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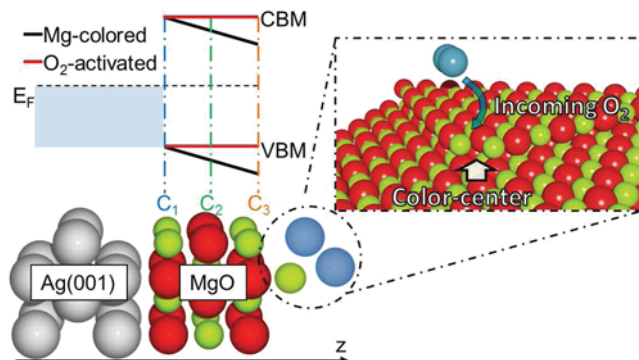
## Band Bending in Mg-Colored and O<sub>2</sub>-Activated Ultrathin MgO(001) Films

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**ABSTRACT:** Ultrathin MgO films grown on Ag(001) have been investigated using X-ray and ultraviolet photoemission spectroscopies for oxide films successively exposed to Mg and O<sub>2</sub> flux. Studying work functions and layer-resolved Auger shifts allows us to keep track of band profiles from the oxide surface to the interface and reveal the charge-transfer mechanisms underlying the controlled creation of Mg-induced surface color centers and the catalytic enhancement of O<sub>2</sub> activation. Our results demonstrate that one can intimately probe the catalytic properties of metal-supported ultrathin oxide films by studying the electronic band alignment at interfaces.



### INTRODUCTION

Metal-supported ultrathin oxide films have been widely studied both experimentally and theoretically in the field of heterogeneous catalysis due to their pivotal role in controlling charging mechanisms, adsorption properties, and catalytic activation of metal adatoms and molecules.<sup>1–6</sup> The favored electron tunneling through the oxide film brought by the ultrathin limit opens new catalytic pathways for charge-transfer mechanisms, as demonstrated by the charging of Au atoms and the spontaneous activation of carbon monoxide and molecular oxygen upon adsorption on ultrathin MgO(001) films.<sup>4,5,7–9</sup>

The catalytic properties of metal-supported ultrathin oxide films are also strongly impacted by the presence of defects. For example, depositing metal adatoms on oxides is known to boost the reactivity via electron transfer to adsorbed molecules.<sup>10</sup> It has been further shown that metal clusters deposited on defect-rich MgO films were catalytically more active than the stoichiometric oxide.<sup>11–14</sup> Theoretical calculations have provided strong support that this was related to neutral and singly charged oxygen vacancies at the surface of MgO, also known as the  $F_s^-$  and  $F_s^+$  surface color centers.<sup>15,16</sup>

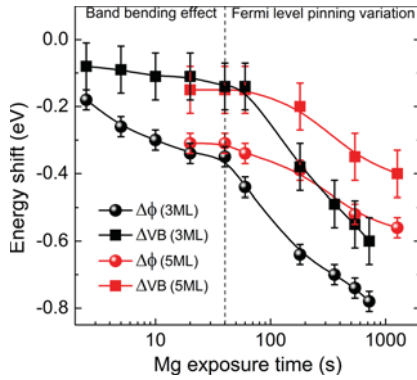
Therefore, ultrathin, colored oxide films can, in principle, allow us to reach optimal catalytic properties. However, the design of such advanced catalytic devices is not only facing the complexity of ultrathin limit and controlled defect generation but also requires a detailed knowledge of electronic structures at interfaces that can host a variety of entangled charge transfer mechanisms. Recent electron paramagnetic resonance (EPR) experiments have reported the ability of producing  $F_s^+$  color centers on 20 ML-thick MgO(001) single-crystalline films by the deposition of small amounts of Mg at low temperatures.<sup>17</sup> These centers have been attributed to unpaired electrons

trapped at morphological defects of the MgO surface in contrast with color centers prepared by electron bombardment that originate in electron trapping within oxygen vacancies.<sup>18</sup>

We report spectroscopic evidence of the controlled creation of Mg-induced color centers on ultrathin MgO films grown on Ag(001). We further show that they serve as host for subsequent O<sub>2</sub> activation by favoring superoxide O<sub>2</sub><sup>-</sup> anion formation. The associated charge-transfer mechanisms are identified by using X-ray and ultraviolet photoemission spectroscopy (XPS-UPS) for monitoring, layer-by-layer, the impact of Mg and O<sub>2</sub> treatments on the electronic band profiles, that is, on band bending.

### EXPERIMENTAL METHODS

All experiments were performed in a multichamber ultrahigh vacuum (UHV) system with base pressures below  $2 \times 10^{-10}$  mbar. A molecular beam epitaxy chamber equipped with reflection high-energy electron diffraction (RHEED) is interconnected with XPS-UPS analysis chamber via a transfer chamber. The (001)-oriented Ag single crystal was cleaned by several cycles of Ar<sup>+</sup> ion bombardment and annealing at 670–720 K for 30 min. The cycles were repeated until a good RHEED pattern and a reproducible work function resulted. The final stable value of the work function of Ag(001) was  $\phi_m = 4.38 \pm 0.05$  eV. The 3 monolayers (ML) thick MgO films (one monolayer is defined as one-half MgO lattice parameter, i.e., 1 ML = 2.105 Å) were grown on the prepared Ag(001) surface by evaporation of Mg in O<sub>2</sub> background atmosphere



**Figure 1.** Variations of the work function  $\Delta\phi$  and MgO-valence band position  $\Delta VB$  as a function of the Mg exposure time for 3 ML MgO/Ag(001) (black lines) and 5 ML MgO/Ag(001) (red lines) samples. The evolutions are plotted in semilogarithmic representation.

(oxygen pressure =  $5 \times 10^{-7}$  mbar) at 453 K with a cube-on-cube epitaxy with respect to the Ag(001) substrate. The Mg ( $2.4 \times 10^{13}$  atoms/(cm<sup>2</sup> s)) and O<sub>2</sub> postexposures have been performed at a substrate temperature of 513 and 300 K, respectively.

The measurements were carried out using XPS and UPS measurements. The kinetic energy of the emitted electrons has been measured by employing a hemispherical analyzer (Omicron EA125) with a five-channels detection system. Al K $\alpha$  was used as the X-ray source, and He-I resonance ( $h\nu = 21.22$  eV) line provided the UPS source for photoemission experiments. The total energy resolutions were, respectively, 0.80 and 0.15 eV for XPS and UPS. The work function of the dielectric system ( $\phi_m^*$ ), defined as the energy of the vacuum level ( $E_{vac}$ ) with respect to the Fermi level of the MgO/Ag(001) system ( $E_F$ ), is determined from the low-energy cutoff ( $E_{cut}$ ) of the secondary photoelectron emission:  $\phi_m^* = h\nu - (E_F - E_{cut})$ . To this end, the samples were biased at  $-8$  V to make the measurements of the very low-energy region of the spectrum reliable.

## RESULTS AND DISCUSSION

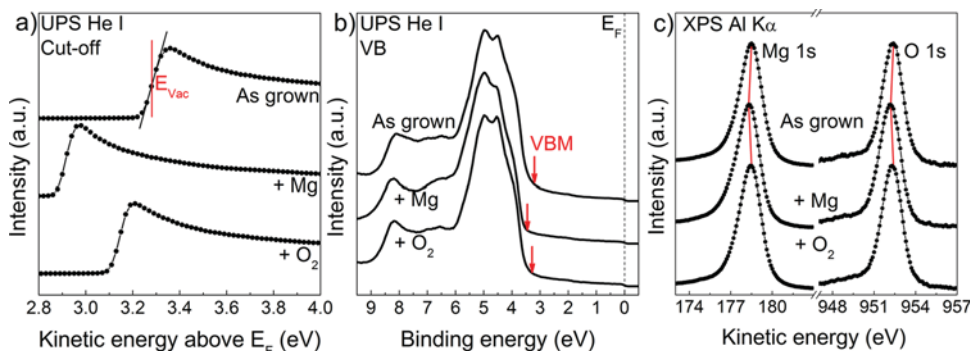
In Figure 1 are shown the variations of the work function ( $\Delta\phi$ ) and of the MgO-valence band position ( $\Delta VB$ ) as a function of the Mg exposition time for two samples consisting of 3 and 5

ML of MgO grown on Ag(001). The methods used for the determination of the work function and valence band (VB) position extracted from our UPS spectra are given in ref 19. The semilogarithmic representation together with the comparison of the energy shifts obtained for the 3 and 5 ML-thick samples allows us to highlight a two-step evolution of the considered quantities as a function of the Mg exposure time. Whereas during the first 20 s the work function shifts are the same for 3 and 5 ML and about two times higher than the VB shifts, further exposition to Mg leads to similar changes of  $\Delta\phi$  and  $\Delta VB$  but whose magnitude is highly dependent on the MgO thickness. For 5 ML, the shifts observed after 720 s of Mg exposure reach  $\sim -0.25$  eV, that is, two times lower than those obtained for the 3 ML thick sample.

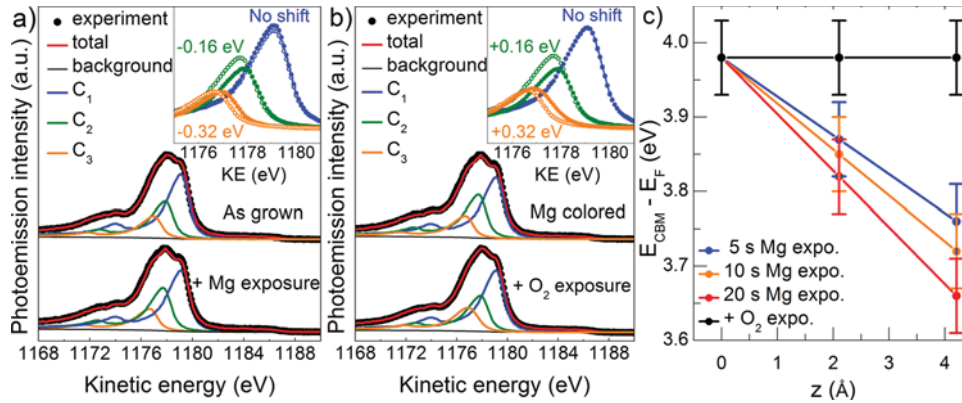
Thus our observations reveal two distinct regimes in the Mg diffusion and adsorption mechanisms as a function of both the Mg flux and the MgO thickness. Because in the first phase  $\Delta\phi$  and  $\Delta VB$  show the same variations and values irrespective of the MgO thickness, we conclude that we are facing a surface mechanism. The measured  $\Delta\phi$  essentially reflects a band bending effect resulting from a positive charge accumulation at the surface of the dielectric layer.

In the second phase, between 40 and 720 s of Mg exposures, we observe similar changes of  $\Delta\phi$  and  $\Delta VB$  for both the 3 and 5 ML thick samples, thus demonstrating that the Fermi level pinning position progressively changes at the MgO/Ag(001) interface. As discussed in ref 20, this is driven by the Mg atom incorporation at the MgO/Ag(001) interface. Note that, in this phase, the variation of the valence band and work function for 3 ML is twice the one observed for 5 ML-MgO. This thickness dependence suggests that the probability that an Mg atom reaches the MgO/Ag(001) interface is inversely proportional to the film thickness.

Figure 2 shows the low-energy cutoff of the secondary photoelectrons emission (Figure 2a), the He-I UPS spectra of the valence band (Figure 2b), and the Al K $\alpha$  spectra of the Mg 1s and O 1s core levels (Figure 2c) for the MgO(3 ML)/Ag(001) sample and after successive exposures to a Mg and O<sub>2</sub> flux. The Mg and O<sub>2</sub> exposures have been performed for 20 and 240 s, respectively. The energy reference is  $E_F$ , the low-energy cutoff positions is taken at the maximum slope of their rising edges, and the MgO valence band maximum (VBM) is obtained via a linear extrapolation of the leading edge of the



**Figure 2.** (a) Low-energy cutoff of the secondary photoelectrons emission for the MgO(3 ML)/Ag(001) sample (top) and after successive exposures to a Mg (middle) and O<sub>2</sub> (bottom) flux. The Mg and O<sub>2</sub> exposures have been performed during 20 and 240 s, respectively. The energy reference is taken at  $E_F$  and the low-energy cutoff positions were taken at the maximum slope of their rising edges, as indicated on the top spectrum. (b) He-I UPS spectra showing the valence band position of the MgO reference sample (top) and after Mg (middle) and O<sub>2</sub> exposures (bottom). The energy reference is taken at the Fermi level  $E_F$ . The VBM positions are indicated by the red arrows. (c) Al K $\alpha$  spectra of the Mg 1s and O 1s core levels of the MgO reference sample (top) and after Mg (middle) and O<sub>2</sub> exposures (bottom). The red lines indicate the induced core level shifts.



**Figure 3.** (a) Photoemission spectra of the Mg  $KL_{23}L_{23}$  Auger transition of the as-grown (top) and Mg-exposed (bottom) 3 ML-thick MgO samples. The inset shows the interface ( $C_1$ ), subsurface ( $C_2$ ), and surface Auger ( $C_3$ ) components of the fitted Auger spectra obtained on the reference (full markers) and Mg-exposed (empty markers) samples. (b) Photoemission spectra of the Mg  $KL_{23}L_{23}$  Auger transition of the 3 ML Mg-exposed MgO film before (top) and after (bottom)  $O_2$  exposure. The inset shows the interface ( $C_1$ ), subsurface ( $C_2$ ), and surface Auger ( $C_3$ ) components of the fitted Auger spectra obtained on the Mg-colored (empty markers) and  $O_2$ -exposed (full markers) samples. Best fit and layer-by-layer decomposition are also shown. (c) Evolution of the band bending through the MgO oxide layer extracted from the layer-resolved photoemission measurements for the different Mg exposure times and after  $O_2$  exposure (black line). The position of the conduction band minimum (CBM) with respect to the Fermi level is taken as reference. The 3 ML MgO/Ag(001) clean sample is initially considered to be under flat-band condition.

valence band (VB) spectrum<sup>19</sup> that consists of overlapped MgO  $O2p$  states (main structures at 4.5 and 8 eV) and Ag 4d states of the metal substrate (around 4.5 eV).<sup>21</sup> The Mg flux exposure leads to a work function reduction of the metal/oxide system by  $0.37 \pm 0.05$  eV (Figure 2a) and to a VBM shift of  $0.16 \pm 0.10$  eV toward higher binding energy (Figure 2b), in close agreement with the Mg 1s ( $-0.20 \pm 0.05$  eV) and O1s ( $-0.20 \pm 0.05$  eV) core-level kinetic energy shifts (red lines Figure 2c). As we can see, further exposures to  $O_2$  now reveal that the work-function shifts can be inverted ( $+0.23 \pm 0.05$  eV) and the VBM ( $+0.16 \pm 0.05$  eV) and core-level shifts ( $+0.20 \pm 0.05$  eV and  $+0.16 \pm 0.05$  eV for Mg 1s and O1s, respectively) fully canceled.

Let us now focus on the evolution of the Mg  $KL_{23}L_{23}$  Auger emission through the different surface treatments. For 3 ML of MgO, the Mg  $KL_{23}L_{23}$  Auger transition is indeed a very powerful spectroscopic probe thanks to its layer-by-layer resolution originating in image potential screening at the metal/oxide interface.<sup>20</sup> It has been, for example, exploited to access, layer-by-layer, the atomic structure of MgO ultrathin films<sup>22</sup> as well as the unoccupied interface and surface states of the MgO/Ag(001) system.<sup>23</sup>

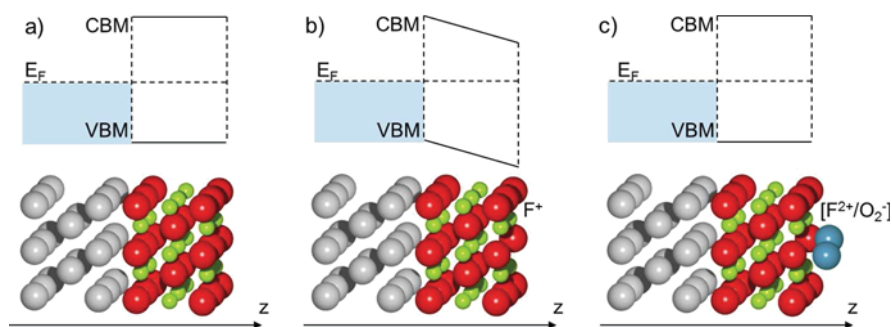
In the present study, the layer-resolved Mg  $KL_{23}L_{23}$  Auger emission further allows us to keep track of band profiles from the oxide surface to the interface. Figure 3a shows normal-emission Mg  $KL_{23}L_{23}$  Auger spectra (vertically shifted for clarity) of the 3 ML MgO sample before and after exposition to the Mg atomic flux. Figure 3b then compares the Auger spectra obtained after Mg (top spectrum) and  $O_2$  (bottom spectrum) exposures. The 3 ML reference spectrum is fitted by three Auger components  $C_1$ ,  $C_2$ , and  $C_3$ , with maxima situated at 1179.1, 1177.8, and 1176.8 eV and which correspond to Auger electron emission from the first, second, and third MgO plane above the metal-oxide interface.

The band bending induced by an Mg exposure of 20 s is clearly revealed in Figure 3a through the evolution of the  $C_1$ ,  $C_2$ , and  $C_3$  positions. The interface component shows no shift after Mg exposure whereas  $C_2$  and  $C_3$  have lowered kinetic energies by 0.16 and 0.32 eV, respectively. Also note that for all Mg exposure times (Figure 3c) the  $C_3$  shift is twice the one of

$C_2$  and similar to the one of the low-energy cutoff. This demonstrates that the work function shift induced by Mg exposures is fully connected to an increased linear and downward band bending in the oxide film and that no bulk defects are present.

Therefore, we conclude that Mg exposures lead to the creation of positively charged donor-type electronic states at the oxide surface. The induced surface charges depend on the Mg exposure times and the linear band bending demonstrates that the electron transfer occurs between the MgO surface and the metal substrate (Figure 4a,b). As demonstrated in electron paramagnetic resonance (EPR) experiments, the deposition of small amounts of Mg at low temperatures on 20 ML-thick MgO(001) single-crystalline films leads to the creation of  $F_s^+$  color centers.<sup>17</sup> Identically, we assign our Mg-induced donor-type surface states to  $F_s^+$  color centers that have been shown to concern unpaired electrons trapped at morphological defects of the MgO surface. Next, assuming only  $F_s^+$  centers of uniform surface density and a relative dielectric constant of 8.0 for MgO, a classical capacitor model can be used to estimate the amount of Mg-induced charged color centers. Considering a MgO thickness of 3 ML and a work function variation  $\Delta\phi$  of  $-0.37$  eV after 20 s of Mg exposure, the density of  $F_s^+$  centers can be estimated to be  $2.7 \times 10^{13}/\text{cm}^2$ .

Interestingly, the measured work function variation  $\Delta\phi$  obtained for the 5 ML-thick MgO film after 20 s of Mg exposure is identical to the one obtained for the MgO(3 ML)/Ag(001) sample (Figure 1). This implies a lower density of  $F_s^+$  centers induced on the 5 ML-thick MgO film. It has been demonstrated that the ultrathin limit strongly impacts the stability of oxygen vacancy-induced  $F_s$  and  $F_s^+$  centers on MgO/Ag(100).<sup>24</sup> In particular, the doubly occupied  $F_s$  center has been shown to be metastable with respect to the paramagnetic  $F_s^+$  center thanks to favored charge transfer with the Ag substrate. Bearing in mind that the present Mg-induced band bending is linear and reflects electron charge transfer between the MgO surface and the metal substrate, we conclude that Mg exposures of ultrathin MgO/Ag(001) films lead to a controlled creation of surface  $F_s^+$  centers that is enhanced by the ultrathin limit and associated tunneling effects.



**Figure 4.** Schematic band diagrams of the MgO(3 ML)/Ag(001) as-grown sample (a) and after Mg (b) and O<sub>2</sub> exposures (c). The as-grown sample is initially considered to be under flat-band condition. The Mg exposure results in F<sup>+</sup> center creation at the MgO surface and in a downward band bending (b), whereas the O<sub>2</sub> exposure leads to [F<sup>2+</sup>/O<sub>2</sub><sup>-</sup>] complex formation that gives rise to an upward band bending, indicating an electron transfer from the metal to the complexes (c). E<sub>F</sub>, VBM, and CBM, respectively, stand for Ag Fermi level, valence band maximum, and conduction band minimum. The red, green, and gray atoms correspond to oxygen, magnesium and silver, respectively. The O<sub>2</sub> molecule involved in the [F<sup>2+</sup>/O<sub>2</sub><sup>-</sup>] complex is shown in blue.

The analogy between electron-induced (oxygen vacancies) and Mg-induced color centers is finally pushed forward through the surface reactivity of Mg-induced color centers with O<sub>2</sub>. From the Mg *KL*<sub>23</sub>*L*<sub>23</sub> Auger spectrum obtained after O<sub>2</sub> exposure (Figure 3b), we indeed see that whereas the C<sub>1</sub> still does not exhibit energy shift the C<sub>2</sub> and C<sub>3</sub> components appear at higher kinetic energy than the ones of the spectrum measured directly after Mg exposure. More precisely, the positive shifts of C<sub>2</sub> and C<sub>3</sub> correspond to a full cancellation of the Mg-induced downward band bending (black curve in Figure 3c), as observed by UPS (Figure 2). Thus, after the O<sub>2</sub> exposition, the charge transfer between the MgO surface and the metal substrate is close to the one of the as-grown sample. The exposition of the sample to an O<sub>2</sub> molecular atmosphere therefore induces an electron transfer from the interface toward the surface of the MgO/Ag(001) sample.

Theoretical work and EPR measurements have shown that activated molecular oxygen O<sub>2</sub><sup>-</sup> (also called superoxide radical anions) form spontaneously on thin Mo-supported MgO(001) films for oxide thicknesses lower than 8 ML.<sup>8,9</sup> Also, it has been theoretically found that O<sub>2</sub><sup>-</sup> anions (with -0.72e) form on MgO(2 ML)/Ag(001)<sup>7</sup> by charge transfer from both the Ag substrate (0.18e) and the underlying MgO layer (0.54e) to the O<sub>2</sub> molecule. We can therefore expect that the adsorption and activation of molecular oxygen on terraces of MgO ultrathin films substantially reduces the band bending (giving an upward bending) of the MgO layer and increases the work function of the metal-dielectric system.

Besides, whereas O<sub>2</sub> molecules do not adsorb on regular MgO(001) surfaces of bulk monocrystals,<sup>25</sup> the presence of surface F<sub>s</sub> and F<sub>s</sub><sup>+</sup> centers allows us to create O<sub>2</sub><sup>-</sup> surface species<sup>26</sup> through the formation of O<sub>2</sub><sup>-</sup>/F<sub>s</sub><sup>+</sup> and O<sub>2</sub><sup>-</sup>/F<sub>s</sub><sup>2+</sup> complexes by electron transfer from the color centers to the empty level of the adsorbed molecule.<sup>27,28</sup> We propose that such kind of complexes is also created at the surface of the MgO ultrathin film during the interaction of the colored MgO(001) surface with O<sub>2</sub>. Their formation gives rise to an upward band-bending that demonstrates that an electron transfer from the metal to the complexes takes place during the exposure phase to molecular oxygen (Figure 4b,c). The fact that MgO recovers the flat band condition implies that the surface is close to neutrality condition in the final stage (Figure 4c).

Thus as Mo-doped CaO<sup>29</sup> or Au chains supported on MgO/Mo thin films,<sup>10</sup> we believe that our Mg-colored 3 ML-thick

MgO/Ag(001) system favors O<sub>2</sub> activation through an enhanced charge transfer mechanism between the oxide surface and the metal support. Further ab initio calculations should confirm our spectroscopic findings and provide a deeper understanding of adsorption processes of O<sub>2</sub> on the low-coordinated surface sites of metal-supported MgO(001)-ultrathin films.

## CONCLUSIONS

We probed charge-transfer mechanisms acting at interfaces of Ag-supported ultrathin MgO films successively exposed to Mg and O<sub>2</sub> flux by photoemission spectroscopy and we reported evidence of controlled creation of Mg-induced color centers and catalytic enhancement of O<sub>2</sub> activation. Our findings give new insight for studying electronic properties of dielectric/metal systems such as work functions, defect states, and Schottky barrier heights and their relationship to catalytic processes.

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### Notes

The authors declare no competing financial interest.

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