NUMERICAL PREDICTION OF UNDRAINED RESPONSE OF PLATE ANCHORS UNDER COMBINED TRANSLATION AND TORSION Hamidreza Nouri¹, Giovanna Biscontin² and Charles P. Aubeny, M. ASCE³

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5 ABSTRACT

6 The undrained pure translational and torsional capacity of anchors and the plate response un-7 der eccentric translational loading is investigated in this paper using three-dimensional finite 8 element (3D-FE) analysis. Plastic limit analysis is adopted to establish a benchmark solution 9 for ultimate translational resistance with satisfactory agreement with the FE values which 10 confirms the developed numerical model. Although plate thickness has a marked impact on 11 the maximum shear and torsion resistance, the shape of failure envelope is minimally affected 12 by thickness. A simple three-degree-of-freedom interaction equation is curve-fitted to FE 13 failure datapoints. Representative interaction relationships are introduced for square and rec-14 tangular plates.

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16 KEY WORDS: plate anchor; torsional capacity; translational (shear) resistance; plastic limit

17 analysis; finite element analysis; failure envelope.

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18 **1- Introduction**

19 The offshore industry is expanding farther into very deep water, leading to a number of economi-20 cal and technical challenges. Floating structures anchored to the seabed using catenary and taut-21 wire mooring systems are generally more technically feasible and cost effective than gravity based 22 platforms in these deep water environments. Anchorage for floating systems can include piles, 23 caissons and various types of plate anchors. This paper addresses performance of the latter alterna-24 tive. Under normal operating conditions, plate anchors are mainly subjected to pull-out loading in 25 an intended plane of loading, for example, the X-Z plane in Fig. 1(a). However, extreme storm events may cause partial failure with the loss of one or several mooring lines resulting in substan-26 27 tial changes in the orientation of the intact lines, as well as to the applied resultant force [Fig. 28 1(b)]. Under these conditions the anchor must resist a general system of loads, with some compo-29 nents of load acting outside the plane of intended loading.

30 A full description of behavior of an anchor plate in a partially failed mooring system requires ade-31 quate understanding of undrained response of anchor plates subjected to general loading condi-32 tions involving six degrees of freedom. This entails characterizing anchor behavior for six components of uniaxial loading (3 translational and 3 rotational) as well as the interaction effects 33 34 amongst all possible load combinations. Much of the research in this area to date has focused on 35 in-plane loading conditions involving forces in the X and Z directions and moments about the Y-36 axis (Fig. 1) (O'Neill et al. 2003, Murff et al. 2005). This paper investigates anchor behavior for 37 loading in the X-Y plane in Fig. 1, which will be termed 'shear-torsional' loading. The loading 38 comprises 'shear' forces H_x and H_y and a torsion T; i.e., a moment acting about the Z-axis. All 39 loads act within the plane of the plate. The study presented herein does not attempt to fully charac-40 terize anchor behavior for general 6 degree of freedom loading conditions. Rather, it investigates a 41 critical component for a more general model to be developed in the future.

The plate anchors considered in this paper are idealized as rectangular plates with specified aspect ratios (*W/L*) and thickness t (Fig. 1). Some plate anchors, such as suction embedded plate anchors (SEPLAs) match this idealized geometry to a reasonable degree. By contrast, drag embedded plate anchors usually have relatively complex geometric configurations to facilitate embedment. The findings of this paper can provide useful insights into the behavior of drag anchors, but the inherent limitations of the analyses arising from the simplified geometry must be carefully considered.

48 O'Neill et al. (2003) used plane strain finite element (2D-FE) analysis to investigate the behavior 49 of rectangular and wedge-shaped strip anchors subjected to combined in-plane translational, verti-50 cal, and rotational loading. They also developed plastic upper bound solutions to evaluate pure 51 parallel, normal, and rotational capacity factors for strip plates to validate the FE results. Using 52 two-dimensional finite element (2D-FE) analysis, they produced failure loci and plastic potentials 53 from considerations of fluke-soil interaction and then developed a design method to predict drag 54 anchor trajectory during embedment. Murff et al. (2005) used 2D-FE to develop a failure locus for the ratio of plate length to thickness (L/t) of 7 and 20 with different plate roughness. Yang et al. 55 56 (2010) employed 3D-FE, to study the behavior of infinitely thin plate anchors subjected to loading 57 in all six degrees of freedom. They also introduced a plasticity solution for failure envelope of a 58 plate under combined translation-torsion and developed formulas to estimate pure sliding and tor-59 sional bearing capacity. This plasticity solution predicts the behavior of an infinitely thin plate, but 60 provided only an approximate solution for a plate with finite thickness. Nouri (2013) modified the 61 plastic limit solution originally developed by Yang et al. (2010) to predict more accurately the tor-62 sional collapse load and failure envelopes of a plate anchor with finite thickness. Researchers have 63 also conducted 3D-FE analysis on other types of offshore foundations under general loading con-64 dition, such as for mudmats (Feng et al. 2014) or OMNI-max anchor (Wei et al. 2015). In the current study, we use 3D-FE analysis to evaluate the translational (parallel to plate) and tor-65

66 sional ultimate bearing capacity of plates, as well as the interaction of these two modes of loading.

Results of the FE analysis are also used to validate analytical solutions to predict bearing capacities and to confirm whether the postulated failure mechanism for the analytical solution is correct.

69 **2- Material Properties for Numerical Model**

70 The commercial software ABAQUS (HKS, 2008) is used in this study. The plate is assumed to be 71 deeply embedded, with localized plastic flow forming around the plate anchor and not extending 72 to the surface, resulting in capacity factors that are not affected by overburden and soil weight 73 (Song et al. 2008; Wang et al. 2010). Thus, the soil is assumed to be a weightless material without 74 loss of accuracy. Since the size of the plastic zone around the plate anchor at failure is smaller for 75 shear and torsion than for uplift and rotation, the assumption of deep embedment is valid even for 76 shallower embedment depths. This assumption also implies that the soil and plate are fully bond-77 ed. Thus, "no separation" is assumed as the normal (i.e. perpendicular to interface) contact behav-78 ior for the plate-soil interface in conjunction with "rough" for the tangential contact property to 79 realistically simulate a fully bonded condition. The interaction between soil and plate is modeled 80 using surface to surface contact pairs in which the plate outer surface is chosen as a "master sur-81 face" and the soil surface in contact with the plate as a "slave surface". The active degrees of free-82 dom of the nodes on the slave surface are constrained to the master surface nodes through rela-83 tionships which define the tangential and normal interaction at the nodes of these two surfaces.

84 Clay under undrained conditions is modeled as a linear elastic perfectly plastic material, with 85 yielding determined by the Von Mises failure criterion with undrained shear strength, s_{u} . The Young's modulus of the soil, E, is given by a modulus ratio of $E/s_u = 500$, and Poisson's ratio is 86 87 taken as 0.49 to simulate no volume change for undrained clay in total stress analysis. The ulti-88 mate capacity of the plate is not affected by the pre-failure elastic behavior of the soil (Chen and 89 Liu 1990). In addition, in studies on plate uplift bearing capacity it was observed that as the soil 90 rigidity (E/s_u) increases, the anchor displacement required to mobilize the maximum capacity re-91 duces, while the maximum capacity does not change (Song et al. 2008; Wang et al. 2010). There-

92 fore, the Young's modulus of the soil is assumed constant in these studies. Since the analysis of 93 the structural behavior of the anchor is not within the scope of this work, the plate is modeled as a rigid body with Young's modulus 10¹⁰ times that of the soil, and Poisson's ratio of 0.30. The 94 analysis is conducted using small strain displacement control, since the ultimate translational and 95 96 torsional capacities are developed within an anchor movement of 0.1 times the anchor length. A 97 reference node is defined in the center of the plate and prescribed displacements are applied at the 98 reference point to perform displacement controlled analysis. Standard boundary conditions are al-99 so applied to the model: the base is fixed in all directions, while the vertical boundaries are fixed 100 in the lateral direction and free to move vertically. In all the models with a plane of symmetry, the 101 nodes on the plane of symmetry are only allowed to move in the plane.

102 **3- Challenges of the Numerical Modeling**

Unrealistic stress distributions arise from the displacement discontinuities that take place at the edges of a rigid structure penetrating a softer material (Van Langen and Vermeer 1991) resulting in higher overall resistance. The unfavorable effect on the numerical results is more severe for buried structures and materials with no volume change. There are mainly two approaches to overcome such loss of numerical accuracy:

108 1- Mesh refinement

109 2- Contact pairs along the corners of the plate

As plane strain 2D analyses take much less time and fewer memory resources, they are a reliable and fast way to evaluate possible solutions for this problem. 2D-FE analyses can also be adopted to develop an optimal mesh size and to verify properties and length of the contact pairs in soil along the plate corners through sensitivity analysis. The plate anchor is idealized as a strip with length L and thickness t in 2D space.

115 **3-1-** Mesh refinement around the sharp edges

116 As the detrimental effect of the stress concentration is more pronounced in thinner plates, a plate 117 with thickness of t = L/20, is considered for the preliminary 2D analysis to calculate the ultimate 118 translational (parallel to plate) bearing capacity of the strip plate. Several 2D models with coarse, 119 fine, and very fine mesh with first order 4-node plane strain quadrilateral hybrid elements, 120 CPE4H, are created. All the soil elements are fully integrated hybrid type to accommodate the 121 constant volume constraint in undrained conditions. The whole model dimensions are $9L \times 9L$, in 122 which L is the length of the plate. The model is divided into several zones with different mesh 123 densities to minimize the degrees of freedom of the model and avoid unnecessary computational 124 effort. The finest zone extending to t = L/20 around the plate consists of square elements with di-125 mension of L/80, L/160, L/320 for the coarse, fine, and very fine meshes, respectively. A fine 126 mesh with second order hybrid square elements (CPE8H) with dimension of $L/160 \times L/160$ is also 127 created to assess the performance of higher order elements in predicting ultimate bearing capacity. 128 Only a small difference can be observed in the maximum parallel bearing capacity of the plate for 129 the fine and very fine meshes in Fig. 2. The results from the fine mesh (elements of $L/160 \times L/160$) 130 with second order elements match exactly the values of parallel bearing capacity from the very fi-131 ne mesh with first order elements. Refining the mesh improves the results and overcomes the is-132 sues of stress concentration to some extent, but the bearing capacity values do not change signifi-133 cantly after some degree of mesh refinement. The ultimate shear capacities calculated using the 134 very fine mesh or higher order elements are still 13% higher than the value obtained from the up-135 per bound plastic limit analysis. In addition, a 3D fine mesh with elements of size $L/160 \times L/160$ 136 would require significant computational effort, making the numerical analysis practically impossi-137 ble to conduct. Thus, for a reasonable computational effort it seems that excessive mesh refine-138 ment fails to provide sufficiently accurate FE results for shear bearing capacity of plate anchors deeply embedded in undrained clay. 139

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- 3-2- Contact interfaces in soil along the corners of the plate

141 An alternative solution to enhance the stress distribution on the plate edges is using contact inter-142 faces near the corners of the plate. Rowe and Davis (1977) introduced potential rupture lines near 143 the edge of an anchor to ensure accurate modeling of the singularity at the anchor tip. This ap-144 proach attempted to overcome the inhibition of free plastic flow inherent in the usual stiffness 145 formulation of the FE method by permitting the formation of velocity discontinuities in the re-146 gions of high stress and velocity gradient near the tip of the anchor plate (Rowe and Davis, 1977). 147 Van Langen and Vermeer (1991) introduced zero-thickness interface elements to enhance the un-148 realistic stress non-uniformity on surfaces with singular points.

149 Using the same general concept, and taking advantage of *ABAQUS* capabilities in defining various 150 types of contact relationships, this study adopts interfaces placed vertically and horizontally in the 151 soil in the vicinity of plate corners to accommodate for vertical and horizontal displacement dis-152 continuities. Each interface is defined as a couple of master and slave surfaces in the soil that can 153 slide independently from each other according to a surface to surface contact relationship (contact 154 pairs). A coefficient of friction ($\mu = 0.001$) and a maximum contact shear stress equal to the un-155 drained shear strength of the soil ($\tau_{max} = s_u$) are adopted to simulate frictional behavior. No separa-156 tion between the master and slave surface nodes is also assumed in the direction normal to the in-157 terface.

As shown by O'Neill et al. (2003) and confirmed through the 2D-FE analysis, in-plane horizontal movement of the rigid plate creates a plastic zone in the soil around the plate that extends for a distance from the plate approximately equal to plate thickness, *t*. Contours of deformation in the *x* direction for a plate under pure sliding in the 2D-FE analysis shown in Fig. 3 also confirm the extents of this plastic zone. The 3D-FE analysis also shows a similar zone of thickness *t* to be affected by the torsional rotation (Fig. 4). The finest zone wrapping around the plate contributes the most to the number of elements, i.e. the computational effort in the model.

165 Thus, this region extends for a distance from the plate equal to the plate thickness, t, to minimize 166 the number of elements. The length of the contact is also adopted equal to plate thickness. The fi-167 nal 2D mesh and the contacts placed horizontally and vertically in the vicinity of the plate corners 168 are shown in Fig. 5. The final mesh is divided into six different zones with element size decreas-169 ing as the zones become closer to the plate. Mesh tie constraints are used to constrain each active 170 degree of freedom of the slave node on the finer zone surface to the master node on the coarser 171 zone surface. FE results for the ultimate parallel bearing capacity are compared with the plastic 172 limit analysis (PLA) value presented by O'Neill et al. (2003). For a plate with t = L/20 and 173 μ =0.001, the FE-derived value of the plane strain shear bearing capacity factor is 2.77, which ex-174 ceeds the PLA result, 2.75, by less than 1%. The effect of the interface coefficient of friction, μ , 175 on the ultimate shear capacity results is evaluated for coefficient values ranging from 0.001 to 176 1000 (Fig. 6). For higher values of friction coefficient the ultimate capacity increase to about 2.85 177 and remains constant for coefficient values larger than one $(\mu > 1)$, which shows a minor impact of 178 the friction coefficient on the results. Thus, the friction coefficient of μ =0.001 is adopted for the 179 interface in all the following analysis. FE results for the ultimate parallel bearing capacity also in-180 dicate that the mesh is adequately fine, contact length and properties are reasonable and the equiv-181 alent 3D mesh could give satisfactory results.

182 **4- 3D finite element analysis**

The 3D mesh is created with the same approach used for the 2D mesh (Fig. 7). The 3D model is created for anchors idealized as square and rectangular (W/L = 2.0) plates with thickness t = L/20, L/14, L/10, and L/7. The model dimensions are $9L \times 9L \times 5L$ and $9L \times 12L \times 5L$ for the square and rectangular plates, respectively. The 3D mesh consists of 8-node full integration 3D hybrid first order elements, C3D8H. Similarly to the 2D analyses, the Von Mises constitutive model with the same elastic-plastic properties is used for the soil. The contact definition of the 2D model is used at the plate corners in the soil and plate-soil interface. The length of the contact and the depth of the fine 190 mesh zone around the plate are equal to the plate thickness in all models (Fig. 4). Taking ad-191 vantage of the symmetric nature of shear-torsion loading of the plate anchor, only half of the FE 192 mesh is simulated. Standard boundary conditions are applied to the model boundaries and plane of 193 symmetry.

194 **4-1- Ultimate load capacity**

In-plane translational (parallel to plate) displacements are applied in the direction parallel to the smaller (*L*) and larger (*W*) sides of the plate (*x* and *y* directions respectively) and torsion is prescribed as the rotation of the plate about the *z* axis. The FE results for the pure shear and pure torsional loading are reported in Tables 1 and 2, respectively, for square and rectangular (*W*/*L* = 2) plates, in their normalized forms, $N_{sx,max}$, $N_{sy,max}$, and $N_{t,max}$, where:

$$N_{sx,\max} = \frac{H_{x,\max}}{LWs_u} \tag{1}$$

$$N_{sy,\max} = \frac{H_{y,\max}}{LWs_u} \tag{2}$$

$$N_{t,\max} = \frac{T_{\max}}{L^2 W s_u} \tag{3}$$

where H_x , and H_y , are the horizontal (shear) reaction force in *x* and *y* directions and *T* is the torsional reaction moment of the rigid anchor plate about the *z* axis.

The FE ultimate shear resistance is compared with benchmark solutions: (1) the resistance factor for an infinitely thin plate, which provides the minimum values of ultimate shear resistance $(N_{sx,max}=N_{sy,max}=2 \text{ for } H_{x,max}=H_{y,max}=2s_uWL)$; (2) A PLA formulation, developed by Yang et al. (2010) to evaluate the plate thickness impact on the ultimate shear capacity (Table 1):

206
$$N_{sx,max} = \frac{H_{x \max}}{s_u LW} = 2\alpha + 2\left[\alpha \frac{L}{W} + N_e\right] \frac{t}{L}$$
(1)

207
$$N_{sy,max} = \frac{H_{ymax}}{s_u LW} = 2\alpha + 2\left[\alpha + N_e \frac{L}{W}\right] \frac{t}{L}$$
(2)

where N_e is the simple plane strain bearing capacity factor, equal to 7.5 (O'Neill et al. 2003). The adhesion factor, α , for the fully bonded condition is assumed equal to unity. The above expression for pure *x*-shear capacity factor, $N_{sx,max}$, for a strip plate ($L/W \approx 0$) reduces to $N_{s,max}=2(1+N_e.t/L)$ as

211 proposed by O'Neill et al. (2003) for ultimate parallel plate capacity in 2D plane strain condition.

212 Whereas the PLA analyses under-predict FE shear capacity factors for plate thickness of t = L/20, 213 and L/14, they over-predict FE capacity factors for t = L/10 and L/7 for both square and rectangu-214 lar plates. This could be attributed to the fact that the shear bearing capacity of thicker plates is 215 less sensitive to stress concentration on plate corners because the ultimate shear bearing capacity 216 of thick plates is larger than that of thinner ones. Overall, the 1-3% difference between PLA and 217 FE results for ultimate shear capacity is minimal. The thickness has a marked impact on the plate 218 shear resistance, resulting in an approximately 40 to 120% increase for square plates of t=L/20 and 219 L/7 respectively, compared to the t=0 plate. The typical normalized load-displacement responses 220 for square and regular plates under in-plane y-shear are also depicted in Fig. 8.

Table 2 summarizes the FE derived maximum torsional resistance. Torsion-rotation plots are also included in Fig. 9. Yang et al. (2010) also proposed a PLA solution for pure torsion, which tends to over-predict resistance for plates of finite thickness (Nouri 2013), but their formulation for an infinitely thin plate ($N_{t,max0}$) provides a lower bound value for ultimate torsional resistance:

$$N_{t,\max 0} = T_{\max 0} / (s_u W L^2) =$$

$$225 \qquad \alpha \frac{(W/L)^2}{6} \left\{ \frac{\sin \theta_o}{\cos^2 \theta_o} + ln \left[\tan \left(\frac{\pi}{4} + \frac{\theta_o}{2} \right) \right] \right\} + \alpha \frac{L/W}{6} \left\{ \frac{\cos \theta_o}{\sin^2 \theta_o} - ln \left[\tan \left(\frac{\theta_o}{2} \right) \right] \right\}$$
(3)

where $\theta_o = \tan^{-1}(L/W)$. For square and rectangular plates this analytical solution yields 0.76 and 1.19, respectively. Torsional resistance also increases markedly with plate thickness by approximately 40 to 90% for *t*=*L*/20 and *L*/7, respectively, versus the zero thickness plate.

229 **4-2- Failure envelopes**

230 In order to construct the interaction curve or failure locus, two types of displacement control 231 methods are used: the swipe test and the probe test. In the swipe test approach failure locus is ob-232 tained by a single analysis in two separate steps. In the swipe test the plate is firstly displaced in 233 the direction of the degree of freedom (DoF) under examination from zero until the ultimate ca-234 pacity in that direction is reached. In the second step, the displacement is imposed in the second 235 DoF until the ultimate capacity in the new direction is fully developed. In the finite element analy-236 sis once the ultimate force is reached in the first step, there will be no more increase in the reac-237 tion force along that DoF, thus no further expansion in overall plate failure locus as the movement 238 of the plate progresses. The disadvantage of this method is that, due to the elasto-plastic "yield-239 ing" occurring within the "failure" locus, the swipe test tracks a load path marginally inside the 240 actual overall plate failure envelope (Bransby and Randolph 1998).

241 In the probe test or fixed displacement method, a single point on the failure envelope is identified 242 for each fixed ratio of the prescribed displacements. Therefore, the failure locus can be created by 243 conducting a number of finite element analyses with different displacement ratios. The prescribed 244 fixed displacement ratio gives rise to a load path beginning at the origin with an initial gradient 245 based on the elastic stiffness of the material. The path gradient decreases as approaching failure 246 envelope to follow the interaction curve and stops where there is no further increase in the forces 247 developed in each intended degree of freedom (Bransby and Randolph 1998). This method gives 248 an accurate trace of the failure envelope, but requires several FE analyses.

249

4-2-1- Combined in-plane shear and torsion

Fig. 10 shows examples of the swipe and probe failure envelopes for square and rectangular plates of t = L/7 subjected to in-plane y-translation and torsion $(N_{sy} - N_t)$. The rest of the analyzed cases are not included for brevity. Note that $du_y/d\phi$ is the prescribed fixed displacement ratio for each probe path, defined as the ratio of plate displacement increment in the y direction, to the increment of the plate in-plane rotation. The probe curve exceeds the shear-torsion (i.e. shear prescribed first, 255 torsion next) swipe envelope. This difference increases significantly if a torsion-shear swipe test 256 (i.e. torsion first, shear next) is carried out instead. In the torsion-shear swipe test, at the moment 257 of full development of the torsional ultimate capacity, the soil elements in regions close to the 258 middle of the plate and the plate sides are not completely plastic yet since the maximum dis-259 placements occur in the region next to the corners of the plate (See Fig. 4). Thus, the stress state of 260 the soil elements in the center is farther from failure at the end of the first stage of pure torsion. In 261 contrast, in the shear-torsion swipe test, all the soil surrounding the plate is fully plastic at the end 262 of full mobilization of the shear resistance. All the swipe tests follow the shear-torsion sequence. 263 As Fig. 10 suggests and as it was discussed earlier in the paper, the probe test approach offers 264 more accurate values since the swipe envelope travels although marginally but still inside the 265 probe curve. Therefore, probe test values are adopted and referred as FE results hereafter in the 266 paper.

267 The failure envelope (e.g. Fig. 10) can be described by two basic components:

268 1- Maximum plate resistance under pure shear ($H_{x,max}$, $H_{y,max}$) and torsion (T_{max}) which quanti-269 fies the size of the envelope (discussed in previous section).

270 2- The shape of the failure envelope defined by the shear or torsional resistance normalized 271 by the corresponding maximum value (e.g. $N_{sy}/N_{sy,max} - N_t/N_{t,max}$ or $H_y/H_{y,max} - T/T_{max}$).

Fig. 11 summarizes the FE probe results in the normalized space for all the studied cases. Interestingly, the plate thickness does not seem to affect the interaction behavior of the plate under oneway planar shear and torsion for the range of analyzed plate thicknesses, which is typical in practice. The interaction response is instead affected by plate aspect ratio and the direction of translation. Thus, a best fit for each aspect ratio and eccentric translation can be adopted as a representative failure envelope (Fig. 11).

The FE derived envelopes are compared with a PLA analytical solution, which was originally developed by Nouri et al. (2014) for surface foundations (i.e. zero embedment) under combined inplane translation and torsion. The formulation is slightly modified for a deeply embedded plate of zero thickness under the same loading combination. Adopting a PLA solution for zero plate thickness ($PLA_{t=0}$) has several advantages:

- It makes the PLA formulation easier and more manageable for excel spreadsheet applications
and provides a straightforward benchmark solution for FE failure envelopes.

- Insensitivity of the normalized interaction response to plate thickness (Fig 11), makes the PLA_{t=0} a simple yet reliable benchmark solution to evaluate the shape of failure envelope in normalized space for other plate thicknesses. Details of the proposed PLA formulation are summarized in Appendix A.

Fig. 12 compares the PLA_{*t*=0} failure envelopes with FE values. Overall, the PLA_{*t*=0} compares very well with the FE failure points for the square plate, slightly overestimates the resistance of rectangular plates under eccentric shear in *x*-direction, and underpredicts the bearing capacity for rectangular plates subjected to sheary-torsion. Besides providing a beneficial benchmark, PLA_{*t*=0} potentially offers a simplified, yet effective, approach for finite thickness plates, if combined with the ultimate resistance factors for shear ($N_{sx.max}$ and $N_{sy.max}$) for $t\neq 0$ (Nouri et al. 2014).

295 Shear capacity reduction with eccentricity

296 Plotting the plate shear resistance versus the eccentricity offers an alternative and more intuitive 297 portrayal of the plate shear capacity reduction with associated torsion. Fig. 13(a) shows examples 298 of FE shear capacity reduction plots versus eccentricity for rectangular and square plates of t = L/7. 299 The y-sliding capacity reduction plots for square and rectangular plates are equivalent to sheary-300 torsion failure envelopes in Fig. 10. Since the shear capacity in Fig. 13(a) is expressed as 301 $N_s = H/WLs_u$, the plots include both the size and shape elements for constructing a failure envelope. 302 Similarly to Fig. 11, Fig. 13(b) summarizes all the FE results with a best fit load reduction curves, but in $(H/H_{max} - e/L)$ space. Fig. 13(b) indicates that plates of higher aspect ratio are more effec-303 304 tive in resisting eccentric loads than square plates, irrespective of the orientation of the plate with respect to shear. However, eccentric sliding in *x*-direction perpendicular to the longer side of the plate tends to reduce the shear capacity more significantly than eccentric *y*-sliding. The same trend was also observed in PLA capacity reduction curves for surface foundations (Nouri et al. 2014).

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4-2-2- Co-planar shear (combined shearx-sheary)

FE analyses are also conducted to determine the interaction response of square and rectangular (*W/L=2*) plates of t = L/20 and L/7 under combined co-planar shear. Typical FE examples of swipe and probe envelopes for N_{sx} - N_{sy} are summarized in Fig. 14 for rectangular plate of t = L/7.

The shape of the co-planar normalized failure envelope in the $H_x/H_{x,max}-H_y/H_{y,max}$ space (Fig. 15) is approximately the same for all plate thicknesses and aspect ratios. Therefore, a single curve for the FE data could be used to describe the normalized interaction response of plates with various thicknesses and geometries. The normalized co-planar shear ($H_x/H_{x,max}-H_y/H_{y,max}$) failure envelope can be described using the following equation:

$$f\left(H_{x},H_{y}\right) = \left(\frac{H_{x}}{H_{x,max}}\right)^{hx} + \left(\frac{H_{y}}{H_{y,max}}\right)^{hy} - 1 = 0$$

$$\tag{4}$$

317 Curve fitting yields interaction factors of hx=hy=2.5. For a zero thickness plate moving along any 318 arbitrary in-plane direction, the plate shear resistance factor is always constant and equal to 2 re-319 gardless of the direction of plate translation. A constant resistance factor results in a circular locus 320 with radius of 2 in the N_{sx} - N_{sy} space and a circle of radius 1 in $H_x/H_{x,max}$ - $H_y/H_{y,max}$ space (i.e. 321 hx=hy=2.0 as shown in Fig. 15). This concept has also been validated through PLA by Nouri 322 (2013). Comparison of circular (t=0) and FE shearx-sheary envelopes in Fig. 15 suggests that even 323 a slight increase in plate thickness increases the plate shear resistance against the reduction in-324 duced by another shear force applied in the perpendicular direction. This increase tends to be in-325 sensitive to plate thickness at least for $t \le L/7$ or changes in plate aspect ratio.

4-3- Interaction equation

327 Eq. (4) for concentric in-plane shear forces can be generalized for eccentric planar loads:

$$f(H_x, H_y, T) = \left[\left(\frac{H_x}{H_{x, \max}} \right)^{hx} + \left(\frac{H_y}{H_{y, \max}} \right)^{hy} \right]^h + \left(\frac{T}{T_{\max}} \right)^{mz} - 1 = 0$$
(5)

328 $f(H_x, H_y, T) = 0$ represents an ellipsoid in the $H_x/H_{x,max} - H_y/H_{y,max} - T/T_{max}$ normalized space. 329 Since failure envelopes in the normalized space seem to be insensitive to plate thickness, the in-330 teraction factors of hx, hy, h, and mz are determined by curve fitting Eq. (5) to all FE best fit fail-331 ure envelopes for the square and rectangular plates in Figs. 11 and 15. Table 3 summarizes the 332 calculated interaction factors.

333 **5- CONCLUSIONS**

The purpose of this study is to evaluate pure translational and torsional ultimate plate capacity, as well as to study the behavior of deeply embedded square and rectangular plates of finite thickness under combined co-planar shear and torsion using 2D and 3D finite element (FE) simulations. This study indicates the following:

1. Very fine meshes and higher order elements do not sufficiently eliminate stress concentration
and overprediction of capacity for a deeply embedded plate in undrained condition. Application of contacts in the vicinity of the plate corners in soil improves bearing capacity estimates. The ultimate sliding and torsional capacities, as well as the interaction factors, are reported for a range of plate thicknesses and aspect ratios common in practice (Tables 1-2 and
Figs. 8-9).

2. Plate shear resistance begins to reduce even at small eccentricity (e.g. 5% reduction for e/L=0.1 for square plate as shown in Fig. 13). Shear capacity reduces at a higher rate for larger eccentricity (e.g. 50% shear capacity reduction for e=0.5L or eccentric shear force at the edge of square plate).

348 3. Reduction in shear resistance is affected slightly by the plate thickness but essentially de 349 creases with plate aspect ratio, since square plates are more vulnerable to reduction of sliding

capacity compared to rectangular plates (Fig. 13). In general, an eccentric shear load parallel
to the shorter side of the plate induces greater capacity reduction in rectangular plates.

4. Thickness has a marked effect on a plate shear and torsional resistance. For perspective, the shear resistance factor equal to $N_s=2$ for a zero thickness plate, increases to about 4.4 for a square plate with t=L/7 (more than 100% increase). Similarly, the torsional resistance factor (N_t) increases by about 90% for square and rectangular plates with t=L/7.

5. The shape of the shear-torsion failure envelope in the normalized space $(H/H_{max} - T/T_{max})$ is insensitive to plate thickness (Fig. 11). This facilitates the application of analytical solutions developed for infinitely thin plates, PLA_{t=0}, (Nouri et al. 2014) to estimate the shape of the failure envelope, which shows a reasonable agreement with FE values (Fig. 13). As shown in Fig. 11, the shape of the shear-torsion failure envelope is affected by aspect ratio and the direction of the eccentric shear force.

362 6. The shape of the co-planar shear (shearx-sheary) failure envelope is generally insensitive to 363 plate geometry and thickness, at least for the cases examined (Fig. 15).

364 7. A simple relationship is proposed to describe the failure envelope in the normalized space 365 $(H_x/H_{x,max} - H_y/H_{y,max} - T/T_{max})$. Taking advantage of the insensitivity of the shape of the fail-366 ure envelope to thickness, a representative envelope and the associated interaction factors are 367 sufficient for practical design applications.

368

369 APPENDIX A. PLA virtual work formulation for zero thickness plate

For a horizontal load *H* applied at a distance *e* from the center of the plate (Fig A.1), equating external virtual work \dot{W} to internal energy dissipation yields to:

372
$$H = \dot{D} / (\rho + e) \dot{\beta}$$
(A.1)

373 where $\dot{\beta}$ is a virtual angular velocity of the plate about a rotation center located a distance ρ from 374 the center. The rate of internal energy dissipation \dot{D} is calculated by integrating the soil resistance 375 times the local velocity over the plate area:

376
$$\dot{D} = 2s_u \dot{\beta} \int_{-W/2}^{W/2} \int_{-L/2}^{L/2} \sqrt{(\rho + x)^2 + y^2} dx dy$$
 (A.2)

The dissipation term (\dot{D}) in the proposed solution is two times the \dot{D} for a surface foundation (Nouri et al. 2014), since the energy is dissipated along both top and bottom faces of the plate. This factor of 2 in the dissipation term is the only difference between this derivation and the PLA formulation for surface foundations (Nouri et al. 2014).

381 A least upper bound collapse load is obtained by minimizing H with respect to ρ , which leads to:

382
$$H = \frac{1}{\dot{\beta}} \frac{\partial \dot{D}}{\partial \rho}$$
(A.3)

383 The collapse load (*H*) for the case of zero plate thickness becomes:

384
$$H = 2s_u \int_{-W/2}^{W/2} \int_{-L/2}^{L/2} \frac{\rho_{opt} + x}{\sqrt{(\rho_{opt} + x)^2 + y^2}} \, dx \, dy \tag{A.4}$$

385 where
$$\rho_{opt}$$
 is the distance to the optimal center of rotation corresponding to a least upper bound.
386 Evaluation of the integral yields the following closed-form expression for *H*:

$$H = 2s_{u} \left[a_{1}^{2} \ln \left| \frac{b_{1} + W/2}{b_{1} - W/2} \right| - a_{2}^{2} \ln \left| \frac{b_{2} + W/2}{b_{2} - W/2} \right| + W(b_{1} - b_{2}) \right]$$

$$a_{1} = \rho_{opt} + L/2$$

$$387 \qquad a_{2} = \rho_{opt} - L/2$$

$$b_{1} = \sqrt{a_{1}^{2} + W^{2}/4}$$

$$b_{2} = \sqrt{a_{2}^{2} + W^{2}/4}$$
(A.5)

The eccentricity, *e*, associated with any arbitrarily selected ρ is calculated from Eq. (A.1) with \dot{D} and *H* evaluated and substituted from Eqs. (A.2) and (A.4), respectively. Analytical integration of Eq. (A.2) is unwieldy; however, reduction to single integration is possible to facilitate simple design spreadsheet calculations:

$$\dot{D} = 2s_{u}\dot{\beta}\int_{-W/2}^{W/2} \left[(a_{1}c_{1} - a_{2}c_{2}) + y^{2} \ln \left| \frac{a_{1} + c_{1}}{a_{2} + c_{2}} \right| \right] dy$$

$$a_{1} = \rho_{opt} + L/2$$

$$a_{2} = \rho_{opt} - L/2$$

$$c_{1} = \sqrt{a_{1}^{2} + y^{2}}$$

$$c_{2} = \sqrt{a_{2}^{2} + y^{2}}$$
(A.6)

392

393 Eq. (A.6) can be integrated between the limits -W/2 to W/2 using classical numerical integration394 formulas.

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