

## **Temperature - Time Curves of Complete Process of Fire Development**

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S. E. MAGNUSSON and S. THELANDERSSON

TEMPERATURE-TIME CURVES OF COMPLETE PROCESS OF FIRE DEVELOPMENT

# ACTA POLYTECHNICA SCANDINAVICA

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Temperature—Time Curves of Complete Process of Fire Development

Theoretical Study of Wood Fuel Fires in Enclosed Spaces

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## Table of contents

	rincipal notations	5
1.	Introduction	8
2.	Equation of heat balance of process of fire development	17
3.	Study of combustion gas flow and heat flow through openings in enclosed spaces	27
	Flow conditions in vertical rectangular opening in enclosed space Deduction of maximum rate of combustion determined by rate of air supply in	28
	accordance with [5]	30
	in enclosed space during the whole process of fire development Modification of treatment in cases where enclosed spaces are provided with several	31
	openings which differ in height	
	zontal openings in roofs	34
4.	Description of programme for digital computer	39
5.	Calculation of time graphs of rate of combustion for some full-scale tests described in literature	42
	Test Series A. Tests carried out by Sjölin. Calculation of time graphs of rate of combustion	44
	Test Series B. Tests made by Kawagoe [5]. Calculation of time graphs of rate of combustion	46
	Test Series C. Tests published by Ödeen [18]. Calculation of time graphs of rate of combustion	48
	Test Series D. Tests made at the National Swedish Institute for Materials Testing.  Calculation of time graphs of rate of combustion	55
	Summary	57
6.	Determination of general time graphs of quantity of energy released per unit time during different phases of process of fire development	61
7.	Calculation of time graphs of temperature of combustion gases for characteristic types of enclosed spaces	64
8.	Summary	76

Acknowledgements	79
References	80
Appendix 1. Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases. Test Series A	84
Appendix 2. Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases. Test Series B	89
Appendix 3. Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases. Test Series C	91
Appendix 4. Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases. Test Series D	100
Appendix 5. Calculated time graphs of temperature of combustion gases for seven types of enclosed spaces differing in opening factor and in bounding structures	103
Appendix 6. Calculated time graphs of temperature of combustion gases represented in tabular form for seven types of enclosed spaces differing in opening factor and in bounding structures	129

## Principal notations

$\boldsymbol{A}$	Area of vertical openings in the en-	•
	closed space	$m^2$
$A_h$	Area of horizontal openings in the en-	
,	closed space	$m^2$
$A_t$	Total bounding surface area of the en-	
	closed space	$m^2$
$(A \cdot \sqrt[]{H})$	Air flow factor (Ventilation factor)	$m^{5/2}$
$(A \cdot \sqrt{H}/A_t)$	· •	$m^{1/2}$
В	Width of an opening in the enclosed	
_	space	m
G	Volume of combustion gases produced	
_	per unit weight of fuel	$m^3 \cdot kg^{-1}$
$G_{0}$	Volume of combustion gases (expressed	
	in Nm³) produced per unit weight of	
	fuel	$\mathrm{Nm^3\cdot kg^{-1}}$
H	Height of vertical opening in the en-	
	closed space	m
H'	Height of vertical opening below the	
TT#	neutral zone level in the enclosed space	m
H''	Height of vertical opening above the	
_	neutral zone level in the enclosed space	m
I	Enthalpy	kcal $\cdot$ m <sup>-3</sup>
$I_C$	Heat energy released per unit time	
-	during combustion	kcal · h⁻¹
$I_{\mathcal{B}}$	Heat energy stored per unit time in the	
	gas volume which is contained in the	
-	enclosed space	kcal · h <sup>−1</sup>
$I_L$	Heat energy withdrawn per unit time	
	from the enclosed space owing to the	
7	replacement of hot gases by cold air	kcal · h⁻¹
$I_R$	Heat energy withdrawn per unit time	
	from the enclosed space by radiation	
	through openings in the enclosed space	$kcal \cdot h^{-1}$

$I_W$	Heat energy withdrawn per unit time from the enclosed space through wall,	
	roof or ceiling, and floor structures	kcal · h <sup>−1</sup>
L	Quantity of air consumed per unit weight of fuel during combustion	$Nm^3 \cdot kg^{-1}$
M	Quantity of combustible material	kg
P	Static pressure	kg⋅m <sup>-2</sup>
$Q_{ m out}$	Rate of flow of the outgoing gases	
2 out	through a vertical opening in the en-	
	closed space	kg⋅h <sup>-1</sup>
$Q_{ m in}$	Rate of flow of the incoming air through	_
æm	a vertical opening in the enclosed space	kg⋅h <sup>-1</sup>
$Q_h$	Rate of flow of the outgoing gases	
2"	through a horizontal opening in the	
	enclosed space	kg⋅h <sup>-1</sup>
Q	Rate of flow of air supplied to the en-	
	closed space by means of fans	$m^3 \cdot s^{-1}$
R	Rate of combustion	kg of wood per
		unit time
$R_{\text{max}}$	Maximum rate of combustion deter-	kg of wood per
	mined by the rate of air supply	unit time
T	Duration of the fire defined as the du-	
	ration of the flame phase	h
W	Heat value of the fuel	kcal · kg <sup>-1</sup>
c	Specific heat	$kcal \cdot m^{-3} \cdot {}^{\circ}C^{-1}$
$C_{P}$	Specific heat of the combustion gases	$kcal \cdot m^{-3} \cdot {}^{\circ}C^{-1}$
g	Acceleration of gravity	$m \cdot s^{-2}$
h	Difference in level between the centre of	
	a vertical opening and a horizontal	
	opening	m
h'	Difference in level between the neutral	
	zone and a horizontal opening	m $Mcal \cdot m^{-2}$ of bound-
q	Fire load	
	II. J	ing surface area
r	Hydraulic radius	cm m ⋅ h <sup>-1</sup>
$v_y, v_z$	Velocity of flow	ш.п.
$v_h$	Velocity of flow through a horizontal	m · h-1
4	opening in the enclosed space Time co-ordinate	h
t	Position co-ordinate	m
x	Coefficient of heat transfer at a surface	X11
$\alpha_i$	exposed to fire (internal surface)	$kcal \cdot m^{-2} \cdot h^{-1} \cdot {}^{\circ}C^{-1}$
	exposed to me (internal surface)	Roux III II

$\alpha_u$	Coefficient of heat transfer at a surface	<b>;</b>
	not exposed to fire (external surface)	kcal · m -2 · h -1 · °C-1
γ	Weight per unit volume	kg⋅m <sup>-3</sup>
$arepsilon_{res}$	Resultant emissivity for radiation be-	
	tween flames, combustion gases, and a	
	surface exposed to fire (internal surface)	
$arepsilon_{fl}$	Emissivity of flames	
$\varepsilon_i$	Emissivity of a surface exposed to fire	
િ	Temperature	°C
$\vartheta_{0}$	Temperature of the outside air	°C
$\vartheta_g$	Temperature of the combustion gases	°C
$artheta_i$	Temperature of a surface exposed to	C
	fire (internal surface)	°C
$\vartheta_u$	Temperature of a surface not exposed	C
	to fire (external surface)	°C
$\Delta \vartheta$	Temperature difference between the	C
	combustion gases and the outside air	$^{\circ}\mathrm{C}$
λ	Thermal conductivity	$kcal \cdot m^{-1} \cdot h^{-1} \cdot {}^{\circ}C^{-1}$
$\mu$	Coefficient of contraction	real III - II - I
ho	Density	ka . m = 3
$ ho_{0}$	Density of the outside air	kg·m <sup>-3</sup>
$ ho_{g}$	Density of the combustion gases	kg·m <sup>-3</sup>
	, and a strong gases	kg⋅m <sup>-3</sup>

### 1. Introduction

The efforts made during the past decade in the field of structural fire engineering research have paved the way for differentiated, functionally correct structural fire engineering design carried out on the basis of theoretical calculations. This was rendered possible by investigations which can on the whole be classified in one or several of the main groups enumerated below. At the same time, these groups may be regarded as the essential stages or steps in an appropriate procedure for fire engineering design of load-bearing and separating structures [1].

- (a) Determination of the characteristics of the fire load in an enclosed space under exposure to fire.
- (b) Study of the variations in the development of energy, in the requisite air supply, and in the evolution of gases, with the time in the course of a fire. Determination of the temperature of the combustion gases in the enclosed space as a function of the time.
- (c) Determination of the thermal properties of the materials used for structures in the temperature range which is of interest in connection with fires.
- (d) Determination of the non-stationary temperature fields which are produced in a fire-exposed structure on the assumption that the temperature-time curve for the combustion gases is given, cf. (b).
- (e) Determination of the structural behaviour and the load-bearing capacity of a fire-exposed structure on the basis of the temperature fields defined under (d), and with the help of the available information on those changes in the strength and deformation characteristics of the materials which take place under such conditions.

The object of the present investigation is to make a close study of the stage (b) in order to determine the complete temperature-time curve for the gaseous products of combustion under different conditions, and in particular the temperature-time curve in the cooling phase, 1) for fires of the wood fuel type in enclosed spaces.

<sup>1)</sup> The characteristics of the different phases of the process of fire development are represented in Fig. 1. The term "cooling phase" will be used in this publication to designate the smoulder phase and the cooling phase taken together.

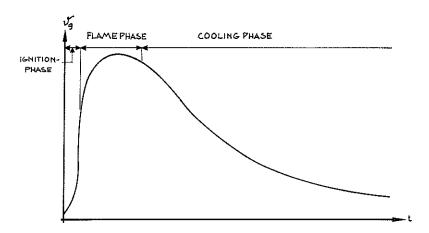


Fig. 1. Phases of the process of fire development as defined in the present publication. The definition in [1] distinguishes between a smoulder phase and a cooling phase, which are here regarded as a single phase designated by the term "cooling phase".

The present state of research in this field is clearly reflected in the sections dealing with fire protection in the Swedish Building Regulations 1967 (abbreviated SBR 67) and in the Draft Specification "Aluminium Structures". In comparison with the relevant regulations which are in force in most other countries, the Swedish rules represent substantial progress on the road to judicious structural fire engineering design. This is primarily due to the fact, that, when the designer has to choose that temperature-time curve which characterises the process of fire development, and which must serve as a basis for all theoretical structural fire engineering design, these rules enables the designer to be guided by all the results, which have been obtained from research in this field during recent years.

On an international plane, it is found that standard temperature-time curves for the process of fire development have been adopted in several countries [1]. If the fuel supply is unlimited, then the agreement between these curves is relatively close, see Fig. 2 a. Under practical conditions, when the fuel supply is limited, the standard specifications used in various countries stipulate that the variation in the temperature with the time shall be in conformity with the standard curve during a certain definite period of time, which is designated by the term "duration of the fire", T, and is defined as the duration of the flame phase. A comparison of the relations between the duration of the fire and the fire load which are employed in various countries is represented in Fig. 2 b [1]. This comparison shows very great differences in an assumption which is fundamental for structural fire engineering design. The wide dispersion between the curves reproduced in Fig. 2 b indicates that there exists no univalued relation between the fire load and the duration of the fire. Concerning

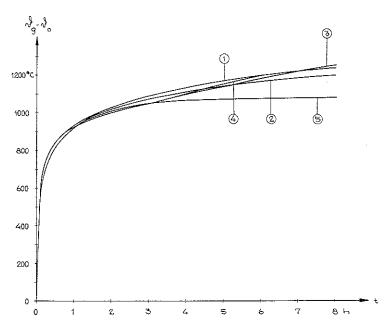


Fig. 2a. Standard curves used in some countries to represent the variation in the temperature,  $\vartheta$ , in an enclosed space exposed to fire with the time, t. The symbol  $\vartheta_0$  denotes the temperature in the enclosed space at the time t=0.

- 1. ISO/TC 92; INSTA 28/2; DIN 4102-62.
- 2. EMPA, Switzerland.
- 3. ASTM 119 (1953), USA.
- 4. V 1076 (1955), Netherlands; BS 476 (1953), United Kingdom.
- 5. A 1304, Japan.

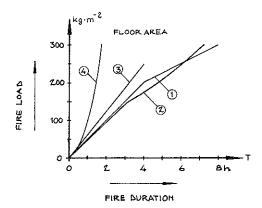


Fig. 2b. Relations between the fire load, in kg of wood per m<sup>2</sup> of floor area, and the duration of the fire, in h, which are stipulated in the standard specifications used in Sweden (Curve 1), United States of America (Curve 2), and United Kingdom (Curve 3), as well as in a Swiss draft standard specification (Curve 4).

the cooling phase of the process of fire development, it should be noted that it is as a rule completely disregarded.

The relevant Swedish standard specifications provide the designer with three alternative methods of design. Just as in other countries, it is permissible to carry out the design in a roughly simplified and stereotyped manner by using Curve 1 in Fig. 2 a as a point of departure. For the flame phase, this curve gives the temperature of the combustion gases,  $\mathfrak{F}_g$ , in the enclosed space in accordance with the equation

$$\theta_q - \theta_0 = 1325 - 430 e^{-0.2t} - 270 e^{-1.7t} - 635 e^{-19t}$$
 (1.1)

where t is the time, in hours, and  $\vartheta_0$  is the temperature in the enclosed space at the time t=0. The differences in the combustion characteristics of various fuels, or the fact that the rate of combustion varies within wide limits with the dimensions of the openings in the enclosed space, are not taken into account in this equation. The above-mentioned curve is closely in agreement with that temperature-time curve which is recommended by the ISO for fire tests on building components.

Alternatively, for certain definite types of fire loads and enclosed spaces, the designer may use a method which is simplified, but is nevertheless more differentiated, in comparison with the design procedure outlined in the above. The applicability of this alternative method presupposes that it is possible to comply with the two necessary conditions which are stated in what follows. In the first place, it is required to demonstrate that the characteristics of the fire load in respect of rate of combustion and radiation are approximately in accordance with those which apply in the case of wood fuel. In the second place, it is stipulated that the opening factor of the enclosed space, which is given by the expression  $A\sqrt{H}/A_t$ , where A is the total opening area of windows and doors, in m<sup>2</sup>, H is a weighted average of the vertical dimensions of these openings, in m, and  $A_t$  is the total area of the surfaces bounding the enclosed space, in m2, shall be known during all phases of the process of fire development. If these two conditions are satisfied, then it is allowed to carry out the design on the basis of a specific curve representing the variation in the temperature of the combustion gases in the enclosed space,  $\vartheta_a$ , with the time. For the flame phase, this curve is determined by the opening factor, see Fig. 2 c, in the course of the duration of the fire, T, which is defined by the equation

$$T = qA_t/(25A\sqrt{H}) \quad \text{min} \tag{1.2}$$

where q is the fire load, in Mcal·m<sup>-2</sup> of bounding surface area. The dash-line curve represents the INSTA curve expressed by Eq. (1.1).

Finally, and generally, in the cases where the quantities of combustible materials which constitute the fire load, as well as the rate of combustion, are accurately known, the above-mentioned two Swedish specifications allow

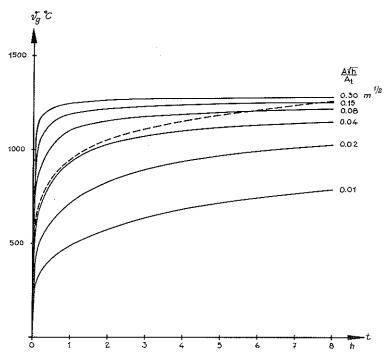


Fig. 2c. Variation in the temperature of the combustion gases,  $\vartheta_g$ , with the time, t, at different values of the opening factor,  $A \cdot \sqrt{H}/A_t$ . Curves published in the Swedish Building Regulations 1967 (SBR 67) and in the Swedish draft specification "Aluminium Structures". The dash-line curve is the standard curve calculated by means of Eq. (1.1).

the fire resistance of a building component to be determined on the basis of the variation in the temperature of the combustion gases with the time, calculated with the help of the formula which is known as the equation of heat balance. This equation, which constitutes a fundamental description of the energy balance of the process of fire development, and hence also serves as a basis for Eq. (1.2), will be discussed at some length in Chapter 2.

What has been said up to this point relates only to the ignition and flame phases of the process of fire development. As regards the cooling phase, it is stipulated in a summary manner merely that the time graph of the temperature of the combustion gases shall be chosen so as to be linear, and that the rate of decrease in the gas temperature shall be taken to be 10°C·min<sup>-1</sup>, unless other assumptions can be demonstrated to be more correct. If the design is carried out by means of the second or third alternative method, each of which is functionally realistic, then this implies that two phases of the same continuous process are represented in such a way that their descriptions are entirely different in the degree of accuracy as well as in the extent to which the actual conditions are taken into account. It is obvious that this gives

rise to a considerable unbalance in the basis for design. To show how important it is that the cooling phase should also be described in a differentiated manner, it may be useful to give two examples of structural members which are characterised by low and high thermal inertia, respectively.

For an enclosed space, where the fire load is  $q = 12 \text{ Mcal} \cdot \text{m}^{-2}$  of bounding surface area, and the opening factor is  $A\sqrt{H}/A_t = 0.08 \text{ m}^{0.5}$ , Fig. 3 represents a calculated temperature-time curve for a steel column exposed to fire and characterized by the ratio  $F_s/V_s = 100 \text{ m}^{-1}$  and by  $\varepsilon_r = 0.5$ , where  $F_s$  is the total bounding surface area of the column, which is equal to its fire-exposed area, in  $m^2$ ,  $V_s$  is the steel volume of the column, in  $m^3$ , and  $\varepsilon_r$  is the resultant emissivity for heat transfer from the flames and the combustion gases to the

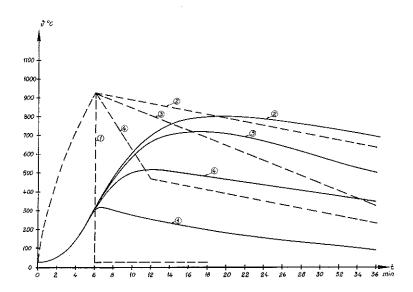


Fig. 3. Variation in the temperature,  $\vartheta_s$ , of a steel column with the time, t, calculated on the basis of temperature-time curves, which differ in shape during the cooling phase (dash-line curves).

 $F_s/V_s = {
m Ratio}$ , in m<sup>-1</sup>, of the fire-exposed surface area, i.e. the total bounding surface area, of the column to the steel volume of the column.

 $q = 12 \text{ Mcal} \cdot \text{m}^2$ .

- 1. Instantaneous cooling. Rate of decrease in temperature  $\infty^{\circ}C \cdot h^{-1}.$
- 2. Linear rate of decrease in temperature, 600°C·h<sup>-1</sup>.
- 3. Linear rate of decrease in temperature, 1200°C·h<sup>-1</sup>.
- 4. Linear rate of decrease in temperature, 6 min  $\leq t \leq 12$  min, 4600°C·h<sup>-1</sup>;  $t \geq 12$  min, 600°C·h<sup>-1</sup>.

 $F_s/V_s = 100 \text{ m}^{-1} \ \varepsilon_r = 0.5.$ 

steel column. The characteristics of the flame phase for the temperature-time curve of the enclosed space have been chosen in conformity with the Swedish Building Regulations 1967. In Fig. 3, the full-line curves represent the temperature of the steel,  $\theta_s$ , and the dash-line curves show the temperature of the combustion gases,  $\theta_g$ , on the basis of the four alternative assumptions concerning the cooling phase of the process of fire development which are stated in what follows.

- (1) After the duration of the fire  $T = qA_t/(25A\sqrt{H}) = 6$  min, the temperature of the combustion gases drops instantaneously to ordinary room temperature.
- (2) The temperature of the combustion gases decreases in accordance with the Swedish Building Regulations 1967 at a linear rate of  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ .
- (3) The temperature of the combustion gases decreases at a linear rate of  $20^{\circ}\text{C} \cdot \text{min}^{-1}$ .
- (4) The temperature of the combustion gases is assumed to vary in a more realistic manner, that is to say, it drops to half their maximum temperature during the first 6 min of the cooling phase, and then decreases at a linear rate of  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ .

At the end of the flame phase, the temperature of the steel is 303°C. After that, the temperature of the steel continues to increase during the cooling phase of the process of fire development to its respective maximum values corresponding to the four alternative assumptions, viz., 303, 799, 719, and 518°C.

Fig. 4 shows the effects on the load-bearing capacity of a reinforced concrete slab which are produced by different slopes of the linear cooling phase of the time-temperature curves. Each T-value on the horisontal axis corresponds to a specific time-temperature curve. T is the duration of the flame phase, and depends on the fire load. The distance from the centre lines of the reinforcing bars to the fire-exposed surface of the slab, is assumed to be 2 cm. The emissivity of the flames is taken to be 0.7. The temperature in the enclosed space is supposed to vary with the time during the flame phase in accordance with Eq. (1.1). The temperature of the combustion gases is assumed to decrease at linear rates of 5, 10, and  $20^{\circ}\text{C} \cdot \text{min}^{-1}$ , or to drop instantaneously to ordinary room temperature ( $\infty^{\circ}\text{C} \cdot \text{h}^{-1}$ ). For a static load which causes failure at a temperature of the reinforcing bars  $\vartheta_{scr} = 450^{\circ}\text{C}$ , we obtain a fire resistance period,  $t_{fr}$ , which varies from 0.52 to 0.82 h, and for  $\vartheta_{scr} = 500^{\circ}\text{C}$ , the corresponding variation in the fire resistance period ranges from 0.72 to 1.01 h.

These examples show that it is necessary to calculate the fire resistance period of a building component so as to take account of that reduction in its load-bearing capacity, or in its separating capacity, which occurs during the cooling phase. Moreover, they indicate that those temperature-time curves for enclosed spaces which are to serve as a basis for such calculations should also be differentiated and as realistic as possible in the cooling phase.

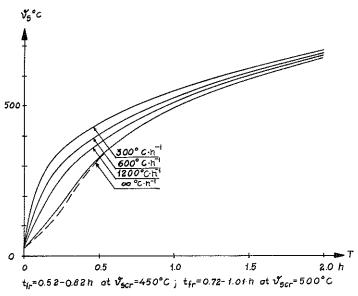


Fig. 4. Relation between the calculated maximum temperature,  $\vartheta_s$ , of the reinforcing steel in a concrete slab, 18 cm in thickness, exposed to fire on one side, and the duration, T, of the flame phase of the process of fire development. This relation is represented for different values of the fire load and for different slopes of the linear cooling phase of the temperature-time curve.

Cooling phase taken into account.

---- Cooling phase not taken into account.

Concrete cover 2 cm.

 $\varepsilon_{fl}$  = Emissivity of the flames=0.7.

 $t_{fr}$  = Fire resistance period.

 $\theta_{sc_r}$  = Critical temperature of the steel, i.e. the temperature at which the reinforcement fails.

 $t_{fr} = 0.52 \text{ to } 0.82 \text{ h at } \theta_{sc_r} = 450 \,^{\circ}\text{C}.$ 

 $t_{fr} = 0.72 \text{ to } 1.01 \text{ h at } \vartheta_{scr} = 500^{\circ}\text{C}.$ 

The object of the present investigation is therefore to evolve a method which shall be applicable to different combinations of the values of the air flow factor (ventilation factor),  $A\sqrt{H}$ , and the fire load, as well as the type of material of the structures bounding the enclosed space under exposure to fire of the wood fuel type. This shall enable that the temperature-time curves for the enclosed space to be calculated by means of a theoretical procedure so as to cover the whole process of fire development, and shall thus make it possible to carry out judicious structural fire engineering design on the basis of the variation in the temperature with the time during all phases of the fire.

In connection with a general treatment of the equation of heat balance, Chapter 2 deals with the problems which are met with when a theoretical calculation of the temperature-time curve for the combustion gases is extended so as to comprise the cooling phase. Chapter 3 describes the methods which have been used to tackle these problems, and the modifications of the equation of heat balance which have been necessary for this purpose. In Chapter 4, an account is given of the computer programme which has been prepared for the calculations, and which is represented in the form of a flow chart. In Chapter 5, the full-scale tests which have served as a basis for the present investigation are subjected to comparative theoretical analysis. The time graphs of the rate of combustion which have been determined with the help of these theoretical analyses are presented in Chapter 6. Finally, in Chapter 7, these graphs are used as a basis for the calculation of complete temperature-time curves for combustion gases in enclosed spaces which vary in the values of the opening factor and the fire load, as well as in type of material employed in the structures bounding the enclosed space.

# 2. Equation of heat balance of process of fire development

The papers published by Kawagoe and Sekine [2], as well as by Ödeen [3], in the early 1960ies have made it possible to carry out theoretical calculations of temperature-time curves for combustion gases in the flame phase of the process of fire development to a degree of accuracy that is sufficient for practical purposes. These three authors have studied the energy balance during the process of fire development. The quantity of energy released per unit time, just as the volume of combustion gases evolved, during combustion were assumed to be known. With the help of the calculation of the quantity of energy which was lost per unit time by conduction and radiation from the enclosed space through its bounding structures, it was possible to deduce an equation of heat balance, and to solve it so as to obtain the temperature of the combustion gases. The treatment of this problem was based on the simplified assumptions which are reproduced in what follows.

- (a) The temperature in the interior of the whole enclosed space is uniform at any given instant.
- (b) The coefficient of heat transfer to the interior bounding surfaces of the enclosed space is uniform at every point.
- (c) The heat flow through the bounding structures of the enclosed space is one-dimensional and, except for the window and door openings, if any, uniformly distributed.

Kawagoe and Sekine, as well as Ödeen, confined themselves throughout their papers to a study of the flame phase of the process of fire development, and the equation which they have deduced cannot be applied directly to the cooling phase. Primarily, the calculation of the temperature-time curve for the combustion gases during the cooling period requires an analytical investigation of two fundamental sub-problems, which have been but little studied up to the present time. In the first place, it is necessary to determine the quantity of energy liberated per unit time when this quantity, as is the case in the cooling phase, is no longer determined by the rate of air supply. In the second place, it is required to investigate the thermodynamic conditions which are encountered when the rate of combustion is no longer limited by the dimensions of the openings in the enclosed space. The equation of energy balance which has been deduced by Kawagoe and Sekine and by Ödeen, as

well as that extension of the theory which is required for the calculation of the temperature-time curve in the cooling phase of the process of fire development, will be dealt with in what follows.

The above-mentioned equation expresses for any given instant ,t, the balance between the respective quantities of heat energy generated and lost per unit time in the enclosed space under consideration. In its complete form, this equation is

$$I_C = I_L + I_W + I_R + I_B$$
 where (2.1)

 $I_C$  = the heat energy released per unit time during combustion,

 $I_L$  = the heat energy withdrawn per unit time from the enclosed space owing to the replacement of hot gases by cold air,

 $I_W$  = the heat energy withdrawn per unit time from the enclosed space through wall, roof or ceiling, and floor structures,

 $I_R$  = the heat energy withdrawn per unit time from the enclosed space by radiation through the openings in the enclosed space,

 $I_B$  = the heat energy stored per unit time in the gas volume which is contained in the enclosed space.

The terms entering into the above equation are schematically illustrated in Fig. 5.

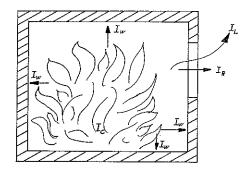


Fig. 5. Schematic illustration of the terms entering into the equation of heat balance.

## Term $I_B$

In comparison with the quantities of energy which are involved in fires, the quantity of energy which can be stored in the gas volume contained in the enclosed space is of minor importance. Therefore, the term  $I_B$  can to a close approximation be put equal to zero.

Term  $I_C$ 

In the case of fires of the wood fuel type, the term  $I_c$ , which expresses the quantity of heat released per unit time during combustion, is a factor which is difficult to determine in the equation of heat balance. In order to obtain this term, Kawagoe and Sekine, as well as Ödeen, chose as a point of departure the rate of combustion, expressed in terms of kilogrammes of wood per unit time, which they multiplied by a heat value.

This gives the equation

$$I_C = R \cdot W \quad \text{kcal} \cdot h^{-1}$$
 (2.2) where

R = the rate of combustion, in kg of wood per h, W = the heat value, in kcal per kg of wood.

With the exception of a small number of calculations which were based on a triangular variation in the rate of combustion, R, with the time, Ödeen assumed that R is constant, and is arbitrarily chosen, and that the quantity of heat, W, released during the flame phase is 4120 kcal per kg of wood. Thus, the heat energy liberated per unit time during combustion,  $I_C$ , is supposed to remain constant until the fuel has burnt up, and the cooling phase is characterized by the fact that no additional energy is supplied to the enclosed space. This description of the quantity of energy developed per unit time is applicable to burning of liquid fuels, but does not take account of the real combustion characteristics of wood fuels. As a rule, these fuels give off 30 to 50 per cent of the total quantity of energy after the end of the flame phase.

The investigation made by Kawagoe and Sekine is more differentiated in respect of the characterization of the process of fire development. As Kawagoe had shown in an earlier paper [5], during that period of the fire when the rate of combustion reaches a maximum, i.e. during the flame phase, the rate of air supply to the enclosed space, and hence also the maximum rate of combustion,  $R_{\rm max}$ , are proportional to the air flow factor,  $A \sqrt{H}$ .

If the areas are expressed in  $m^2$ , and the maximum rate of combustion,  $R_{\text{max}}$ , is expressed in kg of wood per min, then we have the approximate relation

$$R_{\text{max}} = 5.5 \cdot A \sqrt{H}$$
 kg of wood per min (2.3)

Furthermore, in the papers published by Kawagoe and Sekine, the quantity of heat liberated during the flame phase, W, is stated to be 2575 kcal per kg of wood. This value was obtained by reducing the nominal heat value of wood so as to take account of the degree of incomplete combustion. The degree of incomplete combustion was estimated with the help of those analyses of the composition of the combustion gases which were carried out during fire

tests [5]. As regards the cooling phase, Kawagoe and Sekine had made a few isolated comparative calculations based on a polygon-shaped time graph of the rate of combustion, and then found that the temperatures obtained when the cooling phase was characterized by a linear decrease in temperature at a rate of 7 or  $10^{\circ}\text{C} \cdot \text{min}^{-1}$  were much too high.

Accordingly, if the results of the investigations made by Ödeen, as well as by Kawagoe and Sekine, are to be applied to fires where the fuel is of the wood type, then the calculations have to be confined to the flame phase of the process of fire development.

When the treatment of this problem is extended so as to comprise the cooling phase of the process of fire development, Eq. (2.3) is not generally applicable. For the quantity of energy released per unit time during combustion, Eq. (2.3), in combination with Eq. (2.2), gives only the theoretical upper limit, which is determined by the available rate of air supply. During the cooling phase, the energy liberated per unit time will be governed by other factors. For this reason, and since no systematic investigation has so far been made in order to determine the relations between the three quantities which are of interest in this connection, viz., the reduction in the weight of fuel, the quantity of energy developed per unit weight of fuel, and the requisite rate of air supply, the quantity of energy released per unit time during the cooling phase of the process of fire development had to be determined by means of the method described in what follows. The procedure in calculation for the solution of the equation of heat balance was programmed for a CD 3600 computer. A study of the literature was carried out in order to examine the available publications on full-scale tests. A number of these tests were selected in the cases where the reported data were so complete as to enable numerical treatment. After that, the computer was used to calculate the temperature-time curve on the basis of an assumed form of the time graph of the rate of combustion for the complete process of fire development. The time graph of the rate of combustion was then varied until the agreement between the experimental and theoretical temperature-time curves was as close as possible. The only absolute requirement to be fulfilled in this connection was that the total quantity of energy liberated during the whole process of fire development should be equal to the total energy of combustion of the fuel. When an adequate range of variation in the opening factor and in the fire load was considered to have been covered, the time graphs of the rate of combustion obtained in this way were systematized. For a given fire load and a given opening factor, it was then possible to assume that the curve showing the variation in the rate of combustion with the time was known on the basis of this systematization. The investigation referred to in the above is described in Chapter 5, and its results, expressed in terms of time graphs of the rate of combustion in a dimensionless form, are presented in Chapter 6.

In connection with the treatment of the term  $I_c$ , it should be pointed out that the fire load must be described in combustion engineering terms in such a way that it may be associated with the equation of heat balance, Eq. (2.1).

In most countries, the fire load is expressed in terms of the quantity of wood that is equivalent to it in heat value per unit *floor area*. This characterization must be replaced by a parameter which has a physical significance when it is treated in calculations. This has been done in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures", which stipulate that the fire load shall be stated as that total quantity of heat, q, in Mcal·m<sup>-2</sup> of total bounding surface area of the enclosed space exposed to fire, which is liberated on the assumption of complete combustion of all the combustible material contained in the enclosed space. A still more refined description of the fire load, which should take into account the variation with the time in the quantity of energy released by combustion, as well as the emissivities of the flames and the combustion gases, is an urgently recommended subject for research in this field.

Term  $I_L$ 

The term  $I_L$  in the equation of heat balance expresses that quantity of heat which is withdrawn per unit time from the enclosed space owing to the replacement of hot gases by cold air through the openings in the enclosed space. For the determination of  $I_L$ , Ödeen, as well as Kawagoe and Sekine, used the equation

$$I_L = R \cdot G_0 \cdot (\vartheta_g - \vartheta_0) \cdot c_p \tag{2.4}$$
 where

R = the rate of combustion, in kg of wood per h,

 $G_0$  = the volume of combustion gases produced by the fire, in Nm<sup>3</sup> · kg<sup>-1</sup> of fuel,<sup>1</sup>)

 $c_p$  = the specific heat of the combustion gases, in kcal·Nm<sup>-3</sup>·°C<sup>-1</sup>,

 $\theta_q$  = the temperature of the combustion gases, in °C,

 $\theta_0$  = the temperature of the air outside the enclosed space, in °C.

Eq. (2.4) states that the term  $I_L$  is put equal to the heat content of the combustion gases produced by the fire with reference to that of the outside air. On account of the difference in density between the cold outside air and the hot gases in the interior of the enclosed space, an exchange of heat by convection takes place in the openings of the enclosed space. The rate of this heat exchange determines the maximum value of the rate of combustion so far as the supply of oxygen is concerned. Therefore, Eq. (2.4) can be used as an expression for  $I_L$  when the rate of combustion is determined by the rate

<sup>1)</sup> Nm<sup>3</sup>=normal cubic metre=the quantity of a gas which occupies a volume of 1 cubic metre at 0°C and 760 mm barometric pressure.

of air supply, in spite of the fact that this equation in itself does not describe I<sub>L</sub>, but expresses the heat content of the combustion gases produced by the fire. However, if the rate of combustion is limited by factors other than the rate of air supply, e.g. by the available quantity of fuel, or if the combustion is completed, then it is obvious that Eq. (2.4) does not hold good. Accordingly, a theoretical treatment of the cooling phase of the process of fire development requires an expression for  $I_L$  which is more generally applicable, and which is based on the rate of air exchange. This problem has been studied in a thesis for degree of Master of Engineering prepared by Ahlquist and Thelandersson [6], in the Division of Structural Mechanics and Concrete Construction, Lund Institute of Technology, Lund, Sweden. These authors based their study on the assumption that the static pressure distribution in the enclosed space varies linearly from the floor to the ceiling or roof, and that there exists a level (the neutral layer) at which the pressure in the enclosed space is equal to the pressure outside the enclosed space. This simplified model has been used by several authors, e.g. by Kawagoe [5] in the theoretical deduction of Eq. (2.3), and by Thomas [7] in studies which dealt with venting in the course of fires. In this connection, Kawagoe showed by means of a number of tests that the assumed pressure was actually applicable to a close approximation. On the assumption that the openings in the enclosed space are vertical only, Kawagoe used Bernoulli's equation to determine the quantities of outgoing combustion gases and incoming cold air as functions of the difference in temperature and the position of the neutral layer. On the basis of the condition that the difference between the quantities of gases flowing into and out of the enclosed space shall be equal to the difference between the quantities of gases produced and consumed by combustion, Ahlquist and Thelandersson calculated the position of the neutral layer as a function of the temperature and the rate of combustion. In this calculation, it was assumed that the rate of combustion may vary from zero to a maximum value, which is dependent on the dimensions of the openings. Furthermore, it was assumed that the liberation of a certain definite quantity of energy is associated with the consumption of the same quantity of air and the production of the same quantity of combustion gases, irrespective of the rate of combustion.

After the position of the neutral layer had been determined in this way, it was possible to obtain an expression for  $I_L$  at different values of the temperature and the rate of combustion. In the above-mentioned thesis [6], the treatment was also extended so as to comprise the modifications which are necessary when the enclosed space is provided with a vent. The deduction of the equations which are required for this purpose is reproduced in its main features in Chapter 3 of the present publication. Moreover, this chapter also contains a summary treatment of the case where the roof of the enclosed space is provided with horizontal openings.

Term  $I_w$ 

The term  $I_W$  denotes the quantity of heat which is withdrawn per unit time from the enclosed space through the structures bounding this space. The term  $I_W$  is determined by solving the general equation of heat conduction in the one-dimensional case under non-steady flow conditions so as to take into account those thermal properties of the materials which are dependent on the temperature, the evaporation of occluded water, and the possible structural transformations in the materials entering into the bounding structures. This equation is

$$c \cdot \gamma \cdot \frac{\partial \vartheta}{\partial t} = \frac{\partial}{\partial x} \left( \lambda_x \cdot \frac{\partial \vartheta}{\partial x} \right) \tag{2.5}$$

where

c = the specific heat of the wall material,

 $\gamma$  = the weight per unit volume of the wall material,

 $\lambda_x$  = the thermal conductivity of the wall material,

9 = the temperature in the interior of the wall material,

t =the time,

x = the position co-ordinate.

The above equation is solved by means of a numerical procedure which had been described by *Odemark* [8], among others, and which has subsequently been further developed by *Ödeen*, and others, [9]. The walls, ceiling, roof, and floor structures which bound the enclosed space are divided into n layers having the thickness  $\Delta x_k$  each, and the equation of heat balance is written for each one of these layers. If the temperature at the centre of the layer k at the time t is denoted by  $\theta_k$ , and that at the time  $t + \Delta t$  is designated by  $\theta_k + \Delta \theta_k$ , then the application of this procedure gives the relations

$$\varphi_{1} \cdot \frac{\Delta \vartheta_{1}}{\Delta t} = \psi_{1}(\vartheta_{g} - \vartheta_{1}) - \psi_{2}(\vartheta_{1} - \vartheta_{2})$$

$$\vdots$$

$$\varphi_{k} \cdot \frac{\Delta \vartheta_{k}}{\Delta t} = \psi_{k}(\vartheta_{k-1} - \vartheta_{k}) - \psi_{k+1}(\vartheta_{k} - \vartheta_{k+1})$$

$$\vdots$$

$$\vdots$$

$$\varphi_{n} \cdot \frac{\Delta \vartheta_{n}}{\Delta t} = \psi_{n}(\vartheta_{n-1} - \vartheta_{n}) - \psi_{n+1}(\vartheta_{n} - \vartheta_{0})$$

$$(2.6)$$

where

 $\theta_q$  = the temperature of the gases in the enclosed space,

 $\theta_0$  = the temperature of the outside air,

$$\varphi_k = \Delta x_k \cdot c(x, \vartheta) \cdot \gamma$$

$$\psi_1 = \frac{1}{\frac{1}{\alpha_i(\vartheta)} + \frac{\Delta x_1}{2 \cdot \lambda(x, \vartheta)}}$$

$$\psi_k = \frac{1}{\frac{\Delta x_{k-1}}{2 \cdot \lambda(x, \vartheta)} + \frac{\Delta x_k}{2 \cdot \lambda(x, \vartheta)}}$$

$$\psi_{n+1} = \frac{1}{\frac{\Delta x_n}{2 \cdot \lambda(x, \vartheta)} + \frac{1}{\alpha_n(\vartheta)}}$$

where

 $\alpha_i(\theta)$  = the coefficient of heat transfer at the internal surface,

 $\alpha_u(\theta)$  = the coefficient of heat transfer at the external surface,

 $\lambda(x, 9)$  = the thermal conductivity at the section x,

 $c(x, \theta)$  = the specific heat at the section x,

 $\gamma$  = the weight per unit volume at the section x.

The coefficient of heat transfer,  $\alpha_i$ , at the internal surface exposed to fire may be supposed to consist of two components, viz., first, a radiation component, which is markedly predominant at the high temperatures in question, and second, a convection component, which can be chosen with adequate accuracy so as to be constant, and to be equal to 20 kcal·m<sup>-2</sup>·h<sup>-1</sup>·°C<sup>-1</sup> [9]. By applying the Stefan-Boltzmann law, this gives, for  $\alpha_i$ , the relation

$$\alpha_{i} = \frac{4.96 \cdot \varepsilon_{res}}{\vartheta_{g} - \vartheta_{i}} \left[ \left( \frac{\vartheta_{g} + 273}{100} \right)^{4} - \left( \frac{\vartheta_{i} + 273}{100} \right)^{4} \right] + + 20 \text{ kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$$
(2.7a)

where

 $\vartheta_i$  = the temperature of the internal surface,

 $\varepsilon_{res}$  = the resultant emissivity for radiation between flames, combustion gases, and the internal surface.

The resultant emissivity,  $\varepsilon_{res}$ , is determined from the formula

$$\frac{1}{\varepsilon_{res}} = \frac{1}{\varepsilon_{fl}} + \frac{1}{\varepsilon_i} - 1 \tag{2.7b}$$

where

 $\varepsilon_{\text{fl}}$  = the emissivity of the flames,

 $\varepsilon_i$  = the emissivity of the surface exposed to fire.

Properly speaking, Eq. (2.7b) represents the emissivity for radiation between two parallel surfaces, but it was considered to be the best available approximation in the case of radiation between flames and a surface exposed to fire

According to [9], the coefficient of heat transfer,  $\alpha_u$ , at the external surface, which is not exposed to fire, can be represented by the approximate expression

$$\alpha_u = 7.5 + 0.028 \cdot \theta_u \quad \text{kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$$
 where

 $\vartheta_u$  = the temperature of the external surface.

The system of differential equations of the first order, Eq. (2.6), is solved numerically (see Chapter 4), and then the term  $I_W$  is given by the relation

$$I_{W} = A_{t} \cdot \psi_{1} \cdot (\theta_{q} - \theta_{1}) \tag{2.9}$$

where

 $A_t$ =the total area of the surfaces bounding the enclosed space.

If the structures bounding the enclosed space consist of different materials, or if they differ in thickness, as is usually the case in practice, then the above-mentioned operations are carried out separately for each type of structure, and after that the term  $I_W$  is obtained from the expression

$$I_{W} = \sum_{j} I_{W, j} = \sum_{j} A_{j} \cdot \psi_{1, j} \cdot (\vartheta_{g} - \vartheta_{1, j})$$
(2.10)

The present section, which deals with the term  $I_W$ , is based in its entirety on the publications of Kawagoe and Sekine as well as Ödeen referred to in the above.

Term  $I_R$ 

Kawagoe and Sekine calculated the term  $I_R$  from the following formula, which is a generalization of the Stefan-Boltzmann law:

$$I_R = A \cdot (E_g - E_0) \tag{2.11}$$

where

A = the area of the opening,

$$E_g = 4.96 \cdot \left(\frac{9_g + 273}{100}\right)^4$$

$$E_0 = 4.96 \cdot \left(\frac{9_o + 273}{100}\right)^4$$

This formula is applicable to the whole duration of the process of fire development, and is used in its unchanged form for the calculations in the present publication.

# 3. Study of combustion gas flow and heat flow through openings in enclosed spaces

As has been shown in the section dealing with the term  $I_L$  in Chapter 2, the expression given by Eq. (2.4) is applicable only during the flame phase of the process of fire development. In that section, the term  $I_L$  was determined on the basis of a maximum rate of combustion. If this expression is to be extended so as to be valid for the cooling phase of the fire also, then this requires that the rates of gas and air flow through the openings in the enclosed space shall be determined directly, and that these rates shall then be used as a point of departure for determining the quantity of heat which is withdrawn per unit time from the enclosed space. Similar problems, which were defined in thermodynamic terms, and which related to fires in enclosed spaces, have been treated by Kawagoe [5], among others. The primary prerequisites to such a treatment are the assumptions that the pressure distribution in a vertical direction is linear, and that there exists a neutral layer or zone, i.e. a level at which the static pressure in the interior of the enclosed space is equal to the atmospheric pressure outside the enclosed space. From these assumptions, Kawagoe deduced the expression for the maximum rate of combustion,  $R_{\text{max}}$ , which is given by Eq. (2.3). On the assumption that the position of the neutral zone is the unknown variable, he used the Bernoulli equation to calculate the respective quantities of gases and air which flow out of and into the enclosed space per unit time. After that, he determined the position of the neutral zone from the condition that the rate of flow of the incoming gases shall be equal to the rate of flow of the gases which are consumed by combustion, and that the rate of flow of the outgoing gases shall be equal to the rate of flow of the gases which are produced by combustion. Finally, on the assumption that the quantity of air consumed per unit weight of fuel is known, he calculated the maximum rate of combustion,  $R_{\text{max}}$ .

In this chapter, a similar analysis will be carried out in what follows. The purpose of this analysis is to determine the term  $I_L$  by an expression which is more general than that given by Eq. (2.4). A detailed deduction will be presented for the case where the enclosed space is provided with one or several vertical openings which are equal in height. After that, we shall deal with the modifications which are required in the applications which involve openings of other types.

## Flow conditions in vertical rectangular opening in enclosed space

The interchange between the gaseous products of combustion and the combustion air takes place because the density of the hot gases is lower than that of the cold air outside the enclosed space. On the assumption that the temperature in the whole enclosed space is uniform, and that there exists a neutral zone, the velocities of gas and air flow can be determined theoretically. After that, if the dimensions of the opening are known, it is possible to calculate the respective rates of flow, i.e. the masses of the outgoing gases and the incoming air per unit time.

The velocity distribution in a vertical rectangular opening is schematically represented in Fig. 6.

The difference in static pressure between the outside and the inside is equal to zero at the level of the neutral zone. Accordingly, if use is made of the notations given in Fig. 6, the pressure difference,  $P_y$ , above the neutral zone i

$$P_{y} = (\rho_0 - \rho_q) \cdot y \tag{3.1a}$$

and the pressure difference,  $P_z$ , below the neutral zone is

$$P_z = -(\rho_0 - \rho_a) \cdot z \tag{3.1b}$$

where

 $\rho_0$  = density of the outside air,

 $\rho_g$  = density of the combustion gases.

The density of the gaseous products of combustion is assumed to be equal to the density of the air at the same temperature [10].

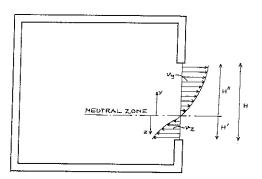


Fig. 6. Gas flow in an enclosed space provided with a vertical opening.

A = Area of the vertical opening in the enclosed space.

H = Height of the vertical opening in the enclosed space.

H'' = Height of the vertical opening above the neutral zone level.

H' = Height of the vertical opening below the neutral zone level.

 $v_y$  = Velocity of gas flow above the neutral zone level.

 $v_z$  = Velocity of gas flow below the neutral zone level.

From Bernoulli's theorem we obtain the following expressions for the variation in the velocity of flow with the distance from the neutral zone

$$v_{y} = \sqrt{2g \cdot y \cdot \frac{\rho_{0} - \rho_{g}}{\rho_{g}}} \tag{3.2a}$$

and

$$v_z = \sqrt{2g \cdot z \cdot \frac{\rho_0 - \rho_g}{\rho_0}} \tag{3.2b}$$

Then the rate of flow of the outgoing gases is

$$Q_{\text{out}} = \mu \cdot B \cdot \rho_g \cdot \int_0^{H''} v_y \cdot dy \tag{3.3a}$$

and the rate of flow of the incoming air is

$$Q_{\rm in} = \mu \cdot B \cdot \rho_0 \cdot \int_0^{H'} v_z \cdot dz \tag{3.3b}$$

where

 $\mu$  = the coefficient of contraction,

B = the width of the opening.

By substituting Eqs. (3.2a) and (3.2b) in Eqs. (3.3a) and (3.3b), respectively, we can directly calculate  $Q_{\text{out}}$  and  $Q_{\text{in}}$ . We find

$$Q_{\text{out}} = \frac{2}{3} \cdot \mu \cdot B \cdot (H'')^{3/2} \cdot \sqrt{2g \cdot \rho_g \cdot (\rho_0 - \rho_g)}$$
(3.4a)

and

$$Q_{\rm in} = \frac{2}{3} \cdot \mu \cdot B \cdot (H')^{3/2} \cdot \sqrt{2g \cdot \rho_0 \cdot (\rho_0 - \rho_g)}$$
 (3.4b)

The position of the neutral zone is determined by the equation of gas interchange in the enclosed space. This equation states that the difference between the rates of flow of the outgoing gases and the incoming air shall be equal to the difference between the rates of flow of the gases which are produced and consumed by combustion.

The mass of air contained in the enclosed space is assumed to be constant during the whole period of time under consideration. That total error, referred to the whole duration of the process of fire development, which is caused by this assumption in the calculation of heat flow is not greater than the heat content of the volume of air in the enclosed space. In comparison with the quantities of heat which are associated with fires, and in view of the other approximations which have been made in connection with the application of Eq. (2.1), this error may be regarded as negligible.

# Deduction of maximum rate of combustion determined by rate of air supply in accordance with [5]

On the assumption that the rate of combustion, R, is determined by the rate of air supply  $(R = R_{max})$ , we have

$$Q_{\text{out}} = R_{\text{max}} \cdot G_0 \cdot \rho_0 \tag{3.5a}$$

and

$$Q_{\rm in} = R_{\rm max} \cdot L \cdot \rho_0 \tag{3.5b}$$

where

 $G_0$  = the volume of combustion gases, in Nm<sup>3</sup>, produced by the combustion of 1 kg of fuel,

L = the volume of air, in Nm<sup>3</sup>, consumed by the combustion of 1 kg of fuel.

By substituting Eqs. (3.4a) and (3.4b) in Eqs. (3.5a) and (3.5b), respectively, we can calculate the position of the neutral zone, i.e. H' and H''. After that, by substituting H' in Eq. (3.5b), we obtain, for  $R_{\text{max}}$ , the expression

$$R_{\text{max}} = \kappa(\Delta \theta) \cdot A \cdot \sqrt{H}$$
 (3.6)

where

 $\kappa(\Delta \theta)$  is a coefficient, which depends on  $\Delta \theta$ 

and

$$\Delta \theta = \theta_g - \theta_0$$

The values of  $\kappa(\Delta \theta)$  is calculated for two fuels. First, for fires of the wood fuel type, which are most characteristic of actual fires, because wood usually constitutes the predominant fire load. Second, for fires, where the fuel consists of alcohol. The numerical values used in these calculations are given in what follows [10].

$$\mu$$
=0.7  
 $G_0$ =4.86 Nm<sup>3</sup> · kg<sup>-1</sup> for wood  
6.22 Nm<sup>3</sup> · kg<sup>-1</sup> for alcohol  
 $L$ =3.98 Nm<sup>3</sup> · kg<sup>-1</sup> for wood  
5.23 Nm<sup>3</sup> · kg<sup>-1</sup> for alcohol

The relation between  $\kappa$  and  $\Delta \vartheta$  is represented in Fig. 7. This graph shows that the variation in the value of  $\kappa$  with the temperature is very slight in the temperature range which is met with in fires. For practical applications, Kawagoe put the value of  $\kappa$  for wood fires equal to 330 kg·h<sup>-1</sup>·m<sup>-5/2</sup>, irrespective of the temperature. The value in question was used in deducing Eq. (2.3). This equation has been verified experimentally by several authors in model tests as well as in full-scale tests, see [5]. Eqs. (3.6) and (2.3) are applicable only to one or several openings of equal height, H. In these equations, A denotes the sum of the areas of the individual openings.

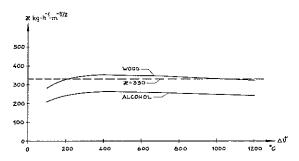


Fig. 7. Relation between the coefficient  $\varkappa$  in the equation  $R_{\max} = \varkappa A \sqrt[3]{H}$  and the temperature difference,  $\Delta \theta$ , between the combustion gases and the outside air.

Wood fires and alcohol fires.

# Determination of quantity of heat, $I_L$ , withdrawn per unit time through openings in enclosed space during the whole process of fire development

In order that an expression may be valid throughout the duration of the process of fire development, it is necessary to presuppose that the rate of combustion, R, may assume all values in the interval extending from zero to the maximum value which is given by Eq. (2.3). Therefore, R can be written

$$R = a \cdot 330 \cdot A \cdot \sqrt{H}$$
 kg · h<sup>-1</sup>

where

$$0 \le a \le 1$$

The balance between the respective rates of flow of the outgoing gases and the incoming air is given by the relation

$$Q_{\text{out}} - Q_{\text{in}} = G_0 \cdot R \cdot \rho_0 - L \cdot R \cdot \rho_0 \tag{3.8}$$

After substitution of  $Q_{\text{out}}$  and  $Q_{\text{in}}$  from Eqs. (3.4a) and (3.4b), respectively, and after simplification, we obtain

$$\left(\frac{H''}{H}\right)^{3/2} \cdot \sqrt{\rho_g(\rho_0 - \rho_g)} - \left(1 - \frac{H''}{H}\right)^{3/2} \cdot \sqrt{\rho_0(\rho_0 - \rho_g)} =$$

$$= \frac{(G_0 - L) \cdot \rho_0 \cdot 330a}{\frac{2}{3}\mu \cdot \sqrt{2g}} \tag{3.9}$$

With the help of Gay-Lussac's law of volumes for gases

$$\rho_0 = \rho_g \left( 1 + \frac{\Delta \vartheta}{273} \right) \tag{3.10}$$

the ratio H''/H can be determined for different values of a and  $\Delta\theta$  from Eq. (3.9).

The quantity of heat which is withdrawn per unit time from the enclosed space can be written

$$I_L = Q_{\text{out}} \cdot c_p \cdot \frac{\Delta \vartheta}{\rho_0} \tag{3.11a}$$

where

 $c_p$ =the specific heat of the outgoing gases, in kcal·m<sup>-3</sup>·°C<sup>-1</sup>. By substituting Eq. (3.4a) in Eq. (3.11a), we get

$$I_{L} = \varphi(\Delta \theta) \cdot c_{p} \cdot \Delta \theta \cdot A \cdot \sqrt{H}$$
(3.11b)

where

$$\varphi(\Delta \theta) = \frac{2}{3}\mu \sqrt{2g} \frac{\sqrt{\frac{\Delta \theta}{273}}}{1 + \frac{\Delta \theta}{273}} \cdot \left(\frac{H''}{H}\right)^{3/2}$$
(3.11c)

If the ratio H''/H is determined from Eq. (3.9), then Eq. (3.11c) yields  $\varphi(\Delta\theta)$  for different values of a and  $\Delta\theta$ . For combustion of wood fuel, this relation is represented in Fig. 8, which is based on the values of  $\mu$ ,  $G_0$ , and L given on p. 30. It is seen from this graph that the family of curves in question

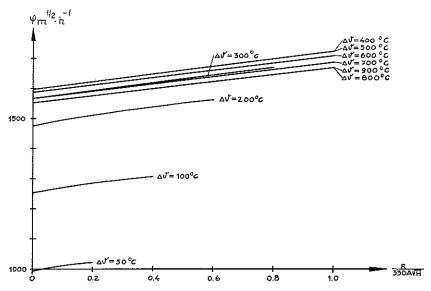


Fig. 8. Relation between the coefficient  $\varphi$  in the equation  $I_L = \varphi \cdot c_p \cdot \Delta \vartheta \cdot A \cdot \sqrt{H}$  and the ratio  $a = R/330 \cdot A \cdot \sqrt{H}$  for various values of the temperature difference  $\Delta \vartheta$ .

consists of approximately parallel straight lines. Therefore, we can write

$$\varphi = \varphi_0 + 120 \cdot a \tag{3.12}$$

where

 $\varphi_0$  is the value of  $\varphi$  for a=0.

Table 1 gives  $\varphi_0$  for various values of  $\Delta \theta$ .

For  $\Delta 9 > 300$  °C, the value of  $\varphi_0$  is nearly independent of the temperature, and may be assumed to range from 1500 to 1600.

If the rate of combustion, and hence the factor a, are known, then  $I_L$  can be calculated for any instant, t, during the process of fire development. The assumptions which have been chosen for the deduction of the expression for the term  $I_L$  will be briefly discussed in what follows.

In spite of the fact that steady-state conditions have been assumed in the calculation of the flow through the opening, the results obtained from this calculation can also be applied under non-steady-state conditions, because a change in a position of equilibrium almost immediately gives rise to the establishment of a new position of equilibrium [7].

Table 1. Relation between  $\varphi_0$  and  $\Delta\vartheta$ .

 Δϑ, °C	$\phi_0,$ $m^{1/2} \cdot h^{-1}$	Δϑ, °C	$\varphi_0,$ $m^{1/2} \cdot h^{-1}$	
10	515	500	1597	
50	991	600	1587	
100	1254	700	1567	
200	1476	800	1551	
300	1567	900	1552	
400	1595	1000	1510	

The calculation of  $I_L$  requires that the specific heat of the outgoing gases shall be known. The air content of these gases is dependent on the rate of combustion. When a=0, the outgoing gases consist of air alone, and when a=1, they consist of gaseous products of combustion only. However, the difference in the specific heat between the air and the combustion gases is very slight, and the specific heat may therefore to a close approximation be regarded as independent of the rate of combustion.

Furthermore, in the calculation of  $\varphi$ , it is assumed that the values of  $G_0$  and L remain constant during the whole process of fire development, irrespective of the rate of combustion. This assumption has not been verified by any physical considerations, but if we examine the right-hand member of Eq. (3.9), then we find that the effect produced by an error in the difference between  $G_0$  and L on the value of  $\varphi$  is comparable to that of an error in a, that is to say, this effect is very slight.

In his treatment which relates to the flame phase only, Kawagoe has assumed that the temperature is uniform in the whole enclosed space. In the present publication, this assumption has been extended so as to be applicable during the whole process of fire development. The assumption that the variation in temperature in a vertical direction is relatively slight when the intensity of the fire decreases has been confirmed by the full-scale tests which are described in Chapter 5 of this publication. In most of these tests, the dispersion in the temperature measured at different points in the enclosed space during the cooling phase was found to be smaller than during the flame phase.

# Modification of treatment in cases where enclosed spaces are provided with several openings which differ in height

The deduction carried out in the preceding section is applicable only in the cases where the air is supplied to the enclosed space through one or several openings which are equal in height, and which have a common neutral zone. If the enclosed space is provided with several openings which differ in height, then a corresponding deduction can be made in each individual case. For the determination of the maximum rate of combustion,  $R_{\text{max}}$ , Yokoi [13] has described an approximate method. It consists in the determination of a fictitious air flow factor, which is used in the original formula, Eq. (2.3). The fictitious air flow factor is determined from the expression

$$(A \cdot \sqrt{H})_{\text{fict}} = \sum_{i} A_{i} \cdot \sqrt{H_{i}}$$
(3.13)

For some cases, Kawagoe [14] has compared the values of  $R_{max}$  which were obtained from Eq. (3.13) with those which were determined by means of accurate calculations. He found that Eq. (3.13) gives values which are sufficiently accurate for practical uses when the differences in the height and in the vertical position of the openings are not too great. The Swedish Building Regulations 1967 recommended another acceptable approximation, namely, that a weighted average of the heights of the individual openings should be used as a value of H in the calculation of the air flow factor.

After a fictitious value of the air flow factor has been determined from Eq. (3.13), this value can be used instead of  $A \cdot \sqrt{H}$  in Eq. (3.11b) where the factor  $\varphi$  is calculated from Eq. (3.12) as before.

## Modification of treatment in cases where enclosed spaces are provided with horizontal openings in roofs

In the preceding two sections, it was assumed that all openings in the enclosed space are vertical. In the present section, we shall expound a more general theory which makes it possible to take account of the presence of

horizontal openings in the roof of the enclosed space. We suppose the enclosed space to be in conformity with Fig. 9.

We assume that there exists a neutral zone in the enclosed space, and that the level of this zone is not higher than the upper edge of the vertical opening, and not lower than its lower edge. A condition which is prerequisite to this assumption will be stated further on in the present section. For the vertical opening, the rate of flow of the outgoing gases,  $Q_{\text{out}}$ , is obtained from Eq. (3.4a), and the rate of flow of the incoming air,  $Q_{\text{in}}$ , is computed from Eq. (3.4b). The velocity of the gases which flow out through the horizontal opening is (Bernoulli's theorem)

$$v_h = \sqrt{2gh' \cdot \frac{\rho_0 - \rho_g}{\rho_g}} \tag{3.14}$$

This formula has been verified experimentally in connection with studies of venting fires [7]. The rate flow of the outgoing gases through the horizontal opening is

$$Q_h = \mu \cdot A_h \cdot \nu_h \cdot \rho_g \tag{3.15}$$

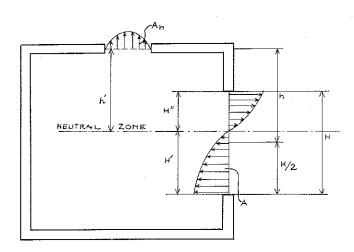


Fig. 9. Gas flow in an enclosed space provided with a vertical opening and a horizontal opening.

A =Area of the vertical opening.

 $A_k$  = Area of the horizontal opening.

H =Height of the vertical opening.

H'' = Height of the vertical opening above the neutral zone level.

H' = Height of the vertical opening below the neutral zone level,

h = Vertical distance from the centre of the vertical opening to the level of the horizontal opening.

h' = Vertical distance from the neutral zone to the level of the horizontal opening.

At a maximum rate of combustion,  $R = R_{\text{max}}$ , the equation of mass balance of gases requires that

$$Q_{\text{out}} + Q_h = G_0 \cdot R_{\text{max}} \cdot \rho_0 
Q_{\text{in}} = L \cdot R_{\text{max}} \cdot \rho_0$$
(3.16)

These two relations can be used to determine the position of the neutral zone which is modified in view of the presence of the horizontal opening, and then the maximum rate of combustion,  $R_{\text{max}}$ , can be calculated.  $R_{\text{max}}$  is a function of the term  $\frac{A_h \cdot \sqrt{h^2}}{A \cdot \sqrt{H}}$  at a given temperature. The value of  $R_{\text{max}}$  varies slightly with the temperature. If we write

$$R_{\text{max}} = 330 \cdot (A \cdot \sqrt{H})_{\text{fict}} \tag{3.17}$$

then  $(A \cdot \sqrt{H})_{\text{fict}}$  can be determined from the alignment chart in Fig. 10, which is entered at the value of  $A_h \cdot \sqrt{h}/A \cdot \sqrt{H}$ .

It is to be expected that the value of  $(A \cdot \sqrt{H})_{\text{fict}}$  determined in this manner may be used to an adequate degree of accuracy in the same way as the air flow factor,  $A \cdot \sqrt{H}$ , to characterize a fire. Accordingly, for an enclosed space with horizontal openings in the roof, the opening factor is given by the expression  $(A \cdot \sqrt{H})_{\text{fict}}/A_t$ , where  $(A \cdot \sqrt{H})_{\text{fict}}$  is determined from the alignment chart in Fig. 10. The term  $I_L$  in the equation of heat balance, Eq. (2.1), of the process of fire development is obtained by analogy with Eq. (3.11b) from

$$I_{L} = \varphi(\Delta \theta) \cdot c_{p} \cdot \Delta \theta \cdot (A \cdot \sqrt{H})_{\text{fict}}$$
(3.18)

where  $\varphi(\Delta \theta)$  is given by Eq. (3.12).

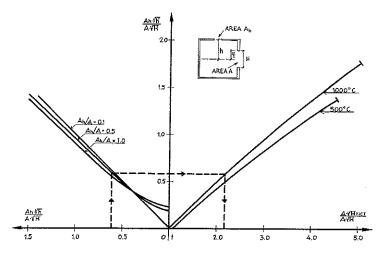


Fig. 10. Alignment chart for the calculation of the value of the modified air flow factor,  $(A\sqrt{H})_{\text{fict}}$ , on the basis of known geometrical data. For notations, see Fig. 9.

This procedure presupposes that the flow through the horizontal opening in the roof is not predominant. Consequently, the factor  $A_h \cdot \sqrt{h'}/A \cdot \sqrt{H}$  has an upper limit at which the above model of the flow conditions ceases to be relevant. This upper limit is

$$\frac{A_h \cdot \sqrt{h'}}{A \cdot \sqrt{H}} = \begin{cases} 1.76 \text{ at } 1000^{\circ}\text{C} \\ 1.37 \text{ at } 500^{\circ}\text{C} \end{cases}$$

At this limit, the neutral zone is on a level with the upper edge of the vertical opening, and then h' is identical with the vertical distance from the level of the horizontal opening in the roof to the upper edge of the vertical opening. Tests [7] have indicated that the model used is relevant up to this upper limit. For values of  $A_h \cdot \sqrt{h'}/A \cdot \sqrt{H}$  which are higher than this limit, all gaseous products of combustion will be vented through the horizontal opening in the roof. This is sometimes intentional if it is desired that the spread of fire to adjoining rooms should be prevented by venting the fire [15]. If the air and the combustion gases flow in the main through horizontal openings, then the flow becomes unstable, and it is difficult to represent it by a simple theoretical model [16].

Example showing how to use alignment chart in Fig. 10

Calculate the maximum rate of combustion,  $R_{\text{max}}$ , during the flame phase at 1000°C in the enclosed space characterized by the following data:

$$A = 2 \text{ m}^2$$
,  $H = 1 \text{ m}$ ,  $A_h = 1 \text{ m}^2$ ,  $h = 1.5 \text{ m}$ .  
 $\frac{A_h \cdot \sqrt{h}}{A \cdot \sqrt{H}} = \frac{1 \cdot \sqrt{1.5}}{2 \cdot \sqrt{1}} = 0.61$ ;  $\frac{A_h}{A} = 0.5$ 

The dash line in the alignment chart gives

$$(A \cdot \sqrt{H})_{\text{flet}} = 2.18 \cdot A \cdot \sqrt{H} = 4.36 \text{ m}^{5/2}$$
  
 $R_{\text{max}} = 330 \cdot 4.36 \text{ kg} \cdot \text{h}^{-1} = 1440 \text{ kg} \cdot \text{h}^{-1}$ 

This alignment chart can also be used for enclosed spaces where the horizontal opening in the roof is replaced by a ventilation duct. In such cases, the height h is replaced by the height of the gas column (the static head), with the reduction of the losses due to friction, which can be expressed in terms of the equivalent loss in static head. In an ordinary flat or office equipped with common ventilators made of non-combustible materials, the effect of the ventilation ducts is usually negligible.

If the enclosed space is ventilated through air inlets and outlets by means of a fan installation, then the corresponding fictitious air flow factor can be calculated in an analogous manner. If the quantity of gases exhausted per

unit time is  $Q_{\text{out,}}^{\nu}$  in kg·h<sup>-1</sup>, and the quantity of air supplied per unit time is  $Q_{\text{in,}}^{\nu}$  in kg·h<sup>-1</sup>, then, for  $R = R_{\text{max,}}$  we have the conditions

$$\begin{array}{c}
Q_{\text{out}} + Q_{\text{out}}^{\nu} = G_0 \cdot R_{\text{max}} \cdot \rho_0 \\
Q_{\text{in}} + Q_{\text{in}}^{\nu} = L \cdot R_{\text{max}} \cdot \rho_0
\end{array} \right}$$
(3.19)

from which  $R_{\max}$  and  $(A\cdot \sqrt{H})_{\mathrm{fict}}$  can be determined by analogy with the above.

### 4. Description of programme for digital computer

The integration of the system of equations given by Eq. (2.6) was carried out by using the Runge-Kutta method in a modified form which has been suggested by Merson [22]. This modified method enables the computer to choose that interval of integration,  $\Delta t$ , which is required in order to ensure a certain definite degree of accuracy. To integrate the above-mentioned system of equations in the time interval from t to  $t+\Delta t$ , this system was evaluated in five individual operations. A determination of the temperature of the combustion gases is required for each one of these operations. According to Eqs. (2.1), (2.10), (2.11), and (3.11), the equation of heat balance of the process of fire development can be written

$$I_C = I_L + I_W + I_R \tag{4.1}$$

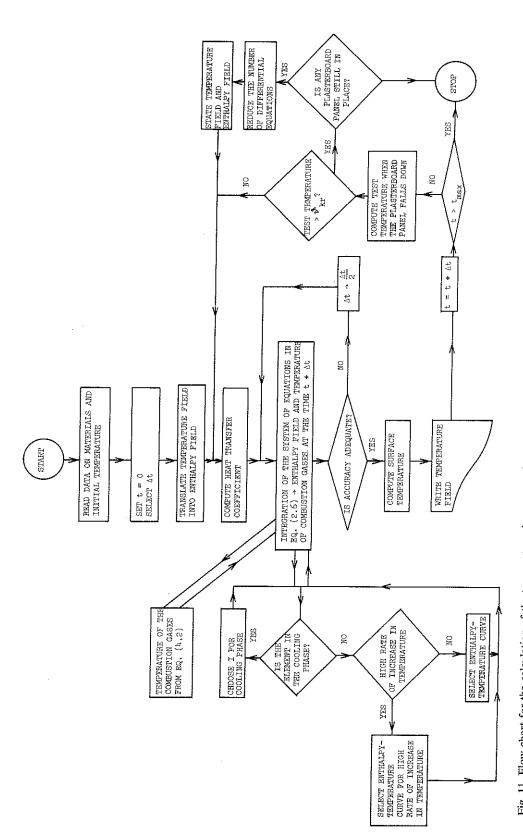
where

$$\begin{split} I_L &= \varphi(\Delta \vartheta) \cdot c_p \cdot A \sqrt{H}(\vartheta_g - \vartheta_0) \\ I_W &= \sum_j I_{W,j} = \sum_j A_j \psi_{1,j}(\vartheta_g - \vartheta_{1,j}) \\ I_R &= A \cdot (E_g - E_0) \end{split}$$

If  $I_R$ ,  $\varphi$  and  $c_p$  are calculated on the basis of that value of the temperature of the combustion gases which has been obtained from the next preceding determination, then Eq. (4.1) can be solved for the temperature of the combustion gases,  $\vartheta_q$ , in an explicit form

$$\vartheta_{g} = \frac{I_{C} + \varphi(\Delta \vartheta) \cdot A \sqrt{H} \cdot \vartheta_{0} + \sum_{j} A_{j} \cdot \psi_{1, j} \cdot \vartheta_{1, j} - I_{R}}{\varphi(\Delta \vartheta) \cdot c_{p} \cdot A \sqrt{H} + \sum_{j} A_{j} \psi_{1, j}}$$
(4.2)

which can be substituted in the system of equations represented by Eq. (2.6). The programme for the computer has been prepared in such a way as to be applicable to enclosed spaces bounded by structures which were assumed to be of up to three different types. Two of these structures were supposed to be homogeneous, whereas the third might be divided into two or three layers consisting of different materials, e.g. plasterboard panels, mineral wool, and brick.



enclosed space are supposed to comprise, first, a roof or/and ceiling and a floor made of homogeneous material, and second, partitions which consist of a Fig. 11. Flow chart for the calculation of the temperature-time curve of the process of fire development in an enclosed space. The structures which bound the load-bearing frame made of steel studs and covered on the inside as well as on the outside with two plasterboard panels, 13 mm in thickness each.

The specific heat of most structural materials, e.g. concrete having a non-negligible moisture content or plasterboard panels, varies discontinuously with the temperature when these materials are subjected to physical or chemical transformations. Therefore the programme used the enthalpy I, in kcal·m<sup>-3</sup>, of the different materials as a dependent variable in the calculation of the temperature fields.

Fig. 11 represents a schematic flow chart which shows the programming procedure in the calculation of the combustion gases in an enclosed space where the floor and the roof or/and ceiling are made of concrete, while all the walls consist of a lightweight frame built of steel which are covered on the inside as well as on the outside with two plasterboard panels, 13 mm in thickness each. This type of wall exhibits two characteristic features, viz., first, experiments have shown that a plasterboard panel exposed to fire disintegrates when the temperature of the panel has reached a certain definite value, and second, the relation between the temperature and the enthalpy is dependent on the rate of temperature rise. Furthermore, it was necessary to choose different enthalpy-temperature curves depending on whether the temperature of the plasterboard panel in question was assumed to be increasing or decreasing. For further particulars, reference is made to the description of the calculations for the Type G enclosed space in Chapter 7.

## 5. Calculation of time graphs of rate of combustion for some full-scale tests described in literature

As has been mentioned in Chapter 2, the present chapter will deal with the comparative calculations which have led to a determination of the variation in the rate of combustion, expressed in kcal·h<sup>-1</sup>, with the time. The method of successive approximations employed for this purpose consisted in making calculations which were based on different forms of the time graph of the rate of combustion. These calculations were repeated until they resulted in that curve which corresponded to the closest agreement between experimental and calculated curves representing the variation in the temperature of the combustion gases with the time. This method required a certain systematization of the description which represents the variation in the quantity of energy released by combustion with the time. What can be assumed to be known to a sufficient degree of accuracy in this connection is solely the total quantity of energy that can be liberated during the whole process of fire development, i.e. the fire load. Furthermore, it can be assumed that Eq. (2.3)

$$R_{\text{max}} = 330 A \sqrt{H}$$
 kg of wood per h

expresses the maximum rate of combustion, in kg of wood per h. In order that a theoretical determination of the temperature of the combustion gases may be possible, it is moreover necessary to determine the relation between released energy and weight loss of fuel, W, in

$$I_C = R \cdot W \quad \text{keal} \cdot h^{-1}$$

This determination is rendered difficult by the fact that the combustion of the gases formed by pyrolysis and that of solid wood fuel constitute a complicated process which involves a series of chemical reactions in different phases, see [4]. In respect of energy conditions, some of these reactions are endothermic, others exothermic. So far as the Authors know, no systematic investigation has been made up to now in order to carry out a quantitative analysis of the liberation of energy during the individual phases of the process of fire development.

These considerations have necessitated certain assumptions which concern the form of the curve showing the variation in the rate of combustion with the time. These assumptions are stated in what follows. The quantity of energy

liberated per unit time during the ignition phase was supposed to increase according to a polygonal function of the time to a level which corresponds to the rate of combustion during the flame phase. The determination of this level was based on Eq. (2.3),  $R_{\rm max} = 330 A \sqrt{H}$  kg of wood per h. When the quantity of energy released per kg of wood fuel during the flame phase was assumed to range from 2500 to 2800 kcal, it was found that the calculated temperature-time curves were closely in agreement with the results of the full-scale tests in respect of the maximum temperature and the duration of the flame phase. In order to adapt these assumptions to the temperature-time curves which have been published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures", and which are used in Sweden as a basis for calculations of the fire endurance of structural components in conformity with standard specifications, the quantity of heat evolved during the flame phase was supposed to be 2575 kcal · kg<sup>-1</sup> of wood, since this value had been employed for calculating the above-mentioned temperature-time curves. Cf. Chapter 2. The value in question has originally been stated by Kawagoe, who used it as a measure of the quantity of energy that is liberated by incomplete combustion of 1 kg of wood. By analysing the composition of the combustion gases in fire tests, see e.g. [5], it has been found that they contain considerable quantities of carbon monoxide, and this indicates that the combustion is not complete. It seems that the analysis of the combustion gases was performed during the flame phase, and the value 2575 therefore applies to this phase only. During the cooling phase, the weight loss of fuel per liberated energy unit is considerably less than during the flame phase. This means, that, when the whole process of fire is considered, the energy released by the combustion of 1 kg of wood must be higher than 2575 keal even if the combustion is incomplete. As a rule, there was scant basis for an accurate prediction of how the rate of completeness of combustion varied during the different phases. Consequently, in our comparative analysis the nominal heat value of wood, ranging from 3500 to 4500 kcal · kg<sup>-1</sup>, was used as the energy liberated during the whole course of fire. This means that combustion is assumed to be complete throughout the present publication. The resulting time-temperature curves for the combustion gases therefore as a rule are found to be more in agreement with the maximum temperature curves than with the mean temperature curves obtained in the full scale tests.

Consequently, during the flame phase, the value of the rate of combustion will be constant,  $330 \cdot A\sqrt{H} \cdot 2575$  kcal·h<sup>-1</sup>. The cooling phase is characterized by a rate of combustion which decreases in conformity with a polygonal time graph in such a way that the slope of each individual side of the polygon is dependent on the duration of the flame phase. In every case, the area between the rate of combustion curve and the time axis must equal the fire load, expressed in energy units. The use of a polygonal time graph of the rate of

combustion makes it easier to check the agreement between the total quantity of energy liberated during the process of fire development and the fire load which is given from the outset.

The factors which must be taken into account in a comparative theoretical calculation of the temperature of the combustion gases in fire tests are enumerated in what follows.

- (1) Characteristics of fuel: Quantity, moisture content, porosity factor (hydraulic radius), distribution in the enclosed space.
- (2) Geometric characteristics of the enclosed space: Opening factor and its variation with the time (e.g. in the case where the fire burns through a door), shape of openings, cross-sectional area of ventilation ducts (if any).
- (3) Characteristics of the structures bounding the enclosed space: Structural design, thermal properties and temperatures of disintegration (if any) of the materials entering into the structures, emissivity characteristics of the surfaces.

By studying the available literature on full-scale wood fire tests, it was found, first, that the number of published tests is relatively small, and second, that those data which are so detailed as to render possible a comparative theoretical calculation have been stated only for a few tests out of this number. For this reason, it was necessary to confine the comparative calculations to only four test series reported in the literature or in other sources. These test series are enumerated below.

- (1) Test series A. Tests carried out by Sjölin, and dealt with in a thesis for a L. Techn. degree, [17].
  - (2) Test series B. Tests made by Kawagoe, and published in [5].
- (3) Test series C. Tests performed by Ödeen at the Royal Institute of Technology, Stockholm, and described in his doctoral thesis [18].
- (4) Test series D. Tests directed by *Pettersson* and *Ödeen*, which were carried out in a test house of the Atomic Energy Co., Ltd. (AB Atomenergi) at Studsvik, Sweden.

## Test series A. Tests carried out by Sjölin. Calculation of time graphs of rate of combustion

The tests made by Sjölin were undertaken in order to study the spread of fire, and the process of fire development, in rooms and combinations of rooms exposed to ignition at a single point by heat radiation emitted by an explosion of nuclear weapons. The fire tests were carried out in a test house which was provided with concrete floor structures and concrete or lightweight concrete

<sup>&</sup>lt;sup>1</sup> Not until too late in the publication of this paper did the authors learn about the full scale fire tests carried out at Fire Research Station, Boreham Hood, London [11, 12].

This is the more regrettable as these experiments make an excellent basis for comparative theoretical calculations.

walls. The test house was designed in such a way that various enclosed spaces might be formed by individual rooms, or by several rooms connected together. The available model scales were 1 to 1, 1 to 2, and 1 to 4. The variables recorded in these tests were expressed as functions of the time, and comprised the temperature and the velocity of gas flow at characteristic points, the intensity of radiation, the composition of the gases, and the rate of combustion, which was determined by continuous weighing of the quantity of fuel in the enclosed space.

The fire load in all these tests consisted of authentic furniture. This was an extraordinarily valuable feature of the tests, seeing that all the other full-scale tests which are dealt with in the present publication were made by using fire loads of the wood crib type.

Seven of the fire tests included in this test series were found to be suited for the present theoretical study. In the other tests, the ignition did not cause the fuel to take fire. Table 2 shows the scope of the seven tests under consideration.

Table 2. Test series A.

Test No.	Type of room	Window area, m <sup>2</sup>	Opening factor, $ \frac{A \cdot \sqrt[4]{H}}{A_t} $ $ m^{1/2} $	Fire load, kg·m <sup>-2</sup> of bounding surface area	Remarks
1	В	1.16	0.0237-0.06	3.5	(The fire burnt through a door, 1.6
2	В	1.16	0.0237-0.06	4.4	$m^2$ in area, during the time interval from $t=0$ min. to $t=6$ min.
3	L	1.16	0.0160-0.0356	4.9	(The fire burnt through a door, 1.6
4	L	1.88	0.0278-0.0486	5.6	$m^2$ in area, during the time interval from $t=8$ min. to $t=12$ min.
5	L	1.88	0.0548	5.0	•
6	L	2.95	0.068	5.7	
7	B+L	3.20	0.040	8.1	

B = Furnished two-person bedroom, 10.4 m<sup>2</sup> in floor area.

L = Furnished living room, 18.8 m<sup>2</sup> in floor area.

B+L =Combination of B and L.

The curves representing the variation in the temperature of the combustion gases with the time in the above-mentioned seven tests, as well as the corresponding curves obtained by calculations with the help of automatic computer, are reproduced in Appendix 1. In the test No. 7, the enclosed space consisted of two contiguous rooms, which communicated through an open door. The partition between these two rooms was considered in the theoretical calculations to be an enclosed structure which possessed a heat-absorbing capacity.

For this test, the curves representing the temperature of the combustion gases as a function of the time are shown separately for the bedroom and the living-room.

All the fires in these tests were characterized by a protracted process of ignition, which was followed by a rapid transition to the flame phase. For the calculated curves, the time was put equal to zero at the instant when the flame phase began, i.e. when the fuel took fire, in the actual fire tests. The time graphs of the rate of combustion which were finally obtained from the calculations are also reproduced in the respective diagrams. In the tests Nos. 1 to 4, the opening factor was changed during the process of fire development because the fire had burnt through a door. This change was taken into account in the calculations, and constituted the cause of the somewhat unusual shape of the time graphs of the rate of combustion which refer to these four tests. In these cases, the curves were plotted in such a way as to relate the rate of combustion to that value of the opening factor which was obtained after the fire had burnt through the door. In the test No. 3, the fact that the fire has burnt through a door is reflected very clearly in the curve, which shows that the temperature of the combustion gases remained constant at 500 to 600°C, and then rapidly rose to about 800°C when the fire burnt through the door.

Since the calculations were based on the opening factor and on the total energy content of the fuel, it was possible to choose the time graphs of the rate of combustion in such a manner that the agreement between the observed and calculated time graphs of the temperature of the combustion gases was very close in all the tests except the test No. 7. In the test No. 7, the calculated time graph of the combustion gas temperature was compared with the corresponding observed curves for the living-room as well as for the bedroom. The agreement between these curves was relatively close in the second case, but not in the first, where the curve is slightly displaced in time with reference to the curve for the bedroom. Moreover, when use was made of the opening factor which was determined geometrically, the temperatures obtained for the flame phase were found to be somewhat too low. However, it is not correct to regard the above-mentioned two rooms as a single enclosed space, since there existed quite a considerable difference in temperature between these rooms. Nor are the two rooms in question to be regarded as two separate enclosed spaces, since a certain heat exchange took place between these rooms.

## Test series B. Tests made by Kawagoe [5]. Calculation of time graphs of rate of combustion

In [5], Kawagoe has described a large number of fire tests which had been carried out in Japan. In this investigation, he primarily studied the relation between the reduction in the weight of fuel per unit time and the dimensions

of the openings in the enclosed space. Among other things, he also deduced the equation which is reproduced in Eq. (2.3) in the present publication. The variables measured in these tests were the reduction in the weight of fuel, the temperature of the combustion gases, the composition of the combustion gases, the gas velocities, the intensity of radiation, and the pressure distribution in the window openings. On account of the above-mentioned main purpose of Kawagoe's investigation, it is only in three of these tests that the results of measurements, the geometric data, and the data on the materials entering into the structures which bounded the enclosed space are presented in such a way as to make it possible to carry out theoretical calculations of the type under consideration. These three tests were performed in a test house which was provided with walls made of hollow concrete blocks and with concrete floor and roof structures. The other test data are given in Table 3.

Table 3. Test series B.

	Fire load,	Bounding surface	Window dimensions,	Opening $fac_t$ or, $A \cdot \sqrt{H}$
Test No.	wood	area, m²	width $\times$ height, $m \times m$	$\frac{A_t}{m^{1/2}}$
1	400	48	0.93×1.8	0.0467
2	900	48	$0.93 \times 1.8$	0.0467
3	1000	48	$0.93 \times 1.8$	0.0467

Each hollow concrete block used for the walls comprised a single large cavity, without any subdivisions. The volume of the cavity was estimated at 30 to 40 per cent of the total volume of the block. In the calculations, the walls were considered to be composed of two different structures. One of them consisted of concrete alone, while the other comprised three layers, viz., concrete, air-filled cavity, and concrete, respectively. The second structure represented that part of the wall surface which corresponded to the cavities of the hollow concrete blocks, while the first structure was equivalent to the remaining part of the wall surface.

The observed and calculated temperature-time curves, as well as the time graphs of the rate of combustion used in the calculations, are reproduced in Appendix 2. The fire loads in the tests Nos. 2 and 3 were relatively high and the duration of the fires in these tests was therefore long. In order that the calculated values should agree with the values observed in these tests, it was necessary to choose a comparatively flat slope for the ascending branch of the time graph of the rate of combustion.

The data on the cooling phases in the tests Nos. 2 and 3 reported in [5] are not complete, and this is the reason why the curves relating to these tests break off at such an early stage.

### Test series C. Tests published by Ödeen [18]. Calculation of time graphs of rate of combustion

The tests described by Ödeen in [18] were carried out in a tunnel building of an approximately semi-circular shape, which had been specially constructed for this purpose. It was provided with a concrete wall, 20 cm in thickness, its total bounding surface area was 75 m², and its total enclosed volume was 46 m³. A fan system made it possible to regulate and to measure the quantity of air which was supplied per unit time of the fire. In addition, a vent, 0.5 m² in cross-sectional area, for conveying the combustion gases to the outside air was provided in the upper part of each end wall of the tunnel.

A series of fire tests using fir wood as fuel has been carried out in this test building. A study was made of the effects produced on the process of fire development by the factors which are enumerated in what follows.

- (1) The volume of air supplied per unit time to the tunnel, Q, in  $m^3 \cdot s^{-1}$ .
- (2) The quantity of combustible material (fire load), M, in kg of wood.
- (3) The hydraulic radius of the fuel, r, in cm. This factor expresses the ratio of the total volume of the fuel to its total bounding surface area.

The scope of the test series using wood fuel is shown in Table 4, which was extracted from [18].

In order that the results of these tests may be compared by means of calculations with those of fire tests in ordinary enclosed spaces, where the rate

Table 4. Test series C.

Test No.	Fire load, M, kg of wood	Rate of air supply, $Q$ , $m^3 \cdot s^{-1}$	Hydraulic radius, r, cm	Moisture content of fuel, per cent	Energy content of fuel, Mcal
1	270	1.0	_	9	1129
2	675	2.0	1.0	17	2565
3	675	1.0	1.0	17	2565
4	675	0.7	1.0	22	2468
5	675	1.5	1.0	22	2468
6	675	MIN	1.0	21	2501
7	675	1.0	1.7	21	2501
8	675	0.7	1.7	21	2501
9	675	1.0	0.6	21	2501
10	270	0.7	1.0	21	1000
11	405	0.7	1.0	22	1481
12	405	0.7	2.4	28	1440
13	135	0.7	1.0	28	481
14	945	1.0	1.6	16	3659
15	1350	2.0	1.4	17	5130
16	405	0.7	0.4	17	1539

of air supply is determined by the openings in the enclosed space, the quantity of air supplied per unit time, Q, must be converted into an air flow factor,  $A\sqrt{H}$ , or into an opening factor,  $A\sqrt{H}/A_t$ . In this connection, it is necessary to take account of the air flow which may possibly enter into the tunnel through the outlets for combustion gases at low values of Q, and may therefore increase the value of the air flow factor.

In an ordinary enclosed space provided with a vertical opening, the rate of flow of the incoming air,  $Q_{in}$ , is given according to Eq. (3.4b), by the relation

$$Q_{\rm in} = 2/3 \ \mu B(H')^{3/2} \sqrt{2g \cdot \frac{(\rho_0 - \rho_g)}{\rho_0}} \ \text{m}^3 \cdot \text{s}^{-1} =$$

$$= 2/3 \cdot 0.7 \cdot A \sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left( \frac{\Delta \theta}{273} - \frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 3/3 \cdot 0.7 \cdot A \sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left( \frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$

$$= 2/3 \cdot 0.7 \cdot A \sqrt{H} \sqrt{2g} (H'/H)^{3/2} \left( \frac{\Delta \theta}{1 + \frac{\Delta \theta}{273}} \right)^{1/2}$$
(5.1)

The notations used in Eq. (5.1) have been explained in Chapter 3. With the help of Eqs. (3.9) and (3.10), we obtain, for a=1, the relation between H''/H and  $\Delta \theta$  shown in Table 5.

Table 5. Relation between H''/H and  $\Delta\vartheta$ .

Δθ	H''/H	Δϑ	H''/H	
400	0.61	800	0.64	
500	0.62	900	0.65	
600	0.63	1000	0.66	

Since H' = H - H'', we can substitute H'/H in Eq. (5.1). For different values of the temperature  $\vartheta_g = \vartheta_0 + \Delta \vartheta$ , this yields the simplified expression

$$Q_{\rm in} = K \cdot A \sqrt{H} \tag{5.2a}$$

where K varies with the temperature of the combustion gases,  $\theta_g$ , in conformity with Table 6, where  $\theta_0$  was put equal to zero.

Table 6. Relation between K and  $\vartheta_a$ .

 $\vartheta_g$	K	$\vartheta_g$	K
400	0.40	900	0.38
500	0.40	1000	0.38
600	0.40	1100	0.37
700	0.39	1200	0.37
800	0.39		

In the case under consideration, the rate of flow of the incoming air,  $Q_{\rm in}$ , was equal to the quantity of air supplied per unit time by the fan system. Hence it follows that this quantity of air can be replaced by a fictitious air flow factor,  $(A\sqrt{H})_{\rm fict}$ , by means of Eq. (5.2a).

If the constant K is put equal to its value at 600 °C, then we obtain

$$Q = 0.40(A\sqrt{H})_{\text{fict}} \tag{5.2b}$$

By substituting  $A_t = 75 \text{ m}^2$ , we get

$$(A\sqrt{H}/A_t)_{\text{fict}} = 0.033Q \tag{5.3}$$

This relation between the fictitious opening factor,  $(A\sqrt{H}/A_t)_{\text{fict}}$ , and the rate of air flow, Q, shall include the effect of the air which may possibly enter into the tunnel through the paraboliform outlets for combustion gases, 0.5 m<sup>2</sup> in cross-sectional area each. In the calculations, these outlets were assumed to be replaced by two rectangular openings having a base B=1.13 m and a height H=0.44 m each.

If the rate of flow of the incoming air through these openings is put equal to  $Q_{\rm in}$ , then  $Q_{\rm in}$  at a maximum rate of combustion is given by the equation

$$Q_{\rm in} + Q\rho_0 = LR_{\rm max}\rho_0$$

For notations in this equation, see Chapter 3. By substituting  $Q_{in}$  from Eq. (3.4b), we find

$$2/3 \sqrt{2gH' \frac{\rho_0 - \rho_g}{\rho_0}} \mu B \cdot H' \rho_0 + Q \rho_0 = L R_{\text{max}} \rho_0$$
 (5.4)

where  $R_{\text{max}}$  is determined from Eq. (3.5a).

$$Q_{\text{out}} = G_0 \cdot \rho_0 \cdot R_{\text{max}}$$

or

$$2/3 \cdot \sqrt{2gH''} \frac{\rho_0 - \rho_g}{\rho_g} \cdot \mu \cdot B \cdot H'' \cdot \rho_g = G_0 \cdot \rho_0 \cdot R_{\text{max}}$$
 (5.5)

By substituting B=1.13 m, H=0.44 m,  $\mu=0.7$ ,  $G_0=4.86$  Nm<sup>3</sup>·kg<sup>-1</sup>, L=3.98 Nm<sup>3</sup>·kg<sup>-1</sup> and  $R_{\text{max}}$  from Eq. (5.5) into Eq. (5.4) we obtain

$$0.83(H'/H)^{3/2} + 0.722Q = 0.38(H''/H)^{3/2}$$
(5.6)

In the calculation of  $\rho_g$ , the temperature of the combustion gases was assumed to be 600°C.

For those values of Q which are of interest in this connection, i.e.  $Q \ge 0.7$  m<sup>3</sup> · s<sup>-1</sup> in conformity with Table 4, Eq. (5.6) has no solution in the interval  $0 \le H'/H \le 1$ . This means that the flow is possible in an outward direction only.

In the test No. 6, the fans were switched off, and the exchange of air took place through the outlets for combustion gases alone. If these outlets are supposed, as before, to be approximately represented by rectangular openings having a base of 1.13 m and a height of 0.44 m each, then we obtain an opening factor  $A\sqrt{H}/A_t = 0.0088$  m<sup>1/2</sup>, which corresponds to a rate of air flow  $Q_{\min} = 0.26$  m<sup>3</sup>·s<sup>-1</sup> according to Eq. (5.3). Cf. the average value of the rate of air flow, 0.19 m<sup>3</sup>·s<sup>-1</sup>, which has been computed by Ödeen for the whole process of fire development.

Thus, in the tests under consideration, the fictitious opening factor,  $(A\sqrt{H}/A_t)_{\text{fict}}$ , and hence the maximum rate of combustion, are determined approximately in conformity with the above from Eq. (5.3).

For the 16 tests comprised in this series, Table 7 gives, first, the opening factor calculated by means of the above relations, and second, the opening factor,  $(A \cdot \sqrt{H}/A_t)_{\rm exp}$ , that has proved to give theoretical results which are in agreement with the experimental values. The test results, the theoretical results, and the time graphs of the rate of combustion are shown in Appendix 3. As regards the time graphs of the combustion gas temperature, it is to be noted that the six full-line curves represent the temperatures at different

Table 7. Theoretical and experimental values of the opening factor.

Test series C [18].

Test No.	Fire load, M, kg of wood	Hydraulic radius, $r$ , cm	Rate of air supply, Q, m <sup>3</sup> ·s <sup>-1</sup>	Opening factor, theoretical value, $(A \overline{H}/A_t)_{\text{theor}}$ calculated from the formula $A \overline{H}/A_t=0.0334 \cdot Q$	Opening factor, experimental value, $(A\sqrt[3]{H}/A_t)_{\rm exp}$	$\frac{(A\sqrt{H}/A_t)_{\text{exp}}}{(A\sqrt{H}/A_t)_{\text{theo}}}$
4	675	1,0	0.7	0.023	0.035	1,52
8	675	1.7	0.7	0.023	0.015	0.65
10	270	1.0	0.7	0.023	0.015	0.65
11	405	1.0	0.7	0.023	0.020	0.87
12	405	2.4	0.7	0.023	0.012	0.52
13	135	1.0	0.7	0.023	0.005	0.22
16	405	0.4	0.7	0.023	0.060	2.60
1	270	_	1.0	0.033	0.023	0.70
3	675	1.0	1.0	0.033	0.043	1.30
7	675	1.7	1.0	0.033	0.037	1.12
9	675	0.6	1.0	0.033	0.051	1.54
14	945	1.6	1.0	0.033	0.037	1.12
5	675	1.0	1.5	0.050	0.055	1.10
2	675	1.0	2.0	0.067	0.060	0.90
15	1350	1.4	2.0	0.067	0.060	0.90
			about			
6	675	1.0	0.25	0.009	0.010	1,11

points in the enclosed space. The fine dash-line curve summarises the values recorded in the radiation measurements, and the heavy dash-line curve is the calculated curve. The variation in the rate of combustion with the time,  $I_c$ , is represented in terms of  $330 \cdot (A\sqrt{H})_{exp} \cdot 2575 \text{ kcal} \cdot h^{-1}$  put equal to unity. The rate of flow of the incoming air, Q, conveyed by the fan system was constant during the whole process of fire development. The radiation measurements have provided certain indications for choosing the instant at which the rate of combustion had decreased to zero. The difference between the temperature at the level of the floor surface and the average temperature in the other parts of the enclosed space has been taken into account. In the test series under review, this difference in temperature has probably been increased owing to the fact that air was supplied to the enclosed space by means of fans at the floor surface level. The temperature difference was taken into consideration by assuming that the coefficient of heat transfer at the floor surface was equal to 80 per cent of the corresponding coefficient for the other surfaces. In all cases when the rate of burning was controlled by the fuel bed and not by the ventilation it was taken into account that heat energy was withdrawn from the enclosed space by that part of the incoming air which did not take part in the combustion.

As may be seen from Table 7, the positive as well as negative differences between the value of the opening factor,  $(A\sqrt{H}/A_t)_{\rm exp}$ , determined from the test results, i.e. the actual maximum values of the rate of combustion, and the corresponding values obtained on the assumption that the rate of combustion is limited by the rate of air supply, were found to be great in some tests. However, the calculated curves were as a rule closely in agreement with the observed values. The agreement between the maximum temperature and the duration of the flame phase indicates that those values of the quantity of energy released per unit time which were used for the theoretical calculations were on the whole correct.

Furthermore, Table 7 shows two other factors among those which, in addition to the air flow factor, determine the rate of combustion. These factors are the hydraulic radius and the amount of fuel (the fire load). The effect of the first-mentioned factor can be demonstrated, for instance, by a comparison between the tests. Nos 16 and 12. The values of the air flow factor, as well as those of the fire load, in these two tests were equal, whereas the respective values of the hydraulic radius were 0.4 and 2.4 cm. In consequence of the difference in the hydraulic radius, the actual maximum rate of combustion in the test No. 16 was about 5 times as high as in the test No. 12. This may roughly be explained by the simplified study of the mechanism of combustion in what follows.

In a wood fuel consisting of comparatively large pieces of wood, pyrolysis takes place in several forms at the same time. In a certain definite inner zone

of the wood, where the temperature is relatively low, say, below about  $250^{\circ}$ C, the reactions are endothermic, whereas in the outer zones, where the temperature is higher, the reactions are exothermic, and in certain cases, e.g. in the secondary pyrolysis of tar products, markedly exothermic. The combustion of the products of pyrolysis generates heat, which increases the temperature of the fuel by conduction, and hence renders possible an exothermic decomposition in the inner zone. In the test No. 16, the wood fuel consisted of concrete form timber, for which the ratio of the volume to the exposed surface was 0.4 cm. Accordingly, if the width, the length, and the thickness of a piece of wood are denoted by b, l, and t, respectively, then we have

$$r = 0.4 \text{ cm} = \frac{b \cdot l \cdot t}{2l(b+t)}$$

If we put b=t (square cross section), then we obtain a thickness of 1.6 cm, and if we set  $b\gg t$ , then we get a thickness of 0.8 cm. In view of the small thickness, in combination with the mechanism of heat return to the fuel described in the above, it is probable that a few minutes after the fuel has taken fire the whole quantity of fuel is in a state of active exothermic pyrolysis. Tests [23] have shown that the progression of the charred layer on a wooden beam exposed to fire is about 0.6 mm · min<sup>-1</sup>. This value is applicable to a firwood beam exposed to fire in conformity with a standard temperature-time curve. The variation in the rate of carbonization with the intensity of the process of fire development is a problem which appears to be wholly unexplored at the present time. But if the above-mentioned value, 0.6 mm · min<sup>-1</sup>, is assumed to be correct, then this implies that the whole amount of fuel would be charred in the course of 5 to 10 min. Since from one half to two thirds of the total quantity of energy is liberated during the flame phase, this can explain the intense release of energy immediately after the fuel has taken fire.

The test No. 12 shows that when the fire-exposed surface area diminishes below a certain definite limit, the quantity of energy which can be developed per unit time is determined by the rate of progression of the charred layer, and not by the air flow factor. In this test, the hydraulic radius was 2.4 cm. If the above-mentioned value,  $0.6 \text{ mm} \cdot \text{min}^{-1}$ , which is probably too high in view of the low temperature during the process of fire development in this case, is used as a measure of the progression of the charred layer, then this corresponds to a maximum rate of combustion of  $405 \cdot 0.6 \cdot 10^{-2}/0.24 = 10.1 \text{ kg} \cdot \text{min}^{-1}$ . The value of the maximum rate of combustion computed from the formula  $R = 5.5 A \sqrt{H}$  is 9.7 kg·min<sup>-1</sup>, and this implies that the quantity of energy released per unit time during a fire is determined by the surface area exposed to fire, and not by the rate of air supply, at least during certain phases of the fire.

Furthermore, a closer study of Table 7 also demonstrates the marked effect produced by the amount of fuel on the rate of combustion. This can likewise be explained by means of the mechanism of return of the heat evolved by combustion which has been outlined in the above. In the tests Nos. 13, 10, 11, and 4, the values of the hydraulic radius, r, were equal, just as those of the rate of air supply Q, whereas the fire load, M, was varied. For the values of the fire load M=135, 270, 405, and 675 kg, which approximately corresponded to 7, 15, 20, and 35 Mcal·m<sup>-2</sup> of bounding surface area of the enclosed space, the respective values of the maximum rate of combustion calculated from the test results were found to be equal to 0.22, 0.65, 0.87, and 1.5 times the rate of combustion which was computed on the basis of the air flow factor. In [18], it is stated that the flame phase was slightly developed in the tests Nos. 10 and 11, whereas the fuel did not take fire at all in the test No. 13.

Moreover, it is seen from Table 7 that the variation in the rate of combustion with the hydraulic radius seems to decrease as the opening factor or the fire load increases.

Finally, it may be useful to touch on the question to what extent these analytic fire tests, which were carried out under conditions that were idealized so far as possible, and which exhibited characteristics of combustion that in several tests markedly differed from those predicted by the theory, can be utilized as a basis for predicting the behaviour of more conventional fires. An ordinary fire in an enclosed space is governed to a varying degree by two feed-back mechanisms. In the first place, the lower density of the combustion gases forces unconsumed air by natural convection towards the flames, and hence increases the rate of combustion, as well as the evolution of combustion gases. In the second place, part of the heat generated by combustion returns to the fuel, and increases the rate of energy release. So long as these mechanisms are negative, the combustion remains stable. Cf. e.g. the test No. 13. An essential difference between Ödeen's tests described in [18] and an ordinary fire is that the return mechanism which governs the rate of air supply was eliminated in the tests. The effect of this circumstance is difficult to determine, but it may be mentioned for comparison that all the theoretical and experimental results in the test series (A, B, and D) where the exchange of air was self-regulated were found to be closely in agreement if the rate of combustion was assumed to be determined by the rate of air supply. However, a comparison of the time graphs of the rate of combustion for the test series A to D shows that, in the cases where differences were present, the results obtained from the test series C deviated from those of the other test series only in respect of the maximum quantity of energy released per unit time. If the time graphs of the rate of combustion are represented in terms of 330  $\cdot$   $(A)\overline{H})_{\rm exp}$   $\cdot$  2575 kcal · h-1 put equal to unity, then the curves for this test series are found

to be closely in agreement with those for the test series A, B, and D when the fire load and the opening factor are given. All the same, since a dispersion in the values of the quantity of energy liberated per unit time has been observed in the test series C even when the tests were identical in respect of the rate of air supply and the wall material, it should be noted that this dispersion indicates the need for determining the factors which, in addition to those mentioned in the above, govern the rate of release of energy.

## Test series D. Tests made at the National Swedish Institute for Materials Testing. Calculation of time graphs of rate of combustion

Under the direction of the Fire Engineering Laboratory of the National Swedish Institute for Materials Testing, a test house for model-scale and full-scale fire studies has been erected at the Studsvik Test Station of the Atomic Energy Co., Ltd. The primary object of the investigations carried out in this house was to study the spread of fire and smoke along the exterior walls and along the ventilation ducts in the case of fire in an individual enclosed space in a multi-storeyed building. Extensive measurements of the reduction in the weight of fuel, the temperature, the intensity of radiation, the gas flow, and the composition of combustion gases were made in these tests, and the test results are therefore well suited for theoretical comparisons.

Fig. 12 shows the test house, which was three storeys high, and which consisted of a load-bearing steel frame clad with lightweight concrete elements. The results of the tests carried out up to now have not yet been published, but the test programme and the test equipment are described in [20] and [21]. The Authors of the present publication were afforded an opportunity to acquaint themselves with the results of the first four full-scale tests. The data for these tests are reproduced in Table 8.

The fires were initiated in the lowermost storey, which was connected with the outside air by means of a vertical ventilation duct by-passing the storeys

Table 8. Test series D.

Test No.	Fire load, kg of wood	Window dimensions, width×height m×m	Bounding surface area, m <sup>2</sup>	Air flow factor, $A \cdot \sqrt{H}$ , $m^{5/2}$	Air flow factor, fictitious value, $(A   \overline{H})_{\text{fict}}$	Moisture content of fuel, per cent	Moisture content of lightweight concrete, per cent
1	350	1.3×1.3	75	1.93	2.10	7.7	4.0
2	200	$1 \times 1$	75	1.00	1.10	8.8	3.1
3	115	$0.8 \times 0.8$	75	0.57	0.69	7.4	2.1
4	1150	$2\times2$	75	5.63	5.65	9.7	2.3

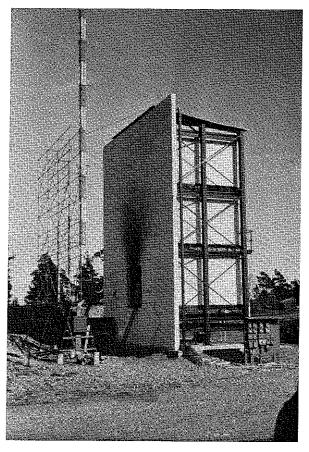


Fig. 12. Photograph of the test house used in the test series **D** to study the spread of fire and smoke along the external walls and through a vertical ventilation duct in the case of fire in a certain definite storey in a multi-storey building.

situated above. Calculations showed that the effect of the ventilation duct on the air flow factor was not to be disregarded, and a fictitious air flow factor,  $(A\sqrt{H})_{\text{fict}}$ , was therefore computed in accordance with the principles stated in Chapter 3. The magnitude of the corresponding correction can be seen from Table 8. The time graphs of the temperature of the combustion gases which were obtained from these tests, and which are expressed in terms of the mean value of the temperatures observed during the tests at 21 points in the enclosed space, are reproduced in Appendix 4. In order to give an idea of the dispersion about this mean value, the diagram relating to the test No. 2 shows the temperature-time curves for a point situated 45 cm above the floor and a point located 45 cm below the ceiling. Furthermore, the four diagrams in Appendix 4 also comprise the calculated curves which represent the varia-

tion in the temperature of the combustion gases with the time, as well as the time graphs of the rate of combustion in a dimensionless form which were used in the calculations.

It proved possible to bring the calculated and observed curves into close agreement by choosing the time graphs of the rate of combustion which were similar in shape in all the tests, and were based on a maximum rate of combustion  $I_c = 330 \cdot A\sqrt{H} \cdot 2575$  kcal·h<sup>-1</sup>. The value of the air flow factor used in the calculations was the value  $(A\sqrt{H})_{fict}$ , which was corrected so as to take account of the effect of the ventilation duct.

### **Summary**

All the comparative calculations dealt with in the present chapter were based on the assumption that the energy conditions during the process of combustion can be characterized by an ignition phase in which the quantity of energy released per unit time increases from a zero value in accordance with a polygonal function of the time to a value that is given by the air flow factor. This phase is followed by a flame phase, during which the rate of combustion was supposed to be constant. After that the rate of combustion decreases to zero as a polygonal function of the time in the course of the cooling phase, during which the slopes of the individual sides of the polygon vary in a marked manner with the fire load. The higher the fire load, the slower the decrease in the rate of combustion. Of course, these assumptions give a simplified picture of the variation in the liberation of energy per unit time during the process of fire development. Thus, for most types of fire loads, it is to be expected that the plane part of the curve which represents the rate of combustion during the flame phase is rather to be regarded as the mean value of the quantity of energy released per unit time. This is illustrated in Fig. 13, which represents the variations in the observed rate of combustion (expressed in terms of the rate of reduction in the weight of fuel) and in the oxygen content of the combustion gases during the test No. 1 in the test series A. On account of technical difficulties in measurements, the values of the rate of reduction in the weight of fuel were somewhat uncertain immediately after the fuel had taken fire, as the rate of combustion was then liable to very wide variations. However, these values were confirmed by the fact that the oxygen content of the combustion gases in the enclosed space exhibited corresponding variations. Even when the curve which represents the variation in the rate of reduction in the weight of fuel with the time is known, our present knowledge of the relation between the rate of reduction in the weight of fuel and the rate of release of energy during the different phases of the process of fire development does not make it possible to determine the quantity of energy liberated per unit time during combustion. For the test No. 3 in the test series D, Fig. 14 shows three theoretical temperature-time curves calculated on the basis

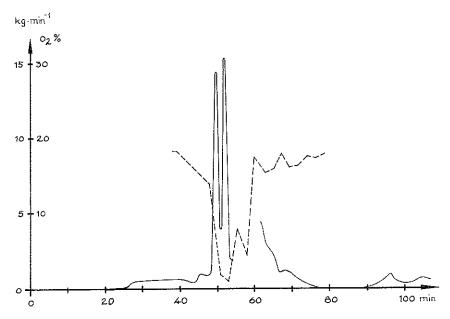


Fig. 13. Variation in the rate of combustion, in kg·min<sup>-1</sup>, with the time, t, determined by measuring the reduction in the weight of fuel in one of the tests comprised in the test series A, see Chapter 5. Furthermore, this figure also shows the time graph of the oxygen content of the combustion gases in the same test (dash-line curve). During the interval from the 55th to the 60th minute, the weighing of the fuel was disturbed by the fact that parts of the ceiling of the enclosed space fell down.

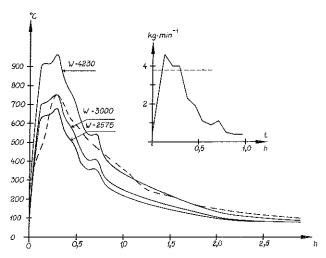


Fig. 14. Time graphs of the temperature of the combustion gases calculated for three different heat values of the fuel, viz., W=2575, 3000, and 4230 kcal·kg<sup>-1</sup> of wood, on the basis of an observed time graph of weight loss of fuel (inset). An experimental temperature-time curve is represented by the dash-line curve.

of a measured rate of weight loss. It was assumed that the heat value W corresponding to 1 kg of weight loss was constant during the whole process of fire development and equal to 2575,3000 and 4230 kcal, respectively. The dash line curve represents the values observed in the test. The observed rate of reduction in the weight of fuel is represented in a separate diagram in fig. 14. As seen from this diagram, the maximum rate of reduction in the weight of fuel is closely in agreement with the theoretical value  $R = 5.5 \times A\sqrt{H} = 3.80 \text{ kg} \cdot \text{min}^{-1}$  (horisontal dash line). It is also seen from fig. 14 that the assumption of a constant value of W obviously is incorrect. A comparison between the curves indicates that the quantity of energy released per unit weight of fuel, at least during the first part of the cooling phase is greater that that during the flame phase.

Accordingly, the relative values of the rate of combustion obtained from the above comparative calculations cannot be directly expressed in terms of the rate of reduction in the weight of fuel, in kg·min<sup>-1</sup>. If the quantity of energy liberated per unit time is expressed directly in kcal·min<sup>-1</sup>, then this obviates the difficulty of determining the heat of combustion of the wood fuel during the various phases of the process of fire development.

A prerequisite in most of the calculations has been that the fire process is controlled by ventilation. This condition is far from being generally realized when it comes to actual fires. As a rule, a combination of small ventilation area and high fire load gives a process of fire development where the rate of burning is proportional to the air flow factor. If the fire load is low and the ventilation area large, the combustion will proceed as if in the open. This means that the rate of burning depends on the fire load density (fire exposed surface) and that an increase in ventilation will not result in a corresponding increase in rate of burning.

It is, however, impossible to say in advance if the process of fire development will be ventilation controlled or not even if the air flow factor and the fire load are known. The orientation and the distribution of the fuel in the enclosed space and the thickness or the porosity of the fuel will be a decisive factor in each particular case. An assumption that the rate of burning is determined by the air flow factor ought to give time-temperature curves which are on the safe side in practically every case. If such an assumption is made and the combustion in spite of this happens as in the open, i.e. is fuel bed controlled, the result will be a fire process of lower maximum temperature and, at least in some cases, of longer duration. The longer duration will not increase the severity of the fire to a corresponding degree. This is due to the fact that part of energy released by the combustion will be withdrawn from the enclosed space by the surplus air. In this way the temperature of the combustion gases will be lower compared to the case when the duration is the same but the process controlled by the ventilation.

To sum up, a comparative theoretical analysis of the results obtained from some thirty full-scale fire tests of the wood fuel type has been carried out in Chapter 5. The calculations made for this purpose covered relatively wide variations in fire load, opening factor, and hydraulic radius, as well as in the thermal properties of the structures bounding the enclosed space. As a result of these calculations, the time graph of the quantity of energy released per unit time may be assumed to be known within this range of variation.

# 6. Determination of general time graphs of quantity of energy released per unit time during different phases of process of fire development

In order to afford a basis for the calculation of the curve which represents the variation in the temperature of the combustion gases with the time during the process of fire development under varying conditions, it is necessary to systematize the time graphs of the rate of combustion which have been obtained in Chapter 5. A detailed investigation has been made of these graphs in order to find out how they vary with the fire load and with the opening factor. This investigation indicated the possibility of the simplification outlined in what follows.

If the ratio of the fire load, which is given from the outset, to the air flow factor,  $A\sqrt{H}$ , is constant, that is to say, if the duration of the fire is constant, then the time graph of the rate of combustion, expressed in a relative form in terms of the maximum rate of combustion,  $330 \cdot A\sqrt{H} \cdot 2575$  kcal·h<sup>-1</sup>, put equal to unity, is independent of the opening factor. This implies, for instance, that an enclosed space where the opening factor is  $A\sqrt{H}/A_t = 0.01$  m<sup>1/2</sup> and the fire load is q = 5 Mcal·m<sup>-2</sup> can be characterized by the same graph of the rate of combustion as an enclosed space where the opening factor is 0.04 m<sup>1/2</sup> and the fire load is 20 Mcal·m<sup>-2</sup>. Accordingly, the results of the calculations in Chapter 5, which show how the quantity of energy released per unit time varies with the time, can be represented by a graph which comprises a separate curve for each value of the ratio  $qA_t/A\sqrt{H}$ . In this connection, it is convenient to introduce the duration of the fire, T, defined as the duration of the flame phase, cf. Eq. (1.2),

$$T = qA_t/(1500A \cdot \sqrt{H})$$
 h

as the variable at which the graph shall be entered. In this formula for calculating the duration of the flame phase, the product of the constant 330 in the expression  $R=330 \cdot A\sqrt{H}$  kg·h<sup>-1</sup> and the heat value of the wood fuel, i.e. 4.5 Mcal·kg<sup>-1</sup>, has been put equal to 1500. For the values of the duration of the fire defined in this way, T=0.1, 0.2, 0.3, 0.5, 0.75, 1.0, 1.5 and 2.0 h, which correspond to the respective fire loads (150, 300, 450, 750, 1125, 1500, 2250, and 3000)  $\cdot A\sqrt{H}/A_t$  Mcal·m<sup>-2</sup> of the total surface area bounding

the enclosed space, Fig. 15 shows the variation in the rate of combustion with the time.

In order to make it easier to check the agreement between the total quantity of energy liberated during the process of fire development and the fire load, which is given from the beginning, and since a more accurate representation would be illusory, considering the character of the available data, the curve form has been assumed to be polygonal, just as in the comparative calculations in Chapter 5. In Fig. 15,  $330A\sqrt{H} \cdot 2575$  kcal·h<sup>-1</sup> has been put equal to unity. The respective areas between the above-mentioned curves and the axis of time shall therefore be  $\frac{1}{330 \times 2575} (150, 300, 450, 750, 1125, 1500, 2250, and 3000)$  area units. In the relevant Swedish regulations, the quantity T, determined by the relation

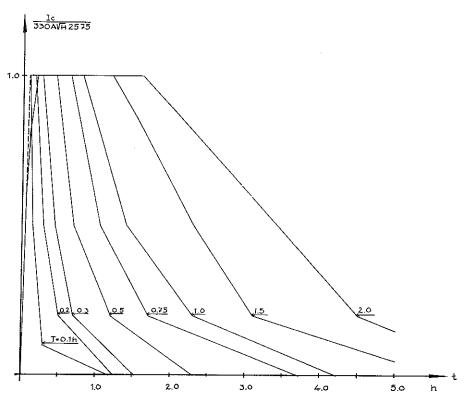


Fig. 15. Time graphs of the energy released per unit time in the process of combustion,  $I_c$ , expressed in a relative form by putting  $330 \cdot 2575 \cdot A \sqrt{H}$  equal to unity. The eight curves shown in this figure correspond to different values of the duration of the fire defined as the duration of the flame phase by the expression  $T = q \cdot A_t / (1500 \cdot A \cdot \sqrt[3]{H})$ . The dash-line portion of the curve for the ignition phase belongs to the curves relating to the lowest four values of the duration of the fire.

### $T = qA_t/(25A\sqrt{H})$ min

designates the instant which marks the end of the flame phase and the beginning of the linear cooling phase. On the basis of the experiences derived from the comparative theoretical analyses, the instant at which the rate of combustion begins to decrease has been chosen so as to be slightly anterior to the instant defined by T. For the values of the duration of the fire T=6, 12, 18, and 30 min, the time graphs of the rate of combustion during the ignition phase have been given a slightly different shape, which implies that the fuel takes fire within a shorter period of time.

# 7. Calculation of time graphs of temperature of combustion gases for characteristic types of enclosed spaces varying in opening factor and in fire load

The time graphs of the rate of combustion represented in a relative form in Fig. 15 for fires of the wood fuel type in enclosed spaces are utilized in the present chapter as a basis for the calculation of complete time graphs of the temperature of the combustion gases. This is done for varying values of the opening factor and the fire load in enclosed spaces of the seven types dealt with in what follows, which differ in respect of the bounding structures.

### Type A enclosed space

Bounding structures.

All the surfaces which bound the enclosed space are supposed to consist of a material, 20 cm in thickness, whose thermal properties are characterized by the average values given below, which apply to structural materials of such types as concrete, brick, and lightweight concrete.

Thermal conductivity,  $\lambda = 0.7 \text{ kcal} \cdot \text{m}^{-1} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$ .

Product of the specific heat and the weight per unit volume,

$$c \cdot \gamma = 400 \text{ kcal} \cdot \text{m}^{-3} \cdot {}^{\circ}\text{C}^{-1}$$
.

The same data on the properties of materials had also been used for the calculation of those temperature-time curves for the flame phase of the process of fire development which have been published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures".

### Type B enclosed space

Bounding structures.

Concrete, 20 cm in thickness.

Thermal conductivity,  $\lambda = 1.4 \cdot e^{-0.001} \cdot \theta \text{ kcal} \cdot \text{m}^{-1} \cdot \text{h}^{-1} \cdot {}^{\circ}\text{C}^{-1}$  [1]. Enthalpy,  $I_s$ , see Fig. 17.

#### Type C enclosed space

Bounding structures.

Lightweight concrete, 20 cm in thickness.

Weight per unit volume,  $\gamma = 500 \text{ kg} \cdot \text{m}^{-3}$ .

Thermal conductivity,  $\lambda$ , see Fig. 16.

Enthalpy, I, see Fig. 17.

The specific heat and the weight per unit volume of the lightweight concrete are assumed to be independent of the temperature. Consequently, the enthalpy-temperature curve is rectilinear. The variation in the thermal conductivity,  $\lambda$ , with the temperature is based on a determination which has been made in connection with the test series D described in Chapter 5.

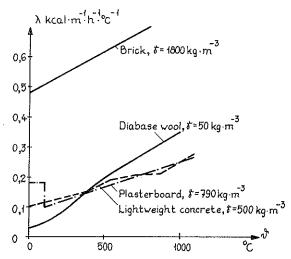


Fig. 16. Relations between the thermal conductivity,  $\lambda$ , and the temperature,  $\vartheta$ , used in the calculations for brick, diabase wool, plasterboard, and lightweight concrete.

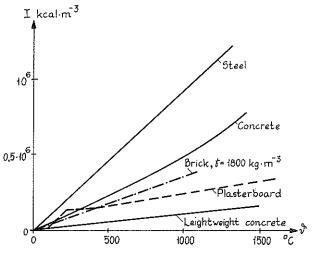


Fig. 17. Relations between the enthalpy, I, and the temperature,  $\vartheta$ , used in the calculations for steel, concrete, brick, plasterboard, and lightweight concrete.

### Type D enclosed space

Bounding structures.

Concrete, 50 per cent of the total bounding surface area.

Lightweight concrete, 50 per cent of the total bounding surface area.

Thicknesses, weight per unit volume, and thermal properties as in the Type B and Type C enclosed spaces, respectively.

### Type E enclosed space

Bounding structures.

Lightweight concrete, 50 per cent of the total bounding surface area. Thickness, weight per unit volume, and thermal properties as in the Type C enclosed space.

Concrete, 33 per cent of the total bounding surface area.

Thickness and thermal properties as in the Type B enclosed space. Other structural components, 17 per cent of the total bounding surface area, enumerated in the order from the interior to the exterior:

Plasterboard panel, 13 mm in thickness.

Weight per unit volume,  $\gamma = 790 \text{ kg} \cdot \text{m}^{-3}$ .

Diabase wool, 10 cm in thickness.

Weight per unit volume,  $y = 50 \text{ kg} \cdot \text{m}^{-3}$ .

Brickwork, 20 cm in thickness.

Weight per unit volume,  $\gamma = 1800 \text{ kg} \cdot \text{m}^{-3}$ .

Thermal properties of plasterboard, diabase wool, and brick, see Figs. 16, 17, and 18.

The enthalpy-temperature curve chosen for brick is based on a value of the specific heat which is supposed to be independent of the temperature. In reality, the specific heat of brick slightly varies with the temperature [1]. However, in the present case, the values of the temperature rise in the brickwork, which is most remote from the surface exposed to fire, are so low that

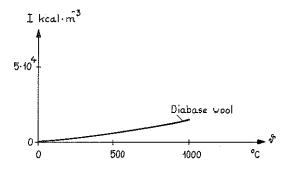


Fig. 18. Relation between the enthalpy, I, and the temperature,  $\vartheta$ , used in the calculations for diabase wool.

the effect of this variation may be disregarded. The variation in the enthalpy of diabase wool with the temperature is based on that relation between the specific heat and the temperature which was published in [1]. The thermal properties of plasterboard are based on the curves which were published in [19]. It is assumed that the plasterboard panels will not fall down or disintegrate.

The values of the thermal conductivity,  $\lambda$ , of brick and diabase wool were taken from [1].

### Type F enclosed space

Bounding structures.

Sheet steel, 2 mm in thickness, 80 per cent of the total bounding surface area.

Concrete, 20 cm in thickness, 20 per cent of the total bounding surface

Thermal properties as in the Type B enclosed space.

Curves representing the variations in the enthalpy and in the thermal conductivity of sheet steel with the temperature, see Figs. 17 and 19, respectively.

This type of enclosed space corresponds to a storage space, or the like with a sheet steel roof, sheet steel walls, and a concrete floor.

### Type G enclosed space

### Bounding structures.

Concrete, 20 per cent of the total bounding surface area.

Thickness and thermal properties as in the Type B enclosed space.

Other structural components, 80 per cent of the total bounding surface area, enumerated in the order from the interior to the exterior:

Two plasterboard panels,  $2 \times 13$  mm in thickness.

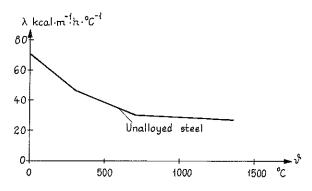


Fig. 19. Relation between the thermal conductivity,  $\lambda$ , and the temperature,  $\vartheta$ , used in the calculations for unalloyed steel.

Weight per unit volume,  $\gamma = 790 \text{ kg} \cdot \text{m}^{-3}$ .

Cavity, 10 cm in width.

Two plasterboard panels,  $2 \times 13$  mm in thickness.

Thermal properties of plasterboard, see Figs. 16 and 20.

This structure represents a type of partition which is becoming more and more common, and which consists of two plasterboard panels on each side, supported on steel stud framing. It is assumed that the steel studs have no thermal conductivity and no thermal absorptivity.

The test results published in [19] have shown that plasterboard panels which are not fibre-filled disintegrate when their temperature on the side that is not exposed to fire reaches about 550°C. However, this does not apply to the outermost, i.e. the fourth, plasterboard panel. This panel is in contact with the air, which has a temperature of 20°C. Therefore, this panel never reaches a surface temperature of 550°C. In fact, tests have demonstrated that a plasterboard panel in this position disintegrates when the temperature at its centre rise to about 750°C. These criteria have been used in calculating the time graph of the combustion gas temperature for enclosed spaces of the type in question. The calculations were discontinued when they had been carried out to the instant at which the fire was expected to burn through the wall, that is to say, after all four plasterboard panels had disintegrated.

As may be seen from Fig. 20, which was taken from [19], the variation in the enthalpy of plasterboard with the temperature is dependent on whether

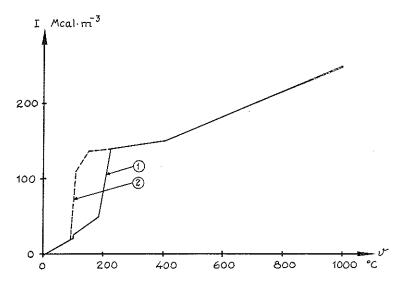


Fig. 20. Relation between the enthalpy, I, and the temperature,  $\vartheta$ , used in the calculations for plasterboard. Curve 1: High rate of temperature rise. Curve 2: Low rate of temperature rise.

the rate of temperature rise is high or low. This circumstance was taken into account in the calculations by assuming a high rate of temperature rise for that plasterboard panel which is nearest to the fire and a low rate of temperature rise for all the other plasterboard panels. Since the structural transformations of plasterboard require an additional quantity of heat that cannot be recovered, the enthalpy-temperature curve during the cooling period is not identical with that during the heating period.

In the calculations, it was assumed that the variation in the enthalpy with the temperature during the cooling period is represented by a straight line which corresponds to a constant value of the product  $c \cdot \gamma = 200 \text{ kcal} \times \text{m}^{-3} \cdot {}^{\circ}\text{C}^{-1}$ .

For the surfaces exposed to fire in the enclosed spaces of all the types dealt with in the present chapter, the coefficient of heat transfer,  $\alpha_i$ , was calculated by means of Eq. (2.7a), where it was assumed that  $\varepsilon_{fi} = 0.7$  and  $\varepsilon_i = 0.8$  for the Type A to the Type E enclosed spaces, and for the Type G enclosed space. These values give  $\varepsilon_{res} \sim 0.60$ . For the Type F enclosed space, the calculations were based on three values of  $\varepsilon_i$ , viz., 0.1, 0.4, and 0.8. Hence, for  $\varepsilon_{fi} = 0.7$ , the respective values of  $\varepsilon_{res}$  were found to be 0.1, 0.35, and 0.6. For the exterior surfaces of the structures bounding the enclosed space, the coefficient of heat transfer,  $\alpha_u$ , for the Type A to Type E enclosed spaces, and for the Type G enclosed space, was supposed to vary with the surface temperature,  $\vartheta_u$ , in accordance with Eq. (2.8), while its value for the Type F enclosed space was supposed to vary in conformity with the relation

$$\alpha_{u} = 7.5 + \frac{4.96 \cdot \varepsilon_{res}}{\theta_{u} - \theta_{o}} \left[ \left( \frac{\theta_{u} + 273}{100} \right)^{4} - \left( \frac{\theta_{o} + 273}{100} \right)^{4} \right] \text{kcal} \cdot \text{m}^{-2} \cdot \text{h}^{-1}. \quad ^{\circ}\text{C}^{-1}$$
(7.1)

where  $\varepsilon_{res}$  was chosen in the same way as in the above-mentioned calculation of  $\alpha_i$ .

The variation in the specific heat,  $c_p$ , of the combustion gases with the temperature is shown in Fig. 21 [3]. For the rest, the calculations were based on the principles which have been stated in Chapters 2 and 4.

The input values used in the calculations were only the opening factor  $A \cdot \sqrt{H}/A_t$ , in  $m^{1/2}$ , and the fire load, q, in Mcal· $m^{-2}$  of bounding surface area. If the radiation term,  $I_R$ , is disregarded, then Eq. (2.1) becomes independent of dimensions. In other words, if the values of  $A \cdot \sqrt{H}/A_t$  and q are given, then the result will be independent of the terms A, H, and  $A_t$  entering into the opening factor. However,  $I_R$  is proportional to the total opening area, A, and, in order that  $I_R$  may be taken into account, it is necessary to specify the above-mentioned terms. This has been done in the calculations dealing with test results in Chapter 5. For the determination of the total

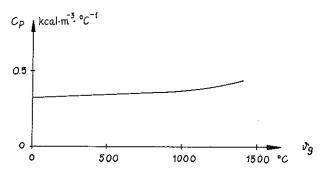


Fig. 21. Relation between the specific heat of the combustion gases,  $c_p$ , and their temperature,  $\vartheta_q$ , used in the calculations.

opening area, A, in the calculations described in the present chapter, it was supposed that the dimensions of the enclosed space were the same as those which had been assumed as a basis for the curves published in the Swedish Building Regulations 1967, viz., a total bounding surface area  $A_t = 10,000 \text{ m}^2$  and a square opening. Consequently, the ratio of the total opening area to the total bounding surface area,  $A/A_t$ , was in all cases lower than those values which can be expected to be met with in ordinary buildings, and all the results will therefore be on the safe side. The value  $A_t = 10,000 \text{ m}^2$  was used only to determine the value of the ratio  $A/A_t$ , which was then substituted in the term  $I_R$  in the equation of heat balance. For a value of the opening factor  $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$ , Table 9 gives the respective values of the ratio  $A/A_t$  which correspond to  $A_t = 10,000 \text{ m}^2$  and  $A_t = 1 \text{ m}^2$  on the assumption that the opening is square, or that it has a height H = 1 m.

Table 9. Values of the ratio  $A/A_t$ .

	Square enemine II 1		
_	Square opening H=1 m		
$A_t = 1 \text{ m}^2$	0.075	0.04	
$A_t = 10,000 \text{ m}^2$	0.012	0.04	

In order to illustrate the consequence of this variation in the ratio  $A/A_t$ , Fig. 22 shows the temperature-time curve for an enclosed space characterized by an opening factor  $A \cdot \sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$ , a fire load  $q = 30 \text{ Mcal} \cdot \text{m}^{-2}$ , as well as by the two extreme values of the ratio  $A/A_t$ , i.e. 0.012 and 0.075. As is seen from this graph, the effect of the difference in the value of the ratio  $A/A_t$  is practically negligible.

Temperature-time curves have been calculated for each one of the seven enclosed spaces, Types A to G, for varying values of the duration of the fire and the opening factor. As regards the relation between the duration of the fire and the fire load, reference is made to Chapter 6. For the Type A to the

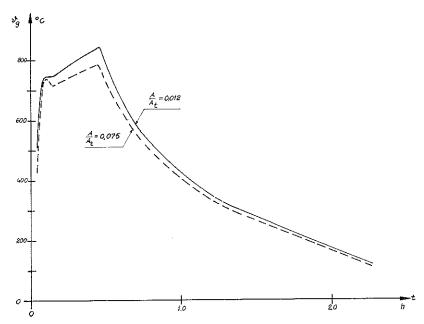


Fig. 22. Calculated temperature-time curves for the Type A enclosed space. Opening factor  $A\sqrt[4]{H}/A_t = 0.04 \text{ m}^{1/2}$ . Fire load  $q = 30 \text{ Mcal} \cdot \text{m}^{-2}$  of bounding surface area. Ratio  $A/A_t$  set equal to 0.012 and 0.075, respectively.

Type E enclosed space, as well as for the Type G enclosed space, the temperature-time curves were computed on the basis of 6 different values of the opening factor, viz.,  $A \cdot \sqrt{H}/A_t = 0.01$ , 0.02, 0.04, 0.06, 0.08, and 0.12 m<sup>1/2</sup>. For each of these values the curves were computed for 8 different values of the duration of the flame phase of the fire for which the time graphs of the rate of combustion have been constructed in Chapter 6. This means that 48 temperature-time curves have been obtained for each type of enclosed space. In the case of the Type F enclosed space, the calculations were carried out for 3 different values of the resultant emissivity. Furthermore, for each one of these values, use was made of 5 different values of the opening factor, viz., 0.01, 0.02, 0.04, 0.08, and  $0.12 \text{ m}^{1/2}$ . For each one of these combinations of values, the curves were calculated on the basis of 5 different values of the fire load corresponding to 5 different values of the duration of the flame phase, viz., 0.1, 0.3, 0.5, 1.0, and 2.0 h, if computed by means of Eq. (2.1). All these curves are shown in Appendix 5. The curves are denoted by the symbols Al to G6, where the letter A refers to the Type A enclosed space, etc. All these curves are represented in an approximate form after smoothing out the irregularities which were caused by the polygonal shape of the time graphs of the rate of combustion. Furthermore, in order to render the graphs in Appendix 5 more readily legible, that part of each one of the curves which represents the ignition phase was based on that ascending branch of the time graph of the rate of combustion which corresponds to the lower four values of the fire load. The exact results of the calculations are reproduced in tabular form in Appendix 6. In the case of the Type G enclosed space, where the calculation of the temperature-time curves was carried out with reference to the instants when the plasterboard panels fell down, each one of these instants is marked with a circle on the corresponding curve. When a plasterboard panel falls down, this corresponds in the calculations to an instantaneous temperature drop in the enclosed space, as may be seen from the relevant curves.

In practical design, it should be possible to proceed in three steps, viz., first, to choose that type of enclosed space which is most closely similar in respect of the thermal properties of the bounding structures to the case under consideration; second, to determine the opening factor and the fire load; and third, to interpolate linearly between the values given in the tables in Appendix 6. If, instead of using this procedure, the designer chooses a curve which is determined without interpolation so as to be on the safe side, that is to say, if he chooses the next higher values of the opening factor and the fire load, then this will probably not involve errors which are too great. In order to afford a basis for the choice of the type of enclosed space, the temperature-time curves which correspond to an opening factor  $A \cdot \sqrt{H}/A_t = 0.04$ m1/2 and to a fire load of 60 Mcal·m-2 of bounding surface area are represented in Fig. 23 for the Type A to the Type F enclosed spaces. For comparison, this graph also reproduces the standard ISO temperature-time curve and the curve for an opening factor of 0.04 m<sup>1/2</sup> published in the Swedish Building Regulations 1967.

A comparison between the temperature-time curves which correspond to the different types of enclosed spaces in Fig. 23 shows that the maximum difference in the maximum temperature amounts to about 400°C. The Type C enclosed space, which is provided with lightweight concrete bounding structures, exhibits markedly higher temperatures than the other types of enclosed spaces comprised in the present calculations. The lowest maximum temperature was obtained in the case of the Type F enclosed space ( $\varepsilon_{res} = 0.60$ ), which is equipped with bounding structures made of sheet steel, 2 mm in thickness. However, it is seen from Fig. 23 that the resultant emissivity for radiation between the flames and a sheet steel surface produces a substantial effect on the magnitude of the maximum temperature. If  $\varepsilon_{\rm res}$  is supposed to change from 0.6 to 0.1, then the corresponding difference in the maximum temperature is slightly over 200°C, other conditions being equal. Therefore, it is important that the heat transfer conditions, and particularly the resultant emissivity, should be accurately determined in the calculation of the temperature-time curve for an enclosed space of this type.

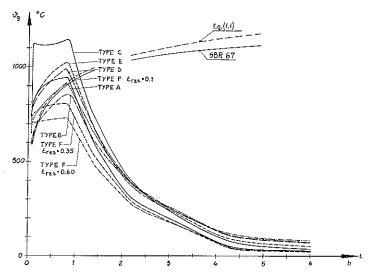


Fig. 23. Temperature-time curves for the Type A to the Type F enclosed spaces. Opening factor  $A\sqrt[4]{H}/A_t$ =0.04 m<sup>1/2</sup>. Fire load q=60 Mcal·m<sup>-2</sup> of bounding surface area. Furthermore, this graph also shows the respective temperature-time curves calculated by means of Eq. (1.1) and determined in conformity with the Swedish Building Regulations 1967 (SBR 67) for an opening factor  $A\sqrt[4]{H}/A_t$ =0.04 m<sup>1/2</sup>.

With the exception of the Type F, it is the Type B enclosed space, which is bounded by concrete walls, that exhibits the lowest maximum temperature. This is due to the relatively high thermal conductivity and the great heat capacity of the concrete. On the other hand, since the large quantity of heat that is stored in the concrete is partly transferred back to the enclosed space during the cooling period, comparatively high temperatures are obtained in the course of the cooling phase.

The temperature-time curves published in the Swedish Building Regulations 1967 and in the Draft Specification "Aluminium Structures" relate to "enclosed spaces bounded by wall, roof or ceiling, and floor structures which are made of brickwork, concrete, or lightweight concrete as a material that is predominant in thermal respects". As has previously been mentioned, these curves had been calculated on the basis of those characteristics of the bounding structures which were used to describe the Type A enclosed space. Moreover, for guidance, the comments on the Draft Specification "Aluminium Structures" also comprise temperature-time curves for an enclosed space which is bounded by walls made of mineral wool. In addition, it is shown how the temperature-time curve is influenced by a concrete wall, 20 cm in thickness, which is situated in the interior of the enclosed space. By examining Fig. 23, it will readily be understood that a further differentiation of the above-mentioned Swedish standard temperature-time curves according to the thermal characteristics of the bounding structures would be desirable.

In order that the temperature-time curves for the cooling phase which have been determined in the present publication might be compared with the corresponding Swedish standard curves, Fig. 24 shows the temperature-time curves for the Type A enclosed space calculated on the basis of an opening factor  $A \cdot \sqrt{H}/A_t = 0.04 \, \text{m}^1/^2$ , together with the Swedish standard temperature-time curves for the cooling phase determined on the assumption that the rate of temperature decrease is  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ . The ascending branches of the curves are identical because the curves for the Type A enclosed space are based on the same assumptions as the standard curves. The linear temperature-time curves for the cooling phase, which start from the ascending branch at the time T, calculated by means of Eq. (1.2), are represented by dash lines in Fig. 24. As is seen from this graph, the calculated curves result in a rate of cooling which is higher or lower than the standard rate of temperature decrease,  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ , according as the duration of the process of fire development is shorter or longer, respectively.

Thus, if the duration of the fire is comparatively short, then an application of the temperature-time curves which have been computed in the present publication gives considerably more favourable results, i.e. lower temperatures, than the standard rules which are at present in force in Sweden. For instance, the temperature-time curve for the cooling phase in fires of short duration is a decisive factor in determining the temperatures of unprotected steel structures, as has already been shown in an example which was adduced under

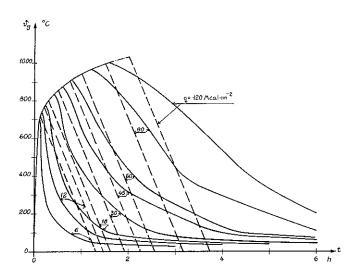


Fig. 24. Temperature-time curves for the Type A enclosed space. Opening factor  $A\sqrt{H}/A_t = 0.04 \text{ m}^{1/2}$ . (Full-line curves.) Curves for the cooling phase corresponding to a rate of decrease in temperature of  $10^{\circ}\text{C} \cdot \text{min}^{-1}$  stipulated in the relevant Swedish regulations. (Dash-line curves.)

the heading "Introduction". However, if the duration of the fire is relatively long, then the calculated temperature-time curves are less favourable, i.e. they give higher temperatures, than the curves stipulated in the Swedish regulations. All the same, since a given structure that is exposed to fire has already been subjected to the action of high temperatures for a long time before the beginning of the cooling phase, the difference between these curves in the latter case will have a comparatively slight effect on practical design.

## 8. Summary

The point of departure of the present investigation was the fact that the results obtained in recent years from research in the field of structural fire engineering had made it possible to carry out reliable calculations of the load bearing and separating capacity in the design of structural components exposed to fire. Such design calculations must be based on the knowledge of the temperature-time curve which covers the whole process of fire development. However, further progress towards realistic structural fire engineering design was impeded by the circumstance that it was not possible to make a theoretical determination of the temperature-time curve for the cooling phase of the process of fire development under known external conditions. Up to now, the research in this field has evolved methods for calculating the variation in the temperature with the time during the flame phase of the process of fire development on the basis of known external conditions, whereas the cooling phase has not been dealt with in this connection.

In consequence of this gap in our knowledge, the methods for determining the temperature-time curves for the flame phase and the cooling phase stipulated in the Swedish Building Regulations 1967 are widely different in degree of accuracy. The determination of the temperature-time curve for the flame phase shall be based on the fire load which characterizes the case under consideration, as well as on the shape and the dimensions of the openings in the enclosed space. For the cooling phase, on the other hand, it is stipulated only that the rate of temperature decrease shall be set equal to  $10^{\circ}\text{C} \cdot \text{min}^{-1}$ , irrespective of the actual conditions which characterize the case in question. This undifferentiated characterization of the cooling phase is particularly unfavourable to structures which possess a low thermal inertia, e.g. non-insulated or slightly insulated load-bearing steel structures. It was therefore considered to be urgently required to undertake an investigation in order to find out whether a theoretical determination of the temperature-time curve for the cooling phase would be possible.

The theoretical calculations in the present publication are founded on a basic equation of heat balance in an enclosed space which has been deduced by Kawagoe and Sekine, as well as by Ödeen. This equation states that the quantity of heat,  $I_c$ , which is released per unit time during the process of combustion is at any instant equal to the sum of the quantities of heat which

are withdrawn per unit time in different ways from the enclosed space. Heat is ordinarily abstracted from the enclosed space by heat transfer through the structures which bound the enclosed space (term  $I_W$  in the equation of heat balance), by radiation through the openings in the enclosed space (term  $I_R$ ), and by the replacement of combustion gases by cold air (term  $I_L$ ).

In order to extend the range of application of the equation of heat balance so that it might cover the whole process of fire development, it was necessary to solve two fundamental problems. In the first place, the quantity of heat released per unit time had to be determined as a function of the time for the entire process of fire development. In the second place, the expression for  $I_L$  which had been deduced previously, and which was applicable to the flame phase only, had to be extended and supplemented.

The study of the last-mentioned problem resulted in an expression for  $I_L$  which was based on the magnitude of the heat transfer by convection through the openings in the enclosed space. The rates of gas and air flow involved in this process were calculated in two steps, viz., first by determining the velocity distribution of gas flow in a vertical opening by which two masses of gas differing in density are separated from each other, and second, by satisfying the condition that the net exchange of gases between the enclosed space and its surroundings shall be equal to the difference between the quantity of gas produced and the quantity of air consumed in the process of combustion. After that, it was possible to determine  $I_L$  directly as the difference in heat content between the outgoing gases and the incoming air. It was found that  $I_L$  was approximately proportional to the temperature of the combustion gases and to the air flow factor  $A \cdot \sqrt{H}$ .

Since no physical basis is available which could enable the quantity of energy liberated per unit time during the process of fire development to be determined as a function of the time, a study of the literature was carried out with a view to an analysis of full-scale fire tests. For the tests where the external conditions were stated in a sufficiently precise manner, comparative calculations of temperature-time curves were made by means of a computer.

A tentatively chosen time graph of the quantity of energy liberated per unit time was used for this purpose. The time graph in question was varied until the agreement between the observed and calculated temperature-time curves became as close as possible. The only requirement to be fulfilled in this connection was that the total quantity of energy released during the whole process of fire development should be equal to that which was available in the fuel from the outset. When all those tests which were suited for this study had been examined, it was possible to systematize the results of the study in such a way that the time graph of the quantity of energy released per unit time during the process of fire development might generally be assumed to be known.

This procedure was primarily justified by the consideration that an error, if any, could only be involved in the time graph of the quantity of energy liberated per unit time, since the total magnitude of this quantity is determined by the fire load, i.e. by the quantity of energy which is available from the beginning.

The computer programme which was used for the calculation of the temperature-time curves has a far-advanced general validity. One of the features of this programme is that it affords a possibility of taking into account various factors, viz., first, those thermal properties of the materials entering into the structures bounding the enclosed space which are dependent on the temperature, second, the variations in the dimensions of the openings during the process of fire development, third, the moisture content of the bounding structures, and fourth, the effects of heat-absorbing structures in the interior of the enclosed space. This programme can be used for enclosed spaces which are bounded by structures of up to three different types at the same time, and one of these structures may be built of up to three different materials.

Moreover, a modified programme has been prepared for enclosed spaces provided with plasterboard panel walls, which are assumed under certain definite conditions to disintegrate during the fire.

The time graphs of the quantity of energy liberated per unit time which had been obtained by means of the method outlined in the above were used to calculate the time graphs of the temperature of the combustion gases during the process of fire development. The latter time graphs were computed on the assumption of different values of the fire load and the opening factor for seven types of enclosed spaces which differ in respect of the bounding structures. The results of these calculations are represented in graphs as well as in tables.

In carrying out the comparative theoretical analyses, it was possible to discuss to a limited extent the effects produced on the temperature-time curve of the process of fire development by some quantities which do not directly enter into the equation of heat balance, with the result that their effects on this process must be determined in each individual case. In addition to the size and the shape of the openings, the factors which may be expected to be of importance in this connection comprise, among others, the porosity of the fuel and its distribution in the enclosed space, the moisture content of the fuel, and the magnitude of the fire load. For future research in this field, it may be urgently recommended to make a study of the effects produced by these and other parameters on the quantity of energy released per unit time in the process of combustion, and hence also on the temperature-time curve, which is dependent on this quantity.

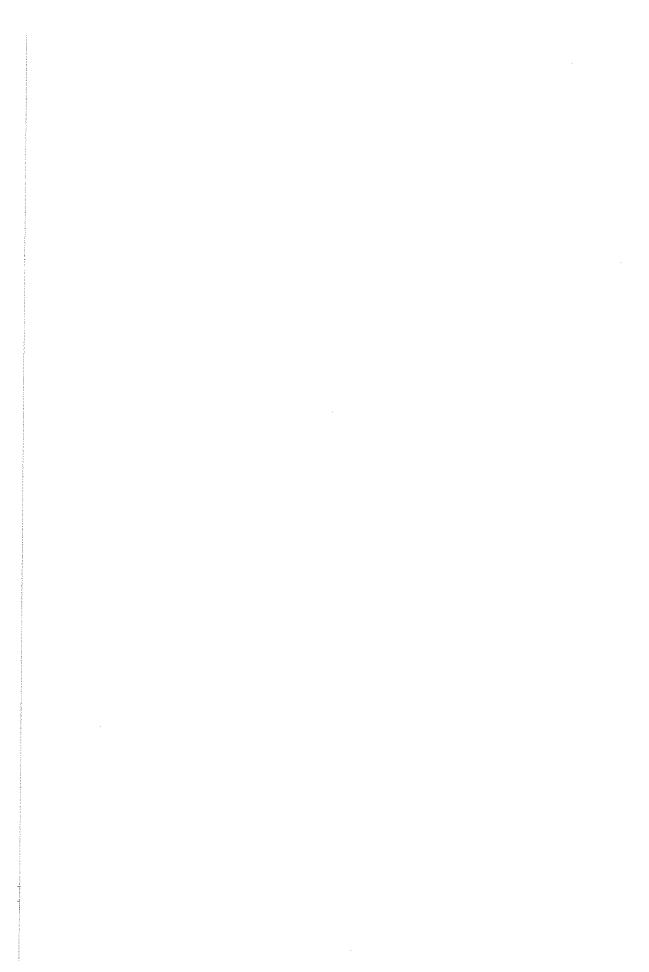
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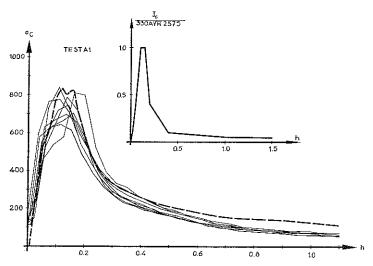
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# Appendix

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series A. See chapter 5.



Test A1

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 34.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 42.2 per cent.

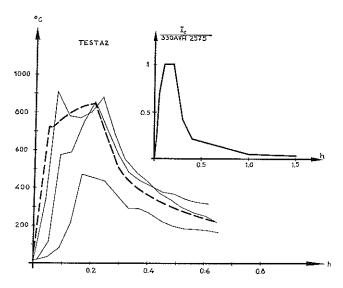
Concrete, 3 cm in thickness+lightweight concrete, 10 cm in thickness, 18.3 per cent.

Window area 4.7 per cent.

Opening factor 0.06 m<sup>1/2</sup> (t > 0.1 h).

Duration of the fire 0.17 h.

Fire load 15.1 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



Test A2

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 34.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 42.2 per cent.

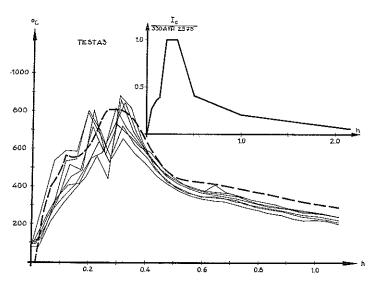
Concrete, 3 cm in thickness+lightweight concrete, 10 cm in thickness, 18.3 per cent.

Window area 4.7 per cent.

Opening factor 0.06 m<sup>1/2</sup> (t > 0.1 h).

Duration of the fire 0.21 h.

Fire load 19 Mcal·m<sup>-2</sup> of bounding surface area.



Test A3

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 38.6 per cent.

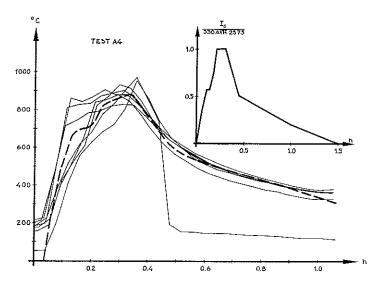
Lightweight concrete, 12.5 cm in thickness, 60.0 per cent.

Window area 1.4 per cent.

Opening factor 0.0356 m<sup>1/2</sup> (t > 0.2 h).

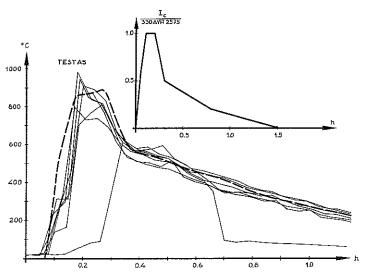
Duration of the fire 0.36 h.

Fire load 19.6 Mcal·m<sup>-2</sup> of bounding surface area.

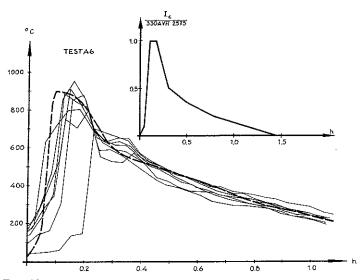


Test A4

Percentages of the total bounding surface area:
Concrete, 20 cm in thickness, 38.3 per cent.
Lightweight concrete, 12.5 cm in thickness, 59.4 per cent.
Window area 2.3 per cent.
Opening factor  $0.0486 \text{ m}^{1/2}$  (t > 0.2 h).
Duration of the fire 0.325 h.
Fire load 22.4 Mcal·m<sup>-2</sup> of bounding surface area.



Test A5
Percentages of the total bounding surface area:
Concrete, 20 cm in thickness, 38.3 per cent.
Lightweight concrete, 12.5 cm in thickness, 59.4 per cent.
Window area 2.3 per cent.
Opening factor  $0.0548~\text{m}^{1/2}$ .
Duration of the fire 0.24~h.
Fire load 20 Mcal·m<sup>-2</sup> of bounding surface area.



Test A6

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 37.8 per cent.

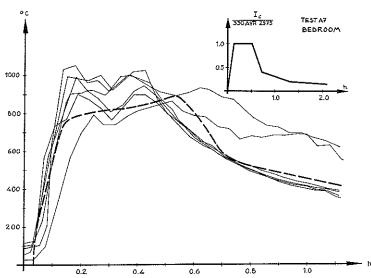
Lightweight concrete, 12.5 cm in thickness, 58.6 per cent.

Window area 3.6 per cent.

Opening factor 0.068 m<sup>1/2</sup>.

Duration of the fire 0.23 h.

Fire load 23 Mcal·m<sup>-2</sup> of bounding surface area.



Test A7a

Living-room.

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 40.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 47.3 per cent.

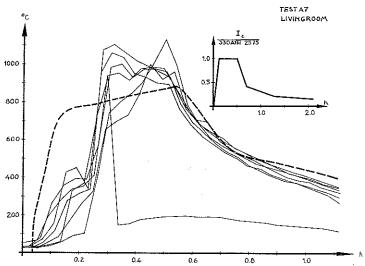
Concrete, 3 cm in thickness+lightweight concrete 10 cm in thickness, 8.3 per cent.

Window area 3.6 per cent.

Opening factor 0.04 m<sup>1/2</sup>.

Duration of the fire 0.32 h.

Fire load 32 Mcal·m<sup>-2</sup> of bounding surface area.



Test A7b

Bedroom.

Percentages of the total bounding surface area:

Concrete, 20 cm in thickness, 40.8 per cent.

Lightweight concrete, 12.5 cm in thickness, 47.3 per cent.

Concrete, 3 cm in thickness+lightweight concrete 10 cm in thickness, 8.3 per cent.

Window area 3.6 per cent.

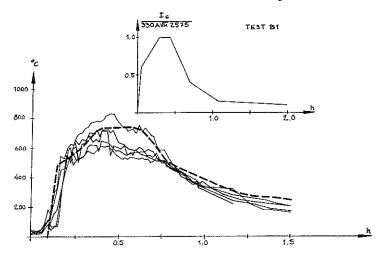
Opening factor 0.04 m<sup>1/2</sup>.

Duration of the fire 0.32 h.

Fire load 32 Mcal·m<sup>-2</sup> of bounding surface area.

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series B. See chapter 5.

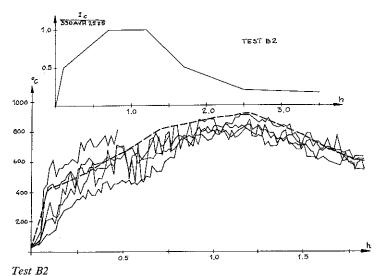


Test B1

Opening factor 0.0467 m<sup>1/2</sup>.

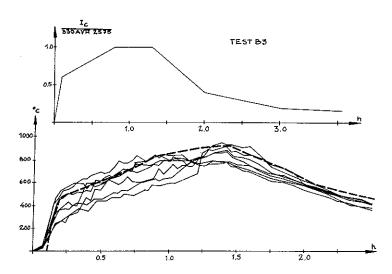
Duration of the fire 0.48 h.

Fire load 33.3 Mcal·m<sup>-2</sup> of bounding surface area.



Opening factor 0.0467 m<sup>1/2</sup>. Duration of the fire 1.07 h.

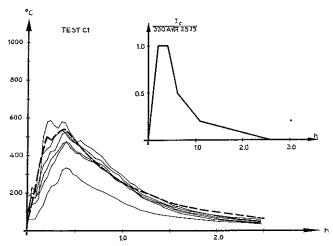
Fire load 75 Mcal·m<sup>-2</sup> of bounding surface area.



Test B3 Opening factor 0.0467 m $^{1/2}$ . Duration of the fire 1.18 h. Fire load 83.5 Mcal·m $^{-2}$  of bounding surface area.

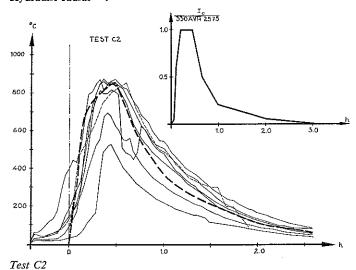
Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series C. See chapter 5.

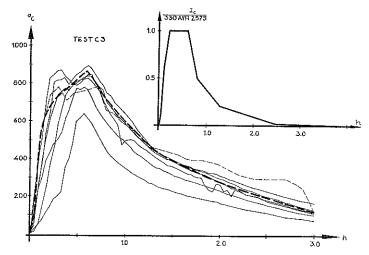


Test C1
Quantity of combustible material 270 kg.
Rate of air supply by fan 1.0 m³/s.
Opening factor 0.0234 m¹ /2.

Hydraulic radius —.



Quantity of combustible material 675 kg. Rate of air supply by fan 2.0 m<sup>3</sup>/s. Opening factor 0.0601 m<sup>1/2</sup>. Hydraulic radius 1.0 cm.



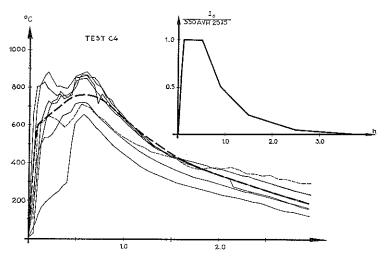
Test C3

Quantity of combustible material 675 kg.

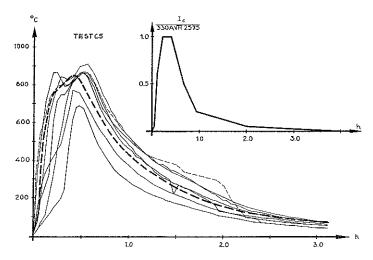
Rate of air supply by fan 1.0 m³/s.

Opening factor 0.0434 m¹/².

Hydraulic radius 1.0 cm.



Test C4 Quantity of combustible material 675 kg. Rate of air supply by fan  $0.7~\mathrm{m}^3/\mathrm{s}$ . Opening factor  $0.0351~\mathrm{m}^{1/2}$ . Hydraulic radius  $1.0~\mathrm{cm}$ .



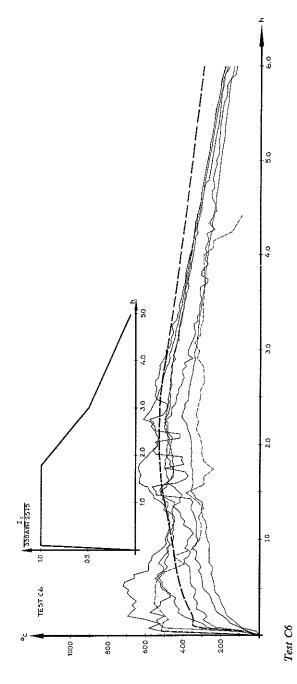
Test C5

Quantity of combustible material 675 kg.

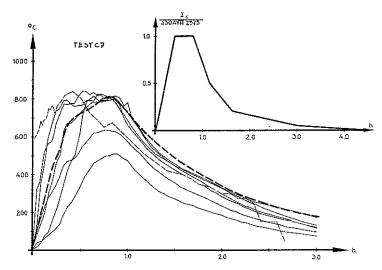
Rate of air supply by fan 1.5 m³/s.

Opening factor 0.0551 m¹/².

Hydraulic radius 1.0 cm.



Quantity of combustible material 675 kg. Rate of air supply by fan 0.25 m³/s. Opening factor 0.01 m<sup>1/2</sup>. Hydraulic radius 1.0 cm.



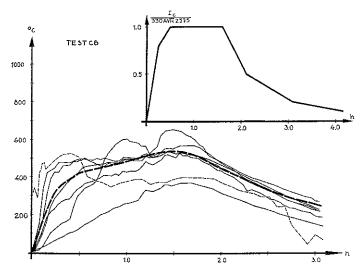
Test C7

Quantity of combustible material 675 kg.

Rate of air supply by fan 1.0 m³/s.

Opening factor 0.0367 m¹/².

Hydraulic radius 1.7 cm.



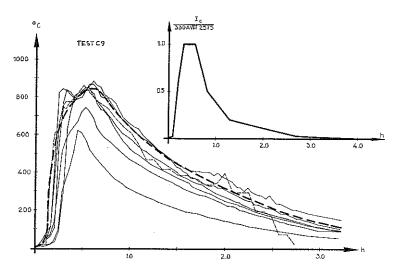
Test C8

Quantity of combustible material 675 kg.

Rate of air supply by fan 0.7 m³/s.

Opening factor 0.015 m¹/².

Hydraulic radius 1.7 cm.



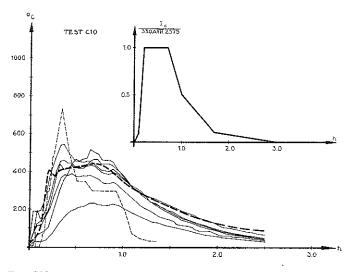
Test C9

Quantity of combustible material 675 kg.

Rate of air supply by fan 1.0 m³/s.

Opening factor 0.051 m¹/².

Hydraulic radius 0.6 cm.



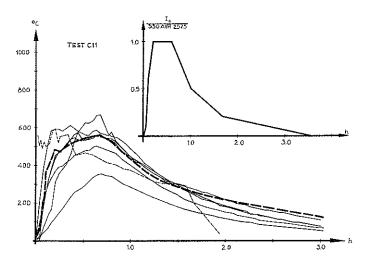
Test C10

Quantity of combustible material 270 kg.

Rate of air supply by fan 0.7 m³/s.

Opening factor 0.015 m¹ 1².

Hydraulic radius 1.0 cm.



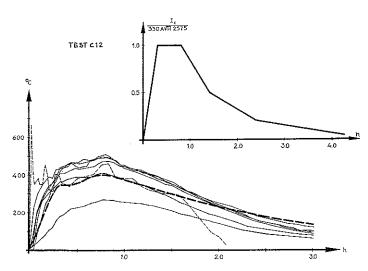
Test C11

Quantity of combustible material 405 kg.

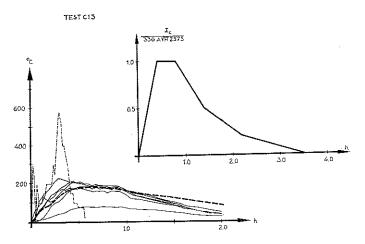
Rate of air supply by fan 0.7 m³/s.

Opening factor 0.02 m¹/².

Hydraulic radius 1.0 cm.



Test C12 Quantity of combustible material 405 kg. Rate of air supply by fan 0.7 m<sup>3</sup>/s. Opening factor 0.012 m<sup>1/2</sup>. Hydraulic radius 2.4 cm.



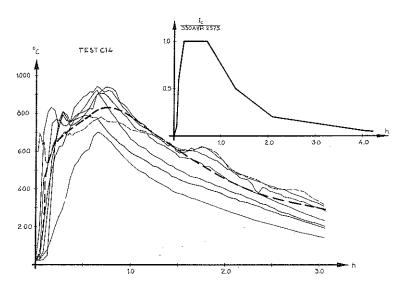
Test C13

Quantity of combustible material 135 kg.

Rate of air supply by fan 0.7 m³/s.

Opening factor 0.005 m¹/².

Hydraulic radius 1.0 cm.



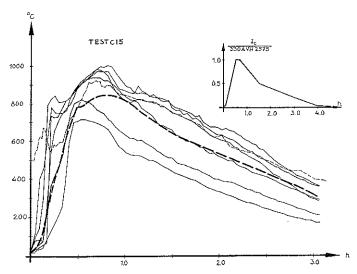
Test C14

Quantity of combustible material 945 kg.

Rate of air supply by fau 1.0 m³/s.

Opening factor 0.0367 m¹/².

Hydraulic radius 1.6 cm.



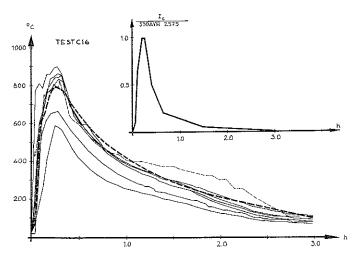
Test C15

Quantity of combustible material 1350 kg.

Rate of air supply by fan 2.0 m³/s.

Opening factor 0.0601 m¹/².

Hydraulic radius 1.4 cm.



Test C16

Quantity of combustible material 405 kg.

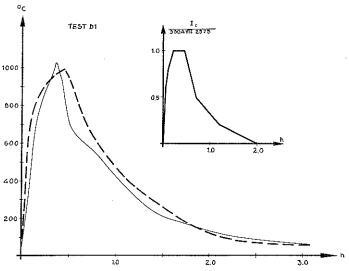
Rate of air supply by fan 0.7 m³/s.

Opening factor 0.06 m¹/².

Hydraulic radius 0.4 cm.

Time graphs of energy released per unit time and corresponding theoretically calculated time graphs of temperature of combustion gases.

Test series D. See chapter 5.



Test DI

Percentages of the total bounding surface area:

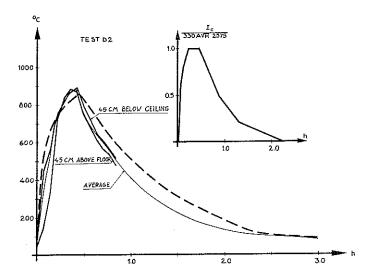
Lightweight concrete, 20 cm in thickness, 97.8 per cent.

Window area 2.2 per cent.

Opening factor 0.028 m<sup>1/2</sup>.

Duration of the fire 0.5 h.

Fire load 20.2 Mcal·m<sup>-2</sup> of bounding surface area.



Test D2

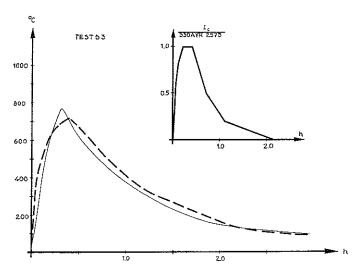
Percentages of the total bounding surface area:
Lightweight concrete, 20 cm in thickness, 98.7 per cent.

Window area 1.3 per cent.

Opening factor  $0.0147 \text{ m}^{1/2}$ .

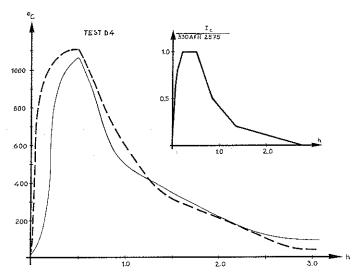
Duration of the fire 0.51 h.

Fire load  $11.2 \text{ Mcal} \cdot \text{m}^{-2}$  of bounding surface area.



Test D3

Percentages of the total bounding surface area:
Lightweight concrete, 20 cm in thickness, 99.15 per cent.
Window area 0.85 per cent.
Opening factor  $0.0092 \text{ m}^{1/2}$ .
Duration of the fire 0.47 h.
Fire load 6.5 Mcal  $\cdot$  m<sup>-2</sup> of bounding surface area.



Test D4

Percentages of the total bounding surface area:
Lightweight concrete, 94.7 per cent.

Window area 5.3 per cent.

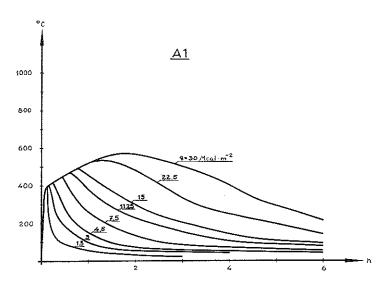
Opening factor 0.075  $\mathrm{m}^{1/2}$ .

Duration of the fire 0.58 h.

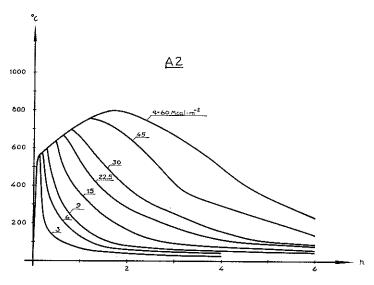
Fire load 65 Mcal· $\mathrm{m}^{-2}$  of bounding surface area.

Calculated time graphs of temperature of combustion gases for seven types of enclosed spaces differing in opening factor and in bounding structures.

See chapter 7.



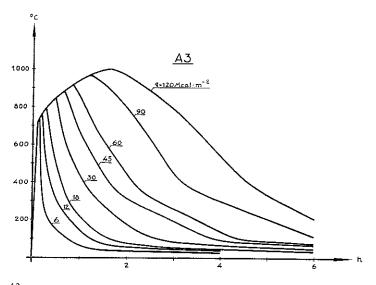
AI Type A enclosed space. Opening factor 0.01 m<sup>1</sup>/<sup>2</sup>.



A2

Type A enclosed space.

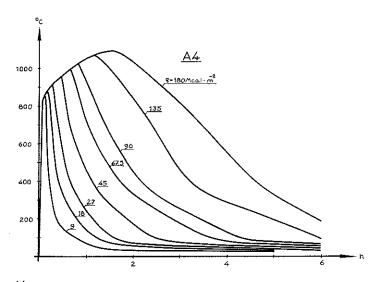
Opening factor 0.02 m<sup>1/2</sup>.



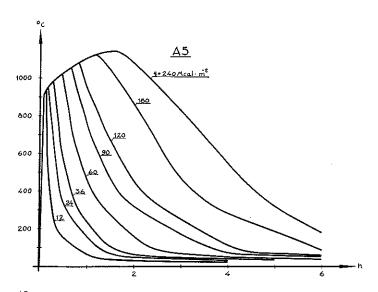
A3

Type A enclosed space,

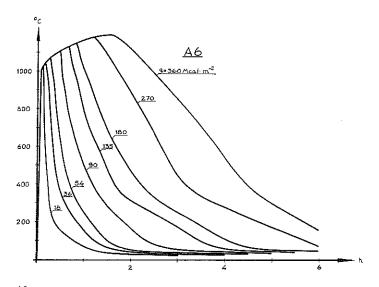
Opening factor 0.04 m<sup>1/2</sup>.



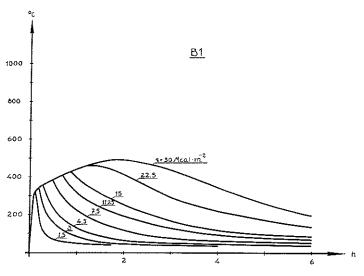
A4 Type A enclosed space. Opening factor  $0.06 \text{ m}^{1/2}$ .



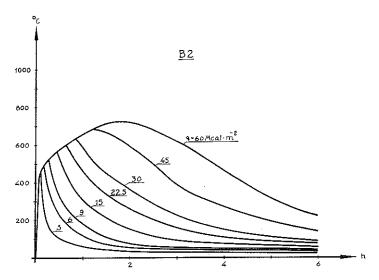
A5 Type A enclosed space. Opening factor 0.08  $m^{1/2}$ .



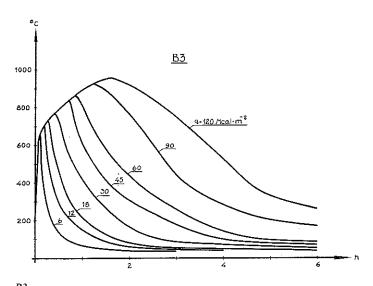
A6
Type A enclosed space.
Opening factor 0.12 m<sup>1</sup>/<sup>2</sup>.



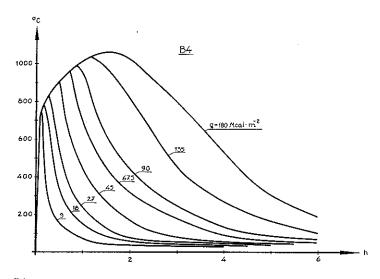
B1 Type B enclosed space. Opening factor 0.01 m $^{1/2}$ .



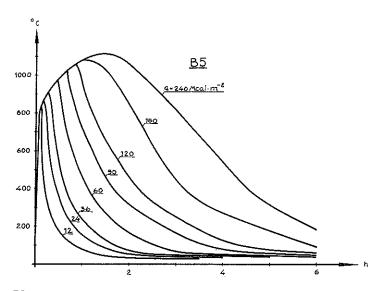
B2
Type B enclosed space.
Opening factor 0.02 m<sup>1/2</sup>.



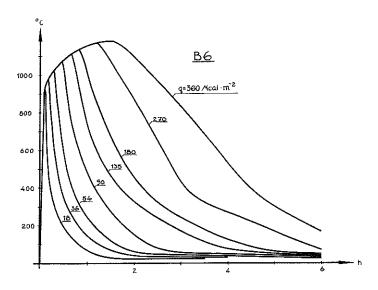
B3 Type B enclosed space. Opening factor 0.04  $m^{1/2}$ .



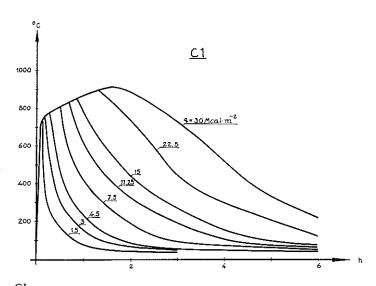
B4 Type B enclosed space. Opening factor 0.06 m<sup>1/2</sup>.



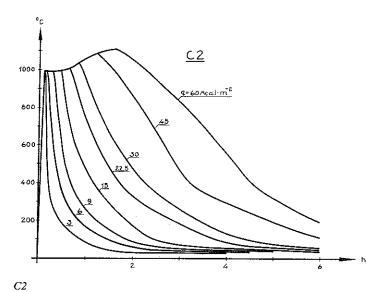
B5 Type B enclosed space. Opening factor 0.08 m<sup>1/2</sup>.



B6 Type B enclosed space. Opening factor 0.12  $m^{1/2}$ .

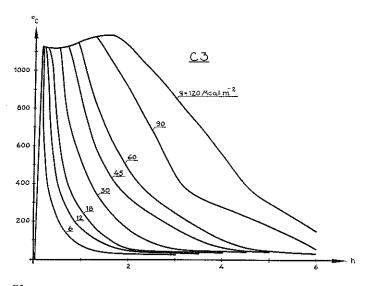


CI Type C enclosed space. Opening factor 0.01  $\mathrm{m}^{1/2}$ .

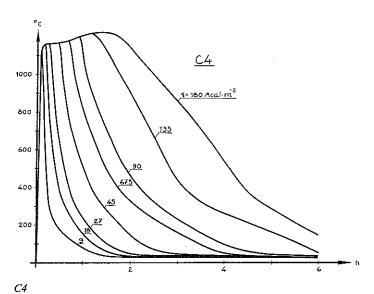


Type C enclosed space.

Opening factor 0.02 m<sup>1/2</sup>.

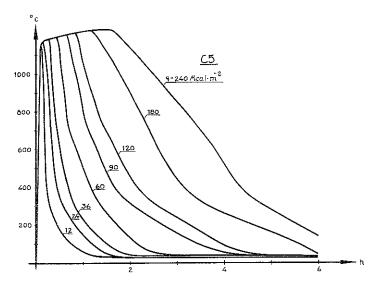


C3 Type C enclosed space. Opening factor 0.04  $\mathrm{m}^{1/2}$ .



Type C enclosed space.

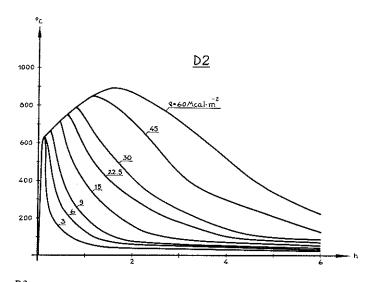
Opening factor 0.06 m<sup>1</sup>/<sup>2</sup>.



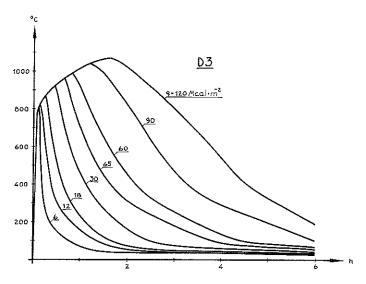
C5

Type C enclosed space.

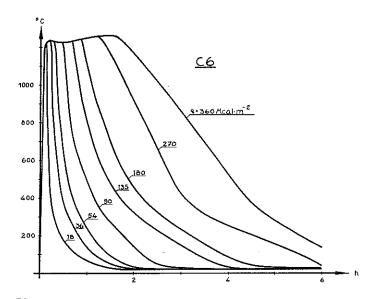
Opening factor 0.08 m<sup>1/2</sup>.



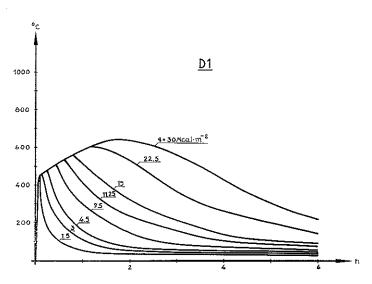
D2 Type D enclosed space. Opening factor 0.02  $\mathrm{m}^{1/2}$ .



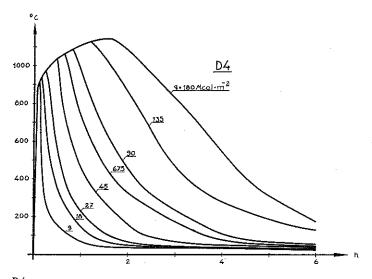
D3 Type D enclosed space. Opening factor 0.04  $\mathrm{m}^{1/2}$ .



C6 Type C enclosed space. Opening factor  $0.12 \ m^{1/2}$ .



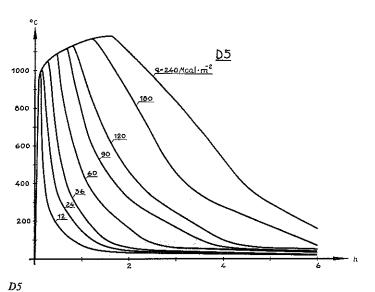
DI Type D enclosed space. Opening factor 0.01  $m^{1/2}$ .



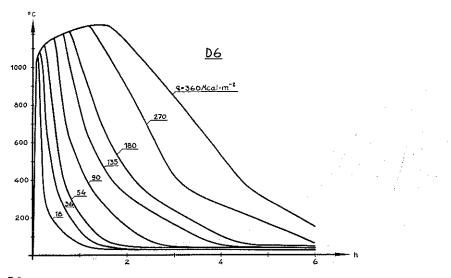
D4

Type D enclosed space.

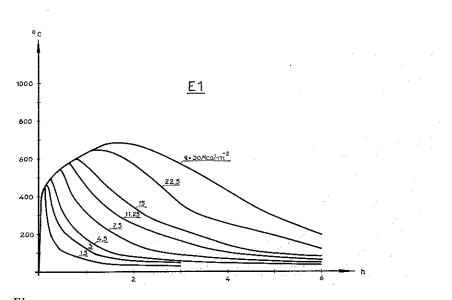
Opening factor 0.06 m<sup>1/2</sup>.



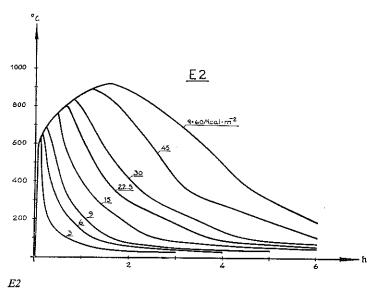
Type D enclosed space.
Opening factor 0.08 m<sup>1/2</sup>.



D6 Type D enclosed space. Opening factor 0.12  $m^{1/2}$ .

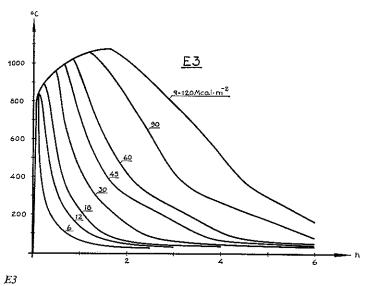


E1 Type E enclosed space. Opening factor 0.01  $m^{1/2}$ .



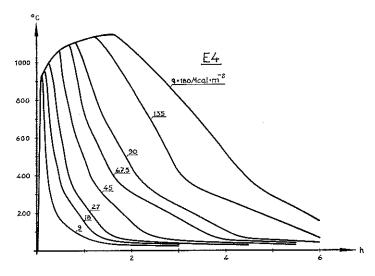
Type E enclosed space.

Opening factor 0.02 m<sup>1/2</sup>.

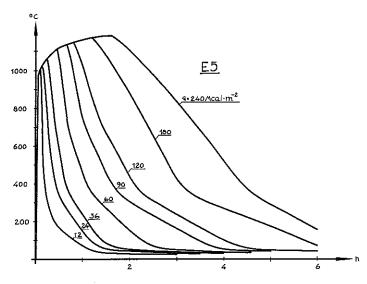


Type E enclosed space.

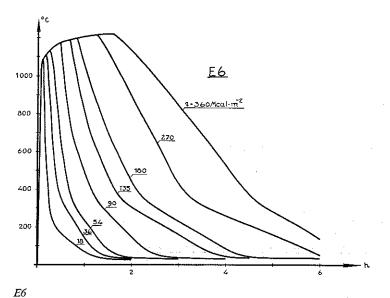
Opening factor 0.04 m<sup>1/2</sup>.



E4 Type E enclosed space. Opening factor  $0.06~\text{m}^{1/2}$ .

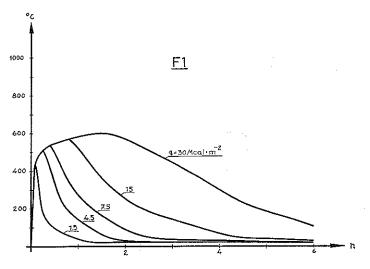


E5 Type E enclosed space. Opening factor  $0.08~\mathrm{m}^{1/2}$ .



Type E enclosed space.

Opening factor 0.12 m<sup>1/2</sup>.

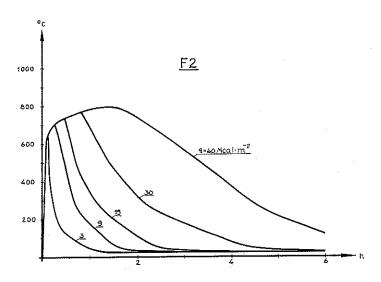


FI

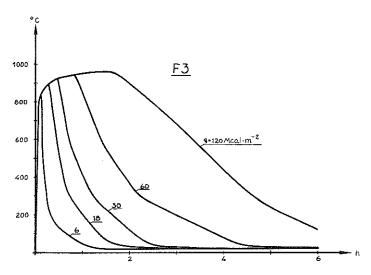
Type F enclosed space.

Resultant emissivity 0.1.

Opening factor 0.01  $m^{1/2}$ .



F2
Type F enclosed space.
Resultant emissivity 0.1.
Opening factor 0.02 m<sup>1/2</sup>.

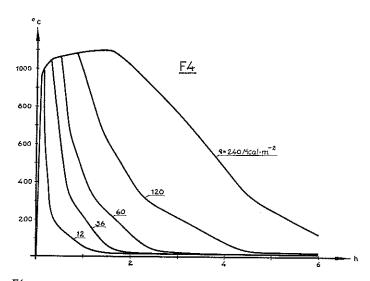


F3

Type F enclosed space.

Resultant emissivity 0.1.

Opening factor 0.04 m<sup>1/2</sup>.

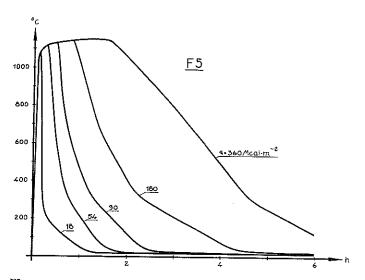


F4

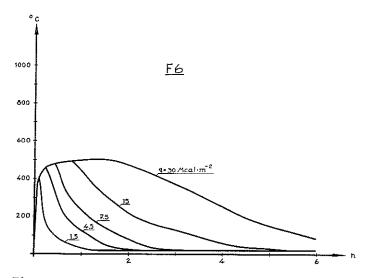
Type F enclosed space.

Resultant emissivity 0.1.

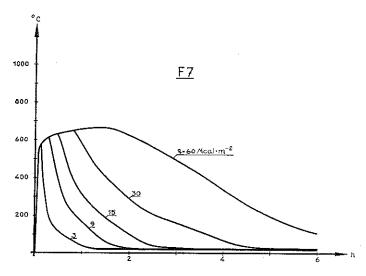
Opening factor 0.08 m<sup>1/2</sup>.



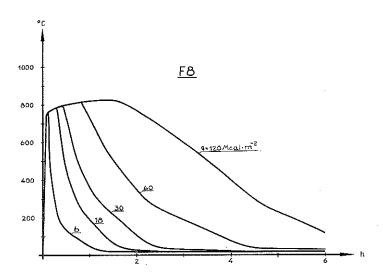
F5
Type F. enclosed space.
Resultant emissivity 0.1.
Opening factor 0.12 m<sup>1</sup>/<sup>2</sup>.



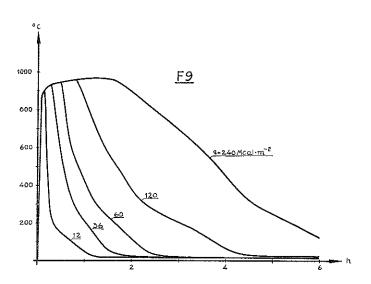
F6
Type F enclosed space.
Resultant emissivity 0.35.
Opening factor 0.01 m<sup>1/2</sup>.



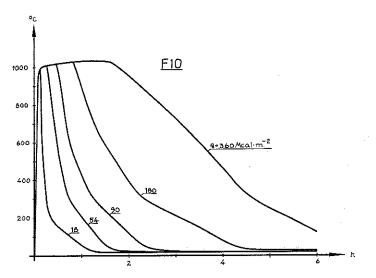
F7 Type F enclosed space. Resultant emissivity 0.35. Opening factor 0.02  $m^{1/2}$ .



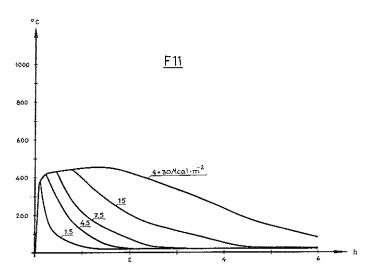
F8 Type F enclosed space. Resultant emissivity 0.35. Opening factor 0.04  $m^{1/2}$ .



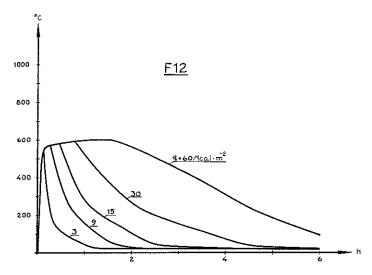
F9
Type F enclosed space.
Resultant emissivity 0.35.
Opening factor 0.08 m<sup>1</sup>/<sub>2</sub>.



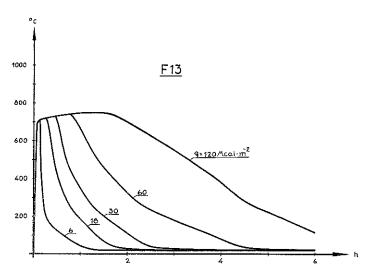
F10
Type F enclosed space.
Resultant emissivity 0.35.
Opening factor 0.12 m<sup>1/2</sup>.



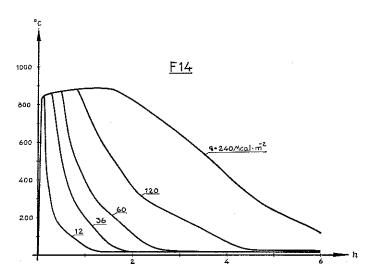
F11
Type F enclosed space.
Resultant emissivity 0.60.
Opening factor 0.01 m<sup>1/2</sup>.



F12
Type F enclosed space.
Resultant emissivity 0.60.
Opening factor 0.02 m<sup>1/2</sup>.



F13 Type F enclosed space. Resultant emissivity 0.60. Opening factor 0.04  $\mathrm{m}^{1/2}$ .

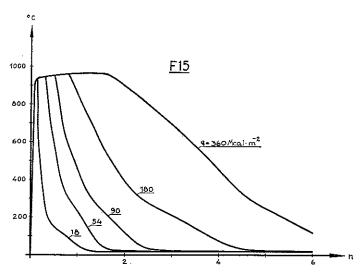


F14

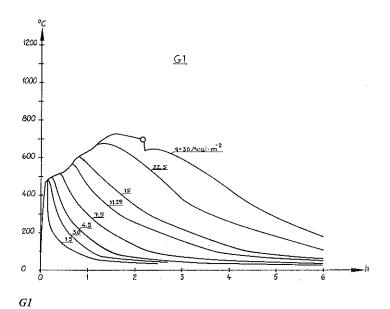
Type F enclosed space.

Resultant emissivity 0.60.

Opening factor 0.08 m<sup>1/2</sup>.



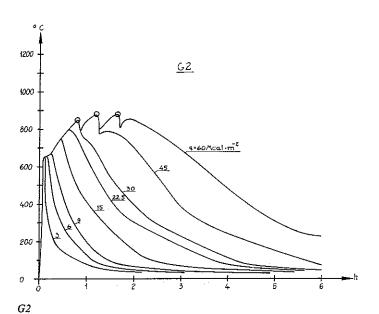
F15 Type F enclosed space. Resultant emissivity 0.60. Opening factor 0.12  $\mathrm{m}^{1/2}$ .



Type G enclosed space.

Opening factor 0.01 m<sup>1/2</sup>.

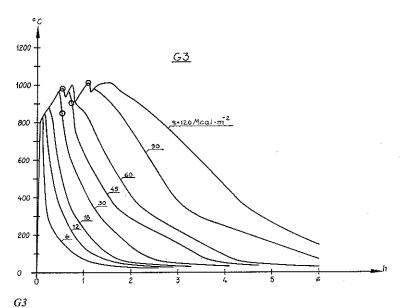
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.02 m<sup>1/2</sup>.

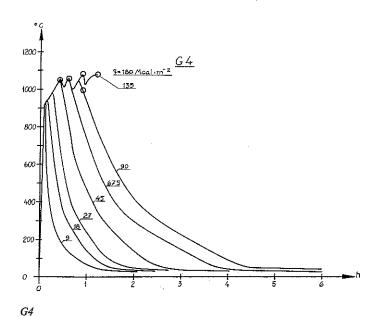
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.04 m<sup>1/2</sup>.

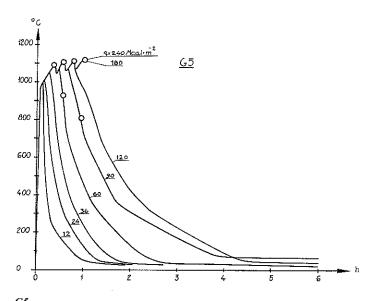
O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.06 m<sup>1/2</sup>.

O=Plasterboard panel falls down.

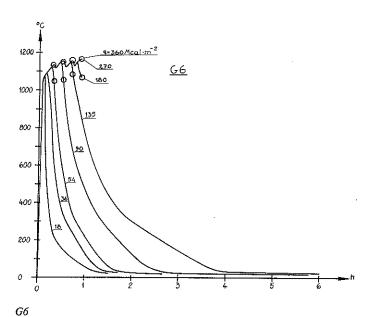


G5

Type G enclosed space.

Opening factor  $0.08 \text{ m}^{1/2}$ .

O=Plasterboard panel falls down.



Type G enclosed space.

Opening factor 0.12 m<sup>1/2</sup>.

O=Plasterboard panel falls down.

## APPENDIX 6

Calculated time graphs of temperature of combustion gases represented in tabular form for seven types of enclosed spaces differing in opening factor and in bounding structures. See chapter 7.



 $\frac{\textbf{A1}}{\text{Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space.}}$  Opening Factor A. $\sqrt{\text{H}}/\text{A}_{\text{t}}$  = 0.01 m<sup>1/2</sup>

0.05	-	τ							
Time h  0.05	T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
10.05   272	q.	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
0.10 395 395 395 395 395 328 328 328 328 328 0.15 228 390 390 390 360 360 360 360 360 360 0.20 196 368 401 401 406 406 406 406 406 6.25 150 313 409 410 405 405 405 405 405 6.30 0.30 98 257 385 421 415 415 415 415 6.35 425 425 425 6.40 94 218 320 437 434 434 434 434 434 434 6.34 434 6.36 6.36		T	e m	ı p	е	r a	a t	u	r e
5.80 35 42 58 78 94 151 231 6.00 34 41 57 77 92 140 216	0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80	395 395 196 99 99 88 79 59 14 40 39 30 31 31 31 31 31 31 31 31 31 31 31 31 31	395 396 397 398 398 399 397 397 397 397 397 397 397 397 397	395 390 409 354 352 269 211 1664 167 168 168 168 169 169 169 169 169 169 169 169 169 169	395 390 401 421 429 437 441 370 355 314 412 370 355 314 328 329 190 160 150 170 83 80 77 77 68 66 65 66 65 67 68 66 67 67 68 68 68 68 68 68 68 68 68 68 68 68 68	328 360 405 405 429 449 463 463 463 463 463 463 463 463 463 463	328 360 405 405 421 446 446 456 421 468 456 439 446 456 439 446 456 439 446 456 456 456 456 456 456 456 456 456	3260 4055 4065 4015 4121 4121 4132 4132 4133 4144 4155 4155 4155 4151 4152 4151 4152 4151 4152 4151 4152 4151	328 3606 405 449 46706 4800 4768 5511 410 555 555 555 555 555 555 555 555 555 5

<u>A2</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.02 m<sup>1/2</sup>

**A3** 

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{\rm H}/\rm A_t$  = 0.04 m  $^{1/2}$ 

	t							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>q</b>	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time h	T	e m	p	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.10 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	504 745 422 360 268 164 162 155 148 120 114 100 87 14 100 87 14 45 43 41 40 33 33 33 31	504 745 747 696 747 696 747 437 438 281 281 281 281 281 281 281 281 281 28	504 745 747 767 7665 513 307 5665 513 307 568 307 568 307 568 307 568 307 568 307 568 307 568 307 568 307 568 307 568 307 568 307 568 308 308 308 308 308 308 308 308 308 30	504 7477684 7684 7684 76826 55273 331666 55554 4383 319666 55555 4383 3197 1130 1098 109666 1097 1098 1098 1098 1098 1098 1098 1098 1098	504 621 7776 802 838 848 877 645 555 480 400 831 832 831 841 852 853 853 854 854 855 855 854 854 855 855 856 857 857 857 857 857 857 857 857 857 857	504 621 7776 808 818 818 818 818 818 818 818 818 818	504 621 681 777 793 802 836 848 912 928 942 942 942 942 942 943 944 944 944 944 944 944 944 944 944	504 621 7776 802 838 848 912 912 913 914 915 915 915 915 915 915 915 915 915 915

<u>A4</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.06 m<sup>1/2</sup>

	L							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	p	е	r a	t	u	r e
h  0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.60 3.80 4.00 4.20	5758 494 496 1751 1669 1511 136 120 104 89 71 148 337 36 333 331 30 29	575 858 861 862 679 863 430 430 363 363 262 247 215 262 247 215 262 247 262 247 262 247 262 247 262 247 262 247 262 247 262 262 262 263 264 265 265 265 265 265 265 265 265 265 265	5758 8619 8798 87669 8798 87669 87669 8777661 8777661 8777661 8777661 8777661 8741 8777661 8741 8741 8741 8741 8741 8741 8741 874	5758 861 8798 8914 928 9547 7694 620 466 404 327 303 281 291 107 107 107 107 107 107 107 107 107 10	575 704 882 889 908 923 934 962 939 872 705 705 705 705 705 705 705 705 705 705	5754 882 936 941 991 991 991 991 1018 991 1018 991 1018 991 1018 1018	575 704 784 882 889 908 949 961 1018 1054 1064 1029 1013 9966 830 756 830 524 381 355 331 308 269	575 704 784 882 890 908 923 949 961 982 10018 1054 1064 1070 1087 1087 1083 1082 943 943 943 956 756 597 541 483
4.40 4.60		32 31	38 37	53 51	73 70	90 85	249 230	423 377
4.80 5.00		31 30	36 35	49 47	67 64	81 77	211 193	348 319
5.20 5.40			35 34	46 44	62 59	74 71	174 155	292 265
5.60			33	43	57	69	135	238
5.80 6.00			32 32	42 41	55 54	6ր 66	114 94	210 185

<u>A5</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.08 m $^{1/2}$ 

	ı							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>q</b>	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	Ŧ	e m	Ď	е	r e	a t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.60 5.60 5.60 5.80 6.00	622 935 532 432 314 181 163 155 139 131 122 106 89 70 44 42 40 39 37 36 35 34 33 31 30 29 27	622 935 937 869 734 575 521 454 386 253 219 186 151 113 70 662 59 55 51 40 38 37 33 31 32 31	622 935 937 957 957 961 953 818 72 961 961 961 961 961 961 961 961 961 961	234	622 766 853 955 981 1050 1050 833 735 652 5514 415 347 3316 2236 917 728 552 447 347 347 347 347 347 347 347 347 347	622 766 853 959 965 1020 1058 1066 1065 897 618 7695 657 454 454 454 454 454 454 454 454 454 4	622 766 853 959 965 981 1050 1058 1056 1081 1092 1111 1077 1058 1038 1004 969 932 853 8774 434 373 302 282 242 223 166 125 104 82	622 767 853 959 965 982 995 1008 1058 1058 1066 1081 1192 1111 1119 1126 1138 1143 1105 1074 1074 1078 873 873 879 1016 873 873 874 875 876 877 878 878 878 878 879 879 879 879 879

<u>A6</u>

Time Graphs of Temperature of Combustion Gases. Type A Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.12 m<sup>1/2</sup>

	t							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>q</b>	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	T	e m	p	e	r s	ı t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.60 3.80 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.50 6.60 6.50 6.60 6.50 6.60 6.50 6.60 6.65 6.70 6.80 6.90 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.10 2.10 3.60 3.80 3.60 3.80 4.10 5.00 5.10	670 1027 581 465 333 186 185 176 168 159 142 133 124 106 88 67 41 39 37 36 31 32 31 30 29	670 1027 1033 951 799 620 556 480 404 324 2275 221 186 149 107 63 50 47 45 39 37 35 33 32	670 1027 1033 1049 981 1063 981 1063 982 774 650 5932 407 244 211 108 69 67 63 60 57 54 40 88 88 88 88 88 88 88 88 88 88 88 88 88	670 1027 1033 1049 1063 1076 1088 1107 681 1098 7681 622 556 422 351 7681 2538 203 150 64 42 21 2538 203 155 44 44 45 45 45 45 46 46 47 47 48 48 48 48 48 48 48 48 48 48 48 48 48	670 847 933 1051 1057 1071 1083 11094 11127 1133 1060 971 873 765 680 628 572 414 358 343 222 220 173 144 186 49 41 42 42 43 44 44 42 43 43 44 44 45 46 46 47 47 48 48 48 48 48 48 48 48 48 48 48 48 48	670 847 933 1057 1071 1083 11094 11127 1133 1150 1162 1001 937 868 794 762 762 762 762 762 763 763 763 763 763 763 763 763 763 763	670 847 933 1051 1057 1071 1083 1103 11127 1133 1139 1150 1166 1173 1178 1128 1106 1043 1003 962 919 874 789 704 613 361 336 313 273 273 273 273 273 273 273 273 273 27	670 847 933 1057 1071 1083 1194 1103 1127 1133 1150 1159 1166 1173 1188 1195 1181 1195 1195 1195 1195 1196 891 891 660 894 403 355 325 297 269 261 261 261 261 261 261 261 261 261 261
0.00				21	20	45	01	1,0

<u>B1</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$ 

•	t							0.0
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>q</b>	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time	T	e m	$\mathbf{p}$	е	r a	t	u	r e
h  0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.80 6.00	237 342 201 172 132 86 85 76 67 67 67 67 59 53 40 33 33 33 33 32	237 332 332 332 332 331 332 333 333 333 333	234398655895388112211111111111111111111111111111111	237 342 339 357 366 378 357 304 126 322 266 1534 121 121 121 121 121 121 121 121 121 12	237 236 314 352 369 372 388 404 387 345 240 255 240 207 209 209 209 209 209 209 209 209 209 209	237 284 314 352 360 372 388 400 400 402 337 313 322 313 227 266 314 410 410 410 410 410 410 410 410 410 4	237 236 314 352 369 372 388 400 410 420 433 456 451 443 443 457 407 363 331 258 217 208 217 218 217 218 217 218 217 218 217 218 217 218 217 218 217 218 218 218 218 218 218 218 218 218 218	237 286 314 352 366 377 388 400 410 410 410 410 410 410 410 410 410

<u>B2</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A. $\sqrt{E}/A_{\rm t}$  = 0.02 m $^{1/2}$ 

	τ								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
ď	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0	
Time h	T	е п	р	е	r a	t	u	r	e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.20 4.40 4.60 4.80 5.00 5.80 6.00	3538 498 288 218 110 105 60 878 988 454 421 40 988 33 33 33 33 33 31	3538 4875 3223 303 2240 455 229 1873 5776 5776 5777 5776 57776 5777 5777 5	354 480 5182 5182 5183 5184 5186 5186 5186 5186 5186 5186 5186 5186	354 487 554 555 554 439 337 319 385 319 319 319 319 319 319 319 319 319 319	3544 5555555555555555555555555555555555	354 4151 5116 5116 5116 5116 5116 5116 51	355 411 510 500 555 560 613 663 664 665 664 665 665 665 665 665 665 665	356 451 550 553 564 555 564 667 716 663 716 663 663 663 663 663 663 663 663 663 6	

<u>B3</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A. $\sqrt{\rm H}/\rm A_t$  = 0.04 m  $^{1/2}$ 

	-							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
Q.	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time h	T	e m	р	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.40 1.50 1.40 1.50 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.40 2.40 2.50 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00	481 675 384 331 250 155 139 1106 94 46 44 40 338 337 331 332 31	481 6758 6555 4304 363 2429 21967 111 818 777 675 660 854 453 453 453 453 453 453 453 453 453 4	481 6758 6758 6758 6758 6758 6758 6758 6758	482 6758 6796 7757 7110 638 452 452 453 453 454 455 455 455 455 455 455 455	482 5576 6988 7727 7769 8082 7335 7355 4854 433 341 3273 3306 3313 3076 3076 3076 3077 3077 3077 3077 307	482 5576 693 688 705 776 802 831 776 802 814 701 803 814 701 803 804 804 805 807 807 807 807 808 807 807 807 807 808 809 807 807 807 808 809 809 809 809 809 809 809 809 809	457 558 606 693 688 709 778 806 822 855 876 893 907 919 889 878 878 856 833 778 862 877 919 878 878 878 878 878 878 878 878 878 87	457866389529988895299999999988887740316942940051168888569400529999999888877403166113354436994005116
0.00		24	Ju	71	00	- 00	TT	200

 $$\underline{B4}$$  Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space Opening Factor A. $\sqrt{H}/A_t$  = 0.06 m  $^{1/2}$ 

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
g .	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	p	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	528 761 378 281 167 159 1150 125 1150 1150 1150 1150 1150 1150	528 766 766 720 766 460 460 460 460 460 460 460 460 460 4	528 761 766 798 818 702 702 715 82 83 75 82 83 75 84 85 85 85 85 85 85 85 85 85 85 85 85 85	528 766 798 818 839 875 890 895 890 895 897 897 897 897 897 897 897 897 897 897	528 632 7036 8055 8059 9359 9456 612 5334 4015 3404 3532 2018 1533 100 9851 857 7666 666 666 666 666 666 666 666 666	528 708 708 808 809 909 909 909 909 909 9	528 632 703 806 805 829 905 942 909 909 909 909 909 909 909 909 909 90	528 632 703 8065 8079 9099 9091 1032 10559 10012 10579 10012 1057 10012 9013 9013 9013 9013 9013 9013 9013 9013

<u>B5</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A  $\sqrt{\pi}/A_{\rm t}$  = 0.08 m<sup>1/2</sup>

	U							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>Q</b>	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	T	e m	р	e	r a		u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.60 5.80 6.00	576 836 480 409 301 177 159 151 135 127 45 42 40 39 38 36 35 34 33	576 836 45 798 45 437 437 437 437 437 437 437 437 437 437	576 834 877 847 178 847 178 847 178 178 178 178 178 178 178 178 178 17	5736 844 8717 899 984 11332 2774 1138 8776 666 5753 1474 1494 1494 1494 1494 1494 1494 1494	576 693 773 871 886 910 935 973 990 1018 1027 968 891 597 552 505 457 408 312 205 180 126 98 177 370 67 661 597 553 57 553 57 553 57 57 57 57 553 57 57 57 57 57 57 57 57 57 57 57 57 57	576 673 871 873 886 9955 9918 1035 9918 1035 9916 605 605 508 488 3304 118 850 773 168 663 118 850 773 168 663 168 663	576 694 773 871 886 935 973 1018 1021 1036 1037 1038 1019 1038 1019 1038 1019 1038 1019 1038 1038 1038 1038 1038 1038 1038 1038	576 694 773 871 886 910 935 973 990 1018 1056 1057 1066 1077 1087 1112 1112 1088 1073 1058 1073 1058 1073 1058 1075 1112 1075 1112 1075 1112 1075 1112 1075 1112 1075 1075 1075 1075 1075 1075 1075 1075

<u>B6</u>

Time Graphs of Temperature of Combustion Gases. Type B Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.12 m<sup>1/2</sup>

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>q</b>	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time h	Т	e m	ą	e	r a	. t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.60 5.60 5.60 6.60	624 940 548 447 324 185 174 166 157 140 132 105 67 43 33 33 31 31	624 940 952 892 7600 541 370 389 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 289 275 276 277 289 277 277 277 277 277 277 277 277 277 27	624 940 952 980 1945 1945 1946 1946 1947 1940 1940 1940 1940 1940 1940 1940 1940	6240 940 9520 100518 10083 100	624 777 867 980 9019 1040 1076 1040 1040 1040 1040 1040 1040 1040 104	624 777 867 980 995 1040 1076 1089 1120 1131 1086 1120 1131 1087 1120 1131 1087 1131 1087 1131 1088 1131 1088 1131 1088 1131 1088 1131 1088 1131 1088 1131 1088 1131 1088 1088	624 777 867 980 995 1040 1060 1076 1133 1152 1159 1165 1116 1093 1159 908 1159 908 780 698 517 423 425 425 425 425 425 425 425 425 425 425	624 777 867 980 995 1040 1060 1133 1152 1159 1165 1171 1185 1185 1185 1185 1185 1185 118

<u>C1</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor A+ $\sqrt{H}/A_{\rm t}$  = 0.01 m<sup>1/2</sup>

	_							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	T	е ш	p	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.40 1.50 1.40 1.50 1.40 2.00 2.40 2.40 2.40 2.40 2.40 3.60 3.80 3.00 3.20 3.40 3.60 3.80 4.00 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00 5.80 6.00	502 778 430 347 255 156 142 137 129 113 101 88 75 60 57 44 44 42 41 39 37 33 32	502 778 769 769 417 368 319 222 221 151 23 90 85 177 168 50 65 50 47 40 40 33 33 33 33 33 33 33 33 33 33 33 33 34 34	502 7769 7719 7719 5688 5779 669 674 674 675 675 675 675 675 675 675 675 675 675	57893794073806553961334577899103466949518555555544777940738065539613345778991034669495185555555447	502 630 696 7776 788 801 808 821 787 801 808 821 7668 404 409 409 409 409 409 409 409 409 409	502 630 696 777 788 808 821 588 836 6436 557 549 431 840 841 840 841 840 841 841 841 841 841 841 841 841 841 841	502 630 67776288 503 67776288 803 803 804 804 804 805 806 807 806 807 807 808 808 807 808 808 808 808 808	502 630 696 777 788 808 821 788 828 845 856 869 890 897 897 897 7720 897 897 897 7720 897 897 897 897 897 897 897 897 897 897

<u>C2</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor  $\Lambda \cdot \sqrt{H}/A_{\rm t}$  = 0.02 m  $^{1/2}$ 

	Ų							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time h	T	e m	р	e	r s	. t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.665 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.40 1.50 1.40 2.60 2.80 3.80 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 6.60	634 992 538 425 303 175 167 160 153 138 130 122 107 91 73 50 44 41 39 36 35 33 31 30 28	634 992 977 889 733 5604 439 373 275 260 2188 73 662 73 45 42 40 33 33 31 31	634 9977 982 9877 988 9877 988 989 989 989 989 989 9	639 977 989 999 990 990 990 990 990 990 990 990	6305 8305 8305 8305 8305 8305 8305 8305 8	634 805 808 988 989 999 999 1001 901 1010 901 1010 901 1010 901 1010 901 1010 901 901	635 886 987 989 989 999 1016 1032 1057 1069 1079 1044 1010 1057	634 805 884 986 981 989 995 997 999 1003 1010 1016 1032 1045 1108 1079 1108 1076 1108 1076 1108 1076 1108 1076 1108 1076 1108 1076 1078 1078 1078 1078 1078 1078 1078 1078

250000000000000000000000000000000000000	*,t								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
<b>Q</b>	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time	T	e	m p	е	r	a t	u	r	е
h									
0.05	726	726	726	726	726	726	726	726	
0.10	1126	1126	1126	1126	916	916	916	916	
0.15	618	1110	1110	1110	1002	1002	1002	1002	
0.20 0.25	478 334	1005 829	1111 1112	1112 1111	1122 1111	1122 1111	1122 1111	1122 1111	
0.30	183	631	1013	1114	1113	1113	1113	1113	
0.35	182	559	902	1115	1115	1115	1115	1115	
0.40	174	480	784	1116	1118	1118	1118	1118	
0.45	166	401	651	1115	1117	1117	1117	1117	
0.50	158	320	593	1005	1115	1115	1115	1115	
0.60	141	290	469	852	1126	1126 1132	1126 1132	1126 1132	
0.65 0.70	132 123	273 256	404 338	770 679	1131 1058	1137	1137	1137	
0.80	106	221	308	621	970	1148	1148	1148	
0.90	88	186	276	557	874	1061	1158	1158	
1.00	67	149	244	491	767	1002	1166	1166	
1.10	41	107	211	424	684	938	1172	1172	
1.20	39	63	178	353	631	870	1178	1178	
1.30 1.40	37	60	144	328	579	797		1182	
1.50	35 34	55 52	107 68	302 278	526 473	717 676	1105 1081	1184 1187	
1.60	32	48	66	254	418	631	1043	1189	
1.70	31	46	62	229	361	589	1003	1146	
1.80	30	43	58	20 <sup>1</sup> 4	345	545	962	1128	
1.90	29	41	54	180	328	500	919	1110	
2.00	28	39	<u>51</u>	155	312	456	875	1091	
2.20		36 34	46 42	100 69	283	363	790	1045 996	
2.60		34 32	39	6 <u>9</u>	255 227	331 300	703 614	996 945	
2.80		30	37	55	199	272	521	890	
3.00		29	35	51	172	244	425	834	
3.20		•	33	47	143	216	362	780	
3.40			31	74.74	112	189	336	722	
3.60			30	41	81	161	312	661	
3.80 4.00			29 28	39	61	131	291	599 535	
4.20		•	20	37 36	56 51	100 66	271 251	535 471	
4.40				34	48	62	231	404	
4.60				33	45	57	211	355	
4.80				32	43	53	193	324	
5.00				31	41	50	173	295	
5.20					39	47	153	267	
5.40					38	45 1.0	132	239	
5.60 5.80					36 35	43 41	110 89	210 184	
6.00					35 34	39	64	155	
					24	37	54	エノノ	

\$\begin{align\*} \text{C4} \\ \text{Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor \$A.\frac{\pi}{H}/A\_t = 0.06 \text{ m}^{1/2} \end{align\*}\$

	•							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time	T	e m	p	е	r s	, t	u	r e
h								
0.05 0.10 0.25 0.30 0.25 0.30 0.45 0.50 0.665 0.70 0.80 0.90 1.00 1.30 1.50 1.60 1.70 1.80 1.90 2.24 2.80 3.40 3.40 4.40 4.80 5.50 6.50	766 1183 650 500 346 185 176 168 159 142 123 105 64 35 32 29 28 27 26 25	766 1183 1165 1053 867 659 581 498 411 323 292 275 220 183 144 101 53 49 45 40 38 36 34 33	766 1183 1165 1165 1166 1060 941 815 674 408 3307 273 240 206 1737 956 54 42 38 333 31 29	766 1183 1165 1166 1166 1166 1166 1166 1166 116	766 962 1052 1178 1166 1167 1166 1167 1166 1177 1181 1100 1003 898 782 693 582 586 469 412 353 338 321 194 166 136 105 71 49 41 39 37 35 34	766 962 1052 1178 1166 1167 1166 1167 1186 1194 1096 1031 1186 1096 1031 1186 1096 1031 1186 1096 1031 1186 1096 1031 1096 1031 1096 1031 1096 1031 1096 1031 1096 1031 1096 1031 1096 1031 1096 1031 1031 1031 1031 1031 1031 1031 103	766 962 1052 1178 1166 1167 1166 1167 1166 1167 1186 1194 1201 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 1214 1158 1132 1107 11064 1102 978 933 885 796 124 224 205 186 146 124 205 186 146 129 51	766 962 1052 1178 1166 1167 1166 1167 1166 1177 1166 1177 1186 1194 1207 1211 1214 1218 1222 1174 1153 1134 1012 957 900 841 784 784 314 286 595 462 201 174

<u>C5</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$  = 0.08 m<sup>1/2</sup>

Т	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
<b>Q</b>	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0	
Time h	T	e m	p	е	r a	t	u	r e	3
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.40 2.40 2.40 2.40 2.40 2.40 2.40 3.60 3.80 4.00 4.00 4.60 4.80 5.00 5.60 5.60 5.80 6.00 5.80 6.00 5.80 6.00	788 1215 668 512 352 186 186 177 168 160 142 132 104 61 32 28 27 26 25 24 23	788 1215 1196 1079 889 675 593 504 416 325 293 275 219 181 141 97 46 42 39 37 35 33 32 31 30	788 1215 1196 1196 1086 1086 963 832 687 410 337 307 3239 205 171 135 40 38 33 31 29 27	788 1215 1196 1196 1197 1196 1197 1073 901 807 705 640 569 422 346 329 244 219 142 386 333 32 30 298 288	788 987 1079 1210 1196 1197 1197 1197 1206 1209 1123 1021 910 790 639 582 525 467 408 334 318 303 275 218 191 163 132 100 66 42 38 36 34 31 30	788 987 1079 1196 1197 1197 1197 1206 973 1219 121	788 987 1079 1210 1196 1197 1197 1296 1297 1299 1213 1219 1225 1229 1231 1233 1174 1120 1076 1032 987 940 891 799 707 609 509 407 344 321 300 280 280 280 280 280 280 280 280 280 2	788 987 1079 1210 1196 1197 1197 1197 1206 1209 1219 1225 1231 1233 1235 1238 1240 1189 1149 1149 1148 1076 1022 965 907 846 724 169 1788 338 388 281 253 196 169 138	

<u>06</u>

Time Graphs of Temperature of Combustion Gases. Type C Enclosed Space. Opening Factor  $A \cdot \sqrt{H}/A_{\rm t} = 0.12~{\rm m}^{1/2}$ 

		U							
T		0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<u>q</u>		18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
Time		T	e	m p	e	r	a t	u	r
h									
0.05		812	812	812	812	812	812	812	812
0.10		1250	1250	1250	1250	1014	1014	1014	1014
0.15		688	1230	1230	1230	1109	1109	1109	1109
0.20		525	1109	1230	1230	1245	1245	1245	1245
0.25		359	912	1230	1230	1230	1230	1230	1230
0.30		186	691	1115	1231	1230	1230	1230	1230
0.35		186	606	987	1230	1230	1230	1230	1230
0.40		177	513	851	1229	1229	1229	1229	1229
0.45		169	421	701	1229	1230	1230	1230	1230
0.50		160	326	632	1100	1231	1231	1231	1231
0.60		141	293	485	918	1237	1237	1237	1237
0.65		132	275	411	820	1239	1239	1239	1239
0.70		122	256	335	713	1148	1242	1242	1242
0.80	-	103	217	303	645	1039	1247	1247	1247
0.90		82	179	270	571	922	1136	1250	1250
1.00		58	138	236	496	796	1063	1253	1253
1.10		28	93	202	420	700	986	1254	1254
1.20		27	38	167	342	641	905	1256	1256
1.30		26	35	130	316	582	820	1192	1257
1.40		25	33	90	290	523	728	1164	1258
1.50			31	40	266	464	683	1136	1259
1.60			30	38	241	403	635	1089 1043	1260 1206
1.70			29	35	215	342	587	996	1185
1.80			. 28	33	190	329	539		1163
1.90			. 27	32	164 137	313	491 442	947 8 <b>9</b> 7	1142
2.00			26	31	78	299	343	802	1087
2.20 2.40					38	271 243	314	707	1031
2.60					35	215	286	607	972
2.80					32	188	258	504	912
3.00					30	159	230	400	849
3.20			٠.		_0	128	202	337	789
3.40						96	174	315	724
3.60						60	145	295	658
3.80						34	113	276	590
4.00	1					32	80	256	521
4.20						<b>5</b> _	37	237	452
4.40		1					31 <sub>4</sub>	216	381
4.60							32	198	332
4.80							31. 、	179	303
5.00							29 `	159	275
5.20							=	137	247
5.40								115	219
5.60								93	191
5.80								68	163
6.00								36	132

<u>D1</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{\rm H}/\rm A_t$  = 0.01 m  $^{1/2}$ 

	· ·							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<b>q</b>	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	T	e m	р	e	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	315 453 261 171 108 100 108 85 76 45 41 41 41 41 41 41 41 41 41 41 41 41 41	315 4456 4161 1169 1194 1197 1196 1196 1196 1197 1197 1197 1197	3153576881339372218884933706830750741864310988765431098875074484444444444444444444444444444444	315357660061685386266986542099614472086653109865442111998888777768665310986555	315 376 411 456 460 472 484 494 509 525 539 481 449 509 525 539 481 368 330 288 221 210 195 1135 1199 975 888 80 78 76	3341560 4784 490 5555555555555555555555555555555555	3156 4116 4560 4560 4560 4560 4560 4560 4560 456	315 316 411 456 472 484 490 509 525 539 556 5791 503 503 503 503 503 503 503 503 503 503

<u>D2</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{H/A_t}$  = 0.02 m<sup>1/2</sup>

<u>D3</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.04 m  $^{1/2}$ 

-	U								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time	T	e m	р	e	r a	t	u	r e	3
h								•	
Time h  0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40		e m  547 797 802 752 638 510 416 358 297 210 181 150 65 241 210 180 77 69 65 51 486 43 42 40 39 37 36 334 34	547 797 802 825 825 851 729 643 552 512 203 210 238 210 2154 210 238 210 247 552 552 552 553 266 238 210 253 266 275 275 275 275 275 275 275 275 275 275	9 547 797 8851 889 887 890 851 860 850 850 850 850 850 850 850 850 850 85	r a 547 660 734 826 837 862 881 898 915 935 966 914 688 622 589 452 406 344 813 286 234 209 185 913 105 894 80 76	547 660 7346 837 862 881 898 915 930 9556 9790 930 887 7251 484 446 337 2557 557 521 444 368 27 308 27 27 29 20 41 29 29 29 29 29 29 29 29 29 29 29 29 29	1 547 660 734 837 862 888 891 901 901 901 901 901 901 901 901 901 9	7 660 734 826 837 862 881 892 973 974 906 1019 1031 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1042 1050 1050 1050 1050 1050 1050 1050 105	2
4.60 4.80		33 32	39 38	53 51	73 70	88 84	231 212	378 348	
5.00		32	37	50	68	81	194	320	
5.20 5.40			37 36	48 47	65 63	77 75	176 157	293 267	
5,60			35 35	47 46	63 61	72	138	261 240	
5.80			34	44	59	70	118	214	
6.00			3 <sup>1</sup> 4	43	57	68	98	188	

<u>D4</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A\* $\sqrt{H/A_t}$  = 0.06 m<sup>1/2</sup>

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	Т	e m	p	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.60 5.80 6.00	613 903 525 429 312 181 170 162 138 129 121 104 88 69 46 44 42 40 38 37 35 34 33 32	613 903 911 853 727 573 519 453 384 268 185 218 150 218 555 559 45 31 32 32 32 32 32 32	613 913 935 936 936 936 936 936 936 936 936 936 936	613 903 914 936 937 958 900 1010 948 832 955 958 966 967 966 967 966 967 967 967 967 967	613 746 831 937 949 1009 1024 1053 1062 1053 1053 1053 1053 1053 1053 1053 1053	613 746 831 937 999 10024 10062 10063 1006	614 746 831 937 949 971 990 1009 1024 1036 1053 1060 1040 1040 1040 1040 1040 1040 1040	614 746 831 937 949 971 990 1009 1024 1036 1053 1062 1070 1136 1115 1130 1136 1147 1108 1092 1076 1076 921 870 817 765 711 870 817 765 536 475 412 365 335 471

<u>D5</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$  = 0.08 m<sup>1/2</sup>

<u>D6</u>

Time Graphs of Temperature of Combustion Gases. Type D Enclosed Space. Opening Factor A+ $\sqrt{\rm H}/\rm A_t$  = 0.12 m  $^{1/2}$ 

T		-							
Time h  0.05     697 697 697 697 697 697 697 697 697	T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
0.05	q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0
0.10		T	e m	р	e	r a	t	u	r e
5.60 36 41 106 205	0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40	697 1056 605 483 343 189 187 178 160 142 133 105 86 65 37 35 34 32 31 30 29 29 28	697 1056 1070 991 834 646 575 494 412 327 294 278 221 184 146 103 56 52 49 44 40 39 37 35 33 30	697 1056 1070 1094 1113 1027 921 805 671 478 410 340 340 308 274 241 207 174 139 101 60 58 54 51 49 47 43 40 38 36	697 1056 1070 11128 11156 11156 11156 11156 11156 11156 11156 11156 11156 11	697 874 965 1093 1104 1125 1135 1154 1159 1171 1176 687 632 776 632 1460 352 1460 352 195 168 138 139 147 145 145 145 145 145 145 145 146 147 147 147 147 147 147 147 147 147 147	6974534 96534 1121576 11	697 874 965 1093 1104 1121 1135 1148 1159 1171 1190 1197 1203 1208 1212 1155 1130 1061 1018 974 929 881 793 703 609 513 413 352 328 306 246 226 227 188 169	697 874 965 1093 1104 1121 1135 1148 1159 1171 1176 1180 1190 1203 1218 1215 1218 1221 1215 1218 1221 1215 1218 1221 1215 1218 1221 1223 1174 1134 1114 1134 1114 1134 1114 1136 1116 897 838 781 721 658 781 721 658 781 721 721 721 721 722 723 724 721 722 723 724 724 725 726 727 727 728 729 729 729 729 729 729 729 729 729 729
6.00 34 39 58 149								•	-

<u>E1</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A  $\sqrt{H}/A_{\rm t}$  = 0.01 m  $^{1/2}$ 

Ţ	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
ď.					-	t	u	r e
Time h	T	e m	g	е	r a	U	u	1 6
0.05	312	312	312	312	312	312	312	312 392
0.10	472	472 465	472 465	472 465	392 428	392 428	392 428	428
0.15	280 242	465	465 482	482	426 486	486	486	486
0.20 0.25	189	384	494	494	488	488	488	488
0.30	130	321	467	506	502	502	502	502
0.35	127	302	435	517	516	516	516	516
0.40	122	275	397	528	529 534	529 534	529 534	529 534
0.45	117 112	242 210	363 330	538. 516	547	547	547	547
0.50 0.60	103	197	288	466	572	572	572	572
0.65	99	197 188	262	435	580	580	580	580
0.70	95	180	236	400	562	588	588	588
0.80	86	163	220	377	534	604	604 620	604 620
0.90	77 68	146 128	204 188	351 323	499 458	584 567	632	632
1.00 1.10	56	108	172	293	425	543	643	643
1.20	52	89	156	263	402	519	656	656
1.30	49	84	139	249	380	491	647	665
1.40	46	80	121	236	358	459	644	675 686
1.50	44 42	76	102 97	223 210	334 308	440 421	640 628	694
1.60 1.70	42 40	73 . 70	91 92	197	283	403	616	686
1.80	39	68	88	184	273	383	600	684
1.90	38	64	85	171	263	363	583	681
2.00	37	6 <u>1</u>	82	158	254	342	566	676
2.20	35	56	76	130	237 220	298 277	528 491	662 644
2.40	3 <sup>1</sup> 4 32	52 49	72 66	110 102	204	259	450	622
2.60 2.80	31.	49 47	61	95	188	242	406	599
3.00	30	45	58	90	172	225	359	572
3.20		43	55	85	155	208	322	546
3.40		41	52	80	138	192	303 287	518 489
3.60		40 39	50 48	75 71	120 108	175 158	272 ·	
3.80 4.00		39 38	47	68	102	140	258	427
4.20		37	45	65	96	121	244	394
4.40		36	կկ	63	92	114	230	359
4.60		35	43	60	87	108	217	330
4.80		34	42	58 56	82 79	103 98	204 191	309 290
5.00 5.20		33	40 40	55	76	90 94	178	272
5.40			39	53	73	89	165	254
5.60			38	52	71	85	152	236
5.80			37	50	68	82	138	217
6.00			36	49	66	79	124	199

<u>E2</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor  $A \cdot \sqrt{E}/A_t = 0.02 \text{ m}^{1/2}$ .

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
<u>g</u>	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0	
Time h	T	e m	p	e	r a	t	u	r e	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	441 637 378 328 249 163 153 146 139 126 119 126 54 50 44 40 38 37 35 33 31 30	441 637 642 669 524 435 460 327 320 275 207 142 207 142 207 207 207 207 207 207 207 207 207 20	4437 4437 4437 4457 4457 4457 4457 4457	441 637 642 668 77 738 668 77 738 67 738 759 759 759 759 759 759 759 759 759 759	442 536 579 674 674 773 760 675 773 608 775 608 451 552 415 805 552 415 810 810 810 810 810 810 810 810 810 810	44368143245552144138930853113728839520	442 536 671 673 671 763 815 7748 815 7768 815 815 815 815 815 815 815 815 815 81	443 536 598 671 674 693 7131 748 763 815 835 853 861 910 919 926 882 910 919 929 901 882 751 497 400 363 338 497 449 400 363 338 284 261 238 215 219	

<u>E3</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$  = 0.04 m  $^{1/2}$ 

	U							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
đ	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0
Time h	T	e m	р	e	r a	t	u	r e
h 0.05 0.10 0.15 0.20 0.25 0.30 0.45 0.50 0.665 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 3.20 3.40 3.60 3.80 4.20 4.40 4.60 4.80 5.00	547 844 478 406 183 1656 183 1656 1123 107 5518 443 183 183 183 183 183 183 183 183 183 18	547 844 831 783 681 555 506 445 282 265 218 82 773 65 65 441 39 37 35 34 33 31 31	547 844 831 8578 870 595 595 449 333 224 333 221 345 346 347 347 347 347 347 347 347 347 347 347	544 831 8578 867 876 877 887 876 877 887 876 877 887 876 877 887 877 87	547 689 759 865 869 939 956 971 997 1006 955 885 711 637 592 547 503 453 322 203 254 203 254 203 254 203 254 265 275 286 275 286 275 286 275 286 275 286 275 286 275 286 286 286 286 286 286 286 286 286 286	547 689 759 865 869 919 956 971 965 917 864 806 741 671 671 871 871 871 871 871 871 871 871 871 8	547 689 759 759 865 895 919 936 1011 1056 1030 1051 1056 1063 1056 1063 1063 1063 1063 1063 1063 1063 106	547 689 759 865 869 919 939 956 1013 1027 1038 1047 1056 1063 1070 1076 1081 1086 1055 1041 1026 1010 971 929 883 835 784 734 682 574 460 354 355 298
5.20			JE	39	52	63	162	272
5.40 5.60	 •			38 37	50 48	60 57	143 123	246 220
5.80 6.00				36 35	46 45	55 52	103 83	195 169

<u>E4</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A+ $\sqrt{\pi}/\Lambda_{\rm t}$  = 0.06 m<sup>1/2</sup>

	L L							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
<u>q</u>	9.0	13.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	ā	е	r a	ı t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00 6.60	615 929 532 445 323 189 170 161 145 126 108 91 149 45 43 40 33 33 33 30	615 929 937 882 759 6546 478 406 329 274 257 2288 153 114 73 606 553 508 441 386 342	615 927 937 9358 9353 9353 9353 9353 9357 6490 9494 9494 9494 9494 9494 9494 9494	615 9237 9378 9013 10503	615 766 854 959 1002 1058 1070 1085 1091 1038 1091 1038 1091 1038 1091 1038 1091 1038 1091 1038 1091 1038 1091 1038 1091 1038 1091 1038 1091 1091 1091 1091 1091 1091 1091 109	615 615 616 617 616 617 617 617 617 617 617 617	615 766 854 976 1024 1058 1070 1085 1091 1107 1112 1129 1134 1092 1070 1017 1011 973 893 850 684 1070 1070 1070 1070 1070 1070 1070 107	615 766 854 960 976 1002 1058 1070 1085 1097 1105 11097 11139 1139 1147 1151 11098 1015 1015 1015 864 805 807 808 809 809 809 809 809 809 809 809 809

<u>E5</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A· $\sqrt{H}/A_{\rm t}$  = 0.08 m<sup>1/2</sup>

${f T}$	0.1	0,2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Tîme h	Т	e m	р	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	663 999 577 468 335 191 181 172 163 145 108 90 44 41 39 37 36 33 31 29 28	663 999 1005 942 805 638 568 494 417 259 223 187 150 109 65 61 56 53 49 41 37 34 32 31 30	663 999 1005 1025 1055 988 897 792 669 608 480 414 210 179 145 108 70 66 62 58 54 43 33 33 33 33 33 33 33 33 33 33 33 33	663 995 1026 105776 1057	663 831 913 1027 1047 1086 1101 1113 1138 1068 571 623 571 408 352 109 767 677 623 140 109 767 140 109 77 109 140 109 140 109 140 109 109 109 109 109 109 109 109 109 10	663 8313 1043 1046 1086 1113 1123 1143 1143 1143 1143 1144 1	663 831 913 1043 1046 11086 1113 1122 1133 1138 11451 1169 1173 11076 1173 11076 1173 11076 1173 11076 1173 11076 1173 11076 1173 11076 1177 11076 1177 11076 1177 11076 1177 1177	663 831 913 1027 1043 1067 1086 1101 1113 1122 1133 1158 1164 1169 1173 1177 1180 1183 1186 1183 1186 1183 1186 1183 1186 1183 1186 1183 1186 1183 1186 1183 1186 1183 1186 1183 1186 1183 1185 1186 1183 1186 1186

<u>E6</u>

Time Graphs of Temperature of Combustion Gases. Type E Enclosed Space. Opening Factor A  $\sqrt{\rm H}/\rm A_t$  = 0.12 m<sup>1/2</sup>

	τ							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	18.0	35.0	5 <sup>4</sup> .0	90.0	135.0	180.0	270.0	360.0
Time h	T	e r	n p	е	r a	t	u	r e
0.05 0.15 0.20 0.25 0.35 0.45 0.65	709 1086 610 494 348 194 192 183 164 145 166 87 668 334 32 31 29 28 27 25	709 1086 1087 1013 860 670 599 428 297 278 260 222 184 550 43 41 39 37 334 34	709 1086 1087 1108 1130 9447 6927 488 417 3074 2417 138 54 54 42 35 33 33 33 33 33 33 33 33 33 33 33 33	709 1086 1087 1108 1130 1148 11619 11675 1069 903 811 709 637 5494 420 245 220 191 145 440 37 333 333 31	709 899 985 1122 1141 1156 1178 1187 1187 1187 1187 1187 1187 118	709 899 985 1124 1156 1178 1187 1190 1103 1103 1103 1103 1103 1103 1103	709 899 985 1106 1122 1141 1155 1166 1178 1190 1205 1205 1201 1215 1160 1215 1160 1215 1160 1215 1160 1215 1215 1215 1215 1216 1216 1217 1217 1217 1217 1217 1217	709 899 985 1106 1122 1141 1155 1166 1178 1190 1205 1209 1215 1218 1220 1221 1218 1222 1224 1175 1153 1162 1113 1062 1007 1007 1007 1009 1009 1009 1009 1009

<u>F1</u>

Time Graphs of Temperatu	e of	Combustion	Gases.	Type F	'Enclosed	Space.	$\epsilon_{ m res}$	= 0.10
Opening Factor A·VH/A_ =	0.01	m <sup>1/2</sup>						

-		τ											
Ť		0.	.1		0.3		0.	5		1.0		2.0	
q		1.	.5		4.5		7.	5	:	15.0		30.0	
Time		T	е	m	p	е	r	а	t	u	r	е	
h	•						,						
0.05		25 43			252 418		25 41	2		252 350		252 350	
0.10 0.15		-29			456		45	6		413		413	
0.20		26	1		489		48			480		480	
0.25		20 1 <sup>1</sup>	8 18		506 483		50 51			498 513		498 513	
0.35		12	9		452		52	8		524		524	
0.40		11	Ţţ		409		53 54			535 540		535 540	
0.45		Š.	13		359 329		51	6		544		544	
0.60		8	1		270		ħħ			553		553	
0.65 0.70			5 0		239 207		40 36		-	557 560		557 560	
0.80		5	9		180		33	2		566		566	
0.90		4	.9 .8		159 140		29 26	9 8		538 510		572 577	
1.00		2	8		122		23			480		581	
1.20					105		20			449		585	
1.30 1.40					88 70		18 17			414 377		589 592	
1.50					50		15	7		355		595	
1.60 1.70					42 37		14 13			334 314		598 585	
1.80					33		11	8		293		576	
1.90 2.00					31 29		10 9	5		273 251		568 560	
2.20					29		6	5		207		538	
2.40							4			186		515	
2.60 2.80							3 3	( 3		170 155		490 463	
3.00							3			140		436	
3.20 3.40				,						125 110		409 382	
3.60										96		354	
3.80										80 64		325 294	
4.00 4.20	.:									47		264	
4.40					·					40		231	
4.60 4.80										37 35		205 189	
5.00										33		174	
5.20 5.40												159 144	
5.60												129	
5.80												113	
6.00												98	

<u>F2</u>

Time Graphs of Temperature	of	Combustion	Gases.	Туре	F	Enclosed	Space.	Eres	=	0.10
Opening Factor $A \cdot \sqrt{H}/A_{\perp} = 0$ .	.02	<sub>m</sub> 1/2						165		

Т	0.1	0.3	0.5	1.0	2.0
<b>Q</b>	3.0	9.0	15.0	30.0	60.0
Time h	т е	m p	e r a	t u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.00 3.20 3.40 3.60 4.00 4.60 4.80 5.00 5.40 5.60 5.80 6.00	382 612 413 354 273 184 161 143 130 120 104 97 90 77 63 48 32	382 612 645 684 699 666 610 547 433 351 308 203 180 156 133 108 84 57 41 38 35 33	382 612 645 684 711 722 734 689 597 545 489 395 351 260 237 218 200 184 113 77 48 31 32	382 508 508 507 716 717 757 757 757 757 759 762 759 759 759 759 759 759 759 759 759 759	382 508 585 676 716 723 755 776 7780 788 790 778 788 790 772 7748 736 603 792 454 414 331 283 221 213 455 114 155 114

162

Time Graphs of Temperature of	f Combustion	Gases.	Type F	Enclosed	${\tt Space.}$	$\epsilon_{ ext{res}}$	= 0.10
Opening Factor $A \cdot \sqrt{H}/A_{\perp} = 0.0$	<sub>04 m</sub> 1/2						

Opening ractor a tar	~t											
T	0.1	L	C	3.0		0.5	5	]	L <b>.</b> 0		2.0	
<b>q</b>	6.0	)	18	3.0		30.0	)	60	0.0	]	L20.0	
Time	T	е	m	P	е	r	a	t	u	r	e	
h												
0.05	522			22		522			22		522	
0.10	815			15 41		815 841			73 63		674 763	
0.15 0.20	525 437			70		870			74		874	
0.25	327		8	92		892		8	84		884	
0.30	206			41		901		8	96		897	
0.35	181 163			67 79		909 916		9	05 13		905 913	
0.40 0.45	150			74		921			19		919	
0.50	139		5	21		855		9	24		924	
0.60	121			<u>1</u> 4		731		9	36		936	
0.65	112			59		661 585		9	39 41		939 941	
0.70 0.80	105 89		20	02 64		525			դկ		944	
0.90	72			34		469		8	77		948	
1.00	54		2	05		413		8	24		951	
1.10	33			78 No		356 296		7	71 13		953 955	
1.20				49 19		272		6	51 51		957	
1.40				88		250		5	83		959	
1.50				54		229		5	48		960	
1.60				<u>կ</u> կ 39		208 188		5. 10	14 79		962 931	
1.70 1.80			:	39 36		168			19 43		917	
1.90				34		147		4	07		903	
2.00				32		124			70		888	
2.20 2.40						79 44		25	94 68		851 811	
2.60						37		2	կկ		765	
2.80						34			21		720	
3.00						32			98		672	
3.20								1.	76 52		627 581	
3.40 3.60								1:	28		533	
3.80								10	02		483	
4.00									76		432	
4.20									46 38		379	
4.40 4.60									30 35		325 284	
4.80									33		261	
5.00									32		238	
5.20											214	
5.40 5.60											192 169	
5.80											146	
6.00											120	

<u>F4</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.10 Opening Factor A· $\sqrt{\rm H}/\rm A_t$  = 0.08 m<sup>1/2</sup>

		. t. . n. t											
T		0.3	÷.	(	3.3		0.9	5	:	1.0		2.0	
đ		12.0	)	36	0.0		60.0	)	120	0.0	2	240.0	
Time h		T e		m	Þ	e	r	a	t	u	r	е	
0.05 0.10 0.15 0.25 0.35 0.45 0.65 0.80 0.45 0.90 1.20 1.30 1.40 1.70 1.80		628 971 620 493 355 208 188 172 1131 96 77 55 30			32 43 72 76 71 76 71 76 76 76 76 76 76 76 76 76 76 76 76 76		628 971 1012 1043 1058 1058 1058 1058 1058 1058 1058 1058		8	37756658277795635110897641472616135672		628 822 1025 1025 1037 1047 1066 1077 1081 1075 1083 1077 1088 1099 1099 1099 1099 1099 1099 1099	

<u>F5</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\varepsilon_{\rm res}$  = 0.10 Opening Factor A·VH/A<sub>t</sub> = 0.12 m<sup>1/2</sup>

			0.1 0.3										
T			0.1					0.5	5	1	0		2.0
đ		18.0 54.0 90.0 180 Temperat					.0		360.0				
Time h			T	е	•		e	r	a	t	u	r	е
0.05 0.10 0.20 0.25 0.30 0.45 0.60 0.65 0.70 0.80 0.10 1.30 1.40 1.50 1.70 1.80 2.20 2.40 2.30 3.40 3.40 3.40 4.40 4.60 5.60 5.60 5.60 6.60			677 1060 514 364 205 189 1754 125 1155 27		100 110 111 100 99 88 66 66 44 22 22 22 11 11 12 13 14 15 15 16 16 16 16 16 16 16 16 16 16 16 16 16	86 06 14 30 32 4 57 34 57 57 57		677 1060 1086 1114 1121 1125 1128 1120 1021 672 603 534 465 393 321 225 202 272 249 225 202 179 128 73 31		67 88 109 111 112 113 113 113 113 113 113 113 113	001707379146806130151851720371403554		677 981 1117 1123 1127 1123 1129 1134 1134 1134 1134 1134 1134 1134 113

<u>F6</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$  = 0.01 m<sup>1/2</sup>

T	0.1		0	•3		0.5	;	1	0		2.0
<b>q</b>	1.5		4	•5		7.5		15	0.8		30.0
Time h	Т е		m	р	е	r	a	t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	241 387 283 243 193 136 102 93 544 434		7 6 1 <sub>4</sub> 3	37 19 17 55 10 11 16 13 13 16 16 18 16 18 16 18 16 18 16 18 18 18 18 18 18 18 18 18 18 18 18 18		241 387 419 447 455 466 470 471 451 451 451 451 466 470 471 451 466 470 471 451 466 470 471 451 466 470 471 471 471 471 471 471 471 471 471 471		32341 4944 4944 4944 4944 4944 4944 4944 4	51 58 58 59 59 59 59 59 59 59 59 59 59 59 59 59		2472 2472 2473

F7

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$  = 0.02 m<sup>1/2</sup>

2.0	.0	1	i	0.5		.3	(		0.1		T	7
60.0	0.0	30	)	15.0		.0	9	3.0			q	q
r e	u	t	a.	r	e	p	m	т е			Time	1
											h	Ì.
364 472 537 614 629 639 647 652 664 665 665 664 665 665 665 665	+1	49 55 66 66 66 66 66 66 66 66 66 66 66 66		364 560 588 614 615 625 630 591 385 331 322 347 321 423 331 321 423 331 321 423 331 321 423 331 321 423 331 331 331 331 331 331 331 331 331 3		36 97 95 55 13	55 66 66 55 55 44 43 33 22 20 20 11 11 11 11 11 11 11 11 11 11 11 11 11		364 560 394 332 257 147 119 90 84 75 84 29	3 2 2 3 3 3 3	0.05 0.10 0.15 0.25 0.25 0.30 0.45 0.45 0.50 0.65 0.70 0.65 0.70 0.10 1.10 1.20 1.30 1.40 1.40 1.50 1.40 1.50 1.40 1.50	00000000000000111111111122222333333344444555555

Time Graphs of	Temperature of	f Combustion	Gases.	Type F	Enclosed	Space.	E	= 0.35
Opening Factor	$A \cdot \sqrt{H}/A_{+} = 0.0$	)4 m <sup>1</sup> /2					162	

		•													
T		0.1		į	0.3	•	0.5		J	.0		2.0			
q		6.0		ı	0.8		30.0	ı	60	0.0		120.0			
Time		T	e	m	р	е	r	æ	t	u	r	e			
h		-	•		P	Ů	-	-	•	-	-	-			
0.05 0.10 0.15 0.25 0.35 0.40 0.50 0.65 0.65 0.665 0.665 0.70 1.20 1.30 1.50 1.70 1.80 1.70 1.80 1.70 1.80 1.70 1.80 1.70 1.80 1.70 1.80 1.90 2.20 2.40 2.80 3.80 3.80 4.40 4.60 4.60 5.50 5.60 5.60 5.60 5.60 5.60 6.60		473 751 407 305 191 168 152 141 132 1157 100 85 69 50 30		7 7 7 7 7 7 6 6 5 5 4 4 3 3 2 2 2 2 1 1 1 1	73137087566986123889420		473 751 763 779 783 790 796 793 737 634 426 471 426 379 328 275 421 178 159 117 73 39 34 31		626 627 77 77 77 77 77 77 77 77 77	888 95 97 99 99 90 90 90 90 90 90 90 90 90 90 90		473 628 6784 7780 805 7780 805 809 811 813 816 817 777 777 7696 816 816 817 777 777 777 777 777 777 777 777 777			

<u>F9</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.35 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$  = 0.08 m<sup>1/2</sup>

	-										
Ţ	0.1		0.3			0.5	;	3	0		2.0
q.	12.0		36.0			60.0	•	120	.0	2	240.0
Time h	Т	е	m	р	e	r	a	t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	593 892 593 464 338 1978 1128 119 110 94 75 28		. 3	92 21 23 37 98 39 88 86 22 23 23 22		593 892 921 923 934 937 940 943 860 738 868 590 258 236 214 121 148 124 72 33		778499999999999999999999999999999999999	336 9 244 7 8 9 4 4 6 5 7 5 7 7 4 00 4 8 00 22 7 7 3 9 5 1 7 00 4 8 00 22 7 7 3 9 5 1 7 00		577843 999999999999999999999999999999999999

<u>F10</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.35 Opening Factor  $\text{A} \cdot \sqrt{\text{H}}/\text{A}_{\rm t}$  = 0.12 m<sup>1/2</sup>

	•				0.0		
T	0.1	0.3	0.5	1.0	2.0		
q	18.0	54.0	90.0	180.0	360.0		
Time	T e	m p	e ra	t u	r e		
h							
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	640 985 631 489 350 197 182 170 160 150 123 114 96 77 53 26	640 985 1000 1003 1010 934 847 746 628 569 447 382 315 281 250 218 187 155 120 83 36 30	640 985 1000 1003 1010 1013 1016 1018 1020 934 800 721 633 572 509 445 379 311 288 266 243 220 198 175 151 125 71 30	640 852 921 1006 1005 1012 1015 1017 1021 1022 1023 1024 1025 946 893 837 776 706 631 595 558 519 480 439 397 311 286 261 236 211 186 160 132 103 73 31	640 852 921 1006 1005 1012 1015 1017 1019 1021 1023 1024 1025 1027 1030 1030 1031 1030 1031 992 979 964 948 909 868 824 777 726 678 628 575 520 464 404 344 404 344 300 276 225 200 175 149		

F11

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{res} = 0.60$ 

Opening Factor  $A \cdot \sqrt{H}/A_t = 0.01 \text{ m}^{1/2}$ 0.5 1.0 2.0 0.1 4.5 7.5 15.0 30.0 1.5 q Ţ r a u r Time e m p h 0.05 0.10 0.15 0.20 408 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 

1.80

1.90

2.00

2.20

2.40

2,60

2.80

3.00

3.20

3.40

3.60

3.80

4.00

4.20

4,40

4.60

4.80

5.00

5.20 5.40

5.60

5.80

6.00

<u>F12</u>

Time Graphs of Temperature	of Combustion	Gases.	Type F Enclosed	Space.	$\varepsilon_{res} = 0.60$
Opening Factor $A \cdot \sqrt{H}/A_{\perp} = 0$	.02 m <sup>1/2</sup>				100

T	0.1		0.3			0.5	i	1	.0		2.0
q	3.0		9	0.0		15.0	)	30	0.0		60.0
Time h	Т	e	m	р	е	r	a.	t	u	r	e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	348 525 374 310 240 159 136 111 104 91 85 80 68 55 41 27		505 56 55 55 55 55 55 55 55 55 55 55 55 55	54 18 93 73 54 34		348 525 547 560 568 571 570 534 426 385 322 254 217 200 186 172 128 112 128 37 31		144 50 50 50 50 50 50 50 50 50 50 50 50 50	71 73 74 77 78 78 78 78 78 78 78 78 78 78 78 78		94724127755555555555555555555555555555555

<u>F13</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.60 Opening Factor A· $\sqrt{\rm H}/\rm A_t$  = 0.04 m<sup>1/2</sup>

opening ractor A vii/	^t	0.0-	7 102									
Т	0.1		C	0.3		0.5	5	1	L.0		2.0	
<b>q</b>	6.0	)	18	3.0		30.0	)	60	0.0	:	120.0	
Time h	T	е	m	Þ	е	r	8.	t	u	r	е	
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	450 708 470 381 288 180 158 144 125 110 97 82 66 48 28		70 70 70 70 70 666 655 553 444 43 30 20 20 18 11 11 11 11 11 11 11 11 11 11 11 11	33 08 35 61 35		450 708 709 710 7114 717 719 668 578 529 474 437 397 309 221 1223 206 189 171 1523 1312 69 36 31		595 695 777 777 777 777 777 777 777 777 777 7	11 16 18 18 18 18 18 18 18 18 18 18 18 18 18		450 597 711 716 728 731 732 733 733 733 733 733 733 733	

<u>F14</u>

Time Gra	aphs (	of	Temperature	of	Combustion	Gases.	Туре 1	F Enclosed	Space.	$\epsilon_{ ext{res}}$	=	0.60	
Opening	Fact	or	$A \cdot \sqrt{H} / A_{+} = 0.$	.08	m <sup>1/2</sup>								

m	^ 1		0	3	0.5			1.0			2.0
T	0.1		0.3 36.0			60.0		120			240.0
<b>q</b>	12.0		36	.0							
Time	T	е	m	$\mathbf{p}$	е	r	a	t	u	r	е
h											
0.05	568		56 84			550 884			68		568
0.10 0.15	840 560		85			845			+1 91		741 791
0.20	440		85			854			56		856
0.25	324		86			860			55		855
0.30	190		80			862			52 3.		862
0.35 0.40	172 160		73 65			865 867		86 86			864 866
0.45	150		55			869		86			868
0.50	142		50	9		787		87			870
0.60	125		40			683		87			872
0.65 0.70	116 108		35 29			621 552		87 87			873 874
0.80	92		26			510		87			875
0.90	73		23	15		460		81	16		880
1.00	51		20			406		77			877
1.10 1.20	27		17 14			350 290		72 67			878 880
1.30			11			270		62			879
1.40			8	2		250		56			880
1.50				9		229		53			885 880
1.60 1.70			7	1		208 188		50 47			852
1.80						167		43			843
1.90						145		40			831
2.00						121 70		36 29			819 787
2.40						31		26			753
2.60								24			717
2.80								22			679
3.00 3.20								20 17			638 600
3.40								15			560
3.60								12	7		517
3.80								10			472
4.00 4.20									1 4		424 374
4.40								2	-		320
4.60											282
4.80											260
5.00 5.20											237 214
5.40											191
5.60										-	168
5.80											143
6.00											116

<u>F15</u>

Time Graphs of Temperature of Combustion Gases. Type F Enclosed Space.  $\epsilon_{\rm res}$  = 0.60 Opening Factor A· $\sqrt{\rm H}/\rm A_{\rm t}$  = 0.12 m<sup>1/2</sup>

T	0.1	0.3	0.5	1.0	2.0
<b>q</b> .	18.0	54.0	90.0	180.0	360.0
Time h	Т е	m p	e r a	t u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.40 3.60 3.80 4.00 4.60 4.60 4.60 4.60 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	613 931 599 468 339 192 177 166 157 148 130 121 112 95 76 52	613 931 933 936 943 870 795 703 594 432 274 214 152 118 82 35 28	613 931 933 945 945 947 948 949 873 752 680 601 547 432 370 305 282 261 239 217 195 173 28	613 818 868 940 937 949 950 952 953 953 892 845 793 735 672 604 424 385 303 256 232 207 183 158 130	618 818 919 919 919 919 919 919 919 919 9

<u>G1</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A+ $\sqrt{H}/\Lambda_{\rm t}$  = 0.01 m<sup>1/2</sup>

	•							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	1.5	3.0	4.5	7.5	11.25	15.0	22.5	30.0
Time h	T	e m	ũ	е	r a	t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.20 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	287 473 290 253 199 135 128 102 97 88 79 55 52 49 46 44 41 38 36 31 30	287 478 478 456 394 305 217 187 180 180 180 180 180 180 180 180 180 180	287 478 494 500 467 3350 467 3350 423 210 1157 1204 88 88 88 87 77 86 46 41 40 83 33 33 33 33 33 33 33 33 33 33 33 33	2871849555155142015551201996306803839630864431098876839630864431098876839630864431098876839630864431098876830864431098876830860383963086443109887683086443109887683086443109887683086443109887683086443109887683086443109887	287 362 440 498 495 501 513 517 521 548 495 573 573 573 396 411 397 213 213 213 196 213 1145 74 66 63 60 74 74 66 67 75 75 75 76 76 76 76 76 76 76 76 76 76 76 76 76	287 340 498 495 517 517 517 517 517 517 517 517 517 51	287 3610 498 4995 5017 5137 5147 5137 5147 5131 620 644 674 6738 664 674 673 674 674 675 674 673 674 674 675 674 675 677 678 679 679 679 679 679 679 679 679 679 679	287 361 498 495 501 513 517 521 583 606 629 629 638 631 716 638 635 706 638 635 706 638 635 706 838 848 7580 838 848 7580 848 848 848 848 848 848 848 848 848 8

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.02 m  $^{1/2}$ 

opening ractor is	t							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q.	3.0	6.0	9.0	15.0	22.5	30.0	45.0	60.0
Time	${f T}$	e m	$\mathbf{p}$	e	r a	t	u	r e
h							•	
0.05	413 675	413 676	413 676	413 676	413 558	413 559	413 550	413 559
0.10 0.15	391	654	654	654	605	602	559 602	602
0.20	335	606	663	663	674	674	674	674
0.25	257	513	672 623	672 680	668	668	668	668 677
0.30	173 167	417 276	570	707	677 684	677 684	677 684	684
0.35 0.40	157	376 344	536	736	730	730	730	730
0.45	149	301	468	753	745	745	745	745
0.50	141 126	258 243	440 368	717 637	766 796	766 796	766 706	766 796
0.60 0.65	119	243	327	580	808	808	796 868	808
0.70	112	213	286	530	783	819	819	819
0.80	100	187 161	260 234	493 458	747 698	849 762	$\frac{849}{814}$	849 814
0.90 1.00	88 75	134	210	430 418	631	738	014 845	845
1.10	62	106	186	373	572	713	845 866	866
1.20	58	76	162	325	538	670	877	877
1.30 1.40	55 51	75 70	138 111	302 281	501 464	626 580	780 787	807 834
1.50	48	67	85	261	425	551	790	856
1.60	46	62	80	241	386	521	772	876
1.70	43 41	59 55	78 75	221 201	345 326	491 461	769 751	814 853
1.80 1.90	39	53	72	182	310	431	724	856
2.00	38	50	68	162	295	399	695	847
2.20	35	46 42	61. 55	120 92	270 245	334 303	635	821 790
2.40 2.60	33 31	42 40	51	83	222	277	575 512	756
2.80	29	37	47	78	199	253	443	719
3.00	28	35	hh	72	176	230	374	681 645
3.20 3.40	27 26	34 32	41. 39	67 61	153 128	207 185	322 297	607
3.60	26	31	37	57	103	163	277	569
3.80		30	36	53	85	139	258	529
4.00		29 28	34 33	50 47	80 73	114 88	241 224	488 444
կ.20 4.40		28	32	45	68	81	208	400
4.60		27	31	43	63	73	193	364
4.80		27	30 30	<u>4</u> 1 40	59 56	67 62	177	342 321
5.00		26	29	38	56 53	59	160 143	301
5.20 5.40			28	37	50	55	126	281
5.60			28	36	48	52	107	261
5.80			27 27	35 34	46 44	50 48	90 73	241 220
6.00			4 ا	<b>⊃</b> +	44	40	13	220

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.04 m<sup>1/2</sup>

	Ü								
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
q	6.0	12.0	18.0	30.0	45.0	60.0	90.0	120.0	
Time h	Т	e	m p	е	r	a t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.60 3.60 4.00 4.20 4.140 4.60 4.80 5.00 5.20 5.40 5.60 5.60 5.60 6.00	545 8481 403 1849 170 161 152 103 71 103 71 103 103 103 103 103 103 103 103 103 10	542 8439 765496 33774 25418 1116 6655518 453 3642 3765498 437748 43779 40779 40779 40779 40779 40779 40779 40779 40779 40779 40779 4	546 842 839 852 876 838 7693 5944 387 268 208 208 208 2148 418 419 419 419 419 419 419 419 419 419 419	542 842 851 857 962 733 865 966 760 760 760 760 760 760 760 760 760 7	545 695 763 855 917 855 917 917 917 914 915 915 917 914 915 915 915 915 915 915 915 915 915 915	546 595 596 597 597 598 597 598 598 598 598 598 598 598 598 598 598	546 6953 8551 938 9959 9959 901 901 902 903 903 903 903 903 903 903 903 903 903	546 695 855 915 939 939 939 931 1009 1011 985 950 968 971 968 968 971 968 971 968 971 968 971 968 971 971 971 971 971 971 971 971 971 971	

<u>G4</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.06 m $^{1/2}$ 

	•							
T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q.	9.0	18.0	27.0	45.0	67.5	90.0	135.0	180.0
Time h	T	e m	p	е	r	a t	u	r
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.66 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.20 2.40 2.60 2.80 3.00 3.20 3.40 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00	593 947 532 4396 184 175 156 138 121 106 86 86 40 37 53 33 31 32 28	593 947 943 863 759 402 3286 249 3104 45 45 42 331 331 298 27 287	59473 59473 5943 5983 5975 5876 5876 5876 5876 5876 5876 5876 58	593 947 947 953 983 10051 998 1051 998 1051 998 1051 998 1051 1051 1051 1051 1051 1051 1051 105	593 786 960 9738 9738 1055 1005 1005 1005 1005 1005 1005 100	593 780 960 973 998 1023 1044 1059 1001 1005 1005 1005 1005 1005 1005	593 780 860 960 973 998 1059 1008 1057 1001 1047 1074 1076	593 780 860 960 973 998 1059 1008 1057 1001 1002 1045 1047 1074 1076

<u>G5</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor  $\text{A.}\sqrt{\text{H}}/\text{A}_{\text{t}}$  = 0.08 m<sup>1/2</sup>

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0
q	12.0	24.0	36.0	60.0	90.0	120.0	180.0	240.0
Time h	T	e m	q ı	е	r a	ı t	u	r e
0.05 0.10 0.15 0.20 0.25 0.30 0.40 0.45 0.50 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.60 1.70 1.80 1.90 2.00 2.40 2.60 2.80 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.40 3.60 3.80 4.90 4.40 4.60 4.80 5.00 5.60	643 1015 574 460 327 188 186 176 167 157 139 130 121 103 85 66 41 38 36 34 32	640 1015 1009 939 802 636 491 255 180 271 255 180 142 101 55 48 43 43 33 33 33 33 33 34 36 34 36 37 37 38 38 38 38 38 38 38 38 38 38 38 38 38	640 1015 1009 1024 1049 985 977 673 610 484 418 350 241 207 173 139 60 55 43 33 31 29	640 1015 1009 1024 1049 1074 1074 1074 1074 1074 1074 1074 1074	641 9446 91246	640 842 918 1038 1058 1067 1069 1067 1069 1067 1069 1069 1069 1069 1069 1069 1069 1069	640 842 918 1027 1038 1062 1092 1058 1067 1069 1069 1117	640 842 918 1027 1038 1062 1090 1046 1067 1092 1058 1067 1069 1097 1117

<u>G6</u>

Time Graphs of Temperature of Combustion Gases. Type G Enclosed Space. Opening Factor A. $\sqrt{H}/A_{\rm t}$  = 0.12 m  $^{1/2}$ 

T	0.1	0.2	0.3	0.5	0.75	1.0	1.5	2.0	
q	18.0	35.0	54.0	90.0	135.0	180.0	270.0	360.0	
Time h	T	e m	מַ	е	r a	, t	u	r	е
0.05 0.10 0.15 0.20 0.25 0.30 0.35 0.40 0.45 0.50 0.60 0.65 0.70 0.80 0.90 1.00 1.10 1.20 1.30 1.40 1.50 1.80 1.90 2.20 2.40 2.60 2.20 2.40 2.60 2.80 3.60 3.80 4.00 4.40 4.60 4.80 5.00 5.20 5.40 5.60 5.80 6.00 5.20 5.40 5.60 5.80 6.00 5.80 6.00	726 1099 615 486 340 188 187 177 168 159 140 131 103 84 63 35 33 31 29 28 27 26	723 1099 1091 1009 853 669 593 508 422 332 274 255 216 140 97 41 40 337 35 33 29 27 25	723 1099 1091 1104 1125 1048 900 670 610 477 407 336 301 267 233 200 167 130 92 46 44 40 37 35 33 30 28 26	723 1099 1091 1104 1118 1134 1153 1052 856 772 674 405 333 308 284 260 236 211 187 163 136 79 39 34 31 28	726 907 988 1107 11136 11131 1158 11137 1085	723 907 988 1104 1117 1136 1157 1120 1141 1158 1131 1147 1163 1148 1071	723 907 988 1104 1117 1136 1157 1120 1141 1158 1131 1147 1163 1148 1169	723 907 988 1104 1117 1136 1157 1158 1131 1147 1163 1148 1169	

