

Discriminating quantum states:  
the *multiple Chernoff distance*

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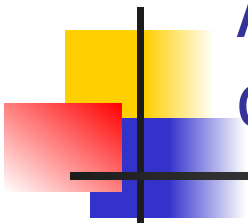
K. Li, *Annals of Statistics* 44: 1661-1679 (2016); arXiv:1508.06624



# Outline

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1. The problem
2. The answer
3. History review
4. Proof sketch
5. One-shot case
6. Open questions



## Accessing quantum systems: quantum measurement

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- Quantum measurement: formulated as positive operator-valued measure (POVM)

$$\mathcal{M} = \{M_i\}_i, \text{ with } 0 \leq M_i \leq \mathbb{1} \text{ and } \sum_i M_i = \mathbb{1};$$

when performing the POVM on a system in the state  $\omega$ , we obtain outcome " $i$ " with probability  $\text{Tr}(\omega M_i)$ .

- von Neumann measurement: special case of POVM, with the POVM elements being orthogonal projectors:  $M_i M_j = \delta_{ij} M_i$ , where  $\delta_{ij}$  is the Kronecker delta.



## Quantum state discrimination (quantum hypothesis testing)

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- Suppose a quantum system is in one of a set of states  $\{\omega_1, \dots, \omega_r\}$ , with a given prior  $\{p_1, \dots, p_r\}$ . The task is to detect the true state with a minimal error probability.

- Method: making quantum measurement  $\{M_i\}_{i=1}^r$ .

- Error probability (let  $A_i := p_i \omega_i$ )

$$P_e(\{A_1, \dots, A_r\}; \{M_1, \dots, M_r\}) := \sum_{i=1}^r \text{Tr} A_i (\mathbb{1} - M_i).$$

- Optimal error probability

$$P_e^*(\{A_1, \dots, A_r\}) :=$$

$$\min \left\{ P_e(\{A_1, \dots, A_r\}; \{M_1, \dots, M_r\}) : \text{POVM } \{M_1, \dots, M_r\} \right\}.$$



## Asymptotics in quantum hypothesis testing

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- What's the asymptotic behavior of  $P_e^* (\{p_1 \rho_1^{\otimes n}, \dots, p_r \rho_r^{\otimes n}\})$ , as  $n \rightarrow \infty$  ?

- Exponentially decay! (Parthasarathy '2001)

$$P_e^* \sim \exp(-\xi n)$$

- But, what's the error exponent

$$\xi = \liminf_{n \rightarrow \infty} \frac{-1}{n} \log P_e^* (\{p_1 \rho_1^{\otimes n}, \dots, p_r \rho_r^{\otimes n}\}) \quad ?$$

It has been an open problem (except for  $r=2$ )!



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Our result:

error exponent = multiple Chernoff distance

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- We prove that

**Theorem** *Let  $\{\rho_1, \dots, \rho_r\}$  be a finite set of quantum states on a finite-dimensional Hilbert space  $\mathcal{H}$ . Then the asymptotic error exponent for testing  $\{\rho_1^{\otimes n}, \dots, \rho_r^{\otimes n}\}$ , for an arbitrary prior  $\{p_1, \dots, p_r\}$ , is given by the multiple quantum Chernoff distance:*

$$\lim_{n \rightarrow \infty} \frac{-1}{n} \log P_e^* (\{p_1 \rho_1^{\otimes n}, \dots, p_r \rho_r^{\otimes n}\}) = \min_{(i,j): i \neq j} \max_{0 \leq s \leq 1} \left\{ -\log \text{Tr} \rho_i^s \rho_j^{1-s} \right\}. \quad (1)$$



## Remarks

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- Remark 1: Our result is a multiple-hypothesis generalization of the  $r=2$  case. Denote the **multiple quantum Chernoff distance** (r.h.s. of eq. (1)) as  $C(\rho_1, \dots, \rho_r)$ , then

$$C(\rho_1, \dots, \rho_r) = \min_{(i,j):i \neq j} C(\rho_i, \rho_j),$$

with the **binary quantum Chernoff distance** is defined as

$$C(\rho_1, \rho_2) := \max_{0 \leq s \leq 1} \{-\log \text{Tr} \rho_1^s \rho_2^{1-s}\}.$$

- Remark 2: when  $\rho_1, \dots, \rho_r$  commute, the problem reduces to classical statistical hypothesis testing. Compared to the classical case, the difficulty of quantum statistics comes from **noncommutativity & entanglement**.





# Outline

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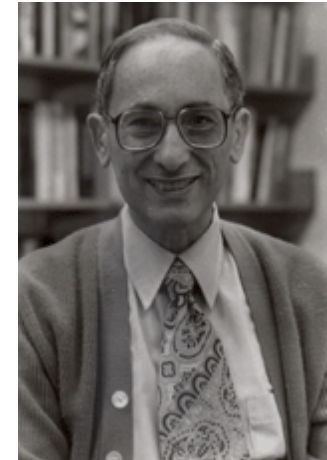
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## Some history review

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- The classical Chernoff distance as the optimal error exponent for testing two probability distributions was given in [H. Chernoff, Ann. Math. Statist. 23, 493 \(1952\)](#).
- The multiple generalizations were subsequently made in
  - [N. P. Salihov, Dokl. Akad. Nauk SSSR 209, 54 \(1973\)](#);
  - [E. N. Torgersen, Ann. Statist. 9, 638 \(1981\)](#);
  - [C. C. Leang and D. H. Johnson, IEEE Trans. Inf. Theory 43, 280 \(1997\)](#);
  - [N. P. Salihov, Teor. Veroyatn. Primen. 43, 294 \(1998\)](#).





## Some history review

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- Quantum hypothesis testing (state discrimination) was the main topic in the early days of quantum information theory in 1970s.

- Maximum likelihood estimation

- for two states: Holevo-Helstrom tests

$$(\{\rho_1 - \rho_2 > 0\}, \mathbb{1} - \{\rho_1 - \rho_2 > 0\})$$

C. W. Helstrom, *Quantum Detection and Estimation Theory*, Academic Press (1976); A. S. Holevo, *Theor. Prob. Appl.* 23, 411 (1978).

- for more than two states: only formulated in a complex and implicit way. **Competitions between pairs** make the problem complicated!

A. S. Holevo, *J. Multivariate Anal.* 3, 337 (1973); H. P. Yuen, R. S. Kennedy and M. Lax, *IEEE Trans. Inf. Theory* 21, 125 (1975).



## Some history review

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- In 2001, Parthasarathy showed exponential decay.  
K. R. Parthasarathy, in *Stochastics in Finite and Infinite Dimensions* 361 (2001).
- In 2006, two groups [Audenaert et al] and [Nussbaum & Szkola] together solved the  $r=2$  case.  
K. Audenaert et al, arXiv: quant-ph/0610027; *Phys. Rev. Lett.* 98, 160501 (2007);  
M. Nussbaum and A. Szkola, arXiv: quant-ph/0607216 ; *Ann. Statist.* 37, 1040 (2009).
- In 2010/2011, Nussbaum & Szkola conjectured the solution (our theorem), and proved that  $C/3 \leq \xi \leq C$  .  
M. Nussbaum and A. Szkola, *J. Math. Phys.* 51, 072203 (2010); *Ann. Statist.* 39, 3211 (2011).
- In 2014, Audenaert & Mosonyi proved that  $C/2 \leq \xi \leq C$  .  
K. Audenaert and M. Mosonyi, *J. Math. Phys.* 55, 102201 (2014).



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## Sketch of proof

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- We only need to prove the achievability part " $\xi \geq C$ ".

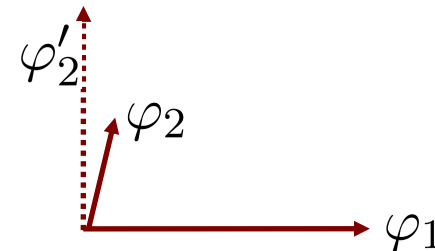
For this purpose, we construct an **asymptotically optimal quantum measurement**, and show that it achieves the quantum multiple Chernoff distance as the error exponent.

- Motivation: consider detecting two weighted pure states.

**Big overlap:** give up the light one;



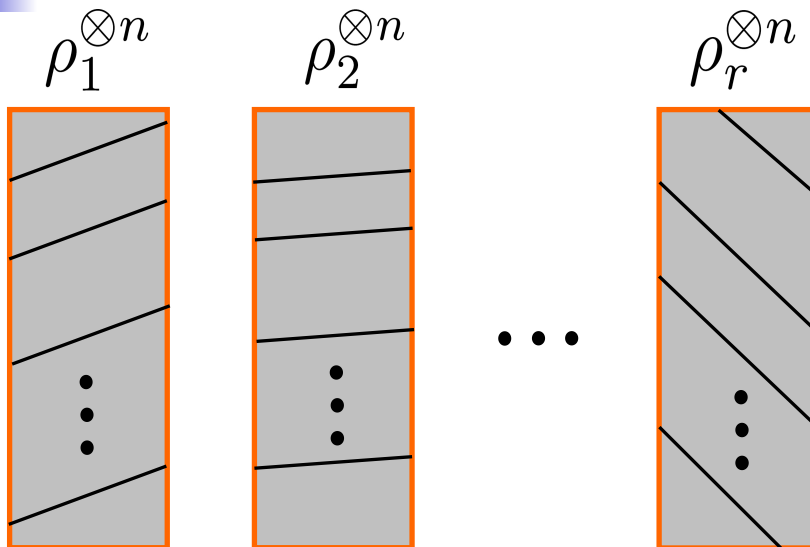
**Small overlap:** make a projective measurement, using orthonormalized version of the two states.





## Sketch of proof

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Spectral decomposition:

$$\rho_i^{\otimes n} = \bigoplus_{k=1}^{T_i} \lambda_{ik}^{(n)} Q_{ik}^{(n)},$$

$$T := \max\{T_i\}_i \leq (n+1)^d$$

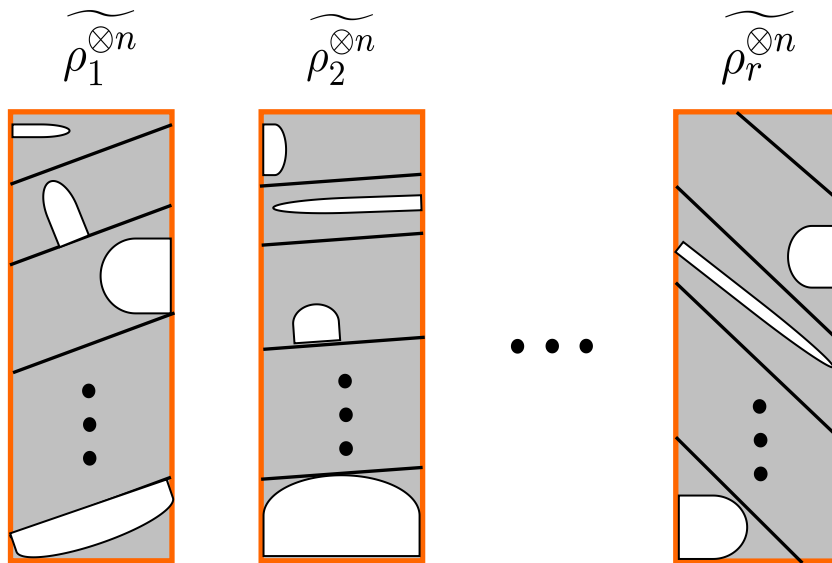
Overlap between eigenspaces:

$$\text{Olap} \left( \text{supp} \left( Q_{ik}^{(n)} \right), \text{supp} \left( Q_{jl}^{(n)} \right) \right)$$

$$:= \max \left\{ |\langle \varphi | \phi \rangle| : |\varphi\rangle \in \text{supp} \left( Q_{ik}^{(n)} \right), |\phi\rangle \in \text{supp} \left( Q_{jl}^{(n)} \right) \right\}$$

# Sketch of proof

"Dig holes" in every eigenspaces to reduce overlaps



$\epsilon$ -subtraction:

$$\text{Let } P_1 P_2 P_1 = \bigoplus_x \lambda_x Q_x$$

$$\text{Define } P_1 \ominus_\epsilon P_2 := P_1 - \sum_{x: \lambda_x \geq \epsilon^2} Q_x$$

$$\widetilde{\rho}_i^{\otimes n} = \bigoplus_{k=1}^{T_i} \lambda_{ik}^{(n)} \widetilde{Q}_{ik}^{(n)}, \quad \text{Olap} \left( \text{supp} \left( \widetilde{Q}_{ik}^{(n)} \right), \text{supp} \left( \widetilde{Q}_{jl}^{(n)} \right) \right) \leq \epsilon$$

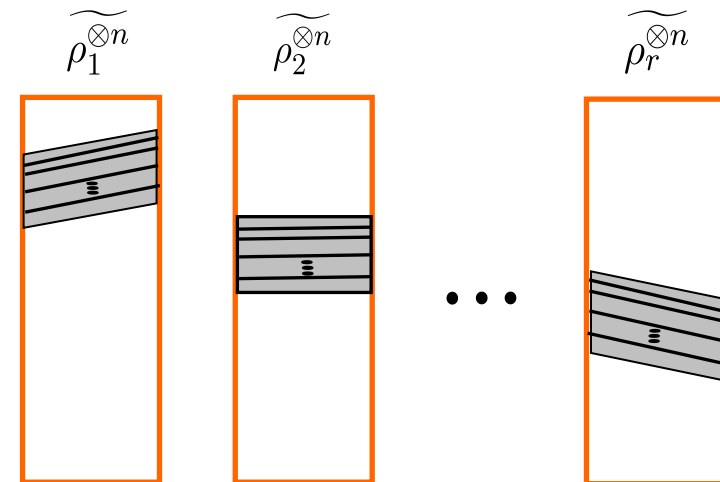




## Sketch of proof

- Now the supporting space of the hypothetical states have small overlaps. For  $i \neq j$ ,

$$\text{Olap} \left( \text{supp} \left( \widetilde{\rho}_i^{\otimes n} \right), \text{supp} \left( \widetilde{\rho}_j^{\otimes n} \right) \right) \leq T\epsilon$$



- The next step is to orthogonalize these eigenspaces
  - Order the eigenspaces according to their eigenvalues, in the decreasing order.
  - Orthogonalization using the Gram-Schmidt process.

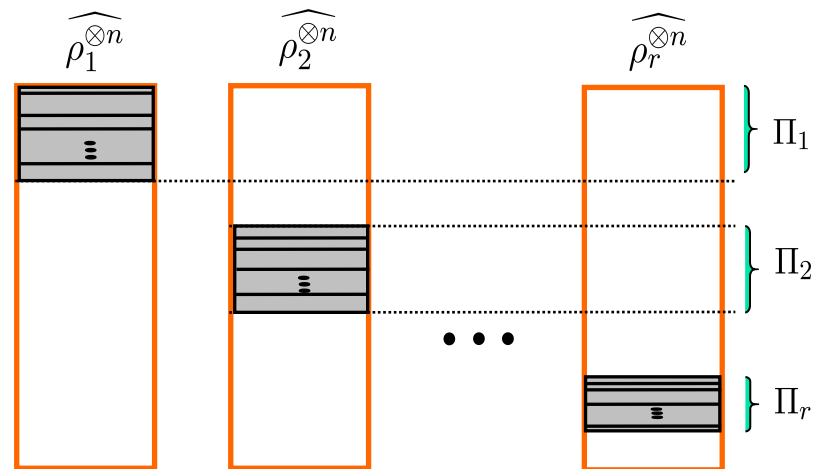
## Sketch of proof

- Now the eigenspaces are all orthogonal.

$$\widehat{\rho}_i^{\otimes n} = \bigoplus_{k=1}^{T_i} \lambda_{ik}^{(n)} \widehat{Q}_{ik}^{(n)}$$

- We construct a projective measurement

$$\left\{ \Pi_i = \bigoplus_k \widehat{Q}_{ik}^{(n)} \right\}_{i=1}^r$$



- Use this to discriminate the original states:

$$P_{succ} = \sum_{i=1}^r p_i \text{Tr} \rho_i^{\otimes n} \Pi_i$$



## Sketch of proof

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$$Q_{ik}^{(n)} \xrightarrow{\text{"digging holes"}} \widetilde{Q}_{ik}^{(n)} \xrightarrow{\text{orthogonalization}} \widehat{Q}_{ik}^{(n)}$$

- Loss in "digging holes":

$$\text{Tr} \left( Q_{ik}^{(n)} - \widetilde{Q}_{ik}^{(n)} \right) \leq \frac{1}{\epsilon^2} \sum_{(j,\ell): \lambda_{j\ell}^{(n)} > \lambda_{ik}^{(n)}} \text{Tr} Q_{ik}^{(n)} Q_{j\ell}^{(n)}$$

- Mismatch due to orthogonalization:

$$\text{Tr} \left[ \widetilde{Q}_{ik}^{(n)} \left( \mathbb{1} - \widehat{Q}_{ik}^{(n)} \right) \right] \leq \frac{1 - (r-1)T\epsilon}{1 - 2(r-1)T\epsilon} \sum_{(j,\ell): \lambda_{j\ell}^{(n)} > \lambda_{ik}^{(n)}} \text{Tr} Q_{ik}^{(n)} Q_{j\ell}^{(n)}$$

- Estimation of the total error:

$$P_e \leq \sum_{(i,k)} \lambda_{ik}^{(n)} \text{Tr} \left[ Q_{ik}^{(n)} \left( \mathbb{1} - \widehat{Q}_{ik}^{(n)} \right) \right] \leq \sum_{(i,k)} \lambda_{ik}^{(n)} \left\{ \text{Tr} \left( Q_{ik}^{(n)} - \widetilde{Q}_{ik}^{(n)} \right) + \text{Tr} \left[ \widetilde{Q}_{ik}^{(n)} \left( \mathbb{1} - \widehat{Q}_{ik}^{(n)} \right) \right] \right\}$$



## Sketch of proof

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$$P_e \leq \underbrace{\left( \frac{1}{\epsilon^2} + \frac{1 - (r-1)T\epsilon}{1 - 2(r-1)T\epsilon} \right)}_{\leq p(n)} \sum_{(i,j):i \neq j} \sum_{k,\ell} \underbrace{\min\{\lambda_{ik}^{(n)}, \lambda_{j\ell}^{(n)}\}}_{\leq \left(\lambda_{ik}^{(n)}\right)^s \left(\lambda_{j\ell}^{(n)}\right)^{(1-s)}} \text{Tr} Q_{ik}^{(n)} Q_{j\ell}^{(n)}$$

$$\leq p(n) \sum_{(i,j):i \neq j} \min_{0 \leq s \leq 1} \left( \text{Tr} \rho_i^s \rho_j^{(1-s)} \right)^n$$

$$\sim \exp \left\{ -n \left( \min_{(i,j):i \neq j} \max_{0 \leq s \leq 1} \left\{ -\log \text{Tr} \rho_i^s \rho_j^{1-s} \right\} \right) \right\}$$





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## Result for the one-shot case

**Theorem** *Let  $A_1, \dots, A_r \in \mathcal{P}(\mathcal{H})$  be nonnegative matrices on a finite-dimensional Hilbert space  $\mathcal{H}$ . For all  $1 \leq i \leq r$ , let  $A_i = \bigoplus_{k=1}^{T_i} \lambda_{ik} Q_{ik}$  be the spectral decomposition of  $A_i$ , and write  $T := \max\{T_1, \dots, T_r\}$ . Then*

$$P_e^* (\{A_1, \dots, A_r\}) \leq 10(r-1)^2 T^2 \sum_{(i,j): i < j} \sum_{k,\ell} \min\{\lambda_{ik}, \lambda_{j\ell}\} \text{Tr } Q_{ik} Q_{j\ell}.$$

- **Remark 1:** It matches a lower bound up to some states-dependent factors:

$$P_e^* (\{A_1, \dots, A_r\}) \geq \frac{1}{2(r-1)} \sum_{(i,j): i < j} \sum_{k,\ell} \min\{\lambda_{ik}, \lambda_{j\ell}\} \text{Tr } Q_{ik} Q_{j\ell}.$$

Obtained by combining [M. Nussbaum and A. Szkola, *Ann. Statist.* 37, 1040 (2009)] and [D.-W. Qiu, *PRA* 77. 012328 (2008)].



## Result for the one-shot case

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- Remark 2: for the case  $r=2$ , we have

$$P_e^* (\{A_1, A_2\}) \leq 10T^2 \sum_{k,l} \min\{\lambda_{1k}, \lambda_{2l}\} \text{Tr } Q_{1k} Q_{2l}.$$

On the other hand, it is proved in [[K. Audenaert et al, PRL, 2007](#)] that

$$P_e^* (\{A_1, A_2\}) \leq \min_{0 \leq s \leq 1} \text{Tr } A_1^s A_2^{1-s}.$$

(note that it is always true that

$$\sum_{k,l} \min\{\lambda_{1k}, \lambda_{2l}\} \text{Tr } Q_{1k} Q_{2l} \leq \min_{0 \leq s \leq 1} \text{Tr } A_1^s A_2^{1-s} . )$$



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## Open questions

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### 1. Applications of the bounds:

$$P_e^* (\{A_1, \dots, A_r\}) \begin{cases} \leq 10(r-1)^2 T^2 \sum_{(i,j):i<j} \sum_{k,l} \min\{\lambda_{ik}, \lambda_{jl}\} \text{Tr } Q_{ik} Q_{jl} \\ \geq \frac{1}{2(r-1)} \sum_{(i,j):i<j} \sum_{k,l} \min\{\lambda_{ik}, \lambda_{jl}\} \text{Tr } Q_{ik} Q_{jl} \end{cases}$$

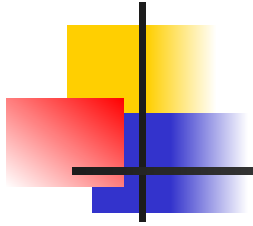
### 2. Strengthening the states-dependent factors

### 3. Testing composite hypotheses:

$$\rho^{\otimes n} \quad \text{Vs} \quad \sum_i q_i \sigma_i^{\otimes n} \quad (\text{or, } \int \sigma^{\otimes n} d\mu(\sigma))$$

K. Audenaert and M. Mosonyi, J. Math. Phys. 55, 102201 (2014).

Brandao, Harrow, Oppenheim and Strelchuk, PRL 115, 050501 (2015).



Thank you !