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Comparison of image restoration methods for lunar epithermal neutron emission mapping

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ABSTRACT

Orbital measurements of neutrons by the Lunar Exploring Neutron Detector (LEND) onboard the Lunar Reconnaissance Orbiter are being used to quantify the spatial distribution of near surface hydrogen (H). Inferred H concentration maps have low signal-to-noise (SN) and image restoration (IR) techniques are being studied to enhance results. A single-blind, two-phase study is described in which four teams of researchers independently developed image restoration techniques optimized for LEND data. Synthetic lunar epithermal neutron emission maps were derived from LEND simulations. These data were used as ground truth to determine the relative quantitative performance of the IR methods vs. a default denoising (smoothing) technique. We review and used factors influencing orbital remote sensing of neutrons emitted from the lunar surface to develop a database of synthetic "true" maps for performance evaluation. A prior independent training phase was implemented for each technique to assure methods were optimized before the blind trial. Method performance was determined using several regional root-mean-square error metrics specific to epithermal signals of interest. Results indicate unbiased IR methods realize only small signal gains in most of the tested metrics. This suggests other physically based modeling assumptions are required to produce appreciable signal gains in similar low SN IR applications.

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1. Introduction

NASA's Exploration Systems Mission Directorate developed the now orbiting Lunar Reconnaissance Orbiter (LRO) to quantify resources for future human activities on the lunar surface (Chin et al., 2007). The operating Lunar Exploration Neutron Detector (LEND), onboard LRO, was included to provide direct geochemical evidence for water on the Moon (Mitrofanov et al., 2008). However, the efficiency in which neutrons can be detected in lunar orbit is limited by the coupled effects of instrument design constraints and the low lunar neutron fluences at orbital altitudes (Feldman et al., 1993). This results in higher uncertainties or lower signal-to-noise (SN) at the detector which is offset by accumulating samples over a surface region using multiple orbits.

Low SN in geochemical maps has been addressed in similar analytical efforts on data from the precursor Lunar Prospector (1998) mission's neutron spectrometer (LPNS) and gamma-ray

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spectrometer (LPGRS) experiments that used image restoration (IR) techniques to deblur and denoise maps (Elphic et al., 2005; Lawrence et al., 2006; Eke, 2001). In this paper we evaluate the potential performance of several IR methods for unbiased restoration of expected lunar neutron emission derived from LEND (Mitrofanov et al., 2008). This includes the results of a single-blind study of IR method performance via direct comparison to synthetic neutron maps and against a default denoising technique designed to quantify signal gains specific to IR deblurring.

While LEND is similar to LPNS in its operations and scientific objectives, there are significant differences in instrument design, expected mapping resolutions and SN that warrant independent evaluation of IR methods. For our single-blind study, four out of twelve invited teams of IR researchers accepted collaboration invitations. Teams configured and submit a total of four IR methods plus a default Gaussian denoising (smoothing) method for evaluation. The blind study was performed in a two-phase process similar to Nesbitt (2004). A preliminary training phase was used to configure and optimize IR methods which provided both "true" and degraded maps. In the subsequent blind study

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phase only "true" maps were provided for restoration (Ivatury and McClanahan, 2009). The primary research objectives are to identify optimal IR methods for LEND map restoration as well as to quantify the expected improvement vs. the default map denoising method. The IR methods and teams include:

- Iterative Gaussian Smoothing (GSFC-IGS), T. McClanahan, NASA/Goddard Space Flight Center.
- Regularized Deconvolution (GSFC-UM-RD), V. Ivatury, Univ. of Michigan/GSFC.
- Weiner (UMD-Fourier), D. Usikov, G. Nandikotkur, G. Milikh, R. Sagdeev, Univ. Maryland (MD).
- Conjugate Gradient (UMD-CG), Usikov, G. Nandikotkur, G. Milikh, R. Sagdeev, Univ. MD.
- 5) Pixon (PIXON), Pixon LLC., R. Peutter.

The duration of the paper reviews the necessary background in orbital lunar neutron detection and LEND specific instrument configuration used to formulate a model of the expected epithermal maps. A short description of the IR methods that were selected is provided as well as a description of the protocols that were used for the two-phase blind study. This review also includes the design and formulation of synthetic LEND maps, training, parametric configuration of methods and IR evaluation metrics. The results section quantifies IR method performances and signal gains against the default denoising (Gaussian smoothing) technique.

2. Background

The LRO mission was created by NASA in response to the Presidents 2004 Vision for Space Exploration and the tantalizing evidence of hydrogen deposits found at the lunar poles by LPNS (Arnold, 1979; Feldman et al., 1999b; Feldman et al., 2000). LRO successfully entered lunar orbit in June 2009 with scientific objectives that include locating and quantifying volatile deposits, e.g. water, and other potential resources that would support future human exploration of the surface. LEND is measuring lunar surface neutron emissions to refine the spatial distributions and quantity estimates of the putative polar H deposits identified by LPNS (Mitrofanov et al., 2007).

Hydrogen on the Moon has been postulated for over 30 years (Arnold, 1979). Though originally thought to be dry, recent reanalysis of Apollo-15 lunar basalts collected at mid-latitudes indicates ambient H concentration may approach 745 ppm $\sim\!0.07\%$ (Saal et al., 2008). However, LPNS low resolution evaluations of the polar region indicated H concentrations may approach $\sim\!3\text{--}5\%$ and that enhanced concentrations may exist in smaller spots at the poles (Feldman et al., 1998a). It is inferred from these results that H enhanced regions are accumulated through the combined effects of a regional H budget driven by variably offsetting factors governing depositional and entrainment processes. H depositional processes include periodic cometary and meteoritic bombardment and/or isotropic influx of solar wind protons (Vondrak and Crider, 2003).

H accumulations can be regionally influenced by various geophysical entrainment effects that constrain the rate of H sublimation. At extremes some regional sublimation processes may be completely attenuated in polar topographic depressions, where persistent low polar inclination $\sim 1.5^{\circ}$ to the sun produces low and permanent shadowing conditions (PSR)'s with persistently low temperatures < 100° K (Neubert et al., 2005; Vasavada et al., 1999). These conditions, over time are postulated to minimize H losses from sublimation and assuming isotropic H influx, may lead to regional H accumulations. As a result, PSR's

have been defined as the primary LEND surface targets (Mitrofanov et al., 2010; McClanahan et al., 2009).

Importantly, the interpretation of orbital neutron measurements only indicates the presence of H, not its geochemical form e.g. water (H_20) , hydroxyl $((OH)^{-1})$ and protons (H^+) . Neutrons are produced in the lunar regolith via isotropic influx of high energy galactic and solar cosmic rays. In this process, subsequent nuclear interactions induce a number of reactions including spallation of neutrons and y-ray's in the top few centimeters of regolith. Neutrons propagate through the regolith interacting with atoms, thereby thermalizing (losing energy) until reaching thermal equilibrium and finally being absorbed in the regolith, or importantly, lost into space as a leakage flux. Rates at which neutrons thermalize (lose energy) in materials are dependent on the material density and atomic cross section of atoms encountered during this process. Compared to most elemental constituents of planetary materials, H has a high cross section to neutrons. As a result, the energy distribution of the neutron leakage flux reflects intensity deficits at epithermal energies commensurate with higher H concentrations (Feldman et al., 1993). Lunar H concentrations are derived from detected neutron distributions cross referenced to radiation and neutron transport modeling that factors regolith elemental compositions with variable H, neutron energies and detector/geometric configurations, e.g. Monte Carlo N-Particle Transport (MCNPX) codes (Pelowitz, 2008).

The areal scale of most polar PSR targets is small < few km diameter and the LPNS uncollimated field of view (FOV) at low altitude is inferred to be larger (~44 km FWHM) than the diameter of many PSR's (Feldman et al., 1999a; Bussey et al., 2003). From the areal disparity in the expected small area of PSR's and the larger LPNS orbital FOV, it is inferred some localized hydrogen deposits may exceed LPNS estimates due to the effective averaging (blurring) of variable epithermal rate regions within the FOV. LEND's design enhances its spatial resolution using a passive neutron collimator containing neutron absorbing materials, 10B and polyethylene, to effectively discriminate surface emissions of epithermal neutrons to within $\pm 5.6^{\circ}$ of the instrument boresite. From its nominal 50 km mapping altitude, this configuration yields an instrument FOV that subtends a ~10 km diameter surface region (footprint). LEND will primarily point nadir and integrate signals at 1 Hz rates using its 10 internal and external detectors. These detectors monitor the energy spectrum of low (thermal), medium (epithermal) and high (fast) energy neutrons emitted from the lunar surface and from the spacecraft. To increase SN, four identical co-aligned nadir pointing epithermal neutron detectors are implemented in the collimator, as well as a total of six detectors positioned internal and external to the detector system for tracking background, instrument and spacecraft induced neutron fluxes (Mitrofanov et al., 2008, 2009).

2.1. Image and epithermal neutron map reconstruction

Image formation is a stochastic process that convolves a true signal \mathbf{f} , degraded by instrumental or environmental blurring \mathbf{w} , with additive noise functions \mathbf{n} , to produce an output image \mathbf{g} , where $\mathbf{g} = \mathbf{f} \otimes \mathbf{w} + \mathbf{n}$ (Peutter et al., 2005; Gonzalez and Woods, 1993). Instrumental blurring processes degrade the original image \mathbf{f} as a convolution \otimes with the instrument PSF, \mathbf{w} and statistical counting noise is added, \mathbf{n} (Poisson). Image restoration techniques aim to reverse the blurring process \mathbf{w}^{-1} and attenuate the noise \mathbf{n} , to obtain an estimate of the original image, $\mathbf{g}' \approx \mathbf{f}$. In high SN applications many types of image restoration method have been shown to yield significantly improved results (Peutter et al., 2005). However, image restoration becomes increasingly problematic with lower SN applications due to the combined effects of

statistical noise being variably spatially correlated. In these conditions spatial frequencies may intersect with real image frequencies and increase the difficulty in estimating an underlying image formation model. Additionally, some restoration processes may amplify noise signals, yielding artifacts e.g. fluctuations (mottling), which strongly resemble real signals (McClanahan et al., 2007; Ugbeme et al., 2007).

Accurate map formation models for orbital neutron remote sensing are also difficult to estimate due to the low signal rates and variances in environmental and instrumental sampling conditions inherent in long duration mapping. Efforts to mitigate these issues and to optimize reconstructions include techniques to extend map formation models by use of other physically based constraints. However, the implementation of physical constraints in the IR process is performed through either parametric biases or to establish prior physical limits for some pixel values (Elphic et al., 2007; Eke et al., 2009). Eke et al. claim signal gains in restoring LPNS data only with the implementation of physically based constraints, e.g. shadowing. Concerns with this approach are that the lunar polar hydrogen budget is presently poorly understood. Implementing these assumptions as restoration priors may neglect either the impact or appropriate weight of other geophysical processes. Results of non-biased IR techniques on LPNS data were equivalent to the same data smoothed with a Gaussian operator, underscoring the importance of accurate priors and understanding the ramifications of potential bias errors in low SN IR applications (No Free Lunch) (Wolpert and Macready, 1997).

Fig. 1 illustrates LEND's neutron map formation process including required functions and parameters to configure the described IR methods. This includes surface emission and instrumental processes e.g. PSF w, background estimation and techniques used to map neutron detections from orbit. Poisson counting statistics, n describe the arrival process of neutrons accumulated by the detector over a given region and defines the expected pixel uncertainty as sgrt(n). A consequence of this long term multisample mapping strategy for IR e.g. LPNS, is that as opposed to traditional higher SN IR models, each detector measurement accumulated to the map is uniquely a function of time, instrument operations, position and background. As a result, new variables are introduced to the map model e.g. altitude, solar flux, position, pointing which impact estimates of the PSF and statistical uncertainties needed for image reconstruction. For instance, altitude variance in combined measurements at a given pixel i,

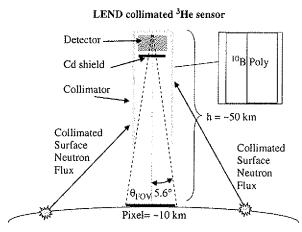


Fig. 1. A single LEND 3 He collimated detector configuration and orbital factors defining a nadir pointing collimation process. Lunar surface flux of neutrons incident the collimator exterior are passively discriminated using neutron absorbing materials. Collimation yields a pixel size $\sim 10 \, \mathrm{km}$ from 50 km altitude (h).

adversely impacts both the instrument blurring function \mathbf{w} and statistical uncertainty \mathbf{n} estimates at i. As a result, accumulation of neutrons from variable sampling conditions to the map incurs additional systematic errors.

3. Methods

The IR performance study was implemented as a two-phase process using four independent research groups, Section 1, GSFC-IGS is defined as the benchmark "denoising only" approach and establishes the null hypothesis, Each IR method uses different numbers and types of parameters which must be configured for a given application domain. To configure and optimize these methods for LEND map restorations, a preliminary set of 12 synthetic LEND epithermal maps were provided, simulating north, south, polar (NP, SP) and randomized PSR spatial and intensity distributions. Training phase data was distributed to teams and included both degraded synthetic LEND maps g and ground truth maps f. Maps for both the training and blind study phases were produced using LRO mission planning ephemeris for the 382 day primary exploration mission with the assumption of continuous nadir pointing of the LEND instrument and 100% duty cycle. This virtual flight process effectively flew the LEND instrument mission over quasi-randomly derived lunar polar neutron emission models (±) 80-90° factoring ephemeris, 1 second integrations and LEND's expected PSF (Mitrofanov et al., 2009).

A consequence of LEND's collimated design is that the exact surface emission points of detected neutrons are only statistically defined within the FOV. This spatial uncertainty constitutes the present basis of LEND's symmetric, 2-D normal PSF, \mathbf{w} , $N(\mu=0,\sigma_{xy}=1.67~\mathrm{km})$ from 50 km altitude. Where, \mathbf{w} defines the spatial probability distribution of subtended pixels. The PSF is defined theoretically via the 5.6° aperture and collimator length, assuming complete discrimination of incoming epithermal neutrons incident to the external side of the collimator. Only epithermal neutron detection processes are assumed with statistical counting uncertainties and no degradation due to motion blur. The map formation process for training and blind studies is illustrated in the flow diagram, Fig. 2. Examples of the LEND map model components are illustrated in Figs. 1 and 3, where the LEND 382 day coverage

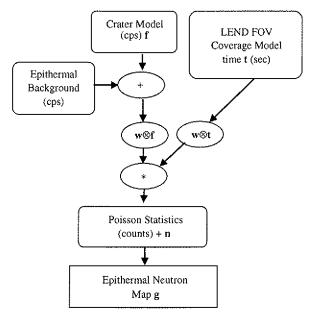


Fig. 2. Flow diagram for generating synthetic lunar epithermal maps.

distribution over the NP region is depicted in Fig. 3a. A synthetic true lunar epithermal map with craters as depressions in the epithermal count rate (dark spots) **f**, is defined in Fig. 3b. Craters delimit sources of enhanced (H) that have lower epithermal flux rates than surrounding background. These regions deviate randomly up to 18% below the expected 0.88 cps collimated rate which is equivalent to the 4.5% suppression seen by LPNS at the lunar north pole. (Feldman et al., 1998b). Crater locations, sizes and intensities were defined in (40) maps randomly, and in known (40) south polar and (40) north polar craters defined by International Astrophysical Union (IAU) conventions (Anderson and Whitaker, 1982; Bussey and Spudis, 2004).

Blurring of the image in Fig. 2 is performed and the product of the blurred time and blurred true image $(\mathbf{t} \otimes \mathbf{w})*(\mathbf{f} \otimes \mathbf{w})$ provides the counts map upon which random Poisson statistical uncertainties \mathbf{n} are added. Factors for randomly deriving the synthetic neutron epithermal maps used in training and single-blind study are listed in Table 1.

3.1. Map restoration methods: iterative Gaussian smoothing (GSFC-IGS)

Gaussian smoothing (GS) has significant precedence in signal denoising applications and has been used to smooth published planetary neutron and gamma-ray geochemistry maps (Mitrofanov et al., 2002; Boynton et al., 2004). Iterative Gaussian Smoothing (IGS) has been shown to be an equivalent process to a diffusion process (Koenderink, 1984; Perona and Malik, 1987). This denoising method systematically smooths data by incrementally attenuating the power in high frequency signals first and progressing towards attenuating the power in lower spatial frequencies Costa and Cesar (2001). In the GSFC-IGS approach, the detected signal, \mathbf{g} is iteratively smoothed j times. For each iteration the map \mathbf{g}_i is convolved with a fixed size, unit volume

Table 1
Synthetic neutron emission map generation factors for training and blind study.

Map size [620 × 620] pixels, pixels 1 km2 ± 80-90°, pole pt [310, 310] +303 km Map evaluation area radius Predefined North Craters (40) 85 (size, position fixed per crater ID) Predefined South Craters (40) 129 (size, position fixed per crater ID) Predefined Random Craters (40) Random Size, Random Position [3, 6, 9, 12, 15, 18] % decrease Crater intensities (random/crater) Accumulated signal scale/382 day [1.0, 1.75, 2.5] x Accum. cover. baseline (random/crater) Background intensity (blobby, variable Mean = 0.88 cps, range 15% of mean smooth) Point spread function (field of view) $N(\mu = 0, \sigma_{x,y} = 1.67)$ Ephemeris 382 days LRO, Beckman-V4, 2007 Spacecraft polar altitude variable [30-60] km, mean = 50 km

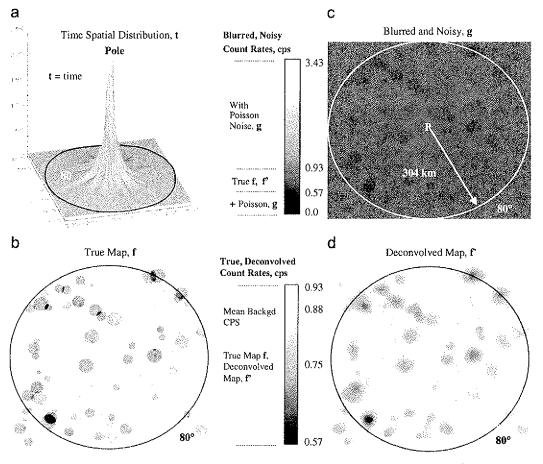


Fig. 3. (a) (Top left) LEND lunar pole accumulated spatial coverage (t=time) distribution 80- 90° for 382 days ephemeris using 1 km² pixels. Simulation projects the LEND FOV in 1 s integrations to a base map. Double lobes in the distribution are due to lunar gravitational anomalies impacting LRO trajectories near the pole (P). (b) (Bottom left) Example true count rate map f/t, time normalized with randomized placement and crater intensity (dark spots, variable intensity) blobby background, μ =0.88 cps, range $\pm 7.5\%$. Range $\pm [0.57, 0.93]$ cps. Craters represent randomly defined, discrete decreases in epithermal count rates, [3-18]%. Area $> 80^\circ$ latitude evaluated by methods=288413 pixels/map, (c) (Top right) Example time normalized blurred and noisy map g/t, after convolution of true map (b) with LEND FOV and added Poisson statistical uncertainties from map used as input to IR methods, Range= $\{0., 3.43\}$ cps. Metrics compare reconstructed maps against the true map (b), (d) (Bottom right) Example of LEND epithermal map after image restoration (PIXON), f'/t. Range= $\{0.57, 0.93\}$.

2-D Gaussian smoothing kernel \mathbf{k} , $\mathbf{g}_{j+1} = \mathbf{g}_{j} \otimes \mathbf{k}$. For each true training map \mathbf{f} , the root-mean-square (RMS) error is quantified at each iteration j comparing true \mathbf{f} and smoothed \mathbf{g}_{j} . For each training map i, the minimal RMS error is determined after each iteration. An average of the IGS RMS error training curves yields [11 iterations] as the optimal iteration setting j to denoise the blind study dataset (GSFC-IGS).

3.1.1. Regularized IR: (GSFC-UM-RD)

The GSFC-UM-RD method optimizes a regularized deconvolution function (DECONVREG) available via the MATLAB® signal processing package. A regularization strategy minimizes the degree of noise amplification incurred from division of small -0, high frequency Fourier coefficients common in Weiner type IR methods (Weiner, 1942). In this approach, inverse filtering is performed using a variant of the least squares method, $\Sigma[\mathbf{g}-\mathbf{w}\otimes\mathbf{f}]^2+\alpha[\mathbf{p}\otimes\mathbf{f}]^2$. Regularization term \mathbf{p} defines a discrete Laplacian distribution convolved with f yielding a smoothed function scaled by α (LRANGE), defining a spatially specific noise minimization constraint. A second parameter in the function is known as the additive noise power, np of the image. np is calculated using a 2-D Fast Fourier Transform (FFT) of g. A center inverse Fourier transform is then applied to the blurred image g. Optimal settings for NP, LRANGE parameters were defined via gridded search of the discretized DECONVREG input parameter space. Optimal input parameter settings [32, 993].

3.1.2. Weiner (UMD-Fourier)

A fast Fourier transform (FFT) based method, developed by (Press et al., 2008; Weiner, 1942) is well known for speed, and has application to signals corrupted by noise and coordinate-invariant point spread function w (Section 3.1). The simplicity of the method lies in the fact that the FFT of the detected image is the product of FFT's of the blurring function w and true image f, $FFT(\mathbf{g}) = FFT(\mathbf{w}) * FFT(\mathbf{f})$ thereby enabling inversion $FFT(\mathbf{f}) = FFT(\mathbf{g}) / FFT(\mathbf{g})$ $FFT(\mathbf{w})$, yielding an estimate of the true image (FFT⁻¹), \mathbf{f}' . However, the inversion becomes problematic in the presence of noise (low SN), and small high frequency coefficients in the denominator that can amplify noise. We address the issue of small FFT coefficient division using additional frequency dependent terms in the denominator [a, b] (a variety of the Wiener filter) as FFT(g)/[FFT(w)+a+b*FFT(f'')]. f'' is an estimate of the true image. This approach optimizes model parameters using the trial images via gridded search of the input parameter space and obtained the following settings [1e-04, 1e-05] (Rosenfeld and Kak, 1982).

3.1.3. Conjugate Gradient (UMD~CG)

Performance of the UMD-Fourier method is dependent on event statistics in the individual pixels. These problems are treated using a Conjugate-Gradient method (CG) (Nguyen et al., 2001). Reconstruction is performed by minimizing the least-square-method (LSM) sum, $\mathbf{S} = \frac{1}{2} (\mathbf{w}^T \mathbf{f} - \mathbf{g})^T \Sigma^{-1} (\mathbf{w}^T \mathbf{f} - \mathbf{g})$ which translates to solving the matrix equation $\mathbf{A}\mathbf{f} = \mathbf{w}^T\mathbf{g}$, where the $\mathbf{A} = \mathbf{w}^T \Sigma^{-1} \mathbf{w}$, and terms \mathbf{w} , Σ , \mathbf{f} and \mathbf{g} are the point spread function, the covariance matrix of event statistics (in our case—a diagonal one), the true(unknown) and convolved images in vectorized form. Although the solution is a simple matrix inversion, its computational complexity, $O(620^2)^2$ precludes processing via conventional computer hardware.

Alternatively, gradient descent methods such as CG can mitigate the inversion limitation by iteratively adjusting method parameters via discrete estimation of parametric gradients. Methods converge in a limited number of iterations assuming initial parametric estimates are proximal to the global solution. For low SN, a regularization term is included, e.g. $\mathbf{A} = \mathbf{w}^T \sum_{i=1}^{n} \mathbf{w} + \mathbf{w}^T \sum_{i=1}$

 γ L, where γ is a positive regularization coefficient for noise suppression, L is a Laplacian circular distribution. The Least Square Method (LSM) sum becomes $\mathbf{S} = \frac{1}{2} (\mathbf{w}^T \mathbf{f} - \mathbf{g})^T \ \Sigma^{-1} (\mathbf{w}^T \mathbf{f} - \mathbf{g}) + \gamma Lt$. Simulations with the trial images yielded optimal reconstructions using [600 iterations, $\gamma = 1000$].

3.1.4. Pixon LLC (PIXON)

The Pixon® method is a picture-information-based method that performs image reconstruction under the constraint that the model for the image should be the simplest model consistent with the data (Puetter et al., 1995, 2005). PIXON gains its power by imposing a minimum complexity constraint on the reconstructed image. Unlike many other competing techniques, the minimum complexity constraint strongly suppresses unwanted image artifacts, as such artifacts are unneeded in fitting the data, and hence are by definition removed. PIXON proceeds by localized smoothing of a minimized set of independent map sub-regions (Pixon elements). Formally, this is described as an integral over the map, $I(y) = \int dz K(y,z) \varphi(z)$, where I is a pseudo-image, K is a positive valued kernel function, φ is a spatial scale adaptive regularizing term sensitive to localized noise estimates and limits high frequency fluctuations at scales < K. PIXON regions are dynamically size adaptive as a function of local uncertainties. Convergence to a solution comprises a parallel optimization seeking largest kernel functions K and minimal least squares fit of the image data for each pixon.

There are several options for implementation and configuring PIXON parameters. For this research *INIT_ANNEAL* and *ANNEAL* parameters were optimized in the training process using a gridded search of the 2-D parameter space [5e-6, 6e-4]. Maps were preprocessed prior to PIXON processing using a spatially adaptive flux preserving filter to identify regionally variable SN for an optimally smoothed result.

4. Results: RMS error + 80-90°

The main goals of this research are to evaluate and quantify the denoising and deblurring properties of IR methods (Section 3.1) on expected LEND epithermal neutron maps. This includes (1) identify IR methods applicable to LEND (2) their parametric configuration (3) potential for induction of artifacts and (4) quantify the error reduction specific to IR vs. the default GSFC-IGS denoising approach. Metrics for method evaluation entail comparing the database of 120 restored maps f against the known synthetic maps f, by averaging the RMS error of map f regions/signals of interest reg, with pixels in regions f, Eq. (1). Secondary metrics include analysis of residual error distributions and pixels quantified as statistical outliers

Region_RMSError =
$$\frac{1}{120} \sum_{i=1}^{120} \sqrt{\frac{1}{n} \sum_{j=1}^{n} (\mathbf{f}_{i,reg}(j) - \mathbf{f}_{i,reg}(j))^{2}}$$
. (1)

Table 2 states each methods mean RMS error results averaged over the 120 blind study maps. Columns indicate region specific evaluations: (1) Global RMS error for \pm 80–90° region followed by the decomposition of each map i into (2) Background, (3) Crater regions. Three best performing methods in each metric are color coded by rank, from lowest to highest RMS error [1=green, 2=yellow, 3=red]. Uncolored cells indicate rank positions > 3rd. Top performing methods: (1) UMD-CG, (2) PIXON and (3) GSFC-IGS. However, PIXON performed nearly as well with first differences only noted at the 7th mean RMS decimal position.

An average of 80% of all map area is allocated to background and 20% to craters, where each region is characterized as follows:

Table 2
Method results: mean RMS error and method rank for 120 maps for 10° estimates 80–90° and method rank in set (1 green (best), 2 yellow, 3 red) (column 1). 10° results and rank using Minkowski L2 metric (Duda et al., 2001) (column 2). Background only mean RMS error (column 3), in crater regions mean RMS error (column 3).

Method \ Metric	10-degree MRMS	10-degree L ² Norm	Background MRMS	Crater MRMS	Category Rank
GSFC-IGS	0.0211180	11.3412	0.0186679	0.0286845	1
GSFC-UM-RD	0.0213130	11.4450	0.0190138	0.0284829	2
UMD CG	0.0210508	11.3052	0.0178985	0.0302985	3
UMD Fourier	0.0217578	11.6848	0.0183534	0.0316352	
PIXON	0.0210573	11.3086	0.0175651	0.0310757	

(2) Background: low amplitude spatial frequencies emulating the expected dynamic range and variation of lunar epithermal surface flux. Best performing methods identified (1) PIXON, (2) UMD-CG and (3) UMD-Fourier. Only small differences were noted between these results < 0.001 cnts/sec, RMS. In background regions the dominant spatial frequency is near the direct current (DC) coefficient. Amplitudes of other higher frequency coefficients are comparatively small so, convergence to the DC baseline is expected for both IR methods and GSFC-IGS. As such, GSFC-IGS performed nearly as well as comparative restoration methods, in 4th place, with small first differences between GSFC-IGS and UMD-CG at the 4th decimal place [0.021118, 0.0210508] cnts/sec. (3) Craters: Table 2, column three defines results of evaluation of localized small spatial scale depressions in the epithermal neutron flux. Cratered regions are of primary interest for NASA's manned lunar exploration programs due to the potential for exploiting localized quantities of (H) as a resource. These regions are inferred to occur primarily in regions of low illumination, but are defined here as cratered regions with variable sized depressions in the lunar surface epithermal neutron flux with variably high spatial frequencies at crater edges.

All method performances were worse than their (1) Global RMS errors, > 0.007 RMS. Relative weighting of cratered and background results are consistent with the contributions factoring area-weights in the 10° latitude band results. Higher RMS errors are seen due to the overlap of higher spatial frequencies in crater rims with noise frequencies at lower latitudes and are considered more difficult spatial frequencies to reconstruct. The ranges of method results indicate performance was consistent across the study, range < 0.003 RMS. Best performers in this category were (1) GSFC-UM-RD, (2) GSFC-IGS and (3) UMD-CG. Importantly, with the exception of GSFC-UM-RD, results indicate that restoration methods did not improve significantly upon GSFC-IGS in cratered regions. Signal gains in cratered regions attributable to IR methods are of significant interest in this study due to the expected higher H concentrations and potentials for human exploration. Results in these higher spatial frequency regions did not identify significant gains vs. the IGS methodology. Additionally, UMD-CG and PIXON, use of a smoothing preprocessing step prior to the restoration step may have incurred additional unrecoverable signal losses in these regions.

4.1. RMS error vs. latitude

Fig. 4 and Table 3 illustrate the RMS error results as a function of latitude. Map SN increases as a function of higher latitude via the combined effects of the time, LRO mission's polar crossing trajectory, altitude and the expected overlapping of the LEND footprint over the poles. As a result the accumulated coverage is highest within $\sim\!3^\circ$ of the pole where the greatest degree of

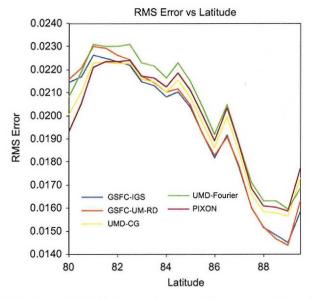


Fig. 4. Five IR methods latitude dependent mean RMS error averaged over 120 blind study images.

ground coverage overlap occurs. The accumulated coverage decreases significantly at latitudes below $\sim 87^{\circ}$, thereby decreasing the SN, Fig. 3a. To evaluate the impact of variable SN on method restorations, the map is decomposed as a function of 20 concentric latitude bands, 0.5° in size (15 pixels wide) in which average RMS error is calculated, Eq. (1). A consequence of this binning of results is that each latitude band area decreases towards the pole yielding unequal weighting of bands in Fig. 4.

Moving from low to high latitude (outside map to center), 80-82° indicates PIXON and UMD-CG methods have superior reconstruction performances. This region constitutes ~36% of the polar region evaluated. At 80°, images are dominated aerially by background regions with exceptions in some included craters in the randomized maps. As a result, due to lower spatial frequencies UMD-CG and PIXON performed best, which was consistent with both the overall and background results described in Table 2. As a fraction of area NP, SP craters are more prominent in the higher latitudes, which decreases the relative distribution of background to cratered pixels in these latitudes. Moving poleward to 81°, all performances are noted to improve, method performances intersect and RMS error decreases due to the higher SN approaching the poles. GSFC-UM-RD, GSFC-IGS results are equivalent, and alternate as the best performers. Trend discontinuities are indicated at 84.5° and 86.5°, which are correlated with the generally higher occurrences of fixed craters for NP, SP maps.

Table 3
Mean RMS errors for 0.5° latitude bands $\pm\,80-90^{\circ}$ indicating method performance and relative rank for restoring 120 blind study epithermal maps. Region from $80-82.5^{\circ}$ constitutes $\sim\!50\%$ of evaluated area.

Latitude	GSFC-IGS	GSFC-UM- RD	U-MD-CG	UMD Fourier	PIXON	Band Area (fraction)	Rank
80	0.0214532	0.0215892	0.0200488	0.0208074	0.0192914	0.09741	1
80.5	0.0216799	0.0220767	0.0210381	0.0219247	0.0205457	0.09251	2
81	0.0226211	0.0230121	0.0223008	0.0230952	0.0220945	0.08753	
81.5	0.0224759	0.0229242	0.0223110	0.0230088	0.0223411	0.08233	
82	0.0223501	0.0226315	0.0222680	0.0230125	0.0223495	0.07746	
82.5	0.0221881	0.0224350	0.0222824	0.0230964	0.0224117	0.07270	
83	0.0214667	0.0217276	0.0215741	0.0222858	0.0217145	0.06750	
83.5	0.0213115	0.0214573	0.0214711	0.0221484	0.0216352	0.06241	
84	0.0208043	0.0210435	0.0210271	0.0216268	0.0212538	0.05756	
84.5	0.0210452	0.0211569	0.0215559	0.0222956	0.0218656	0.05260	
85	0.0203602	0.0204726	0.0208171	0.0214978	0.0211061	0.04755	
85.5	0.0192660	0.0192408	0.0197538	0.0204231	0.0200792	0.04251	
86	0.0181739	0.0182990	0.0186240	0.0191865	0.0189104	0.03734	
86.5	0.0191531	0.0190952	0.0199612	0.0204816	0.0203486	0.03258	
87	0.0177555	0.0178467	0.0184375	0.0188290	0.0187810	0.02751	
87.5	0.0160259	0.0160105	0.0166219	0.0171157	0.0168565	0.02254	
88	0.0151404	0.0151743	0.0158315	0.0163040	0.0160822	0.01741	
88.5	0.0148233	0.0146576	0.0157896	0.0163187	0.0160251	0.01254	
89	0.0145032	0.0143719	0.0156416	0.0159537	0.0158649	0.00751	
89.5	0.0158328	0.0163019	0.0172973	0.0168458	0.0177089	0.00250	

Error trends are elaborated in Table 3 where, the top two performers in each latitude band are highlighted [1=green, 2=red].

4.2. Residuals

Residual error distributions are derived by subtracting the true map from the restored map, r=g'-f. Histograms of r pixels provide the basis for metrics that evaluate the globalized IR reconstruction quality as well as the potential induction of map artifacts. Assuming residuals are normally distributed, a Gaussian fit of the **r** distribution provides both the mean μ_r and variance σ_r^2 terms. Binning of the residual errors was performed at 0.001 cps intervals. As residual pixels are identical in weight and total numbers, ($\sim 3.46 \times 10^7$ total pixels in the study), method residual error histograms have identical area. Direct comparisons of method residual distributions are performed as follows: (1) the residual error distribution mean μ_r deviation from 0. quantifies the existence of bias in reconstructions. (2) Gaussian fit of the r distribution yields, amplitude and variance σ_r^2 terms which are inversely related. Optimally, for best IR performers the fit amplitude is increased and variance is decreased. (3) Artifact induction is inferred by greater numbers of high error pixels. To quantify this error for each method, the integral error in r histogram tails (a) outside $\pm > 3\sigma_r$ is derived. σ_r from the (PIXON) method (lowest variance) was used as the baseline, $\sigma_{\text{PIXON}} = 0.0163$ cps.

Fig. 5 illustrates fits of each methods residual error distribution. Results of the bias study (1) indicated low bias for all methods ≤ 0.0022–0.2% of the mean detector count rate 0.88 cps. Best performing methods were (1) GSFC-UM-RD, (2) GSFC-IGS and (3) UMD-CG. The amplitude/variance metrics are coupled due to their mutual dependence. These results identified: (1) PIXON, (2) UMD-CG and (3) GSFC-IGS as having the highest amplitude/ lowest fit variance. Values for each methods residual error metrics are stated in Table 4.

Table 5 defines method performances in the analysis of residual tail areas. In each case the Tail Hi areas were approximately twice the areas of the Tail Lo areas. This is assumed due to Poisson statistics generating fewer 0. cut-offs for pixel counts. Relative areas in these regions are approximately Tail Lo: 1.3% and Tail Hi: 2.4% of the total pixels processed for all methods. For normally distributed data, each tail above $\pm 3\sigma$ should contain approximately 0.15% of all pixels processed indicating that in all methods tail areas are above the expected area and not normally distributed. Each methods tail areas are integrated above and below the two tail thresholds established using the PIXON baseline positions $\pm 3\sigma_{PIXON}$. For Tail Lo) (1) UMD-CG, (2) UMD-Fourier and (3) PIXON had the lowest tail areas. For Tail Hi) GSFC-UMD-RD had the lowest area followed by GSFC-IGS and UMD-CG. A third metric combining tail areas was established after noting the larger numbers of high intensity pixel

Table 4

Moments from Gaussian fit of residual error. Mean indicates existence of bias as deviation from 0. High amplitude and low sigma indicate higher number of pixels with low error. Rankings indicate all methods displayed low bias < 0.0023 cps, optimal GSFC-UM-RD. PIXON had highest amplitude and lowest variance followed by UMD-CG.

Residual Error Analysis	Method	Amplitude	Mean	Sigma	Category Rank
	GSFC-IGS	748838	-0.0015015	0.0176412	1
	GSFC-UM-RD	735330	-0.0009466	0.0180208	2
	UMD CG	775114	-0.0017858	0.0169150	3
	UMD Fourier	738976	-0.0022218	0.0178038	
	PIXON	797804	-0.0019891	0.0163151	

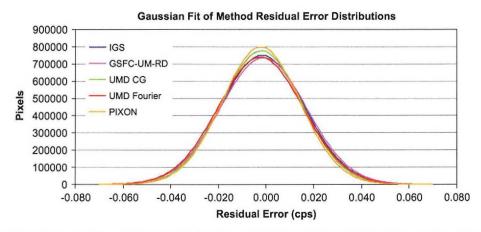


Fig. 5. Gaussian fits of method residual errors. Means indicate low bias for all methods. Variances indicate PIXON has lowest variance (also greatest peak amplitude), however examination of residuals in regions $\pm 3\sigma$ indicate IGS smoothing has the lowest tail area in the high error region.

Table 5
Integration of residual tail areas using (a) Tail Lo=residual errors < -0.048 cps, UMD-CG indicated best performance, (b) Tail Hi residual errors > 0.048 cps, GSFC-UM-RD indicated best performance. Combined tail areas indicated IGS performed best.

Tail Error Analysis	Method	Tail Area Lo	Tails Area Hi	Total Areas	Category Rank	
	GSFC-IGS	431880	728209	1160089	1	
	GSFC-UM-RD	462744	717249	1179993	2	
	UMD CG	393927	802467	1196394	3	
	UMD Fourier	401568	900394	1301962		
	PIXON	405688	844369	1250057		

errors in the Tail Hi) metric. (3) Tail Totals yielded (1) GSFC-IGS, (2) GSFC-UM-RD and (3) UM-CG with the respectively lowest total tail area.

5. Conclusions

This paper described a two-phase blind study evaluation of four IR methods configured to evaluate and quantify restored lunar epithermal neutron flux maps. Background in orbital remote sensing of neutrons and LEND related factors was described with context in image restoration research. Comparison of IR methods was performed using a test database of synthetic LEND maps, derived from signals modeled from existing research and LEND instrument factors. IR methods were directly evaluated against a

traditional smoothing process, GSFC-IGS to quantify the expected gains specific to IR deblurring vs. denoising. A 12-map training phase was included prior to the blind study to familiarize teams with the dynamics of LEND signals and facilitate method optimization. Teams optimized their IR methods on the training dataset via gridded search or simulated annealing methodology by evaluating method performance. GSFC-IGS was trained by establishing the optimal number of smoothing iterations (11) on the training set that minimized the RMS error.

The blind study phase comprised each teams reconstructing the 120 degraded map database using their respective training configurations.

Three types of IR method metrics were described (1) RMS error \pm 80–90°, (2) RMS error vs. latitude and 3) residuals. In (1) only small gains were realized for UMD-CG and PIXON over

GSFC-IGS, < 0.0001 cps RMS. Relative improvements were approximately 2.4% of the expected LEND 0.88 cps collimated count rate. It is inferred this is primarily due to the larger relative distributions of background vs. crater area, where background regions are more homogenous in intensity. However, for cratered regions of interest, GSFC-UM-RD, GSFC-IGS performed slightly better [0.02848, 0.02868] cps than UMD-CG, PIXON [0.0303, 0.0311]. (2) Map latitude analysis indicated method performance varied as a function of the relative distributions of background and cratered regions which varied with latitude. PIXON and UMD-CG performed best in map outer latitudes characterized by lower SN and higher relative numbers of background pixels. GSFC-UM-RD and GSFC-IGS performed better in map polar regions with higher SN and higher relative numbers of cratered pixels, Fig. 4. This result may alternately reflect differences in training method optimization. (3) Residual error analysis via Gaussian fitting of residuals identified PIXON with the best performance at minimizing fit variance / maximizing amplitude [0.016, 797804]. However, the analysis of method tails indicates the GSFC-IGS method best minimized the existence of high error map pixels $> \pm 3\sigma$ vs. reconstructions, < 0.04% tail area difference.

In summary, only small inconsistent signal gains in some metrics was realized by unbiased IR methods over the default GSFC-IGS, specifically in the global metrics. Importantly, in more detailed regional analysis (craters) with higher "true" spatial frequencies, three of four IR methods performance lagged GSFC-IGS in these critical signals. Similarly, all IR methods induced greater numbers of high error pixels (artifacts) vs. the default GSFC-IGS. Finally, IR methods were optimized with ground truth (generally not available in most applications) and yielded only small inconsistent gains vs. GSFC-IGS for the low LEND SN. These results would appear to limit the effectiveness of unbiased IR methods in similar low SN applications. However, methods incorporating other physically based constraints into low SN IR methods may have potential in improving these results but may incur the risk of additional reconstruction errors due to the incorporation of new image modeling assumptions.

Future research will study the use and application of physically based modeling assumptions for LEND map reconstructions.

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