

# Theoretical and Applied Climatology

## MULTIFRACTALITY AND AUTOREGRESSIVE PROCESSES OF DRY SPELL LENGTHS IN EUROPE: AN APPROACH TO THEIR COMPLEXITY AND PREDICTABILITY

--Manuscript Draft--

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<b>Abstract:</b>	<p>Dry spell lengths, DSL, defined as the number of consecutive days with daily rain amounts below a given threshold, may provide relevant information about drought regimes. Taking advantage of a daily pluviometric database covering a great extension of Europe, a detailed analysis of the multifractality of the dry spell regimes is achieved. At the same time, an autoregressive process is applied with the aim of predicting DSL. A set of parameters, namely Hurst exponent, <math>H</math>, estimated from multifractal spectrum, <math>\alpha</math>, critical Hölder exponent, <math>\beta</math>, for which <math>H</math> reaches its maximum value, spectral width, <math>W</math>, and spectral asymmetry, <math>B</math>, permits a first clustering of European rain gauges in terms of the complexity of their DSL series. This set of parameters also allows distinguishing between time series describing fine- or smooth-structure of the DSL regime by using the Complexity Index, CI. Results of previous monofractal analyses also permits establishing comparisons between smooth-structures, relatively low correlation dimensions, notable predictive instability and anti-persistence of DSL for European areas, sometimes submitted to long droughts. Relationships are also found between the CI and the mean absolute deviation, MAD, and the optimum autoregressive order, OAO, of an ARIMA(p,d,0) autoregressive process applied to the DSL series. The detailed analysis of the discrepancies between empiric and predicted DSL underlines the uncertainty over predictability of long DSL, particularly for the Mediterranean region.</p>
<b>Response to Reviewers:</b>	<p>Manuscript TAAC-D-14-00581: Multifractality and autoregressive processes of dry spell lengths in Europe: An approach to their complexity and predictability (Authors: X. Lana, A. Burgueño, C. Serra and M.D. Martínez)</p> <p>Response to Reviewer 1:</p> <p>a) In agreement with recommendations of the Reviewer and the Editor, the Section 3.1 has been notably shortened, being only introduced the basic concepts concerning the</p>

MF-DFA algorithm and the appropriate references (see also lines 162-165). The structure of Section 3.2 (Singularity spectrum) has not been changed, as many definitions of parameters and concepts used after are summarised in this Section. Additionally, in agreement with Reviewer 2, a detailed discussion about the range of qth-order has been added to this Section.

b)Effectively, the interpolation method has been the “inverse distance”. We think that in our case the main problem of obtaining accurate plots of the spatial distribution of fractal parameters is the limited rain gauge density in some areas, as mentioned now in page 5 (lines 118-124).

c)In Section 5, lines 575-580, several alternatives to the AR(p) process are cited as possible improvements on DSL prediction for very long DSL. Certainly, an ARIMA(p,d,q) modelling and others methods based on the Poisson distribution and Monte Carlo algorithms would improve the mentioned very long DSL prediction, but we think they are beyond the scope of the present paper. It is also worth of mention that the relatively simple AR(p) process has led to good results when predicting monthly Western Mediterranean Oscillation index, as the authors of this manuscript have found (manuscript nowadays submitted to the International Journal of Climatology). With respect to these questions, we have to mention that in line 224-225 a mistake concerning the definition of AR(p) has been amended. AR(p) has to be properly defined as an ARIMA(p,d,q) with d=1 and q=0.

Response to Reviewer 2:

a)With the aim of a more complete description of the DSL series, new Figure 1c includes a histogram of the number of DSL, NDSL, for the 267 DSL series. Additionally, two very different examples of DSL series (Vaexjoe, Sweden, and Almeria, Spain) are shown in a new Figure 1d. The corresponding comments are developed in Section 2, lines 136-143.

b)A discussion about the appropriate range of the qth-order is developed in Section 3.2, lines 203-217. The paper suggested by the reviewer (\*), the assumption that multifractal spectra should be fitted to a quadratic function taking as argument the Hölder exponent, and previous experience of the authors about this question have been used as reference points of this discussion.

(\*) Ivanov P. Ch., Nunes Amaral, L.A., Goldenberg A.L., Haulin S., Rosenblum M.G., Stanley, H.E., Struzik, Z.B. (2001). From 1/f noise to multifractal cascades in hearthbeat dynamics. Chaos, 11, 641-652.

c)The process to obtain a configuration of 14 clusters is described with more detail in Section 4.3, lines 340-356. The explanation is partially based on the concept of similarity index, Lij, which is defined and quantified by the new Equation 9.

d)With respect to Figures 8, due to a technical problem with the writing software, the two first figures were plotted without appearing “C” codes designing the cluster number. The problem has been now solved.

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**MULTIFRACTALITY AND AUTOREGRESSIVE PROCESSES  
OF DRY SPELL LENGTHS IN EUROPE: AN APPROACH TO THEIR  
COMPLEXITY AND PREDICTABILITY**

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20 **Abstract**

21 Dry spell lengths, DSL, defined as the number of consecutive days with daily rain amounts  
22 below a given threshold, may provide relevant information about drought regimes. Taking  
23 advantage of a daily pluviometric database covering a great extension of Europe, a detailed  
24 analysis of the multifractality of the dry spell regimes is achieved. At the same time, an  
25 autoregressive process is applied with the aim of predicting DSL. A set of parameters,  
26 namely Hurst exponent,  $H$ , estimated from multifractal spectrum,  $f(\alpha)$ , critical Hölder  
27 exponent,  $\alpha_0$ , for which  $f(\alpha)$  reaches its maximum value, spectral width,  $W$ , and spectral  
28 asymmetry,  $B$ , permits a first clustering of European rain gauges in terms of the complexity  
29 of their DSL series. This set of parameters also allows distinguishing between time series  
30 describing fine- or smooth-structure of the DSL regime by using the Complexity Index, CI.  
31 Results of previous monofractal analyses also permits establishing comparisons between  
32 smooth-structures, relatively low correlation dimensions, notable predictive instability and  
33 anti-persistence of DSL for European areas, sometimes submitted to long droughts.  
34 Relationships are also found between the CI and the mean absolute deviation, MAD, and the  
35 optimum autoregressive order, OAO, of an ARIMA(p,d,0) autoregressive process applied to  
36 the DSL series. The detailed analysis of the discrepancies between empiric and predicted DSL  
37 underlines the uncertainty over predictability of long DSL, particularly for the Mediterranean  
38 region.

39  
40 **Keywords:** DSL series, drought regime, multifractal detrended fluctuation analysis, Hölder  
41 and Hurst exponents, ARIMA process, European pluviometric network.

## 45 1. Introduction

46 Dry spell lengths, DSL, defined as the number of consecutive days with daily rain amounts  
47 lowering a certain threshold, represent a valuable magnitude to analyse several aspects of  
48 drought episodes. Besides studies based on “the standardised precipitation index”, SPI  
49 (McKee et al, 1993, 1995; Hayes et al., 1999; Lana et al, 2001), or “dry days since last rainy  
50 day”, DDSLR (Aviad et al. 2009; Reiser and Kutiel 2010; Lana et al., 2012), DSL have also been  
51 used along last years to characterise several patterns of the drought regime in Europe  
52 (Brunetti et al., 2002; Anagnostopoulou et al., 2003; Kostopoulou and Jones, 2005; Schmidli  
53 and Frei, 2005; Cindrić et al., 2010; Carvalho et al., 2013). The analysis of DSL series would  
54 permit to improve the knowledge about physical mechanisms and time trends governing  
55 drought regimes, particularly concerning southern Europe, as severe drought episodes  
56 frequently occur in the Mediterranean region (Lana et al., 2006, 2008b; Livada and  
57 Assimakopoulos, 2007; Diodato and Bellocchi, 2008; Nastos and Zerefos, 2009; García-Ruiz  
58 et al., 2011; Heinrich and Gobiet, 2011; Zolina et al., 2013). Specifically, the statistical  
59 distribution of DSL in Europe, with special emphasis on expected DSL for several return  
60 periods, and the assignment of statistical parent distributions to spatial clusters of rain  
61 gauges have been recently analysed (Serra et al, 2013, 2014). European partial duration  
62 series of DSL have also been recently studied (Serra et al., 2015).

63 A different approach may be the application of fractal theory. First, the rescaled range  
64 analysis (Turcotte, 1997) leads to obtain the Hurst exponent, which permits characterising  
65 the persistence, anti-persistence or randomness of DSL. Second, the self-affine character of  
66 DSL (Hausdorff exponent) and the possibility of modelling DSL series by fractional Gaussian  
67 noise series can be tested (Mandelbrot and Wallis, 1969; Malamud and Turcotte, 1999). And  
68 third, in agreement with the reconstruction theorem (Diks, 1999), it is possible to quantify  
69 several aspects of the physical mechanism governing DSL, like the complexity and chaotic  
70 behaviour (correlation and strange attractor dimensions), loss of memory (Kolmogorov  
71 entropy) and predictive instability (Lyapunov exponents and Kaplan-Yorke dimension). A  
72 complete analysis of the complexity and predictive instability of European DSL series, based  
73 on the rescaled range analysis and the reconstruction theorem, can be found in Lana et al.  
74 (2010).

75  
76 The multifractal analysis of the DSL series (Kantelhardt et al., 2002) is now proposed through  
77 four parameters: the Hurst exponent,  $H$ , estimated from the multifractal spectral curve,

78  $f(\alpha)$ , the range,  $W$ , of the Hölder exponent,  $\alpha$ , the asymmetry,  $B$ , of  $f(\alpha)$ , and the critical  
79 Hölder exponent,  $\alpha_0$ , for which the maximum of  $f(\alpha)$  is reached. These parameters permit  
80 to characterise the complexity of every DSL series by using the Complexity Index, CI. After  
81 applying the Principal Component Analysis (Jolliffe, 1986; Preisendorfer, 1988) and clustering  
82 algorithms (Kalkstein et al., 1987; Davis and Kalkstein, 1990), DSL series may be spatially  
83 grouped in terms of their degrees of complexity, taking as variables for the classification  
84 those four multifractal parameters. In this way, regions with simple (smooth-structure) or  
85 complex (fine-structure) drought predictability can be distinguished.

86  
87 The description and quantification of the predictability of DSL are complemented by the  
88 ARIMA(p,d,0) autoregressive process (Box and Jenkins, 1976), previously applied in  
89 Climatology, for instance, to the North Atlantic Oscillation, NAO, series (Stephenson et al.,  
90 2000; Mills, 2004). The mean absolute deviation, MAD, permits to quantify the goodness of  
91 fit between empiric DSL and those generated by the ARIMA(p,d,0) process, while the  
92 optimum autoregressive order, OAO, points out the order with minimum MAD. Additionally,  
93 valuable information can be obtained by analysing the discrepancies (residuals) between  
94 empiric and predicted DSL for every analysed rain gauge.

95  
96 The main objectives of this paper are, first of all, to extend the multifractal analysis of DSL  
97 series to Europe. Specifically, areas of simple or complex predictability are well bounded and  
98 comparisons are made with previous results derived from the application of the  
99 reconstruction theorem and the rescaled range analysis (Lana et al., 2010). And second, to  
100 establish relationships between CI derived from multifractal theory and MAD and OAO  
101 values based on ARIMA(p,d,0) autoregressive algorithms, and to highlight advantages and  
102 shortcomings of predictability based on this autoregressive process.

103 The contents of the paper are organised as follows. The database is introduced in Section 2;  
104 the methodology to obtain the multifractal spectrum, the Complexity Index, CI, the spatial  
105 distribution of clusters and a description of the ARIMA(p,d,0) process are detailed in Section  
106 3; results concerning multifractal spectrum parameters, complexity measures, clustering and  
107 autoregressive predictions are developed in Section 4. Finally, Conclusions are summarised  
108 in Section 5.

## 110 2. Database

111 Daily precipitation data for the years 1951-2000 have been compiled from 267 rain gauges in  
112 Europe and neighbouring countries. Most of these series (236) come from the European  
113 Climate Assessment and Dataset (ECA&D), <http://eca.knmi.nl/>. All these series are public  
114 and non-blended and their quality has been analysed by ECA&D (Klein Tank et al., 2002;  
115 Wijngaard et al. 2003; Klok and Klein Tank, 2009). The rest of series come from the *Agencia*  
116 *Estatad de Meteorología* (Spanish Ministry of Environment) and standard homogeneity and  
117 quality controls have been previously applied (Lana et al., 2008a).

118 Figure 1a depicts the spatial distribution of the available stations. Dense rain gauge coverage  
119 is observed in Western Europe, except for Italy, Great Britain and the north of the  
120 Scandinavian Peninsula, where it is not so dense. Results for Turkey and Israel are also  
121 included in the maps, but they only provide a broad approach for these regions.  
122 Interpolation of scarce spatial data could generate computational artefacts leading to some  
123 unrealistic spatial patterns, even taking into account additional data as height above sea  
124 level or orographic effects. All rain gauges have a minimum continuous recording period of  
125 40 years, and the series are complete (50 years) for 102 out of 267 stations. Figure 1b shows,  
126 year by year, the number of stations with complete recordings. Most of data series are  
127 continuous for the period 1955-1990, the number of available records diminishing at the  
128 beginning (1951-1955), and especially at the end (1990-2000) of the recording period. If a  
129 DSL is likely to be incomplete due to lack of record continuity, it is rejected. This does not  
130 constitute a relevant shortcoming on account of the large enough statistical population of  
131 DSL.

132 Thresholds commonly used to define a DSL are 0.1, 1.0, 5.0 and 10.0 mm/day (Kutiel and  
133 Maheras, 1992; Martín Vide and Gómez, 1999; Anagnostopoulou et al. 2003; Serra et al.  
134 2006, 2013, 2014; Lana et al 2008b; Cindrić et al., 2010). The current study is constrained to  
135 the threshold of 0.1 mm/day (the assumed resolution of pluviometers), with daily excess  
136 (shortage) defining a wet (dry) day, thus being obtained a long enough data basis. As a global  
137 description of the 267 DSL series, their numbers of dry spells,  $N_{DSL}$ , range from 1061 up to  
138 3203, with an average of 2621 and a standard deviation of 477. The histogram of  $N_{DSL}$  (Figure  
139 1c) clearly depicts a skewed distribution toward high  $N_{DSL}$ . Two examples of DSL series from  
140 northern and southern Europe are shown in Figure 1d. As expected, the southern Europe  
141 example is characterised by a relatively low  $N_{DSL}$  and a non negligible number of long DSLs.

142 The opposite example could be that corresponding to higher latitudes, with larger  $N_{DSL}$  and  
143 shorter DSLs.

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3 144 The relevance of the DSL concept and its relationship with droughts is manifested, for  
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5 145 example, by searching for the longest DSL,  $L_{max}$ , obtained for every rain gauge along the  
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7 146 recording period. In agreement with [Serra et al. \(2014\)](#), values of  $L_{max}$  exceeding three  
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9 147 months are detected in the southern Mediterranean coast, for latitudes south of  $40^{\circ}N$ . The  
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11 148 high spatial gradient of  $L_{max}$  obtained in the Mediterranean region contrasts with the  
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13 149 notable homogeneity and low  $L_{max}$  values at northern latitudes.

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## 151 3. Methodology

### 152 3.1 Multifractality.

153 Multifractals are complex self-similar objects that consist on differently weighted fractals  
154 with different non-integer dimensions. Thus, the fundamental characteristic of  
155 multifractality is that scaling properties may vary in different regions of the system (Dutta,  
156 2010; Ghosh et al., 2012). If scaling properties are kept, whatever the region, the signal is  
157 known as monofractal. The multifractal detrended fluctuation analysis (MF-DFA) has been  
158 introduced as a reliable characterization of multifractal non-stationary and stationary time  
159 series (Kantelhardt et al., 2002). It is based on the identification of the scaling of the  $q^{th}$ -  
160 order moments depending on the signal length. The MF-DFA surpasses in quality and  
161 simplicity to previous algorithms such as the multifractal box counting (MF-BOX) (Feder,  
162 1988) or the wavelet transform modulus maxima (WTMM) (Muzy et al., 1994) algorithms. A  
163 detailed description of the steps leading to the right application of the MF-DFA algorithm  
164 and the concepts of  $q^{th}$ -order fluctuation function,  $F_q(s)$ , also named standard partition  
165 function, can be found in Burgueño et al. (2013).

166 A variety of continuous multifractal time signals have been analyzed by the MF-DFA,  
167 including climatological series as hourly or daily wind speeds (Kavasseri and Nagarajan, 2005;  
168 Feng et al, 2009), daily temperatures (Lin and Fu, 2008; Yuan et al., 2013; Burgueño et al.,  
169 2013) and global monthly temperature anomalies (Mali, 2014), as well as not so typical  
170 magnitudes as lightning initiation process (Gou et al., 2010).

### 171 3.2 The singularity spectrum

172 According to Kantelhardt et al. (2002), the  $q^{th}$ -order fluctuation function,  $F_q(s)$ , follows a  
173 power law with exponent  $h(q)$ , which is known as the generalized Hurst exponent. The  
174 singularity spectrum,  $f(\alpha)$ , can be related to  $h(q)$  via a Legendre transform:

$$175 \alpha = h(q) + q \frac{d[h(q)]}{dq} \stackrel{\text{Legendre}}{\leftrightarrow} f(\alpha) = q[\alpha - h(q)] + 1 \quad (1)$$

176 where  $\alpha$  is the singularity strength or Hölder exponent, while  $f(\alpha)$  denotes the dimension  
177 of the subset of the series. The multifractal scaling exponent is

$$178 \tau(q) = qh(q) - 1 \quad (2)$$

179  $\alpha$  being expressed as

$$\alpha = \frac{d\tau(q)}{dq} \quad (3)$$

The characteristics of the singularity spectrum  $f(\alpha)$  provide a new way of comparing signals, because it describes the dimensions of subsets of the series characterized by the same singularity strength  $\alpha$ . Designing  $\alpha_0$  as the singularity strength with maximum spectrum, a small value of  $\alpha_0$  means that the underlying process loses fine-structure, that is, it becomes more regular in appearance; conversely, a large value of  $\alpha_0$  ensures more complexity. In this sense, the Hurst exponent depicts a linear relationship with  $\alpha_0$ , as proved by the results. The shape of  $f(\alpha)$  may be fitted to a quadratic function around the position  $\alpha_0$ ,

$$f(\alpha) = A(\alpha - \alpha_0)^2 + B(\alpha - \alpha_0) + C \quad (4)$$

where  $C$  is an additive constant equal to 1. Coefficient  $B$  indicates the asymmetry of the spectrum, being zero for a symmetric spectrum. A right-skewed spectrum,  $B > 0$ , indicates relatively strongly weighted high fractal exponents (with "fine-structure"), while left-skewed shapes,  $B < 0$ , point to lower ones (more regular or smooth-structure). Spectral width,  $W$ , is defined as

$$W = \alpha_2 - \alpha_1 \quad (5)$$

with  $f(\alpha_1) = f(\alpha_2) = 0$ , being  $\alpha_2$  larger than  $\alpha_1$ , and the wider the spectral content, the stronger is the multifractality. In other words, the wider the range of possible fractal exponents, the "richer" is the process in structure. A signal with a high value of  $\alpha_0$ , a wide range  $W$  of fractal exponents, and a right-skewed shape,  $B > 0$ , is more "complex" than one with the opposite characteristics (Shimizu et al., 2002). For monofractal series, the width of the spectrum would be zero and  $h(q)$  would be independent of  $q$ . Hence, from Equation (1), there will be a unique value for both  $\alpha$  and  $f(\alpha)$ ,  $\alpha$  being the Hurst exponent,  $H$ , and  $f(\alpha)$  being equal to 1.

With respect to the appropriate range of the  $q^{\text{th}}$ -order, two questions have to be considered. First, an accurate revision of the goodness of fit of the  $q^{\text{th}}$ -order fluctuation functions,  $F_q(s)$ , to a power law is recommended. And second, a relevant property of the Legendre transform has to be considered. From a pure analytical point of view, the theoretical maximum ( $\alpha_2$ ) and minimum ( $\alpha_1$ ) Hölder exponents for which the multifractal spectrum is zero correspond

208 to  $q \rightarrow -\infty$  and  $q \rightarrow +\infty$  respectively. In consequence, the expected range of the  $q^{\text{th}}$ -orders  
 209 should be  $(-\infty, +\infty)$ . Computational instabilities for very high (positive and negative)  $q^{\text{th}}$ -  
 210 orders may lead to departures from a power law for the  $q^{\text{th}}$ -order fluctuation function.  
 211 Additionally, [Burgueño et al. \(2013\)](#) detected significant departures from the expected  
 212 quadratic function (Equation 4) modeling the multifractal spectrum when high  $q^{\text{th}}$ -orders  
 213 were applied. Considerations made by [Ivanov et al. \(2001\)](#), who assume that the appropriate  
 214 range of  $q^{\text{th}}$ -orders depend on the series length, would be in agreement with the detection  
 215 of  $q^{\text{th}}$ -orders for which the mentioned power law is not well accomplished. Then, the  
 216 verification of spectral contents satisfying Equation 4 is suggested as an additional test to  
 217 decide the optimum  $q^{\text{th}}$ -order range.

### 218 3.3 Autoregressive process

219 The autoregressive integrated moving average ARIMA(p,d,0) model ([Box and Jenkins, 1976](#))  
 220 assumes that

$$221 \Delta^d x(i) = \theta + \mu x(i-1) + \sum_{k=1}^p \delta_k \Delta^d x(i-k) + a_i \quad (i = p+2, \dots, N) \quad (6a)$$

222 Where  $\{x\}$  is a set of  $N$  empirical data,  $\Delta x$  is the set of first differences  $\Delta x(i) = [x(i+1) - x(i)]$ ,  
 223 with  $\Delta^d x(i-k) = [x(i-k+1) - x(i-k)]^d$ ,  $\{\theta, \mu, \delta_1, \dots, \delta_p\}$  are the parameters of the  
 224 autoregressive process,  $\{a\}$  is a noise series and  $d$  is a real number. **Alternatively, the**  
 225 **ARIMA(p,d,0), with  $d = 1.0$ , can be written as**

$$227 x(i) = \theta + \sum_{k=1}^p \delta_k x(i-k) + a_i \quad , i = p+1, \dots, N \quad (6b)$$

228 where time series  $\{x\}$  are directly used instead of first differences. With the aim of avoiding  
 229 singularities in the linear system of equations used to estimate  $\{\theta, \mu, \delta_1, \dots, \delta_p\}$ , parameter  $\mu$   
 230 is implicitly included in parameter  $\delta_1$ . The corresponding [Equation \(6b\)](#) is usually designed as  
 231 autoregression, AR(p). The resulting system of linear equations, disregarding the stochastic  
 232 component  $\{a\}$ , can be represented in terms of matrix formulation by

$$233 Z = AW \quad (7a)$$

234 with  $Z$  the  $\{x(p+1), x(p+2), \dots, x(n)\}$  vector,  $n$  the number of empiric elements belonging to  
 235 series  $\{x\}$ , the  $(n-p-1, p+1)$  matrix  $A$  with elements multiplying parameters  $\{\theta, \delta_1, \dots, \delta_p\}$  and a  
 236  $p+1$  dimension vector  $W$  containing the parameters to be solved from the linear system of

237 Equation (7a). The components of vector  $W$  can be estimated by multiplying Equation (7a)  
238 by the transposed  $A$  matrix,  $A^T$

$$239 \quad A^T Z = A^T A W \quad (7b)$$

240 and remembering that the symmetric matrix  $A^T A$  can be decomposed in two triangular  
241 matrices. Then, it is straightforward to obtain the values of parameters  $\{\theta, \delta_1, \dots, \delta_p\}$  taking  
242 advantage of the Crout's algorithm (Press et al., 1992).

243  
244 A convincing solution of Equation (6b) demands some criterion to decide the optimum  
245 autoregression order, OAO. The decision can be taken by searching for the OAO leading to a  
246 minimum of a convenient goodness of fit index. The selected index is the mean absolute  
247 deviation, MAD (Stephenson et al., 2000),

$$248 \quad MAD = \frac{1.483}{(N-p-1)} \sum_{k=1}^{N-p-1} |x(k) - x^*(k)| \quad (8)$$

249 being  $x(k)$  empiric DSL and  $x^*(k)$  those predicted by parameters derived from the  
250 ARIMA(p,1,0) process.

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## 252 4. Results

### 253 4.1 A representative example of multifractal spectrum

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2 254 As an example, DSL multifractal spectrum parameters for the Fabra Observatory (NE Iberian  
3  
4 255 Peninsula) are shown in [Figures 2a-2d](#). It should be noted that their general features are  
5  
6 256 common to the rest of DSL series obtained from the European pluviometric network records.  
7  
8 257 Nevertheless, it is important to take into account that the different parameter values for  
9  
10 258 every DSL series characterize their specific complexity. [Figure 2a](#) depicts three selected  $q^{\text{th}}$ -  
11  
12 259 order fluctuation functions in a log-log scale. In spite of fluctuations, their evolution with  
13  
14 260 length  $s$  is well fitted by a power-law. [Figure 2b](#) describes the evolution of the generalized  
15  
16 261 Hurst exponents,  $h(q)$ , with  $q$ . Remembering that, for stationary series,  $h(q=2)$  is the Hurst  
17  
18 262 exponent  $H$ , this DSL series is characterized by a strong randomness ( $H$  very close to 0.5). As  
19  
20 263 a general rule, also observed for daily extreme temperature series ([Burgueño et al., 2013](#)), a  
21  
22 264 second order polynomial on  $q$  fits well  $h(q)$ . The multiscaling fractal exponent,  $\tau(q)$ , and  
23  
24 265 Hölder exponent,  $\alpha(q)$ , are shown in [Figures 2c](#) and [2d](#), being worth mentioning signs of  
25  
26 266 linear behavior of  $\tau$  and the clear linear evolution of  $\alpha(q)$ , with  $q$  within the  $\pm 15$  range. The  
27  
28 267 corresponding multifractal spectrum,  $f(\alpha)$ , is shown in [Figure 3](#). It should be underlined  
29  
30 268 that the empiric normalized value of 1.0 is reached for  $\alpha = \alpha_0$ , but the fit of empirical  $f(\alpha)$   
31  
32 269 to a second order polynomial is not perfect. To summarize, DSL series corresponding to the  
33  
34 270 Fabra Observatory are characterized by the following patterns:

- 35 271 1) Strong randomness, with weak signs of persistence, as Hurst exponent exceeds  
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37 272 slightly 0.5.
- 38  
39 273 2) A moderate Hölder exponent width,  $W$ , close to 0.38, with  $\alpha_1 = 0.36$  and  
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41  
42 274  $\alpha_2 = 0.74$ , thus suggesting a moderate multifractality in comparison with other DSL  
43  
44 275 series.
- 45  
46 276 3) A centered multifractal spectrum around  $\alpha_0 \approx 0.55$  and a notable asymmetry (right-  
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48 277 skewed shape) with  $B = +1.5$ . Both parameters indicate a better description of the  
49  
50 278 “fine-structure” than the “smooth-structure” of the DSL series.

### 53 279 4.2 Spatial distribution of multifractal parameters

54  
55 280 With the aim of obtaining a better picture of the multifractal variety of the different DSL  
56  
57 281 series, [Figure 4](#) shows the spatial distribution of the four parameters,  $H$ ,  $\alpha_0$ ,  $W$  and  $B$ ,  
58  
59 282 describing multifractal spectra. Minimum  $\alpha_1$  and maximum  $\alpha_2$  Hölder exponents have not  
60  
61 283 been considered as fundamental parameters, given that they could be substituted by the

284 critical Hölder exponent  $\alpha_0$  and the spectral width  $W$ , as verified after through the Principal  
285 Component Analysis.

286  
287 [Figures 4a and 4b](#) show quite similar  $H$  and  $\alpha_0$  patterns across Europe. It is worth  
288 mentioning some signs of a decreasing tendency from high to low latitudes. The strong  
289 negative gradient with latitude detected in Turkey (probably due to computational artifacts,  
290 attributable to low rain gauge density as mentioned before) and low values obtained in the  
291 SW of the Iberian Peninsula are indicative of a strong anti-persistence. Spatial patterns for  
292 the asymmetry shown in [Figure 4c](#) are difficult to interpret, as areas of “fine-structure” ( $B >$   
293  $0$ ) and “smooth-structure” ( $B < 0$ ) of the DSL regimes are not generally distributed according  
294 to latitudinal or longitudinal factors or vicinity to Atlantic or Mediterranean coasts. The  
295 spectral width,  $W$  ([Figure 4d](#)), varies within the (0.2, 0.6) interval for most of Europe. The  
296 behavior of the DSL would be qualified as close to monofractal only for a very small area of  
297 the Iberian Peninsula, with  $W$  lowering 0.2. In short, it is difficult to obtain simplified spatial  
298 patterns describing the degree of multifractality and “fine/smooth” structure of the DSL  
299 regime. Alternatively, a clustering process ([Davies and Kalkstein, 1990](#)) to group DSL series  
300 with similar patterns is applied in Section 4.3.

301 Before the clustering process, some relationships between  $H$ ,  $\alpha_0$ ,  $\alpha_1$ ,  $\alpha_2$  and  $W$  can be  
302 analyzed. [Figure 5a](#) depicts a clear linear relationship between the Hurst exponent and the  
303 critical Hölder exponent, such a pattern also found for daily extreme temperature regimes  
304 ([Burgueño et al., 2013](#)), which are characterized by strong persistence. Consequently, it  
305 could be proposed that increasing Hurst exponents would reinforce the multifractal  
306 character of the climatic series analyzed. With respect to minimum and maximum Hölder  
307 exponents, whereas  $\alpha_2$  tends to increase whit  $\alpha_0$  ([Figure 5b](#)),  $\alpha_1$  does not depict so clear  
308 evidences of a decreasing tendency. This fact would indicate that an increase of “fine-  
309 structure” of the DSL series would be accompanied by an increase of  $W$ , being reinforced  
310 then the multifractal character and the complex structure of the series. This hypothesis  
311 would be confirmed by some signs of an increasing tendency of  $W$  with  $\alpha_0$ .

### 312 **4.3 Principal Component Analysis**

313 Previous to the clustering process, a revision of ranges corresponding to the Hurst and  
314 critical Hölder exponents,  $H$  and  $\alpha_0$ , spectral asymmetry,  $B$ , and spectral width,  $W$ , is done,  
315

316 these parameters being obtained from the MF-DFA successful application to 258 out of 267  
317 DSL series. The nine rain gauges for which it has not been possible to obtain the multifractal  
318 spectrum are detailed in [Figure 1a](#). Taking into account they are widespread throughout  
319 Europe, the different pluviometric regimes would not be the reason for a successful or  
320 unsuccessful quantification of multifractal parameters. Most of  $H$  and  $\alpha_0$  values are within a  
321 relatively narrow interval from 0.4 to 0.6. Many of the asymmetries  $B$  are within the  $\pm 1.5$   
322 interval and a high number of multifractal spectra have a width  $W$  within the (0.25, 0.55)  
323 range. This distribution of the four multifractal parameters permits a better explanation and  
324 discussion of rain gauge spatial clusters obtained after the application of the Principal  
325 Component Analysis, PCA ([Jolliffe, 1986](#); [Preisendorfer, 1988](#)), and the clustering process  
326 known as average linkage algorithm ([Kalkstein et al., 1987](#); [Davis and Kalkstein, 1990](#)). [Table](#)  
327 [1](#) summarizes the results of the PCA. After the application of the PCA algorithm to the  
328 covariance matrix of original data, the four original variables,  $H$ ,  $\alpha_0$ ,  $B$ ,  $W$ , are substituted  
329 for three principal components, PC1, PC2 and PC3, which explain 99.8% of data variance. A  
330 review of the factor loadings, FL, quantifying the relationship among original variables and  
331 PCs, clearly manifests that  $H$  and  $\alpha_0$  are strongly correlated with PC1,  $W$  with PC2, and  $B$   
332 with PC3. Consequently, the values of the four original variables are substituted for the  
333 factor score values, FSC, associated with the three chosen PCs.

334  
335 [Figure 6](#) describes the spatial distribution of the three FSCs throughout Europe. Only FSC1,  
336 essentially representing  $H$  and  $\alpha_0$ , shows a similar pattern to those of [Figures 4a and 4b](#).  
337 Spatial distributions of FSC2 and FSC3 are much more spotted, in such a way that  
338 establishing any spatial patterns becomes difficult. Consequently, it is very likely that the  
339 clustering process will lead to define homogeneous groups of rain gauges without excessive  
340 spatial coherence. [This clustering process is based, as mentioned before, on the average](#)  
341 [linkage algorithm and, more specifically, on the similarity index,  \$L\_{ij}\$ , concept. This index is](#)  
342 [defined as](#)

$$343 \quad L_{ij} = D_{ij}^2 + v_i + v_j \quad (9)$$

344 [with  \$D\_{ij}\$  the Euclidean distance between centroids of clusters  \$i\$  and  \$j\$ , and  \$v\_i\$  and  \$v\_j\$  the](#)  
345 [corresponding within-group variances.  \$L\_{ij}\$  systematically increases whatever the merged](#)  
346 [pair  \$\(i, j\)\$  of clusters. Then, in each step of the AL process, only the merging of the pair of](#)

347 clusters associated with the minimum increase of  $L_{ij}$  is considered. When an accepted  
348 merging of clusters is associated with a sharp increase on  $L_{ij}$ , the previous cluster  
349 configuration is chosen as the optimum and the clustering process finishes.

350 Figure 7 depicts the evolution of  $L_{ij}$  of the average linkage algorithm with the decreasing  
351 number of clusters. It should be remembered that every time the algorithm merges two  
352 clusters of very different characteristics, a step is clearly observed in the  $L_{ij}$  curve. Looking  
353 at Figure 7, two configurations of 9 or 14 clusters could be accepted. The configuration of 14  
354 clusters is finally chosen. The 9 clusters configuration is rejected by having a large cluster and  
355 a high number of very small, or even singular, clusters. Additionally, the configuration of 9  
356 clusters would neglect valuable spatial variety from the viewpoint of DSL multifractality.

357  
358 The spatial distribution of 14 clusters is shown in Figure 8. The number of rain gauges  
359 belonging to every cluster, their centroid coordinates (average and standard deviation of  
360  $H, \alpha_0, B$  and  $W$ ) and range of every multifractal parameter are given in Table 2. Rain gauges  
361 belonging to clusters C1, C2, C3 and C7 are located at latitudes approximately north of 40°N.  
362 Clusters C4, C5 and C13 are mainly associated with areas of the Iberian Peninsula and the  
363 rest of clusters, all of them with a small number of elements, are spread throughout Europe,  
364 including Mediterranean countries. In agreement with the distribution of clusters, different  
365 European zones share more than one DSL regime. Then, to establish simple relationships  
366 between geographical factors (orography, latitudes, longitudes and vicinity to the Atlantic  
367 Ocean and the Mediterranean Sea, for instance) and DSL regimes becomes quite difficult.  
368 Nevertheless, in accordance with the well differentiated centroids summarized in Table 2,  
369 the classification of rain gauges within specific clusters becomes coherent.

#### 4.4 Complexity Index

372 Another complementary point of view of the DSL regime in Europe consists in the  
373 application of the complexity index, CI (Burgueño et al., 2013), which is defined through the  
374 addition of the standardized  $\alpha_0, z(\alpha_0), B, z(B)$ , and  $W, z(W)$ , parameters, in agreement with  
375 the concept of complexity developed by Shimizu et al. (2002)

$$Z = z(\alpha_0) + z(B) + z(W) \quad (10)$$

377 For a better interpretation of this addition,  $Z$  is also normalized,  $CI$  being defined as



$$CI(Z) = \frac{Z - \langle Z \rangle}{S(Z)} \quad (11)$$

$\langle Z \rangle$  being equal to zero and  $S(Z)$  the standard deviation of  $Z$ . Definition given by Equation (11) permits a straightforward quantification in terms of “high complexity” ( $Z \geq 0$ ) and “low complexity” ( $Z < 0$ ). Table 3 summarizes the number of DSL series belonging to different  $Z$  intervals given in standard deviation units. It is outstanding the concentration of  $CI(Z)$  within the  $\pm 1.0$  interval (in standard deviation units), close to 70%, and the very low number of series with very smooth- (9 out of 258) and very fine-structure (2 out of 258). As a general feature, DSL series associated with fine-structure ( $Z \geq 0$ ) are slightly predominant (close to 56%) in comparison with those related to smooth-structure ( $Z < 0$ ), which represent close to 44% of series. The spatial distribution of  $CI(Z)$  is depicted in Figure 9. Similarly to the spatial distribution of clusters, a clear spatial coherence is not found, being detected small areas sharing DSL series with smooth and fine-structures, especially in Central and Western Europe. Nevertheless, south-western Iberian Peninsula, southern Italy and the Adriatic coast are characterized by smooth-structure of the DSL regime. North-eastern regions of Europe have a predominant (almost unique) fine-structure of the DSL regime.

#### 4.5. Comparisons with monofractal results.

Comparisons with monofractal results (Lana et al., 2010) are aimed to assess if Hurst exponents derived from rescaled range, R/S, analysis and from MF-DFA are coincident and, in addition, to establish some relationship between DSL series characterized by smooth/fine-structure and monofractal properties derived from the reconstruction theorem, namely, correlation dimension, Kolmogorov entropy and Lyapunov exponents.

Hurst exponent values obtained from the R/S analysis (within the 0.2–0.7 interval) are generally slightly shifted from those derived from the MF-DFA (within the 0.1–0.6 range). A good number of Hurst exponents derived from R/S analysis are related to randomness with slight signs of persistence ( $H \approx 0.6$ ) of DSL series, whereas a significant number of Hurst exponents estimated by MF-DFA lying within the (0.4, 0.6) interval points also to predominant randomness, but with weak signs of both persistence or anti-persistence. Anti-persistence is detected, both for R/S and MF-DFA, at south and south-western Iberian Peninsula and, additionally, at southern Italy and Sicily, in agreement with MF-DFA. The highest Hurst exponents are predominantly found at high latitudes (North-Eastern Europe), close to Atlantic Ocean and Baltic Sea. In short, the results concerning Hurst exponents

410 derived from R/S and MF-DFA, and its spatial distributions, could be assumed as roughly  
411 coincident.

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4 413 With respect to possible relationships among monofractal variables and fine/smooth-  
5 structures defined by CI, the incidence of the Kolmogorov entropy should be discarded as  
6 414 their values are scattered throughout Europe without clear spatial patterns. Consequently,  
7 415 the loss of memory of the physical system should be a common feature, but not a  
8 416 discriminating factor to distinguish between smooth and fine-structures. On the contrary,  
9 417 spatial distributions of low values of the correlation dimension (giving an estimation of the  
10 418 minimum number of nonlinear equations required to describe the physical process) and high  
11 419 values of the Lyapunov exponents (quantifying the predictive instability of the physical  
12 420 system) could be associated with smooth-structures. From the point of view of CI, south and  
13 421 south-west of the Iberian Peninsula, south of Italy and Sicily are characterised by smooth-  
14 422 structure. Almost the same areas are linked to the highest values of the first Lyapunov  
15 423 exponent, which strongly governs the predictive instability, and the lowest correlation  
16 424 dimensions. Consequently, it could be proposed the hypothesis that smooth-structures of  
17 425 the DSL regime would be described by a relatively simple system of nonlinear equations.  
18 426 Nevertheless, as a counterpart, the predictive instability would be notable. Given that many  
19 427 areas for the rest of Europe share smooth- and fine-structures, the conjecture of a certain  
20 428 relationship between fine-structure and high correlation dimensions and low Lyapunov  
21 429 exponents is difficult to be ascertained.  
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24 432 It is also worth of mention that, in Mediterranean countries, with quite usual drought  
25 433 episodes (some of them unusually long), smooth-structures are compatible with low  
26 434 correlation dimensions and high Lyapunov exponents, as well as with Hurst exponents  
27 435 indicating anti-persistent behaviours.  
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#### 30 437 **4.6 Autoregression results**

31 438 **Figure 10** illustrates the dependence of MAD and OAO on latitude. A roughly linear decrease  
32 439 of MAD, from 20 days up to approximately 4 days, with increasing latitudes ranging from  
33 440 30°N to 45°N, is clearly observed. For higher latitudes, this almost linear decreasing  
34 441 tendency is smoother, MAD reaching values close to 2 days for the highest latitudes. Even  
35 442 though OAO does not show a so clearly linear evolution, an increasing tendency is found up  
36 443 to latitudes close to 45°N. For higher latitudes, values of OAO remain within the 180-200

444 range, this last autoregressive order being the largest lag analyzed. These changes on  
445 tendencies close to 45°N were also detected by [Serra et al. \(2015\)](#), who analyzed, among  
1 446 other questions, the evolution with latitude of the L-moment coefficient of variation of the  
2 447 longest DSL. For a better understanding of OAO, It has to be underlined that for the same  
3 448 OAO, the required recording period depends on the DSL, especially when comparing dry and  
4 449 wet pluviometric regimes or very different latitudes. As an example, an OAO value equal to  
5 450 150 (150 consecutive dry spells) represents a wide range of recording periods for the  
6 451 different rain gauges of the European database. These periods vary from 10 months  
7 452 (Bjoernoeya, Norway, 74.59°N, 19.02°E) with a wet pluviometric regime and a short average  
8 453 DSL of 2.1 days, up to 5 years (Alicante, Spain, 38.01°N, 0.71°W) with a dry pluviometric  
9 454 regime and a longer average DSL of 12.5 days.  
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21 456 With respect to the distribution of MAD and OAO values, it is worth mentioning that MAD  
22 457 values are mainly concentrated within the 2-5 days interval (197 out of 267 series), without  
23 458 signs of a Gaussian distribution, and most OAO values equal to or exceeding 175 (233 out of  
24 459 267 series). [Figure 11](#) shows the histogram of the residuals of all DSL series, defined as the  
25 460 difference between empiric values and those reproduced with the OAO. These residuals are  
26 461 mainly concentrated within the (-4, 2) days range, representing close to 70% of DSL in  
27 462 Europe, and without signs of a Gaussian distribution.  
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36 464 [Figure 12](#) describes the spatial distribution of MAD and OAO for the 267 DSL series. Both  
37 465 parameters have been classified in five intervals corresponding to 0-20<sup>th</sup>, 20-40<sup>th</sup>, 40-60<sup>th</sup>,  
38 466 60-80<sup>th</sup> and 80-100<sup>th</sup> percentiles of their empirical cumulative distributions. The first class of  
39 467 MAD percentiles (0-20<sup>th</sup>) covers latitudes exceeding 50°N. Most of series with MAD  
40 468 belonging to the second interval are within the (45-55°N, 5-15°E) latitude and longitude  
41 469 ranges. The third interval mostly corresponds to a long fringe (45-60°N, 0-45°E) with a low  
42 470 rain gauge spatial density for longitudes exceeding 15°E. The fourth interval is  
43 471 predominantly associated with latitudes 40-60°N, whatever the longitude. Finally, the fifth  
44 472 interval is mainly linked to the Iberian Peninsula and some Mediterranean coastal rain  
45 473 gauges, with latitudes south of 45°N. The five classes of OAO percentiles are however,  
46 474 scattered throughout Europe, without any clear spatial pattern.  
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60 476 Two examples of MAD values and DSL residuals for very different latitudes and pluviometric  
61 477 regimes are shown in [Figure 13a](#) (Varexjoe, Sweden) and [Figure 13b](#) (Zaragoza, Spain).  
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478 Whereas MAD is close to 3 days for Varexjoe, it attains 7 days for Zaragoza. Additionally, the  
479 standard deviations of the residuals are close to 3 and 7 days respectively. After a detailed  
1 480 review of DSL residuals, it is observed that the percentage of overestimated DSL (negative  
2 481 residuals) for both series is quite similar (close to 66%). Nevertheless, a notable difference  
3 482 appears when the most negative residuals are compared. Whereas the largest DSL  
4 483 overestimation is 5 days for Varexjoe, it is 12 days for Zaragoza. Another similar feature is  
5 484 observed when comparing positive residuals (underestimated DSL) exceeding one standard  
6 485 deviation. In both cases the percentage of underestimations is relatively small (12%).  
7 486 Nevertheless, notable differences are detected once again when comparing the maximum  
8 487 underestimation. Whereas it could be assumed as notably high (30 days) for Varexjoe, it  
9 488 should be considered absolutely inappropriate for Zaragoza (70 days).  
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21 490 [Table 4](#) summarizes the main features of the residuals for the 267 DSL series. It should be  
22 491 accepted that predicted DSL are not biased given that the average of the residuals is very  
23 492 close to zero, whatever the DSL series. Their standard deviation varies within a wide range  
24 493 from 1.7 to 29.5 days. The whole set of European DSL residuals would be then represented  
25 494 by a moderate average standard deviation of 4.6 days. In spite of this small length, at least  
26 495 one of the DSL series would be predicted by the AR(p) process without an excessive error, as  
27 496 manifested by the mentioned maximum standard deviation of 29.5 days. Skewness and  
28 497 kurtosis for every one of the 267 sets of residuals indicate that a Gaussian distribution  
29 498 should be discarded. These empirical skewness and kurtosis notably depart from values  
30 499 expected for a Gaussian distribution, which should be equal to 0.0 and 3.0 respectively.  
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43 501 Percentages  $R(>2\sigma)$  of positive residuals (underestimated DSL) exceeding two standard  
44 502 deviations are not very relevant. Minimum and maximum percentages of 2.9% and 5.3%  
45 503 respectively, as well as an average of 4.7% for all DSL series, summarize the behavior of the  
46 504 largest positive residuals. An excessive overestimation of DSL is quantified by the percentage  
47 505  $R(<-\sigma)$  of negative residuals lowering at least minus one standard deviation. The  
48 506 overestimation of DSL is not a relevant problem, given that for the whole database analyzed,  
49 507 only 4.6% of DSL have been overestimated and the maximum detected ratio for a single DSL  
50 508 series attains 8.3%.  
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60 510 The hypothesis that the positive residuals (underestimated DSL) are higher than the negative  
61 511 residuals (overestimated DSL) is not only suggested by the results given in [Table 4](#), but also  
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512 searching for the maximum positive residual for every one of the 267 DSL series. These  
513 extreme residuals range from 14.7 to 286.8 days. The first length could be considered high,  
1 514 but in some way acceptable. The second length is unacceptable. Conversely, minima of  
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4 515 negative residuals of the 267 series range from  $-30.0$  to  $-2.7$  days.

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7 517 For a better comprehension of advantages and shortcomings of the AR(p) process applied to  
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9 518 the DSL prediction, additional reasons for these large values of residuals have to be found, as  
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11 519 well as for the high standard deviation detected for at least one series of residuals and to  
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13 520 null ratios of residuals lowering  $-\sigma$  or within the range  $(\sigma, 2\sigma)$ . Four DSL series with these  
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15 521 characteristics are found. They come from pluviometric records in South-eastern  
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17 522 Mediterranean, with longitudes within the  $(25.2^{\circ}\text{E} - 35.5^{\circ}\text{E})$  range and low latitudes within a  
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19 523 narrow fringe  $(31.8^{\circ} - 35.3^{\circ}\text{N})$ , where dryness is well-known. All of them have a very high  
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21 524 maximum underestimation of real DSL (from  $-152.9$  to  $-286.8$  days) and maximum  
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23 525 outstanding overestimation of real DSL (from 25.3 to 30.0 days), accompanied by high  
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25 526 standard deviation of residuals, varying from 18.7 to 29.5 days. It is also worth mentioning  
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27 527 that the percentage of overestimated DSL lowering the residual  $-\sigma$  is null for two series, and  
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29 528 almost null for the other two.

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33 530 A possible relationship between CI and MAD is shown in [Figures 14a](#). Most of MAD values  
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35 531 range from 2 to 5 days and are associated with CI within the  $(-2.0, +2.5)$  interval. [Figure 14b](#)  
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37 532 shows signs of a possible relationship between CI and OAO too. Whereas for values of CI  
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39 533 within the  $\pm 2.5$  interval, a concentration of high values of OAO (exceeding 145) is observed,  
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41 534 a positive trend in OAO values lowering 145 is detected for the  $(-4.0, +1.5)$  CI range.

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## 537 5. Discussion and conclusions

538 The present MF-DFA widens the monofractal analysis of the European DSL regime (Lana et  
1 539 al., 2010). The degree of complexity of DSL series derived from daily pluviometric records  
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3 540 has been quantified through multifractal parameters  $H$ ,  $\alpha_0$ ,  $B$  and  $W$ , and their synthesis, the  
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5 541 complexity index, CI.

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8 542 With respect to possible relationships between monofractal and multifractal parameters, it  
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10 543 is worth mentioning that DSL regimes characterized by physical mechanisms with a low  
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12 544 number of nonlinear equations (correlation dimension) and notable predictive instability  
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14 545 (high positive Lyapunov exponents) are generally associated with CI negative values  
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16 546 corresponding to smooth structures. On the contrary, fine structures (positive CI) are related  
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18 547 to mechanisms requiring a high number of nonlinear equations. As a counterpart, the  
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20 548 predictive instability would not be as high as for smooth structures. It is also worth  
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22 549 mentioning that the loss of memory of the system (Kolmogorov entropy) is a common  
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24 550 feature to all DSL regimes, without clear relationships with multifractal parameters.

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27 551 From the point of view of multifractal parameters across Europe, signs of a North-east to  
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29 552 South-west decreasing trend is only observed for the Hurst exponent,  $H$ , and the critical  
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31 553 Hölder exponent,  $\alpha_0$ . For the other two multifractal parameters, spectral width,  $W$ , and  
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33 554 spectral asymmetry,  $B$ , clear spatial patterns are not obtained. Neither the quantification of  
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35 555 DSL complexity, in terms of CI, nor a clustering process, based on a previous PCA, have  
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37 556 permitted a spatially coherent distribution of fine- and smooth-structures. This fact could be  
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39 557 considered a shortcoming of the multifractal analysis. Alternatively, strong dependence of  
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41 558 the DSL regime on topographic parameters (height above sea level and orographic slope for  
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43 559 instance) and on other local and atmospheric dynamic variables (dominant wind direction,  
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45 560 vicinity to orographic barriers such as the Alps and the Pyrenees and vicinity to Atlantic or  
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47 561 Mediterranean seas) could be an explanation to this lack of coherent spatial clustering.

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50 563 With respect to the ARIMA process, several questions concerning the residuals of predicted  
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52 564 DSL should be mentioned. First, they are not distributed according to a Gaussian model.  
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54 565 Second, predicted DSL are not biased given that residual averages are for every one of the  
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56 566 DSL series almost null. Third, whereas the overestimation of DSL is not a very relevant  
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58 567 problem, their underestimation may become very important, especially for pluviometric  
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60 568 series characterized by long DSL. Fourth, similar to spatial patterns of multifractal  
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62 569 parameters, the spatial distribution of MAD is also complex and the OAO is not coherently

570 distributed. And fifth, all DSL residuals characterized by high standard deviations and very  
571 high underestimations belong to pluviometric records associated with very dry regimes at  
1 572 low latitudes in the Mediterranean area. In other words, DSL for dry regimes (usually low  
2 573 latitudes) are more difficult to be predicted than for wet regimes (generally high latitudes).  
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4 574 In fact, DSL belonging to the aforementioned very dry regimes should be qualified as very  
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6 575 hardly predictable in agreement with results obtained from the ARIMA process. An  
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8 576 acceptable prediction for these very dry regimes might be achieved with the ARIMA(p,d,q)  
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10 577 model, applied before for instance to monthly rainfall (Wang et al., 2014). Also, methods  
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12 578 based on conditional Poisson distribution and Monte Carlo algorithms (Jung et al., 2006)  
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14 579 might be considered. In short, the prediction of these extreme DSL would be beyond the  
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16 580 main objectives of this paper.  
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21 582 It is worth noticing the high OAO values for most of DSL series. These high values could be in  
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23 583 agreement with the Kolmogorov entropy, which quantifies the loss of memory of the  
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25 584 physical mechanisms along the time process. Prediction of DSL series affected by a high loss  
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27 585 of memory would require then a notable number of previous steps (DSLs) to obtain  
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29 586 acceptable predictions.  
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33 588 Similarly to comparisons made between monofractality and multifractality, coincidences and  
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35 589 disagreements between multifractal and ARIMA results are worthy to be mentioned. A clear  
36  
37 590 relationship between CI and MAD has not been found. Nevertheless, low MAD, from 2 to 5  
38  
39 591 days, are usually associated with fine and smooth structures with CI ranging from -2 to +2.  
40  
41 592 The highest MAD are detected for CI within the range (-4, -3), corresponding to smooth  
42  
43 593 structures from the multifractal point of view. With respect to a possible relationship  
44  
45 594 between CI and OAO, a great number of DSL series are characterized by OAO exceeding 150  
46  
47 595 and belonging to fine or smooth structures. A few cases of OAO lowering 150 with CI varying  
48  
49 596 from -4 (very smooth structures) to +1 (moderate fine structures) are characterized by a  
50  
51 597 clear linear increasing trend.  
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53 598  
54 599 Finally, latitudinal changes on MAD and OAO are outstanding, roughly at 45°N, thus  
55  
56 600 reinforcing the hypothesis of a certain dependence on latitude. As mentioned in Section 4.6,  
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58 601 a similar behavior has been observed in recent statistical analyses of long DSL.  
59

60 602

	EV(%)	$FL_H$	$FL_{\alpha_0}$	$FL_W$	$FL_B$
PC1	49.1	0.992	0.970	0.180	-0.084
PC2	25.5	0.101	0.205	0.983	0.030
PC3	25.2	-0.035	-0.112	0.032	0.996

**Table 1.** Explained variance, EV, in percentage, of the original data for every one of the three first rotated Principal Components, PC.  $FL_H$ ,  $FL_{\alpha_0}$ ,  $FL_W$  and  $FL_B$  represent the rotated factor loadings for every PC.

Cluster	$\langle H \rangle$	$\sigma_H$	$R_H$	$\langle \alpha_0 \rangle$	$\sigma_{\alpha_0}$	$R_{\alpha_0}$	$\langle W \rangle$	$\sigma_W$	$R_W$	$\langle B \rangle$	$\sigma_B$	$R_B$
C1(53)	0.504	0.055	0.262	0.519	0.056	0.282	0.323	0.055	0.223	-1.038	0.602	2.594
C2(47)	0.490	0.043	0.171	0.506	0.044	0.175	0.398	0.035	0.120	1.296	0.348	1.275
C3(85)	0.495	0.053	0.229	0.528	0.053	0.275	0.432	0.046	0.200	-0.145	0.575	3.126
C4(12)	0.326	0.054	0.163	0.346	0.053	0.160	0.445	0.025	0.097	1.068	0.287	0.915
C5(08)	0.183	0.058	0.157	0.190	0.058	0.164	0.258	0.068	0.186	0.742	0.495	1.523
C6(03)	0.125	----	----	0.122	----	----	0.157	----	----	-0.817	----	----
C7(25)	0.501	0.054	0.232	0.507	0.048	0.225	0.247	0.043	0.148	0.935	0.708	2.824
C8(01)	0.541	----	----	0.651	----	----	0.618	----	----	-1.132	----	----
C9(01)	0.508	----	----	0.407	----	----	0.591	----	----	3.364	----	----
C10(01)	0.469	----	----	0.612	----	----	0.599	----	----	-2.901	----	----
C11(06)	0.334	0.014	0.039	0.365	0.019	0.055	0.402	0.060	0.170	-0.612	0.392	1.133
C12(10)	0.474	0.044	0.123	0.492	0.045	0.131	0.533	0.034	0.126	1.034	0.248	0.709
C13(05)	0.251	0.015	0.036	0.258	0.016	0.037	0.223	0.051	0.113	1.782	0.560	1.305
C14(01)	0.177	----	----	0.197	----	----	0.367	----	----	-0.639	----	----

**Table 2.** Summary of average  $\{\langle H \rangle, \langle \alpha_0 \rangle, \langle W \rangle, \langle B \rangle\}$ , standard deviation  $\{\sigma_H, \sigma_{\alpha_0}, \sigma_W, \sigma_B\}$  and range  $\{R_H, R_{\alpha_0}, R_W, R_B\}$  of the four multifractal parameters representing every spatial cluster. The number of rain gauges belonging to every cluster is given in the first column within parentheses. Standard deviations for clusters with less than four rain gauges are not included.

CI	$Z < -2$	$-2 \leq Z < -1$	$-1 \leq Z < 0$	$0 \leq Z < 1$	$1 \leq Z < 2$	$Z \geq 2$
$N_{DSL}$	9	31	74	115	27	2

**Table 3.** Number of DSL series,  $N_{DSL}$ , with Complexity Index, CI, within several standard deviation intervals.



	$\sigma(\text{days})$	Sk	K	$R(\sigma, 2\sigma)$ (%)	$R(>2\sigma)$ (%)	$R(<-\sigma)$ (%)
<b>Minimum</b>	1.71	1.85	4.43	0.0	2.90	0.0
<b>Maximum</b>	29.48	5.19	31.20	9.21	5.27	8.28
<b>Average</b>	4.64	2.55	10.91	7.03	4.69	4.55

624

**Table 4.** Statistical summary (minimum, maximum and average) of the 267 samples of standard deviations,  $\sigma$ , skewness, Sk, kurtosis, K, and percentages of residuals within the  $(\sigma, 2\sigma)$ ,  $R(\sigma, 2\sigma)$ , exceeding  $2\sigma$ ,  $R(>2\sigma)$ , and lowering  $-\sigma$ ,  $R(<-\sigma)$ , ranges.

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814 **List of Figures**

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1 816 **Figure 1. a)** Spatial distribution of the 267 European rain gauges. Solid circles design rain gauges for  
2 817 which the multifractal spectrum has not been computed. b) Number of available gauges per year  
3 818 along the recording period. c) Histogram of  $N_{\text{DSL}}$ . d) DSL series for Vaexjoe (Sweden) and Almería  
4 819 (Spain).

5 819  
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7 820 **Figure 2. a)** Fluctuation function  $F_q(s)$  for some  $q^{\text{th}}$ -orders, b)  $h(q)$  curve fitted to a polynomial  
8  
9 821 function of  $q^{\text{th}}$ - orders (dashed lines correspond to the Hurst exponent,  $h(q=2)$ ), c) multifractal scaling  
10 822 exponents,  $\tau(q)$ , and d) Hölder exponents,  $\alpha(q)$ , for DSL series corresponding to the Fabra  
11 823 Observatory (NW Mediterranean coast).

12 823  
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14 824 **Figure 3.** Multifractal spectrum (solid circles) and second order polynomial fit (solid line) for DSL  
15 825 series corresponding to the Fabra Observatory (NW Mediterranean coast).

16 825  
17  
18 826 **Figure 4.** Spatial distribution of a) Hurst exponent,  $H$ , b) critical Hölder exponent,  $\alpha_0$ , c) asymmetry,  
19 827  $B$ , and d) spectral width,  $W$ .

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24 829 **Figure 6.** Spatial distribution of FSCs obtained after the PCA applied to parameters  $H$ ,  $\alpha_0$ ,  $B$  and  $W$ .

25 829  
26 830 **Figure 7.** Evolution of the similarity index along the clustering process.

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28 831 **Figure 8.** Spatial distribution of the selected 14 clusters.

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30 832 **Figure 9.** Spatial distribution of (a) negative CI and b) positive CI.

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32 833 **Figure 10.** Dependence of (a) MAD and (b) OAO on latitude.

33 833  
34 834 **Figure 11.** Histogram of DSL residuals for the 267 DSL series.

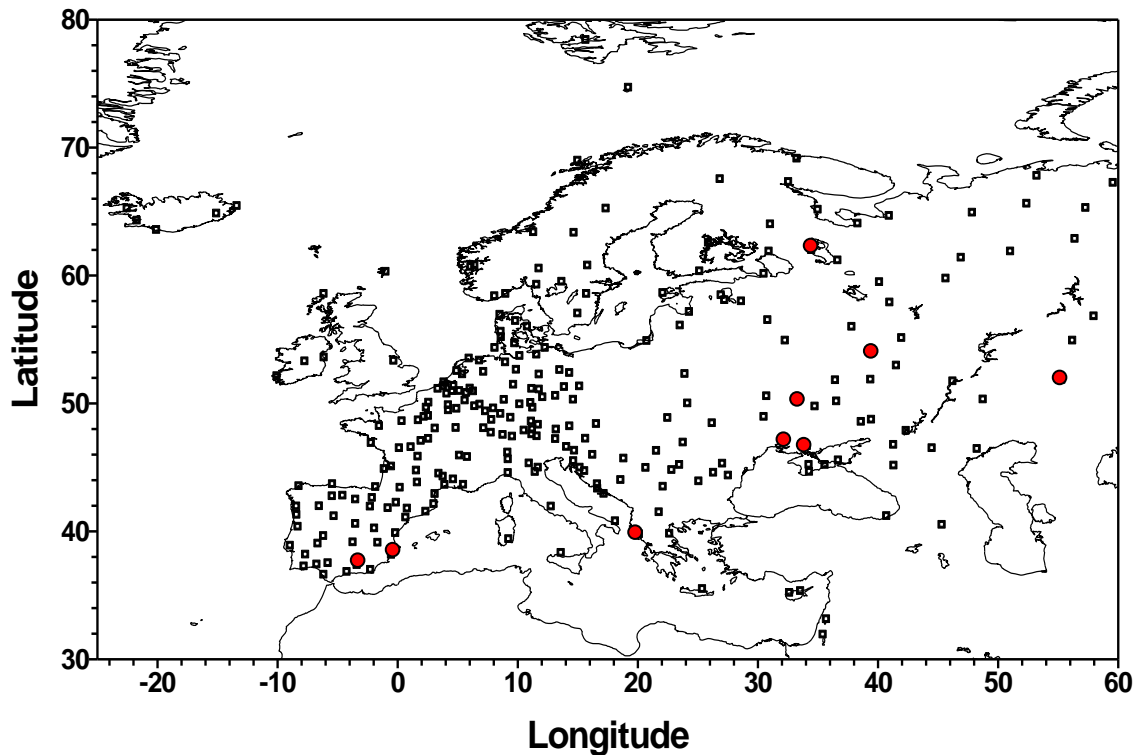
35 834  
36 835 **Figure 12.** Spatial distribution of (a) MAD and (b) OAO, which have been classified as belonging to 0-  
37 836 20<sup>th</sup>, 20-40<sup>th</sup>, 40-60<sup>th</sup>, 60-80<sup>th</sup> and 80-100<sup>th</sup> percentile intervals. The corresponding thresholds of MAD  
38 837 and OAO are included, in increasing order, at the foot of every map.

39 837  
40 838 **Figure 13.** Two examples of DSL residuals for DSL series belonging to (a) Vaexjoe (Sweden) and (b)  
41 839 Zaragoza (Spain). The evolution of the corresponding MAD with OAO is also included.

42 839  
43 840 **Figure 14.** Dependence of (a) MAD and (b) OAO on CI.

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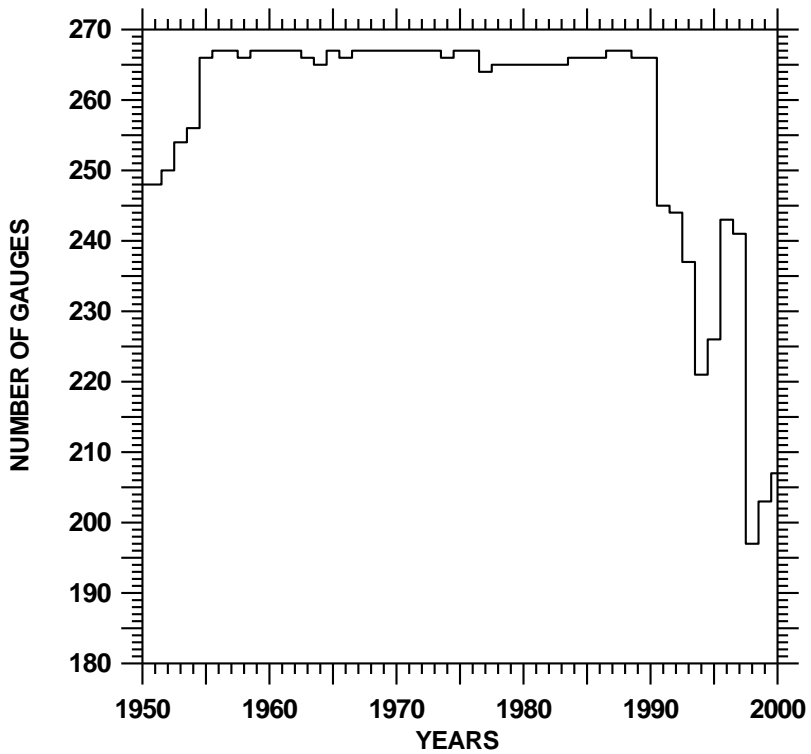
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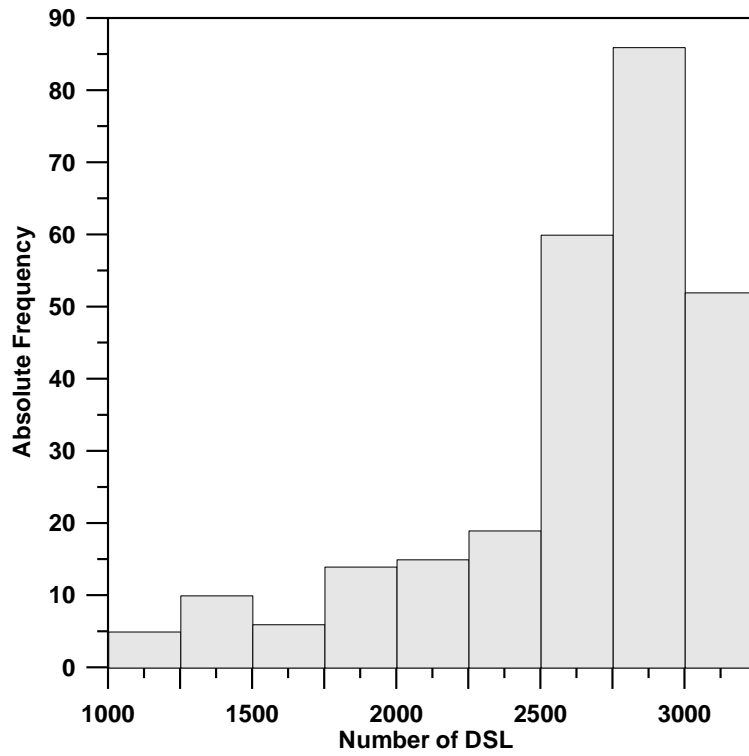
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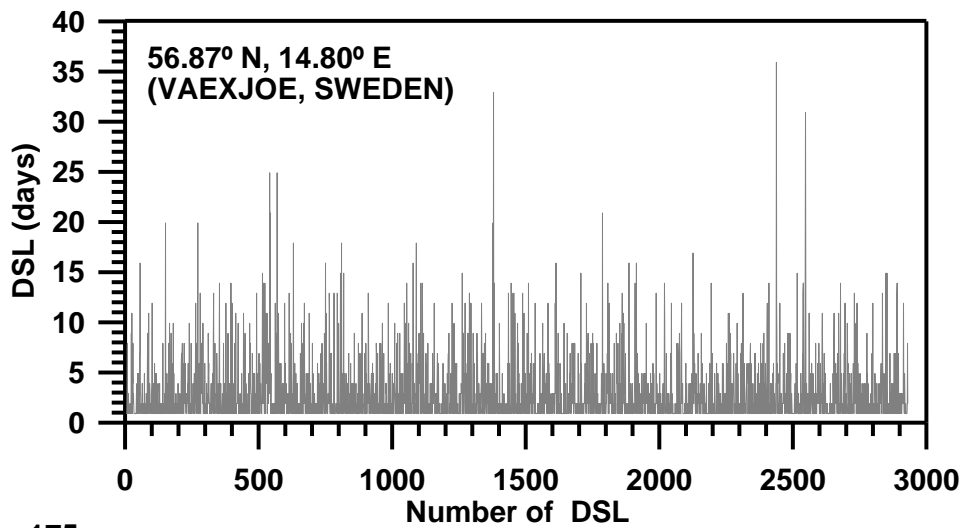
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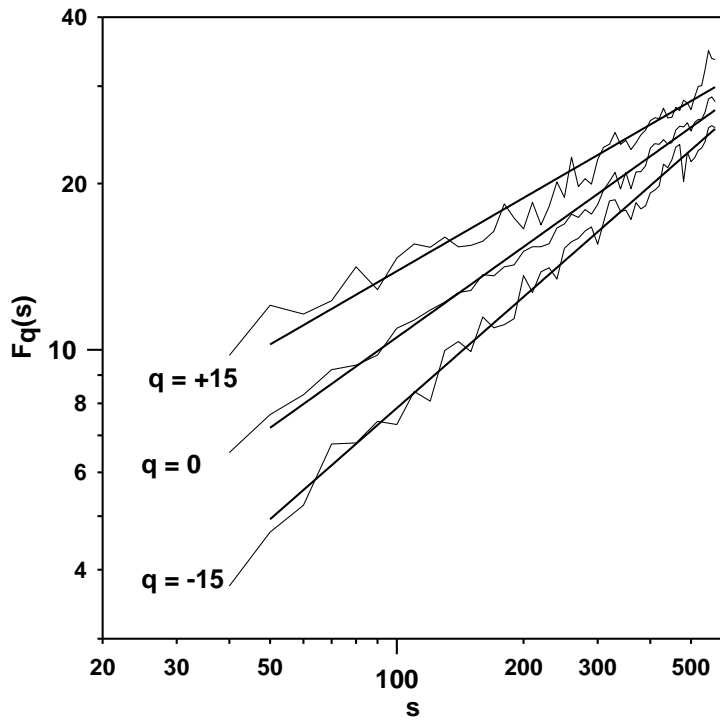
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Figure 1

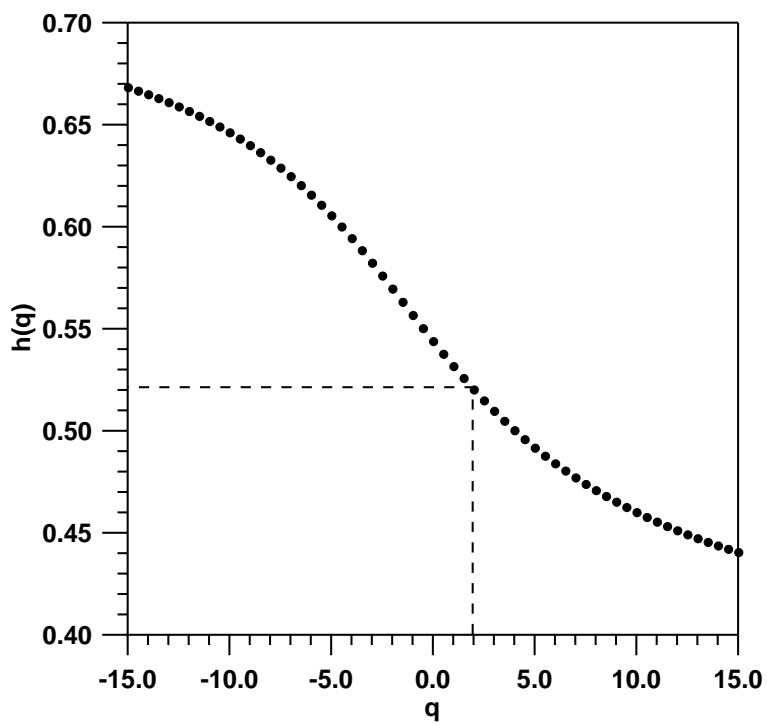
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(b)



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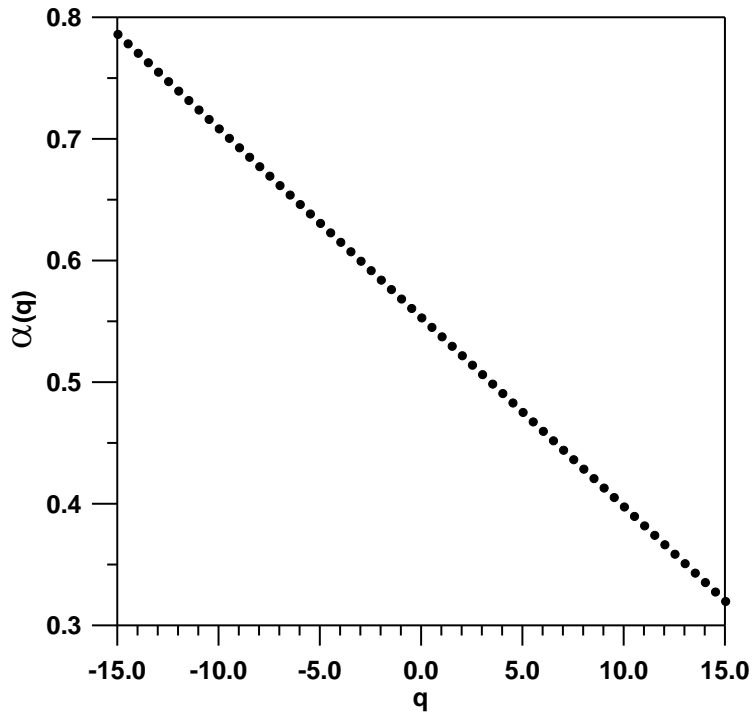
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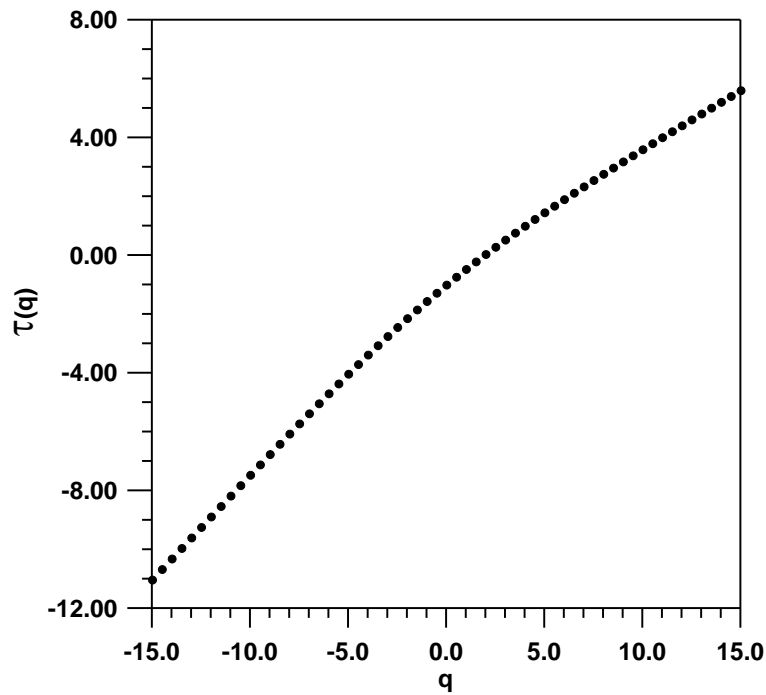
892 (c)

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895 (d)



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900 **Figure 2**

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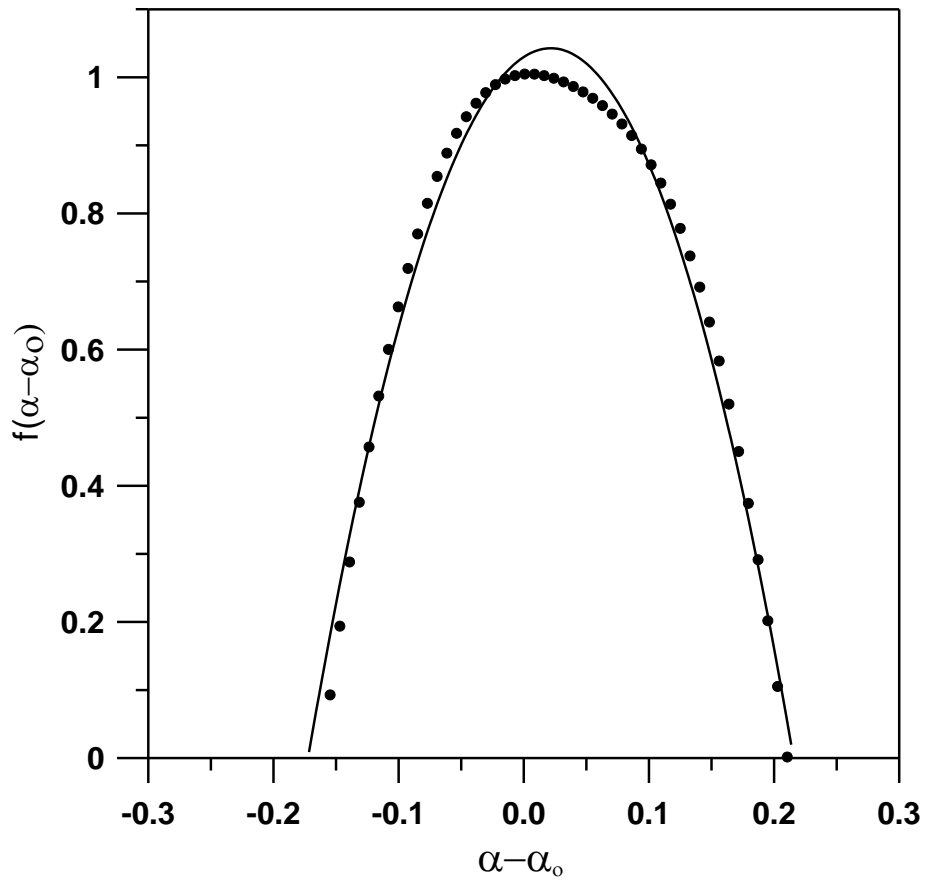
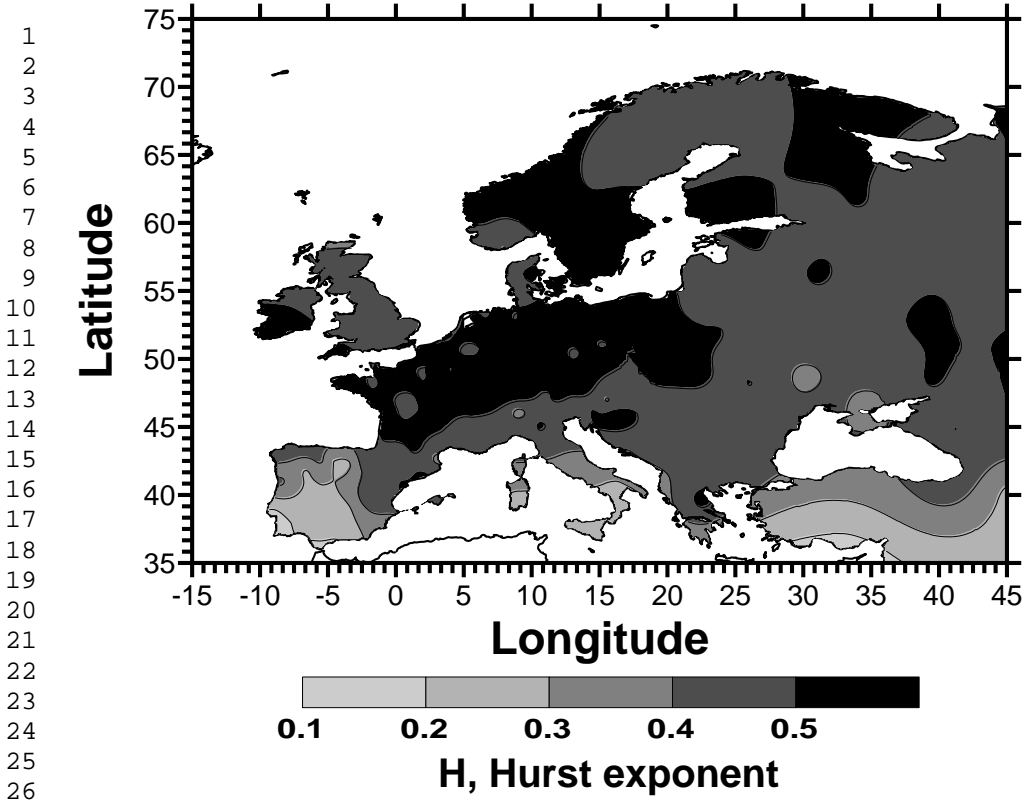
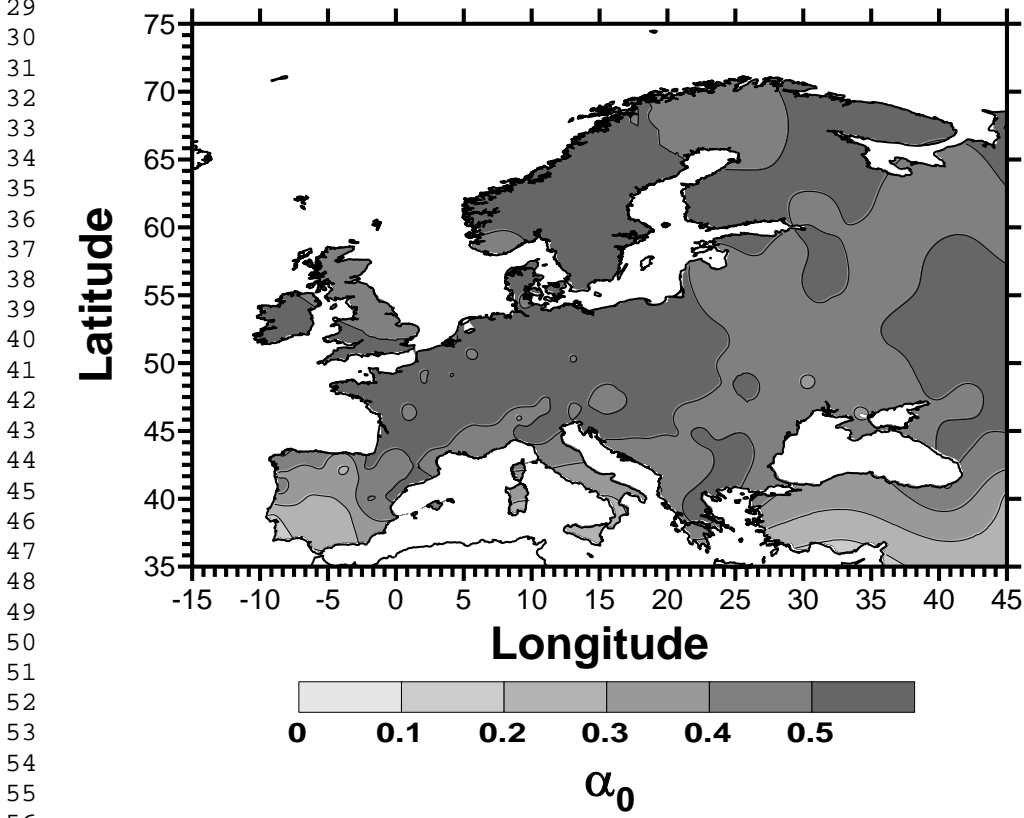


Figure 3

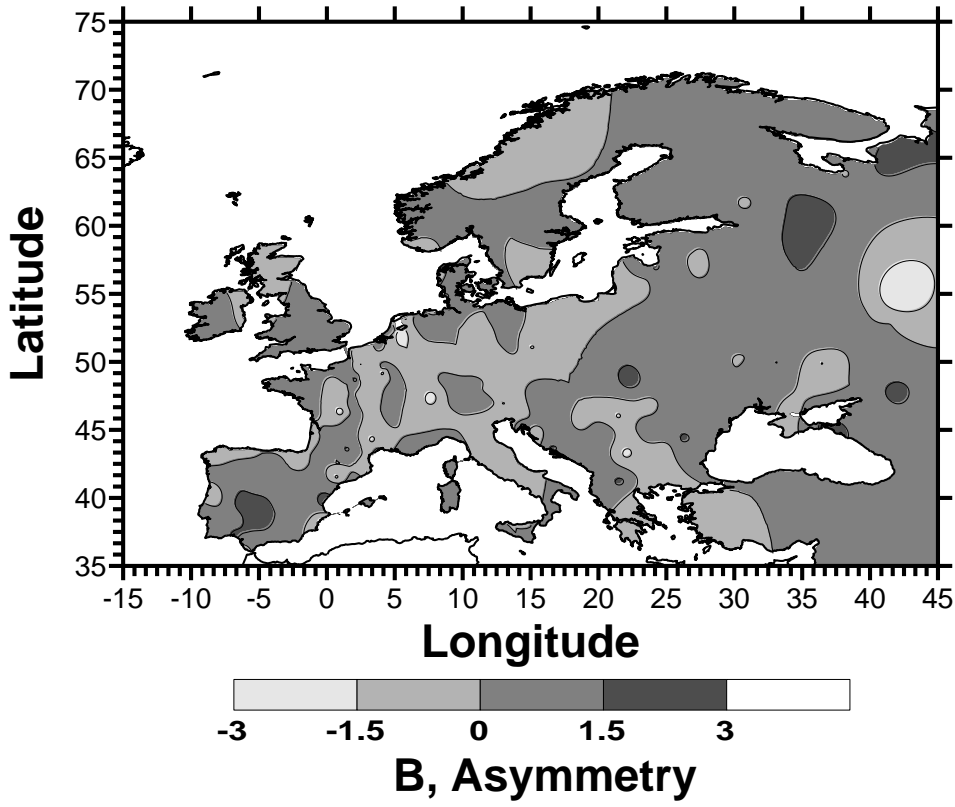
(a)



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(d)

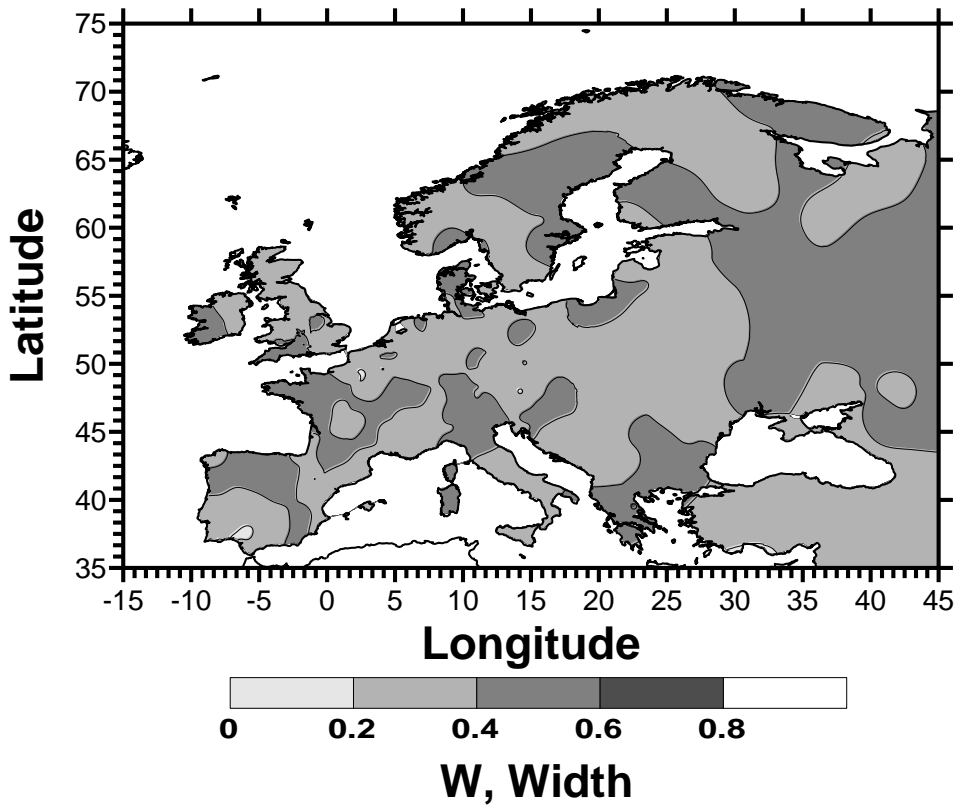
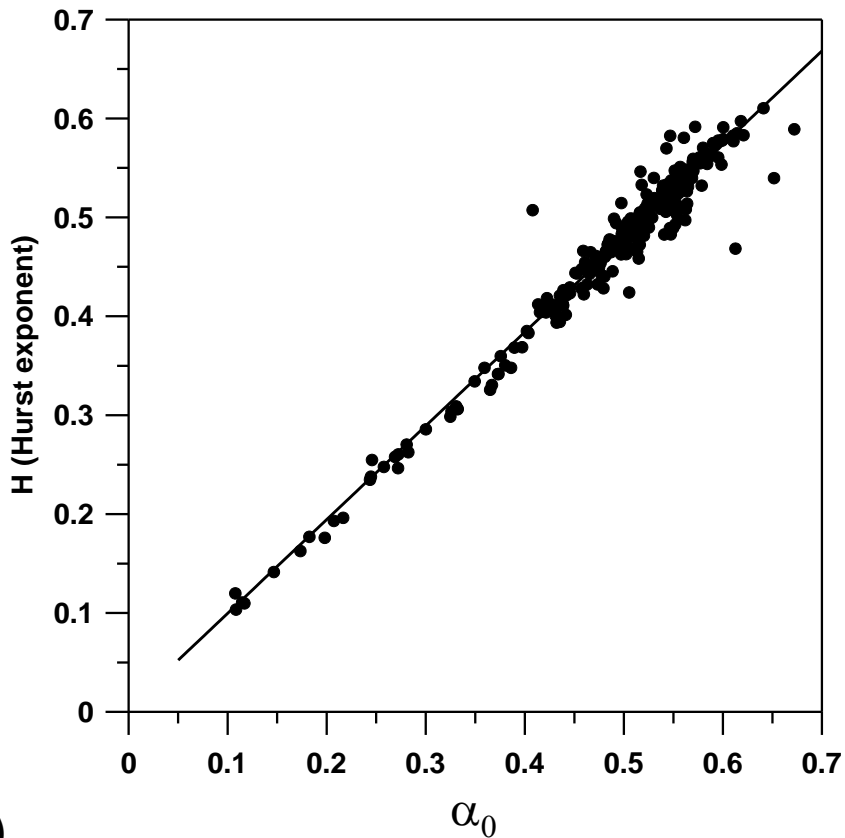


Figure 4

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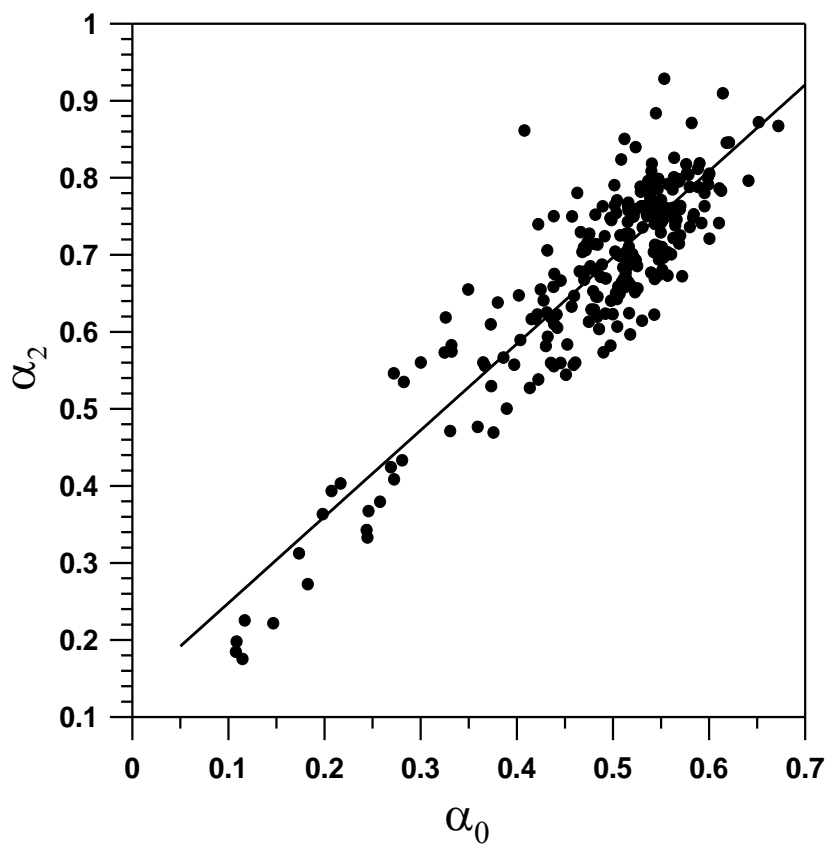


Figure 5

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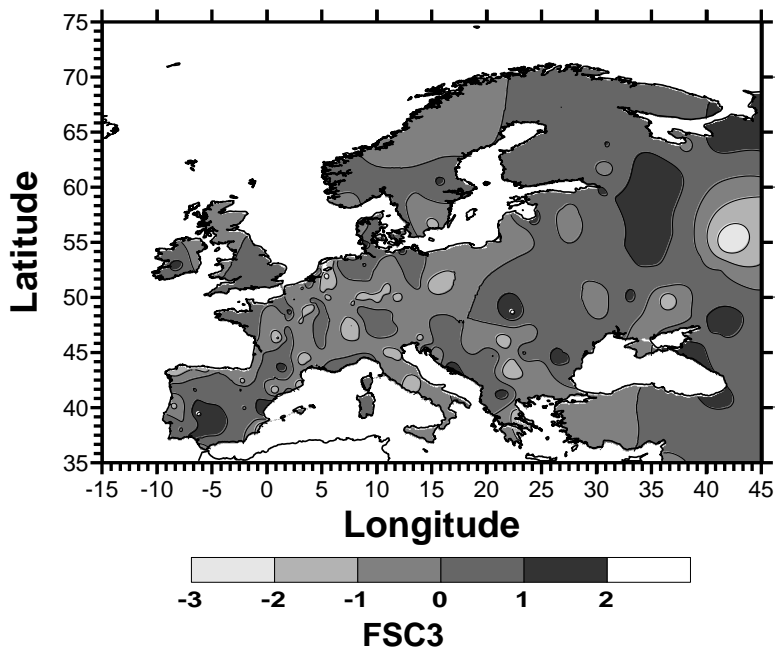
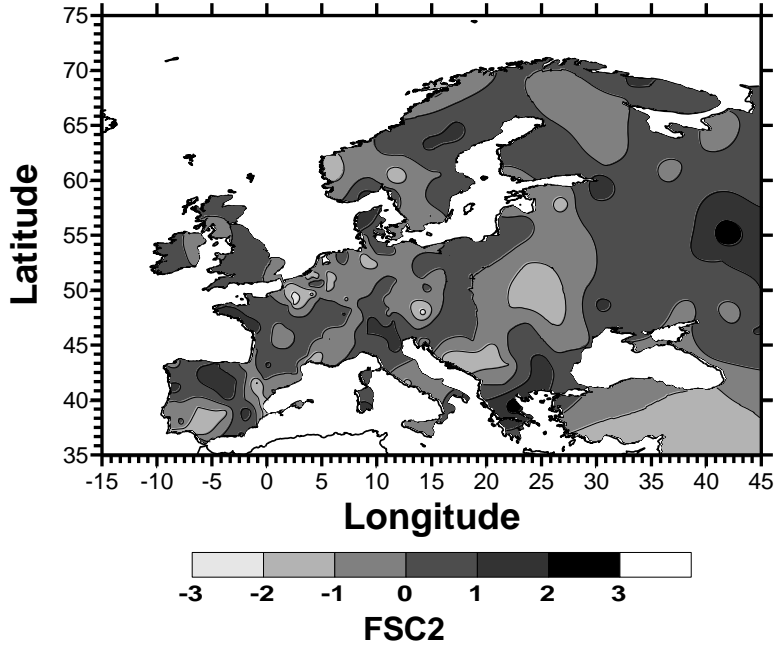
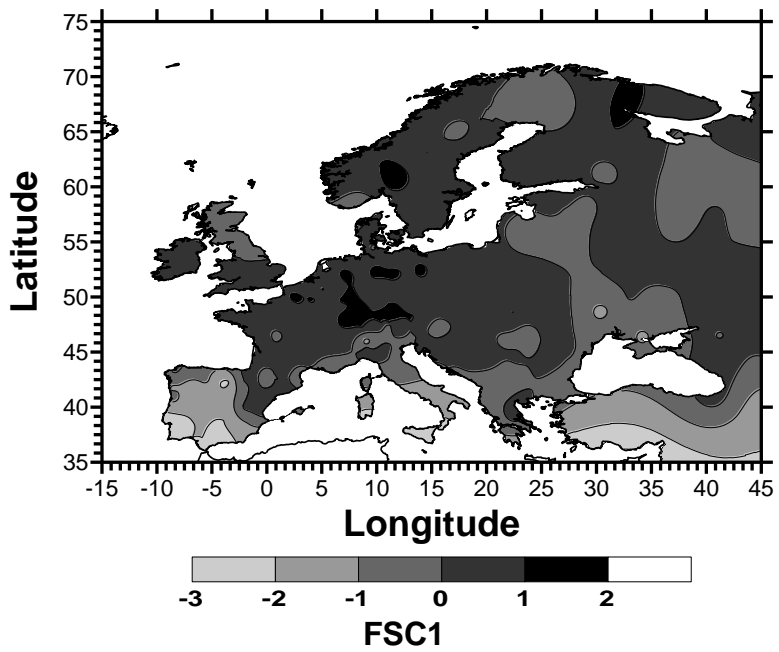


Figure 6



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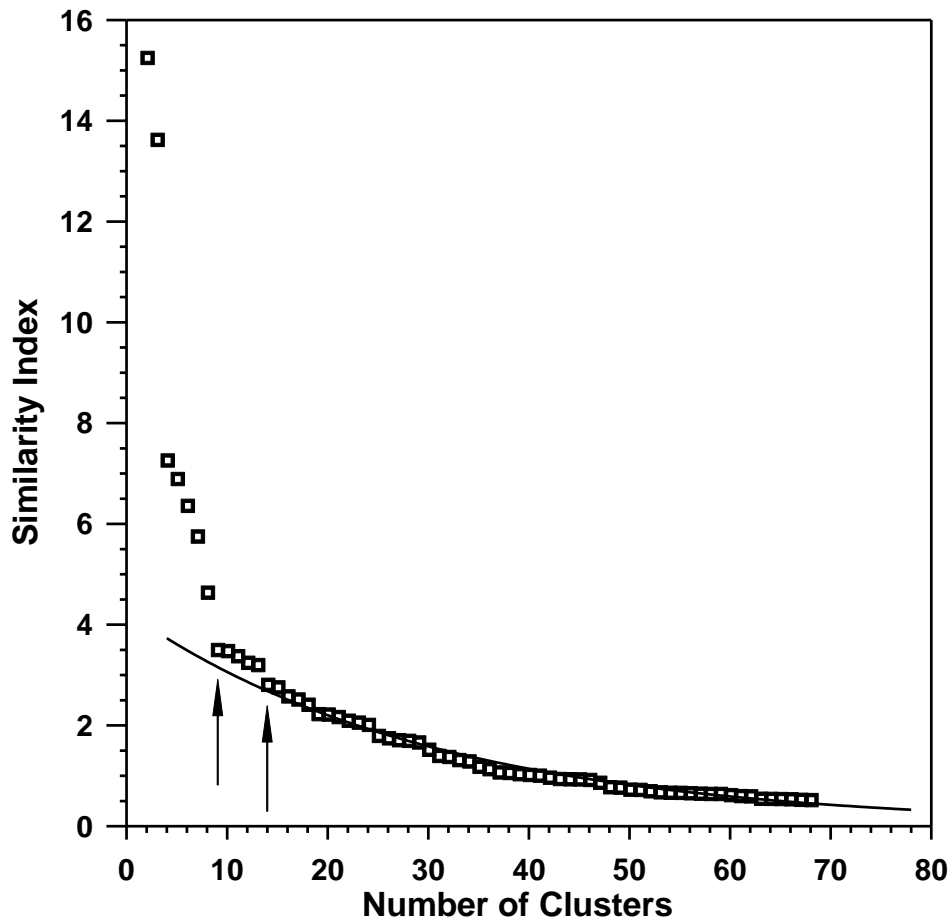


Figure 7

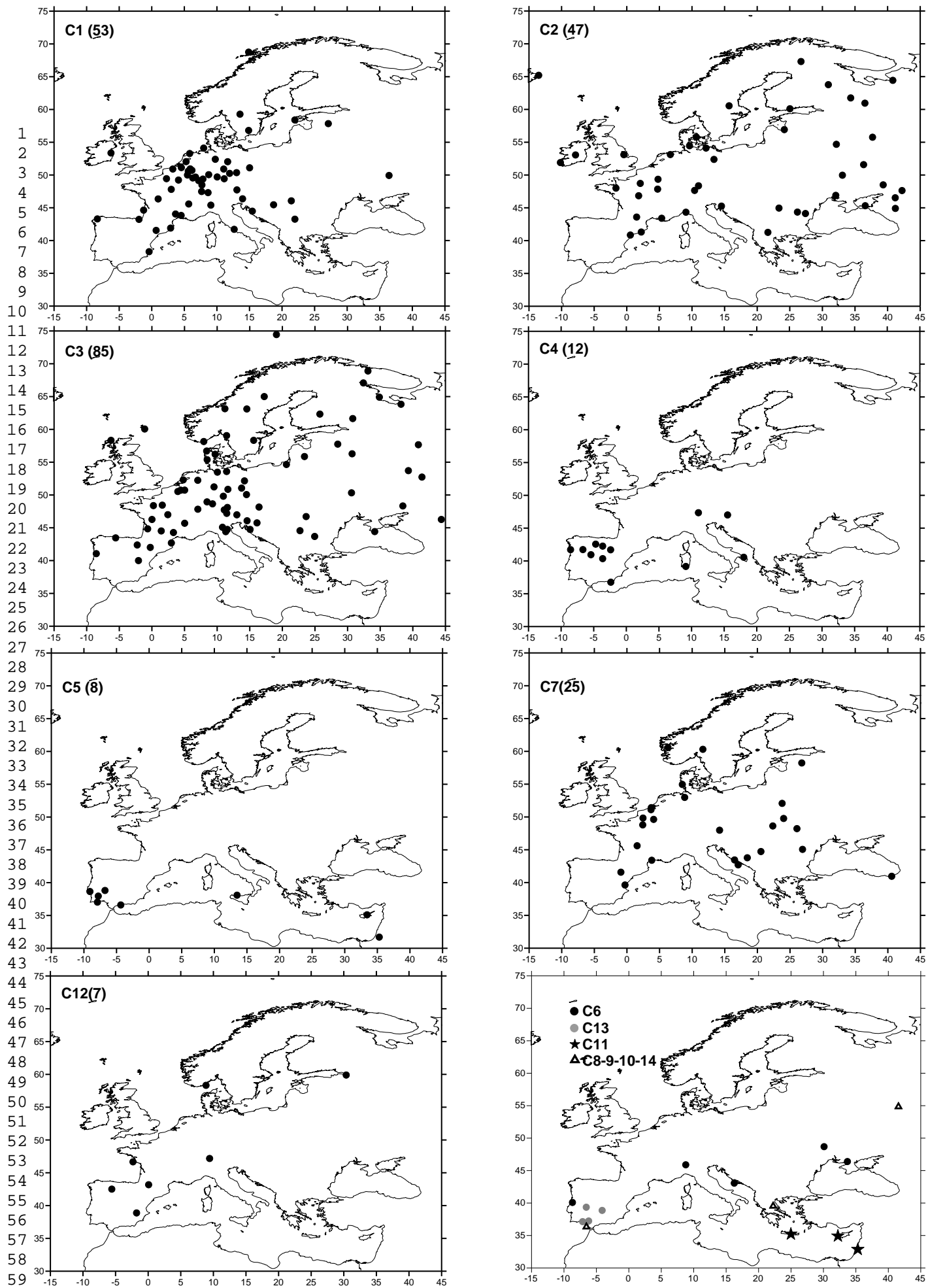
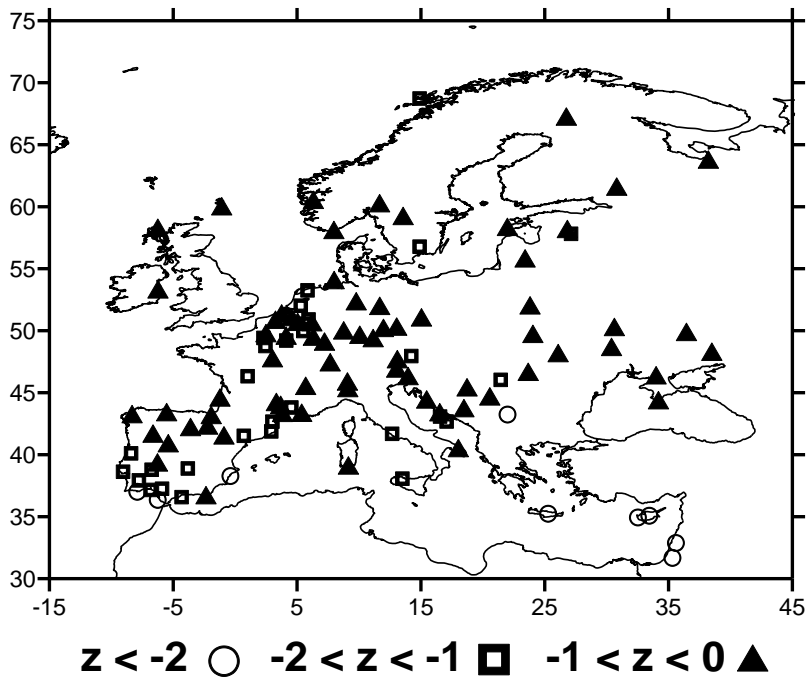


Figure 8

### Complexity Index, CI

a)



b)

### Complexity Index, CI

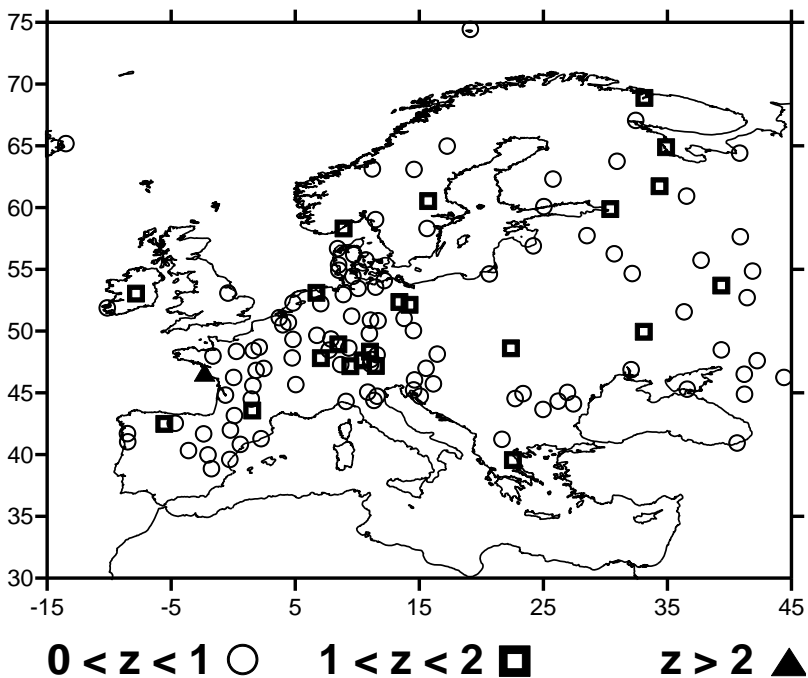
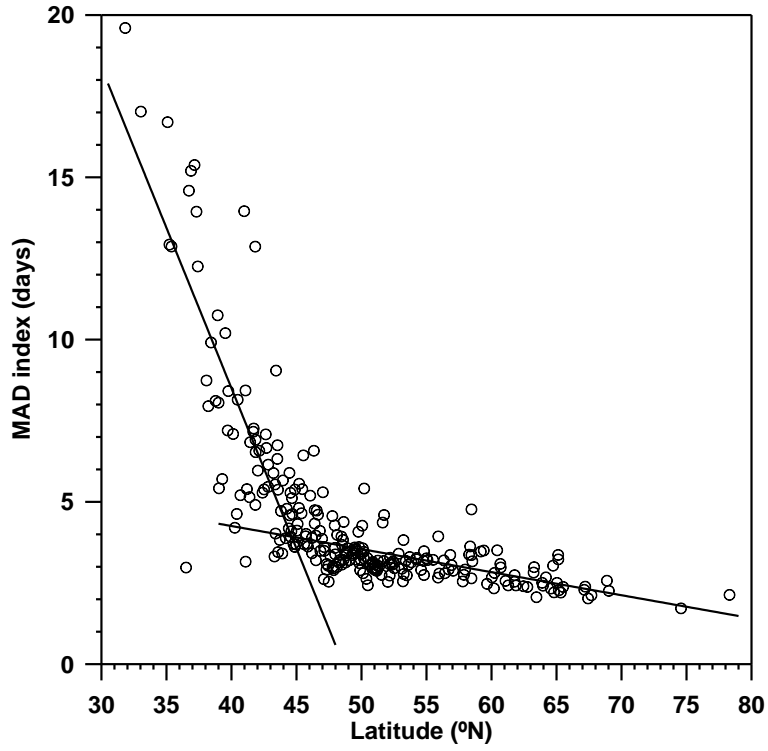


Figure 9

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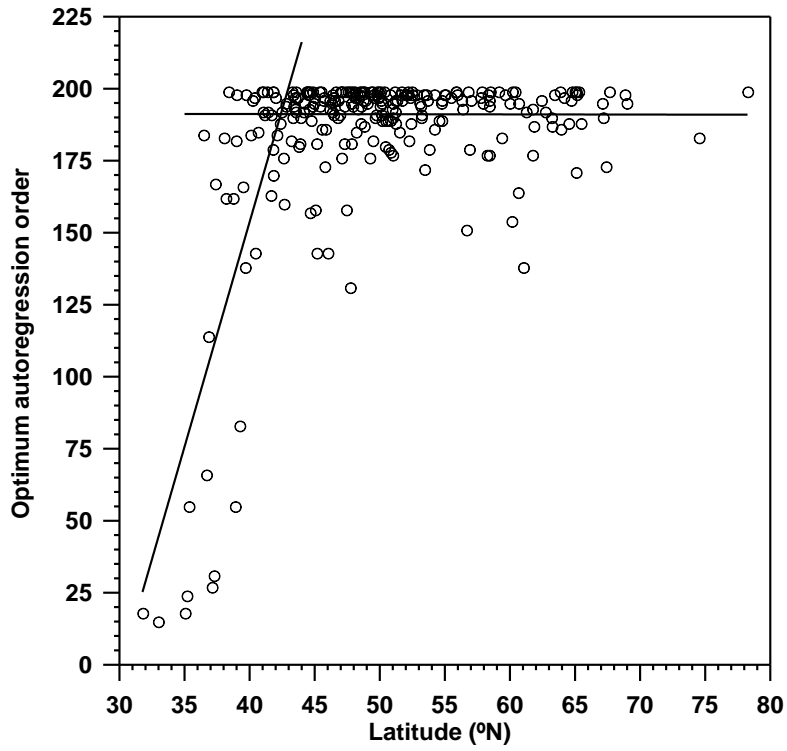
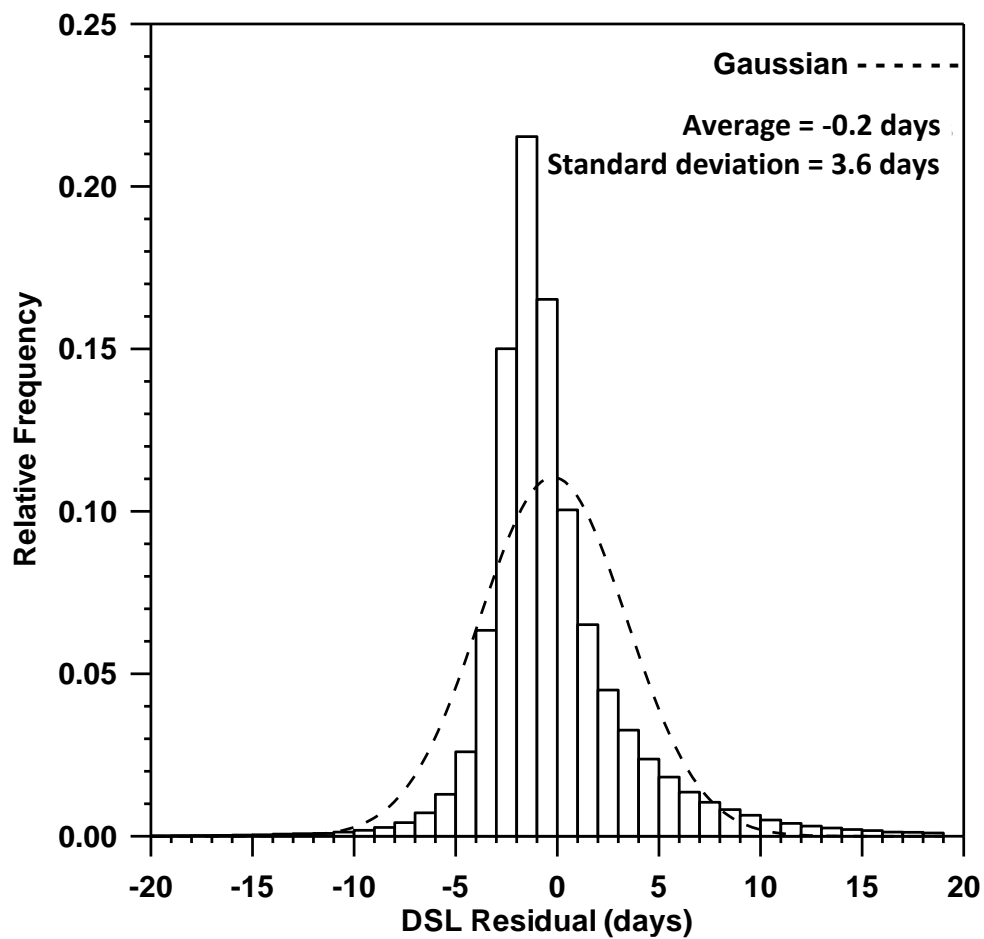


Figure 10

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**Figure 11**

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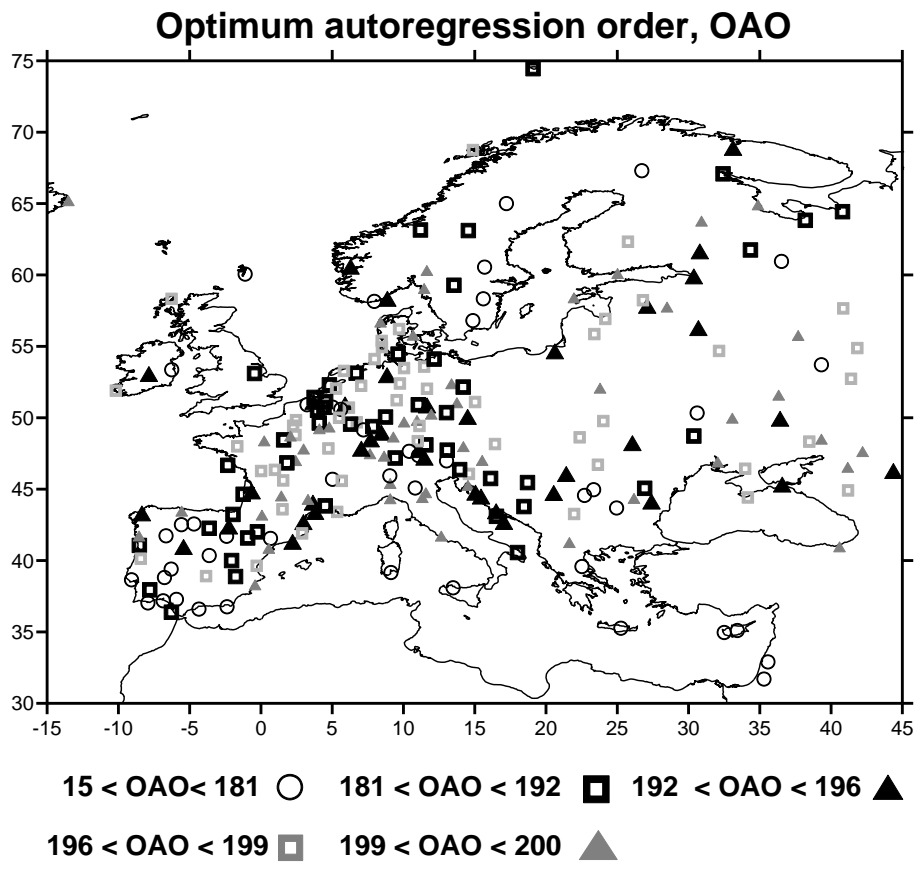
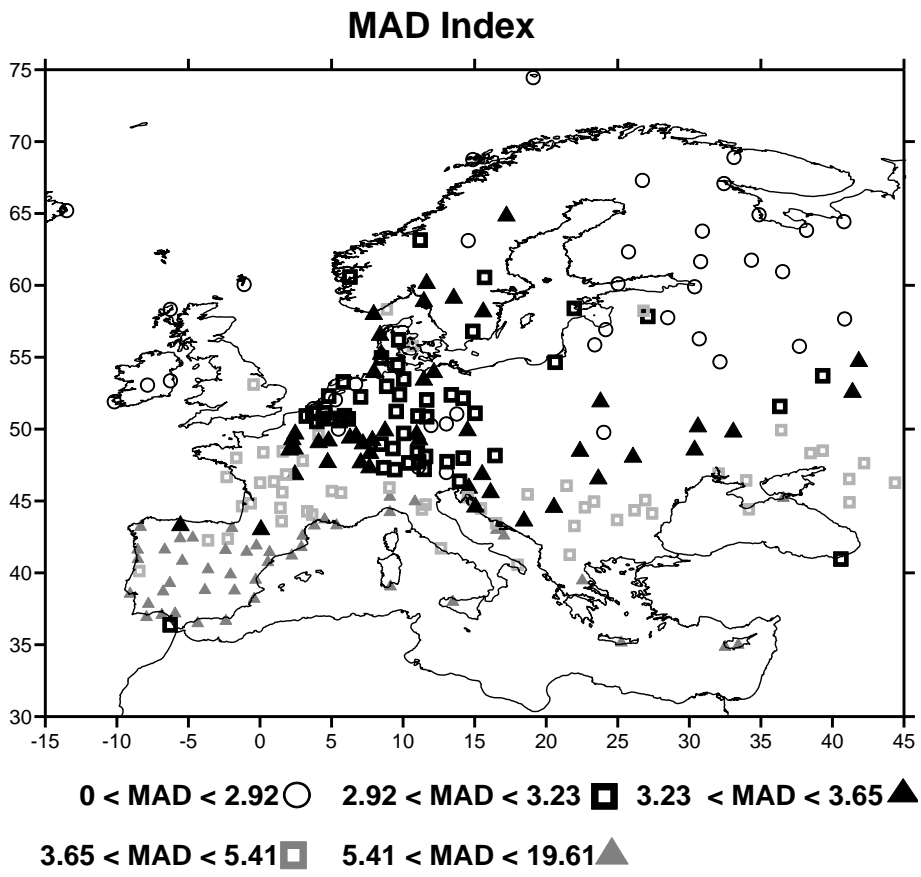
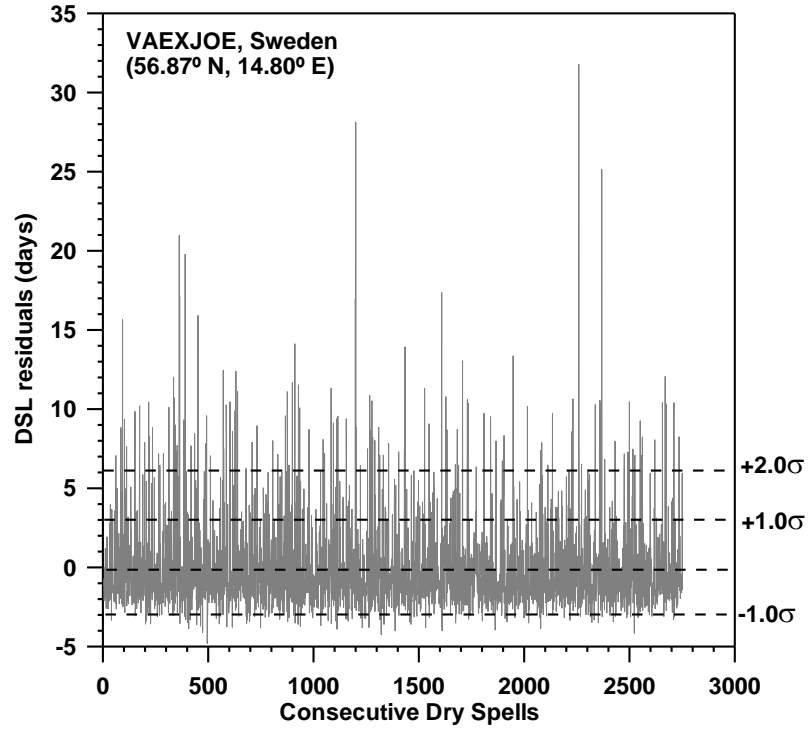
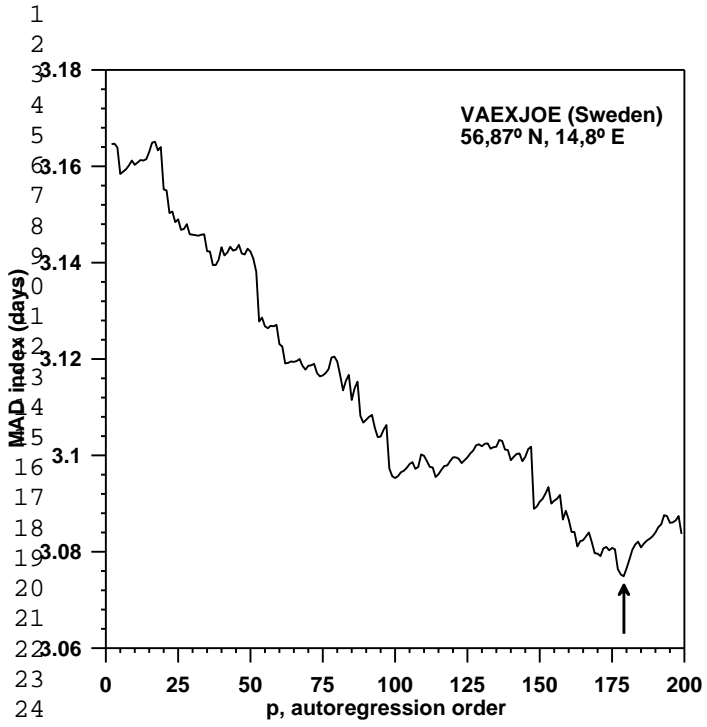


Figure 12

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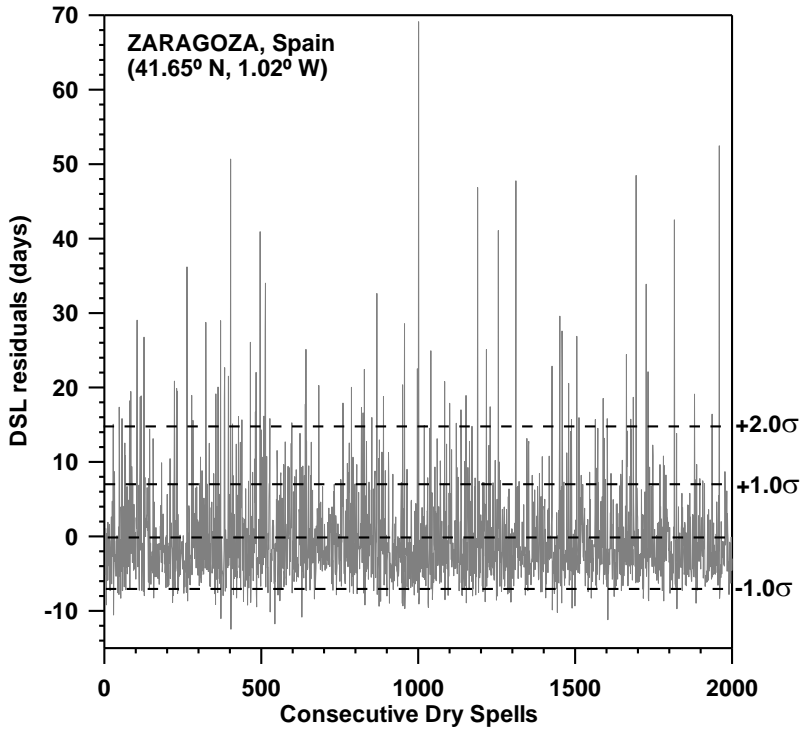
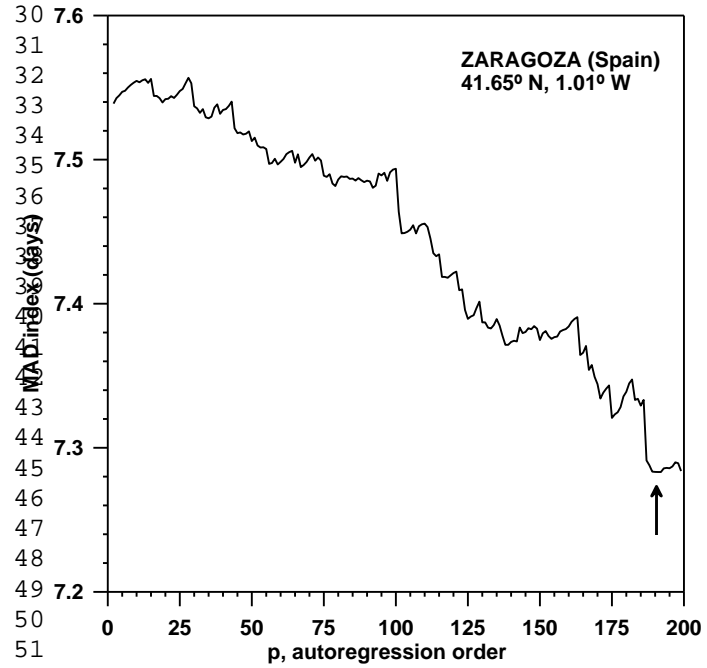
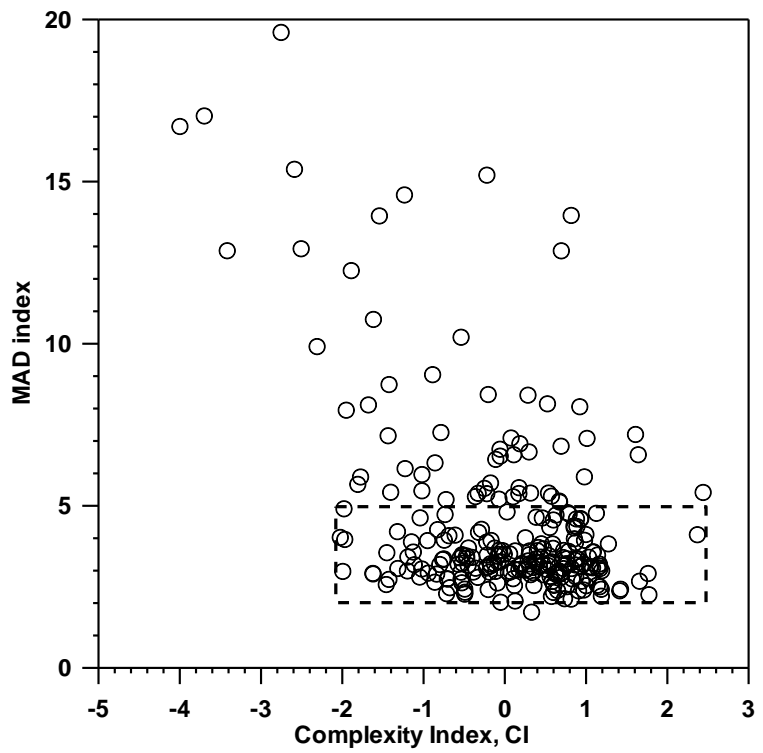


Figure 13

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(b)

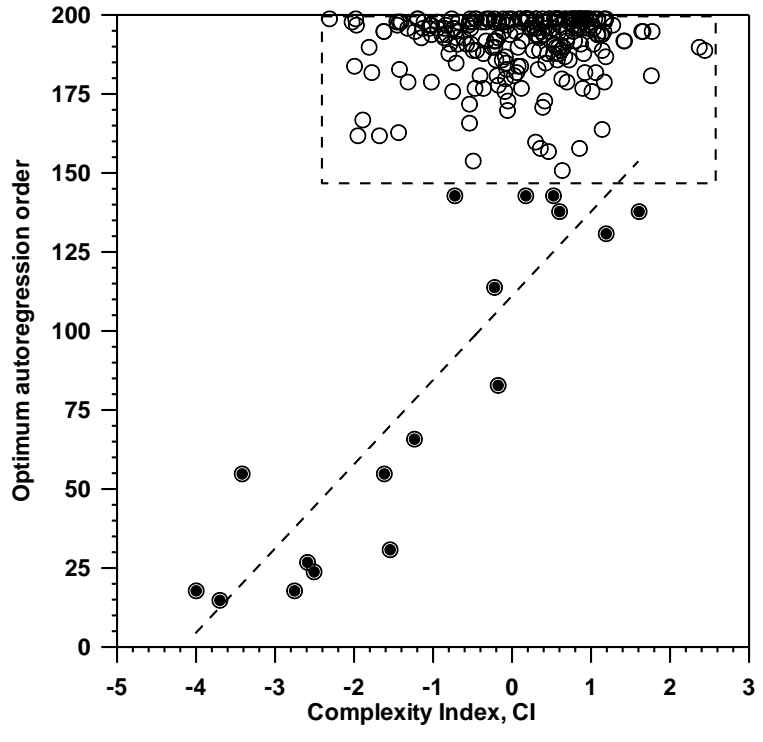


Figure 14



Manuscript TAAC-D-14-00581: Multifractality and autoregressive processes of dry spell lengths in Europe: An approach to their complexity and predictability (Authors: X. Lana, A. Burgueño, C. Serra and M.D. Martínez)

#### Response to Reviewer 1:

- a) In agreement with recommendations of the Reviewer and the Editor, the **Section 3.1** has been notably shortened, being only introduced the basic concepts concerning the MF-DFA algorithm and the appropriate references (see also **lines 162-165**). The structure of **Section 3.2** (Singularity spectrum) has not been changed, as many definitions of parameters and concepts used after are summarised in this Section. Additionally, in agreement with Reviewer 2, a detailed discussion about the range of  $q^{\text{th}}$ -order has been added to this Section.
- b) Effectively, the interpolation method has been the “inverse distance”. We think that in our case the main problem of obtaining accurate plots of the spatial distribution of fractal parameters is the limited rain gauge density in some areas, as mentioned now in **page 5 (lines 118-124)**.
- c) In **Section 5, lines 575-580**, several alternatives to the AR(p) process are cited as possible improvements on DSL prediction for very long DSL. Certainly, an ARIMA(p,d,q) modelling and others methods based on the Poisson distribution and Monte Carlo algorithms would improve the mentioned very long DSL prediction, but we think they are beyond the scope of the present paper. It is also worth of mention that the relatively simple AR(p) process has led to good results when predicting monthly Western Mediterranean Oscillation index, as the authors of this manuscript have found (manuscript nowadays submitted to the International Journal of Climatology). With respect to these questions, we have to mention that in **line 224-225** a mistake concerning the definition of AR(p) has been amended. AR(p) has to be properly defined as an ARIMA(p,d,q) with  $d=1$  and  $q=0$ .

#### Response to Reviewer 2:

- a) With the aim of a more complete description of the DSL series, **new Figure 1c** includes a histogram of the number of DSL,  $N_{\text{DSL}}$ , for the 267 DSL series. Additionally, two very different examples of DSL series (Vaexjoe, Sweden, and Almeria, Spain) are shown in a **new Figure 1d**. The corresponding comments are developed in **Section 2, lines 136-143**.
- b) A discussion about the appropriate range of the  $q^{\text{th}}$ -order is developed in **Section 3.2, lines 203-217**. The paper suggested by the reviewer (\*), the assumption that multifractal spectra should be fitted to a quadratic function taking as argument the Hölder exponent, and previous experience of the authors about this question have been used as reference points of this discussion.
 

*(\*) Ivanov P. Ch., Nunes Amaral, L.A., Goldenberg A.L., Haulin S., Rosenblum M.G., Stanley, H.E., Struzik, Z.B. (2001). From  $1/f$  noise to multifractal cascades in hearthbeat dynamics. *Chaos*, 11, 641-652.*
- c) The process to obtain a configuration of 14 clusters is described with more detail in **Section 4.3, lines 340-356**. The explanation is partially based on the concept of similarity index,  $L_{ij}$ , which is defined and quantified by the **new Equation 9**.
- d) With respect to **Figures 8**, due to a technical problem with the writing software, the two first figures were plotted without appearing “C” codes designing the cluster number. The problem has been now solved.