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# A Validated BEM Model to Analyse Hydrodynamic Loading on Tidal Stream Turbine Blades

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Abstract— This paper details a Blade Element Momentum (BEM) model for a 3 bladed, horizontal axis Tidal Stream Turbine (TST). The code capabilities are tested and validated by applying a range of different turbine parameters and operating conditions, where results are compared to numerous validation datasets.

The model shows excellent agreement to performance and thrust measurements for 3 of the 4 datasets. Additionally, compared to other BEM models, improved correlations are seen at high rotational speeds. The fourth case shows over predictions of around 30% in power at peak operating speed. In this case, CFD studies show better correlation due to the ability to capture detailed flow features around the blade as well as free surface effects, however require 3 to 4 orders of magnitude greater computational cost.

Steady, non-uniform inflow functionality is incorporated into the model, where distributions of thrust and torque along the blade as well as cyclic loads are determined. These show the potential of the model to be used in combination with tools such as stress and fatigue analyses to improve the blade design process.

Keywords— Tidal Stream Turbine (TST), Blade Element Momentum (BEM), performance modelling, non-uniform inflow, blade cyclic loading, hydrodynamic loading

# I. INTRODUCTION

TST technology is currently at commercial scale array deployment phase, with EDF involved in the installation and grid connection of two 2MW rated OpenHydro devices in Brittany, Northern France in 2016.

Improvements in numerical modelling techniques have enabled the analysis of TSTs using complex CFD simulations. The modelling capabilities include performing detailed assessments of performance, dynamic loading, fluid/structure interactions and wake formation to a high degree of accuracy, however this comes at the price of high computational cost and long processing times.

The BEM model is a simple but effective and well understood method of predicting turbine performance and rotor thrust, commonly used in the wind industry and has been more recently adapted for tidal applications with commercial [1] and research-led [2] models. Benefits include significantly reduced

running times and low computational intensity, making BEM models more suited to applications requiring multiple, iterative engineering assessments or when access to high grade computational resources is restricted.

The aim of this paper is firstly to present results of extensive testing of a BEM model developed for TSTs with 3 different scale model turbines, providing validations with experimental measurements and other numerical models. Secondly, the paper aims to demonstrate the effects of non-uniform inflow on blade cyclic loading, and how these are predicted using the model.

The remainder of the paper is structured into 5 main sections: the theory behind the BEM model (II); definition of model inputs (III); results and validation (IV); discussion of results (V); conclusion (VI) and further work (VII).

# II. THE BEM METHODOLOGY

## A. Blade Element Momentum Theory (BEMT)

One dimensional momentum theory models the turbine as an infinitely thin, semipermeable actuator disc bounded by a streamtube (Figure 1a), where flow velocities and pressures at different positions can be related using conservation of mass and Bernoulli's equations. The axial force (thrust) on the disc can then be derived from the change in momentum and pressure differential across the disc:

$$dT = 4\pi \rho U_0^2 a (1-a) r dr$$

Where  $a = \frac{U_0 - U_d}{U_0} = \frac{U_0 - U_\infty}{2U_0}$  is the axial induction factor, and dT is the element thrust (N),  $\rho$  the fluid density (kg m<sup>-3</sup>),  $U_0$  the reference velocity (m s<sup>-1</sup>), r the local element mean radius (m) and dr the local element length (m).

Rotational momentum is gained by the flow in the wake which can be equated to the torque transmitted to the rotor. As this is a function of tangential velocity, the disc is split into a number of annular rings, where torque applied to each ring is expressed as:

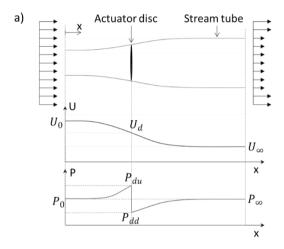
$$dQ = 4\pi \rho a' \Omega U_0 (1 - a) r^3 dr$$

Where  $a' = \frac{\omega}{2\Omega}$  is the tangential induction factor and, dQ is the element torque (N m),  $\omega$  the angular velocity of the wake (rad s<sup>-1</sup>) and  $\Omega$  the angular velocity of the turbine (rad s<sup>-1</sup>).

Blade element theory splits the blade into a number of discrete two-dimensional aerofoil sections, where radial interactions are neglected. Thrust and force causing torque can be resolved as a function of the aerodynamic forces (Figure 1b) and inflow angle using axial and tangential flow velocities:

$$dT = \frac{1}{2}\rho W^2 Bc(C_L \cos \phi + C_D \sin \phi) dr$$
$$dQ = \frac{1}{2}\rho W^2 Bc(C_L \sin \phi - C_D \cos \phi) r dr$$

Where W is the resultant fluid velocity (m s<sup>-1</sup>), B the number of blades, c the blade chord (m),  $C_L$  and  $C_D$  the lift and drag coefficients respectively and  $\phi$  the inflow angle (°).



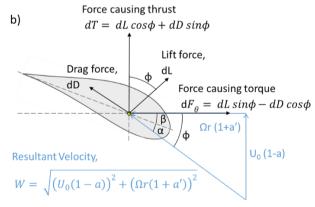


Figure 1 a) velocity and pressure distribution of flow in a streamtube b) flow velocities and aerodynamic forces on a blade element

BEM is a combination of these theories, where it is assumed that the change in momentum is solely accountable from the aerodynamic forces on the blade elements. Coefficients of power  $(C_P)$  and trust  $(C_T)$  on the rotor are calculated for a range of tip speed ratio (TSR), defined as:

$$C_T = rac{\sum_{rhub}^R dT}{rac{1}{2} 
ho A U_0^2}$$
 ,  $C_P = rac{\sum_{rhub}^R dQ \Omega}{rac{1}{2} 
ho A U_0^3}$  ,  $TSR = rac{\Omega R}{U_0}$ 

Where  $A = \pi R^2$  is the swept area of the disc (m<sup>2</sup>).

#### B. Tip and hub losses

The reduction in hydrodynamic efficiency at the blade tips and root due to radial flow are not modelled due to 2D flow assumptions. A correction factor (*F*) is therefore introduced as devised by Glauert, taking Prandtl's approximation of a helical wake as a succession of discs travelling at a velocity between the wake and free stream [3], incorporated into the baseline equations as:

$$F = F_{tip} F_{hub}$$
 Where: 
$$F_{tip} = \frac{2}{\pi} cos^{-1} e^{-\frac{B}{2} \frac{R-r}{R} \frac{1}{sin\phi}}$$
 
$$F_{root} = \frac{2}{\pi} cos^{-1} e^{-\frac{B}{2} \frac{r-r_h}{R} \frac{1}{sin\phi}}$$

#### C. High loading conditions

At high axial induction factors, thrust forces are under predicted due to the streamtube model assuming no interactions with the control volume, leading to unphysical reversal of flow predictions in the wake. In reality, turbulent mixing with the freestream occurs, injecting momentum into the slow moving fluid in the wake. Experiments on flat plates by Glauert shows much higher thrust forces than BEMT at axial induction factors above 0.4, with various best fit lines proposed [3], [4] as seen in Figure 2. When combined with the tip/hub loss factor, a numerical instability occurs due to the gap at transition to the highly loaded regime. Buhl devised a solution to overcome this yielding a smooth transition from the Glauert parabola to the BEMT prediction [5], [6]. He has shown reasonable agreement with experimental data and a fixed boundary condition at a=1 which is analogous to solid flat plate which fully impedes flow. These are set as conditions such that:

When 
$$a \le 0.4$$
:  $C_T = 4Fa(1-a)$ 

When  $a > 0.4$ :  $C_T = \frac{8}{9} + \left(4F - \frac{40}{9}\right)a + \left(\frac{50}{9} - 4F\right)a^2$ 

$$\begin{array}{c} 2.0 \\ 1.5 \\ \hline 1.0 \\ \hline 0.5 \\ \hline 0.0 \\ \hline \end{array}$$

$$\begin{array}{c} 2.0 \\ \hline 1.5 \\ \hline \end{array}$$

$$\begin{array}{c} 0.5 \\ \hline 0.0 \\ \hline \end{array}$$

$$\begin{array}{c} 0.2 \\ \hline 0.4 \\ \hline \end{array}$$

$$\begin{array}{c} 0.5 \\ \hline 0.0 \\ \hline \end{array}$$

$$\begin{array}{c} 0.5 \\ \hline \end{array}$$

Figure 2 Comparison of BEMT predictions of thrust at full range of axial induction with experiments and semi-empirical factors with no tip/hub loss F=1 (left) and F=0.8 (right)

## D. Non-uniform inflow velocity

In order to account for non-uniform inflow velocities, the axial velocity component must be calculated as a function of the relative blade element position in the channel. The BEM model is optimised to calculate the height above the channel base for each element at each azimuth. The resultant velocity is determined using the local axial flow velocity of the element, and used within the BEM loop to calculate the corresponding thrust and torque.

#### E. Blockage correction

The presence of the channel walls in the experiments constrain the flow, resulting in higher velocities around the rotor and a restriction in the wake expansion. This leads to higher forces and power output than seen in 'open water' conditions. These effects have been investigated by [7], developing a one dimensional analysis of flow in a stream tube between two rigid surfaces. This is in agreement with equations presented in [8], where an iterative procedure is proposed to determine correction factors converting bounded flow to 'equivalent open water' values. Extensions of this work include accounting for deformation of a free surface behind the turbine seen in experiments [9].

Experimental results for Cases 1 and 2 presented in their blockage corrected form by the authors, whereas Case 3 recommends application of correction factors to the experimental measurements [10] based on the following:

$$TSR_f = TSR\left(\frac{U_0}{U_f}\right) \quad C_{Tf} = C_T \left(\frac{U_0}{U_f}\right)^2 \quad C_{Pf} = C_P \left(\frac{U_0}{U_f}\right)^3$$

Where denoted values are the equivalent open water, and  $\left(\frac{U_0}{U_f}\right)$ =0.94 is the ratio of bounded flow to equivalent open water velocity found using the iterative procedure defined in [8].

# F. Numerical model

A model written in Python programming language is developed to couple the blade element and momentum theory equations, and incorporates an iterative loop to solve the axial and tangential induction factors. User specified inputs include turbine geometry, flow properties and aerodynamic coefficients of blade elements are read into the code, and post processing capabilities are developed to present turbine performance and thrust curves, as well as individual blade load distributions.

#### III. INPUT DATA

# A. Experimental Parameters

Four datasets are used to validate the model, based on three different scale model experiments for a 3 bladed horizontal axis turbine, performed under different experimental conditions, summarised in Table I. Case 1 (shown in Figure 3) is performed in a cavitation tunnel and therefore is not influenced by free surface effects. Case 2 is the most recent test to the athor's knowledge, performed in a very wide channel (width 5 times the height) and therefore benefits from a low blockage.

TABLE I EXPERIMENTAL PARAMETERS FOR 3 VALIDATION CASES

	1.a	1.b	2.	3.	units
Scale	1/20th		1/60th	1/30th	
Radius	0.4		0.135	0.3	m
Tank	Cavitation tunnel		Flume	Flume	
Velocity	1.30	1.73	0.46	0.55	ms <sup>-1</sup>
Pitch	12	5	0	0	۰
Aerofoil	NACA638 12-24		Gottingen804	NACA4415	
Blockage	17		2.5	19	%
Ref	[8], [11]		[12]	[10], [13]	



Figure 3 Experimental setup of 1/20th scale model TST in a cavitation tunnel for cases 1a and 1b [8]

#### B. Blade Parameters

Figure 4 and Figure 5 show the chord and twist distributions along the blades respectively, for each of the experimental cases. For Case 3, the blade was designed with large chord lengths in order to increase the Reynolds numbers past the transitional region, as justified by [14].

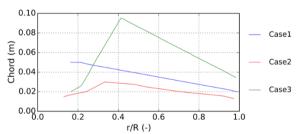


Figure 4 Chord distribution along normalised blade length

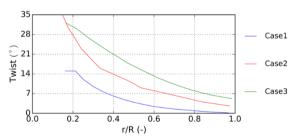


Figure 5 Twist distribution along normalised blade length

#### C. Inflow conditions

The model is optimised to assess non-uniform inflow conditions in the form of steady, 'frozen' velocity profiles. Boundary layer effects due to friction on the channel base can be approximated as a shear profile, where a 1/7<sup>th</sup> power law shows excellent agreement with experimental measurements in a flume shown in Figure 6. This is then used to assess the level of load fluctuations experienced by the blades.

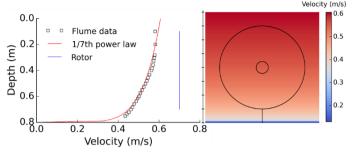


Figure 6 Shear inflow profile with 1/7th power law compared with flume measurements (left) and front contour view (right)

#### D. Aerofoil Coefficients

Aerofoil coefficients are used to calculate the aerodynamic forces within the blade element section of the model. Numerous data sources are available [15] and for Case 2, coefficients are provided based on catalogued data [16], at Reynolds numbers between 20,000 - 30,000.

For Cases 1 and 3, Lift and drag coefficients are generated using XFOIL [4], a linear vorticity function panel method with viscous boundary layer and wake model. Chord based Reynolds numbers for each case were calculated using the following equation:

$$Re_{ch} = \frac{\rho CW}{\mu}$$

Where  $W = \sqrt{U_0^2 + (\Omega r)^2}$  is the resultant velocity at the aerofoil (ms<sup>-1</sup>) and  $\mu$  is the dynamic viscosity (Ns m<sup>-2</sup>). In order to reduce the number of cumbersome analyses for many Reynolds numbers, the rotational velocity was taken at the optimal performance of each turbine, and the local radius taken at 75% of the blade length.

2D static wind tunnel measurements or XFOIL data do not take into account the complex 3D interactions of flow over a rotating blade. The effects of radial forces in rotating foils induces a Coriolis force acting in the direction of the trailing edge, effectively delaying the onset of boundary layer separation. The delayed stall effects can be accounted by applying a Du-Selig model to the and lift [17], and Eggers adjustment to the drag [18]. These are dependent on the local element radius and chord, and therefore varies along the blade length.

Due to the high range of inflow angles to be analysed, angles of attack are able to exceed the point of stall, which is beyond the recommended capabilities of XFOIL. Post stall coefficients of lift and drag are therefore generated using a Viterna extrapolation function [5].

Coefficients against angle of attack at different points along the rotor radius for the Cases 1 and 3 are presented in Figure 7.

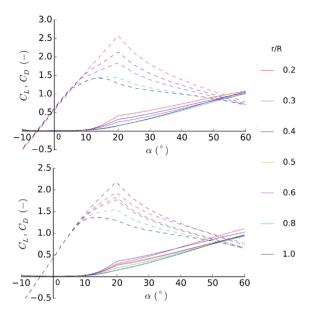


Figure 7 Blade local lift (dashed lines) and drag (solid lines) curves for Case 1 (top): NACA638(12-24), Re=3.0E+05 and Case 3 (bottom): NACA4415, Re=1.5E+05

#### IV. RESULTS

#### A. BEM correction factors

As the BEM model requires correction factors to account for where physical effects are not captured, it is useful to determine the extent at which these are applied for assessments in their influence on the overall thrust and power predictions.

Figure 8 shows the axial induction factor and tip/hub loss correction factors calculated by the model for case 1b. From this, one can see that the model enters the 'highly loaded regime', based on the semi-empirical Buhl factor only at the tip and becomes more evident at higher TSRs. The tip/hub loss correction is significant, indicating large hydrodynamic losses particularly at low TSRs.

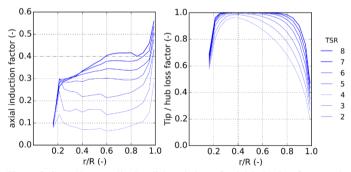


Figure 8 Case 1b normalised radial variation of axial induction factor and tip/hub loss correction factor at varying TSR

#### B. Validation with Coefficients of Power and Thrust

Validations are made with outputs of non-dimensional power and rotor thrust with varying rotational velocities. Figure 9 to Figure 12 shows the comparison of experimental measurements with output from the BEM model in this study as well as other studies within the field of TST modelling.

Figure 9 shows that the present BEM model has excellent agreement with the 1/20<sup>th</sup> scale experimental data [8], and in in line with a study using a research-led BEM model 'SERG' developed at Southampton University [11]. At the higher inflow velocity shown in Figure 10, there is again a strong correlation with measurements, however both models over predict power at high TSRs (>8) of up to 23%.

Figure 11 shows the BEM model performs very well with  $1/60^{th}$  scale parameters, showing again excellent agreement with experiments [12]. The high levels of thrust seen at TSRs >4 is captured by the model, whereas a study using commercial software Tidal Bladed seems to show an under prediction of up to 10%.

Figure 12 shows reasonable agreement with TSRs <3, however for higher TSRs, the model over predicts power and thrust. Results from CFD studies using a Hybrid RANS-BEM (Reynolds Average Navier Stokes) model, as well as a fully blade resolved RANS case [19] both show better agreement, however there is still an over prediction seen at TSRs >3. It was also noted that unlike previous cases, the numerical models also do not capture the optimal TSR where peak performance occurs. At peak TSR, BEM over predicts power by 28% and thrust by 20%.

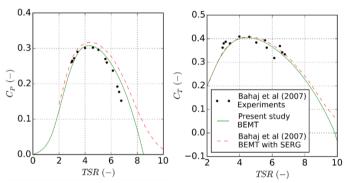


Figure 9 Case 1a validation with 1/20th scale experiments at 1.3ms<sup>-1</sup> with 12° pitch and SERG BEM model

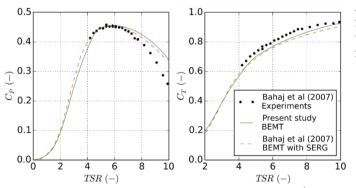


Figure 10 Case 1b validation with 1/20th scale experiments at 1.73ms<sup>-1</sup> with 5° blade pitch and SERG BEM model

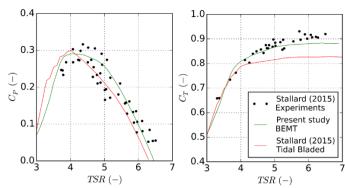


Figure 11 Case 2 validation with 1/60th scale experiments a study using Tidal Bladed

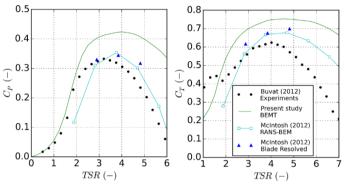


Figure 12 Case 3 validation with  $1/30^{th}$  scale experiments and CFD studies using RANS-BEM and RANS blade resolved cases

#### C. Blade cyclic thrust loads

The average and cyclic thrusts experienced by each blade as the turbine passes through each rotation can also be determined by the model. An example for Case 1b (1/20<sup>th</sup> scale, 1.73ms<sup>-1</sup>, 5° blade pitch) is shown in Figure 13. At higher velocities, the blade experiences higher thrusts, therefore the peak occurs at top dead centre. It can be seen that as the TSR increases, the average thrust increases, as to be expected from the C<sub>T</sub> curves. Additionally, it can also be seen that the variation of thrust forces also increases. At the optimal performance rotational velocity of TSR = 6 from Figure 10, the blade experiences fluctuations of around 20% around the rotor average.

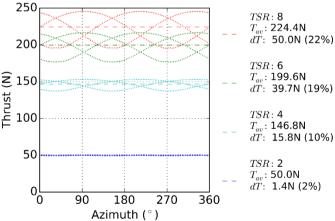


Figure 13 Case 1b at 1.73ms<sup>-1</sup> with 5° blade pitch average and cyclic thrust loads during one turbine rotation for various rotational velocities

#### D. Non-uniform blade loading

Additionally to the thrust levels for individual blades, the model also outputs the distribution of these thrusts along the blade length.

Figure 14 shows the thrust distribution with normalised rotor radius through one turbine rotation at the optimal performance rotational velocity. The peak again occurs at top dead centre, and it is seen that maximum values are seen close to the blade root.

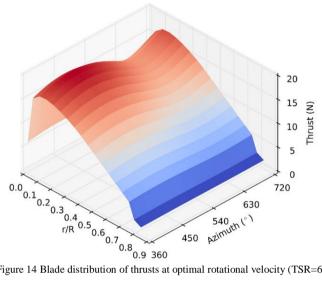


Figure 14 Blade distribution of thrusts at optimal rotational velocity (TSR=6)

#### V. DISCUSSION

#### A. Validation

C<sub>P</sub> and C<sub>T</sub> curves from the model have excellent agreement with measurements for Cases 1 and 2, showing the methodology is well applied to these types of experiments. The strong correlations with other BEM models also show that the theory is well implemented in the code.

Over predictions in power, thrust and peak TSR are seen within Case 3 results at TSRs > 3, which are also seen (although less significantly) in the CFD studies. This is thought to be a result of a combination of causes:

- Numerical errors in the generation of aero foil coefficients;
- Experimental errors in measurements;
- Effects of high blockage correction and high widthwise velocity distributions seen across the channel;
- Lack of consideration into the presence of a free surface (CFD uses a volume of fluid free surface model).

Methods to combat these have been explored, including: CFD studies of 2D static aerofoils have been tested on NACA profiles, and show improved correlation to experimental data. This methodology may lower uncertainty in aerofoil coefficients, however are time intensive to generate. A free surface model developed by [9] could be applied to the model,

however difficulties have been identified when used in combination with the Buhl highly loaded correction factor.

To generate the C<sub>P</sub> and C<sub>T</sub> curves for each case, the present BEMT model took 3 minutes on a single quad core processor, equivalent to 0.05 CPU-hours. In comparison, CFD studies of Case 3 reported a requirement of 12 CPU-hours for the RANS-BEM model, and 100 CPU-hours per turbine rotation of the RANS fully blade resolved model [20]. This indicates the significant computational savings of the present study, and hence its advantages in cases where processing time and ability to run on local machines outweighs the requirement for complex, high detailed simulations with wake characterisation.

#### B. Blade loading

When a uniform inflow velocity is applied, blades experience a uniform thrust down its length. When nonuniform profiles are applied, such as in this case where a shear profile to account for boundary layer effects of the channel base, loading fluctuates as a function of azimuth. High fluctuations of almost 20% at the optimal rotational speed in Case 1b indicates the impact of just one component of the flow on loading conditions. Other inflow profiles could be applied to model actual flow regimes seen in the field, in order to gain cyclic loading information which could be used in a fatigue analysis.

Further assessment of the thrust forces shows that the blade is subjected to non-uniform loading along its length, which are not accounted for when considering average loads. Results are dependent on flow and operating conditions, however results from Case 1b at optimal TSR shows a peak occurs at a specific location along the blade. This shows the potential of using this tool as a methodology for attaining 'hot spots' where loads are concentrated, as well as generally more detailed loading information that can be used in Finite Element Analyses (FEA) to assess stress distributions. The resolution can be improved by splitting the blade into more elements.

It should be noted that this model considers steady, 'frozen' profiles, and is not optimised to consider dynamic inflows, which would be required to assess effects such as turbulence or waves. A time stepping function could be used to apply a series of different inflow profiles over time, however additional functions would need to be applied in order to consider dynamic effects such as inertia.

### VI. CONCLUSIONS

A BEM model developed in this study is applied to numerous scale models of TST and validated with experimental measurements and numerical models. Three out of four cases show excellent agreement with experimental data, and in line with other BEM models. The exception case, over predictions of power and thrust are seen due to errors in aerofoil coefficients, blockage and free surface effects, which are captured better by CFD studies. Despite this, the BEM model ran with 2 to 3 orders of magnitude less computational cost than CFD, indicating its advantage for applications where multiple iterative engineering assessments are made using low grade computational resources.

The model is then applied to assess the thrust and torque distributions along the blades, which can be used as more detailed inputs to FEA analyses as well as to determine localised 'hotspots' where stresses are concentrated. Additionally, fluctuating thrust forces as the turbine rotates through a non-uniform inflow profile are calculated, which can be used to determine cyclic loading conditions for a fatigue analysis. These indicate the potential of using this tool for improving the overall design process of turbine blades.

#### VII. FURTHER WORK

A methodology for attaining aerodynamic coefficients using CFD studies of static aerofoils in 2D flow are to be further developed in order to reduce uncertainty of the BEM inputs. This methodology will then be applied to 'flat-plate' aerofoil profiles, such as those seen in some full scale TSTs designs.

The model is currently under adaptation to perform assessments of the OpenHydro device, as high solidity, hubless and ducted turbine. Significant differences in design will be accounted for by using results from CFD studies and validations with a full scale deployment of a 2MW turbine at the Paimpol-Bréhat site in Jan 2016.

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