General Disclaimer

One or more of the Following Statements may affect this Document

- This document has been reproduced from the best copy furnished by the organizational source. It is being released in the interest of making available as much information as possible.
- This document may contain data, which exceeds the sheet parameters. It was furnished in this condition by the organizational source and is the best copy available.
- This document may contain tone-on-tone or color graphs, charts and/or pictures, which have been reproduced in black and white.
- This document is paginated as submitted by the original source.
- Portions of this document are not fully legible due to the historical nature of some of the material. However, it is the best reproduction available from the original submission.

Produced by the NASA Scientific and Technical Information Program



RENSSELAER POLYTECHNIC INSTITUTE

TROY, NEW YORK

REPRESENTATION OF PROPAGATION PARAMETERS FOR THE PLASMA IN A MAGNETIC FIELD by R.E. Haskell and E.H. Eolt

Pkasma Research Laboratory TECHNICAL REPORT No. 2, June, 1961

This work was partially supported by the Mational Aeronautics and Space Administration under Research Grant No, NsG 48-60. REPRESENTATION OF PROPAGATION PARAMETERS FOR THE PLASMA IN A MAGNETIC FIELD.

In a recent paper the properties of a uniform plasma for the propagation of electromagnetic waves were represented by curves plotted in the complex propagation plane¹. The effect of a d.c. magnetic field can be shown by extending these curves to three dimensional models in the way described below.

The complex propagation plane is represented by coordinates $A = \alpha$ (c/ ω) and $B = \beta$ (c/ ω) where α and β are the real and imaginary parts of the complex propagation constant and (c/ ω) is the ratio of the velocity of light to the frequency of propagation. In reference 1 the parameters plotted were the normalized collision frequency $Z = \beta/\omega$, where β is the collision frequency and the normalized electron density $X = (\omega_p^{-1} \omega)^2$, where ω_p is the plasma frequency. The equations for the curves were derived from a conformal transformation of corresponding curves in the complex dielectric plane. These equations in turn were derived from the wave equation using the dielectric coefficient for a uniform plasma.

A latter communication² indicated that by suitably choosing the normalized parameters to include the electron cyclotron frequency $\omega_{\rm b}$ the γ same curves could be used for the case of a plasma in the presence of a dc magnetic field. The purpose of this note is to point out that the original curves plotted in the complex propagation plane can be extended to three-dimensional models for the case of a dc magnetic field by introducing a third coordinate $Y = \omega_{\rm c}/\omega_{\rm c}$.

Appleton's equation for the complex refractive index of a magneto-ionic medium is given $\delta_1^{\,3}$

$$n^{2} = 1 - \frac{x}{1 - jZ - 1/2} - \frac{y^{2}}{1 - x - jZ} + \frac{1/4}{(1 - x - jZ)^{2}} - \frac{1/2}{(1 - x - jZ)^{2}}$$
(1)

where X and Z are the normalized parameters defined above and where

$$\begin{split} Y_z &= \omega_z / \omega = \text{the normalized cyclotron frequency in the direction} \\ &\quad \text{of propagation, z} \\ Y_y &= \omega_y / \omega = \text{the normalized cyclotron frequency in the y-direction} \\ &\quad \omega_z &= (e/m)B \cos \Theta \\ &\quad \omega_y &= (e/m)B \sin \Theta \\ \Theta &= \text{angle between the magnetic field B and the direction of} \\ &\quad \text{propagation, z} \end{split}$$

For the case of a longitudinal magnetic field, $Y_j = 0$, $Y_z = Y$ and (1) reduces to

$$n^2 = 1 - \frac{X}{1 - jZ + Y}$$
 (2)

Since the square of the refractive index is equal to the dielectric coefficient, one can rationalize the right-hand side of (2) and equate the real and imaginary parts to K_r and K_i , the real and imaginary parts of the dielectric coefficient K. A conformal transformation from the complex dielectric plane to the complex propagation plane yields the following two equations for the normalized collision frequency case and for the normalized electron density case.

$$B^{2} - A^{2} + \frac{2AB(1 \pm Y)}{7} = 1$$
 (3)

$$(A^{2} + B^{2})^{2} - (2 - \frac{X}{1 + Y}) (B^{2} - A^{2}) + (1 - \frac{X}{1 + Y}) = 0$$

$$(A^{2} + B^{2})^{2} - (2 - \frac{X}{1 + Y}) (B^{2} - A^{2}) + (1 - \frac{X}{1 + Y}) = 0$$

$$(4)$$

For the special case of no megnetic field these equations refduce to Equations (12b) and (13b) of reference 1. They are also consistent with the redefined parameters of reference 2. Flots of equations (3) and (4) are shown in Figs.1, where surfaces are formed for different values of Z and X. The plus sign in (3) and (4) corresponds to propagation of the ordinary wave while the minut sign corresponds to propagation of the extraordinary wave.

This work was partially reperted by MSA under research grant no. NsG 49-70 .

R.E. Haskell E.H. Eolt Rensselae: Polytechnic IInstitute Troy, N. Y.

- M.P. Bachynski, T.V. Johnston, and I.F. Shkevofsky, "Electromagnetic Properties of High-Temperature Air", Proc. IRE, Vol. 48, pp. 347-356; March 1960.
- M.P. Bachynski, T.W. Johnston, and J.P. Shkarofsky, "Universal Representation of Electromegnetic Parameters in Presence of D^o Magnetic Fields," Proc. IRE, Vol. 49, pp. 354-355; January 1961.
- 3. Just Rabiliffe. "The Magneto-Joung Theory and its Applicabions to the "chosphere," Cambridge Univ. Press, 1959.



FIGURE 1

COLLISION FREQUENCY MODEL ORDINARY WAVE



FIGURE 2

COLLISION FREQUENCY MODEL EXTRAORDINARY WAVE



FIGURE 3

ELECTRON DENSITY MODEL ORDINARY WAVE



E

Z

FIGURE 4

ELECTRON DENSITY MODEL EXTRAORDINARY WAVE