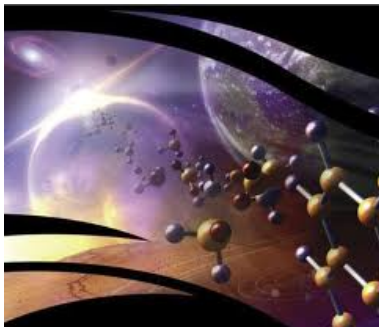


# Pompage chimique de $CH^+$ et $OH^+$

**Octavio Roncero**

Inst. Física Fundamental, CSIC

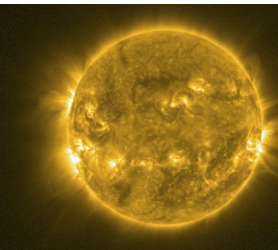
[octavio.roncero@csic.es](mailto:octavio.roncero@csic.es)



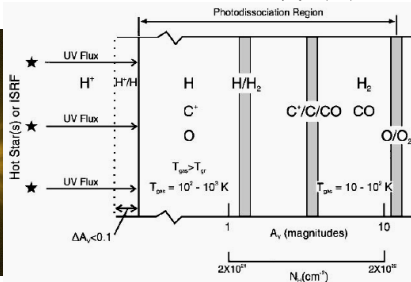
# Outline

- 1 Introduction
- 2  $C^+ + H_2(v,j)$
- 3  $O^+ + H_2(v,j)$
- 4  $OH^+ + H$
- 5 Conclusions

# Molecular Universe



Hollenbach & Tielens, Annu. Rev. Astrophys. ('97)



Stellar atmosphere:

atomic lines

Atom+Atom collisions

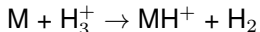
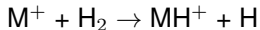
Magnetic fields

Polarization of lines

Hanle effect

PDR

Formation of  $H_2$   
act as shield for other  
molecules



Molecular clouds

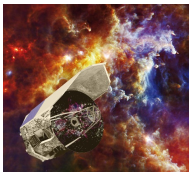
low temperatures

varying densities

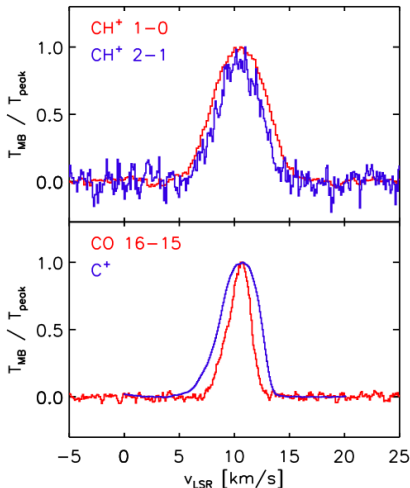
formation of complex  
molecules

# Molecules as probes

- Hydrogen is  $> 70\%$  in mass
- The rest of elements detected in smaller fractions
- Detection of infrared-microwave individual transitions



CH, CH<sup>+</sup>, OH<sup>+</sup>, SH, SH<sup>+</sup>, ...



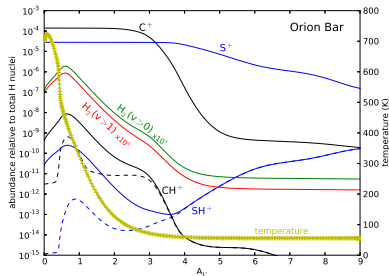
by Nagy, *et al.*, A & A (2013)

# Molecules as probes

- Hydrogen is > 70 % in mass
- The rest of elements detected in smaller fractions
- Detection of infrared-microwave individual transitions
- Flux from different excitations

Beyond Maxwell-Boltzmann distributions  
for short-lived species

**Need for final state formation rate**



M. Agúndez, PDR model  
(Medon code)

Astrophysical conditions  
density, temperature, etc

# Chemical pumping in molecular clouds and PDR's

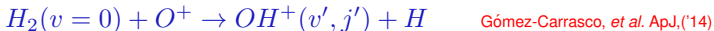
- Exothermic reactions of  $H_2$  with atoms and ions



Gómez-Carrasco, *et al.* ApJ, ('14)

# Chemical pumping in molecular clouds and PDR's

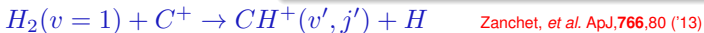
- Exothermic reactions of  $H_2$  with atoms and ions



- Chemistry of  $H_2(v > 0)$  in ISM

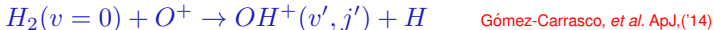
Agúndez, et al. ApJ, 713,662 ('10)

**Initial state dependent chemistry**



# Chemical pumping in molecular clouds and PDR's

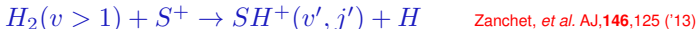
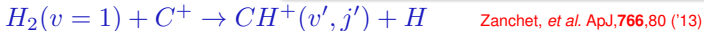
- Exothermic reactions of H<sub>2</sub> with atoms and ions



- Chemistry of H<sub>2</sub>(v > 0) in ISM

Agúndez, et al. ApJ, 713,662 ('10)

**Initial state dependent chemistry**



- The most abundant ion is H<sub>3</sub><sup>+</sup>

- H<sub>2</sub><sup>+</sup> disappears in the exothermic  $H_2 + H_2^+ \rightarrow H_3^+(v) + H$  reaction

6D PES, Sanz-Sanz, et al. JCP, 139, 184302 ('13)

- However  $H_2 + H_3^+ \rightarrow H_3^+ + H_2$  (ortho/para conversion, deuteration)

9D PES, Aguado, et al. JCP, 133, 024306 ('10)

QCT biased statistical model, Gomez-Carrasco, et al. JCP, 137, 094303 ('12)

- H<sub>3</sub><sup>+</sup> is very reactive with other species:  $H_3^+ + O \rightarrow H_2 + OH^+$



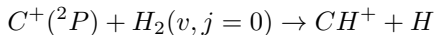
# Outline and acknowledgements

- Reactive collisions:  $A+BC \rightarrow AB+C$ 
  - Endothermic from excited vibrational states  $C^+ + H_2(v)$   
*A. Zanchet, B. Godard, N. Bulut, P. Halvick and J. Cernicharo*
  - Exothermic from excited vibrational states  $O^+ + H_2(v)$   
*S. Gómez-Carrasco, B. Godard, F. Lique, N. Bulut  
J. Kloss, A. Aguado, F.J. Aoiz, J. F. Castillo  
J. R. Goicoechea, M. Etxaluze and J. Cernicharo*
- Inelastic vs. exchange collisions:
  - $OH^+ + H \rightarrow H + OH^+(v'J')$  in two PES's  
*N. Bulut and F. Lique*

# Outline

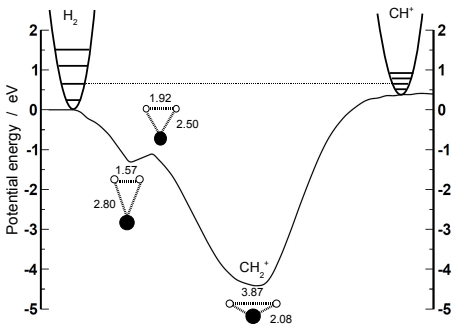
- 1 Introduction
- 2  $C^+ + H_2(v,j)$**
- 3  $O^+ + H_2(v,j)$
- 4  $OH^+ + H$
- 5 Conclusions

# $C^+ + H_2$ : PES

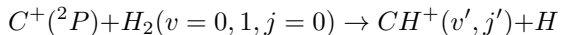


Ground state PES by  
Stoeklin & Halvick, PCCP('05)

- Deep well  $\approx 4.5$  eV
- Endothermic by  $\approx 0.38$  eV
- No Quantum dynamical calculation.



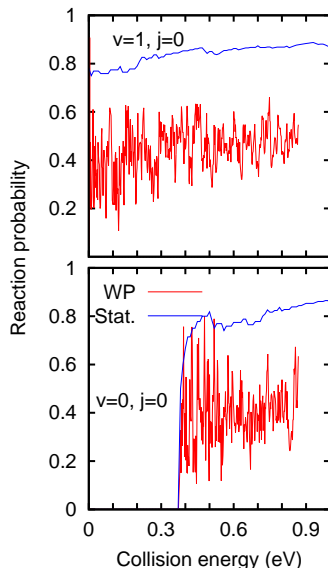
# C<sup>+</sup> + H<sub>2</sub>: Total Reaction Probabilities for $J=0$



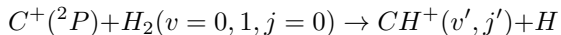
Full dim. Wave packet dynamics

reactant Jacobi coord. with MADWAVE3

- Dense grids: about 105 Mpoints (including hecity  $\Omega$ )
- Many resonances: CH<sub>2</sub><sup>+</sup> complex  
long propagation times
- For H<sub>2</sub>(v=0): threshold at  $\approx 0.38$  eV
- For H<sub>2</sub>(v=1): no threshold
- Statistical model **does not** work properly



# $C^+ + H_2$ : Integral Reaction Cross section

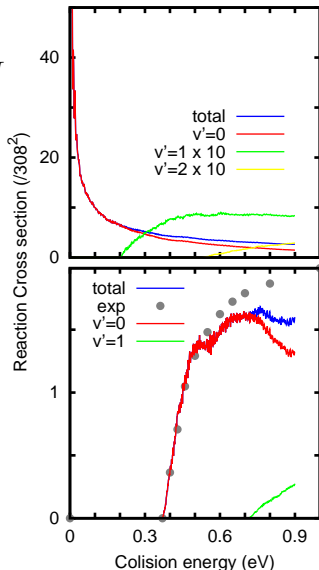


$J=0, 1, 2, \dots, 20$

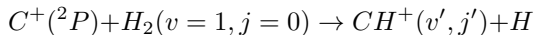
$J=20, 25, 30, \dots, 50$

J-shifting interpolation for  $J > 20$

- Experimental data:  
Gerlich, Disch & S. Scherbarth, JCP ('87)  
in arbitrary units
- Inclusion of electronic partition function:  
0.407 at 300K  
with the  $C^+(^2P_{1/2,3/2})$  spin-orbit splitting  
 $64 \text{ cm}^{-1}$



# C<sup>+</sup> + H<sub>2</sub>: State-to-state rate constants



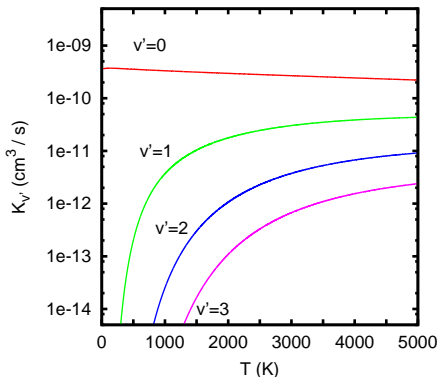
- No threshold for  $j' < 8$

Langevin extrapolation

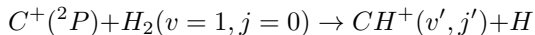
to avoid WP errors

- State-to-state coefficients

$$K_{vj \rightarrow v'j'}(T) = \sqrt{\frac{8}{\pi\mu(k_B T)^3}} \times \int_0^\infty E dE \sigma_{vj \rightarrow v'j'}(E) e^{-E/k_B T}$$



# $C^+ + H_2$ : State-to-state rate constants



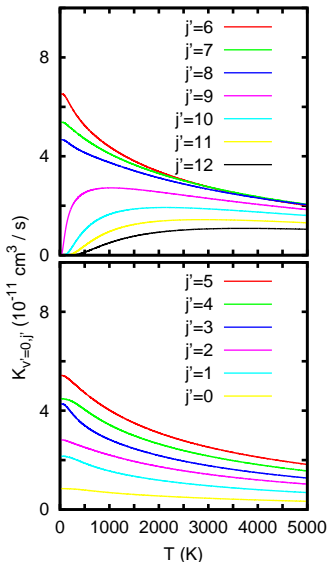
- No threshold for  $j' < 8$

Langevin extrapolation

to avoid WP errors

- State-to-state coefficients

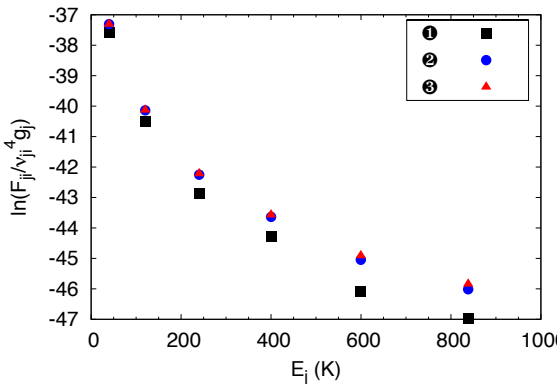
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# Astrophysical models

- PDR model of Orion bar (Madex code)
- Intensity of some lines of  $CH^+(j' > 0)$
- Improves previous models
- Still discrepancies

Zanchet *et al.*, *ApJ* ('13) 766:88



**Chemical pumping increases flux from higher  $j$ 's**

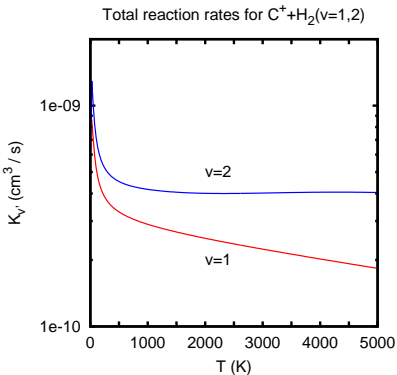


# Possible improvements

- Higher  $v > 1$  contributions still need to be included in PDR models

Zanchet *et al.*, AJ ('13) 146:125

also  $S^+ + H_2(v > 2)$



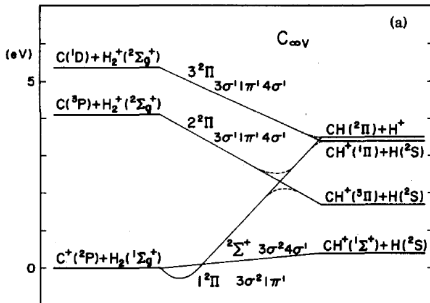
# Possible improvements

- Higher  $v > 1$  contributions still need to be included in PDR models

- Only  $^2P_{1/2}$  reacts

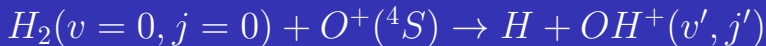
$$Q_e(T) = \frac{2}{2 + 4e^{-91,2/T}}$$

- Non-adiabatic and spin-orbit transitions

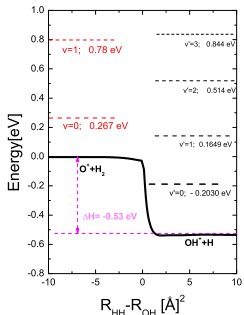


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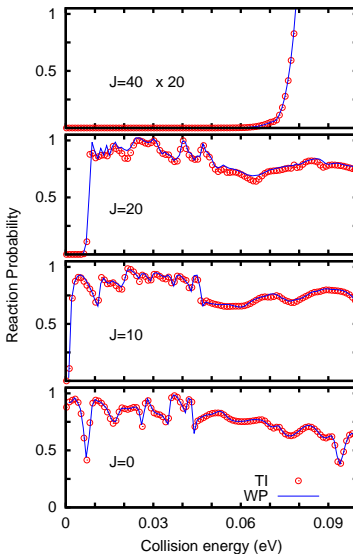


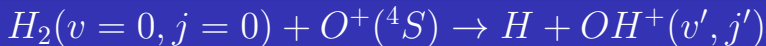
PES: Martínez, *et al.*, JCP ('04)



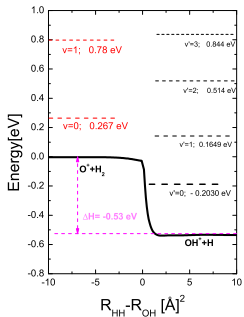
$$P^J(E) = \sum_{v'j'} |S_{v,j \rightarrow v'j'}|^2 \rightarrow$$

**Good agreement with TI-ABC**  
 **$E \geq 0.1$  meV !!**





PES: Martínez, *et al.*, JCP ('04)



$$P^J(E) = \sum_{v'j'} |S_{v,j \rightarrow v'j'}|^2 \rightarrow$$

# J-shifting approach

- J-shifting approach: Bowmann ('85)

$$P_J(E) = P_{J=0}(E^*) \quad \text{with} \quad E^* = E - BJ(J+1)$$

overestimates reaction probabilities for  $J \gg 0$

- J-shifting interpolation:

$$P_J(E) = \frac{J - J_1}{J_2 - J_1} P_{J_1}(E_1) + \frac{J_2 - J}{J_2 - J_1} P_{J_2}(E_2) \quad \text{with} \quad J_1 < J < J_2$$

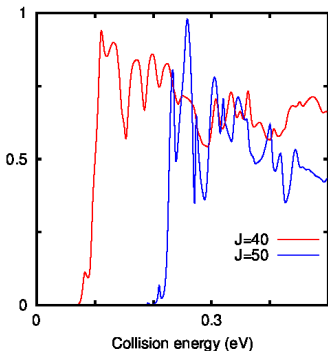
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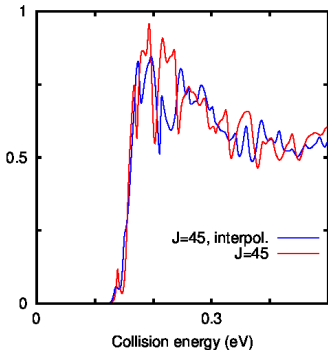
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overestimates reaction probabilities for  $J \gg 0$

- J-shifting interpolation:





# Cross section: $H_2(v = 0, j = 0) + O^+$

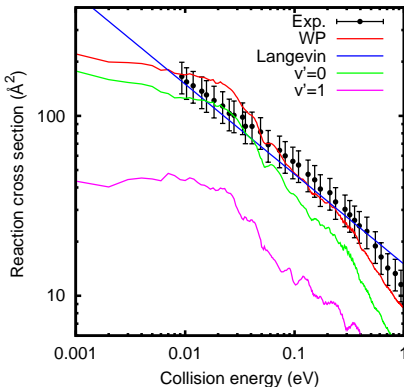
$$\sigma_{vj} = \frac{\pi}{k^2} \sum_J (2J + 1) P^J(E)$$

Experiment:

Burley, Ervin, Armentrout, ('87)

Langevin model works for total  $\sigma_{vj}$

but not for individual  $\sigma_{vj \rightarrow v'j'}$  !!



# State-2-state rates: $H_2(v = 0, j = 0) + O^+$

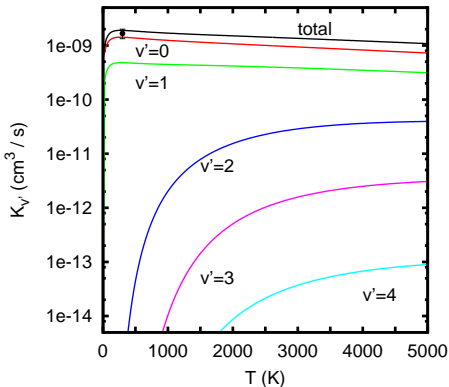
Gomez-Carrasco *et al.*, ApJ 794:33 ('14)

$$K(v, j, v', j')(T) = \left[ \frac{8}{\pi \mu (k_B T)^3} \right]^{1/2}$$

$$\times \int_0^\infty E \sigma_{vj \rightarrow v'j'}(E) e^{-E/k_B T} dE$$

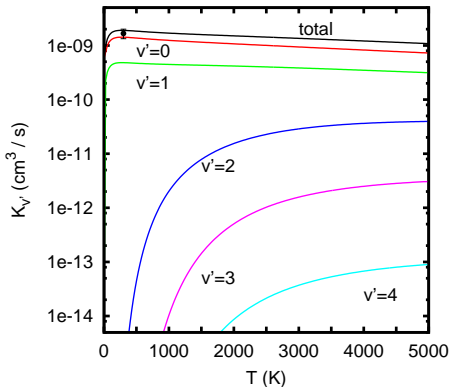
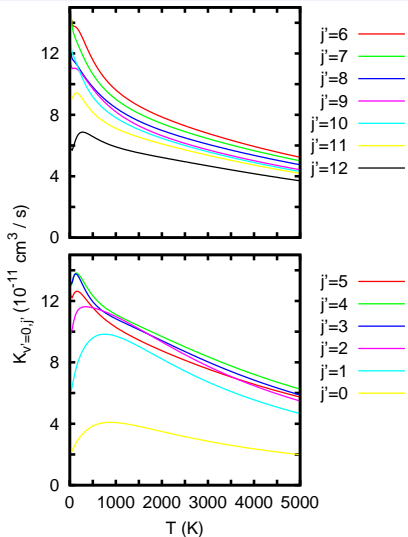
Experiment:

Burley, Ervin, Armentrout, ('87)

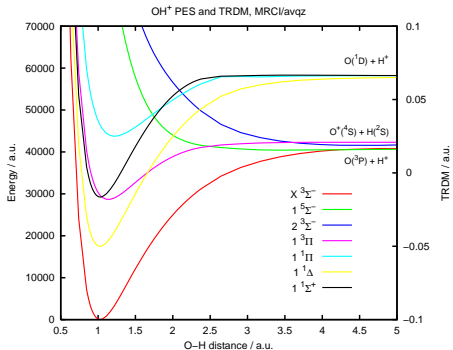


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Gomez-Carrasco *et al.*, ApJ 794:33 ('14)

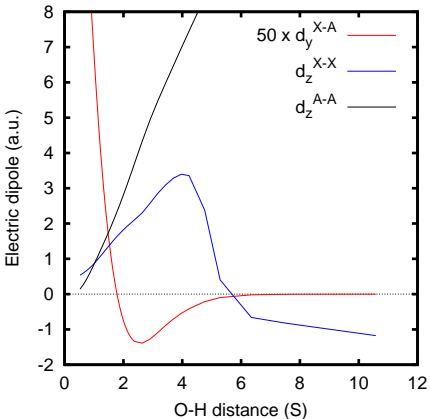


# $OH^+(X^3\Sigma^-, A^3\Pi)$ radiative rates



- $OH^+(X^3\Sigma^-)$ : Hund's case b
- $OH^+(A^3\Pi)$ : Hund's case a
- Reasonable agreement with experiment  
*Merer et al. ('75)*

# $OH^+(X^3\Sigma^-, A^3\Pi)$ radiative rates



## Radiative Lifetimes of A: $\tau_b$

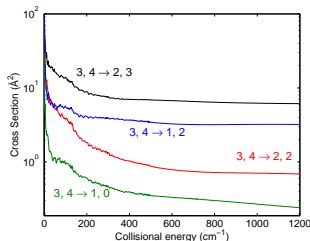
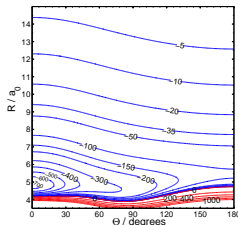
$v$	$\tau_v$ (ns)	$\tau_v^a$ (ns)	$\tau_v^b$ (ns)
0	2524	2410	$2400 \pm 300$
1	2665	2560	
2	2820	2930	
3	3004		
4	3233		
5	3534		
6	3960		
7	4637		
8	5961		
9	9559		
10	16118		

<sup>a</sup> Merchan *et al.* ('91)

<sup>b</sup> Möhlman *et al.* ('78)

# $OH^+(X^3\Sigma^-)$ collisional rates with He

- New PES for  $He + OH^+(X^3\Sigma^-)$   
F. Lique and J. Kloss
- TI-CC calculations of inelastic rates  
F. Lique and J. Kloss
- $H + OH^+$  and  $H_2 + OH^+$  rates were scaled



# $OH^+(X^3\Sigma^-)$ in dense and hot PDR

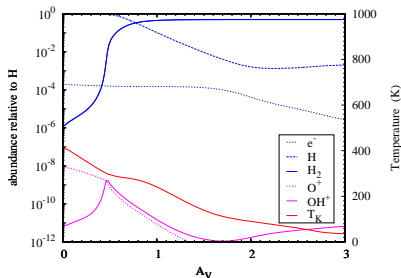
Meudon PDR chemical model under 3 conditions:

B. Godard

(a) only non-reactive collisions

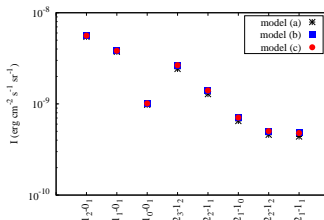
(b) chemical pumping, products according to Boltzmann distribution at 2000K

(c) chemical pumping using WP state to state rates



$$\chi = 10^4 \text{ and } n = 10^4 \text{ cm}^3$$

Gomez-Carrasco *et al.*, ApJ 794:33 ('14)



-  $N < 3$  driven by  $OH^+ + H$  collisions  
 - Chemical pumping for  $N > 3$  (?)

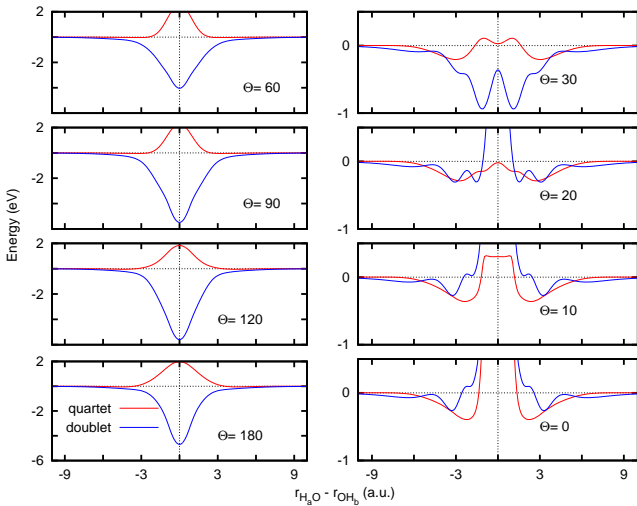
# Outline

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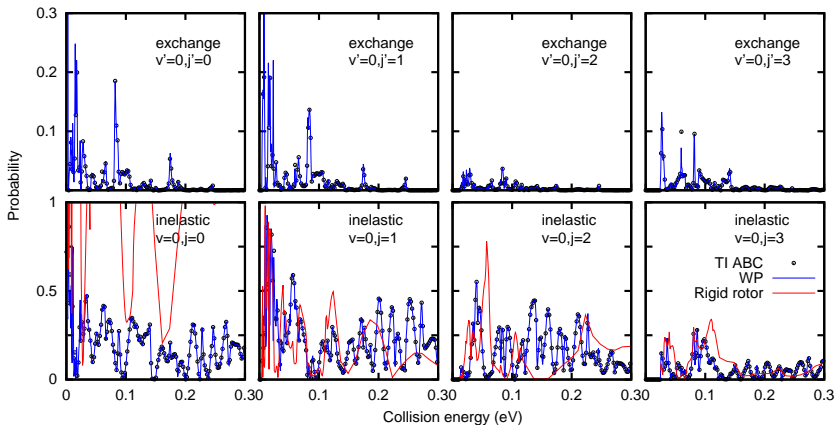




# MEP's for exchange: two mechanisms

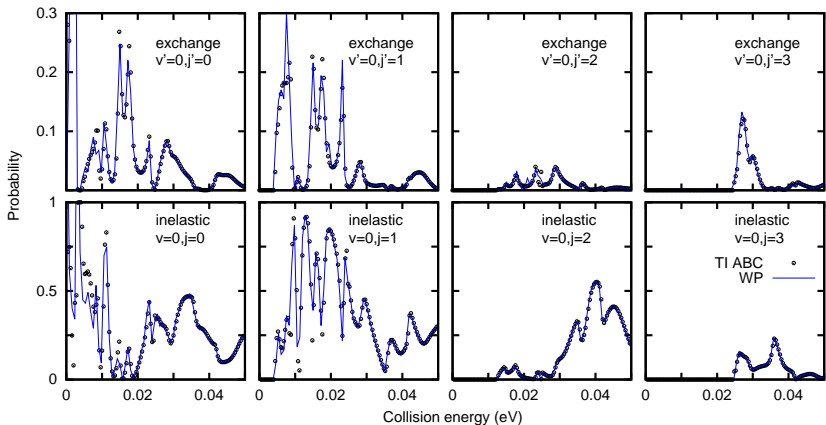


# Quadruplet: $H + OH^+(v = 0, j = 0, J = 0)$



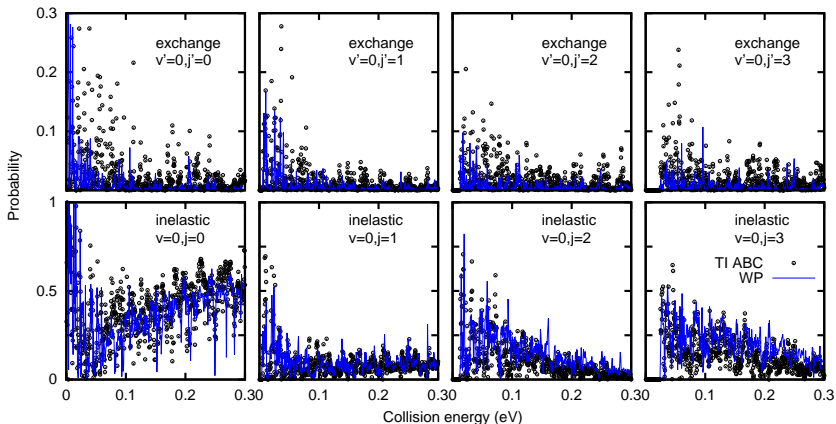
- Good agreement WP vs. TI-ABC
- Rigid rotor of the same order

# Quadruplet: $H + OH^+(v = 0, j = 0, J = 0)$



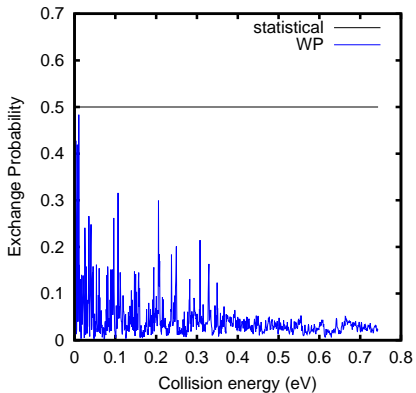
- Good agreement WP vs. TI-ABC, even at rather low energies!!
- Rigid rotor of the same order

# Doublet: $H + OH^+(v = 0, j = 0, J = 0)$



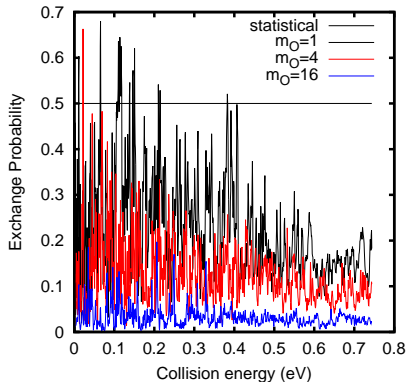
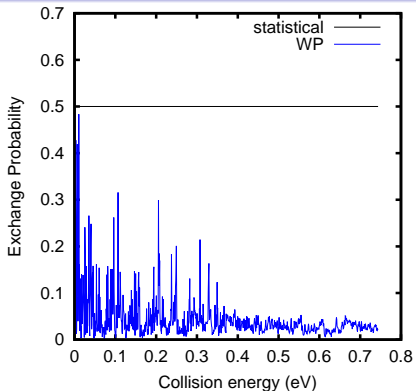
- Many resonances
- Comparison difficult using different coordinates
- Even ABC has problems when  $D_{H-OH^+} \leq D_{OH^+}$

# Is the exchange statistical in the doublet state?



**Non statistical  
energy transfer inefficient**

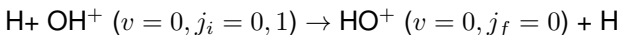
# Is the exchange statistical in the doublet state?



**Non statistical  
energy transfer inefficient**

**Energy redistribution increases  
- as energy decreases  
- and mass difference reduces**

# Cross sections: inelastic and exchange for quadruplet



## For quadruplet:

8  $10^4$  iterations

J=0,5,10,15,20,25,30,40,..., 110

$\Omega_{max} = 15$

$j_i = 0, 1$

## For doublet:

2.5  $10^5$  iterations and denser grids

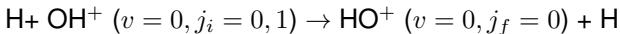
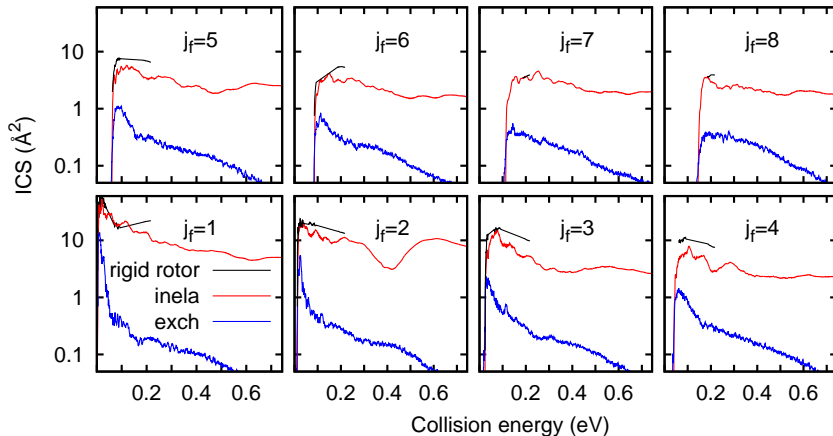
Calculations still in progress

$\Omega_{max} = 19$

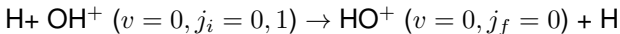
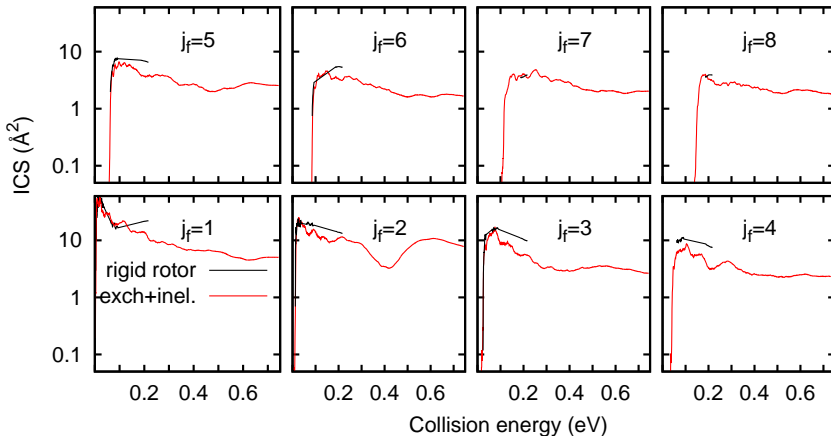
$j_i = 0$



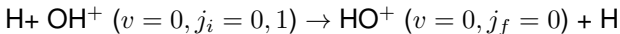
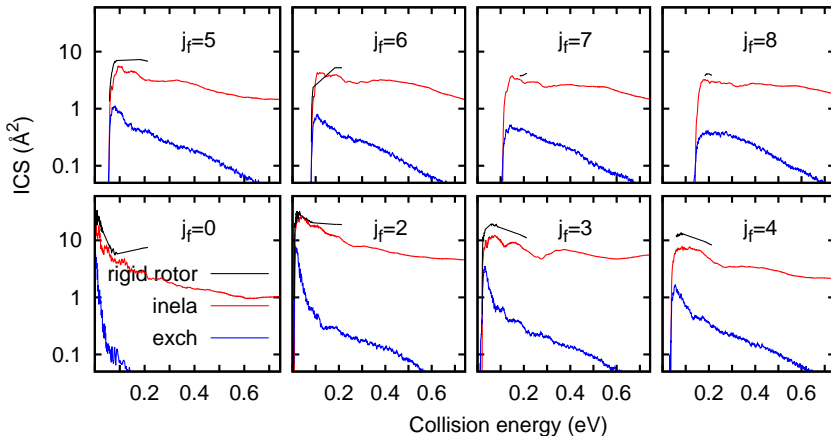
# Cross sections: inelastic and exchange for quadruplet

 $j_i = 0$ 

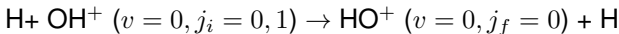
# Cross sections: inelastic and exchange for quadruplet

 $j_i = 0$ 

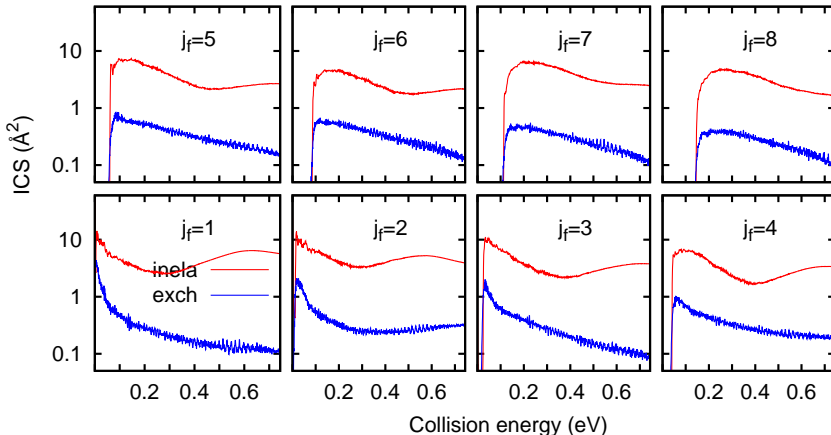
# Cross sections: inelastic and exchange for quadruplet

 $j_i = 1$ 

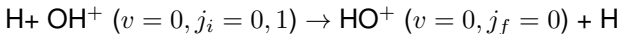
# Cross sections: inelastic and exchange for doublet



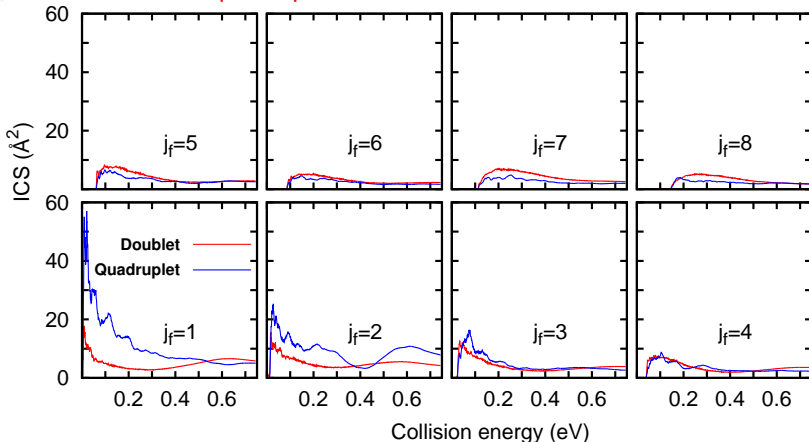
$j_i = 0$ : Preliminary



# Cross sections: inelastic and exchange for doublet



$j_i = 0$ : Doublet vs. quadruplet



# Outline

- 1 Introduction
- 2  $C^+ + H_2(v,j)$
- 3  $O^+ + H_2(v,j)$
- 4  $OH^+ + H$
- 5 Conclusions**

# Conclusions and perspectives

- Chemical pumping improves the description of the flux of relatively high rotational excitations of hydrides
- For low rotational excitations, inelastic scattering with H and H<sub>2</sub> becomes dominant
- Collisions with H(<sup>2</sup>S) imply several electronic states for open shell hydrides like OH<sup>+</sup>(<sup>3</sup>Σ<sup>-</sup>)
- Exchange reactions have a considerably lower cross section than inelastic collisions
- In the quadruplet case, with a relatively shallow well, the agreement between rigid rotor and exact (including inelastic+exchange) is very good.  
This has to be checked for the doublet state with a deep insertion well.
- A proper electronic partition function need to be included in open shell systems
- Also the recoupling with the electronic spin should be done to provide results comparable with the observations