

# Operator algebras and data hiding in topologically ordered systems

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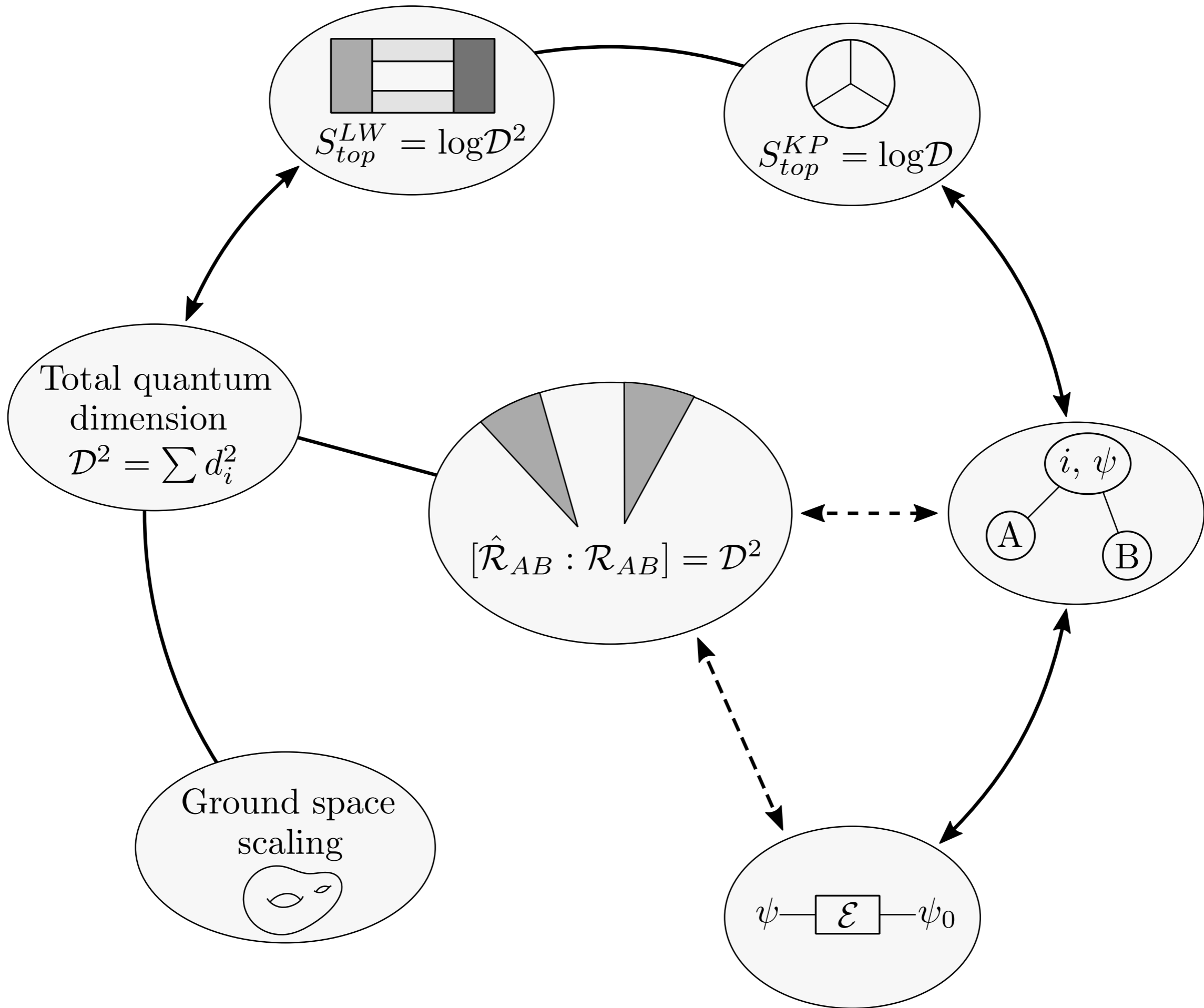
Tobias Osborne

UC Davis & RWTH Aachen

arXiv:1608.02618

9 October 2016

QMath 13



# Topological order

# Topological phases

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Quantum phase outside of Landau theory

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Describes all properties of the anyons, e.g.  
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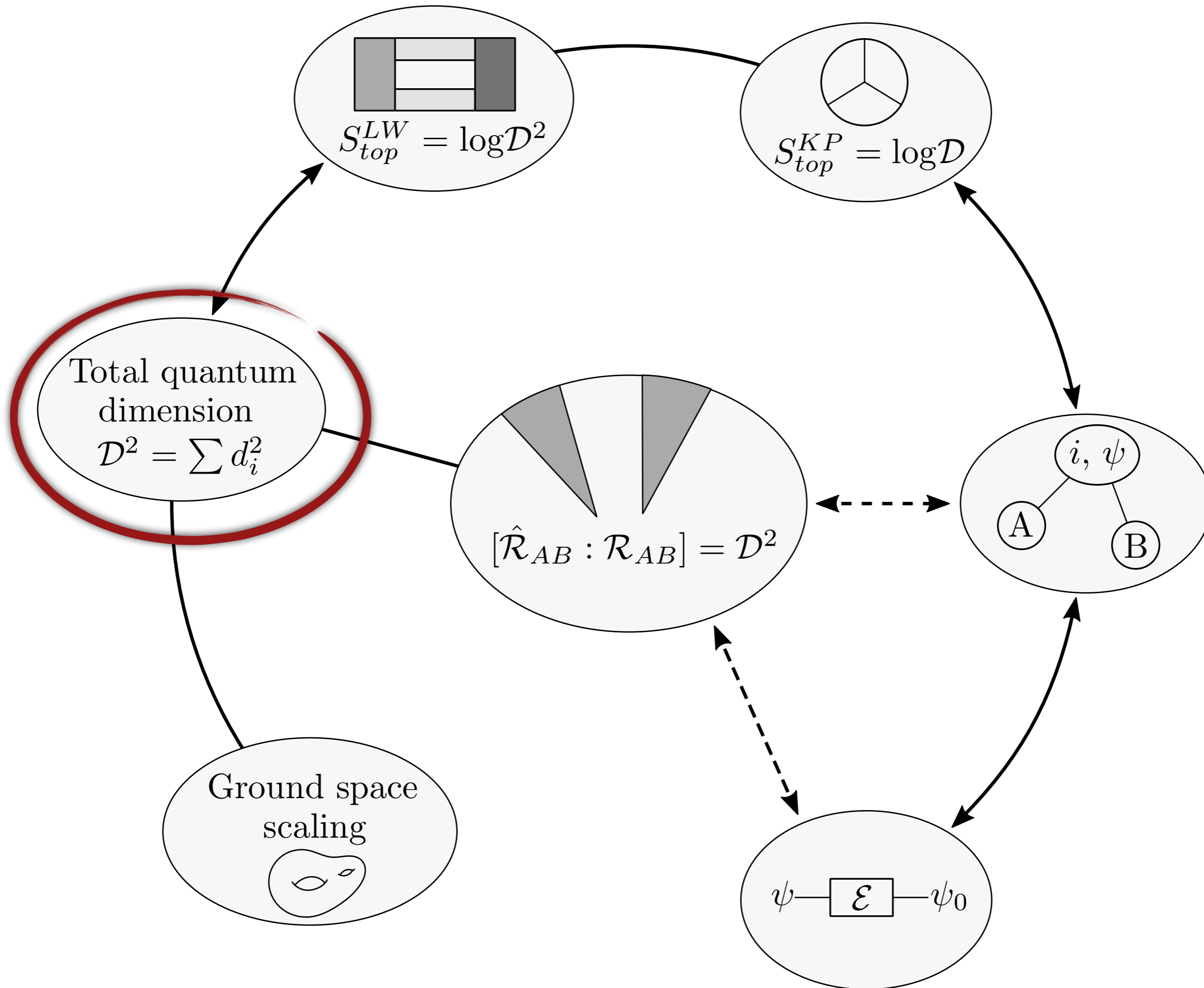
Irreducible objects  $\rho_i \Leftrightarrow$  anyons

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Irreducible objects  $\rho_i \Leftrightarrow$  anyons

**Quantum dimension**  $\mathcal{D}^2 = \sum_i d(\rho_i)^2$



# Topological entanglement entropy

Area law for top. ordered states:

$$S_{\Lambda} = \alpha |\partial\Lambda| - \gamma + \dots$$

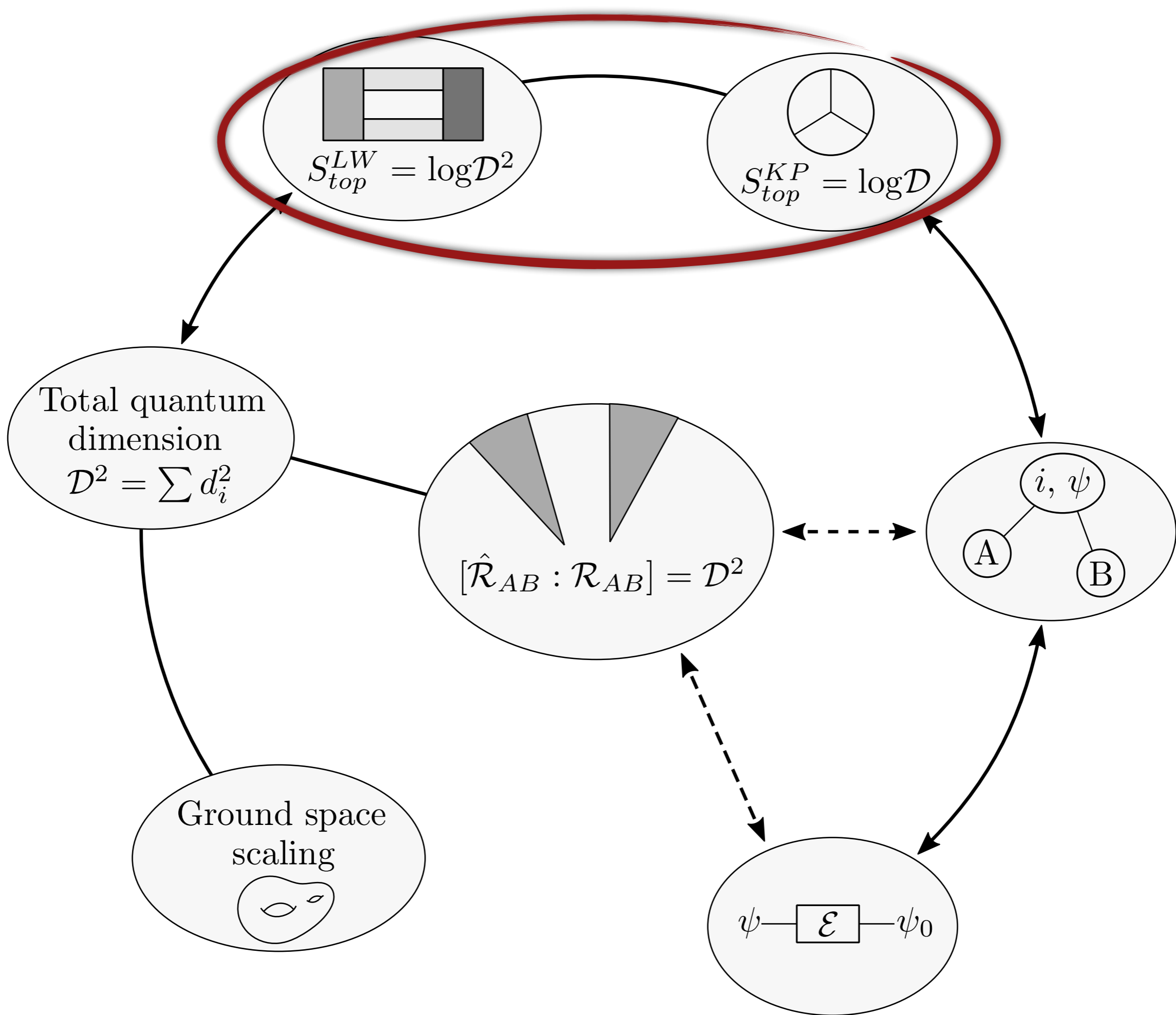
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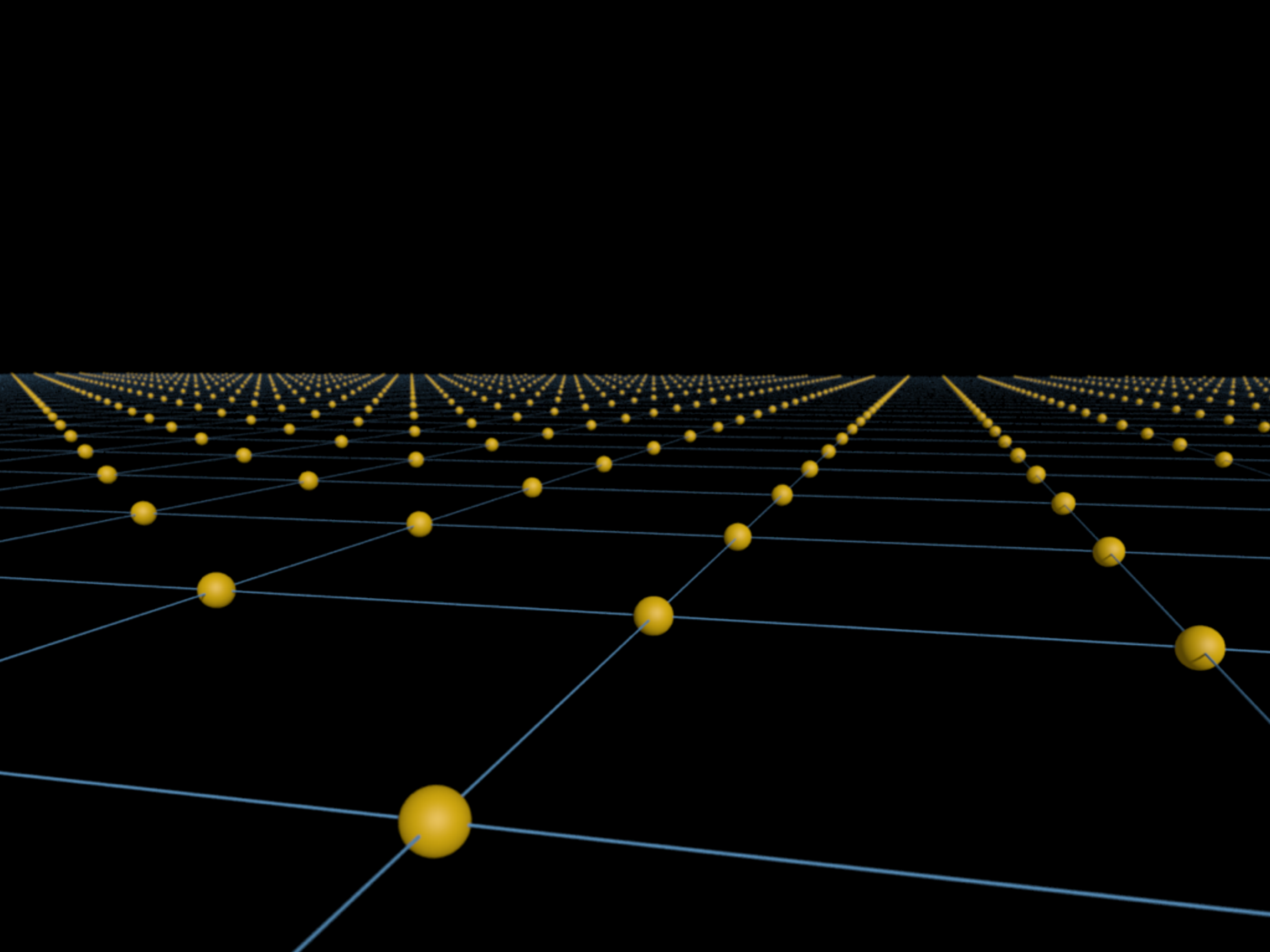
$$S_{\Lambda} = \alpha |\partial\Lambda| - \gamma + \dots$$

Universal constant:  $\gamma = \log \mathcal{D}$





# Technical framework



Quasi-local algebra  $\mathfrak{A} = \overline{\bigcup_{\Lambda} \mathfrak{A}(\Lambda)}^{\|\cdot\|}$


$$\mathfrak{A}(\Lambda) = \bigotimes_{x \in \Lambda} M_d(\mathbb{C})$$

and **local Hamiltonians**  $H_\Lambda \in \mathfrak{A}(\Lambda)$


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ground state representation  $\pi_0$


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# Example: toric code

$\omega_0 \circ \rho$  is a **single excitation state**

$$\rho(A) := \lim_{n \rightarrow \infty} F_{\xi_n} A F_{\xi_n}^*$$

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$\pi_0 \circ \rho$  describes  
observables in  
presence of  
**background charge**



# Quantum dimension



$$\mathcal{R}_A = \pi_0(\mathfrak{A}(A))''$$


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$$\mathcal{R}_{AB} = \mathcal{R}_A \vee \mathcal{R}_B$$


$$\hat{\mathcal{R}}_{AB} = \pi_0(\mathfrak{A}((A \cup B)^c))'$$

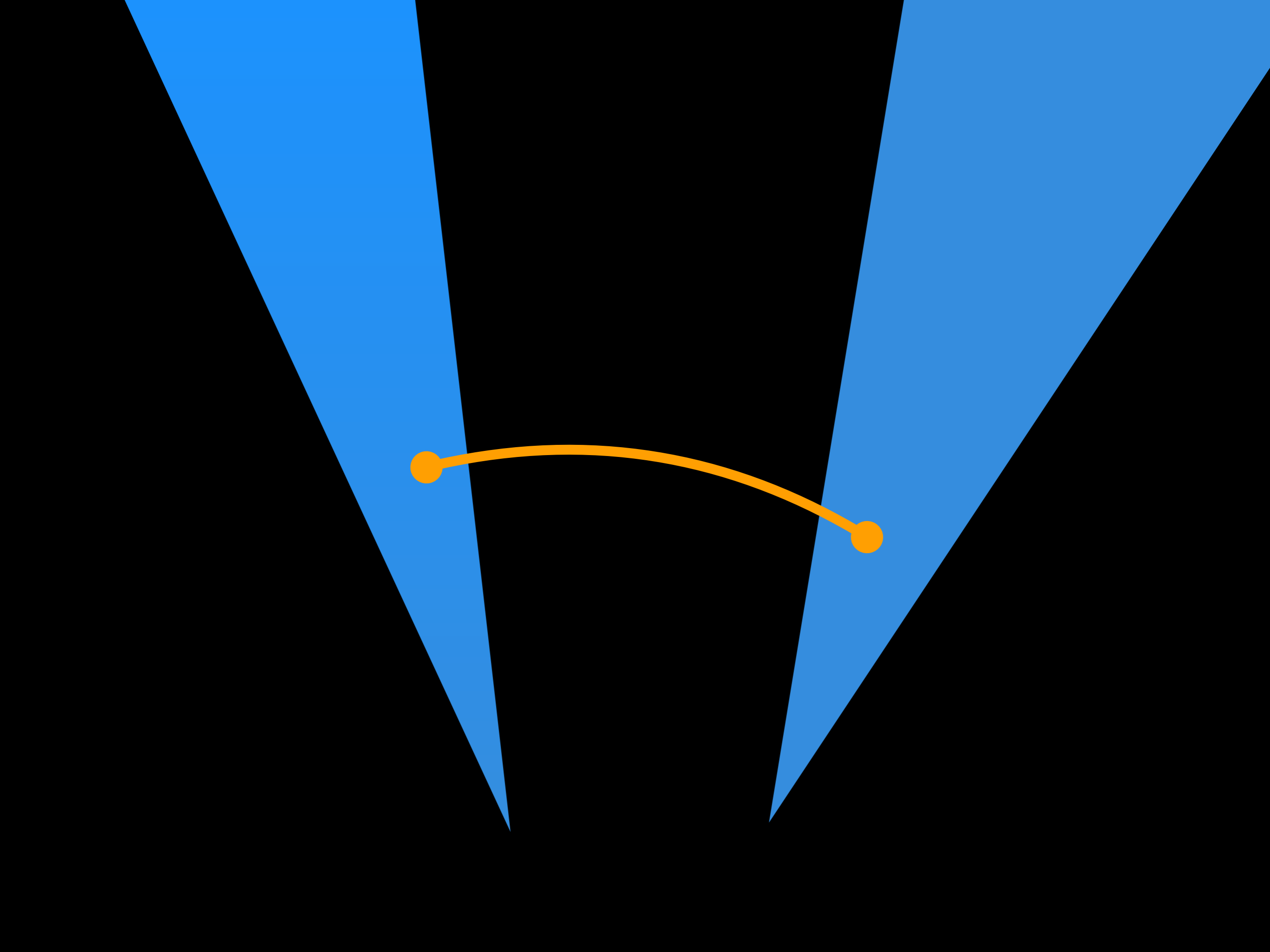
Locality:  $\mathcal{R}_{AB} \subset \hat{\mathcal{R}}_{AB}$

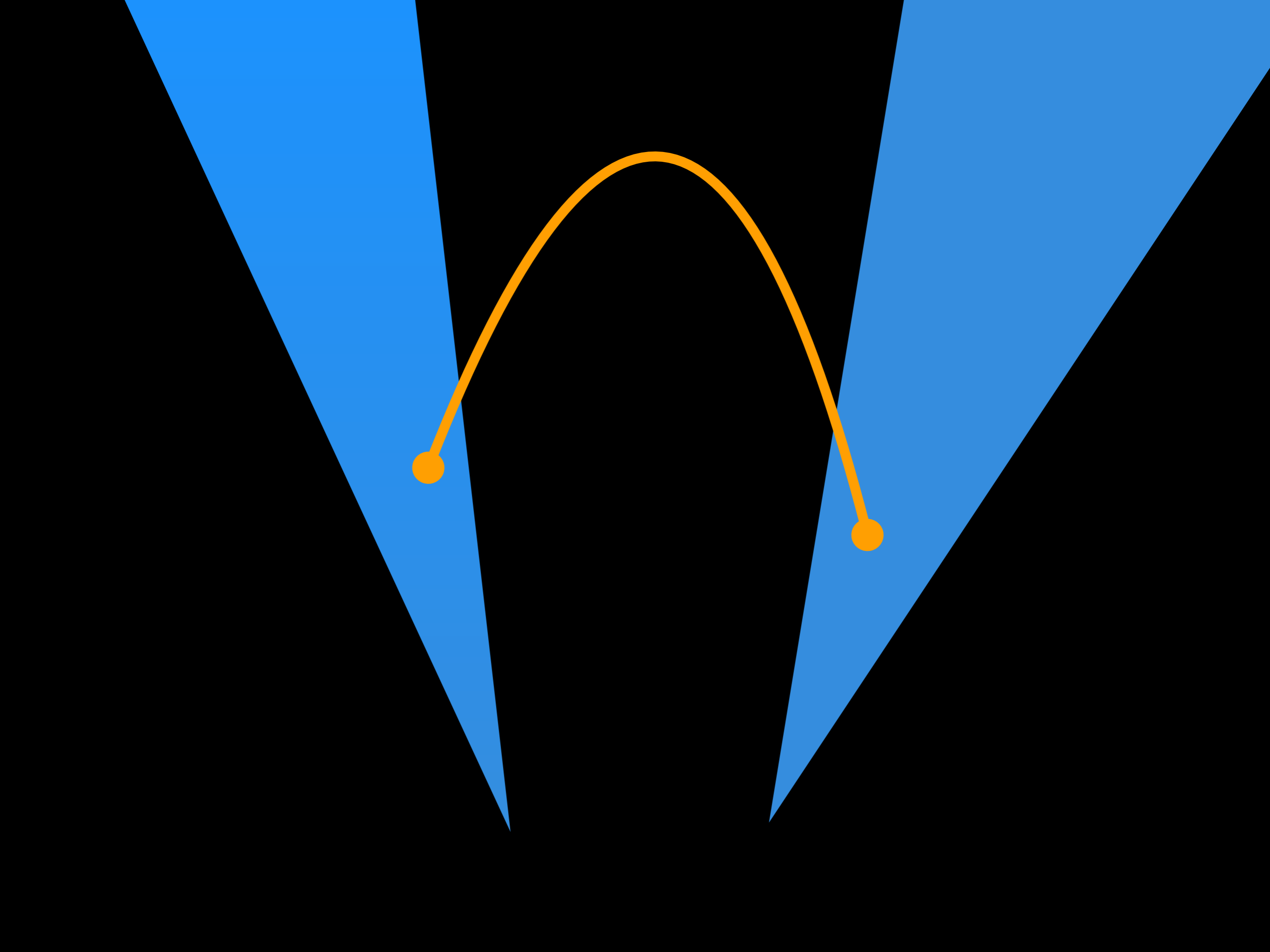
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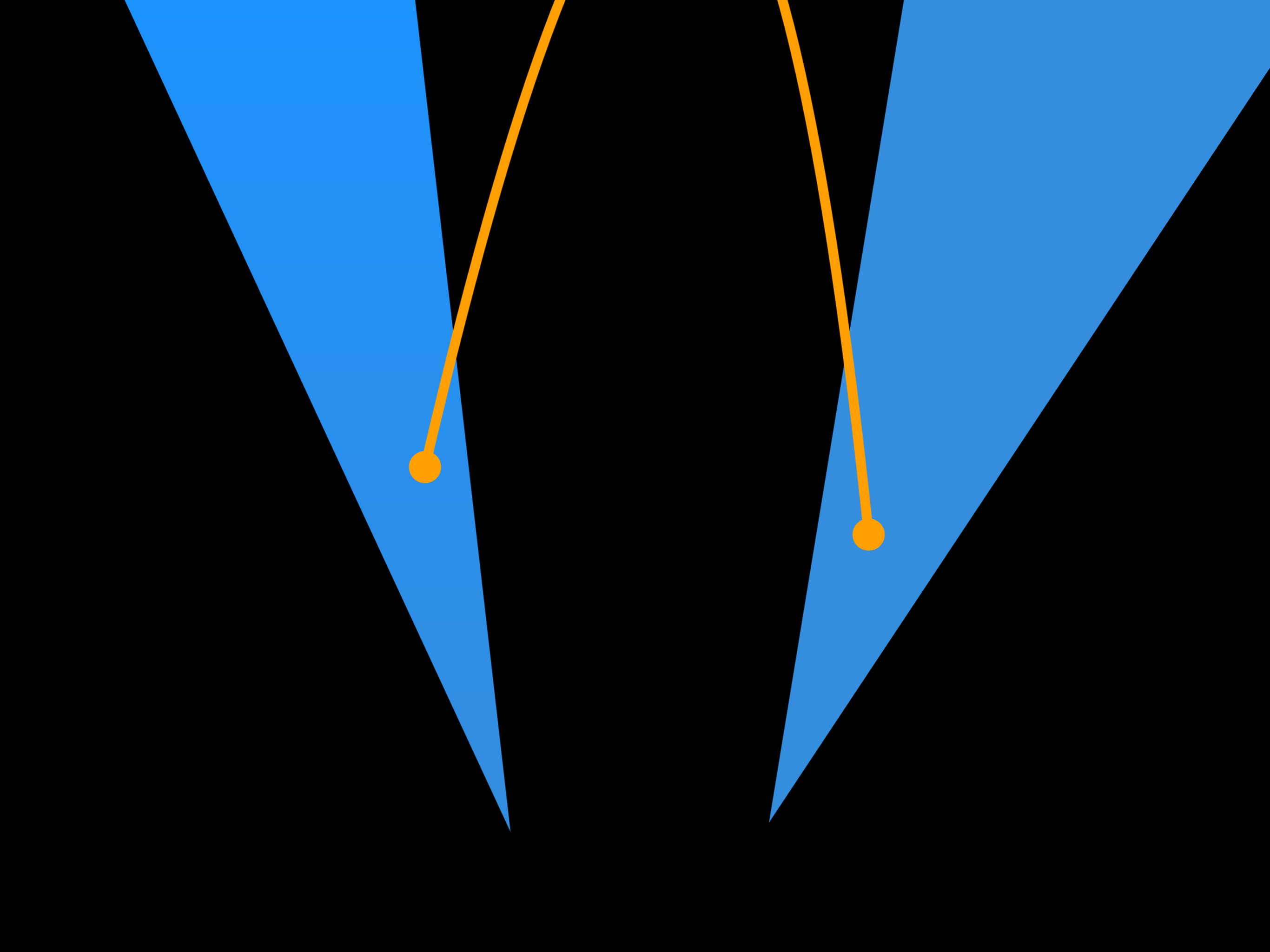
but:

$$\mathcal{R}_{AB} \subsetneq \hat{\mathcal{R}}_{AB}$$











**Weak-operator limit is in  $\hat{\mathcal{R}}_{AB}$**

Jones-Kosaki-Longo index  $[\widehat{\mathcal{R}}_{AB} : \mathcal{R}_{AB}]$

**Weak-operator limit is in  $\widehat{\mathcal{R}}_{AB}$**

## Theorem

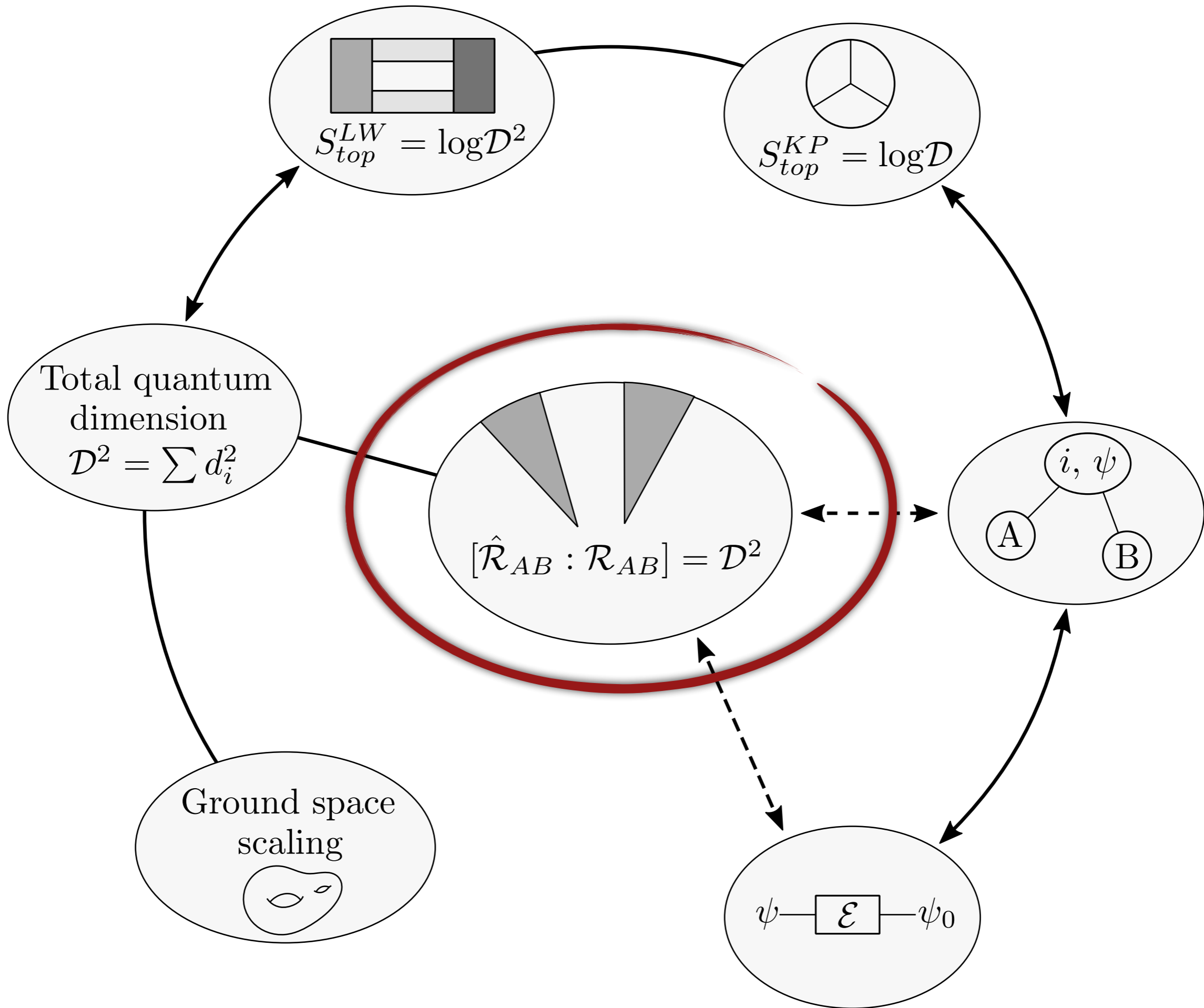
The number of excitation types is bounded by

$$\mu_{\pi_0} = \sup_{A \cup B} [\widehat{\mathcal{R}}_{AB} : \mathcal{R}_{AB}]$$

If all excitations have conjugates,  $\mu_{\pi_0}$  is equal to the **total quantum dimension**.

PN, J. Math. Phys. '13

Kawahigashi, Longo & Müger, Commun. Math. Phys. '01



# Data hiding



# A data hiding task

Alice

Bob

Eve

# A data hiding task

Alice

Bob

Eve

Operations in  $\hat{\mathcal{R}}_{AB}$  are invisible to Eve

# A data hiding task

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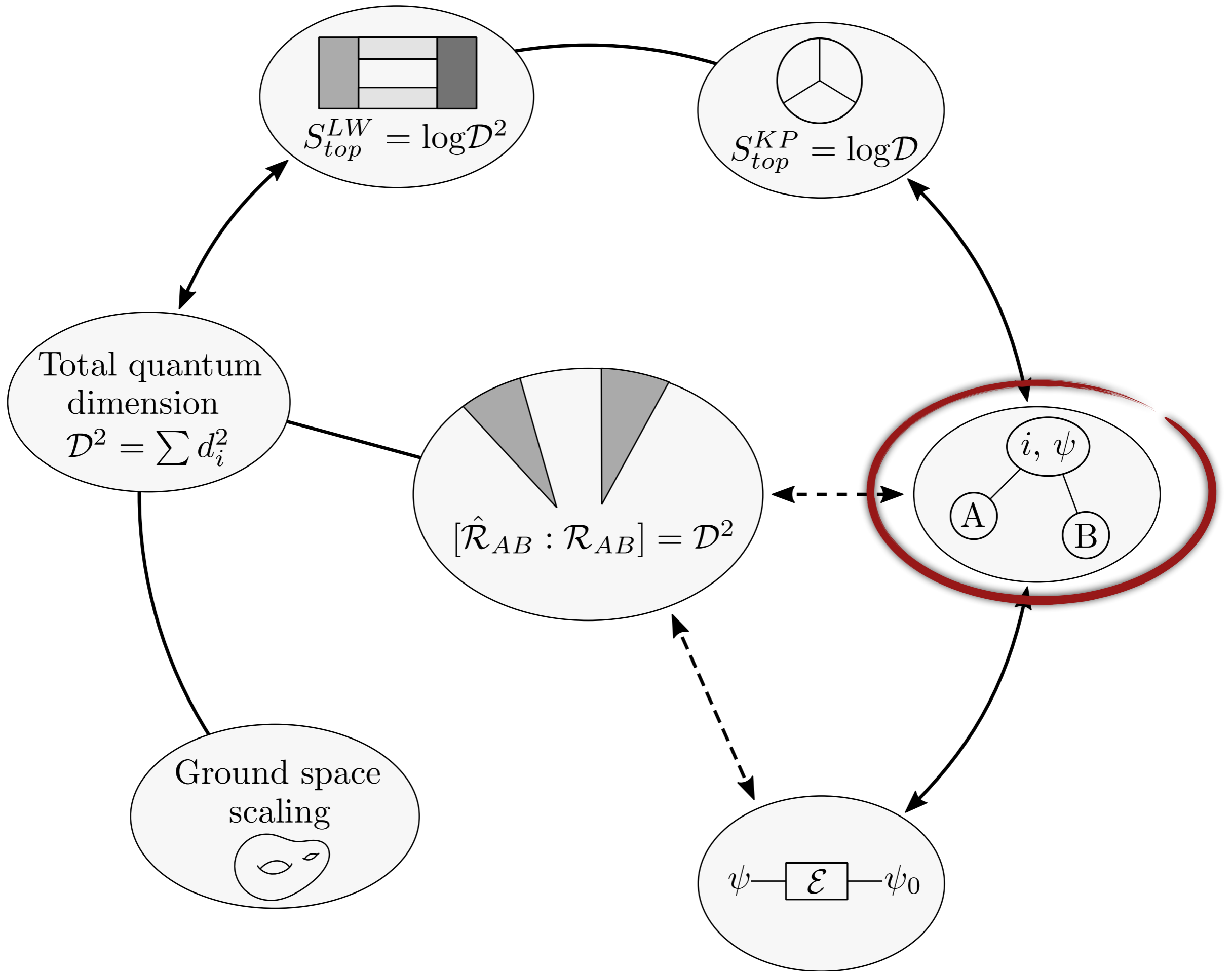
Eve

and can be used to create charge pairs

# A data hiding task

Similar conclusion: TEE  
as a **secret sharing capacity**

Kato, Furrer & Muraio, Phys. Rev. A., '16



# Distinguishing states

Alice prepares a mixed state  $\rho$ :

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Can Bob recover  $\{p_i\}$ ?



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Generalisation of **Shannon information**

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In general not exactly:

$$\begin{aligned}\chi(\{p_i\}, \{\rho_i\}) &:= S(\rho) - \sum_i p_i S(\rho_i) \\ &= \sum_i p_i S(\rho_i, \rho)\end{aligned}$$

Generalisation of **Shannon information**

# Optimal strategy

Want to compare  $\hat{\mathcal{R}}$  and  $\mathcal{R}$ :

$$H_\phi(\hat{\mathcal{R}}|\mathcal{R}) = \sup_{(\phi_i)} \left( \sum_i [S(p_i\phi_i, \phi) - S(p_i\phi_i \upharpoonright \mathcal{R}, \phi \upharpoonright \mathcal{R})] \right)$$

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$\Delta_\chi$

Shirokov & Holevo, arXiv:1608.02203



# A quantum channel

For finite index inclusion  $\mathcal{R} \subset \hat{\mathcal{R}}$

$$\mathcal{E} : \hat{\mathcal{R}} \rightarrow \mathcal{R}, \quad \mathcal{E}(X^*X) \geq \frac{1}{[\hat{\mathcal{R}} : \mathcal{R}]} X^*X$$

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quantum channel, describes the  
**restriction** of operations

# Quantum dimension and entropy

$$\log[\widehat{\mathcal{R}} : \mathcal{R}] = \sup_{\phi: \phi \circ \mathcal{E} = \phi} H_{\phi}(\widehat{\mathcal{R}} | \mathcal{R})$$

Hiai, J. Operator Theory, '90; J. Math. Soc. Japan, '91

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Completely different methods from Kato/Furrer/  
Murao, PRA **93**, 022317 (2016)

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- > Different methods compared to Kato et al.
- > No finite dimensional analogue to index
- > Can use powerful methods from mathematics
- > Right framework to study stability?

