# EFFECT OF BACKGROUND INTENSITY ON RESOLUTION 

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ABSTRACT The paper discusses the effect of background intensity on the resolving power Tables, illustrated by graphs, have been given for the variation of resolving power with background intensity in case of liabry-Perot etalon, prism, grating and reflecting echelon, and when the instrumentai width is negligible

## INTRODUCTION

This paper discusses the effect of background intensity on the resolving power of spectroscopic instruments, on the Rayleigh's criterion of resolution of spectral lines. The iwo cases, viz., when the instrumental width is negligible as compared to the Doppler line-width and vice versa have been distinguished.

RESOLVING POWER WIIEN INSTRUMENTAL WIDTH IS NEGLIGIBLF

The intensity distribution of a spectral linc of wave number $v_{0}$ due to Doppler effect is given by

$$
I^{\prime}=I_{0} e^{-\beta\left(v-v_{0}\right)^{2}}
$$

where $\beta=\frac{\mu c^{2}}{2 R T v_{0}}{ }^{2}, \mu$ being the mass of radiant atoms.
The intensity distribution of another spectral line of wave number $v_{0}+\Delta v$ and same intensity is

$$
I^{\prime \prime}=I_{0} e^{-\beta\left(v-v_{0}-\Delta_{v}\right)^{2}}
$$

if $\Delta v$ is small. ( $\beta$ same for both lines).
Putting $\sqrt{ } \beta\left(\nu-v_{0}\right)=x$ and $\sqrt{ } \bar{\beta} \Delta v=a$, the resultant intensity pattern, in the presence of a background intensity $k I_{0}$ is given by

$$
\begin{equation*}
I=k I_{0}+I_{0} e^{-x^{2}}+I_{0} e^{-(x-a)^{2}} \tag{1}
\end{equation*}
$$

Oldenburg (1922) gives the position of maximum as

$$
x_{\max }=\frac{a}{e^{a^{2}}+1-2 a^{2}}
$$

In this section we will discuss cases with $a^{2}>3.524$ and we may assume, as a good approximation, $x_{\max } \approx 0$.
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Hence the intensity maxima and minimum ( $x=a / 2$ ) are given by
and

$$
\begin{array}{ll}
I_{\max } / I_{0}=1+k+e^{-a^{2}} & \ldots \\
I_{\min } / I_{0}=k+2 e^{-a^{2} / 4} & \ldots \\
\text { (2) }
\end{array}
$$

Kayleigh's criterion for resolution states

$$
I_{\text {min }} / I_{\text {max }}=0.8
$$

Hence we get

$$
\begin{equation*}
k=4\left(\mathrm{I}+e^{-a^{a}}\right)-\mathrm{IO} e^{-n^{2} / 4} \tag{4}
\end{equation*}
$$

We find that $k=0$ when $a^{2}=3.524$. If the resolving power be denoted by $R$ and its value, when $k=0$ by $R_{0}$ we have

$$
\begin{equation*}
R / R_{0}=\sqrt{\left(3.524 / a^{2}\right)} \tag{5}
\end{equation*}
$$

Table 1 gives $R / R_{0}$ for a few values of $k$.

## Table I

Variation of $R / R_{\mathrm{v}}$ with $k$ when instrumental width is negligible.

| $a^{2}$ | 3.524 | 4000 | 4400 | 4.800 | 5.600 | $6.4 \% 0$ | 7.500 | 9.200 | 12.003 | $\alpha$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\ldots R / R_{0}$ | 100 | 0939 | 0.895 | 0.857 | 0793 | 0.742 | 0.671 | 0619 | 0532 | 0.00 |
| $k$ | 000 | 039 | 0.72 | 1.02 | 155 | 1.99 | 250 | 3.00 | 3.50 | 4.00 |

RESOLVING POWER OFGRATING, RIVLECTING ECHIIION AND PRISM

The intensity of a spectral line diffracted by a grating or reflecting echelon is given by

$$
I^{\prime}=B \frac{\sin ^{2} N \beta}{\sin ^{2} \beta^{2}}
$$

where $N$ is the number of lines of the grating or the number of steps in the reflecting echelon and $2 \beta$ the phase difference between two adjacent beams. If the intensity maximum be denoted by $I_{0}$ we have

$$
\begin{equation*}
I^{\prime} / I_{0}=\frac{B}{B N^{2} \sin ^{2} \beta} \sin ^{2} N \beta \tag{6}
\end{equation*}
$$

where $x=N \beta$, and $\beta$ is small.
The above expression also represents the intensity distribution in a prism if $x=\pi l \sin \theta_{i} \lambda$.

The intensity distribution of another line of the same intensity and an angular separation corresponding to $\Delta x=a$ is represented by

$$
\begin{equation*}
J^{\prime \prime} / I_{0}=\frac{\sin ^{2}(x-a)}{(x-a)^{2}} \tag{7}
\end{equation*}
$$

The resultant intensity distribution of the two lines when the background intensity is $k I_{0}$ is given by

$$
I / I_{0}=k+\frac{\sin ^{2} x}{x^{2}}+\frac{\sin ^{2}(x-a)}{(x-a)^{2}}
$$

The intensity maxima ( $x \approx 0$ or $a$ ) and central minimum ( $x=a / 2$ ) are given by
and

$$
\begin{equation*}
I_{\max }=i+k+\frac{\sin ^{2}-a}{a^{2}} \tag{9}
\end{equation*}
$$

$$
\begin{equation*}
I_{\min }=k+2 \frac{\sin ^{2} a / 2}{(a / 2)^{2}} \tag{IO}
\end{equation*}
$$

Applying Rayleigh's criterion for resolution we get

$$
\begin{equation*}
k=4^{\prime} 1+\frac{\sin ^{2} a}{a^{2}}-10 \frac{\sin ^{2}(a / 2)}{(a / 2)^{2}} \tag{II}
\end{equation*}
$$

when $a=1.006 \pi, k=0$. Hence if $R_{0}$ denotes the resolving power for $k=0$ we have

$$
\begin{equation*}
R / R_{0}=1.006 \pi / a \tag{12}
\end{equation*}
$$

Table II gives $R / R_{0}$ for various values of $k$.
Tabre II
Variation of $R / R_{0}$ with $k$ for grating, reflecting echelon and prism.

| $a$ | $1.006 \pi$ | $1.10 \pi$ | $1.2 .0 \pi$ | $130 \pi$ | $1.40 \pi$ | $151 \pi$ | $1.60 \pi$ | $2.00 \pi$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $k$ | 0.00 | 0.77 | 1.55 | 225 | 2.83 | 3.28 | 3.60 | 1.03 |
| $R / R_{0}$ | 100 | 091 | 0.84 | 077 | 072 | 066 | 0.63 | 0.50 |

For $k>_{4}$ we get imaginary values for $a$ showing thereby that no resolution is possible. Physically it means that even single lines are not visible because $I_{\min } / I_{\max }>0.8$.

## RESOLVING POWER OF FABRY-PEROT ETALON

The intensity of a spectral line in the order $n_{0}+n$, where $n$ is small and $n_{0}$ an integer, is given in the case of Fabry-Perot etalon by

$$
I^{\prime}=I_{0} /\left\{\mathrm{I}+F \sin ^{2} \pi\left(n_{0}+n\right)\right\}=I_{0} /\left(\mathrm{I}+x^{2}\right)
$$

where $x=\pi n F^{1 / 2}, F$ being the coefficient of fineness.
The intensity of another spectral line of equal intensity separated by an order $\Delta n$ is given by

$$
I^{\prime \prime}=I_{0} /\left\{\mathrm{I}+F \sin ^{2} \pi\left(n_{0}+n-\Delta n\right)\right\}-I_{0} /\left\{\mathrm{I}+(x-a)^{2}\right\}
$$

where $a=\pi F^{1 / 2} \Delta n$.
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The resultant intensity pattern is

$$
\begin{equation*}
I / I_{n}=k+\left[\mathrm{I} /\left(\mathrm{I}+\mathrm{I}^{2}\right)\right]+\left[\mathrm{I} /\left\{\mathrm{I}+(\mathrm{x}-a)^{2}\right\}\right] \tag{13}
\end{equation*}
$$

The maximum ( $x \approx 0$ or $a$ ) and minimum $(x=a / 2)$ of the resultant intensity pattern, when the background intensity is $k I_{0}$, are given by
and

$$
\begin{array}{rlr}
I_{\text {min }} / I_{0}=k+2 /\left\{\mathrm{I}+a^{2} / 4\right\} & \ldots & (\mathrm{I} 4) \\
I_{\text {max }} / I_{0} & =k+\mathrm{I}+I_{i}\left(1 \mathrm{I}+a^{2}\right) & \ldots \\
(15)
\end{array}
$$

Rayleigh's criterion for resolution requires
or

$$
I_{\mathrm{mun}}=0.8 I_{\mathrm{max}}
$$

$$
\begin{gathered}
(4-k) a^{4}-(5 k+16) a^{2}-(4 k+8)=0 \\
a^{2}=(5 k+16)+\sqrt{(5 k+16)^{2}+16(4-k)(k+2)} \\
2(4-k)
\end{gathered}
$$

or
since $a$ is real.
The resolving power of the etalon is given by

$$
\begin{equation*}
R=\lambda / d \lambda-n_{0} / \Delta n=\frac{\pi n_{0} F^{1 / 2}\{2(4-k \cdot\} 1 / 2}{\left[(5 k+16)+\sqrt{\left.(5 k+16)^{2}+16(1-k)(k+2)\right]^{\frac{1}{2}}}\right.} \tag{16}
\end{equation*}
$$

The resolving power with $k=0$ (Meissner, 1941) is

$$
R_{0}=\mathrm{I}_{.49 n_{0}} F^{1 / 2}
$$

Hence $R / R_{11}=2.109\left[2(1-k) /\left\{(5 k+16)+\sqrt{ }(5 k+16)^{2}+16(4-k)(k+2)\right\}\right]^{\frac{1}{2}}$

Table III gives $R / R_{0}$ for some values of $k$.

## Table III

Variation of $R / R_{0}$ with $k$ for Fabry-Perot ctalon.

| 0.0 | 0.5 | 10 | 15 | 2.0 | 2.5 | 30 | 35 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $/ R_{0}$ | 1.00 | 0.89 | 0.77 | 0.67 | 057 | 0.47 | 0.38 | 0.26 |

For higher values of $k$ which give high values for $a$, eqn. ( 13 ) docs not hold because $\sin a$ differs appreciably from $a$ for large values of $a$.

## 1)ISCUSSION

The variation of $R / R_{0}$ with $k$, in the three cases discussed earlier is represented graphically in figure I . It can be seen that the background intensity has an important bearing on the choice of an instrument for a particular investigation. For example, a reflecting echelon, having half the resolving power as a Febry Perol ctalon for $k=0$ is superior to it for $k=3.5$.


Fig. 1
Variation of $R / k_{\mathrm{e}}$ with $k$
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