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# Impact of nitrogen incorporation on pseudomorphic site-controlled quantum dots grown by metalorganic vapor phase epitaxy

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We report on some surprising optical properties of diluted nitride  $\text{InGaAs}_{1-\varepsilon}\text{N}_\varepsilon/\text{GaAs}$  ( $\varepsilon \ll 1$ ) pyramidal site-controlled quantum dots, grown by metalorganic vapor phase epitaxy on patterned GaAs (111)B substrates. Microphotoluminescence characterizations showed antibinding exciton/biexciton behavior, a spread of exciton lifetimes in an otherwise very uniform sample, with unexpected long neutral exciton lifetimes (up to 7 ns) and a nearly zero fine structure splitting on a majority of dots. © 2010 American Institute of Physics. [doi:10.1063/1.3481675]

In the past decades researches in fundamental physics, electrostatics, and the more recently arisen technologies of quantum information and quantum computing have regarded semiconductor quantum dots (QDs) as ideal candidate for both theoretical studies and practical applications. The interest QDs have attracted is amplified by the possibility of controlling and manipulating their electro-optical properties. Emission energies can be tuned by modifying the dot dimensions and/or the dot-barriers band discontinuity.<sup>1</sup> A fine tuning of the exciton/biexciton binding energy can be achieved through a size scalability mechanism or via the application of electro-magnetic fields.<sup>2</sup> Strain induced, electric and magnetic fields can be used to reduce the built-in asymmetries which produce a fine structure splitting (FSS) in the cascade biexciton/exciton pair, critical for a reliable source of entangled photons.<sup>3</sup>

More recent investigations on self-assembled  $\text{In}(\text{Ga})\text{As}/\text{GaAs}$  QDs have suggested tuning mechanisms of the dot properties through strain engineering.<sup>4</sup> Employing strain-controlled schemes in these systems would overcome the main limiting factor which still inhibits a further redshift in the emission wavelength. The compressive strain effect, originating from the high difference between GaAs substrate and  $\text{In}(\text{Ga})\text{As}$  dot layer lattice constant, limits in fact an increase in the thickness and/or indium concentration. Moreover, the asymmetric strain distribution is one of the causes of the lack of degeneracy of the cascade biexciton/exciton photon pair, hampering the generation of polarized entangled photons.

It has been demonstrated that a small quantity of nitrogen into  $(\text{In})\text{GaAs}$  layers can partially compensate the compressive strain,<sup>5</sup> reducing relaxation processes and the consequent degradation of the emission properties. Incorporating nitrogen remains though a real issue, both for molecular beam epitaxy (MBE)<sup>6</sup> and metalorganic vapor phase epitaxy (MOVPE) techniques; the actual low incorporation efficiency requires particular growth conditions, based mainly on low temperatures and low group V source partial pressures,<sup>7</sup> which lead to a generally poorer crystallographic quality of the grown structures.

We present in this letter unforeseen optical properties of pseudomorphic pyramidal site controlled  $\text{InGaAs}_{1-\varepsilon}\text{N}_\varepsilon/\text{GaAs}$  QDs grown by MOVPE on (111)B GaAs substrates.

These systems have shown both high uniformity and spectral purity on par with self assembled samples. Recently a simple wavelength tuning mechanism was achieved by changing the dot indium concentration.<sup>8,9</sup> The strain in the  $\text{InGaAs}$  layers nevertheless limited the QD emission wavelength to  $\sim 880$  nm (for a nominal Indium concentration of 45%) and delivered poorer spectral quality in the samples with higher In content.

To limit these deteriorating effects, we started growing and investigating diluted nitride QD systems. We report here on the properties of a specific sample structure (fully reproducible in its properties) in the limit of ultra nitrogen dilution. We anticipate that, although no substantial redshift was observed in the emission spectra from the specimen here analyzed (for other  $\text{InGaAsN}/\text{GaAs}$  systems, grown under different conditions, we have a more significant increase in the emission wavelength), the small concentration of nitrogen incorporated in the dot changed other optical features of the system; an unexpected antibinding biexciton/exciton behavior was observed and a significant percentage of neutral excitons show emission decay times up to 7 ns, while most of them exhibit lifetimes in the range 1–4 ns. Moreover a FSS remarkably lower than precedent nitrogen-free  $\text{InGaAs}/\text{GaAs}$  dots was measured.

Diluted nitride  $\text{InGaAs}_{1-\varepsilon}\text{N}_\varepsilon$  QDs were grown by MOVPE with standard metalorganic sources into  $7.5 \mu\text{m}$  pitch pyramidal recesses, chemically etched onto GaAs (111)B substrates and acting as nucleation seeds during the growth of the dot structure.<sup>10</sup> The sample design is similar to what published in Ref. 9, with the only difference being the 0.5 nm thick  $\text{In}_{0.25}\text{Ga}_{0.75}\text{As}_{1-\varepsilon}\text{N}_\varepsilon$  dot layer embedded between GaAs barriers. The QD layer was grown at  $730^\circ\text{C}$ , with an  $\text{AsH}_3/\text{III}$  ratio of 750 and a flux ratio Unsymmetrical-DimethylHydrazine(U-DMHy)/ $\text{AsH}_3$  of 2/3.

Although state of art diluted nitride QD systems are generally grown at relatively low temperatures<sup>7</sup> to ease the nitrogen incorporation,<sup>11</sup> our findings suggest that, a low percentage of nitrogen can be incorporated into our dot systems at temperatures as high as  $730^\circ\text{C}$ . Further studies with the goal to better understand the dependence of the nitrogen incorporation on the temperature (and V/III ratios) are necessary and ongoing on other specimens.

In Fig. 1 the layer sequence of a typical pyramidal QD structure is shown through an atomic force microscopy (AFM) image of a representative 1.5 nm thick

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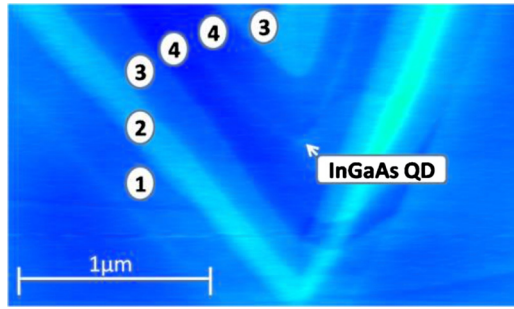


FIG. 1. (Color online) AFM image of a typical pyramidal QD scanned in cross-section. The sequence of the layers is numbered: 1. GaAs buffer; 2.  $\text{Al}_{0.8}\text{Ga}_{0.2}\text{As}$  etch stop cladding; 3.  $\text{Al}_{0.55}\text{Ga}_{0.45}\text{As}$  barriers; and 4. GaAs barriers. The 1.5 nm thick  $\text{In}_{0.35}\text{Ga}_{0.65}\text{As}$  dot is also indicated.

$\text{In}_{0.35}\text{Ga}_{0.65}\text{As}/\text{GaAs}$  dot, cleaved along one of the (110) directions and scanned in cross section. The dot, buried along the central axis of the inverted pyramid, is labeled in the same figure. The contrast between AlGaAs and pure GaAs is given by the presence of a thin oxide layer on AlGaAs alloy, while the stress released after cleaving the sample between InGaAs and the confining GaAs makes the dot layer identification clear.

After growth a processing known as surface etching was performed on the sample in order to chemically remove irregular surface faceting and nanoclusters whose emission can hide the signal coming from the QD.<sup>12</sup> We optically characterized the sample in top-view geometry via  $\mu$ -photoluminescence measurements at low temperature ( $\sim 10$  K), under nonresonant conditions. A continuous power excitation laser emitting at 658 nm was used for power dependence and exchange splitting measurements. A pulsed laser, with central wavelength at 656 nm, a repetition rate of 40 MHz, and a pulse duration of 400 ps, was employed for time-resolved investigations.

A preliminary analysis over a large number of single pyramidal QDs ( $\sim 60$ ) showed an average emission wavelength of the neutral exciton equal to  $\sim 853$  nm, with a standard deviation of 1.57 nm. As already anticipated, this represents only a small redshift if compared to previously studied, nitrogen-free, 0.5 nm thick  $\text{In}_{0.25}\text{Ga}_{0.75}\text{As}/\text{GaAs}$  QDs.<sup>9</sup> More importantly power dependence measurements revealed an unexpected antibinding biexciton energy. The assignment of the peaks was conducted on the base of combined power dependence, lifetime measurements and FSS characterization. In Fig. 2 the dynamics exciton/biexciton is illustrated for five different excitation power levels. At weak excitation levels the pure exciton peak at lower energy is dominant; the biexciton peak at higher energy becomes visible and more intense as the excitation power increases. From this power dependence evolution a clear antibinding behavior can be observed. This tendency was systematically found on all the dots analyzed, and the *average* value of the antibinding energy was found to be  $\sim 3.3$  meV.

Time-resolved measurements confirmed the nature of the biexciton and neutral exciton assignment: they systematically exhibited lifetimes which were, respectively, half the other, within a small variance, as summarized in Fig. 3(a). Most of excitons decayed with a time constant of 1–4 ns, but a non-negligible ( $< 7\%$ ) amount of neutral excitonic states, although emitting in the very same wavelength range, presents surprisingly longer lifetimes, up to 7 ns. In Figs. 3(b)

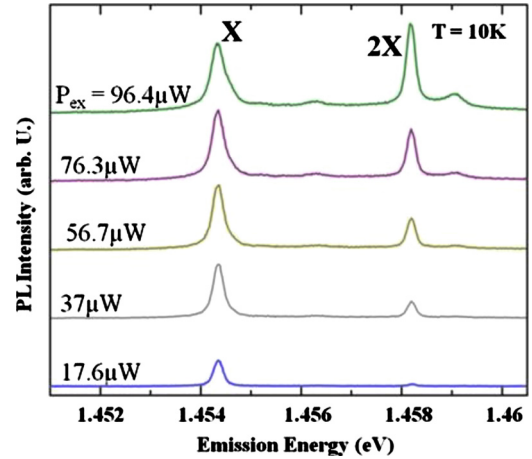


FIG. 2. (Color online) Low temperature excitation power dependence evolution of excitons and biexcitons emitted from a pyramidal  $\text{In}_{0.25}\text{Ga}_{0.75}\text{As}_{1-\epsilon}\text{N}_{\epsilon}/\text{GaAs}$  QD. Exciton (x) and biexciton (2X) peaks are labeled and their dynamics show a clear antibinding behavior.

and 3(c) representative emission decays of the pair exciton/biexciton emitted from two different pyramidal QDs are plotted in logarithmic scale. Correspondent lifetimes were extrapolated by fitting the data and they are indicated in the same figure with the exponential fit. Further analysis and more systematic characterizations are still ongoing to identify the cause of the lifetime values disparity.

Polarization dependent measurements<sup>13</sup> revealed extremely low values of FSS overall the sample, being the

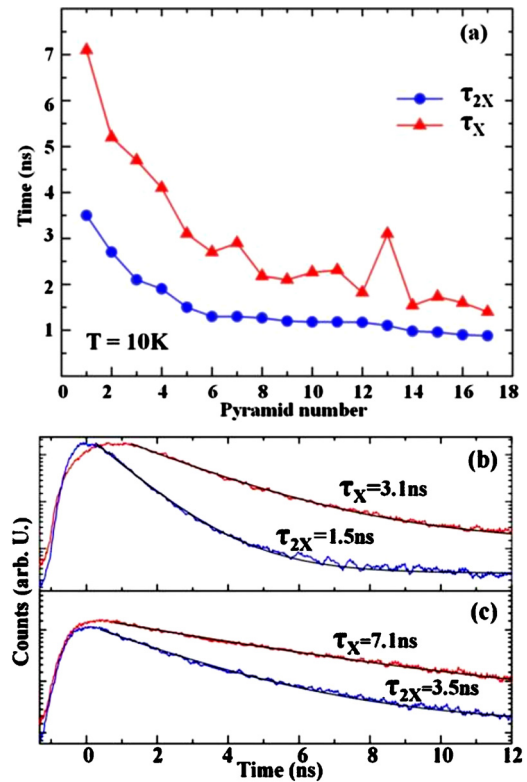


FIG. 3. (Color online) (a) Exciton and bi-exciton lifetimes measured for some representative InGaAsN/GaAs QDs at low temperature. For each dot excitons present a lifetime twice the biexciton one. [(b)–(c)] Exciton and bi-exciton emission decays, with relative exponential fit curves (smooth continuous lines), in logarithmic scale, as function of time measured from two different InGaAsN/GaAs QDs. Their lifetimes are indicated next to the correspondent curve.

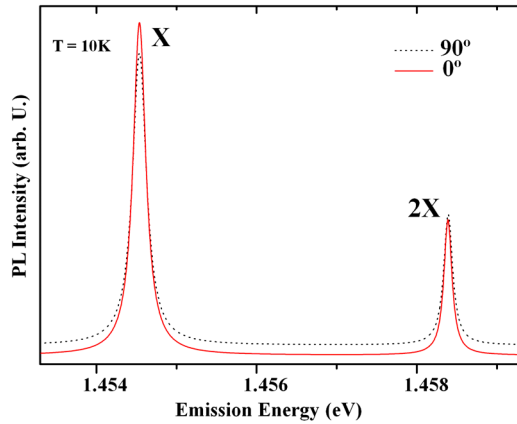


FIG. 4. (Color online) Lorentzian fit of normalized vertically (dotted trace) and horizontally (continuous trace) polarized emission spectra measured from a representative InGaAsN/GaAs QD. The two orthogonal linearly polarized emissions overlap, indicating a nearly zero value of the FSS, limited by the setup resolution ( $\sim 2\text{--}3 \mu\text{eV}$ ).

highest measured splitting  $3.4 \mu\text{eV}$ . Figure 4 shows an exemplifying (Lorentzian fit) spectrum of linearly polarized (vertical, dotted trace, and horizontal, continuous trace) neutral excitons and biexcitons emitted from a single In<sub>0.25</sub>Ga<sub>0.75</sub>As<sub>1- $\epsilon$</sub> N <sub>$\epsilon$</sub> /GaAs QD. The differently polarized emission spectra are not distinguishable, suggesting that the FSS value is limited by the combined system and fitting procedure resolution ( $< 3 \mu\text{eV}$ ). Compared to our nitrogen-free QDs, the introduction of nitrogen in the structure seems to suggest a remarkable improvement in terms of asymmetries of the crystallographic structure. Although the substrate orientation (111) and the patterning process assure a three-fold  $C_{3v}$  symmetry, ideally sufficient to guarantee a nearly zero excitonic anisotropic exchange splitting,<sup>14</sup> we always observe in our counterpart nitrogen-free samples a non-negligible FSS ( $> 10 \mu\text{eV}$ ).<sup>15</sup>

In conclusion we have encountered a number of unforeseen optical features in In<sub>0.25</sub>Ga<sub>0.75</sub>As<sub>1- $\epsilon$</sub> N <sub>$\epsilon$</sub> /GaAs pyramidal site-controlled QD samples. While further investigations are needed to determine the exact incorporation and localization of nitrogen atoms in the QD system (e.g., nonlinear incorporation of nitrogen due to In segregation,<sup>16</sup> distribution of the atoms and their possible diffusion into the confining barriers),<sup>17</sup> the exposure to U-DMHy during growth resulted in a significant change in the QD properties.

An antibinding biexciton was observed: this leads us to consider an hypothetical engineering mechanism of the biexciton binding energy (from negative to positive), based on a modulation of the U-DMHy flux during growth. Time-resolved measurements emphasized the presence of a significant percentage of dots with long lifetimes. Exchange splitting analysis revealed an extremely symmetrical nature of our QDs, being the FSS measured on different dots nearly zero. Causes of in-plane asymmetry reduction and/or strain redistribution need to be analyzed in detail. It has been proposed by pseudopotential calculations and confirmed by structural microscopy that the fine structure in MBE grown self-assembled QDs is mainly originated from a nonuniform incorporation of Indium in the islands.<sup>18</sup> This suggests that even in the case of MOVPE growth the interplay between U-DMHy and other precursors might lead to a different incorporation dynamics of In atoms, changing the segregation

effects in the dot layer. Nevertheless a fine control of the exchange splitting via tuning of the nitrogen incorporation can be conceived and interesting future opportunities are open for the generation of polarized entangled photons.

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<sup>11</sup>The difficulty related to nitrogen incorporation during MOVPE growths is amplified by the lack of data in the literature on (U-)DMHydrazine incorporation on (111) surfaces, and similarly on parameters like the excited atomic nitrogen diffusion length or its incorporation probability in the (111)B direction during MOVPE.

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