudes in winds and temperatures can exceed those of the diurnal tide (e.g. Burrage et al., 1995; Avery et al., 1989; Manson et al., 1989; Vincent et al., 1989; Jacobi et al., 1999). The migrating tides (i.e., sun-synchronous diurnal and semidiurnal oscillations) dominate. But the non-migrating tides are also found to be important in the region. Based on observations and supported by modeling studies, a large body of evidence has accumulated over the years that shows that the non-migrating tides represent a major part of the phenomenology that characterizes the mesosphere and lower thermosphere (e.g., Manson et al., 1989; Miyahara et al., 1999; Forbes et al., 1995; Talaat and Lieberman, 1999, 2010).

The phase progression of observed non-migrating tides suggests that they propagate up from below [e.g., Talaat and Lieberman, 1999]. However, the sources for these tides in the troposphere and stratosphere are not sufficient to reproduce the observed amplitudes in the mesosphere [e.g., Miyahara et al., 1993; Hagan et al., 1997]. Non-migrating tides are also generated through non-linear interactions between migrating tides and planetary waves (e.g., Teitelbaum and Vial, 1991; Forbes et al., 1995; Miyahara et al., 1997; Hagan and Roble, 2001), and Mayr et al. (2003, 2005) demonstrated that gravity wave interactions contribute significantly to the non-linear coupling involved.

We report here that the Numerical Spectral Model (NSM) produces, without solar excitation of tides, non-migrating semidiurnal tides in the upper mesosphere at highlatitudes – we refer to as pseudo tides. (This term was coined by Walterscheid et al. [1986] to describe a semidiurnal wave at high latitudes that was much larger than the migrating tide.) The present paper is the follow-on of a study of wind oscillations with periods around 10 hours, which have been observed with different measurement techniques in the upper mesosphere (e.g., Manson and Meek, 1990; Hernandez, 1992, 1993; Ruster and Sivjee, 1994; Forbes et al., 1999; Portnyagin, 2000; Wu et al., 2002). Lamb waves and "normal modes" were invoked to explain the measurements, but inertio gravity waves, nonlinear coupling between tides and planetary waves were also proposed as possible excitation mechanisms. Such oscillations are generated in the NSM with zonal wave numbers m = 0 - 2 (Mayr et al., 2004). With periods around 10 hours, close to 12 hours, the waves have much in common with the semi-diurnal tide. Based on classical tidal theory (Longuett-Higgins, 1968), close examination shows that the oscillations are related to Class-I waves, which have also been referred to as Inertio Gravity Waves (Andrews et al., 1987; Volland, 1988). Since the waves occupy low zonal wave numbers, Mayr et al. refer to them as planetary-scale inertio gravity waves (PIGW). The PIGW are not produced in the NSM when the planetary waves (PW) are turned off. When the diurnal tides are turned off, the PIGW are still generated as shown in Fig. 1. In that case, the amplitudes (Fig. 1b) are comparable to those produced with tides (Fig. 1a), but the seasonal variations differ considerably.

We found that the PIGW spectrum without tidal excitation is highly variable and can broaden around 10 hours to produce 12-hour oscillations. As seen in the following model results for m = 1, the generated pseudo tides contribute significantly to the variability of the wind field to produce, at times, relatively large semidiurnal amplitudes in the upper mesosphere at high latitudes.

# 2. Numerical Spectral Model

The Numerical Spectral Model (NSM) extends from the Earth' surface into the thermosphere, and its design and applications have been discussed in the literature (e.g., Chan et al., 1995; Mengel et al., 1995; Mayr et al., 1997, 1998). For the zonal mean, the model is driven by solar heating due to ultraviolet radiation in the mesosphere and stratosphere taken from Strobel (1978). The tides are driven exclusively by the migrating thermal excitation sources in the troposphere and stratosphere (Forbes and Garret, 1978). Extreme ultraviolet heating is applied in the thermosphere based on the efficiency provided by Torr et al. (1980). Nonlinear interactions between migrating tides and planetary waves produce the non-migrating tides. The planetary waves are generated through the instabilities that arise in the zonal-mean temperature and wind fields. The radiative loss is described in terms of Newtonian cooling adopted from Zhu [1989].

An integral part of the NSM is that it incorporates the Doppler Spread Parameterization (DSP) for small-scale gravity waves (GWs) (Hines, 1997a, b). The DSP has also been applied successfully in a variety of other global-scale models (e.g., Akmaev, 2001a, b; Manzini et al., 1997) and has been discussed extensively in the literature.

In the NSM discussed here, tropospheric heating is applied for the zonal mean to reproduce qualitatively the observed zonal jets near the tropopause and the accompanying latitudinal temperature variations. This heat source generates through instabilities tropospheric planetary waves, in particular those for zonal wave numbers m = 1, 2. To capture the GW effects and nonlinear interactions, the time step used in these runs is 5 min. The data are recorded every hour, and a Fourier spectral analysis is performed in time (over 10, 5, or 3 day running windows) and longitude at each latitude and altitude to extract wave amplitude and phase.

## 3. Comparison of Pseudo Tide with Solar Non-migrating Tides

Under the steady source of solar heating, the regular non-migrating tides reveal persistent seasonal variation. The present pseudo tide is solely generated by variable planetary waves, which are excited by instabilities and influenced by background zonal wind and temperature variations. The tide thus manifests itself mainly during relative short periods of time. This variability is apparent in the following model results that are analyzed and presented with variable running windows.

Figure 2a shows the maximum amplitude m = 1 spectra of meridional winds at 100 km for westward and eastward propagating waves, computed with standard excitation of the solar migrating tides. Periods between 8 and 14 hours are shown, and a Mercator projection is applied to expand the regions at high latitudes. The amplitudes are computed with running windows (RW) of 10 (a), 5 (b), and 3 (c) days (sliding every day) applied to a time span of 6 months after December solstice. In the Polar Regions, the 12-hour westward propagating non-migrating tide dominates, and the maximum amplitudes increase from 50 to 62 m/s as the RW is lowered from 10 to 3 days. Also shown in the spectra are oscillations with periods around 10 hours, which have been identified as planetary-scale inertio gravity waves (PIGW) (Mayr et al., 2004) earlier discussed. As the RW is shortened, the amplitudes of the PIGW broaden and merge into the 12-hour semidiurnal tide.

The meridional wind spectra generated without excitation of the migrating tides are shown in Fig. 3, analogous to Fig. 2. As expected, and in contrast to Fig 2, the sharp spectral peaks of the non-migrating 12-hour tide are not generated. But there is considerable spectral power around 11 hours, which leaks into the 12-hour oscillation with wind amplitudes that increase from about 12 to 50 m/s near the South Pole as the RW is shortened from 10 to 3 days. Figure 3 clearly shows that the semidiurnal oscillations generated without the solar tidal excitation are an important part of the non-migrating tides in Fig. 2. Generated solely with planetary waves (PW) excited by instabilities, however, the spectrum of the semidiurnal pseudo tide is much more variable than that produced with the solar tide. That variability can produce, at times, relatively large semidiurnal amplitudes in the upper mesosphere.

In Figure 4 we examine the seasonal variations of the 12-hour westward propagating oscillations for m = 1 meridional winds at 100 km, which are produced with tidal excitation. The model results are presented with 4-m/s contour intervals, and the lowest 2 levels are suppressed to eliminate clutter. As in Fig. 2 for the corresponding spectra, different running windows (RW) are applied, and the amplitudes are presented for the 3rd model year from 24 to 36 months, varying with latitude and Mercator projection. (The spectral analyses for m = 1 in Fig. 2 apply to the time span from 24 to 30 months and covers all the seasonal variations, albeit in both hemispheres.) As seen in Fig. 2 for the westward propagating oscillations, the semidiurnal tides peak in the Polar Regions. With tidal excitation, the peak amplitudes around 50 m/s tend to occur during the fall and spring seasons, but considerable oscillations are also generated throughout the year including the summer and winter months.

Analogous to Fig. 4, we present in Figure 5 the seasonal variations generated without tidal excitation. They occur mainly around equinox during the spring and fall seasons in

both hemispheres. With RW of 10 days (5b), the amplitudes are small, generally less than 15 m/s, compared to values around 50 m/s with tidal excitation in Fig. 4a. For the 5-day window (5b) the amplitudes increase to 30 m/s, versus 60 m/s in Fig. 4b. And with the 3-day RW (5c), the equinoctial maxima of the 12-hour pseudo tide have amplitude maxima around 50 m/s, which are comparable to the values generated with tidal excitation (Fig. 4c). Consistent with the spectral analysis in Fig. 3, the variable 12-hour pseudo tide can produce relatively large oscillations in the upper mesosphere. During short periods of time in the spring and fall seasons, the pseudo tides (Figs. 5b and 5c) thus contribute significantly to the large amplitudes that are generated in the semidiurnal wind field (Figs. 4b and 4c).

In Figure 6, we examine the phase relationship with altitude of the m = 1 meridional winds at 84° South, which are computed with (6a) and without (6b) excitation of the migrating tides. The plot covers a period of 3 days during a time span after March of the 3rd model year, when the 12-month pseudo tide has large amplitudes as seen in Fig. 5c. (The contours appear rugged because the model results were recorded in 5-km intervals.) For the run with solar forcing, with contours of 6 m/s, the oscillation velocities grow with altitude and attain values as large as 71 m/s at 105 km. The periods are around 12 hours, varying from day to day. In addition, the vertical structure also shows some variability in wavelength, which is generally 20-30 km at this high latitude. In comparison, the pseudo tide (Fig. 6b) attains a peak amplitude of 60 m/s at 100 km. The periods are close to 12 hours, and the vertical wavelengths are about 20 km. Although the oscillations are not generated with solar heating, the waves have properties closely resembling semidiurnal tides. During the short time period depicted, the large oscillations in Fig. 6b contribute significantly to the variability of the velocity field (Fig. 6a) produced with standard excitation of the solar tides.

## 4. Discussion and Conclusion

The numerical results discussed are generated in the NSM with small-scale gravity waves (GW), applying Hines' Doppler spread parameterization (DSP). As shown from a numerical experiment presented in Fig. 3 of Mayr et al. (2011), the GWs amplify the diurnal tide and decrease its vertical wavelength. We believe that this wave source is playing a central role in exciting the m=1pseudo tide. Due to the positive nonlinear feedback associated with critical level absorption, oscillations can be generated with periods that are favoured by the internal dynamical properties of the system. An example for this mechanism is the quasi-biennial oscillation (QBO), whose period is controlled by the 2-year time constant for the eddy viscosity (e.g., Mayr et al., 2010). And in the present case of the semidiurnal pseudo tide at polar latitudes, the Coriolis force would favour the excitation of the 12-hour periodicity.

In general, ground-based observations have difficulty resolving the diurnal or semidiurnal field into non-migrating and migrating components due to insufficient geographic coverage. Measurements taken near the poles are the exception, since different sites and/or different viewing angles from a single site allow longitudinal variations to be quantified. As a result, ground-based observations have for years detected the presence of non-migrating tides in Antarctica (e.g., Hernandez, 1993; Forbes, 1995; Portnyagin et al., 1998) and Arctic (e.g., Wu et al. 2003; Murphy et al., 2006).

Recently, Iimura et al. [2009, 2010] examined the non-migrating semidiurnal tide over the Antarctica and Arctic respectively using data from TIMED/TIDI. The observed dominant non-migrating modes were found to be the westward propagating tides with m = 1 (and m = 3). As shown in Figs. 3, 5, the m = 1 semidiurnal pseudo tides generated in the NSM would contribute significantly to the observed oscillations at high latitudes. (The TIDI data also show a sizable stationary m = 0 tide, but there is no such pseudo tide in the model.) Iimura et al. found that, as a function of latitude, the amplitude of the m = 1 tide increases towards the pole – and our simulations produce a similar morphology (Figs. 3, 5). Typically, the ground-based meteor wind observations show that the vertical wavelength of the m = 1 semidiurnal tide is observed to decrease from ~40 km at 65° to 20 km at the pole. In agreement with the observations, Fig. 6 shows that the wavelength of the model pseudo tide is also close to 20 km at high latitudes.

In the real world and in the present nonlinear model, the sources of the semidiurnal tide are co-mingled and do not produce simple additive effects. Notwithstanding this complexity, we have identified with heuristic computer runs an additional source of variability for the semidiurnal tide. The pseudo tide discussed would likely manifest itself at high latitudes, particularly during periods of enhanced planetary wave activity or during seasons with weak solar forcing of the semidiurnal tides.

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# **Figure Captions**

<u>Figure 1.</u> Computed meridional wind amplitudes for m = 1 westward and eastward propagating planetary-scale inertio gravity waves, which are produced in the NSM with periods between 9 and 10.5 hours at 100 km. The oscillations are generated with (a) and without (b) the excitation of the solar migrating diurnal tides. The amplitudes are of comparable magnitude, but the seasonal variations differ considerably. (Figure are taken from Mayr et al., 2004)

Figure 2. Maximum amplitude spectra of meridional winds at 100 km for westward and eastward propagating m = 1 waves, computed with standard excitation of the solar migrating diurnal tides. Periods between 8 and 14 hours are shown, and the Mercator projection is applied to expand the regions at high latitudes. Westward waves are shaded orange, while eastward waves are blue. Amplitudes are computed with running windows of 10 (a), 5 (b), 3(c) days applied to a time span of 6 months after December solstice. As expected, the non-migrating 12-hour tide dominates, and the spectral amplitudes increase with decreasing windows.

Figure 3. Same as Figure 1 but computed without the solar excitation of the migrating diurnal tides. With decreasing windows, the amplitudes of the periods around 10 hours broaden and leak apparently into the 12-hour oscillation, to produce maximum westward wind amplitudes of 12 (a), 30 (b) and 47 (c) m/s. The pseudo tides thus become increasingly more important when compared with the non-migrating semidiurnal amplitudes in Fig. 2, which have amplitudes of 50 (a), 60 (b), and 65 (c) m/s.

Figure 4. Seasonal variations of westward propagating oscillations for m = 1 meridional winds at 100 km, which are produced with solar excitation of diurnal tides, applying running windows of 10 (a), 5 (b) and 3 (c) days as in Figure 2. The results are presented for the 3<sup>rd</sup> model year from 24 to 36 months, and the Mercator projection is applied to expand the regions at high latitudes. With 4-m/s contour intervals, the lowest 2 levels are suppressed to eliminate clutter. Wind oscillations are generated throughout the year, but the largest amplitudes tend to develop during the spring and fall seasons in both hemispheres.

Figure 5. Same as Figure 4 but computed without the solar excitation of the migrating diurnal tides. The oscillations are generated only around the spring and fall seasons, and the maximum amplitudes increase with decreasing widows, i.e. 14 (a), 31 (b), 47 (c) m/s, in agreement with the varying spectra in Fig. 3. The pseudo tides thus contribute increasingly to the variability of the non-migrating tides in Fig. 4.

Figure 6. Meridional winds varying with altitude for m = 1 at  $84^{\circ}$  South, positive/red and negative/blue (6 m/s contours), computed with (a) and without (b) excitation of the migrating diurnal tides. The oscillations cover a short period of 3 days during a time span after March of the  $3^{rd}$  model year when the 12-hour pseudo tide has large amplitudes as seen in Fig. 5c.





Figure 2









