

# ATCA detections of massive molecular gas reservoirs in dusty, high- $z$ radio galaxies

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## ABSTRACT

Observations using the 7-mm receiver system on the Australia Telescope Compact Array have revealed large reservoirs of molecular gas in two high-redshift radio galaxies: HATLAS J090426.9+015448 ( $z = 2.37$ ) and HATLAS J140930.4+003803 ( $z = 2.04$ ). Optically, the targets are very faint, and spectroscopy classifies them as narrow-line radio galaxies. In addition to harbouring an active galactic nucleus the targets share many characteristics of sub-mm galaxies. Far-infrared data from *Herschel*-Astrophysical Terahertz Large Area Survey suggest high levels of dust ( $>10^9 M_{\odot}$ ) and a correspondingly large amount of obscured star formation ( $\sim 1000 M_{\odot} \text{ yr}^{-1}$ ). The molecular gas is traced via the  $J = 1 \rightarrow 0$  transition of  $^{12}\text{CO}$ , its luminosity implying total  $\text{H}_2$  masses of  $(1.7 \pm 0.3) \times 10^{11}$  and  $(9.5 \pm 2.4) \times 10^{10}$  ( $\alpha_{\text{CO}}/0.8$ )  $M_{\odot}$  in HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively. Both galaxies exhibit molecular line emission over a broad ( $\sim 1000 \text{ km s}^{-1}$ ) velocity range and feature double-peaked profiles. We interpret this as evidence of either a large rotating disc or an on-going merger. Gas depletion time-scales are  $\sim 100 \text{ Myr}$ . The 1.4-GHz radio luminosities of our targets place them close to the break in the luminosity function. As such they represent ‘typical’  $z > 2$  radio sources, responsible for the bulk of the energy emitted at radio wavelengths from accretion-powered sources at high redshift, and yet they rank amongst the most massive systems in terms of molecular gas and dust content. We also detect 115-GHz rest-frame continuum emission, indicating a very steep high-radio-frequency spectrum, possibly classifying the targets as compact steep spectrum objects.

**Key words:** galaxies: high-redshift – radio continuum: galaxies – radio lines: galaxies.

## 1 INTRODUCTION

Radio galaxies are one of several classes of object which harbour an active galactic nucleus (AGN; Krawczynski & Treister 2013).

They are characterized by their extended radio emission, which due to its high luminosity allows such objects to be detected out to very large distances. This property has allowed them to be used as cosmological probes since the discovery over half a century ago that the brightest radio sources in the sky were associated with galaxies at cosmologically significant distances (Minkowski 1960). Identifying high-redshift radio galaxies (HzRGs) is an effective way to

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pinpoint the sites of the most massive galaxies undergoing formation in the early Universe (Eales et al. 1997; Jarvis et al. 2001; Miley & De Breuck 2008), as well as cluster environments at various stages of evolution (Stevens et al. 2003; Wylezalek et al. 2013; Hatch et al. 2014; Cooke et al. 2016). The present-day end-points of these systems are thought to be the giant elliptical galaxies which reside at the centres of rich clusters (McLure et al. 1999; Orsi et al. 2016).

One of the most effective ways to study galaxies and galaxy groups at high redshift is to examine the cold gas in and around them via observations of molecular or atomic fine structure lines. A thorough review of this subject is given by Carilli & Walter (2013). The molecular phase of the interstellar medium (ISM) is dominated by  $\text{H}_2$ ; however, the usual observational tracers are the  $J \rightarrow (J - 1)$  rotational transitions of  $^{12}\text{C}^{16}\text{O}$  (hereafter CO), as this is the next most abundant molecule. The high excitation temperature of  $\text{H}_2$  means it can only be observed directly in e.g. the presence of strong shocks; however, the low excitation temperature of 5.5 K that the CO molecule has means that its ground-state transition can be used to trace even the normal, cold ISM within a galaxy. The 115.27-GHz rest frequency means that at the peak cosmological epoch of galaxy assembly, between redshifts of 1 and 3, the line is observable in the frequency ranges of many ground-based cm- and mm-wave facilities.

Spatially and spectrally resolving this line allows much to be inferred about the dynamics of high-redshift systems (e.g. Hodge et al. 2012). The star-forming potential of a galaxy can be established as the cold gas represents the fuel reservoir for this process, which when coupled with a star formation rate (SFR) estimate allows gas consumption time-scales to be determined (e.g. Daddi et al. 2010). Thermodynamic modelling of the spectral energy distribution (SED) described by multiple  $J$  transitions allows physical conditions of the ISM to be constrained (Obreschkow et al. 2009).

There has been a steady stream of studies of individual high-redshift systems via their CO lines since the first detection of the line at high redshift nearly a quarter of a century ago (IRAS F10214+4724 at  $z = 2.3$ ; Brown & Vanden Bout 1991; Solomon, Radford & Downes 1992). Such studies have now looked back in cosmic time to the very first galaxies to form in the Universe (Wang et al. 2013). The last half-decade has seen a significant increase in the quantity and quality of the data being recorded in this field, largely due to significant upgrades to the bandwidth and sensitivity of existing radio interferometers such as the Karl G. Jansky Very Large Array, the Plateau de Bure Interferometer (PdBI) and the Compact Array Broad-band Back-end (CABB; Wilson et al. 2011) for the Australia Telescope Compact Array (ATCA), and the deployment of the Atacama Large Millimetre/submillimetre Array (ALMA). Broad bandwidths have unlocked the potential for blind molecular line scans for the first time, both for single objects (Harris et al. 2012; Vieira et al. 2013) and for searches in extragalactic deep fields over significant cosmological volumes (Lentati et al. 2014; Walter et al. 2014).

Radio galaxies being signposts for massive systems at high- $z$  makes them prime targets for follow-up molecular line observations, the first tentative CO detection in an HzRG being reported nearly 20 years ago by Scoville et al. (1997). Since then numerous studies have shown that radio galaxies have molecular gas properties similar to those of quasars and, like quasars, often exhibit similarities to sub-mm galaxies (SMGs), e.g. Papadopoulos et al. (2000). Typical inferred  $\text{H}_2$  masses exceed  $10^{10} M_{\odot}$ . Multiple components are often seen, corroborating the scenario that

radio galaxy triggering is associated with merger activity or proto-cluster formation (Greve, Ivison & Papadopoulos 2004; De Breuck et al. 2005; Ivison et al. 2012, 2013), a picture further supported by the often-disturbed optical morphologies seen at lower redshifts (Ramos Almeida et al. 2011). Several studies have reported alignments between radio jets and the peaks of the CO emission (Klamer et al. 2005; Emonts et al. 2014). The latter study reports CO peaks beyond the boundary of the radio emission, suggesting the radio jets play a role in stimulating star formation and metal enrichment. Molecular line emission has also been detected in the outer regions of radio galaxies (Nesvadba et al. 2009), likely associated with the Ly  $\alpha$  haloes which extend for tens to hundreds of kiloparsecs beyond the central radio galaxy (Jarvis et al. 2003; Wilman et al. 2004). A spectacular example of this is the ‘Dragonfly’ galaxy at  $z = 2$ , where 60 per cent of the total CO content is associated with tidal features within the halo (Emonts et al. 2015).

Here, we add to the existing body of work on characterizing the molecular gas in high-redshift systems, presenting ATCA observations of CO ( $J = 1 \rightarrow 0$ ) in a pair of HzRGs. The 1.4-GHz radio luminosities place the two targets close to the break in the radio luminosity function, distinguishing them from existing molecular line studies which have typically targeted more powerful HzRGs (Klamer et al. 2005; Emonts et al. 2011). Our targets therefore represent the ‘typical’ radio source at  $z > 2$ , responsible for the bulk of the energy density emanating at radio wavelengths from high accretion rate radio galaxies. Throughout the paper the assumed cosmological parameters are  $H_0 = 67.74 \pm 0.46 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_M = 0.3089 \pm 0.0062$  and  $\Omega_{\Lambda} = 0.6911 \pm 0.0062$  (Planck Collaboration XVI 2014).

## 2 TARGETS

Our targets were selected from the sample of HzRGs studied by Virdee (2013), who investigated the radio properties of sources selected from the *Herschel*-Astrophysical Terahertz Large Area Survey (ATLAS) (H-ATLAS; Eales et al. 2010)/Galaxy And Mass Assembly (GAMA; Driver et al. 2011) fields. A sample of seven radio sources were selected for follow-up observations, including long-slit spectroscopy with the Intermediate-dispersion Spectrograph and Imaging System on the William Herschel Telescope in order to determine redshifts, and very long baseline interferometry (VLBI) observations with the European VLBI Network (EVN) in order to disentangle the core, jet and star formation components of the radio emission. The sample was constructed based on the fact that the far-infrared (FIR) and radio emission suggested that the sources harboured both an AGN, exhibited a significant amount of star formation and were at  $z > 1$ . The systems selected are also optically very faint. The spectroscopic follow-up by Virdee (2013) revealed emission line properties which classified the target sources as narrow-line radio galaxies. Candidate lines were fitted to the spectral emission peaks and secondary lines were then searched for at their expected positions in order to confirm the candidate line. Our two targets only had single lines in each spectrum, which were assumed to be the bright Ly  $\alpha$  line, on the basis which other typically bright lines (e.g. C IV, O II) should have had secondary lines visible in the spectrum.

FIR luminosities measured from the *Herschel*-ATLAS data imply SFRs of  $\sim 1000 M_{\odot} \text{ yr}^{-1}$  and high levels of dust (Section 4.4).

We targeted two of these sources for ATCA follow-up to search for CO line emission as their redshifts placed the ground-state transition of CO in the 7-mm band of the ATCA. The properties of these

**Table 1.** Properties of HATLAS J090426.9+015448 and HATLAS J140930.4+003803. Parameters from existing observations are listed, together with the results derived in the subsequent sections of this paper, as listed in the final column. For the CO lines, each quantity is listed for the two components present in each spectrum, as labelled on Figs 1 and 2, as well as for the system as a whole. It is not possible to conclusively differentiate between a major merger and a single massive molecular gas disc scenario for the two targets with the data in hand. The individual component measurements and the total measurements are listed in order to capture both possibilities, although we note that existing high-resolution observations of similar systems suggest that a merger is the most likely of the two (Genzel et al. 2003; Ivison et al. 2010a).

Quantity	Units/Epoch	HATLAS J090426.9+015448		HATLAS J140930.4+003803		Reference
RA	J2000	09 <sup>h</sup> 04 <sup>m</sup> 26 <sup>s</sup> .80		14 <sup>h</sup> 09 <sup>m</sup> 30 <sup>s</sup> .23		
Dec	J2000	+01°54′48″.73		+00°38′04″.54		
$z_{\text{Ly}\alpha}$		2.372		2.044		Virdee (2013)
$S_{100}$	mJy	2.7 ± 46.8		113.1 ± 52.4		H-ATLAS
$S_{160}$	mJy	−57.7 ± 66.8		−50.07 ± 70.1		H-ATLAS
$S_{250}$	mJy	45.8 ± 7.0		38.5 ± 6.8		H-ATLAS
$S_{350}$	mJy	51.1 ± 8.0		47.1 ± 8.1		H-ATLAS
$S_{500}$	mJy	45.1 ± 8.8		22.0 ± 8.9		H-ATLAS
$S_{500}$	mJy	45.1 ± 8.8		22.0 ± 8.9		H-ATLAS
$L_{1.4\text{GHz}}$	$10^{26} \text{ W Hz}^{-1} \text{ sr}^{-1}$	0.99 ± 0.07		6.25 ± 0.14		Section 4.1
$L'_{\text{CO}}$	$\text{K km s}^{-1} \text{ pc}^2$	$(2.11 \pm 0.37) \times 10^{11}$		$(1.18 \pm 0.30) \times 10^{11}$		Section 4.4
$M_{\text{H}_2}$	$(\alpha_{\text{CO}}/0.8) M_{\odot}$	$(1.69 \pm 0.29) \times 10^{11}$		$(9.45 \pm 2.36) \times 10^{10}$		Section 4.4
$L_{\text{FIR}} (8\text{--}1000 \mu\text{m})$	$L_{\odot}$	$(11.75 \pm 1.62) \times 10^{12}$		$(7.59 \pm 1.22) \times 10^{12}$		Section 4.4
$q_{\text{IR}}$		1.92 ± 0.13		0.81 ± 0.10		Section 4.4
SFR	$M_{\odot} \text{ yr}^{-1}$	1250 ± 160		810 ± 130		Section 4.4
$T_{\text{dust}}$	K	30 ± 4		31 ± 6		Section 4.4
$L_{\text{dust}}$	$L_{\odot}$	$(11.5 \pm 1.6) \times 10^{12}$		$(7.9 \pm 1.1) \times 10^{12}$		Section 4.4
$M_{\text{dust}}$	$M_{\odot}$	$(2.3 \pm 1.1) \times 10^9$		$(1.3 \pm 0.7) \times 10^9$		Section 4.4
$M_{\text{dyn}} \sin^2(i)$	$M_{\odot}$	$<(5.67 \pm 0.17) \times 10^{12}$		$<(9.24 \pm 0.27) \times 10^{11}$		Section 4.3
		Component 1	Component 2	Component 1	Component 2	Figs 1,2
$\nu_{\text{centre}}$	GHz	34.1475 ± 0.0040	34.2397 ± 0.0022	37.8293 ± 0.0038	37.8708 ± 0.0056	Section 4.2
$\nu_{\text{peak}}$	mJy beam <sup>−1</sup>	0.352 ± 0.145	0.5821 ± 0.0575	0.436 ± 0.077	0.374 ± 0.061	Section 4.2
$\nu_{\text{FWHM}}$	km s <sup>−1</sup>	173 ± 80	414 ± 46	228 ± 63	318 ± 103	Section 4.2
$z_{\text{CO}}$		2.3756 ± 0.0004	2.3666 ± 0.0002	2.0471 ± 0.0003	2.0437 ± 0.0004	Section 4.2
$D_{\text{L}}$	Gpc	19.652 ± 0.293	19.562 ± 0.291	16.397 ± 0.242	16.364 ± 0.242	Section 4.2
$I_{\text{CO}}$	Jy km s <sup>−1</sup>	0.153 ± 0.095	0.605 ± 0.090	0.249 ± 0.082	0.299 ± 0.109	Section 4.4
$L'_{\text{CO}}$	$\text{K km s}^{-1} \text{ pc}^2$	$(4.27 \pm 2.66) \times 10^{10}$	$(1.68 \pm 0.25) \times 10^{11}$	$(5.37 \pm 1.78) \times 10^{10}$	$(6.43 \pm 2.35) \times 10^{10}$	Section 4.4
$M_{\text{H}_2}$	$(\alpha_{\text{CO}}/0.8) M_{\odot}$	$(3.42 \pm 2.13) \times 10^{10}$	$(1.35 \pm 0.20) \times 10^{11}$	$(4.30 \pm 1.32) \times 10^{10}$	$(5.15 \pm 1.88) \times 10^{10}$	Section 4.4
$M_{\text{dyn}} \sin^2(i)$	$M_{\odot}$	$<(1.4 \pm 1.3) \times 10^{11}$	$<(8.1 \pm 1.8) \times 10^{11}$	$<(1.1 \pm 0.6) \times 10^{11}$	$<(2.1 \pm 1.4) \times 10^{11}$	Section 4.3

two sources are given in Table 1, including the results derived in subsequent sections of this paper, as noted in the final column.

### 3 OBSERVATIONS AND DATA REDUCTION

We made use of the ATCA in its most compact H75 configuration, the northern spur of the array being necessary due to the proximity of our targets to the celestial equator. Observations<sup>1</sup> were conducted over five nights for each target, avoiding antenna elevations below 30° in order to keep the system temperature low. The 7-mm receivers were used, with the CABB configured to deliver 2 × 2 GHz basebands, each consisting of 2,048 × 1 MHz channels. For this project, we made use of only one of the basebands, namely those which were expected to contain the CO line. For HATLAS J090426.9+015448, the frequency coverage was 33.416–35.464 GHz, and for HATLAS J140930.4+003803 it was 37.100–39.148 GHz.

The strong source PKS 1253–055 (3C 279) was observed once per night in order to calibrate the bandpass, and we set the flux density scale using the standard calibrator PKS B1934–638. Nearby (<2° separation), compact secondary calibrator sources PKS B0906+015 for HATLAS J090426.9+015448 and PKS

B1356+022 for HATLAS J140930.4+003803) were observed for 2 min after every 10 min of target observation in order to calibrate the complex antenna gains.

#### 3.1 Calibration

The 10 nights’ worth of data were initially loaded into the MIRIAD package (Sault, Teuben & Wright 1995), using ATLOD, which was set to automatically flag channels which are known to contain radio frequency interference as well as the band edges where the gain drops due to the filter responses. Autocorrelations were also discarded at this stage. Antenna positions were updated using ATFIX. Following this step, the data were converted to Measurement Set format using CASA’s (McMullin et al. 2007) IMPORTMIRIAD task, and the baseband relevant to the CO emission was extracted using the SPLIT task. The 10 Measurement Sets were examined using PLOTMS and any obviously bad data were removed. Antenna 6 was also discarded outright. It forms baselines with the other five antennas which are approximately 4.4 km in length, the next longest baseline for the H75 configuration being 89 m. Such a large gap in the ( $u$ ,  $v$ ) plane coverage does not usefully contribute to a higher resolution image without discarding most of the sensitivity provided by the inner baselines. The BANDPASS task was used to derive normalized, time-independent solutions for each 1-MHz frequency channel

<sup>1</sup> Project code: C2847

from the scans of 1253–055. A single frequency-independent complex gain correction was derived for each scan of the secondary calibrator using the `GAINCAL` task. The flux density scale of these solutions were then corrected using the `FLUXSCALE` task and a model for PKS B1934–638. The bandpass and flux-calibrated complex gain solutions were then applied to the target sources, using a linear interpolation in time for the latter set of corrections, via the `APPLYCAL` task. Note that the `CASA SETJY` model for PKS B1934–638 does not reliably extend to these frequencies. The polynomial spectral model provided by the ATCA Calibrator Database<sup>2</sup> was implemented instead. The `SETJY` task was mimicked by producing per-channel point source models of appropriate brightness using `CASA`'s `CL` tool, and filling the `MODEL` column of the Measurement Set using the `FT` task. The total on-source observing time following the removal of bad data was 16.13 and 22.28 h for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively.

### 3.2 Imaging and continuum subtraction

Following calibration, a full-band continuum image of each target was made using multifrequency synthesis (MFS) imaging. No deconvolution was applied initially. The dirty map was inspected and it revealed a compact feature at the expected position of each target. This step was necessary because the declinations of our targets result in point spread functions (PSFs, or dirty/synthesized beams) with pronounced north–south sidelobe features which rise to up to 95 per cent of the peak value. Any deconvolution must therefore be done with extreme caution to avoid sidelobes being interpreted as genuine emission and biasing the true flux density measurement. Having identified the central peak and compared the associated sidelobes to those of the PSF, a shallow (100 clean iterations) deconvolution was allowed to proceed, and within a region which contained only the central peak. The resulting images are presented in Section 4.1. The restoring beams used in the continuum images are 2D Gaussians with full width at half-maximum (FWHM; major axis  $\times$  minor axis) extents of  $19'' \times 15''$  (position angle =  $-57^\circ$ , east of north) and  $16'' \times 14''$  (position angle =  $-82^\circ$ ) for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively.

In order to search for spectral line emission, the continuum emission was subtracted using the `CASA`'s `UVCONTSUB` task. This task estimates the continuum emission by fitting polynomials (in this case a simple first-order polynomial) to the real and imaginary parts of the visibility spectrum (e.g. Sault 1994). A solution was generated for each 10 min scan of the target, and the resulting model was subtracted from the visibilities. Following this, image cubes were made, with no deconvolution, and spectra were extracted at the peak position of the continuum emission. Incremental averaging in frequency was used to boost the signal-to-noise ratio of the line emission. A spectral averaging factor of 6 was found to reveal a spectrum which had both high significance and adequate velocity resolution. Having identified the region of the spectrum over which line emission was present, the continuum subtraction was repeated, excluding from the fit the regions where the line emission was known to be. This was mainly done out of caution, and made no discernible difference, as one would expect from a low-order polynomial fit and a line width which is narrow compared to the total bandwidth. The resulting spectra are presented in Section 4.2. The velocity resolutions of each 6-MHz channel are 53 and 46 km s<sup>-1</sup> for

HATLAS J090426.9+015448 and HATLAS J140930.4+003803 respectively; however, note that Hanning smoothing was applied to the plotted spectra to emphasize the line emission, which doubles the effective velocity resolution of each channel in the figures.

## 4 RESULTS AND DISCUSSION

### 4.1 Continuum detections at 115-GHz rest frame

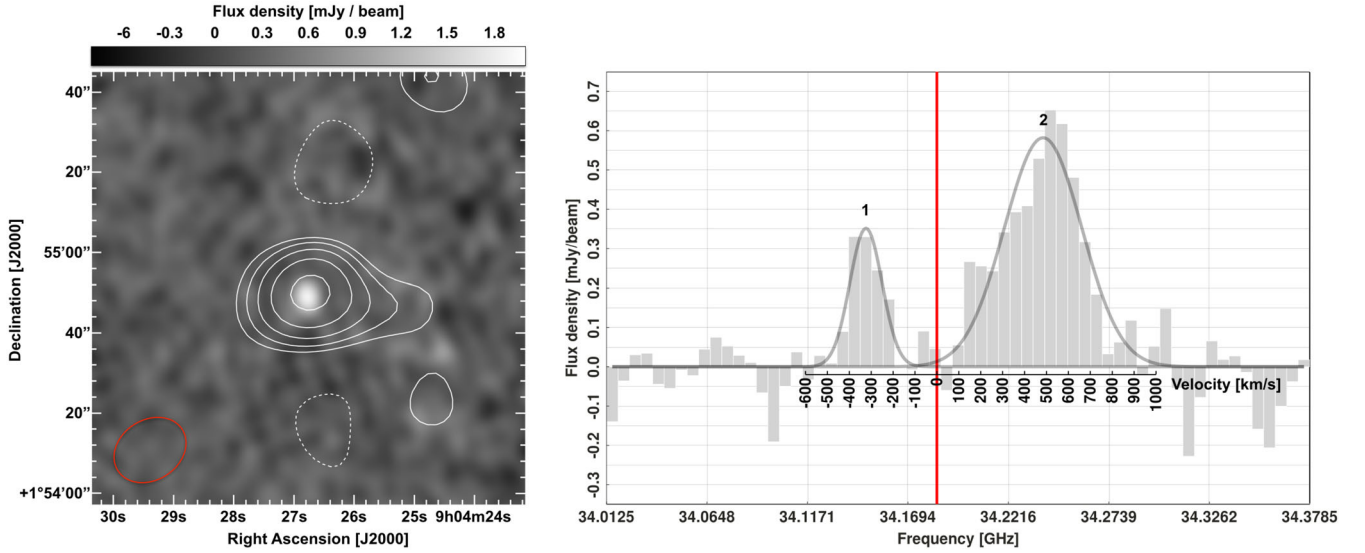
Both of our targets are detected in the ATCA continuum maps, presented in the left-hand panels of Figs 1 and 2. These images show the ATCA contours (115 GHz rest-frame) overlaid on the corresponding 1.4 GHz images from the Faint Images of the Radio Sky at Twenty-cm (FIRST; Becker, White & Helfand 1995) survey. Contour and grey-scale levels are provided in the figure captions. Component fitting to the ATCA detections was performed using the `CASA` task `IMFIT`, which determined the emission to be point-like for both sources. The fitted position for HATLAS J090426.9+015448 is at right ascension  $09^{\text{h}}04^{\text{m}}26^{\text{s}}.711 \pm 0^{\text{s}}.068$  and declination  $+01^{\circ}54'49''.67 \pm 0''.58$ , with a peak flux density of  $160 \pm 16 \mu\text{Jy beam}^{-1}$ . For HATLAS J140930.4+003803, the fitted right ascension and declination are  $14^{\text{h}}09^{\text{m}}30^{\text{s}}.211 \pm 0^{\text{s}}.020$  and  $+00^{\circ}38'04''.62 \pm 0''.23$ , with a peak flux density of  $310 \pm 12 \mu\text{Jy beam}^{-1}$ . Note that the quoted positional errors are those derived from the statistical uncertainties in the component fitting procedure and do not contain any astrometric reference frame errors; however, the latter should be negligible. No self-calibration has taken place so the astrometric accuracy due to the calibration will be tied to the position of the phase calibrator.

Radio continuum spectra for both of the targets are shown in Fig. 3, which features both the NRAO VLA Sky Survey (NVSS) measurement as well as the 327-MHz Giant Metrewave Radio Telescope (GMRT) detection from the catalogues of Mauch et al. (2013). The 1.4-GHz radio luminosities ( $L_{1.4\text{GHz}}$ ) in Table 1 are derived from the NVSS (Condon et al. 1998) flux density measurements and the mean luminosity distances listed in Table 1.

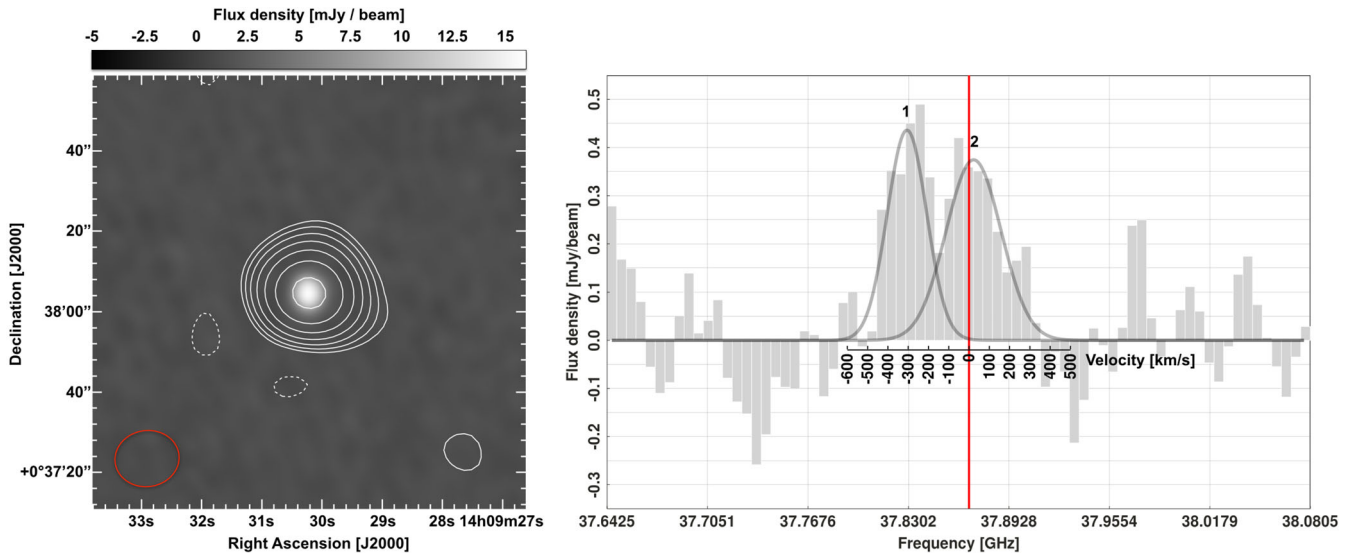
The spectra for both objects exhibit a break at high frequencies. While the spectral break may be explained in terms of the higher frequency emission being core-dominated and the lower frequency emission containing a mixture of core and extended synchrotron emission (e.g. Whittam et al. 2013, 2016), the spectra of these two objects are atypical for HzRGs. Klammer et al. (2006) conducted a search for HzRGs based a spectral steepness selection, finding that the 89 per cent of their sample have radio spectra which are well described by a simple power law, with only 11 per cent exhibiting mild spectral curvature. The two HzRGs studied via their CO lines by Emonts et al. (2011) also display no spectral curvature between 1.4 and 30 GHz (observed frame). Herzog et al. (2016) also report that a high fraction ( $\sim 75$  per cent) of the radio SEDs of infrared-faint radio sources (IFRS) can be modelled by a simple power law, extending as far as 105 GHz (observed frame). The radio properties of IFRS are reported to be consistent with the general population of radio-loud AGN at high redshift, with similar fractions of compact steep spectrum (CSS) and gigahertz-peaked spectrum (GPS) sources. CSS and GPS sources have concave radio spectra with steep spectral slopes either side of a peak (see O'Dea 1998 for a review). Any extended radio emission associated with these sources exists in a small region, well within the host galaxy, either due to youth or environmental confinement (e.g. Marr et al. 2014). If the spectral peak for our two targets, lying somewhere below 1 GHz, gives way to a low-frequency turnover they could reasonably be classified as faint, high-redshift CSS sources.

<sup>2</sup> <http://www.narrabri.atnf.csiro.au/calibrators/>





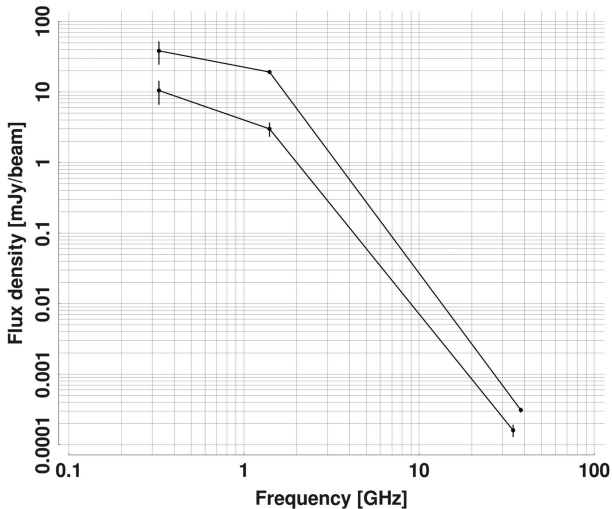
**Figure 1.** HATLAS J090426.9+015448 ( $z = 2.372$ , marked with vertical line). Left-hand panel: the ATCA-detected radio continuum emission from HATLAS J090426.9+015448. The grey-scale shows the 1.4-GHz FIRST detection and the contours show the ATCA-detected radio emission (SNR = 14) at 34.44 GHz (115-GHz rest frame). Contour levels are  $22 \times (1, \sqrt{2}, 2, 2\sqrt{2}, 4) \mu\text{Jy beam}^{-1}$ . There is a single negative contour at  $-22 \mu\text{Jy beam}^{-1}$ . The  $1\sigma$  background noise level in the ATCA map is  $14 \mu\text{Jy beam}^{-1}$ . The extended feature to the west of the peak is not believed to be significant, a result of residual phase errors. The restoring beam used following deconvolution is a Gaussian with a major axis of 19 arcsec, a minor axis of 15 arcsec and a position angle of  $-57^\circ$  east of north, as indicated by the ellipse in the lower left. Right-hand panel: histogram showing the CO ( $J = 1 \rightarrow 0$ ) spectrum of HATLAS J090426.9+015448. The vertical line is the Ly  $\alpha$  redshift (Virdee 2013) and the velocity of the molecular gas with respect to this is shown on the panel. The solid lines show the two Gaussian components fitted to the peaks as described in Section 4.4. The  $1\sigma$  image noise values, measured in each plane of the image cube over the channels shown above, has a mean value of  $0.21 \text{ mJy beam}^{-1}$  with a standard deviation of  $0.02 \text{ mJy beam}^{-1}$ . The noise in the plotted spectrum appears considerably lower due to the effects of the Hanning smoothing.



**Figure 2.** HATLAS J140930.4+003803 ( $z = 2.044$ , marked with vertical line). Left-hand panel: the ATCA-detected radio continuum emission from HATLAS J140930.4+003803. Grey-scale and contours are as per Fig. 1. The 38.124-GHz (115-GHz rest-frame) ATCA detection has an SNR of 18. Contour levels are  $20 \times (1, \sqrt{2}, 2, 2\sqrt{2}, 4, 4\sqrt{2}, 8) \mu\text{Jy beam}^{-1}$ , with a single negative contour at  $-16 \mu\text{Jy beam}^{-1}$ . The  $1\sigma$  noise level in the ATCA map is  $16 \mu\text{Jy beam}^{-1}$ . The restoring beam has a major axis of 16 arcsec, a minor axis of 14 arcsec and a position angle of  $-82^\circ$  east of north, as shown in the lower left. Right-hand panel: histogram showing the CO ( $J = 1 \rightarrow 0$ ) spectrum of HATLAS J140930.4+003803. The vertical line is the Ly  $\alpha$  redshift (Virdee 2013). Properties of the two Gaussian components fitted to the peaks are described in Section 4.4. The  $1\sigma$  image noise values, measured in each plane of the image cube over the channels shown above, has a mean value of  $0.22 \text{ mJy beam}^{-1}$  with a standard deviation of  $0.02 \text{ mJy beam}^{-1}$ . Again, the Hanning smoothing makes the noise in the plotted spectrum appear much lower.

The 1.4–35 GHz (observed frame) spectral indices ( $\alpha$ , where flux density  $S \propto \nu^\alpha$ ) as measured from Fig. 3 are  $-3.1 \pm 0.1$  and  $-3.34 \pm 0.03$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively. Even for a steep-spectrum

radio galaxy (e.g. Sadler et al. 2014), this is excessively steep. Returning to our previous possible explanation for the observed spectral break, the continuum spectrum may be rendered artificially steep if the existing lower frequency measurements



**Figure 3.** Radio continuum spectra of HATLAS J090426.9+015448 and HATLAS J140930.4+003803. In increasing frequency order, the measurements are from the 325-MHz GMRT survey of Mauch et al. (2013), the 1.4-GHz NVSS survey (Condon et al. 1998) and this work.

are capturing both core and extended emission, and our ATCA observations are sensitive only to a flatter spectrum core contribution. Inverse-Compton losses due to the scattering of electrons by cosmic microwave background photons may also be important at this redshift (Murphy 2009) as a mechanism for suppressing the synchrotron emission. In the event that any extended jet-driven emission is compact (as it would be for CSS or GPS sources), an artificial steepening of the measured spectrum would be more likely occur for multifrequency radio observations with (approximately) matched angular resolution which is insufficient to separate core and jet structures, as we have used here. It is noteworthy that there is no discrepancy between the FIRST and NVSS measurements for HATLAS J140930.4+003803; however for HATLAS J090426.9+015448, the peak flux density measurement from FIRST is 28 per cent lower than the NVSS measurement, suggesting possible extended emission on scales greater than those corresponding to the synthesized beam of FIRST (50 kpc at  $z = 2.37$ ). The EVN imaging presented by Virdee (2013) resulted in no significant detection of emission associated with the latter. For the former, there was a tentative detection of two components, with a separation of 0.8 arcsec (6.8 kpc projected on the plane of the sky at  $z = 2.04$ ) and a combined integrated flux density of  $\sim 12$  mJy, consistent with the CSS classification.

With the data in hand, the only solid conclusion that we may draw is that the radio spectra of these two objects are intriguing. Further observations at intermediate frequencies, as well as sub-327 MHz would be needed to accurately constrain the shape and peak of the continuum spectra for these two sources, and with angular resolution sufficient to separate any extended structures from core emission. Observations of the GAMA fields with the Low Frequency Array are pending as part of its campaign to observe the *Herschel*-ATLAS fields (Hardcastle et al. 2016).

#### 4.2 CO ( $J = 1 \rightarrow 0$ ) detections

The spectra of HATLAS J090426.9+015448 and HATLAS J140930.4+003803 extracted from the continuum-subtracted data at the pixel coincident with the peak of the continuum emission are shown in the right-hand panels of Figs 1 and 2, respectively.

There is strong line emission present, confirming the redshift of what was assumed to be the Ly  $\alpha$  line as measured by Virdee (2013), at the expected frequency for CO ( $J = 1 \rightarrow 0$ ). Note that Virdee (2013) provides no estimate of the uncertainty in the Ly  $\alpha$  redshift.

Velocity-integrated (moment zero) maps formed from the continuum-subtracted channels containing the line emission did not reveal any plausible extended emission. This is perhaps unsurprising given the low angular resolution of the observations; however, without the benefits afforded by MFS, the deconvolution of the narrow channel cubes was not straightforward. Many sidelobe-like features remained, possibly exacerbated by residual, low-level phase calibration errors. The Ly  $\alpha$  redshift is marked on each spectrum with a vertical line, and the velocity of the molecular gas with respect to this is also shown on the figure. Velocity resolution in each 6-MHz channel is 53 and 46 km s $^{-1}$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively, although the effective velocity resolutions of the plotted spectra are twice these values due to the application of Hanning smoothing. The CO lines are well modelled by pairs of Gaussians in flux-density/frequency space, and these fits (to the non-smoothed data) are also shown on the figure. The RMS of the pixel values is measured over the spatial dimensions of each 6 MHz channel in the data cube, and this corresponding quantity is attached to each value in the line spectrum as an estimate of the uncertainty. The mean and standard deviations of these error spectra are provided in the captions of Figs 1 and 2. Note that with the application of Hanning smoothing to the plotted spectra, the apparent noise properties will be deceptively low. Adjacent channels become correlated depending on the width of the filter function, resulting not only in a coarser effective channel resolution but a significantly broader noise equivalent bandwidth for each channel.

The best-fitting parameters and their associated uncertainties for the four components are provided in Table 1, with the FWHM frequency values ( $\nu_{\text{centre}}$ ) being converted to a velocity FWHM ( $\nu_{\text{FWHM}}$ ) about the component centre, and the central frequency also being converted to a redshift ( $z_{\text{CO}}$ ) for CO ( $J = 1 \rightarrow 0$ ). The total velocity-integrated line flux for each component ( $I_{\text{CO}}$ ) is also given. These parameters will feature again in the discussion in Sections 4.3 and 4.4.

#### 4.3 Velocity structure

Frequency structure in atomic and molecular lines which are not spatially resolved can still reveal much about the gas velocities in galaxies and galaxy groups, as well as provide dynamical mass limits. Probably the best-studied examples are local extragalactic systems observed via their neutral hydrogen (H I) lines (e.g. Koribalski et al. 2004). H I lines in disc-dominated systems exhibit characteristic inclination-dependent double-peaked profiles. The profiles tend towards single peaks as the inclination angle<sup>3</sup> ( $i$ ) approaches  $0^\circ$  and the extent of the profile becomes dominated by the velocity dispersion internal to the disc, assuming the (typical) case that dispersion velocity is small compared to the rotation velocity. H I double-peaked profiles are often asymmetric, possible reasons for which include warped discs or otherwise asymmetric rotating gas distributions. Profiles consistent with discs are also seen in numerous high-redshift CO observations, a set of prime examples of which are presented by Daddi et al. (2010). Their CO ( $J = 2 \rightarrow 1$ ) observations with the PdBI at 92 GHz reveal significant discs of

<sup>3</sup> We follow the convention that  $i = 0^\circ$  for a face-on system.

cold molecular gas in a sample of star-forming galaxies at  $z \sim 1.5$ . The velocity profiles of the sample show single- and double-peaked structures, the latter exhibiting both symmetric and highly asymmetric profiles. Cases of highly irregular high-redshift CO line profiles have also been observed (e.g. Riechers et al. 2008), suggestive of highly disrupted systems with numerous components. Numerous studies present substantial evidence that on-going mergers are responsible for the observed CO line profiles of many systems (e.g. Frayer et al. 2008; Engel et al. 2010; Ivison et al. 2010a), including examples which see evidence of mergers between gas-rich discs (Ivison et al. 2011, 2013). Gravitational lensing plays a part in many high-redshift systems for which CO lines have been observed, and this can also skew the velocity profile due to differential lensing (Deane et al. 2013; Deane, Obreschkow & Heywood 2015).

The line structure in HATLAS J090426.9+015448, shown in the right hand panel of Fig. 1, exhibits an asymmetric double-peaked profile which covers a total velocity width of approximately  $1100 \text{ km s}^{-1}$ . This is consistent with the broadest CO line widths observed in SMGs, e.g. Harris et al. (2012). The vertical line on Fig. 1 is the redshift of the Ly  $\alpha$  line measured by Virdee (2013), which lies at the centre of the two CO peaks. If the Ly  $\alpha$  redshift is coincident with the central AGN, then a plausible explanation for the line structure is that the radio galaxy HATLAS J090426.9+015448 is surrounded by a large, asymmetric disc of molecular gas. The line profile is similar to that of SMM J02399–0136 as reported by Genzel et al. (2003), who interpreted it as a single, rapidly rotating disc. Ly  $\alpha$  lines can however be offset from cold gas tracer lines by approximately hundreds of  $\text{km s}^{-1}$  (e.g. Willott et al. 2015), and in a manner which is different from object to object due to the resonant absorption on the blue side of the Ly  $\alpha$  line. The systemic redshift of the AGN could well be centred on one of the CO peaks, and the double profile could be the result of two separate galaxies in the process of merging. SMM J02399–0136 was indeed later revealed to contain multiple merging objects (Ivison et al. 2010a). Higher resolution optical spectroscopy of the Ly  $\alpha$  line or preferably longer wavelength observations of non-resonant lines would be needed to more accurately measure the systemic AGN redshift.

Our second target HATLAS J140930.4+003803 also exhibits two distinct peaks in the CO line; however, in this case, the profile is much more symmetric. The centre of component 2 (as per the labelling on Fig. 2) is coincident with the measured Ly  $\alpha$  redshift. Again if the Ly  $\alpha$  line is coincident with the central AGN, the line profile can be explained by an on-going merger, with component 2 representing a molecular gas reservoir centred on the AGN and component 1 representing a second system. A second plausible explanation is that the CO line profile is the result of a single, regular, rotating disc of molecular gas, and either the Ly  $\alpha$  redshift is slightly offset from the systemic redshift of the AGN or the AGN itself is offset from the centre of the disc.

We can place upper limits on the dynamical masses in both systems, for both the merger and single disc scenario. The dynamical mass of a rotating disc ( $M_{\text{dyn}}$ ) in solar masses can be estimated as

$$M_{\text{dyn}} \sin^2(i) = 4 \times 10^4 \Delta v^2 R, \quad (1)$$

where  $i$  is the inclination angle,  $\Delta v$  is the FWHM of the line width in kilometre per second and  $R$  is the outer radius of the disc in kiloparsec (Neri et al. 2003). Attempts to deconvolve the PSF and estimate the angular extent of the sources result only in upper limits of 14 arcsec and 6 arcsec for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively. The assumed cosmological model and the redshifts of the sources result in angular diameter distances which translate to angular scales of  $8.37 \pm 0.25$

and  $8.56 \pm 0.25 \text{ kpc arcsec}^{-1}$  for the two targets, and the product of these and the angular size upper limits give upper limits to the value of  $R$ . We note that the typical extents of similar high- $z$  sources observed in the ground-state of CO are  $\sim 5\text{--}15 \text{ kpc}$  (Ivison et al. 2011; Riechers et al. 2011; Hodge et al. 2012), approximately an order of magnitude smaller than the constraints we are able to place on our targets using the current data. Crude dynamical mass upper limits can be estimated for each pair of components in each source by using the fitted velocities listed in Table 1. For HATLAS J090426.9+015448, we obtain values of  $M_{\text{dyn}} \sin^2(i) < (1.4 \pm 1.3) \times 10^{11}$  and  $< (8.1 \pm 1.8) \times 10^{11} M_{\odot}$  for components 1 and 2, respectively. The corresponding values for components 1 and 2 in HATLAS J140930.4+003803 are  $< (1.1 \pm 0.6) \times 10^{11}$  and  $< (2.1 \pm 1.4) \times 10^{11} M_{\odot}$ . The total dynamical mass limits are  $< 5.7 \times 10^{12} M_{\odot}$  and  $< 9.2 \times 10^{11} M_{\odot}$ , assuming velocity widths of 1100 and 665  $\text{km s}^{-1}$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively.

These (very conservative) dynamical mass limits for the individual components are consistent with existing measurements in other systems. One of the best examples of a dynamical study of a high- $z$  CO disc is the SMG GN20 (J123711.89+622211.8). Hodge et al. (2012) comfortably resolve a clumpy 14-kpc molecular disc in this system and derive a dynamical mass of  $(5.4 \pm 2.4) \times 10^{11} M_{\odot}$ . Wang et al. (2013) report values of  $M_{\text{dyn}} \sin^2(i)$  of the order of  $10^{10}\text{--}10^{11} M_{\odot}$  in their sample of  $z \sim 6$  quasars (QSOs). SMGs tend to have broader linewidths and therefore higher dynamical mass estimates than QSOs, and as noted by Coppin et al. (2010), this may be expected from orientation arguments when targeting optically selected QSOs under the assumption of a unified AGN model. However, there are overlaps and exceptions in both populations of sources. Invoking a unified model would imply that HzRGs should have a broader range of linewidths and dynamical masses, similar to the typical SMG population. Ultimately, it must be stressed that our dynamical mass limits are simply not that stringent due to the relatively coarse angular resolution afforded by the ATCA's H75 configuration.

The possibility that the CO spectra also represent an AGN-driven outflow should be considered. Velocities of  $\sim 1000 \text{ km s}^{-1}$  are seen in many systems and via many spectral line tracers, with kinematics suggestive of both bipolar and shell-like flows. These flows are thought to be driven by radio jets or shocks induced by radiation pressure originating close to the black hole, transporting dust and metal-rich material out to vast distances from the central AGN (Prochaska & Hennawi 2009; Ivison et al. 2012). Outflows typically manifest themselves as broad wings on the line profile, and require multiple component Gaussian fits (e.g. Nesvadba et al. 2008; Maiolino et al. 2012). No such features are apparent in our current data; however, higher angular resolution CO mapping coupled with matched resolution radio continuum imaging could investigate whether the molecular material in our two targets is being entrained by radio jets from the AGN (e.g. Emonts et al. 2014). We note again the tentative detection of two components in the EVN imaging of HATLAS J140930.4+003803 presented by Virdee (2013), possibly indicative of a jet.

#### 4.4 Luminosities, molecular gas masses and dust masses

The luminosities of the CO line in  $\text{K km s}^{-1} \text{ pc}^2$  can be determined by the following relationship, as derived by Solomon & Vanden Bout (2005):

$$L'_{\text{CO}} = 3.25 \times 10^7 I_{\text{CO}} v_{\text{centre}}^{-2} D_{\text{L}}^2 (1 + z_{\text{CO}})^{-3}, \quad (2)$$

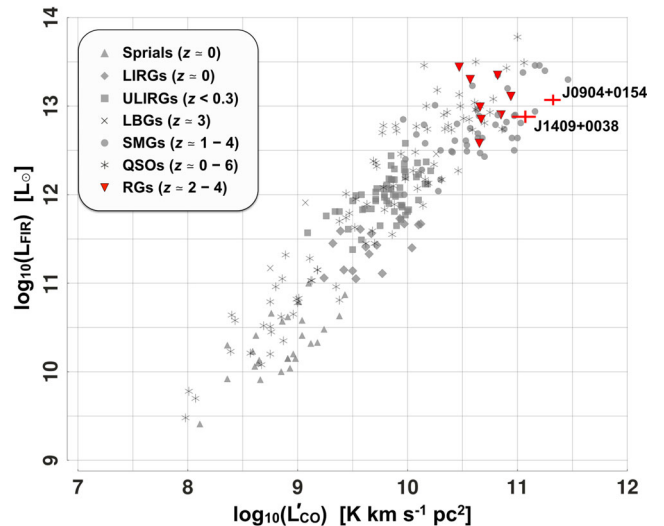


where  $I_{\text{CO}}$  is the velocity-integrated CO line flux in  $\text{Jy km s}^{-1}$ ,  $\nu_{\text{centre}}$  is the observing frequency in GHz,  $D_L$  is the luminosity distance in Mpc and  $z$  is the redshift of the line. CO line luminosities were derived on a per-component basis as listed in Table 1. We assume that the total CO line luminosity is the sum of the two components, also listed in the table.

Molecular ( $\text{H}_2$ ) gas masses are derived from  $L'_{\text{CO}}$  via the conversion factor  $\alpha_{\text{CO}}$ , which has units of  $M_{\odot} (\text{K km s}^{-1} \text{pc}^2)^{-1}$  (for a review, see Bolatto, Wolfire & Leroy 2013). Determining the value of  $\alpha_{\text{CO}}$  is an active area of research, and the choice of conversion factor is the principal source of uncertainty in high- $z$  molecular gas mass estimates. The traditional range of conversion factors ranges from 4, applicable to Giant Molecular Clouds (GMCs) in the Milky Way, to 0.8 in ultraluminous infrared galaxies (ULIRGs), the latter value traditionally also used for starbursting systems at high redshift (Downes & Solomon 1998). There are several theoretical models for describing  $\alpha_{\text{CO}}$ , suggesting dependences on gas temperature, dynamical state and metallicity (e.g. Narayanan et al. 2012). Appropriate observational constraints allow  $\alpha_{\text{CO}}$  to be estimated on a per-source basis, and this has been described by several studies using dynamical or radiative transfer modelling, resolved gas surface density measurements and estimates based on dust mass. Values all tend to lie within the aforementioned 0.8–4 range: Ivison et al. (2011) report values in the range 0.9–2.3 for a sample of SMGs, the colour-selected galaxies studied by Daddi et al. (2010) have conversion factors in the range  $3.6 \pm 0.8$  and Genzel et al. (2012) derive values of  $1.7 \pm 0.4$  for a sample of high- $z$  star-forming galaxies, corroborating the strong metallicity dependence predicted by theoretical modelling.

Adopting the standard SMG value of  $\alpha_{\text{CO}} = 0.8 M_{\odot} (\text{K km s}^{-1} \text{pc}^2)^{-1}$  results in the molecular gas masses ( $M_{\text{H}_2}$ ) derived from the CO luminosities listed in Table 1, for each component and the total for each galaxy. The total  $\text{H}_2$  masses place HATLAS J090426.9+015448 and HATLAS J140930.4+003803 amongst the most massive high- $z$  systems. The limits on  $M_{\text{dyn}}$  derived in Section 4.3 are not at odds with the  $M_{\text{H}_2}$  values, in which the latter does not exceed the former, although we emphasize again the loose constraints which we can place on  $M_{\text{dyn}}$  with the current data. The molecular gas in a typical high- $z$  SMG tends to be a significant fraction of the dynamical mass;<sup>4</sup> Tacconi et al. (2006, 2008) find fractions in the range 20–60 per cent. Higher resolution observations [with the Very Large Array (VLA), ALMA or an extended ATCA configuration] would provide tighter constraints on the dynamical masses of HATLAS J090426.9+015448 and HATLAS J140930.4+003803, as well determine whether each target was a single large disc or a merger of multiple components. Such observations would then unlock independent estimates of  $\alpha_{\text{CO}}$  (e.g. Ivison et al. 2013).

It is informative to compare  $L'_{\text{CO}}$  to the FIR luminosity  $L_{\text{FIR}}$ , the former being a proxy for the star formation potential of a galaxy and the latter being a good indicator of the SFR. We estimated the FIR properties of our sample as follows. First, we fit both galaxies using an isothermal modified blackbody model (e.g. Hildebrand 1983) with fixed emissivity index ( $\beta = 1.82$ , following Smith et al. 2013) and accounted for the *Herschel* response curves in the 100, 160, 250, 350 and 500  $\mu\text{m}$  bands. We fit the model to the photometry on a grid of temperatures between 10 and 60 K, recording the best-fitting luminosity each time. To estimate the associated uncertainties, we



**Figure 4.** FIR luminosity ( $L_{\text{FIR}}$ ) against CO luminosity ( $L'_{\text{CO}}$ ) for HATLAS J090426.9+015448 and HATLAS J140930.4+003803 (labelled crosses showing the uncertainties in the two parameters), in the context of a range of sources. A sample of CO-detected radio galaxies (RGs) is shown, as well as spiral galaxies, luminous infrared galaxies (LIRGs), ultraluminous infrared galaxies (ULIRGs), Lyman-break galaxies (LBGs), sub-millimetre galaxies (SMGs) and quasars (QSOs). The diagonal line shows the fit to the SMG and ULIRG population by Bothwell et al. (2013). Measurements in this plot are from various studies as collated by Heywood et al. (2013), and include additional HzRG points from Emonts et al. (2014).

created 500 Monte Carlo realizations of each galaxy by varying the photometry within the errors, and used half the difference between the 16th and 84th percentiles of the resulting luminosity distribution for each galaxy to estimate the associated uncertainties.

HATLAS J090426.9+015448 and HATLAS J140930.4+003803 are shown on the  $L'_{\text{CO}}-L_{\text{FIR}}$  diagram in Fig. 4, showing the  $L_{\text{FIR}}$  values,  $(11.7 \pm 1.5) \times 10^{12} L_{\odot}$  and  $(7.6 \pm 1.2) \times 10^{12} L_{\odot}$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively, as listed in Table 1. The targets are marked by crosses on Fig. 4, the extent of which show the  $1\sigma$  error bars in  $L'_{\text{CO}}$  and  $L_{\text{FIR}}$ . The two sources are placed in the context of existing observations of radio galaxies (inverted triangles) as well as populations of other source types as indicated in the figure legend. These values are derived from various studies as collated by Heywood et al. (2013), with additional HzRG points from Emonts et al. (2014). The diagonal line shows the fit to the SMG and ULIRG population by Bothwell et al. (2013). Note that some authors have claimed that this distribution features two distinct populations of galaxies, namely a ‘main sequence’ group and a starbursting group, although others have strongly disputed this (Ivison et al. 2011). We refer the reader to Carilli & Walter (2013) for further discussion.

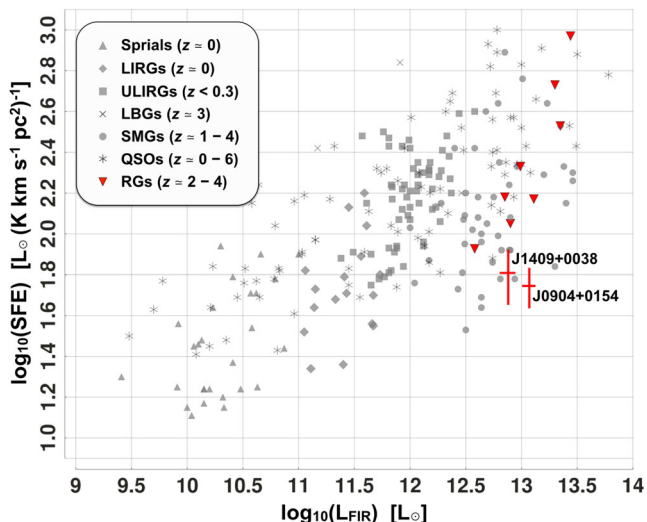
It is clear from their position on this diagram that for their given FIR luminosities HATLAS J090426.9+015448 and HATLAS J140930.4+003803 have very high CO luminosities both in terms of the general distribution, but in particular for the radio galaxy population. The star formation efficiencies (SFEs) for the targets, defined as

$$\text{SFE} = \frac{L_{\text{FIR}}}{L'_{\text{CO}}}, \quad (3)$$

are  $1.74 \pm 0.09$  and  $1.81 \pm 0.12 L_{\odot} (\text{K km s}^{-1} \text{pc}^2)^{-1}$ . This places them at the lower end of the range for SMGs, QSOs and HzRGs in

<sup>4</sup> Determining the ratio  $M_{\text{H}_2}/M_{\text{dyn}} > 1$  was early evidence that the GMC conversion factor was inappropriate for high- $z$  systems.





**Figure 5.** Star formation efficiency ( $L_{\text{FIR}}/L'_{\text{CO}}$ ) against FIR luminosity ( $L_{\text{FIR}}$ ) for HATLAS J090426.9+015448 and HATLAS J140930.4+003803 (labelled crosses showing the uncertainties in the two parameters). Additional points are as per Fig. 4.

the distribution shown in Fig. 5, which shows the SFE against  $L_{\text{FIR}}$  for the same galaxy populations plotted on Fig. 4.

The FIR modelling can also be used to provide estimates of the SFR, albeit with an unknown and unmodelled contribution to the FIR luminosity provided by AGN heating. Using the Kennicutt (1998) scaling and a Chabrier (2003) initial mass function yields  $1250 \pm 160 M_{\odot} \text{ yr}^{-1}$  and  $810 \pm 130 M_{\odot} \text{ yr}^{-1}$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively, as listed in Table 1.

The slope of the far-infrared/radio correlation (FIRC) is captured by the dimensionless luminosity ratio

$$q_{\text{IR}} = \log_{10} \left[ \frac{(S_{\text{IR}}/3.75 \times 10^{12})}{S_{1.4\text{GHz}}} \right], \quad (4)$$

where  $S_{\text{IR}}$  is the FIR flux integrated over rest-frame wavelengths of 8–1000  $\mu\text{m}$  in units of  $\text{W m}^{-2}$  and normalized to a frequency of  $3.75 \times 10^{12} \text{ Hz}$  and  $S_{1.4\text{GHz}}$  is the monochromatic,  $k$ -corrected radio flux density in units of  $\text{W m}^{-2} \text{ Hz}^{-1}$ . This definition follows e.g. Ivison et al. (2010b), who report a median  $q_{\text{IR}}$  value of  $2.44 \pm 0.02$  for a sample of radio-selected galaxies out to  $z \sim 3$ . Similar values are reported in numerous studies for various selection methods (e.g. Jarvis et al. 2010; Bourne et al. 2011; Mao et al. 2011; Magnelli et al. 2015), often reporting moderate evolution of the FIRC with redshift. We compute  $q_{\text{IR}}$  for our two targets using the properties listed in Table 1, and  $k$ -correcting  $S_{1.4\text{GHz}}$  using spectral indices derived from the 327 MHz and 1.4 GHz measurements, and the mean of the CO redshifts. This leads to  $q_{\text{IR}} = 1.92 \pm 0.13$  and  $0.81 \pm 0.10$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively. These lie somewhat below the typical FIRC, particularly HATLAS J140930.4+003803, likely due to an excess of radio emission which is driven by the AGN rather than by star formation.

Since the isothermal models employed in the FIR modelling above do not include a contribution to the dust luminosity from very small grains which dominate at mid-infrared wavelengths, we estimate the dust luminosities using an M82 template from Polletta et al. (2007), which is more luminous by around 0.36 dex than the best-fitting isothermal estimates. We therefore apply this correction

to the isothermal dust luminosities and derive  $L_{\text{dust}} = (11.5 \pm 1.6) \times 10^{12} L_{\odot}$  for HATLAS J090426.9+015448 and  $(7.9 \pm 1.1) \times 10^{12} L_{\odot}$  for HATLAS J140930.4+003803. For the purpose of calculating dust masses, the dust emissivity is normalized to  $\kappa_{850 \mu\text{m}} = 0.77 \text{ g}^{-1} \text{ cm}^2$  (Dunne et al. 2000), and we use the best-fitting isothermal model to each of the 500 realizations of the galaxy photometry to estimate the dust mass, with uncertainties again derived according to half the difference between the 16th and 84th percentiles of the resulting distribution. We estimate the total dust mass in each system to be  $(2.3 \pm 1.1) \times 10^9 M_{\odot}$  and  $(1.3 \pm 0.7) \times 10^9 M_{\odot}$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively.

These dust masses are high, comparable to the highest dust masses reported in SMGs (Michałowski, Hjorth & Watson 2010; Rowlands et al. 2014). The typical SMG value for  $\alpha_{\text{CO}}$  is justifiably applicable to our targets. Combining the SFR values listed in Table 1 with the molecular gas masses derived above yields gas depletion time-scales of  $135 \pm 29 \text{ Myr}$  and  $117 \pm 35 \text{ Myr}$  for HATLAS J090426.9+015448 and HATLAS J140930.4+003803, respectively, assuming that the SFR is sustained. This places the targets at the upper end of the depletion time-scale distribution for QSOs, HzRGs and SMGs (e.g. Tacconi et al. 2008; Emonts et al. 2014; Jones et al. 2016).

## 5 CONCLUSIONS

We have used the ATCA to observe the CO ground-state ( $J = 1 \rightarrow 0$ ) spectra of two HzRGs. In the context of existing studies of HzRGs, our targets have relatively low 1.4-GHz radio luminosities, are close to the break in the luminosity function and therefore represent typical systems which are undergoing high rates of accretion at high redshift. Despite this, our targets are amongst the most CO-luminous systems yet observed, with CO luminosities in excess of  $10^{11} \text{ K km s}^{-1} \text{ pc}^2$ , and total inferred molecular gas masses of  $\sim 10^{11} M_{\odot}$ . The targets have high FIR luminosities, with detections in all five *Herschel* bands. SED fitting to the *Herschel* photometry yields FIR luminosities of  $\sim 10^{13} L_{\odot}$ , and dust masses of  $\sim 10^9 M_{\odot}$ . These systems not only harbour an AGN, but share many characteristics of SMGs, with large amounts of dust obscuring a star formation event which is producing stars at a rate of  $\sim 1000 M_{\odot} \text{ yr}^{-1}$ . If sustained, this would deplete the molecular gas reservoirs in  $\sim 100 \text{ Myr}$ . The resolved line spectra both show double-peaked profiles with total velocity ranges of  $\sim 1000 \text{ km s}^{-1}$ , indicative of either an on-going merger, or a very large rotating disc of molecular material surrounding the AGN. HzRGs are thought to be the progenitors of the most massive present-day elliptical galaxies, and the merger scenario is consistent with the evolutionary picture for such objects. Existing observations with superior angular resolution tend to resolve similar systems into multiple components, and this, along with the upper limits to the dynamical masses implied by a single-disc scenario, suggests that the merger interpretation is most likely to be the correct one. Molecular line observations with higher angular resolution (ideally coupled with high-resolution radio continuum maps to determine the orientation of the molecular gas with respect to any radio jets) would reveal the true dynamical state of these systems. An extended ATCA configuration, the VLA or ALMA (or some combination thereof) are all viable choices of instrument for such observations. The radio continuum spectra of our targets are intriguing, peaking somewhere below 1 GHz (rest frame) with a break in the spectrum at high radio frequencies which gives rise to a very steep spectral decline. The spectral shapes are broadly consistent

with the objects being faint CSS sources; however, the extreme steepness of the high-frequency radio spectrum may be caused by insufficient angular resolution of the existing radio continuum observations.

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