

Parton fragmentation in the vacuum and in the medium

Mini-proceedings Workshop ECT*, Trento, Feb. 25 - 29, 2008

S. Albino¹, F. Anulli², F. Arleo³, D. Besson⁴, W. Brooks⁵, B. Buschbeck⁶, M. Cacciari⁷, E. Christova⁸, G. Corcella⁹, D. d'Enterria¹⁰, J. Dolejší¹¹, S. Domdey¹², M. Estienne¹³, K. Hamacher¹⁴, M. Heinz¹⁵, K. Hicks¹⁶, D. Kettler¹⁷, S. Kumano¹⁸, S.-O. Moch¹⁹, V. Muccifora²⁰, S. Pacetti^{20,21}, R. Pérez-Ramos¹, H.-J. Pirner¹², A. Pronko²², M. Radici²³, J. Rak²⁴, C. Roland²⁵, G. Rudolph²⁶, Z. Rúriková²⁷, C. A. Salgado²⁸, S. Sapeta²⁹, D. H. Saxon³⁰, R. Seidl³¹, R. Seuster³², M. Stratmann³³, M. J. Tannenbaum³⁴, M. Tasevsky¹¹, T. Trainor¹⁷, D. Traynor³⁵, M. Werlen³, and C. Zhou³⁶

¹*Institut für Theoretische Physik, Universität Hamburg, 22761 Hamburg, Germany*

²*Sezione di Roma dell'INFN, I-00185 Roma, Italy*

³*LAPTH, Annecy-le-Vieux, France*

⁴*Kansas University, USA*

⁵*Dept. Física y Centro de Estudios Subatómicos, Univ. Técnica F. Santa María, Valparaíso, Chile*

⁶*OA Wien, Austria*

⁷*LPTHE, UPMC Université Paris 6, CNRS Paris, France*

⁸*Institute for Nuclear Research and Nuclear Energy, BAS, Sofia, Bulgaria*

⁹*Museo Storico della Fisica / Centro E. Fermi Roma, and Scuola Normale Superiore / INFN, Pisa, Italy*

¹⁰*CERN, CH-1211 Geneva 23, Switzerland*

¹¹*Charles University, Institute of Particle & Nuclear Physics, 180 00 Prague 8, Czech Republic*

¹²*Institut für Theoretische Physik, Universität Heidelberg, D-69120 Heidelberg, Germany*

¹³*IPHC, Strasbourg, France*

¹⁴*Bergische University Wuppertal, Germany*

¹⁵*Yale University, Physics Department, WNSL, New Haven, CT 06520, USA*

¹⁶*Ohio University, USA*

¹⁷*CENPA 354290, University of Washington, Seattle, 98195, USA*

¹⁸*Inst. of Particle and Nuclear Studies, KEK, 1-1, Ooho, Tsukuba, Ibaraki, 305-0801, Japan*

¹⁹*DESY, Platanenallee 6, D-15738 Zeuthen, Germany*

²⁰*INFN Laboratori Nazionali di Frascati, via E. Fermi, Frascati, Italy*

²¹*Centro E. Fermi, Roma, Italy*

²²*Fermilab, P.O. Box 500, MS 318, Batavia, IL, 60510, USA*

²³*INFN, Sezione di Pavia, I-27100 Pavia*

²⁴*University of Jyväskylä, Finland*

²⁵*MIT, Cambridge, USA*

²⁶*Inst. f. Astro- und Teilchenphysik, Universität Innsbruck, Austria*

²⁷*DESY, Hamburg, Germany*

²⁸*Dept. Física de Partículas and IGFAE, Univ. Santiago de Compostela*

²⁹*M. Smoluchowski Institute of Physics, Jagellonian University, 30-059 Cracow, Poland*

³⁰*University of Glasgow, Scotland*

³¹*RIKEN Brookhaven Research Center, Upton, NY 11973, USA*

³²*University of Victoria, Canada*

³³*Radiation Laboratory, RIKEN, 2-1 Hirosawa, Wako, Saitama, Japan*

³⁴*Brookhaven National Laboratory, Upton, NY 11973, USA*

³⁵*QMWC, Dept. of Physics, London E1 4NS, England*

³⁶*McGill University, Montréal, Canada*

ABSTRACT

We present the mini-proceedings of the workshop on “Parton fragmentation in the vacuum and in the medium” held at the European Centre for Theoretical Studies in Nuclear Physics and Related Areas (ECT*, Trento) in February 2008. The workshop gathered both theorists and experimentalists to discuss

the current status of investigations of quark and gluon fragmentation into hadrons at different accelerator facilities (LEP, *B*-factories, JLab, HERA, RHIC, and Tevatron) as well as preparations for extension of these studies at the LHC. The main physics topics covered were: (i) light-quark and gluon fragmentation in the vacuum including theoretical (global fits analyses and MLLA) and experimental (data from e^+e^- , $p-p$, $e-p$ collisions) aspects, (ii) strange and heavy-quark fragmentation, (iii) parton fragmentation in cold QCD matter (nuclear DIS), and (iv) medium-modified fragmentation in hot and dense QCD matter (high-energy nucleus-nucleus collisions). These mini-proceedings consist of an introduction and short summaries of the talks presented at the meeting.

Contents

1	Introduction	2
2	Theoretical aspects of parton fragmentation in the vacuum: global fits and MLLA	4
3	Experimental aspects of parton fragmentation in the vacuum	12
4	Strange and heavy quarks fragmentation: theory and experiment	20
5	Parton fragmentation in cold nuclear matter: theory and experiment	26
6	Parton fragmentation in hot & dense QCD matter: theory & experiment	29
A	List of participants	35
B	Programme	35

1 Introduction

The transition from coloured quarks and gluons to colourless hadrons – the so-called fragmentation or hadronization process – is a Quantum Chromodynamics (QCD) phenomenon with many important theoretical and phenomenological implications for the physics at high-energy colliders. At large hadron fractional momenta $z = p_{\text{hadron}}/p_{\text{parton}} \ll 1$, the fragmentation functions (FFs) of partons into hadrons or into photons obey DGLAP evolution equations and are obtained from global analyses (fits) of various experimental hadron and photon production data. The FFs are, for example, a basic ingredient for the calculation of the production of high transverse-momentum particles at collider energies within perturbative QCD. At small z , successful QCD resummation techniques (e.g. the Modified Leading Logarithmic Approximation, MLLA) have been developed to understand and describe the evolution of a highly-virtual time-like partons into final hadrons. Most of the available data used in the study of the fragmentation process comes from $q\bar{q}(g)$ production at e^+e^- colliders (LEP, SLC), although data from deep-inelastic $e-p$ scattering (DIS) at HERA has also been used especially in the heavy-quark sector. New high-precision flavour-identified hadron data from the *B*-factories (BELLE, BABAR), from DIS (HERA), and from hadronic colliders (RHIC, Tevatron), can further help to constrain the FFs and, *in*

fine, to improve the perturbative calculation of hadron and photon production at the LHC (in a way similar to the recent developments in improved global-analyses of parton distribution functions). In the near future, the vast kinematical range opened in p - p collisions at LHC energies and luminosities will open up new channels for the study of fragmentation functions.

The study of possible modifications of the parton fragmentation processes in heavy-ion collisions is also an important tool for the determination of the thermodynamical and transport properties of the dense QCD matter produced at RHIC and LHC energies. Since the discovery of the suppression of high- p_T hadrons in central Au-Au collisions at RHIC, a lot of effort has been devoted to the understanding of the propagation of partons in QCD media (cold nuclear matter or quark-gluon-plasma) as well as to understand such a “medium-modified” fragmentation mechanism. These phenomenological studies - which often use techniques originally developed “in the vacuum” (e.g. MLLA) - are well supplemented by a wealth of new data from DIS on nuclear targets (HERMES at HERA, and CLAS at JLab) and in heavy-ion collisions (RHIC).

A few months before the start of the LHC, it seemed a timely moment to have a workshop, gathering both theorists and experimentalists, to discuss the current status of investigations concerning the fragmentation mechanisms of gluons, light-quarks and heavy-quarks into hadrons and photons, as well as to discuss their implications for the upcoming p - p and A-A experimental programmes at the LHC. A meeting was organized at the ECT* (Trento) from February 25–29, 2008, with partial financial support from the center. The meeting had 50 participants whose names and institutes are listed in Appendix A. There were 43 presentations of various lengths. Ample time was left for discussions after each talk. The talks and discussions were organized around the following main topics:

Light-quark and gluon fragmentation in the vacuum ($\bar{e}e$, p - p , e - p collisions): global fits analyses and MLLA.

Strange and heavy-quark fragmentation.

Parton fragmentation in cold QCD matter.

Medium-modified fragmentation in hot and dense QCD matter.

This mini-proceedings include a short summary of each talk including relevant references, the list of participants and the workshop programme. We felt that such a format was more appropriate a format than full-fledged proceedings. Most results are or will soon be published and available on arXiv. Most of the talks can also be downloaded from the workshop website:

http://arleo.web.cern.ch/arleo/ff_vacuum_medium_ect08

We thank the ECT* management and secretariat, in particular Cristina Costa, for the excellent organization of the workshop and all participants for their valuable contributions. We believe that this was only the first workshop of this kind and look forward to similar meetings in the future.

FRANÇOIS ARLEO, DAVID D’ENTERRIA

2 Theoretical aspects of parton fragmentation in the vacuum: global fits and MLLA

Vacuum fragmentation functions: a common interface?

François Arleo, David d’Enterria

A lot of effort has been invested over the past few years in order to gather all existing sets of Parton Distributions Functions (PDFs) under a common interface. This idea was developed in the Les Houches Workshop 2001 and later gave birth to the LHAPDF interface [1, 2], aimed to be the successor of PDFLIB [3]. The essential points of LHAPDF, summarized in [1], are:

PDF sets including their uncertainties should be easily handled, such that any user can quantify the influence of PDF uncertainties on a given observable;

QCD evolution codes can be compared easily;

PDF sets are given under a parametrized form at a soft scale Q and not through heavy grids (yet this is now possible in practice for older sets);

LHAPDF is modular and therefore flexible enough to facilitate the inclusion of new evolution codes or new FF sets.

As shown in the workshop, there is also an important ongoing activity in the FF parametrizations community with e.g. the AKK08, DSS, and HKNS sets recently released, which complement and refine previously existing analyses (AKK, BFG, BKK, KKP, Kretzer, ...). All the existing sets with their main characteristics are summarized in Table 2.

Name	Ref.	Species	Error	z_{\min}	Q^2 (GeV ²)
AKK	[4]	$\pi, K, K_s^0, p, \bar{p}, \Lambda, \bar{\Lambda}$	no	0.1	2 – 4 μb
AKK08	[5]	$\pi, K, K_s^0, p, \bar{p}, \Lambda, \bar{\Lambda}$	yes	0.05	2 – 4 μb
BKK	[6]	$\pi^+ + \pi^-, \pi^0, K^+ + K^-, K^0 + \bar{K}^0, h^+ + h^-$	no	0.05	2 – 200
BFG	[7]	γ	no	10^{-3}	2 – 1.2 μb
BFGW	[8]	h	yes ¹	10^{-3}	2 – 1.2 μb
CGRW	[9]	π^0	no	10^{-3}	2 – 1.2 μb
DSS	[10, 11]	π, K, p, \bar{p}, h	yes ²	0.05-0.1	1 – 10^5
DSV	[12]	polarized and unpolarized Λ	no	0.05	1 – 10^4
GRV	[13]	γ	no	0.05	1
HKNS	[14]	$\pi, \pi^0, K, K^0 + \bar{K}^0, n, p + \bar{p}$	yes	0.01 – 1	1 – 10^8
KKP	[15]	$\pi^+ + \pi^-, \pi^0, K^+ + K^-, K^0 + \bar{K}^0, p + \bar{p}, n + \bar{n}, h^+ + h^-$	no	0.1	1 – 10^4
Kretzer	[16]	$\pi, K, h^+ + h^-$	no	0.01	0.8 – 10^6

Table 1: Main characteristics of fragmentation function (FF) sets obtained from global fit analyses.

¹Errors on the parameters and their correlation are given.

²Typical uncertainties for “truncated” energy fractions $\int_0^1 z D_i^h(z; Q^2) dz$ studied for all flavors and hadron species using the Lagrange multiplier method.

Just like parton densities, the first error analyses in the fragmentation functions have also just appeared. In this context, it appears sensible to examine whether an effort similar to LHAPDF should be performed for the different FF sets. Ideally, a common interface would facilitate the systematic comparison between the different sets, in order to explore for instance the theoretical uncertainties of NLO calculations for hadron production in p - p collisions at the LHC associated with the FFs (which are usually much larger than the uncertainties due to the PDFs). As a very first step, a simple wrapper program should be developed to call any parton-to-hadron FF for any given species. Also, the latest produced FF parametrizations should be made available in a common public web-page such as the existing one from Marco Radici [17].

Albino-Kniehl-Kramer FFs: Improvements from new theoretical input and experimental data

Simon Albino

We present results from an update (AKK08) [5] of the previous AKK fragmentation function (FF) sets [4] for charged pions, charged kaons, (anti)protons, neutral kaons and (anti)lambdas with additional theoretical and experimental input. We incorporate hadron mass effects [18], and fit the hadron mass in the case of the e^+e^- calculation which also has the effects of subtracting out other low p_T and small- x effects beyond the fixed order approach, such as higher twist, small- x logarithms, etc. In the case of the baryons, the fitted masses are about 1% above the true masses, which is consistent with scenarios in which the baryons are produced mainly from direct partonic fragmentation with a small contribution from decays from slightly heavier resonances. A greater excess is found for the pion mass, suggesting contributions to the sample from decays of heavier particles such as $\rho(770)$. Charged and neutral kaon masses are significantly below their true masses. A possible explanation for this is that there are significant contributions from complicated decay channels, such that the direct partonic fragmentation approach is insufficient. For this reason we do not impose SU(2) isospin symmetry of u and d quarks between charged and neutral kaons, which we do for pions. We implement large- x resummation in the quark coefficient function of e^+e^- reactions using the results from Ref. [19], since this is a simple improvement which modifies the cross section over the whole range in x that we constrain. Large- x resummation is also implemented in the DGLAP evolution of the FFs [5]. This results in a significant improvement in the fit for charged kaons, (anti)protons and (anti)lambdas, while χ^2 is essentially unchanged for neutral kaons and charged pions.

In addition to the constraints of the previous AKK fit, we impose further constraints on the charge-sign unidentified FFs from the data for single inclusive production of identified particles ($p_T \geq 2\text{GeV}/c$) from RHIC [20, 21, 22, 23, 24], the Tevatron [25] and electron-positron reactions below the Z pole mass and in the range $0.05 < x < 0.1$. While the untagged measurements from electron-positron reactions provide excellent constraints for the sums of the charge-sign unidentified FFs for quarks of the same electroweak charges, they do not constrain the remaining degrees of freedom at all. As in the previous AKK fit, these were constrained using quark tagged data, while in the new AKK fit additional constraints are provided by the data from RHIC. These data are also much more sensitive to gluon fragmentation and impose (exclusively) new constraints on the charge-sign asymmetry FFs. Normalization errors were treated as systematic effects, i.e. these errors were incorporated via a correlation matrix. Their weights were fitted analytically and independently of the fit in order to further ascertain the quality of the fit, and their magnitudes were typically found to lie in the reasonable range of 0 – 2.

While the results for the fitted masses suggest that the baryons are the best candidates for studying direct partonic fragmentation, there unfortunately exist some inconsistencies between the calculation and the measurements of the inclusive production of these particles at RHIC: The description of the STAR

data for $\Lambda=\bar{\Lambda}$ fails, while the contribution from the initial protons' valence d quarks to the charge-sign asymmetry for $p=\bar{p}$ from STAR is negative. Furthermore, while the contributions to the production from the fragmentation of the valence and sea quarks of the initial protons for the data from RHIC exhibited the expected behaviour, the contribution from the valence d quark fragmentation to the charge-sign asymmetry in (anti)proton production is negative. All these issues would be better understood in the context of an error analysis of the FFs.

We compared our sets to the recent DSS [10] / DSV [12] and HKNS [14] ones, and typically found reasonable agreement for favoured FFs but not for unfavoured, and large discrepancies exist at large- x in some cases. Future hadron identified data from BABAR, CLEO, HERA and RHIC will help to clarify these issues and significantly reduce the large uncertainties in much of the FF degrees of freedom.

De Florian-Sassot-Stratmann (DSS) global QCD analysis of fragmentation functions

Marco Stratmann

We present new sets fragmentation functions for pions, kaons [10], protons, and unidentified charged hadrons [11] obtained in NLO combined analyses of single-inclusive hadron production in electron-positron annihilation, proton-proton collisions, and deep-inelastic lepton-proton scattering. At variance with previous fits, the present analyses take into account data where hadrons of different electrical charge are identified, which allow to discriminate quark from anti-quark fragmentation functions without the need of non-trivial flavor symmetry assumptions. The resulting sets are in good agreement with all data analyzed. The success of the global analysis performed here, including observables other than e^+e^- annihilation, stands for an explicit check of factorization, universality, and the underlying framework of perturbative QCD.

Increasingly accessible hadron production data from proton-proton collisions and hadron multiplicities coming from semi-inclusive DIS should not be disregarded as they offer a crucial piece of complementary information that reduces significantly the uncertainties of the resulting fragmentation functions. For instance, we find that the new DIS pion and kaon multiplicities provided by the HERMES experiment effectively constrain the separation between favored and unfavored distributions, a separation that was either not implemented in previous sets or it was based on certain assumptions. The most recent RHIC results provide stringent constraints on the gluon fragmentation function and, in general, on the large z behavior of the other distributions.

The implementation of the χ^2 minimization in our global analysis is numerically fast and efficient and can be straightforwardly expanded to any future set of hadron production data. With the help of the Mellin moment technique [26], the entire analysis was consistently performed at NLO accuracy without resorting to often used approximations for NLO hard scattering cross sections. For completeness we also provide LO sets. An extensive use of the Lagrange multiplier technique [27] is made in order to assess the typical uncertainties in the extraction of the fragmentation functions and the synergy from the complementary data sets in our global analysis.

Hirai-Kumano-Nagai-Sudoh (HKNS) fragmentation functions and proposal for exotic-hadron search

Shunzo Kumano

Fragmentation functions and their uncertainties are determined for pions, kaons, and protons by a global χ^2 analysis of charged-hadron production data in electron-positron annihilation and by the Hessian method for error estimation [14]. The results indicate that the fragmentation functions, especially

gluon and light-quark fragmentation functions, have large uncertainties at small Q^2 . There are large differences between widely-used functions by KKP (Kniehl, Kramer, and Pötter) [15], AKK (Albino, Kniehl, and Kramer) [4], and Kretzer [16]; however, they are compatible with each other and also with our functions if the uncertainties are taken into account. We find that determination of the fragmentation functions is improved in next-to-leading-order (NLO) analyses for the pion and kaon in comparison with leading-order ones. Such a NLO improvement is not obvious in the proton. Since the uncertainties are large at small Q^2 , the uncertainty estimation is very important for analyzing hadron-production data at small Q^2 or p_T ($Q^2; p_T^2 \sim M_Z^2$) in lepton scattering and hadron-hadron collisions. A code is available for general users for calculating the obtained fragmentation functions [14].

Next, it is proposed that fragmentation functions should be used to identify exotic hadrons [28]. As an example, fragmentation functions of the scalar meson $f_0(980)$ are investigated. The f_0 meson is considered as a candidate for an exotic hadron beyond the usual $q\bar{q}$ configuration because its strong-decay width is much larger than the experimental one [29] if it is an ordinary $q\bar{q}$ meson. The radiative decay $f_0 \rightarrow \pi^0 \gamma$ was proposed to find its internal structure [30], and subsequent measurements indicated tetra-quark or $K\bar{K}$ structure. However, its structure has not been clearly determined yet. Here, we investigate the possibility of ordinary $q\bar{q}$, $s\bar{s}$, tetra-quark, $K\bar{K}$ molecule, or glueball via the fragmentation functions of f_0 . It is pointed out that the second moments and functional forms of the u - and s -quark fragmentation functions can distinguish the tetra-quark structure from $q\bar{q}$. By the global analysis of $f_0(980)$ production data in electron-positron annihilation, its fragmentation functions and their uncertainties are determined. It is found that the current available data are not sufficient to determine its internal structure, while precise data in future should be able to identify exotic quark configurations.

Main sources of uncertainty in quark and gluon fragmentation functions into hadrons and photons

Monique Werlen

A precise knowledge of the QCD hadron and photon production is necessary while searching for e.g. a low mass Higgs signal at LHC and Tevatron. Uncertainties on this background must therefore be asserted with care, in particular those due to the measurement of fragmentation functions. In the calculation of inclusive hadron and γ cross sections in pp collisions with the PHOX package – a set of NLO parton-level event generators for large p_T PHOTon, hadron and/or jet X-sections [31] – uncertainties in fragmentation functions (FFs) do matter compared to those from scales and parton distribution functions (PDF) at RHIC and LHC energies [32]. FFs are obtained from fits to inclusive cross sections assuming a functional form at a given scale M_{F0} and evolving it to the scale of the data. The sources of uncertainties are both experimental and theoretical. Experimental uncertainties are due to statistical errors, in particular at large momentum fraction z of the hadron in e^+e^- data. But also systematic effects contribute such as in the data normalization or in the extrapolation (evolution) needed to cover the full z (Q^2) range. Matching theory and data due to binning and cuts settings may also be an issue. Theoretical uncertainties come from the choice of the functional form for the selected z range, the choice of scales, the order of the theory (leading, next-to-leading NLO, resummation) and further parameters like PDFs and α_s .

As shown for the first time in [9], NLO fits to $e^+e^- \rightarrow \pi^0 X$ data from PEP and PETRA constrain the quark fragmentation into neutral pions while those to $pp \rightarrow \pi^0 X$ from ISR and UA2 constrain the gluon fragmentation. BFGW [8] sets of FF into unidentified charged hadrons have been obtained from LEP and PETRA data choosing optimized scales: a rough approximation of the large- z resummation ([19]). Gluon parameters are quite sensitive to the functional form chosen for the quark distribution. A full statistical error analysis is performed while estimation of the “theoretical error” as previously detailed can be obtained by comparison to other FF sets. Important discrepancies with the BKK set [6] have been

shown, specially at large z while the gluon FFs at $z > 0.5$ is found much higher, in accordance with the UA1 data, than in the Kretzer set [16]. Large scale instabilities (specially at low \sqrt{s} and low p_T) affect the phenomenology of inclusive production of pion [33]. Theoretical estimates rely on extrapolations of the FFs outside ($0.75 < x < 0.9$) of the region where they are actually constrained ($0.1 < z < 0.7$) by the data [34]. A preliminary study [35] with JETPHOX shows that the p_T imbalance hadron-jet correlation cross-sections at RHIC energy may be used to constrain the FFs into hadrons in the high z region. For $p_T(h) > 25$ GeV/c and $p_T(jet) > 30$ GeV/c, the three set (BFGW, KKP, Kretzer) give quite different predictions: BFGW is systematically higher than KKP while Kretzer is a factor two lower.

Thanks to recent data from PHENIX and D0 spanning a large $x_T = 2p_T = \sqrt{s}$ range, inclusive prompt photons cross sections from $\sqrt{s} = 23$ GeV to $\sqrt{s} = 1.96$ TeV are now well understood in the NLO QCD framework [36]. Photons can be produced either directly or via a parton-to-photon FF. The latter becomes important at low p_T and high \sqrt{s} . Uncertainties on the parton-to-photon FFs are mainly due to the uncertainty on the gluon FF as quantified by the BFG sets I and II [7]. BFG-I reduces the cross section by up to 10% (a factor 2.5) at $p_T = 3$ GeV at RHIC (LHC) energy compared to BFG II [32]. As shown in [35] the photon-jet correlation at RHIC energy may also be used to constrain the photon FF. Fixing the jet and varying the photon momenta allows for a direct measurement of the photon fragmentation at low z , a region barely accessible in the LEP experiments.

In the calculation of inclusive hadron and photon cross sections in pp collisions with the PHOX programs – NLO event generators (parton level) for large p_T PHOton (hadron or jet) cross-sections (X-sections) [31] –, uncertainties on fragmentation functions (FF) do matter compared to those from scales and parton distribution functions (PDF) at RHIC and LHC energies [32]. There are various sources of uncertainties in the determination of FF from fits to inclusive cross sections where one assumes a functional form at a scale M_{F0} and evolves it to the scale of the data. Experimental uncertainties are due to statistical errors (FF with a large momentum fraction z of the parton in e^+e^- data), systematic, normalization, the probed z and Q^2 range as well as the sensitivity to quark and gluon FF. Matching theory and data (binning, cuts) may also be an issue. Theoretical uncertainties come from the choice of the functional form, the z range of assumptions, the choice of scales, the order of the theory (leading, next-to-leading NLO, resummation) and further parameters like PDFs and α_s .

NLO fits to $e^+e^- \rightarrow \pi^0 X$ data from PEP and PETRA constrain the quark fragmentation into neutral pions while NLO fits to $pp \rightarrow \pi^0 X$ from ISR and UA2 constrain the gluon fragmentation [9]. BFGW [8] sets of FF into unidentified charged hadrons have been obtained from LEP and PETRA data choosing optimized scales (a rough approximation of the large z resummation performed later in [19]). Gluon parameters were quite sensitive to the functional form chosen for the quark distribution. A full statistical error analysis is performed while estimation of the “theoretical error” embedded in the parametrizations and which come from various theoretical assumptions can be obtained by comparing with different sets: There are important discrepancies, specially at large z , with the BKK set [6] while the very low gluon FF at $z > 0.5$ from the Kretzer set [16] seems unfavored by UA1 data. Large scale instabilities (specially at low \sqrt{s} and low p_T) affect the phenomenology of inclusive production of pion [33]. Theoretical estimates rely on extrapolations of the FF outside ($0.75 < x < 0.9$) of the region where they are actually constrained by the data ($0.1 < z < 0.7$) from which they are extracted [34]. A preliminary study [35] with JETPHOX shows that hadron-jet correlation at RHIC energy may be used to constrain the FF into hadron in the high- z region. With $p_T(h) > 25$ GeV/c and $p_T(jet) > 30$ GeV/c, BFGW is systematically higher than KKP [15] while Kretzer is almost a factor two lower than BFGW.

Thanks to recent data from PHENIX and D0 spanning a large $x_T = 2p_T = \sqrt{s}$ range, inclusive prompt

photons cross sections from $\sqrt{s} = 23 \text{ GeV}$ to $\sqrt{s} = 1.96 \text{ TeV}$ are now well understood in the NLO QCD framework [36]. Photons can be produced either directly or via a parton-to-photon FF. The later become important at low p_T and high \sqrt{s} . Uncertainties on the photon FF are mainly due to the uncertainty on the gluon FF as quantified by the BFG sets I and II [7]. BFG I reduce the cross section by up to 10% (a factor 2.5) at $p_T = 3 \text{ GeV}/c$ at RHIC (LHC) energy compared to BFG II [32]. A preliminary study [35] with JETPHOX shows that photon-jet correlation at RHIC energy may be used to constrain the photon FF. With the momentum of the jet fixed, varying the momentum of the non isolated photon allows for a direct access to the photon fragmentation at low z , a region barely accessible in the LEP experiments.

Time-like splitting functions at NNLO in QCD

Sven-Olaf Moch

We have employed relations between space-like and time-like deep-inelastic processes in perturbative QCD to calculate the next-to-next-to-leading order (NNLO) contributions to the time-like non-singlet and singlet quark-quark and gluon-gluon splitting functions for the scale dependence (evolution) of the parton fragmentation distributions $D_f^h(x; Q^2)$. The evolution equations read

$$\frac{d}{d \ln Q^2} D_i^h(x; Q^2) = \int_x^1 \frac{dz}{z} P_{ji}^T(z; \alpha_s(Q^2)) D_j^h\left(\frac{x}{z}; Q^2\right); \quad (1)$$

where x denotes the fraction of the momentum of the final-state parton f carried by the outgoing hadron h , Q^2 is a time-like hard scale and summation over $j = q; \bar{q}; g$ is understood. The time-like splitting functions P_{ji}^T governing Eq. (1) admit an expansion in powers of the strong coupling α_s ,

$$P_{ji}^T(x; \alpha_s(Q^2)) = a_s P_{ji}^{(0)T}(x) + a_s^2 P_{ji}^{(1)T}(x) + a_s^3 P_{ji}^{(2)T}(x) + \dots; \quad (2)$$

where we normalize to $a_s = \alpha_s(Q^2)/(4\pi)$. The leading-order (LO) terms in Eq. (2) are identical to the space-like case of the initial-state parton distributions, which is the so-called Gribov-Lipatov relation [37]. Also at the next-to-leading order (NLO), the functions $P_{ji}^{(1)T}(x)$ are related to their space-like counterparts by a suitable analytic continuation, see e.g. [38]. Following these ideas, we have derived at NNLO the three distinct non-singlet functions $P_{ns\xi}^{(2)T}$ (with $\xi = q; \bar{q}; g$) [39] and the diagonal singlet quantities $P_{qq}^{(2)T}$, $P_{gg}^{(2)T}$ [40] from the corresponding space-like results [41] and [42] using two independent methods. One approach relied on mass factorization and the structure of the infrared singularities, the other implemented the concept of universal (kinematics independent) splitting functions as conjectured in [43]. The off-diagonal three-loop time-like quantities $P_{gq}^{(2)T}$, $P_{qg}^{(2)T}$ are presently still unknown.

Theory of (extended) dihadron fragmentation functions

Marco Radici

Dihadron Fragmentation Functions (DiFF) describe the probability that a quark hadronizes into two hadrons plus anything else, i.e. the process $q \rightarrow h_A h_B X$. They can appear in lepton-lepton, lepton-hadron and hadron-hadron collisions semi-inclusively producing final-state hadrons. A thorough investigation of their formal properties was performed in the last decade, identifying the leading [44] and sub-leading twist [45] terms, and recognizing the possibility of separating the diagonal and interfering contributions for the $(h_A; h_B)$ pair being in different relative partial waves [46].

At present, the most important application deals with DiFF as analyzers of the fragmenting quark spin and regards the extraction of the transversity distribution h_1 in the nucleon [47], whose knowledge

is a basic test of QCD in the nonperturbative domain [48]. The extraction proceeds via the measurement of a single-spin asymmetry, that has been recently performed by the HERMES [49] and COMPASS [50] collaborations. Information on DiFF can be extracted by the $e^+e^- \rightarrow (h_{A1}h_{A2})(h_{B1}h_{B2})X$ process [51], which is currently being measured by the BELLE collaboration [52], or by models [53].

Recently, DiFF have been recognized as a crucial ingredient for canceling all collinear singularities in the NLO calculation of the cross section for $e^+e^- \rightarrow h_A h_B X$ [54]. When DiFF depend only on the fractional energies of the two final hadrons, there is no way to distinguish between the usual $q \rightarrow h_A h_B X$ process and the mechanism $q \rightarrow q_A q_B \rightarrow h_A h_B X$, where the initial parton branches into two other partons q_A, q_B ; each one subsequently hadronizing in one single hadron. The net outcome is that evolution equations for DiFF contain an inhomogeneous term involving single-hadron fragmentation functions [54].

However, most of the experimental information about DiFF consists of spectra in the invariant mass M_h^2 of the hadron pair. We keep the explicit dependence upon M_h^2 inside DiFF and we define the so called extended DiFF (extDiFF). Using the jet calculus technique, we show that this (soft) scale breaks the degeneracy between $q \rightarrow h_A h_B X$ and $q \rightarrow q_A q_B \rightarrow h_A h_B X$ mechanisms, producing for the extDiFF the usual DGLAP evolution [55]. Therefore, we deduce that the NLO cross section for the $e^+e^- \rightarrow h_A h_B X$ process, when differential also in M_h^2 and for the hadron pair belonging to the same jet, can be described as the factorized convolution of extDiFF and of the same kernel used for the single-hadron fragmentation case. Finally, we show that it is possible to separately study the evolution of each single term in the expansion in partial waves, and we give a preliminary numerical example.

Fragmentation functions in e^+e^- collisions at low Q^2

David Kettler

We report measurements of transverse momentum p_T spectra for ten event multiplicity classes of p - p collisions at $\sqrt{s} = 200$ GeV at STAR. By analyzing the multiplicity dependence we find that the spectrum shape can be decomposed into a part with amplitude proportional to multiplicity and described by a Lévy distribution on transverse mass m_T , and a part with amplitude proportional to multiplicity squared and described by a Gaussian distribution on transverse rapidity y_T . The functional forms of the two parts are nearly independent of event multiplicity. The two parts can be identified with the soft and hard components of a two-component model of p - p collisions. This analysis then provides the first isolation of the hard component of the p_T spectrum as a distribution of simple form on y_T .

We analyze the energy scale dependence of fragmentation functions from e^+e^- collisions using conventional momentum measures x_p and ξ_p and rapidity y . We find that replotting fragmentation functions on a normalized rapidity variable results in a compact form precisely represented by the beta distribution, its two parameters varying slowly and simply with parton energy scale Q . Dijet multiplicities are used to constrain these parameters at lower energies. The resulting parameterization enables extrapolation of fragmentation functions to low Q in order to describe fragment distributions at low transverse momentum p_T . These results compared favorably to conventional representations of fragmentation functions.

Finally, we analyze the minimum bias two-particle correlations of p - p collisions. Cuts are used to isolate the jet-like peak in the angular correlations and the corresponding two-particle y_T correlations are compared to the fragmentation function parametrization derived from e^+e^- collisions.

Jet fragmentation at MLLA and beyond

Redamy Pérez-Ramos

The hadronic k_T -spectrum inside a high energy jet is determined including corrections of relative magnitude $O(\overline{\alpha_s})$ with respect to the Modified Leading Logarithmic Approximation (MLLA), in the limiting spectrum approximation (assuming an infrared cut-off $Q_0 = \Lambda_{QCD}$) and beyond ($Q_0 \notin \Lambda_{QCD}$). The results in the limiting spectrum approximation are found to be in impressive agreement with preliminary measurements by the CDF collaboration, unlike what occurs at MLLA, pointing out small overall non-perturbative contributions. Within the same framework, 2-particle correlations inside a jet are also predicted at Next-to-MLLA and compared to previous MLLA calculations. MLLA corrections, of relative magnitude $O(\overline{\alpha_s})$ with respect to the leading double logarithmic approximation (DLA), were shown to be quite substantial for single-inclusive distributions and 2-particle correlations [56,57]. Therefore, it appears legitimate to wonder whether corrections of order $O(\overline{\alpha_s})$, that is next-to-next-to-leading or next-to-MLLA (NMLLA), are negligible or not.

The starting point of this analysis is the MLLA evolution equation for the generating functional of QCD jets [58]. Together with the initial condition at threshold, it determines jet properties at every energies. At high energies one can represent the solution as an expansion in $\overline{\alpha_s}$. Then, the leading (DLA) and next-to-leading (MLLA) approximations are complete. The next terms (NMLLA) are not complete but they include an important contribution which takes into account energy conservation and an improved behavior near threshold. Some results for such NMLLA terms have been studied previously for global observables and have been found to better account for recoil effects. They were shown to drastically affect multiplicities and particle correlations in jets: this is in particular the case in [59], which deals with multiplicity correlators of order 2, and in [60], where multiplicity correlators involving a higher number of partons are studied; in particular, the higher this number, the larger turn out to be NMLLA corrections.

The present study makes use of this evolution equation to estimate NMLLA contributions to our differential observables. It presents the complete calculations of the single inclusive k_T distribution leading to the results published in [61], and extends them to 2-particle correlations inside a high energy jet. The obtained agreement between the CDF results and the NMLLA distributions over the whole k_T -range is excellent. The NMLLA calculation is in particular able to capture the shape of CDF spectra at every jet hardness Q . Conversely, predictions at MLLA prove only reliable at not too large k_T . The domain of validity of the predictions has been enlarged to larger k_T (and thus to larger x since Y is fixed) computing from MLLA to NMLLA accuracy. This agreement further supports the Local Hadron Parton Duality (LPHD) hypotheses [62].

The MLLA, NMLLA equations for correlators were analyzed and solved iteratively [57]. This allowed us to generalize the result previously obtained by Fong and Webber in [63] that was valid in the vicinity of the maximum of the single inclusive parton energy distribution (“hump”). In particular, we have analyzed the regions of moderately small x above which the correlation becomes “negative” ($C < 0$). This happens when suppression because of the limitation of the phase space takes over the positive correlation due to gluon cascading. This region turned out to be larger in NMLLA than in MLLA. Also, the correlation vanishes ($C = 0$) when one of the partons becomes very soft ($x = 1-x = Y = \ln E \Theta = Q_0$). The reason for that is dynamical rather than kinematical: radiation of a soft gluon occurs at *large angles* which makes the radiation coherent and thus insensitive to the internal parton structure of the jet ensemble. Qualitatively, our MLLA and NMLLA result agrees better with available OPAL data than the Fong–Webber prediction. There remains however a significant discrep-

ancy, markedly at very small x . In this region non-perturbative effects are likely to be more pronounced. They may undermine the applicability *to particle correlations* of the local parton-hadron duality considerations that were successful in translating parton level predictions to hadronic observations in the case of more inclusive *single particle energy spectra* [58].

3 Experimental aspects of parton fragmentation in the vacuum

Colour coherence and a comparison of the fragmentation of gluon and quark jets

Klaus Hamacher

Coherence is an important property of gluon radiation in the hadronic final state of inelastic interactions and a direct consequence of the quantum nature of strong interactions. Theoretical calculations observing coherence predict the evolution scales of jets inside of multi-hadronic events. Typically these scales are transverse-momentum-like and fulfill Lorentz-invariance in contrary to the often used jet energies. A clear demonstration of colour coherence is obtained from the measurement of the ratio of the soft hadron multiplicity in narrow cones oriented perpendicular to the plane of a three-jet event and the axis of a two jet event [64,65]. This ratio is predicted as the product of the colour factor ratio $C_A=C_F$ and an event scale r_γ . The r_γ -dependence allows to compare the amount of coherent soft radiation due to the gluon and the quark colour charge. The prediction without any free parameter describes the data very well. A destructive $1=N_c^2$ suppressed term included in r_γ is required to describe the data. The colour factor ratio when fitted to the data is $C_A=C_F = 2.211 \pm 0.014(\text{stat.}) \pm 0.053(\text{syst.})$. The reason for this good agreement of the leading order prediction with the data lies in the soft, large angle hadrons used as testing probes. These hadrons due to their large wavelength are insensitive to higher order structures and also to finite energy and leading particle effects.

A precision study of the overall multiplicity of three-jet events [66,65] also relies on properly defined scales in order to describe the topology dependence of the multiplicity. Two different choices are available as the division of the event multiplicity in a $q\bar{q}$ and a gluon part is ambiguous. The most natural choice describes the data whereas the choice minimising the amount of hadrons assigned to the gluon jet (chosen by OPAL) fails to agree. The colour factor ratio was fit to the slope of the scale dependence of the multiplicity yielding $C_A=C_F = 2.261 \pm 0.014(\text{stat.}) \pm 0.036(\text{exp.}) \pm 0.052(\text{theo.}) \pm 0.041(\text{clus.})$. This result together with a measurement of the β -function of the thrust [67] which agrees to the one of the coupling in NLO restrict the gauge group of strong interactions to SU(3). From the measurement the multiplicity of gluon-gluon colour singlet events can be extracted. This multiplicity rises about twice as fast with the energy scale as in the $q\bar{q}$ -case clearly demonstrating the higher colour charge of the gluon.

The precise agreement of the above discussed measurements to the colour factor expectation leaves little room for further differences in the particle production from gluons and quarks. The measurements mainly admit differences for leading particles. Such are indeed observed for baryons [68], where an excess of about a factor two (i.e. $2 \cdot A \in C_F$) is observed for leading particles. This excess is understood within the string fragmentation model due to the higher number of possible splitting processes into baryons. Cluster models here fail if no fundamental splittings of gluons into diquark-anti-diquark pairs are foreseen.

It has been found in Monte Carlo studies that gluon fragmentation functions obtained from three-jet events and hypothetical gluon-gluon events disagree in specific parts of phase-space [M. Tasevsky].

Therefore, the results obtained from three-jet events were quantified as “biased”. This distinction seems unfounded as three-jet events form the really accessible source of data. Moreover, if the observed differences were essential they need to be brought into accord with the QCD factorisation theorem. It turns out, however, that the discrepancies are consequence of hadronisation and the necessities of the measurement process [69]. The reconstruction of the parton kinematics due to confinement must proceed by a direct identification of the jet with the parton kinematics. However, the hadronisation transition then leads to a smearing of the parton kinematics. Typically this smearing vanishes $\propto 1/E_{jet}$. The parton energy inferred from the jets enters in the definition of the scaled momentum $x_h = p_{hadron}/p_{g=q}$ and in turn especially strongly varying fragmentation functions are influenced by error propagation. In fact the effect is stronger for the gluon compared to the quark fragmentation function due to the stronger fall-off, as well as due to the somewhat increased hadronisation smearing for gluon jets. A numerical calculation shows ratios of up to a factor two between the smeared, measurable fragmentation function and the underlying partonic one for gluons at high x_h and small E_{jet} . The difference detected for gluon fragmentation [M. Tasevsky] can be fully explained by the hadronisation/measurement process. It is important to note that similar problems unavoidably exist in any measurement involving jets, e.g., when gluon fragmentation is measured in $p\bar{p} \rightarrow 2$ jets. The hadronisation smearing cannot be properly included in an experimental unfolding as it will depend on the fragmentation function to be measured. Still, it is possible to use such measurements if the hadronisation smearing is included in the phenomenological fitting procedure of the fragmentation functions.

ALEPH results on quark and gluon fragmentation

Gerald Rudolph, for the ALEPH collaboration

Parton fragmentation is studied in e^+e^- annihilation into hadrons in the energy range 91 - 207 GeV using the ALEPH detector at LEP. Scaling violation in the flavour-inclusive x -distribution of charged particles has been observed which is, for $x > 0.1$, qualitatively reproduced by global NLO parametrisations [70,71]. The data are better described by the QCD Monte Carlo (MC) models. The energy variation of $\langle N_{ch} \rangle$ is well explained by the QCD-MC models and by a two-parameter 3NLO calculation.

Inclusive x distributions of a variety of identified light mesons and baryons have been measured at the Z pole [72, 73]. These data supplemented with p_T and event shape distributions were used to tune the free parameters of the QCD-MC programs JETSET, ARIADNE and HERWIG [72, 74]. The long-standing discrepancy at $p_{T, out} > 1$ GeV/c can be slightly improved by using the new p_T -ordered parton shower [75]. The fragmentation functions of D^+ and B mesons as well as the relative production rates of the 0^- and higher meson spin states in the c and b sector have been measured [76, 77, 78]. To simulate Bose-Einstein correlations in hadronic Z decays, the parameters λ and σ of the BE_{32} model of PYTHIA have been determined [79].

The gluon and quark fragmentation functions have been measured from symmetric 3-jet events at two scales using b -tagging techniques [80, 81]. As expected, the gluon fragmentation function is steeper in x than the quark fragmentation function. The failure of the QCD-MC models to describe the $x > 0.4$ region of the gluon fragmentation function, first noted in [81, 82], is confirmed with high statistics in 3-jet events of general topologies with high statistics (preliminary). The rate of neutral gluon jets with a central rapidity gap is found to be higher than expected from JETSET or ARIADNE [74]. The excess is located at low invariant mass (0.8-2.2 GeV/c², preliminary), as first observed by DELPHI [83]. The colour-reconnected versions of these generators predict a much higher rate and thus can be excluded [74].

Colour flux studies in quark and gluon fragmentation in the DELPHI experiment at LEP

Brigitte Buschbeck, for the DELPHI collaboration

The occurrence of colour reconnection between two quark-antiquark (colour triplet) strings [84], Bose-Einstein correlations of their decay products [85], and the possible occurrence of colour octet fragmentation with the formation of gluonic systems in the fragmentation of gluons [86] is studied in three contributions.

The first contribution investigates colour reconnection directly in $e^+e^- \rightarrow W^+W^- \rightarrow (q_1\bar{q}_2)(q_3\bar{q}_4)$, i.e. between the colour strings spanned by each W. The second contribution investigates Bose-Einstein correlations (BEC) between the decay products of the two different W bosons. Both effects are determined by the space-time overlap of the two colour flux tubes or their decay products. The data are compatible with some reconnection effects, however with a large statistical error. Bose-Einstein correlations between the decay products of different W's are seen by DELPHI (significance 2.4σ), however they show up weaker than predicted by a Monte Carlo which includes full inter W correlations. Since all other LEP2 experiments do not show any signal, it is argued, that BEC may be diminished or even absent between the hadronisation products of different strings. This can be a valuable tool to study the string structure in hadron-hadron and heavy ion reactions.

The third contribution investigates the leading hadrons in gluon fragmentation in 3-jet ($e^+e^- \rightarrow q\bar{q}g$) events [86]. It shows that gluon jets produce more neutral leading systems than predicted by the LUND string models (JETSET, ARIADNE). This is however predicted, if gluons (g) fragment in some cases as a colour octet field which is then neutralised by a gg pair. The effect is enhanced by demanding a rapidity gap and amounts at a rapidity gap $\Delta y = 1.5$ of about 10%. The concentration of the excess of the neutral systems at low invariant mass could be an indication for a hitherto undetected fragmentation mode of the gluon via the formation of a gluonic system.

Fragmentation functions of quark and gluon jets as measured by OPAL

Marek Tasevsky, for the OPAL collaboration

Scaling violations of quark and gluon jet fragmentation functions are studied in e^+e^- annihilations at $\sqrt{s} = 91.2$ and 183–209 GeV using data collected with the OPAL detector at LEP. The scale dependence of the flavour inclusive, $udsc$ and b fragmentation functions from unbiased jets is measured at $\sqrt{s} = 45.6$ and 91.5–104.5 GeV. Biased jets are used to extract the flavour inclusive, $udsc$ and b , and gluon fragmentation functions in the ranges $Q_{jet} = 4\text{--}42$, $4\text{--}105$ and $4\text{--}70$ GeV, respectively, where Q_{jet} is the jet energy scale. Three methods are used to extract the fragmentation functions, namely the b -tag and energy-ordering methods for biased jets, and the hemisphere method for unbiased jets. The results obtained using these methods are found to be consistent with each other. The $udsc$ jet results above the scale of 45.6 GeV, the gluon jet results above 30 GeV (except for the scale of 40.1 GeV), and the b jet results at all scales except 45.6 GeV represent new measurements. The results of this analysis are compared with existing lower energy e^+e^- data and with previous results from DELPHI and OPAL. The overall consistency of the biased jet results with the unbiased jet results suggests that Q_{jet} is a generally appropriate scale in events with a general three-jet topology. The scaling violation is observed to be positive for lower x_E and negative for higher x_E , for all the types of fragmentation functions. The gluon jet fragmentation function exhibits stronger scaling violation than that of $udsc$ jets.

The bias of the procedure used to construct biased jet fragmentation functions is estimated by study-

ing hadron level Monte Carlo generator events. In explaining the observed differences between biased and unbiased jet results, we note the effects of non-negligible masses of hadrons and b -quarks at low scales. Due to the considerable bias found for the gluon jet fragmentation functions in the region of $x_E > 0.6$, precautions should be taken when comparing the biased gluon jet results with theory. The data are compared to the predictions of NLO calculations. In a wide range of scaled momentum x_E , all calculations satisfactorily describe the data for the $udsc$ jet fragmentation functions. The description is worse and the spread between the predictions larger for the b and gluon jet fragmentation functions, in particular in regions of very low and high x_E . The data are also compared with predictions of three Monte Carlo models, PYTHIA 6.125, HERWIG 6.2 and ARIADNE 4.08. A reasonable agreement with data is observed for all models, except for high x_E region with small scales ($\sqrt{s} = 14$ GeV) in case of the $udsc$ and gluon jet fragmentation functions. The charged particle multiplicities of $udsc$, b and inclusive hadronic events are obtained by integrating the measured fragmentation functions. All values are found to be in agreement with previous measurements, where available.

The first experimental study to use the jet boost algorithm, a method based on the QCD dipole model to extract properties of unbiased gluon jets from $e^+e^- \rightarrow q\bar{q}g$ events, has also been presented. We test the jet boost algorithm using the Herwig Monte Carlo QCD simulation program, comparing the results of this method to those derived from unbiased gluon jets defined by hemispheres of inclusive gg events from a color singlet point source. We find that the results of the jet boost algorithm for the multiplicity distribution are in close correspondence to those of the gg hemispheres for jet energies E_g larger than about 5 GeV. For the fragmentation functions, the results of the two methods agree to good precision for $E_g \lesssim 14$ GeV. Therefore, the fragmentation function of unbiased gluon jets have been measured only at two points, namely at 14.24 and 17.72 GeV.

Parton fragmentation studies in *BABAR*

Fabio Anulli, for the BABAR collaboration

The *BABAR* detector [87], operating at the asymmetric B -factory PEP-II, has been optimized for CP violation studies in B meson decays. However, the high luminosity and the high detector performances, allow also precise measurements of different aspects of strong interactions. We present preliminary measurements of inclusive momentum spectra of a variety of light mesons and baryons, namely π^+ , K^+ , η , and $p=\bar{p}$, at an energy of $\sqrt{s} = 10.54$ GeV, below the $B\bar{B}$ production threshold. The data cover the full scaled momentum range with high precision (few percent relative), allowing sensitive tests of QCD calculations and fragmentation models. Comparison with results from higher energies shows significant scaling violation at high scaled momentum values, for pions and kaons, while data for $p=\bar{p}$ do not present any clear difference between 10 and 90 GeV. The same spectra have been measured also in hadronic decays of the $\Upsilon(4S)$, providing significant inputs to model the inclusive properties of B meson decays.

We present also a measurement of the inclusive spectra of the lightest charmed baryon, the Λ_c^+ [88]. The scaled momentum distribution is in agreement with the existing data and shows a maximum at $x_p \approx 0.6$, followed by a sharp decrease. The total measured production rate is $N_{\Lambda_c^+} = 0.057 \pm 0.002(\text{exp}) \pm 0.015(\text{BF}) \Lambda_c^+$ per hadronic event, where the dominant uncertainty is given by the branching fraction into the reconstructed mode, $\Lambda_c^+ \rightarrow pK^-\pi^+$. The production of $\Lambda_c^+\bar{\Lambda}_c$ pair has been measured at a rate about four times higher than expected from the measured single Λ_c^+ production rate; the event topology is not consistent with four-baryons production, while an average of about four popcorn mesons are produced in association with the $\Lambda_c^+\bar{\Lambda}_c$ pair. This result is compatible with long distance baryon number conservation, as already pointed out by the CLEO experiment [89].

The very high luminosity at the B -factory allows the study of e^+e^- annihilation into exclusive, low-multiplicity final states. Such processes dominate at low \sqrt{s} , but their cross sections fall much faster than $1/s$, and are several order of magnitudes smaller at $\sqrt{s} \approx 10$ GeV. Their measurement in terms of total cross section and internal structures, provide a rich testing ground for QCD. As an example, we show here the first observation of e^+e^- annihilation in two hadrons via two-virtual-photon-annihilation (TVPA), in the processes $e^+e^- \rightarrow \rho^0\rho^0 \rightarrow \pi^+\pi^-\pi^+\pi^-$ and $e^+e^- \rightarrow \phi\rho^0 \rightarrow K^+K^-\pi^+\pi^-$, whose final states have positive charge conjugation, $C = +1$, and cannot be produced by a single γ [90]. Annihilation of e^+e^- pairs at low \sqrt{s} is studied in *BABAR* via initial state radiation, as shown in detail in [S. Pacetti].

Light-quark and spin-dependent fragmentation functions at BELLE

Ralf Seidl, for the BELLE collaboration

The BELLE experiment [91] at the asymmetric e^+e^- collider KEKB [92] has taken a large amount of data which can be used for the study of fragmentation functions. In particular the more than 60 fb^{-1} of continuum data contain only $u\bar{d}s\bar{c}$ quark-antiquark pairs and no B meson decays. In leading order, the normalized yield of inclusive identified hadrons can be directly related to the sum of fragmentation functions of all accessible quarks and antiquarks. With the statistics available systematic uncertainties will be dominating. While the momentum scale and smearing are very well understood the amount of misidentified hadrons has to be corrected. For this purpose particles where the type of the final state particles is known from the decay $D \rightarrow \Lambda$ are used to obtain the particle detection efficiencies and fake rates. Results on the unpolarized fragmentation functions are expected soon.

A second type of fragmentation functions studied in BELLE treats the fragmentation of transversely polarized quarks. The most prominent of these fragmentation functions is the Collins function [93]. Nonzero results from a 29 fb^{-1} continuum data sample have been published [94]. The statistics of these results have been increased by nearly a factor of 20 as for these measurements also the data taken under the $\Upsilon(4S)$ could be included due to a vanishing B -meson background in this analysis.

A global analysis of the published BELLE Collins asymmetries and the semi-inclusive DIS results from HERMES [95] and COMPASS [96] containing the Collins effect in conjunction with the quark transversity distribution were successfully performed by Anselmino and collaborators [97]. The obtained transversity distribution seems to be smaller than expected from bounds and the Lattice QCD predictions. However many aspects, such as the evolution of the Collins function are still open questions. Another spin dependent fragmentation function, the interference fragmentation function is also being studied at BELLE and results are expected soon.

Fragmentation studies at CLEO

David Besson, for the CLEO collaboration

Over the last two decades, CLEO has made a number of measurements testing fragmentation of quarks in $q\bar{q}$ events vs. digluon fragmentation vs. three-gluon fragmentation at center of mass energies around 10 GeV. Not reported in my talk are older results on suppression of decuplet baryons to octet baryons in inclusive $q\bar{q}$ production on the continuum, inclusive pion, kaon and proton fractions in 3-gluon production vs. $q\bar{q}$ production. Reviewed results include:

1. The older correlated $\Lambda_c = \bar{\Lambda}_c$ enhancement observed on the continuum [89], now being redone with higher statistics by BABAR, which showed the original enhanced baryon-antibaryon production

2. Similar $\Lambda=\bar{\Lambda}$ correlated production [98], which showed a similar effect, even after correcting for c -meson cascades. As with the previous process, JETSET 7.4 underestimates this enhancement by about a factor of 3.
3. Inclusive deuteron production in gluonic decays of the narrow $b\bar{b}$ Υ resonances [99], which indicates a 3-fold enhancement of per-event deuteron production in gg events compared to $q\bar{q}$ events.
4. A comparison of particle production quark vs. gluon fragmentation, using $q\bar{q}$ (γ) ISR events on the continuum to tag $q\bar{q}$ events of a particular Q^2 vs. $gg\gamma$ events from the narrow resonances to tag gg fragmentation at the same Q^2 . After confirming the large excess of baryons in ggg production relative to $q\bar{q}$ production, (about double the JETSET 7.4 prediction), the photon-tagged analysis indicates only a modest (and less momentum-dependent) baryon enhancement in gg events, and better consistency with JETSET 7.4.
5. A comparison of direct photon production in quarkonium decays. This, in turn, depends on the color-octet and color-singlet characteristics of the recoil gg states. Newer data on the direct photon momentum spectrum from the ψ (presented first at this conference) are consistent, once accounting for two-body radiative decays, with older data from the Υ [100]. This is somewhat surprising, given the larger relativistic corrections in the charmonium system. We also note that the rate of $gg\gamma=ggg$ is about a factor of 50% larger than the value that extrapolates directly from the Υ , taking into account the running of α_s from the Υ to the ψ mass region.
6. A comparison of the average charged multiplicity in two-gluon vs. $q\bar{q}$ events [101], which showed that, at our center of mass energies, the average multiplicities are consistent with unity, to within 4%.
7. Inclusive production of charm from the χ_b states ($J = 1$; decaying into $q\bar{q}$, with a nearly on-shell gluon). The level of D^0 production observed (approximately 10%) is consistent with a recent color-octet+color-singlet calculation of Braaten *et al.* [102].

Particle production and fragmentation studies at H1

Daniel Traynor, for the H1 Collaboration

Results from H1 were presented for the average charged track multiplicity and for the normalised distribution of the scaled momentum, x_p , of charged final state hadrons as measured in deep-inelastic $e-p$ scattering at high Q^2 in the Breit frame of reference [103]. The analysis covers the range of photon virtuality $100 < Q^2 < 20000 \text{ GeV}^2$. The results are compared with e^+e^- annihilation data and with various calculations based on perturbative QCD using different models of the hadronisation process.

The results broadly support the concept of quark fragmentation universality in $e-p$ collisions and e^+e^- annihilation. A small multiplicity depletion compared to e^+e^- is observed at low Q which can be attributed to higher order QCD processes occurring as part of the hard interaction in $e-p$ scattering but not in e^+e^- annihilation. At high Q a large depletion is also observed. In the low and high Q regions, where the comparison to e^+e^- is poor, the Monte Carlo models are able to provide a better description of the data. The results are compared with NLO QCD calculations as implemented in the CYCLOPS program. All three parametrisations of the fragmentation functions used in this program fail to describe the scaling violations seen in the data.

Results from a study of mini jets in deep-inelastic electron proton scattering were presented with the aim of finding evidence for hadronic activities in excess to those expected from the primary interaction [104]. The analysis is performed separately for the inclusive jet sample, which defines the towards

direction, and for a dijet sample where the second jet is required to have an azimuthal angle larger than 140 degrees with respect to the leading jet, which defines the away region. The dijet sample is split into two samples which are enhanced in direct photon and resolved photon processes. The transverse region between the toward and away regions should be sensitive to additional hadronic activity. The results are compared to various QCD based models. The analysis covers the range $5 < Q^2 < 100 \text{ GeV}^2$, the leading jet is required to have a $p_{T,jet} > 5 \text{ GeV}$ as are both dijets, the mini jets must have a $p_T > 3 \text{ GeV}$.

An overall good description of the data in the “toward” and “away” regions is given by all models for both the inclusive sample and the dijet sample, which proves that the models are able to describe the primary process. In the transverse region the predictions of the models without multiple interactions are generally undershooting the data for the inclusive sample and for the dijet sample with resolved photon events. The inclusion of multiple interactions improve the agreement significantly but not completely.

Particle production at ZEUS

David H. Saxon, for the ZEUS collaboration

Results are presented and compared to theory and models on a range of topics using part or all of the full HERA luminosity of 0.5 pb^{-1} [105]:

Charged Multiplicities and scaled momenta. The use of the Breit frame for comparison to e^+e^- results and the choice of scale are explained. The benefits of using the hadronic centre of mass frame are shown. Scaled momentum distributions are shown for $10 < Q^2 < 40980 \text{ GeV}^2$ and compared to MC and MLLA predictions based in LEP fits. Scaling violation is presented.

Strange Particle Production. \bar{K}^0 ; Λ and $\bar{\Lambda}$ distributions are presented. Compared to expectation there is an excess of Λ 's in resolved photoproduction.

Anti-deuteron and antiproton production. Signals identified by $dE=dx$ are shown. \bar{p} and p rates are equal whereas \bar{d} production is suppressed.

Charm fragmentation and F_2^c . D^{*+} ; D^+ ; D^0 ; D_s and Λ_c^+ signals are presented. Fragmentation parameters are extracted and are mostly consistent with e^+e^- results, with an excess of Λ_c .

Excited charm and charm-strange mesons. Signals are seen for a range of mesonic states. Angular distributions indicate the likely spins.

Baryons decaying to strange particles. Signals are seen for $\Lambda(2286)$; $\Lambda(1520)$; $\Xi(1320)$; $\Sigma^2(1385)$ and other states.

K_s^0 Bose-Einstein correlations. The formalism is recalled and comparison made to K^+K^- . The confusion caused by $f_0(980)$ production is described.

Jets and MLLA studies at the Tevatron

Alexandre Pronko, for the CDF collaboration

The evolution of jets is driven by the emission of gluons at very small transverse momenta with respect to the jet axis. Analytical predictions are based on Next-to-Leading Log Approximation (NLLA) [58] calculations supplemented with the hypothesis of Local Parton-Hadron Duality (LPHD) [62]. NLLA provides an analytical description of parton shower formation, while LPHD states that the hadronization process takes place locally, and, therefore, properties of hadrons and partons at the end of the parton

shower are closely related. Detailed studies of jet fragmentation allow one to better understand the relative roles of perturbative parton showering and non-perturbative hadronization in shaping the main jet characteristics. The Tevatron data presents a unique opportunity to verify the validity and consistency of the NLLA+LPHD approach on a broader range of jet energies than that available at other machines. Overlap of the energy scales between Tevatron and e^+e^- experiments allows for a direct comparison of experimental results obtained in different environments.

The inclusive momentum distributions of charged particles, $\frac{1}{N_{jets}} \frac{dN}{d\xi}$, in jets were measured for dijet events in a wide range of dijet masses, 80-600 GeV/c² [106]. The analysis was done for particles in restricted cones around the jet direction ($\theta_C = 0.28, 0.36, 0.47$). The data were found to agree with theoretical calculations in the NLLA+LPHD framework, in the region where NLLA is indeed applicable. A fit of the shape of the distributions yields the NLLA cutoff scale $Q_{eff} = 230 \pm 30$ MeV and the ratio of charged hadrons to the number of partons produced in a jet $K_{LPHD}^{charged} = 0.56 \pm 0.10$.

The charged particle multiplicities in quark and gluon jets, N_q and N_g , were obtained by comparing charged particle multiplicities in two data samples: dijet data and photon+jet data [107]. These two samples have a different quark/gluon jet content (60% for dijet and 20% for γ +jet events), which allows for a measurement of the inclusive properties of gluon and quark jets. The evolution of measured charged particle multiplicities in gluon and quark jets with jet hardness scaling variable $Q=2E_{jet} \tan(\theta_c/2)$ agree well with trends predicted by recent 3NLLA calculations [108] if Q_{eff} is set to 230 MeV [106] and normalization constant is let to be only free parameter. We also measured the ratio of charged particle multiplicities in gluon and quark jets, $r = N_g/N_q$. At $Q = 19$ GeV, we obtain $r = N_g/N_q = 1.64 \pm 0.17$, in agreement with 3NLLA calculations [108].

The measurement of the two-particle momentum correlations in jets is based on dijet events with the invariant mass, M_{jj} , in the range from 66 to 563 GeV/c² [109]. The correlation function was introduced [63] in terms of the parameter $\xi = \ln(E_{jet}/p_{hadron})$ and was defined as a ratio of 2- and 1-particle inclusive momentum distributions: $C(\Delta\xi_1; \Delta\xi_2) = \frac{D(\xi_1, \xi_2)}{D(\xi_1)D(\xi_2)} = c_0 + c_1(\Delta\xi_1 + \Delta\xi_2) + c_2(\Delta\xi_1 - \Delta\xi_2)^2$. The results were obtained for charged particles within a restricted cone with an opening angle of $\theta_c = 0.5$ radians around the jet axis. Overall, the data and theory show the same trends over the entire range of dijet energies. The NLLA fits [63] to the evolution of correlation parameters c_1 and c_2 with jet energy scale allowed to extract the value of the parton shower cutoff scale $Q_{eff} = 137^{+0.85}_{-0.69}$ MeV. The comparison to Monte-Carlo revealed that both Pythia Tune A and Herwig 6.5 reproduce correlation in data fairly well.

The measurement of the intrinsic k_T (transverse momenta of particles with respect to jet axis) distributions is of particular interest because it allows to probe softer particle spectra than previously measured observables. The analysis was based on the same dijet events that were used in the described above measurement of the two-particle momentum correlations in jets. To compare shapes of the k_T distributions in data with the recent theoretical calculations [56], the predictions were normalized to agree with data in the bin $-0.2 < \ln(k_T) < 0.0$. Within the validity region [56], the agreement between data and theory is fairly good. The inclusion of the higher order corrections in theoretical calculations [61] improves the agreement over the wider range of $\ln(k_T)$.

4 Strange and heavy quarks fragmentation: theory and experiment

Evidence for power corrections in heavy quark fragmentation in e^+e^- collisions

Matteo Cacciari

I report on the findings of [110], which compare QCD theoretical predictions for heavy flavoured meson fragmentation spectra in e^+e^- annihilation with data from CLEO [111], BELLE [112] and ALEPH [76]. Several ingredients are included in the calculation: next-to-leading order initial conditions, evolution and coefficient functions; soft-gluon resummation at next-to-leading-log accuracy; a matching condition for the crossing of the bottom threshold in evolution [113] implemented at next-to-leading order accuracy; and important initial-state electromagnetic radiation effects in the CLEO and BELLE data.

It is found that with reasonably simple choices of a non-perturbative correction to the fixed-order initial condition for the evolution, the data from CLEO and BELLE can be fitted with remarkable accuracy. The main result of the analysis is however that the fitted fragmentation function, when evolved to LEP energies, does not reproduce fully well the D fragmentation spectrum measured by ALEPH. Large non-perturbative power corrections to the coefficient functions of the meson spectrum are speculated to be needed in order to reconcile CLEO/BELLE and ALEPH results. The data do not allow, however, to distinguish between a standard $C=Q^2$ correction with a fairly large coefficient ($C \sim 5 \text{ GeV}^2$) or an unexpected linear term, $C=Q$, with a more reasonable coefficient ($C \sim 0.5 \text{ GeV}$).

Heavy quark fragmentation with an effective coupling constant

Gennaro Corcella

The transition of partons into hadrons, for the time being, cannot be calculated from first principles QCD, but is usually described in terms of phenomenological models, containing parameters which are to be fitted to experimental data. Alternatively, one can model non-perturbative corrections by means of an effective strong coupling constant, which does not exhibit the Landau pole any longer and includes absorptive effects due to parton branching [114].

The talk discusses bottom- and charm-quark fragmentation in e^+e^- annihilation, using the effective-coupling model to include hadronization effects. Perturbative quark production is treated in the framework of perturbative fragmentation functions [115], which factorizes the production of a heavy quark as the convolution of a coefficient function, describing the emission off a massless parton, and a perturbative fragmentation function, associated with the transition of a light parton into a massive quark. The perturbative fragmentation function follows the DGLAP evolution equations, whose solution allows one to resum the large logarithm of the heavy-quark mass exhibited by the fixed-order calculation. In our computation, we use NLO $\overline{\text{MS}}$ coefficient functions and NLO initial condition of the perturbative fragmentation function. The DGLAP evolution equations are implemented in the non-singlet sector, to NLL accuracy, and threshold resummation is performed in the NNLL approximation, in both coefficient function and initial condition. The effective coupling constant, evaluated to NNLO accuracy, is the only source of non-perturbative corrections. In order to account for the theoretical uncertainty on our predictions, we vary the parameters entering in the perturbative computation, such as renormalization and factorization scales, the b -quark mass and $\alpha_S(m_Z)$, within conventional ranges.

Following [116], we consider data on B -hadron energy distributions at the Z^0 pole, collected by

LEP and SLD experiments. In particular, the ALEPH collaboration reconstructed only B mesons, while OPAL and SLD had also a small fraction of b -flavoured hadrons. We find that our calculation, provided with the effective-coupling model, is able to give a reasonable description of the considered data, within the experimental and theoretical uncertainties. We also compare our model with data in Mellin moment space, measured by the DELPHI collaboration, and obtain fair agreement with the first five experimental moments, within the errors. We then investigate data on charmed hadrons from ALEPH and B -factory experiments, along the lines of [117]. Our result is that the effective-coupling model is able to acceptably reproduce the ALEPH D^+ spectrum, while serious discrepancies are present when comparing with D and D^0 data from CLEO and BELLE. In fact, other analyses had previously shown that, even with a flexible non-perturbative fragmentation function, provided with three tunable parameters, one was not able to reproduce at the same time both ALEPH and B -factory D -data. Nonetheless, we still manage to describe quite well the first ten Mellin moments of all considered data from ALEPH, CLEO and BELLE.

Towards a model independent approach to kaon fragmentation functions

Ekaterina Christova³

We derive different measurable quantities in kaon production in $e^+e^- \rightarrow K+X$ and in Semi-inclusive DIS (SIDIS) $eN \rightarrow eK+X$ that measure certain Non singlet (NS) combinations of quark fragmentation functions (FF) for kaons [118]. These quantities are obtained in a model independent way and hold in any order of QCD. In SIDIS, if charged K are measured we have, $\sigma_N^{K^+K} = \alpha_N^{K^+} \alpha_N^K$:

$$\sigma_p^{K^+K} = (4u_V + D_u + d_V + D_d)^{K^+K} (1 + (\alpha_s=2\pi))C_{qq} \quad (3)$$

$$\sigma_d^{K^+K} = (u_V + d_V) (1 + (\alpha_s=2\pi))C_{qq} + (4D_u + D_d)^{K^+K} \quad (4)$$

that present two measurables for the two unknown FFs: $D_u^{K^+K}$ and $D_d^{K^+K}$. In particular, this allows one to test the usually made assumption $D_d^{K^+K} = 0$. Recently very precise data on unpolarized SIDIS was presented by HERMES [119]. If both K and K_s^0 are measured, then the combination $K^+ + K = 2K_s^0$, both in e^+e^- and SIDIS kaon production, always measures the same NS combination of FFs $(D_u + D_d)^{K^+K}$, $(\sigma^{K^+K} + 2\sigma^{K_s^0} = \sigma^{K^+} + \sigma^K + 2\sigma^{K_s^0}; \tilde{q} = q + \bar{q})$:

$$d\sigma^{K^+K} + 2\sigma^{K_s^0} = 6\sigma_0 (\hat{e}_u^2 + \hat{e}_d^2) (1 + \alpha_s C_q) D_u^{K^+K} \quad (5)$$

$$d\sigma_p^{K^+K} + 2\sigma_s^{K_s^0} = [(4\tilde{u} + \tilde{d}) (1 + \alpha_s C_{qq}) + \alpha_s g C_{gq}] D_u^{K^+K} \quad (6)$$

$$d\sigma_d^{K^+K} + 2\sigma_s^{K_s^0} = [(\tilde{u} + \tilde{d}) (1 + \alpha_s C_{qq}) + \alpha_s g C_{gq}] D_u^{K^+K} \quad (7)$$

These results can be used to determine the kaon FFs. Being model independent, they allow also one to test the parametrizations already obtained, but at different assumptions. In addition, eqs. (5)-(7) allow to compare the FFs obtained in e^+e^- at rather high $Q^2 \sim m_s^2$, eq. (5), with those of SIDIS at quite low Q^2 , eqs. (6) and (7). Note that as $D_u^{K^+K}$ is a NS, its Q^2 evolution does not involve any other FFs.

Heavy Quark Fragmentation Functions in e^+e^- annihilation at $\sqrt{s} = 10.5$ GeV

Rolf Seuster, for the BELLE collaboration

Over the years, QCD was found to be a reliable theory for many predictions. It is however restricted to regions of phase space where the scale of the hard interaction was (much) larger than any other scale in the particular event. In pushing the limit to its extreme, one learns more about the theory, its failures

³Work supported by I3HP EU network.

and its successes. One limit that can be challenged at the very successfully operated b -factories is the measurement of the heavy quark fragmentation function. At b -factories, charm quarks are the heaviest quarks for which such measurements can be done, as b 's are produced at threshold. For the measurement presented here, the scaled momentum $x_p = \frac{p_{candidate}}{p_{max}}$ is used throughout. The BELLE detector [91] is an experiment operated at the asymmetric e^+e^- collider KEKB [92] at a center of mass energy of about 10.6 GeV. For this analysis, about 450 events below and 352 events at the $\Upsilon(4S)$ resonance were used, corresponding to 15 fb^{-1} and 88 fb^{-1} of integrated luminosity. As the contribution from B hadron decay will only occur below $x_p = 0.5$, combining on- and off-resonance data above this threshold decreases significantly the statistical uncertainty. It allows for a very precise and detailed description of the FF of charmed quarks into charmed hadrons, as shown in [112]. The difference of the on- and off-resonance data below $x_p = 0.5$ is solely due to the contribution from B decays. The numbers reported in [112] are in good agreement with the world average. These high precision data will have to be tested against Monte Carlo generators. Such a comparison will improve the theoretical understanding of the underlying physics, as well as the description of the data by the MC itself. One example for such comparisons is the determination of how good commonly used FFs describe the data and what parameters of these functions to choose. Via a reweighting technique, special MC events were generated at many parameter points for many fragmentation functions. In general, the best agreement was found for the Bowler and the Lund FFs. The parameters at which these functions showed the best agreement can be found in [112]. The commonly used Peterson *et al.* fragmentation function showed by far the overall worst performance, not agreeing with the high precision data at all. For details, see [112].

Another example, for which such comparisons could be made for are ratios of fragmentation functions. In the ratio, their systematic uncertainties will partially cancel. Whereas for the x_p dependent ratio of $x_p(D^+) = x_p(D^+)$, the Monte Carlo can be tuned to give a perfect description of the data, the ratios $x_p(D_s) = x_p(D^+)$ and $\Lambda_C = x_p(D^+)$ are completely mis-modelled in the Monte Carlo. No parameter will improve the description significantly. No conclusive reason for this discrepancy is known.

Charmed particles production in $e^+e^- \rightarrow c\bar{c}$ at 10.6 GeV

Simone Pacetti, for the BABAR collaboration

During the last few years many new charmed mesons and baryons have been discovered. Some of these, especially in the charmed-strange mesons sector, were unexpected. The theoretical predictions based on potential models and the heavy quark effective theory (HQET) [120] fail in describing mesons like the $D_{sJ}(2317)^+$, $D_{sJ}(2460)^+$, $D_{sJ}(2860)^+$ and $X(2690)^+$. We focused on the analyses of such states:

$D_{sJ}(2317)^+$ and $D_{sJ}(2460)^+$ have been observed in the decays $D_{sJ}(2317)^+ \rightarrow D_s^+ \pi^0$ [121] and $D_{sJ}(2460)^+ \rightarrow D_s^+ \gamma$, $D_s^+ \pi^0 \gamma$, $D_s^+ \pi^+ \pi^-$ [121]. These analyses have been realized over the data collected by the BABAR experiment in inclusive $c\bar{c}$ production at 10.6 GeV, with an integrated luminosity of 232 fb^{-1} . Masses, widths, and decay modes of these mesons do not fit the HQET predictions based on the $c\bar{s}$ structure, hence there are speculations that both $D_{sJ}(2317)^+$ and $D_{sJ}(2460)^+$ are some type of exotic state, such as four-quark or molecular states [122].

$D_{sJ}(2860)^+$ and $X(2690)^+$ have been observed in their decay into $D^{0(+)}K_{(S)}^+$, with $D^0 \rightarrow K^+ \pi^-$, $K^+ \pi^+ \pi^0$, and $D^+ \rightarrow K^+ \pi^+ \pi^+$ [123]. This analysis is realized in 240 fb^{-1} of data collected by the BABAR experiment. The evidence of the lightest structure is less significant, however it is compatible with a similar structure observed by BELLE [124].

Before the advent of b factories the spectrum of charmed baryons comprised only 12 states [125]. In the last years BABAR found 4 new charmed baryons and confirmed 3 other states observed by BELLE.

The Ω_c is a $c\bar{s}s$ baryon with $J^P = 3/2^+$, it is the first excited state of the Ω_c family. It has been observed by *BABAR* in the radiative decay: $\Omega_c \rightarrow \Omega_c \gamma$, in $c\bar{c}$ production with an integrated luminosity of 231 fb^{-1} [126]. The mass difference with respect to the ground state Ω_c is in agreement with lattice calculations [127] and theoretical expectations [128].

The $\Delta(2940)^+$ is a udc baryon, it has been observed by *BABAR* decaying in $D^0 p$ [129]. It is the first observation of a charmed baryon which decays in a light baryon, i.e.: the proton. The data for this analysis refer to an integrated luminosity of 287 fb^{-1} .

Two new charmed-strange baryons $\Xi_c(3055)^+$ and $\Xi_c(3123)^+$, belonging to the qsc -family ($q = u, d$) have been discovered by the *BABAR* experiment, studying the $\Lambda_c^+ K \pi^+$ system, obtained in inclusive $c\bar{c}$ production at 10.6 GeV with 384 fb^{-1} [130]. The decay in the $\Lambda_c^+ K \pi^+$ final state goes through different intermediate states, namely: $\Xi_c(3055)^+ \rightarrow \Sigma_c(2455)^{++} K \rightarrow \Lambda_c^+ \pi^+ K$ and $\Xi_c(3123)^+ \rightarrow \Sigma_c(2520)^{++} K \rightarrow \Lambda_c^+ \pi^+ K$. No neutral partners have been observed.

ISR Physics at *BABAR*

Simone Pacetti, for the BABAR collaboration

Using the initial state radiation (ISR) technique in fixed center-of-mass energy machines, like the flavour factories, we may mimic a typical e^+e^- annihilation experiment with an energy scan in the center of mass. At the Born approximation, in the differential cross section for processes like $e^+e^- \rightarrow \gamma_{\text{ISR}} + \text{hadrons}$, the function describing the radiation of a photon (γ_{ISR}) by one of the initial leptons and the hadronic cross section factorizes. This allows us to perform precise measurements of hadronic cross sections at low energy (from threshold up to 4-5 GeV). The knowledge of $R = \sigma(e^+e^- \rightarrow \text{hadrons})/\sigma(e^+e^- \rightarrow \mu^+\mu^-)$ is crucial to compute the hadronic contributions to the muon anomalous magnetic moment [131] (low-energy values), and to the running electromagnetic coupling α_{QED} [132] (high-energy values). In addition we perform a better parameter determination for many ρ -, ω -, and ϕ -excited states ($J^{PC} = 1^{--}$).

The light-meson final states measured by the *BABAR* experiment in the 1-3 GeV region are: $\pi^+\pi^-\pi^0$, $2(\pi^+\pi^-)$, $K^+K^-\pi^+\pi^-$ [133], $3(\pi^+\pi^-)$, $2(\pi^+\pi^-\pi^0)$, $K^+K^-2(\pi^+\pi^-)$ [134].

A new narrow ($\Gamma \approx 60 \text{ MeV}$) vector meson called $X(2175)$ has been discovered decaying in $\phi f_0(980)$ [135] and confirmed in the $\phi\eta$ channel [136, 137].

We have analyzed the reaction $e^+e^- \rightarrow KK\pi$ [137] in two samples: $K^+K^-\pi^0$ and $K_S^0K^-\pi^+$. The study of the $K_S^0K^-\pi^+$ asymmetric Dalitz plot allowed, for the first time, the identification of the two isospin components and their relative phase. We performed a global fit, exploiting five different data sets, to describe the total $KK\pi$ cross section in terms of intermediate resonances. We determined parameters for ϕ - and ρ -excited states, with a special care to the OZI-suppressed $\phi\pi^0$. The $\phi\pi^0$ cross section, with a peak value of $\approx 0.2 \text{ nb}$, shows two resonant structures: a broad bump at $\approx 1.6 \text{ GeV}$, compatible with a ρ recurrence [125] and a narrower resonance at $\approx 1.9 \text{ GeV}$, whose parameters are in agreement with those of the ‘‘dip’’ observed in multi-hadronic final states [138].

The *BABAR* experiment measured with unprecedented accuracy the differential cross section for $e^+e^- \rightarrow p\bar{p}$ [139]. In particular, by studying the angular distribution, *BABAR* provides also a time-like measurement for the ratio $\mathcal{G}_E^p/\mathcal{G}_M^p$ of the electric and magnetic proton form factors. We observed, for the first time, a displacement from the ‘‘scaling behavior’’: $\mathcal{G}_E^p/\mathcal{G}_M^p \neq 1$, by measuring $\mathcal{G}_E^p/\mathcal{G}_M^p$ near threshold. The angular distribution provides also a powerful tool to estimate the two-photon exchange contribution in the annihilation $e^+e^- \rightarrow p\bar{p}$. Using the *BABAR*

data this contribution turns out to be compatible with zero [140]. In addition, the $p\bar{p}$ cross section has two very interesting features: a non-vanishing threshold value, which is however expected as a consequence of the $p\bar{p}$ Coulomb interaction, and the presence of structures (“steps”), not yet understood, at $\sqrt{s} = 2.2$ and 3.0 GeV.

Other baryon-antibaryon final states like: $\bar{\Lambda}\Lambda$, $\bar{\Sigma}^0\Sigma^0$ and $\bar{\Lambda}\Sigma^0$ [141], have been measured by *BABAR* with the ISR technique. In all these cases the cross section data have finite threshold values, as in the $p\bar{p}$ cross section, which are not well understood. Indeed, having always neutral baryons, no Coulomb interaction effect can be used to cancel the vanishing of the phase space.

Study of charm fragmentation in DIS at H1

Zuzana Růriková, for the H1 collaboration

The process of charm fragmentation is studied in deep-inelastic D meson production as measured by the H1 detector at HERA [142]. Two different methods are used for the definition of the observable approximating the fraction of the 4-momentum of the D meson with respect to the charm quark. The c -quark momentum is approximated in the γ p -rest-frame either by the momentum of the jet including the D meson or by the momentum of a suitably defined hemisphere which includes the D meson.

The visible phase space of this analysis is given by cuts on the photon virtuality $2 < Q^2 < 100$ GeV², event inelasticity $0.05 < y < 0.7$ and D meson phase space $1.5 < p_T(D) < 15$ GeV/c, $|\eta(D)| < 1.5$. In addition a D jet with $E_T > 3$ GeV in γ p -rest-frame is required in order to have hard scale for the events. The fractional energy distributions of D mesons, corrected for detector effects, are used to extract parameters for the Kartvelishvili and Peterson fragmentation functions within the framework of leading order + parton shower QCD models RAPGAP 3.1 and CASCADE 1.2 with string fragmentation and particle decays as implemented in PYTHIA, as well as for the next-to-leading order QCD calculation in the fixed flavour scheme HVQDIS [143].

The fragmentation parameters extracted using both observables are in good agreement with each other. Furthermore the value of Peterson parameter ϵ , extracted for the PYTHIA parameter setting [144], is in agreement the ALEPH tuned value of $\epsilon = 0.04$, supporting the hypothesis of fragmentation universality between e^+e^- and $e-p$. Finally, the fragmentation of charm produced close to the kinematic threshold is studied. The description of this sample by the QCD models and NLO calculation is not good, and the extracted fragmentation parameters are significantly different from the parameters fitted to the nominal sample with a hard scale. This can be interpreted as a hint of inadequacies of the QCD models and possibly of insufficient flexibility of the simple parametrisations used for the fragmentation function in the phase space region close to kinematic threshold.

Inclusive $K_S^0 K_S^0$ resonance production in $e-p$ collisions at HERA

Changyi Zhou, for the ZEUS Collaboration

The lightest glueball is predicted to have $J^{PC} = 0^{++}$ [145, 146] and a mass in the range 1450–1750 MeV/c². Thus, it can mix with $q\bar{q}$ states from the scalar meson nonet, which have $I = 0$ and similar masses. In the literature, the state $f_0(1710)$ is frequently considered to be a state with a possible glueball or tetra-quark composition (see for example the reviews [147, 148]).

Inclusive $K_S^0 K_S^0$ production in $e-p$ collisions dominated by photoproduction with exchanged photon virtuality, Q^2 , below 1 GeV², at HERA has been studied with the ZEUS detector using an integrated

luminosity of 0.5 fb^{-1} from HERA-I and HERA-II combined. K_S^0 mesons were identified through the charged-decay mode, $K_S^0 \rightarrow \pi^+ \pi^-$, and reconstructed with various track quality cuts for selection. The $K_S^0 K_S^0$ invariant mass distribution was reconstructed by combining two K_S^0 candidates. Enhancements in the mass spectrum have been observed attributed to the production of $f_2(1270)/a_2^0(1320)$, $f_2^0(1525)$ and $f_0(1710)$. The three states $f_2(1270)$, $a_2^0(1320)$ and $f_2^0(1525)$ are all of $J^P = 2^+$ and so their interference is seen in the total cross-section. The intensity is the modulus-squared of the sum of these three amplitudes plus the incoherent addition of $f_0(1710)$ and the non-resonant background. The amplitudes for $f_2(1270)$, $a_2^0(1320)$ and $f_2^0(1525)$ production were fixed at the SU(3) ratios of $5 : 3 : 2$ as expected for production via an electromagnetic process [149, 150]. Very competitive measurements on peak position and width for $f_2^0(1525)$ and $f_0(1710)$ are done with interference fit and the overall fit describes the data very well. Complete systematic checks have been performed. The final values with statistical and systematical uncertainties were compared well with the PDG values [125].

RHIC constraints on fragmentation functions for strange hadrons in p - p collisions at $\sqrt{s} = 200 \text{ GeV}$ *Mark Heinz, for the STAR collaboration*

Perturbative QCD has proven successful in describing inclusive hadron production in elementary collisions. Within the theory's range of applicability, calculations at next-to-leading order (NLO) have produced accurate calculations for the p_T spectra of charged hadrons in different collision systems and energies, which has led to claims of universality of the underlying fragmentation functions (FFs) [15]. With the new mid-rapidity p - p data at $\sqrt{s} = 200 \text{ GeV}$ collected at RHIC, comparisons to pQCD can now be extended to identified strange baryons and mesons [24]. In the last 5 years significant improvements have been made in the field of NLO FF studies. Several groups have updated their parameterizations to include not just $e^+ e^-$ but also (SI)DIS and now p - p data, thus improving the constraints on the FF parameters. In particular the gluon-to-hadron FF was never well constrained by $e^+ e^-$ and DIS data due to the low probability of gluon-jet production. Conversely, at RHIC we are probing a very gluon-jet rich environment and the majority of final states are produced by gluons. In particular baryons are dominated by gluon fragmentation. According to AKK, 90% (10%) of the protons below $p_T = 10 \text{ GeV}/c$ are produced from gluons (quarks). On the other hand, pions are produced in equal parts by gluons and quarks.

Experimentally this fact can be exploited to measure the baryon-to-meson ratios and compare them to LO/NLO calculations. Results by the STAR collaboration confirm that LO models are not able to accurately describe this ratio [22, 24] and underpredict the amount of baryons at low p_T . For strange baryon production (Λ) the first NLO predictions equally fail to describe the data [12]. In an attempt to solve this shortcoming, one group (AKK) [4] readjusted the gluon FF to fit the data, finding $D_g^\Lambda = D_g^p = 3$. Another group (DSS) has for the first time performed separate fits to protons and anti-protons based on STAR data [11], and provided new FFs for pions, kaons and charged hadrons from a global analysis [10]. Another subject which did not receive sufficient attention up until recently are error estimations on the FF parameterizations. One group (HKNS) [14] has now published a global (Hessian) error analysis for charged hadrons and finds that the 4 most ubiquitous FF agree within errors. However this analysis does not include p - p data. The STAR experiment has also measured identified particle spectra in terms of $m_T = \sqrt{p_T^2 + m^2}$ as motivated by earlier studies [151]. After arbitrary normalization of the different particle species, this representation clearly shows different shapes of the spectra for the baryons and mesons for $m_T > 2 \text{ GeV}/c^2$. Interestingly, a simulation of the same observable with PYTHIA (LO pQCD) shows a similar feature when selecting gluon-jet enriched events [75]. Finally, efforts at RHIC are now being focused in the extraction of medium modified FFs. Several approaches are being pursued in parallel, such as di-hadron correlations, 3-particle correlations and also full jet reconstruction.

5 Parton fragmentation in cold nuclear matter: theory and experiment

p_T -broadening in nuclear DIS

Hans-Jürgen Pirner

My work deals with modifications of fragmentation processes in cold nuclear matter and in the hot quark gluon plasma. In the first part a phenomenological analysis of deep-inelastic hadron production in nuclei is performed. The main theoretical picture is a three-stage model which includes partonic propagation, prehadron formation and absorption, and finally hadron formation.

This model has been successfully applied to describe the HERMES data presented in [E. Aschenauer]. The multiplicity ratios of $\pi;K;p;\bar{p}$ produced in e -Ne, He, Kr, Xe collisions as a function of the momentum fraction of the hadrons z_h , and the photon energy ν can be modelled by a prehadron-nucleon absorptive cross section of one third of the hadron cross section. No Q^2 -dependence of the multiplicity ratios is calculated as observed in the data.

The second part emphasizes the p_T -broadening in the hadron spectra as a tool to monitor the path length of the quark in the nucleus. The meson-nucleon interactions have a small elastic cross sections relative to the total cross section, and can be neglected for p_T -broadening. Again the calculated differential Δp_T^2 broadening as function of $z_h, \nu; Q^2$ are compared with data. The observed $\log(Q^2)$ dependence indicates the interleaving of parton radiation and parton-medium scattering in the evolution process.

Both effects can be described with a new equation where a splitting kernel and a scattering kernel are present in the evolution equation. Consequences of this equation for the jet profiles in LHC experiment have been shown and further details will be given in [S. Domdey].

Hadronization of pions and kaons from nuclei using DIS

Kenneth Hicks, for the CLAS Collaboration

The CLAS experiment EG2 was run with a variety of nuclear targets using a 5.5 GeV electron beam from the continuous electron beam accelerator facility (CEBAF) at Jefferson Lab. The goal of this experiment is to measure observables related to the propagation of a quark (struck by the virtual photon from deep inelastic scattering) through cold nuclear matter. These results could be contrasted with quark propagation through hot QCD matter, as in the quark-gluon plasma expected to be formed in relativistic heavy ion collisions (RHIC). Quarks from RHIC are tagged by back-to-back “jets” of high energy hadrons formed in hard parton collisions. In contrast with RHIC results, where the quark propagation is severely damped by passage through hot QCD matter, the preliminary DIS results suggest that the struck quark propagates relatively freely (compared with fully formed mesons) through cold nuclear matter.

The measured quantities of the DIS measurements are the squared four-momentum transfer (Q^2) and energy transfer (ν) from the scattered electron and the energy fraction ($z = E_h/\nu$) and transverse momentum (p_T) of the leading hadron. By comparing the z and p_T distributions of pions and kaons from a deuterium target with that from various nuclear targets, the quark propagation and formation into a hadron can be inferred. Theoretical predictions in a given model provide guidance for interpretation of the results. Ratios of the number of hadrons detected, normalized to the number of DIS events in a given kinematic bin of Q^2 and ν , for a nuclear target divided by the same quantities for a deuterium target are one quantitative measure of the effects of propagation through cold nuclear matter. The subtraction of the average value of p_T^2 seen for a nuclear target minus that for a deuterium target is another measure

which, according to theoretical models, is sensitive to gluonic radiation by the quark before it forms into a hadron. By comparing these two quantitative measures with theoretical calculations over a variety of kinematics, then models of hadronization can be tested. The HERMES experiment has already published the nuclear attenuation ratios for several types of hadrons at DIS kinematics in the range $8 < \nu < 23$ GeV. The CLAS data are for a lower range, with $2.6 < \nu < 4.3$ GeV corresponding to shorter formation times. At HERMES kinematics, calculations suggest that the quark propagates through the full nucleus, followed by hadron formation outside of the nuclear radius. At CLAS, we expect to see hadron formation inside the radius of larger nuclei.

Preliminary results for positive pions produced at CLAS for the EG2 experiment indicate a plateau for the difference Δp_T^2 as the nuclear size increases (targets C to Pb). This suggests that, at these kinematics, the average hadron formation length is shorter than the radius of Pb. In contrast, nuclear attenuation ratios cannot be explained by theoretical models unless some finite propagation of a “pre-hadron” (with a smaller cross section than fully-formed hadrons) is assumed. Attenuation ratios for neutral kaons appear to be smaller than those for π^+ particles, but results are still preliminary. In conclusion, the CLAS data provide a new kinematic window for hadronization at the expected length scale near to the radius of heavy nuclei. Statistical uncertainties are small due to an experimental luminosity of about 100 more than for the HERMES data. Both π^+ and K_s^0 nuclear attenuation ratios and Δp_T^2 are expected to be submitted for publication before the end of 2008.

Characterising hadronization with nuclear DIS

Valeria Muccifora

A review on the recent progress in the study of the parton propagation, interaction and hadronization in *cold* nuclear matter has been presented. It has been pointed out that the cleanest environment to address nuclear modification of hadron production is the nuclear Deep Inelastic Scattering: it allows to experimentally control many kinematic variables; the nuclear medium (i.e., the nucleus itself) is well known; the hadron multiplicity in the final state is low, allowing for precise measurements. Moreover, the nucleons act as femtometer-scale detectors of the hadronizing quark, allowing to study its space-time evolution into the observed hadron.

The experimental highlights from deep inelastic lepton nuclear scattering have been shown, in particular results from the HERMES experiment have been discussed. At HERMES, nuclear semi-inclusive deep-inelastic scattering is used to study the quark propagation and hadronization via single hadron multiplicity ratio [152], double hadron production [153] and broadening of the hadron transverse momentum p_T [154]. While hadron multiplicity ratio is sensitive to all stages of the hadronization process, i.e. to the production time t_p of the colorless $q\bar{q}$ system and to the formation time t_f of the final hadron, the p_T -broadening has been shown to be sensitive to t_p because, once the $q\bar{q}$ is formed ($t > t_p$), no further broadening occurs as inelastic interaction is suppressed.

The most recent theoretical frameworks for describing the interaction of energetic partons and space-time evolution of the hadronization process in nuclear DIS have been discussed. In particular models based mainly on parton energy loss effect [155, 156] have been compared with models based mainly on the interaction of the $q\bar{q}$ system in the medium and that are sensitive to the time evolution of the hadronization in the medium [157, 158]. It has been pointed out the importance to study parton propagation and interaction in *cold* nuclear matter for the interpretation of data in A-A collisions. In particular the connections between the kinematic in DIS and in nucleon-nucleon collisions have been discussed and the equivalent phase spaces covered by different experiments have been compared [159].

Parton propagation and hadron formation: present status, future prospects

Will Brooks

The process of hadronization is unique to QCD, reflecting the non-Abelian nature of the strong interaction. For pragmatic descriptions of scattering phenomena, hadronization is captured by the empirically determined fragmentation functions. Until recently, study of the parton propagation and hadron formation processes was limited to asymptotic properties of the final state, such as hadron multiplicities; any microscopic information on the space-time development of the process on the femtometer scale was lost. The microscopic space-time development of hadronization can now be accessed using the nuclear medium as an analyzer. The interactions with the nuclear medium can be used to refine the understanding of the mechanisms at work at 10^{-15} m, yielding new insights into such fundamental processes as gluon emission by a quark and the enforcement of color confinement through hadron formation. Extraction of confinement quantities such as the lifetime of the deconfined quark now appears to be possible. Cross comparison of new data in a variety of reactions will potentially provide an exciting, coherent picture, although several important points concerning the interaction with the medium still need to be clarified. Foremost among these for the semi-inclusive DIS data is the question of whether the observed hadron attenuation is primarily due to prehadron absorption, or due to medium-simulated gluon emission. Further, the extent to which information derived from cold nuclear matter can be used to interpret data from hot nuclear matter continues to be a subject of discussion. In particular, medium-induced gluon emission is broadly believed to be involved in both cases, offering the possibility of constraining the interpretation of the high-energy heavy-ion data using the well-understood properties of cold nuclear matter.

In addition to semi-inclusive DIS, Drell-Yan data offer another venue in which to study propagating quarks in cold nuclear matter, without the additional complication of hadron formation in the final state. In this talk, transverse momentum broadening in Drell-Yan and semi-inclusive DIS were compared, with significant differences between the Drell-Yan data from HERMES [E. Aschenauer, V. Muccifora] and Jefferson Lab [K. Hicks]. While several possibly relevant aspects of these differences are understood, a comprehensive picture awaits further theoretical development. It should be noted that the extremely high precision of both of the latter two, quite new, data sets is unprecedented and these data should stimulate a new wave of theoretical activity. The JLab data for positive pions comes as a fully three-dimensional function of Q^2 , v , and z , while the HERMES results include both π^- and an exploratory look at kaons. Further, the increase of Δp_T^2 in the JLab data appears to saturate at high nuclear mass. This is due to the lifetime of the deconfined quark being shorter than the dimensions of the heaviest nuclei, and a quantitative extraction of this lifetime within model assumptions is clearly feasible. An inter-comparison of the HERMES and JLab results for hadron attenuation shows a good consistency. The landmark HERMES data set has more than half a dozen different hadrons, extending to much higher v than the JLab data, and thus provides the best constraint on the v and flavor dependence. The JLab data, while limited thus far to pions and the first K^0 study, nonetheless offers two orders of magnitude more luminosity, thus permitting three dimensional binning of the multiplicity ratio for, e.g., π^+ . Studies of these data thus far do reveal systematic variations of the multiplicity ratio in Q^2 , v , p_T^2 , and z , with the largest variation in z and in p_T^2 , in agreement with HERMES.

While further studies of the existing data are ongoing, the next significant advance in the field will come with the 12-GeV Jlab upgrade. The upgrade of CLAS will permit yet another order of magnitude increase in luminosity, adding baryon fragmentation studies, as well as a broader range in v , Q^2 , and z . The increased luminosity allows access to formation of rarer and heavier mesons such as the ϕ . In total, eleven mesons and eight baryons are included in the existing plans with an 11 GeV e^- beam.

6 Parton fragmentation in hot & dense QCD matter: theory & experiment

Medium-evolved fragmentation functions

Carlos Salgado

Medium-induced gluon radiation is usually identified as the dominant dynamical mechanism underlying the *jet quenching* phenomenon observed in heavy-ion collisions (see e.g. [160] for a recent review on the calculation of the spectrum and the corresponding phenomenological implementation). In its present implementation, multiple medium-induced gluon emissions are assumed to be independent, leading, in the eikonal approximation, to a Poisson distribution [161]. In [162] we have introduced a medium term in the splitting probabilities

$$P^{\text{tot}}(z) = P^{\text{vac}}(z) + \Delta P(z; \mathcal{P}); \quad (8)$$

so that both medium and vacuum contributions are included on the same footing in a DGLAP approach. In eq. (8), $\Delta P(z; \mathcal{P})$ is taken directly from the medium-induced gluon radiation spectrum

$$\Delta P(z; \mathcal{P}) = \frac{2\pi t}{\alpha_s} \frac{dI^{\text{med}}}{dz dt}; \quad (9)$$

computed in the multiple scattering approximation [163]. The improvements include energy-momentum conservation at each individual splitting, medium-modified virtuality evolution and a coherent implementation of vacuum and medium splitting probabilities. Noticeably, the usual formalism is recovered when the virtuality and the energy of the parton are very large. This leads to a similar description of the suppression observed in heavy-ion collisions with values of the transport coefficient of the same order as those obtained using the *quenching weights*. In a previous publication [164] it has been found that this formalism would lead to non-trivial angular dependences of the jet profiles when kinematic constraints are imposed. Similar structures were found at RHIC in two-particle correlations in the away side [165].

Hadronic composition as a characteristics of jet quenching at the LHC

Sebastian Sapeta

There are several mechanisms which may lead to modification of the hadronic composition of jets (hadrochemistry), *e.g.* exchange of quantum numbers (color, flavor, baryon number) between projectile and the medium, recombination of partons from the jet with partons from the medium as well as recoil effects. All the above may require serious modeling. Instead, we consider the framework which takes into account only the exchange of momentum between the medium and the developing partonic cascade.

To calculate single particle distributions of identified hadrons we use the framework of Modified Leading Logarithmic Approximation (MLLA) [166]. This perturbative approach combined with the hypothesis of Local Parton-Hadron Duality (LPHD) was shown to reproduce correctly not only the distributions of all charged particles but also the spectra of identified hadrons such as pions, kaons and protons [62, 167]. To model the medium-modification of jets we supplement the above formalism by the formulation of parton energy loss proposed in [168]. The effects of medium-induced gluon radiation are introduced by enhancing the singular parts of the leading order splitting functions by the factor $1 + f_{\text{med}}$. This leads to the softening of hadron spectra. The model accounts for the nuclear modification factor at RHIC when f_{med} is of the order of 1.

Our main result are the ratios of $K = \pi$ and $p = \pi$ [169]. We observe a significant difference between the ratios in vacuum and medium jets. According to our analysis, in the presence of medium

the ratios increase up to 50% for kaons and 100% for protons. This enhancement seems to be a generic feature for radiative energy loss jet quenching models. In addition we observe only mild dependence on the the jet energy and the jet opening angle. The medium modified hadron spectra may be also superimposed on the high multiplicity environment of heavy ion collisions expected at the LHC. We have checked that due to the characteristically different hadrochemistry of the jet and the background our result concerning the ratios remains virtually unchanged. Hence, we conclude that the hadrochemical composition of jets may be very fragile to the medium effects and could be used as an additional handle to study microscopic mechanisms underlying the jet quenching phenomenon.

QCD evolution of jets in the quark-gluon plasma

Svend Domdey

Based on our recent paper [170], I briefly discuss the effect of scattering for jet quenching of gluonic fragmentation functions in the quark-gluon plasma. In contrast to the common discussion [163, 161, 171] which treats medium-induced gluon emission as the main mechanism, we find that energy loss from scattering will also contribute significantly. The basic idea given in the talk is as follows: The DGLAP equation is build from the probability for splitting due to change of virtuality. The time scale for this process is $\tau_{split} = E/Q^2$. We extend this equation to contain also scatterings with gluons from the medium by constructing a “scattering probability” from the ratio of the time scale for splitting and the scattering mean free path $\tau_{scat} = 1/(\sigma n_g)$, i.e. the number of scatterings which can happen between two splittings. Similar to splitting the medium-modified evolution equation contains a gain and a loss term.

We discuss two slightly different approaches: One is more suitable for analytical calculations, the other one for implementation in a Monte Carlo study. Both methods give very similar results. Even with a perturbative cross section, we find a suppression of the fragmentation function in medium relative to vacuum which is increasing with energy fraction x . For $Q = 100$ GeV and a plasma temperature of $T = 500$ MeV (which may be relevant for LHC) we find a suppression of 10% at $x \sim 0.1$, but a suppression of factor 2 and more for large x . For RHIC conditions, this suppression is more moderate.

Jet reconstruction performance in the ALICE experiment

Magali Estienne

Under extreme conditions of pressure or temperature, nuclear matter is expected to undergo a phase transition to a deconfined medium, the quark gluon plasma (QGP). Such an atypical state of matter can be defined as a very high density medium in which free partons strongly interact. Experimentally, it is possible to re-create the conditions necessary to its formation by colliding heavy ions at ultra-relativistic velocity. It is part of the experimental program which will be performed at the Large Hadron Collider at CERN. ALICE, one of the four LHC experiments, is dedicated to the study of the QGP properties.

In QCD, jets are defined as cascades of consecutive emissions of partons initiated by partons from an initial hard scattering. This process is followed by the fragmentation of partons into clusters of particles emitted in some opposite collimated directions. In heavy ion collisions, the scene of parton fragmentation is modified from vacuum to a dense medium. Before they fragment, partons are expected to lose energy through collisional energy loss and medium induced gluon radiation, dominant process in a QGP, also called “jet quenching” [171]. Direct consequences of this quenching effect are a modification of the jet shape as well as of the parton fragmentation function associated [168]. The ALICE experiment is planning to measure these modifications thanks to its excellent tracking capabilities down to momenta of order of Λ_{QCD} and up to 100 GeV/c and thanks to its large PID capabilities [172].

Different types of jet finding algorithms (cone and k_T) are under study in ALICE for both p - p and Pb-Pb collisions [173, 174]. Because of the large underlying background in a jet cone of radius $R = \sqrt{\Delta\eta^2 + \Delta\phi^2} = 1$ as well as its large fluctuations in the LHC heavy ion environment, some studies have been performed in order, first, to reduce their contributions in each event, and then to subtract the remaining contributions event-by-event in the process of jet finding. For the former, HIJING simulations have shown that the application of a p_T cut of 1 or 2 GeV/c on charged particles and the reduction of the jet cone size to $R = 0.3$ - 0.4 have to be applied. For the later, three methods of subtraction (statistic, cone and ratio) have been tested and compared [172, 175]. With the iterative jet cone finder based on the UA1 cone algorithm and a full detector simulation, a mean cone energy of $\langle E_i \rangle = 45$ GeV and a resolution of $RMS = \langle E_i \rangle \cdot 40\%$ are obtained using only charged particles in the reconstruction of 100 GeV mono-energetic jets in p - p collisions. The inclusion of neutral particles measured in the ALICE calorimeter [176] allows one to increase $\langle E_i \rangle$ to 70 GeV while improving the resolution up to $\sim 30\%$ (for a jet cone radius of $R = 0.4$). A deeper study of background subtraction in Pb-Pb collisions has shown that the jet reconstruction algorithms bias the mean background energy contained in the jet cone toward values higher than the ones found in the same area outside the jet. This bias leads to an over estimation of the reconstructed energy [177]. We obtain a preliminary $\langle E_i \rangle \sim 75$ GeV and a resolution of 30-35% for the case charged + neutral particles for 100 GeV mono-energetic jets in minimum-bias Pb-Pb collisions.

The full jet spectrum has also been simulated and studied in p - p and Pb-Pb collisions. Due to the steeply falling jet production spectrum and the background and signal fluctuations of the reconstructed jet energy, we observe a smearing of the reconstructed spectra in p - p . This effect is increased in Pb-Pb due to the background fluctuations and its bias in the jet finding process. Jets clearly present advantages for the study of jet quenching at the LHC compared to the studies of leading particles which are trigger- and surface-biased. With its reconstruction capabilities, ALICE can reconstruct jets with enough resolution for Pb-Pb studies even though the observed smearing of the reconstructed spectrum, due to the steeply falling shape, will have to be well controlled both in elementary and Pb-Pb collisions.

Parton fragmentation studies in ATLAS

Jiří Dolejš, for ATLAS collaboration

ATLAS installation and commissioning proceeds towards ability to record first pp collisions during the year 2008. The ATLAS heavy ion working group has examined the detector performance in conditions of Pb-Pb collision at 5.5 TeV per nucleon, based on full simulations of the detector response, and continues in detailed studies [178, 179]. The focus of this talk is on the jet studies. A dedicated jet finder package has been developed and integrated into the standard ATLAS simulation framework. It contains the background subtraction, where the background is estimated over the region outside jet candidate, and iteratively uses the jet cone algorithm on the calibrated energy depositions in individual cells of the ATLAS calorimeters. This package has been tuned on PYTHIA jets embedded into HIJING Pb-Pb collision with a hard-cut of high p_T processes. The simulations show a reconstruction efficiency of the cone algorithm steeply rising from about 50% at $E_T = 60$ GeV up to 90% above 90 GeV. The fake fraction is about 0.04 at $E_T = 60$ GeV and decreases with higher E_T . These results were obtained for events with moderately high $dN/d\eta = 2700$. The axis of the jet can be measured with an angular resolution of about 0.05 (both in ϕ and η) for jets with $E_T = 60$ GeV and improves for higher E_T 's. The energy resolution is about 30% at 60 GeV and is flat over the accessible η range.

The precise position resolution enables a reliable extraction of the jet shapes. The distributions of j_T (particle momentum perpendicular to the jet axis) as well as jet fragmentation functions $D(z)$ have also

been studied using the tracks reconstructed in the ATLAS Inner Detector. The tracking performance was found to be sufficiently good for these distributions. The agreement between the true and reconstructed distributions is good, providing that the tracks from the underlying event are subtracted. The fast- K_T algorithm [174] is studied as one of the possible alternatives. Its efficiency and energy resolution is comparable to that of the cone algorithm and more detailed studies continue.

Extracting fragmentation functions with γ -tagged jet events in Pb-Pb with the CMS experiment

Christof Roland, for the CMS Collaboration

The energy loss of fast partons traversing the strongly interacting matter produced in high-energy nuclear collisions is one of the most interesting observables to probe the nature of the produced medium. The collisional and radiative energy loss of the partons will modify their fragmentation functions depending on the path length of the partons in the medium and the medium density. Pb-Pb collisions at $\sqrt{s_{NN}} = 5.5$ TeV at the LHC will open access to the first detailed measurements of the in-medium modifications of fragmentation functions of parton initiated jets. The fragmentation functions can be studied using the γ -jet channel [180, 181], where the initial transverse energy of the fragmenting parton can be determined from the photon transverse energy, since the photon does not strongly interact with the medium.

The CMS experiment features large acceptance electromagnetic and hadronic calorimeters, and a high precision tracking system well suited to nucleus-nucleus collisions at the LHC [182]. To extract fragmentation functions of parton initiated jets in γ -jet events, isolated photons reconstructed in the electromagnetic calorimeter ($\eta < 2.0$) are selected and correlated with back-to-back reconstructed calorimeter jets ($\eta < 2.0$) in Pb-Pb events. For the selected photon jet pairs, charged particles reconstructed in the tracking system ($\eta < 2.5$) that lie within a 0.5 radius cone in η - ϕ around the reconstructed jet axis are selected. The fragmentation function is constructed correlating the transverse energy of the photons, with the transverse momentum of the tracks in the jet cone.

This study based on a full GEANT simulation of a data set of γ -jet events, containing a γ with $E_T > 70$ GeV, embedded in central Pb-Pb events corresponding to a full year of data taking with an integrated luminosity of 0.5 nb⁻¹. The capability to measure jet fragmentation functions in the γ -jet channel with high statistics is demonstrated. The systematic uncertainties are estimated to be on the order of 30% (10%) for quenched events containing a photon with $E_T > 70$ GeV (100 GeV) mostly due to a bias induced by the finite efficiency of the jet reconstruction. Overall the fragmentation function can be reconstructed with high significance for $0.2 < \xi < 5$ for both, $E_T > 70$ GeV and > 100 GeV. Reconstructed γ -jet events can be used to study the dependence of high- p_T parton fragmentation on the medium and will provide sensitivity to the foreseeable changes in the fragmentation functions relative to vacuum fragmentation. This measurement will allow a quantitative test of proposed mechanisms for medium-induced parton energy loss, testing fundamental properties of high-density QCD.

Issues of fragmentation function and medium effects in single inclusive production and two particle correlations in p - p and A-A collisions

Michael Tannenbaum

Hard-scattering of point-like constituents (partons) in p - p collisions was discovered at the CERN-ISR in 1972 in measurements of inclusive single or pairs of hadrons with large transverse momentum (p_T). Due to the steeply falling power-law p_T spectrum of the hard-scattered partons, the inclusive single particle (e.g. π^0) p_T spectrum from jet fragmentation of a parton with \hat{p}_T is dominated by trigger fragments with large z , $0.7 < z < 0.8$, where $z = p_T/\hat{p}_T$ is the fragmentation variable. It was generally

assumed, following Feynman, Field and Fox [183], as shown by data from the ISR experiments, that the p_{T_a} distribution of away side hadrons from a single particle trigger [with p_{T_i}], corrected for $\langle x_i \rangle$, would be the same as that from a jet-trigger and follow the same fragmentation function measured in e^+e^- or DIS. PHENIX [184, 185] attempted to measure the fragmentation function from the away side x_E $p_{T_a}=p_{T_i}$ distribution of charged particles triggered by a π^0 in $p-p$ collisions and showed by explicit calculation that the x_E distribution is actually quite insensitive to the fragmentation function. A new formula for the distribution of an associated away-side particle with transverse momentum p_{T_a} , which is presumed to be a fragment of an away-jet with \hat{p}_{T_a} , with exponential fragmentation function $D(z) = Be^{-bz}$, triggered by a particle with transverse momentum p_{T_i} , presumably from a trigger-side jet with \hat{p}_{T_i} , invariant cross section, $E d^3\sigma = dp^3 = A \hat{p}_{T_i}^n$, was given [184]: $dP_{p_{T_a}} = dx_E \frac{1}{\hat{p}_{T_i}} \frac{\langle m_i \rangle (n-1)}{(1+x_E \hat{x}_h)^n}$. This formula relates x_E , the ratio of the transverse momenta of the measured particles, to $\hat{x}_h = \hat{p}_{T_a} / \hat{p}_{T_i}$, the ratio of the transverse momenta of the away-side to trigger-side jets, where $\langle m_i \rangle$ is the mean multiplicity of particles in the away jet.

Many analyses of the away-jet p_{T_a} distributions in Au-Au collisions are available; but these tend to describe the effect of the medium with the variable $I_{AA}(x_E)$, the ratio of the x_E distribution in A-A collisions to that in $p-p$ collisions, which typically shows an enhancement at low values of x_E and a suppression at higher values of x_E . Such behavior could be explained as a decrease in \hat{x}_h in A-A collisions due to energy loss of the away jet in the medium. Fits of the above formula to the available data were presented to establish whether: a) the away-jets simply lose energy; b) some of the away-jets lose energy, others punch-through without losing energy; etc.

Measurements of direct single γ production in $p-p$ and Au-Au are also discussed. Fragmentation photons in $p-p$ collisions, if they exist, can be eliminated by isolation cuts, although, at the ISR and at RHIC, measurements of correlations to a direct single γ show very little activity, if any, on the same side. In contrast to triggers with π^0 which are the fragments of jets, away-side correlations for direct single γ 's which are direct participants in 2-2 hard scattering should measure the fragmentation function of the away side jet, modulo k_T smearing. This fragmentation function can be plotted conventionally as a function of x_E or as a function of $\xi = \ln(1/x_E)$ as suggested for jets at the LHC [168].

Fragmentation function from direct-photon associated yields at RHIC

Jan Rak

The interest of heavy-ion community in the fragmentation function (FF) is manifold. The detailed knowledge of the FF modification in heavy-ion collisions is important for constraints of various jet-quenching models [168]. Another region of interest is related to the p_T -broadening predicted to accompany the parton interaction with excited nuclear medium [186]. One way of evaluating the p_T -broadening is to measure the parton intrinsic momentum k_T [184]. It has been shown that in order to extract the value of k_T^2 the knowledge of the average trigger particle momentum fraction $\langle x_i \rangle$ and thus FF is vital (see Eq. (22) in [184]). It has been also demonstrated that, despite a common belief, an analysis of charged particle distributions associated with the high- p_T trigger hadron cannot provide sufficient constraint for the experimental determination of FF. The reason for that lack of sensitivity of high- p_T trigger hadron associated spectra to the shape of FF lies in the fact that when events with different associated transfer momenta are selected the trigger particle momentum fraction $\langle x_i \rangle$ changes even though the trigger particle momentum is fixed. In this talk I have presented an analysis of direct-photon associated spectra which should be free from the above mention problem. There is, however, an effect of k_T -smearing which causes the imbalance between the photon-quark momenta. I have presented a method of unfolding the k_T -smearing which allows to recover the shape of FF from photon associated spectra.

Systematics of complete fragment distributions from nuclear collisions

Thomas A. Trainor

The study of fragmentation in RHIC nuclear collisions addresses the central question how formation and evolution of a QCD medium in heavy ion collisions may affect parton scattering and fragmentation. Whereas previous emphasis has been placed on larger parton energy scales described by perturbative methods, we seek a complete description of all fragmentation processes in nuclear collisions. Fragmentation in elementary collisions provides a reference system. To better determine the energy scale dependence of fragmentation functions from e^+e^- , the FFs were replotted on a normalized rapidity variable, resulting in a compact form precisely represented by the beta distribution, its two parameters varying slowly and simply with parton energy scale Q [187]. The new parameterization extrapolates fragmentation functions to small Q and describes the fragmentation process down to small transverse momentum p_T in nuclear collisions at RHIC. Analysis of transverse momentum p_T spectra to $p_T = 6$ GeV/c for 10 multiplicity classes of p - p collisions at $\sqrt{s} = 200$ GeV revealed that the spectrum shape can be decomposed into a part with amplitude proportional to multiplicity and described by a Lévy distribution on transverse mass m_T , and a part with amplitude proportional to multiplicity squared and described by a Gaussian distribution on transverse rapidity y_T [188]. The functional forms of the two parts are nearly independent of event multiplicity, and they can be identified with the soft and hard components of a two-component model of p - p collisions: longitudinal participant nucleon fragmentation and transverse scattered parton fragmentation. Interpretation of the hard component as parton fragmentation is supported by correlation studies [189, 190].

A similar two-component analysis of spectra to $p_T = 12$ GeV/c for identified pions and protons from 200 GeV Au-Au collisions applied to Au-Au centrality dependence revealed that the soft-component reference is again a Lévy distribution on transverse mass m_T , and the hard-component reference is again a Gaussian on y_T , but with added exponential (p_T power-law) tail [191]. Deviations of data from the two-component reference are described by hard-component ratio r_{AA} which generalizes nuclear modification factor R_{AA} . Centrality evolution of pion and proton spectra is dominated by changes in parton fragmentation. The structure of r_{AA} suggests that parton energy loss produces a negative boost Δy_T of a large fraction (but not all) of the minimum-bias fragment distribution (hard component), and that lower-energy partons suffer relatively less energy loss, possibly due to color screening.

A recent analysis of azimuth correlations based on methods in [192, 193] has isolated azimuth quadrupole components from published $v_2(p_T)$ data (called elliptic flow) as spectra on transverse rapidity y_T for identified pions, kaons and lambdas from minimum-bias Au-Au collisions at 200 GeV. The form of the spectra reveals that the three hadron species are emitted from a common source with boost $\Delta y_{t0} = 0.6$. The quadrupole spectra have a Lévy form on m_T similar to the soft components of single-particle spectra, but with significantly reduced (≈ 0.7) slope parameter T . Comparison of quadrupole spectra with single-particle spectra suggests that the quadrupole component comprises a small fraction ($< 5\%$) of the total hadron yield, contradicting the hydrodynamic picture of a thermalized, flowing bulk medium. The form of $v_2(p_T)$ is, within a constant factor, the product of $p_t^0(p_T$ in the boost frame) times the ratio of the quadrupole spectrum to the single-particle spectrum. The form of the spectrum ratio then implies that above 0.5 GeV/c, $v_2(p_T)$ is dominated by the hard component of the single-particle spectrum (minijets). It is thus unlikely that so-called *constituent-quark scaling* attributed to v_2 is relevant to soft hadron production mechanisms (e.g., chemical freezeout). These results suggest instead that “elliptic flow” may result from fragmentation of gluonic quadrupole radiation, not hydrodynamic evolution.

A List of participants

PARTICIPANTS:

S. Albino (*Universität Hamburg*),
F. Anulli (*INFN, Roma*),
E. Aschenauer (*DESY, Zeuthen*),
F. Arleo (*LAPTH, Annecy*),
D. Besson (*Kansas University*),
N. Borghini (*Heidelberg Univ.*),
W. Brooks (*Santa Maria Univ., Valparaiso*),
B. Buschbeck (*OA Wien, Austria*),
M. Cacciari (*LPTHE, Université Paris 6*),
E. Christova (*Inst. Nuc. Res. & Nuc Energy, Sofia*),
G. Corcella (*Centro Fermi Roma, SNS/INFN Pisa*),
D. d'Enterria (*CERN, Geneva*),
J. Dolejší (*Charles University, Prague*),
S. Domdey (*Universität Heidelberg*),
M. Estienne (*IPHC, Strasbourg*),
O. Fochler (*University Frankfurt*),
T. Gousset (*Subatech, Nantes*),
K. Hamacher (*Bergische Univ. Wuppertal*),
M. Heinz (*Yale University*),
K. Hicks (*Ohio University*),
D. Kettler (*Univ. of Washington, Seattle*),
S. Kumano (*KEK, Japan*),
B. Machet (*LPTHE, Univ. Paris 6*),
G. Milhano (*CENTRA-IST, Lisbon*),
S.-O. Moch (*DESY, Zeuthen*),
V. Muccifora (*INFN-LNF, Frascati*),
S. Pacetti (*Centro Fermi Roma, LNF Frascati*),
R. Pérez-Ramos (*Universität Hamburg*),
H.-J. Pirner (*Univ. Heidelberg*),
A. Pronko (*Fermilab, Chicago*),
J. Putschke (*Yale University*),
M. Radici (*INFN - Sezione di Pavia*),
J. Rak (*University of Jyväskylä*),
C. Roland (*MIT, Cambridge, MA*),
G. Rudolph (*Universität Innsbruck*),
Z. Rúriková (*DESY, Hamburg*),
C.A. Salgado (*Univ. Santiago de Compostela*),
S. Sapeta (*Jagellonian U. Cracow*),
D.H. Saxon (*University of Glasgow*),
R. Seidl (*RIKEN BNL, Upton, NY*),
R. Seuster (*Victoria Univ., BC*),
M. Stratmann (*RIKEN, Japan*),
M.J. Tannenbaum (*BNL, Upton, NY*),
M. Tasevsky (*Charles Univ., Prague*),
T. Trainor (*University of Washington, Seattle*),
D. Traynor (*QMUL, London*),
M. Werlen (*LAPTH, Annecy*),
C. Zhou (*McGill Univ., Montreal*)

B Programme

Monday, 25 February 2008

09:00 <i>Welcome and introduction</i> (10')	F. Arleo, D. d'Enterria
09:10 <i>Sources of uncertainty in quark & gluon FFs into hadrons & photons</i> (50')	M. Werlen
10:00 <i>OPAL results on quark, gluon fragmentation</i> (30')	M. Tasevsky
11:00 <i>Colour flux studies in quark & gluon Fragmentation in DELPHI</i> (30')	B. Buschbeck
11:30 <i>ALEPH results on quark and gluon fragmentation</i> (30')	G. Rudolph
12:00 <i>Colour coherence and a comparison of gluon and quark fragmentation</i> (30')	K. Hamacher
14:30 <i>AKK fragmentation functions: latest developments.</i> (40')	S. Albino
15:10 <i>DSS Fragmentation functions</i> (40')	M. Stratmann
16:20 <i>HKNS fragmentation functions and proposal for exotic-hadron search</i> (40')	S. Kumano
17:00 <i>Low-Q^2 fragmentation functions in e^+e^- collisions</i> (40')	D. Kettler

Tuesday 26 February 2008

08:30	<i>Evidence for power corrections in heavy quark fragmentation in e^+e^-</i> (30')	M. Cacciari
09:00	<i>Heavy Quark Fragmentation at 10.6 GeV</i> (30')	R. Seuster
09:30	<i>Initial state radiation at BaBar</i> (30')	S. Pacetti
10:00	<i>Heavy quark fragmentation with an effective coupling constant</i> (30')	G. Corcella
11:00	<i>Charm fragmentation at HERA</i> (30')	Z. Rúriková
11:30	<i>Inclusive light and charmed particle production in BaBar</i> (30')	S. Pacetti
14:00	<i>Discussion: interfacing fragmentation functions ?</i> (30')	F. Arleo, D. d'Enterria
14:30	<i>Experimental constraints on strange quark fragmentation functions</i> (30')	M. Heinz
15:00	<i>Latest developments on parton to kaon fragmentation</i> (30')	E. Christova
15:30	<i>Inclusive $K_s^0 K_s^0$ resonance production in ep collisions at HERA</i> (30')	Changyi Zhou
16:30	<i>Theory of (extended) dihadron fragmentation functions</i> (30')	M. Radici
17:00	<i>Time-like splitting functions at NNLO in QCD</i> (30')	Sven-Olaf Moch

Wednesday 27 February 2008

09:00	<i>Recent developments in parton fragmentation at MLLA and beyond</i> (45')	R. Pérez Ramos
09:45	<i>Studies of Jet Fragmentation at CDF</i> (30')	A. Pronko
10:45	<i>Particle production at ZEUS</i> (40')	D.H. Saxon
11:25	<i>Particle production and fragmentation studies in H1</i> (30')	D. Traynor
11:55	<i>Review on parton fragmentation studies in BaBar</i> (30')	F. Anulli
14:30	<i>Fragmentation Studies at CLEO</i> (30')	D. Besson
15:00	<i>Light quark and spin dependent fragmentation at Belle</i> (30')	R. Seidl
16:00	<i>HERMES: News on fragmentation from nucleons to nuclei</i> (30')	E. Aschenauer
16:30	<i>p_T broadening in nuclear collisions</i> (30')	H.-J. Pirner

Thursday 28 February 2008

09:00	<i>Medium modified Fragmentation Functions: an overview</i> (45')	C. Salgado
09:45	<i>Parton fragmentation studies in ATLAS</i> (30')	J. Dolejší
10:45	<i>Fragmentation functions with γ-jet in Pb-Pb at 5.5 TeV at CMS</i> (30')	C. Roland
11:15	<i>Reconstructing jets in Pb-Pb collisions in ALICE</i> (30')	M. Estienne
11:45	<i>Jet fragmentation studies in the ALICE experiment</i> (30')	J. Putschke
14:30	<i>Systematics of complete fragment distributions in nuclear collisions</i> (30')	T. Trainor
15:00	<i>Hadronic composition as a characteristics of jet quenching at LHC</i> (30')	S. Sapeta
16:00	<i>Fragmentation function from direct-photon associated yields at RHIC</i> (30')	J. Rak
16:30	<i>FFs and medium effects in single & 2-particle production in p-p and A-A</i> (30')	M.J. Tannenbaum

Friday 29 February 2008

09:30	<i>Attenuation of leading hadrons in DIS on nuclei: an overview</i> (45')	V. Muccifora
10:15	<i>Hadron Formation in Semi-Inclusive DIS on Nuclei at HERMES</i> (30')	G. Elbakyan
11:15	<i>$\pi;K$ hadronization from Electroproduction DIS on Nuclei at CLAS</i> (30')	K. Hicks
11:45	<i>Parton propagation & hadron formation: current status, future prospects</i> (30')	W. Brooks
12:15	<i>QCD evolution of jets in the quark-gluon plasma</i> (30')	S. Domdey

References

- [1] W. Giele *et al.*, “The QCD/SM working group: Summary report” (Section 1.2), hep-ph/0204316
- [2] M. R. Whalley, D. Bourilkov, and R. C. Group; arXiv:hep-ph/0508110.
- [3] H. Plochow-Besch. *Comput. Phys. Commun.* (1993) 396.
- [4] S. Albino *et al.*, *Nucl. Phys.* **B725** (2005) 181; *Nucl. Phys.* **B734** (2006) 50
- [5] S. Albino, B. A. Kniehl and G. Kramer, arXiv:0803.2768 [hep-ph].
- [6] J. Binnewies, B. A. Kniehl and G. Kramer, *Phys. Rev.* **D52** (1995) 4947.
- [7] L. Bourhis, M. Fontannaz and J. P. Guillet, *Eur. Phys. J. C2* (1998) 529
- [8] L. Bourhis, M. Fontannaz, J. P. Guillet and M. Werlen, *Eur. Phys. J. C19* (2001) 89
- [9] P. Chiappetta, M. Greco, J. P. Guillet, S. Rolli and M. Werlen, *Nucl. Phys.* **B412** (1994) 3.
- [10] D. de Florian, R. Sassot and M. Stratmann, *Phys. Rev.* **D75** (2007) 114010;
- [11] D. de Florian, R. Sassot, and M. Stratmann, *Phys. Rev.* **D76** (2007) 074033.
- [12] D. de Florian, M. Stratmann, W. Vogelsang, *Phys. Rev.* **D57** (1998) 5811.
- [13] M. Glück, E. Reya, A. Vogt, *Phys. Rev.* **D48** (1993) 116.
- [14] M. Hirai, S. Kumano, T. H. Nagai and K. Sudoh, *Phys. Rev.* **D75** (2007) 094009.
Code available at <http://research.kek.jp/people/kumanos/ffs.html>
- [15] B. A. Kniehl, G. Kramer and B. Potter, *Nucl. Phys.* **B582** (2000) 514
- [16] S. Kretzer, *Phys. Rev.* **D62** (2000) 054001
- [17] See <http://www.pv.infn.it/~radici/FFdatabase/>
- [18] S. Albino, B. A. Kniehl, G. Kramer and W. Ochs, *Phys. Rev.* **D73** (2006) 054020;
- [19] M. Cacciari and S. Catani, *Nucl. Phys.* **B617** (2001) 253.
- [20] I. Arsene *et al.* [BRAHMS Collab.], *Phys. Rev. Lett.* **98** (2007) 252001.
- [21] S. S. Adler *et al.* [PHENIX Collab.], *Phys. Rev. Lett.* **91** (2003) 241803.
- [22] J. Adams *et al.* [STAR Collab.], *Phys. Lett.* **B637** (2006) 161.
- [23] J. Adams *et al.* [STAR Collab.], *Phys. Rev. Lett.* **97** (2006) 152302.
- [24] B. I. Abelev *et al.* [STAR Collab.], *Phys. Rev.* **C75** (2007) 064901.
- [25] D. E. Acosta *et al.* [CDF Collab.], *Phys. Rev.* **D72** (2005) 052001.
- [26] M. Stratmann and W. Vogelsang, *Phys. Rev.* **D64** (2001) 114007.
- [27] See, e.g., D. Stump *et al.*, *Phys. Rev.* **D65** (2002) 014012.
- [28] M. Hirai, S. Kumano, M. Oka, and K. Sudoh, *Phys. Rev.* **D77** (2008) 017504.

- [29] S. Kumano and V. R. Pandharipande, Phys. Rev. D**38** (1988) 146.
- [30] F. E. Close, N. Isgur, and S. Kumano, Nucl. Phys. B**389** (1993) 513.
- [31] See http://www.lapp.in2p3.fr/lapth/PHOX_FAMILY/main.html
- [32] F. Arleo *et al.*, arXiv:hep-ph/0311131.
- [33] P. Aurenche *et al.*, Eur. Phys. J. C**13** (2000) 347
- [34] T. Binoth, J. P. Guillet, E. Pilon and M. Werlen, Eur. Phys. J. C**24** (2002) 245.
- [35] Z. Belghobsi, M. Fontannaz, J.Ph. Guillet, G. Heinrich, E. Pilon and M. Werlen, in preparation
- [36] P. Aurenche, M. Fontannaz, J. P. Guillet, E. Pilon and M. Werlen, Phys. Rev. D**73** (2006) 094007
- [37] V.N. Gribov and L.N. Lipatov, Sov. J. Nucl. Phys. **15** (1972) 438, *ibid.* 675.
- [38] M. Stratmann and W. Vogelsang, Nucl. Phys. B**496** (1997) 41
- [39] A. Mitov, S. Moch and A. Vogt, Phys. Lett. B**638** (2006) 61
- [40] S. Moch and A. Vogt, Phys. Lett. B**659** (2008) 290
- [41] S. Moch, J.A.M. Vermaseren and A. Vogt, Nucl. Phys. B**688** (2004) 101
- [42] A. Vogt, S. Moch and J.A.M. Vermaseren, Nucl. Phys. B**691** (2004) 129
- [43] Y.L. Dokshitzer, G. Marchesini and G.P. Salam, Phys. Lett. B**634** (2006) 504
- [44] A. Bianconi *et al.*, Phys. Rev. D**62** (2000) 034008.
- [45] A. Bacchetta and M. Radici, Phys. Rev. D**69** (2004) 074026.
- [46] A. Bacchetta and M. Radici, Phys. Rev. D**67** (2003) 094002.
- [47] M. Radici, R. Jakob, and A. Bianconi, Phys. Rev. D**65** (2002) 074031.
- [48] V. Barone and P.G. Ratcliffe, *Transverse Spin Physics* (World Scientific, 2003).
- [49] A. Airapetian *et al.* [HERMES collaboration], arXiv:0803.2367 [hep-ex]
- [50] A. Martin [COMPASS Collab.], hep-ex/0702002.
- [51] D. Boer, R. Jakob, and M. Radici, Phys. Rev. D**67** (2003) 094003.
- [52] K. Hasuko *et al.* [BELLE Collab.], AIP Conf. Proc. **675** (2003) 454.
- [53] A. Bacchetta and M. Radici, Phys. Rev. D**74** (2006) 114007 and refs. therein.
- [54] D. de Florian and L. Vanni, Phys. Lett. B**578** (2004) 139.
- [55] F.A. Ceccopieri, M. Radici, and A. Bacchetta, Phys. Lett. B**650** (2007) 81.
- [56] R. Pérez-Ramos and B. Machet, JHEP **04** (2006) 043
- [57] R. Pérez-Ramos, JHEP **06** (2006) 019, and references therein

- [58] Yu.L. Dokshitzer, V.A. Khoze, A.H. Mueller and S.I. Troyan: “Basics of Perturbative QCD”, (Editions Frontières, Gif-sur-Yvette, 1991)
- [59] F. Cuypers and K. Tesima, *Z. Phys.* **C54** (1992) 87
- [60] Yu.L. Dokshitzer, *Phys. Lett.* **B305** (1993) 295.
- [61] F. Arleo, R. Pérez-Ramos and B. Machet, *Phys. Rev. Lett.* **100** (2008) 052002; R. Pérez-Ramos, F. Arleo and B. Machet, arXiv:0712.2212 [hep-ph].
- [62] Ya.I. Azimov, Yu.L. Dokshitzer, V.A. Khoze and S.I. Troian, *Z. Phys* **C27** (1985) 65; Yu.L. Dokshitzer, V.A. Khoze and S.I. Troian, *J. Phys.* **G 17** (1991) 1585
- [63] C.P. Fong and B.R. Webber, *Nucl. Phys.* **B355** (1991) 54; *Phys. Lett.* **B241** (1990) 255
- [64] J. Abdallah *et al.* [DELPHI Collab.], *Phys. Lett.* **B605** (2005) 37.
- [65] M. Siebel, PhD thesis, University of Wuppertal (WUB-DIS 2003-11) and refs. therein, <http://elpub.bib.uni-wuppertal.de/edocs/dokumente/fbc/physik/diss2003/siebel>
- [66] J. Abdallah *et al.* [DELPHI Collab.], *Eur. Phys. J.* **C44** (2005) 311; K. Hamacher, O. Klapp, P. Langefeld and M. Siebel [DELPHI Collab.], CERN-OPEN-2000-134; P. Abreu *et al.* [DELPHI Collab.], *Phys. Lett.* **B449** (1999) 383.
- [67] J. Abdallah *et al.* [DELPHI Collab.], *Eur. Phys. J.* **C29** (2003) 285.
- [68] P. Abreu *et al.* [DELPHI Collab.], *Eur. Phys. J.* **C17** (2000) 207.
- [69] K. Hamacher, *Acta Phys. Polon.* **B36** (2005) 433.
- [70] D. Busculic *et al.* [ALEPH Collab.], *Phys. Lett.* **B357** (1995) 487
- [71] A. Heister *et al.* [ALEPH Collab.], *Eur. Phys. J.* **C35** (2004) 457
- [72] R. Barate *et al.* [ALEPH Collab.], *Physics Reports* **294** (1998) 1
- [73] R. Barate *et al.* [ALEPH Collab.], *Eur. Phys. J.* **C16** (2000) 613
- [74] S. Schael *et al.* [ALEPH Collab.], *Eur. Phys. J.* **C48** (2006) 685
- [75] T. Sjöstrand and P. Skands, *Eur. Phys. J.* **C39** (2005) 129
- [76] R. Barate *et al.* [ALEPH Collab.], *Eur. Phys. J.* **C16** (2000) 597
- [77] A. Heister *et al.* [ALEPH Collab.], *Phys. Lett.* **B512** (2001) 30
- [78] D. Busculic *et al.* [ALEPH Collab.], *Z. Phys.* **C69** (1996) 393
- [79] S. Schael *et al.* [ALEPH Collab.], *Phys. Lett.* **B606** (2004) 265
- [80] D. Busculic *et al.* [ALEPH Collab.], *Phys. Lett.* **B384** (1996) 353
- [81] R. Barate *et al.* [ALEPH Collab.], *Eur. Phys. J.* **C17** (2000) 1
- [82] G. Abbiendi *et al.* [OPAL Collab.], *Eur. Phys. J.* **C37** (2004) 25
- [83] J. Abdallah *et al.* [DELPHI Collab.], *Phys. Lett.* **B643** (2006) 147

- [84] J. Abdallah *et al.* [DELPHI Collab.], Eur. Phys. J. C51 (2007) 249.
- [85] J. Abdallah *et al.* [DELPHI Collab.], Eur. Phys. J. C44 (2005) 161.
- [86] J. Abdallah *et al.* [DELPHI Collab.], Phys. Lett. B**643** (2006) 147.
- [87] B. Aubert *et al.* [BABAR Collab.], Nucl. Instrum. Meth. A, **479** (2002) 1.
- [88] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. D**75** (2007) 012003.
- [89] A. Bornheim *et al.* [CLEO Collab.], Phys. Rev. D**63** (2001) 112003.
- [90] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. Lett. **97** (2006) 112002.
- [91] A. Abashian *et al.* [BELLE Collab.], Nucl. Instrum. Meth. A, **479** (2002) 117.
- [92] S. Kurokawa, E. Kikutani, Nucl. Instrum. Meth. A**499** (2003) 1, and other papers in this volume.
- [93] J. C. Collins, Nucl. Phys. B**396** (1993) 161.
- [94] R. Seidl *et al.* [BELLE Collab.], Phys. Rev. Lett. **96** (2006) 232002.
- [95] A. Airapetian *et al.* [HERMES Collab.], Phys. Rev. Lett. **94** (2005) 012002.
- [96] E.S Ageev *et al.* [COMPASS Collab.], Nucl. Phys. B**765** 31 (2007)
- [97] M. Anselmino *et al.*, Phys. Rev. D**75** (2007) 054032
- [98] Z. Metreveli *et al.* [CLEO Collab.], Phys. Rev. D**66** (2002) 052002
- [99] D. M. Asner *et al.* [CLEO Collab.], Phys. Rev. D**75** (2007) 012009
- [100] D. Besson *et al.* [CLEO Collab.], Phys. Rev. D**74** (2006) 012003
- [101] M. S. Alam *et al.* [CLEO Collab.], Phys. Rev. D**56** (1997) 17
- [102] G. T. Bodwin, E. Braaten, D. Kang and J. Lee, Phys. Rev. D**76** (2007) 054001
- [103] F.D. Aaron *et al.* [H1 Collab.], Phys. Lett. B**654** (2007) 148
- [104] Sakar Osman [H1 Collab.], Proceeds DIS 2007, Munich.
- [105] Reference material can be found in <http://www-zeus.desy.de>.
- [106] T. Acosta *et al.* [CDF Collab.], Phys. Rev. D**68** (2003) 012003
- [107] T. Acosta *et al.* [CDF Collab.], Phys. Rev. Lett. **94** (2005) 171802
- [108] A. Capella *et al.*, Phys. Rev. D**61** (2000) 074009.
- [109] T. Aaltonen *et al.* [CDF Collab.], submitted to Phys. Rev. D; arXiv:0802.3182
- [110] M. Cacciari, P. Nason and C. Oleari, JHEP **0604** (2006) 006
- [111] M. Artuso *et al.* [CLEO Collab.], Phys. Rev. D**70** (2004) 112001
- [112] R. Seuster *et al.* [BELLE Collab.], Phys. Rev. D**73** (2006) 032002

- [113] M. Cacciari, P. Nason and C. Oleari, JHEP **0510** (2005) 034
- [114] D. Shirkov, Nucl. Phys. Proc. Suppl. **152** (2006) 51.
- [115] B. Mele and P. Nason, Nucl. Phys. **B361** (1991) 626.
- [116] U. Aglietti, G. Corcella and G. Ferrera, Nucl. Phys. **B775** (2007) 162.
- [117] G. Corcella and G. Ferrera, JHEP **0712** (2007) 029.
- [118] E. Christova, E. Leader, Eur. Phys. J. **C51** (2007) 825
- [119] A. Hillenbrand, DESY-THESIS-2005-035.
- [120] T. Mannel, arXiv:hep-ph/9611411.
- [121] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D74** (2006) 032007
- [122] T. Barnes, F. E. Close and H. J. Lipkin, Phys. Rev. **D68** (2003) 054006; E. Klempt and A. Zaitsev, Phys. Rept. **454** (2007) 1
- [123] B. Aubert [BABAR Collab.], Phys. Rev. Lett. **97** (2006) 222001
- [124] K. Abe *et al.* [BELLE Collab.], arXiv:hep-ex/0608031.
- [125] W.-M. Yao *et al.*, J. Phys. G **33** (2006) 1, and <http://pdf.lbl.gov>.
- [126] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. Lett. **97** (2006) 232001.
- [127] R. M. Woloshyn, Nucl. Phys. Proc. Suppl. **93** (2001) 38.
- [128] L. Burakovsky, T. Goldman and L. P. Horwitz, Phys. Rev. **D56** (1997) 7124; M. J. Savage, Phys. Lett. **B359** (1995) 189; J. L. Rosner, Phys. Rev. **D52** (1995) 6461; R. Roncaglia *et al.*, Phys. Rev. **D52** (1995) 1722 D. B. Lichtenberg *et al.*, Phys. Rev. **D53** (1996) 6678; A. Zalewska and K. Zalewski, hep-ph/9608240; L. Ya. Glozman and D. O. Riska, Nucl. Phys. A **603** (1996) 326; E. Jenkins, Phys. Rev. **D54** (1996) 4515.
- [129] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. Lett. **98** (2007) 012001
- [130] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D77** (2008) 012002
- [131] M. Davier *et al.*, Eur. Phys. J. **C27** (2003) 497.
- [132] H. Burkhardt and B. Pietrzyk, Phys. Lett. **B513** (2001) 46.
- [133] B. Aubert *et al.* [BABAR], Phys. Rev. **D70** (2004) 072004; Phys. Rev. **D71** (2005) 052001
- [134] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D73** (2006) 052003; W. F. Wang [BABAR Collab.], Nucl. Phys. Proc. Suppl. **169** (2007) 282.
- [135] B. Aubert *et al.* [BABAR], Phys. Rev. **D74** (2006) 091103; Phys. Rev. **D76** (2007) 012008
- [136] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D76** (2007) 092005
- [137] B. Aubert *et al.* [BABAR Collab.], arXiv:0710.4451 [hep-ex].

- [138] P. L. Frabetti *et al.*, Phys. Lett. **B514** (2001) 240; Phys. Lett. **B578** (2004) 290; R. Baldini *et al.*, “FENICE” workshop, Frascati (1998); B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D73** (2006) 052003; P. Lebrun, FERMILAB-CONF-97-387-E, “Hadron 97”, Upton, NY, 25-30 Aug 1997.
- [139] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D73** (2006) 012005
- [140] E. Tomasi-Gustafsson, E. A. Kuraev, S. Bakmaev and S. Pacetti, Phys. Lett. **B659** (2008) 197
- [141] B. Aubert *et al.* [BABAR Collab.], Phys. Rev. **D76** (2007) 092006
- [142] Z. Růriková [H1 Collab.], Procs. DIS2005, <http://www-h1.desy.de/publications/htmlsplit/H1prelim-05-074.long.html>
- [143] B.W. Harris and J. Smith, Nucl. Phys. **B452** (1995) 109; Phys. Lett. **B353** (1995) 535
- [144] S. Schael *et al.* [ALEPH Collab.], Phys. Lett. **B606** (2005) 265
- [145] C. J. Morningstar and M. Peardon, Phys. R173 Rev. **D60** (1999) 034509
- [146] C. Michael and M. Teper, Nucl. Phys. **B314** (1989) 347.
- [147] E. Klempt and A. Zaitsev, Phys. Rept. 454 (2007) 1;
- [148] M. Albaladejo, J. A. Oller, hep-ph/0801.4929.
- [149] M. Althoff *et al.*, Phys. Lett. **B121** (1983) 216 .
- [150] D. Faiman, H.J. Lipkin and H.R. Rubinstein, Phys. Lett. **B59** (1975) 269.
- [151] G. Gatof and C.Y. Wong, Phys. Rev. **D46**, (1992) 997
- [152] A. Airapetian *et al.* [HERMES Collab.], Eur. Phys. J. **C20** (2001) 479; Phys. Lett. **B577** (2003) 37; Nucl. Phys. **B780** (2007) 1
- [153] A. Airapetian *et al.* [HERMES Collab.], Phys. Rev. Lett. **96** (2006) 16230.
- [154] Y. Van Haarlem, A. Jgoun, P. Di Nezza, arXiv:0704.3712
- [155] X.N. Wang and X. Guo, Nucl. Phys. A **696** (2001) 788
- [156] F. Arleo, JHEP **0211** (2002) 044, Eur. Phys. J. **C30** (2003) 213.
- [157] B.Z. Kopeliovich, J. Nemchik, E. Predazzi and A. Hayashigaki, Nucl. Phys. A **740** (2003) 211
- [158] K. Gallmeister *et al.* nucl-th/0701064; nucl-th/07122200.
- [159] A. Accardi, arXiv:0706.3227 [nucl-th].
- [160] J. Casalderrey Solana, C. A. Salgado, arXiv:0712.3443
- [161] C. A. Salgado and U. A. Wiedemann, PRD**68** (2003) 014008
- [162] N. Armesto, L. Cunqueiro, C. A. Salgado and W. C. Xiang, arXiv:0710.3073 [hep-ph].
- [163] U.A. Wiedeman Nucl. Phys. **B588** (2000) 303
- [164] A. D. Polosa and C. A. Salgado, Phys. Rev. **C75** (2007) 041901
- [165] A. Adare *et al.* [PHENIX Collab.], arXiv:0801.4545 [nucl-ex].

- [166] Yu. L. Dokshitzer, V. A. Khoze and S. I. Troian, Adv. Ser. Direct. High Energy Phys. (1988) 241.
- [167] Y. I. Azimov, Yu. L. Dokshitzer, V. A. Khoze and S. I. Troian, Z. Phys. C**31** (1986) 213.
- [168] N. Borghini and U. A. Wiedemann, hep-ph/0506218.
- [169] S. Sapeta and U. A. Wiedemann, arXiv:0707.3494 [hep-ph].
- [170] S. Domdey, G. Ingelman, H.J. Pirner, J. Rathsman, K. Zapp, J. Stachel, arXiv: hep-ph/08023282
- [171] M. Gyulassy *et al.*, arXiv:nucl-th/0302077.
- [172] F. Carminati *et al.* [ALICE Collab.], J. Phys. G **32** (2006) 1295.
- [173] G. Arnison *et al.* [UA1 Collab.], Phys. Lett. B**132** (1983) 214.
- [174] G. Salam and M. Cacciari, Phys. Lett. B**641** (2006) 57.
- [175] S.L. Blyth *et al.*, arXiv:nucl-ex/0609023.
- [176] P. Cortese *et al.* [ALICE Collab.], ECAL Addendum, CERN-LHCC-2006-014 (2006).
- [177] A. Abrahantes, R. Diaz, M. Lopez Noriega, E. Lopez and A. Morsch, ALICE-INT-2008-005.
- [178] Heavy Ion Physics with the ATLAS Detector (LoI) [ATLAS Collab.], CERN/LHCC/2004-009, LHCC I-01322 March 2004
- [179] N. Grau [ATLAS Collab.], proceeds. Quark Matter 2008, J. Phys. G to be submitted.
- [180] X. N. Wang, Z. Huang and I. Sarcevic, Phys. Rev. Lett. **77** (1996) 231.
- [181] F. Arleo, P. Aurenche, Z. Belghobsi, J.-P. Guillet, JHEP **0411** (2004) 009; F. Arleo, JHEP **0609** (2006) 015.
- [182] D. d'Enterria (ed.) *et al.* [CMS Collab.], J. Phys. G **34** (2007) 2307.
- [183] R. P. Feynman, R. D. Field, and G. C. Fox, Nucl. Phys. B**128** (1977) 1.
- [184] S. S. Adler, *et al.* [PHENIX Collab.], Phys. Rev. D**74** (2006) 072002.
- [185] Latest results for Au-Au π^0 suppression in A. Adare *et al.*, [PHENIX Collab.], arXiv:0801.1665, arXiv:0801.4020. Latest results for p-p direct- γ production in S.S. Adler, *et al.* [PHENIX Collab.], Phys. Rev. Lett. **98** (2007) 012002.
- [186] J. Qiu, I. Vitev, Phys. Lett., 2003, B**570** (2003) 161
- [187] T. A. Trainor and D. T. Kettler, Phys. Rev. D**74** (2006) 034012.
- [188] J. Adams *et al.* [STAR Collab.], Phys. Rev. D**74** (2006) 032006.
- [189] T. A. Trainor [STAR Collab.], Proceeds ISMD 2005, AIP Conf. Proc. **828** (2006) 238.
- [190] R. J. Porter and T. A. Trainor, Proceeds of Science (CFRNC2006) 004, (2006).
- [191] T. A. Trainor, arXiv:0710.4504 [hep-ph].
- [192] T. A. Trainor and D. T. Kettler, arXiv:0704.1674 [hep-ph].
- [193] T. A. Trainor, arXiv:0708.0792 [hep-ph].