# GAMMA-RAY SPECTROSCOPY IN THE VICINITY OF ${ }^{108} \mathrm{Zr}^{*}$ 

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The half-lives of $2_{1}^{+}$states were measured for ${ }^{102,104} \mathrm{Zr}$ and ${ }^{106,108} \mathrm{Mo}$ to test a new implementation of a $\mathrm{LaBr}_{3}(\mathrm{Ce})$ array at the RIBF, RIKEN, Japan. The nuclei of interest were produced through the fission of a

[^1]$345 \mathrm{MeV} /$ nucleon ${ }^{238} \mathrm{U}$ beam and selected by the BigRIPS separator. Fission fragments were implanted into the WAS3ABi active stopper, surrounding which, $18 \mathrm{LaBr}_{3}(\mathrm{Ce})$ detectors provided fast $\gamma$-ray detection. Timing between the $\mathrm{LaBr}_{3}(\mathrm{Ce})$ array and plastic scintillators allowed for the measurement of half-lives of low-lying states. The preliminary results, which agree with literature values, are presented along with experimental details.

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## 1. Introduction

The $A \sim 100, Z \sim 40$ region of the nuclear chart has long been known for its sudden onset of static quadrupole deformation at $N \sim 60$ [1]. This was first ascribed to the neutron-proton interactions of the spatially-overlapping spin-orbit partner orbits, $\pi g_{7 / 2}$ and $\nu g_{9 / 2}$ [2]. However, more recent calculations [3] and $g$-factor [4] measurements have underlined the importance of core polarisation and the influence of the low- $\Omega \nu h_{11 / 2}$ orbitals.

The complexity of the factors which drive deformation in the neutron-rich zirconium region require stringent testing. The reproduction of energy levels can provide some evidence that the wave-function employed in calculations is correct, however, observables, such as the reduced transition probabilities serve as a more robust test.

In these proceedings, we present the measurement of the known half-lives of the $2_{1}^{+}$states in ${ }^{102,104} \mathrm{Zr}$ and ${ }^{106,108} \mathrm{Mo}$ through $\beta-\gamma$ spectroscopy. From these, the $B\left(\mathrm{E} 2 ; 2_{1}^{+} \rightarrow 0_{\mathrm{g} . \mathrm{s} .}^{+}\right)$values are computed.

## 2. Experimental set-up

A decay spectroscopy experiment was carried out at the RI Beam Factory (RIBF). The in-flight fission of a ${ }^{238} \mathrm{U}^{86+}$ primary beam of average intensity $6.24 \times 10^{10}$ particles/s accelerated to an energy of $345 \mathrm{MeV} /$ nucleon produced a secondary beam of neutron-rich nuclides. Fission fragments were selected by the BigRIPS spectrometer using the $B \rho-\Delta E-B \rho$ method [5] and identified using TOF $-B \rho-\Delta E$ measurements [6].

The secondary beam was implanted into the WAS3ABi silicon array [7], which detected ion implantations and their subsequent $\beta$-decays. Precise timing of $\beta$-electron emission was achieved using plastic scintillators of 2 mm thickness and area $65 \times 45 \mathrm{~mm}^{2}$ installed upstream and downstream of WAS3ABi. An array of $18 \mathrm{LaBr}_{3}(\mathrm{Ce})$ [8] detectors, as well as the EURICA [9]
array, surrounded WAS3ABi for the purpose of measuring isomeric and $\beta$-delayed $\gamma$-rays. The photopeak efficiency of the $\mathrm{LaBr}_{3}(\mathrm{Ce})$ array at $\sim 150 \mathrm{keV}$ was measured to be $4 \%$.

## 3. Experimental results

Implanted fission fragments were correlated with their $\beta$-decays by requiring that the $\beta$-decay had to occur in the same pixel as an implanted ion within approximately five times the $\beta$-decay half-life of the implanted nuclide. The $\beta$-electron was required to be detected in one of the $\beta$-plastics. As an example, the $\gamma$-ray energy spectrum of ${ }^{106} \mathrm{Mo}$ is shown in the right panel of Fig. 1, the inset shows the background-subtracted time-difference, $\Delta T$, spectrum of the $2_{1}^{+} \rightarrow 0_{\text {g.s. }}^{+}$transition ( $\Delta T$ is the time between a signal in the $\beta$-plastics and a $\gamma$-ray detection in the $\mathrm{LaBr}_{3}(\mathrm{Ce})$ array).



Fig. 1. Left: Particle identification plot with labels indicating the nuclides of interest. Right (a): Energy spectrum measured in the $\mathrm{LaBr}_{3}(\mathrm{Ce})$ array within 1 s of a $\beta$-decay correlated to an implantation of ${ }^{106} \mathrm{Nb}$. (b): Background-subtracted $\Delta T$ projection of the $2_{1}^{+} \rightarrow 0_{\mathrm{g} \text {.s. }}^{+}$transition, the curve (blue) is an exponential fit.

To extract the half-life of the $2_{1}^{+}$states for ${ }^{104,106} \mathrm{Zr}$ and ${ }^{106,108} \mathrm{Mo}$, an exponential fit was carried out on the delayed shoulder of the $\Delta T$ spectrum between 2 and 15 ns . The results, presented in the left of Fig. 2 agree with adopted values $[10,11,13,14]$, with the exception of ${ }^{102} \mathrm{Zr}$. This deviation is tentatively attributed to the influence of the half-life of a $K^{\pi}=4^{-}$state [12]. For the four nuclei under discussion, no delayed component was observed for the feeding $4_{1}^{+} \rightarrow 2_{1}^{+}$transitions. The right panel of Fig. 2 shows the $B\left(\mathrm{E} 2 ; 2_{1}^{+} \rightarrow 0_{\mathrm{g} . \mathrm{S} .}^{+}\right)$values obtained.

The presented method shall be extended to more neutron-rich isotopes in the region, to extend the knowledge of transition probabilities.


Fig. 2. Left: Half-lives of the $2_{1}^{+}$states as a function of neutron-number for Mo and Zr. Right: The corresponding $B\left(\mathrm{E} 2 ; 2_{1}^{+} \rightarrow 0_{\mathrm{g} . \mathrm{s} .}^{+}\right)$transition probabilities. In both, the solid symbols are values determined in this work, open symbols from Ref. [10, 11, 13, 14].

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