

# Letters to the Editor

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## A PRELIMINARY NOTES ON MAGNETIC SUSCEPTIBILITY AND ANISOTROPY IN TETRAHEDRALLY CO-ORDINATED $\text{Co}^{++}$ ION †

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The double chlorides  $\text{CoCs}_2\text{Cl}_4$  and  $\text{CoCs}_3\text{Cl}_5$  both have a slightly distorted tetrahedron of four  $\text{Cl}^-$  ions with the  $\text{Co}^{2+}$  at the centre. The  $\text{Co}^{2+}$  may be considered to be under a predominantly cubic ligand field of the type  $T_d$ , which splits the  $3d^7$   $^4F$  ground state of  $\text{Co}^{++}$  into an orbital singlet  $^4A_2$  and two triplets  $^4T_2$  and  $^4T_1$  each with fourfold degeneracy due to spin so long as the small distortion is neglected. The  $^4A_2$  level lies lowest with  $^4T_2$  and  $^4T_1$  successively at energies  $3500\text{ cm}^{-1}$  and  $6300\text{ cm}^{-1}$  above it (Cotton *et al.*, 1961). Spin-orbit coupling has non-vanishing matrix elements between  $^4T_2$  and  $^4A_2$  and in the presence of the small tetragonal field, caused by the abovementioned distortion taking place along a line joining the opposite edges of the tetrahedron (Powell and Wells, 1935), splits up the  $^4A_2$  level into two Kramer's doublets with separation of the order of  $9\text{ cm}^{-1}$ . (Bowers and Owen, 1955). No mixing of  $^4T_1(F)$  and  $^4T_1(P)$  levels with  $^4A_2$  comes directly except in a very high order through  $^4T_2$ , which can therefore be neglected.

Calculation upto second order of magnetic perturbation  $\beta H(L + 2S)$  yields the expression for principal magnetic susceptibility  $K_i$  ( $i = \parallel$  or  $\perp$  to the tetragonal axis of the ion) which explains well (Table I) the magnetic measurements of mean susceptibility and anisotropy in the range  $300^\circ\text{K}$  to  $90^\circ\text{K}$  by one of us (S.M.) on  $\text{CoCs}_3\text{Cl}_5$ . It is seen that the spin-orbit coupling parameters parallel

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and perpendicular to the tetragonal axis of the ion ( $\zeta_{\parallel}$  and  $\zeta_{\perp}$ ) are reduced to  $-140 \text{ cm}^{-1}$  and  $-120 \text{ cm}^{-1}$  respectively, from the free ion value of  $-180 \text{ cm}^{-1}$  (Laporte, 1928). The corresponding values of orbital reduction factors are  $\kappa_{\parallel} = 0.92$  and  $\kappa_{\perp} = 0.98$ . The cubic field coefficient  $Dq$  is  $350 \text{ cm}^{-1}$ ;  $\Delta$ , the tetragonal field separation of the excited  ${}^4T_2$  level varies from  $322 \text{ cm}^{-1}$  at  $90^\circ\text{K}$  to  $980 \text{ cm}^{-1}$  at  $300^\circ\text{K}$ .

The values of the parameters taken here are consistent with the paramagnetic resonance measurements (Bowers and Owen, 1953) and spectroscopic fine structure results (Cotton *et al.*, 1961).

The anisotropic reduction in  $\zeta$  is (1) due to an anisotropic overlap of the surrounding ligands *s*- and *p*-charge clouds with the central *d*-charge clouds (Bose *et al.*, 1960) and also (2) due to the  $3d-4p$  configurational interaction arising from the non-centrosymmetry of the complex (Bose *et al.*, 1964) The increase in  $\Delta$  from  $322 \text{ cm}^{-1}$  at  $90^\circ\text{K}$  to  $980 \text{ cm}^{-1}$  at  $300^\circ\text{K}$  is due to the thermal expansion or relaxation effects of the lattice. Details of the experimental and theoretical results will be published shortly along with those on  $\text{CoCs}_2\text{Cl}_4$  which have also been completed by us.

TABLE

$$Dq = 350 \text{ cm}^{-1} \quad \zeta = -140 \text{ cm}^{-1} \quad \zeta_{\perp} = 120 \text{ cm}^{-1}, \quad \kappa_{\parallel} = 0.92, \kappa_{\perp} = 0.98$$

Temp $^\circ\text{K}$	$\text{cm}^{-1}$	$p_{\parallel}^2$	$(p_{\parallel}^3 - p_{\perp}^2)$	$g$ values	zero field splitting $\text{cm}^{-1}$
300 $^\circ\text{K}$	980	21.84 (21.81)	1.55 (1.56)		
200 $^\circ$	630	21.16 (21.26)	1.73 (1.72)		
90 $^\circ\text{K}$	322	19.96 (20.00)	2.34 (2.34)	$g_{\parallel} = 2.30$ $(2.32 \pm .04)$ $g_{\perp} = 2.28$ $(2.27 \pm .04)$	8.5 (9.0)

Values in parantheses are the experimental values.

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