Absorption Imaging of Ultracold Neutral Plasmas

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Abstract—We report optical absorption imaging of ultracold neutral plasmas. Imaging allows direct observation of the ion density profile and expansion of the plasma. The frequency dependence of the plasma's optical depth gives the ion absorption spectrum, which is broadened by the ion motion. We use the spectral width to monitor ion equilibration in the first 250 ns after plasma formation. On a microsecond time scale, we observe the radial acceleration of ions resulting from pressure exerted by the trapped electron gas.

Index Terms—Image analysis, plasma measurements, plasmas, spectral analysis.

ULTRACOLD neutral plasmas are systems in which particle temperatures can be 1 K or lower. Fundamental interest in these systems stems from a range of phenomena that occur in the ultracold regime. For instance, it may be possible to form strongly coupled systems in which the electrical interaction energy between the charged particles exceeds the average kinetic energy. This reverses the traditional energy hierarchy that underlies our normal understanding of plasmas based on concepts such as Debye screening and hydrodynamics.

To produce ultracold plasmas, laser-cooled atoms [1], [2] are ionized with photons from a laser whose wavelength is tuned barely above the ionization threshold. Due to the small electron mass, the initial electron kinetic energy E_e approximately equals the difference between the photon energy and the ionization potential, which can be as low as a 100 mK. The initial kinetic energy of the ions is in the millikelvin range. The resulting nonequilibrium plasma evolves rapidly, and many of its characteristics remain to be studied. Experiments [3]–[5] conducted so far reveal surprising dynamics and recombination behavior that show us we are stretching the bounds of plasma physics. Many theory papers have addressed the plasma behavior [6]–[10].

We report results with a new diagnostic tool to study ultracold neutral plasmas: absorption imaging of a strontium plasma using the $Sr^+ 2S_{1/2} \rightarrow 2P_{1/2}$ transition at 422 nm [11]. This technique provides *in situ*, nondestructive measurements and offers excellent spatial, temporal, and spectral resolution. The absorption imaging reported here is particularly well adapted for small, cold, and relatively dilute plasmas that evolve very quickly.

A collimated laser beam, tuned near resonance with the principal transition in the ions, illuminates the plasma and falls on an image intensified charge-coupled device (CCD) camera. Following Beer's law, the optical depth (OD) is defined in terms of

the image intensity without ($I_{background}$) and (I_{plasma}) with the plasma present

$$\mathrm{OD}(x,y) = \mathrm{In}(I_{background}(x,y)/I_{image}(x,y))$$
$$= \int_{-\infty}^{\infty} dz\, n_i(x,y,z)\alpha[\nu, T_i(r)]$$

where $n_i(x,y,z)$ is the ion density, and $\alpha[\nu, T_i(r)]$ is the ion absorption cross section at the image beam frequency, ν. The absorption cross section is a function of temperature due to Doppler broadening, and since we expect the temperature to vary with density, we allow α to vary with position.

Fig. 1 shows a typical absorption image. The intensity of the probe beam is much less than the saturation intensity of the transition. The spatial resolution of typical images is 65 limited by pixel averaging performed to improve the signal-to-noise ratio of the images.

To obtain the absorption spectrum, we plot the integral of the optical depth over x and y as a function of the imaging laser frequency. For longer delay times (t_{delay}) between formation of the plasma and camera exposure, the spectral lines become significantly broader than the natural linewidth of the transition. The dominant contribution to the linewidth beyond the natural linewidth is Doppler broadening, which makes the absorption spectrum a very accurate probe of the ion velocity distribution. From a fit of each spectrum to a Voigt profile using a Lorentzian width equal to the sum of the natural width and the laser linewidth, we extract the root mean square (rms) Doppler broadening, $\sigma_D = \sqrt{(k_B T_{i,\mathrm{eff}}/m_i)}/(\gamma)$, where m_i is the ion mass, and λ is the wavelength of the transition. This provides the effective ion temperature $T_{i,\mathrm{eff}}$ as a function of time. Various effects contribute to the ion kinetic energy and thus to $T_{i,\mathrm{eff}}$.

$T_{i,\mathrm{eff}}$ increases rapidly for $t_{delay} < 200$ ns. Two pieces of information imply that this is thermalization of ions with themselves after creation in a spatially disordered state [6]: 1) the time scale is on the order in which ions respond to perturbations from their equilibrium spatial distribution, which is on the order of the inverse plasma frequency of ions, $\tau_i = \omega_{pi}^{-1} = \sqrt{m_i \epsilon_0 / n_{0i} e^2} \approx 100$ ns, where and are the electric charge and initial peak ion density; 2) the energy is also on the order of $e^2/4\pi\epsilon_0 a k_b$ the amount of potential energy inherent in the initial disorder of the ions. Here, $a = (4\pi n_{0e}/3)^{-1/3}$ is the Wigner–Seitz radius, and is the initial peak electron density.

For $t_{delay} > 200$ ns, $T_{i,\mathrm{eff}}$ continues to increase, but at a slower rate. The slow increase results from an outward radial acceleration of the ions caused by pressure exerted by the gas of trapped electrons [9], [10]. This acceleration will lead to the expansion of the plasma on a 10-μs time scale.

In conclusion, absorption imaging of ultracold neutral strontium plasma is a powerful diagnostic technique. We use this
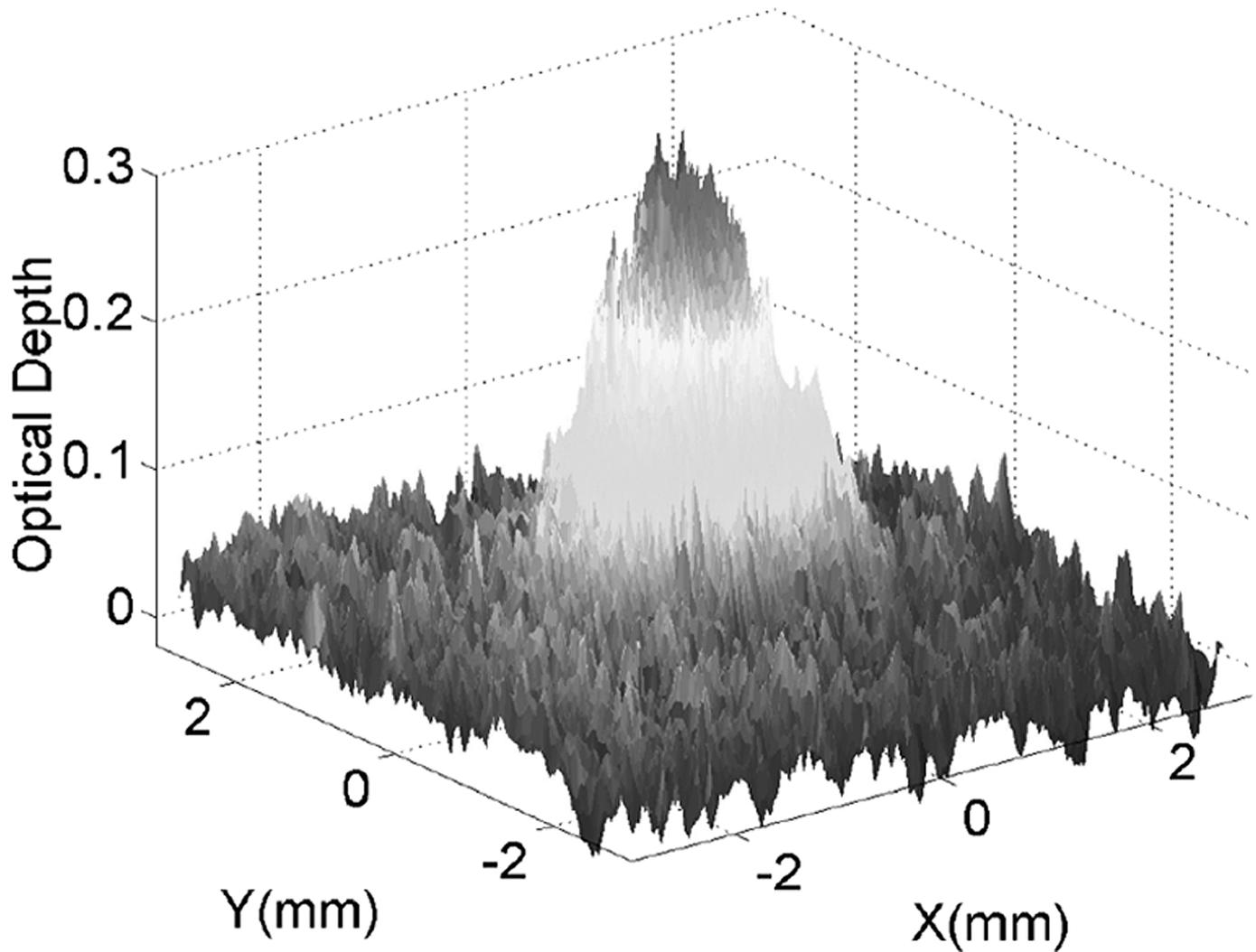
Fig. 1. Optical depth of an ultracold neutral plasma. Delay between the formation of the plasma and image exposure is 85 ns, and the initial peak ion density is $n_{0i} = 2 \times 10^{10}$ cm-3. Plasma contains 7×10^7 ions. Spatial resolution is about $65\,\mu$m, limited by pixel averaging performed to improve the signal-to-noise ratio.

technique to study the ion spectrum at various delay times. The spectrum reveals rapid thermalization of the ions on a submicrosecond time scale as predicted by theory [6], and the acceleration of ions due to electron pressure. In addition, this probe has the potential to study a host of phenomena such as ion collective modes, shock waves, recombination, and particle–particle spatial correlations.

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