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1   **Three dimensional quantification of soil hydraulic properties using X-ray Computed**

2   **Tomography and image based modelling**

3

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17   **Key Points:**

18   • X-ray Computed Tomography is used to image water distribution in soil.

19   • 3D Computed Tomography images are combined with image based modelling.

20   • Unsaturated hydraulic conductivity and the water release characteristic are obtained.

21

1. **Abstract**

We demonstrate the application of a high-resolution X-ray Computed Tomography (CT) method to quantify water distribution in soil pores under successive reductive drying. We focus on the wet end of the water release characteristic (WRC) (0 to -75 kPa) to investigate changes in soil water distribution in contrasting soil textures (sand and clay) and structures (sieved and field structured), to determine the impact of soil structure on hydraulic behaviour. The 3D structure of each soil was obtained from the CT images (at a 10 μm resolution). Stokes equations for flow were solved computationally for each measured structure to estimate hydraulic conductivity. The simulated values obtained compared extremely well with the measured saturated hydraulic conductivity values. By considering different sample sizes we were able to identify that the smallest possible representative sample size which is required to determine a globally valid hydraulic conductivity.

Keywords: Matric potential; soil pores; water release characteristic; X-ray Computed Tomography; image based homogenisation.

Abbreviations:

(3D) - 3-dimensional

(CT) - Computed Tomography

(ROI) - Region of interest

(WRC) – Water release characteristic

45   (WFP) – Water filled pores

46   (AFP) – Air filled pores

47   (REV) -Representative elementary volume

48

49   **Short title for page headings:** Three-dimensional quantification of soil hydraulic properties

50

51   ## 2. Introduction

52   Understanding the dynamic nature of water movement through, and storage within, soil is of

53   fundamental importance for most of its functions. A more detailed knowledge of how soil

54   structure influences soil water distribution and subsequently availability to plants is needed to

55   inform soil management practices and thus contribute to efforts in increasing crop yields for

56   an expanding global population. However, to truly understand how water moves through soil

57   and is retained in soil pores whilst undergoing drying, non-destructive measurements of the

58   soil aggregate geometries and pore structure are needed. The use of X-ray Computed

59   Tomography (CT) to examine the 3D soil pore structure is rapidly gaining pace [*Helliwell et*

60   *al.,* 2014; *Mooney et al.,* 2012]. The ability to segment water from a greyscale CT image

61   remains challenging due to limitations in the phase separation achievable with regular detectors

62   in benchtop CT scanners. Segmentation of fluids in alternative porous media such as glass

63   beads has been previously described in the literature [*Haugen et al.*, 2014; *Iassonov et al.*,

64   2009]. However, due to the more complex, heterogeneous structure of soil and the associated

65   challenges faced during image acquisition this brings, these segmentation techniques cannot

66   always be readily applied to the segmentation of soil phases in X-ray CT images [*Houston et*

67   *al.*, 2013]. Yet, Crestana *et al.* [1985] highlighted the potential of CT for investigating water

68   movement through soil by quantifying the vertical and lateral water flows in 3D. By using dual

3

73 energy tomography Rogasik *et al.* [1999] showed it is possible to quantify soil water, pore

74 space and bulk density. Mooney [2002], compared dry and moist samples and developed an

75 imaging method to separate water from the bulk soil in undisturbed soil cores albeit at a coarse

76 resolution (0.5 mm) using a medical CT scanner. Water flow characteristics were well matched

77 to macropore structure, with *circa* 90% of macropore space active in water transport in a sandy

78 clay texture, compared to *circa* 50% in the sandy loam soil. Wildenschild *et al.* [2005]

79 investigated multiphase flow processes and quantified drainage paths, showing that water is

80 preferentially lost from larger pores and the drainage of the remaining disconnected pores is

81 prevented. Since then Tippkotter *et al.* [2009] detected soil water in macropores and measured

82 water film thickness at a 20 μm resolution. They quantified the water without the use of contrast

83 enhancing agents or comparison of wet and dry scans. Partial volume effects make separating

84 boundaries between phases of interest a challenging task, due to the material of concern often

85 failing to fully fill a voxel [*Ketcham and Carlson*, 2001]. This can lead to individual voxels

86 being misclassified. However, as a microfocus system with a fine resolution was used in the

87 study, the impact of voxel misclassification is minimised [*Clausnitzer and Hopmans*, 2000].

88 The most recent advances in CT imaging technology now allow 3D non-destructive imaging

89 of soils at even higher resolutions, *e.g.* <3 μm spot size [*Zappala et al.*, 2013], allowing water

90 to be observed in smaller pores than previously considered.

91

92 The macroscopic scale flow in saturated and unsaturated porous media are often described

93 using a Darcy's law and Richards' equation [*Hornung*, 1997].   These equations are

94 parameterised by the unsaturated hydraulic conductivity and WRC. From a mathematical

95 perspective Darcy's law can be derived from the underlying Stokes equations using the method

96 of homogenization [*Keller*, 1980]. This approach has been applied to single porosity materials

97 [*Hornung*, 1997; *Keller*, 1980], double porosity materials [*Arbogast and Lehr*, 2006; *Panfilov*,

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105 2000], and porous media containing voids or vugs [*Arbogast and Lehr*, 2006; *Daly and Roose*,

106 2014]. This method is based on the idea that the underlying porous structure is periodic, *i.e.*, it

107 is composed of a set of regular repeating units. By calculating the average fluid velocity for a

108 single representative unit volume due to a unit pressure drop, Darcy's law and the

109 representative hydraulic conductivity can be derived [*Pavliotis and Stuart*, 2008; *Taylor et al.*,

110 1970]. The mathematics underlying the WRC is less well established and homogenization

111 studies to date have assumed that the air water interface is known in advance [*Taylor et al.*,

112 1970].

113

114 Image based modelling can be loosely divided into two categories: pore network modelling

115 and direct simulation [*Blunt*, 2001; *Blunt et al.*, 2013]. The first of these, pore network

116 modelling, refers to the extraction of a representative network from the pore scale geometry

117 [*Fatt*, 1956] and has been widely used to predict averaged properties of packed spheres [*Bryant*

118 *and Blunt,* 1992; *Bryant et al.,* 1993b] and imaged porous media [*Blunt, 2001; Blunt et al.,*

119 *2013; Bryant et al.,* 1993a]. This technique is able to reproduce relative permeability curves

120 and water release characteristics. However, the pore network extraction results in a simplified

121 geometry which may neglect important pore scale phenomena. The alternative technique of

122 direct modelling involves solving equations directly on the imaged geometries. This technique

123 captures the detail of the pore scale geometry down to the resolution limit.   The key

124 disadvantage is that modelling multiphase flow is demanding from a computational point of

125 view. Typically models of this type are based on Lattice Boltzman simulations of two fluids as

126 these are highly parallel and relatively easy to implement [*Gao et al.*, 2012; *Ramstad et al.*,

127 2010]. Other image based modelling studies in porous media include 3D image based

128 simulations of nutrient transport [*Keyes et al.*, 2013] and saturation based flow modelling of

129 water movement in 2-D image slices [*Aravena et al.,* 2014; *Aravena et al.,* 2011].

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130

In this paper we demonstrate the application of a high resolution CT method to quantify water distribution in soil pores under successive reductive drying. Using the structural geometries from the CT images we apply the method of homogenization combined with direct image based modelling to calculate the hydraulic conductivity across a range of matric potentials. This can be seen as analogous to the measurement of a wetting phase relative permeability curve. Specifically the question we aimed to answer is: what is the smallest possible Representative Elementary Volume (REV) required to determine a hydraulic conductivity approximation which is globally valid? In addition the soil moisture content is calculated directly from the CT images at a range of different matric potentials and, hence, the WRC is derived. This approach is chosen, rather than Lattice Boltzmann simulation, as it allows the WRC and hydraulic conductivity to be calculated quickly without the need for time consuming two fluid simulations. From the images we are able to obtain information on soil drainage in both undisturbed field cores and sieved soils. This was used to accurately determine the impact of soil structure on hydraulic behaviour.

145

3. Materials and Methods

147

3.1. Sample preparation

Soil was obtained from The University of Nottingham farm at Bunny, Nottinghamshire, UK (52.52° N, 1.07° W). The soils used in this study were a Eutric Cambisol (Newport series, loamy sand/sandy loam) and an Argillic Pelosol (Worcester series, clay loam). Particle size analysis for the two soils was: 83% sand, 13% clay, 4% silt for the Newport series and 36% sand, 33% clay, 31% silt for the Worcester series. Typical organic matter contents were 2.3%

155 for the Newport series and 5.5% for the Worcester series [*Mooney and Morris, 2008*]. From

156 herein the two soil types are referred to as sand and clay soil. Loose soil was collected from

157 each site in sample bags and field cores (10 mm height, 10 mm diameter) were collected from

158 the immediate soil surfaces after any residues had been cleared. Four replicates were collected

159 for each soil type. Saturated hydraulic conductivity measurements of the field cores were

160 obtained via the standard laboratory method using a constant head device [*Rowell*, 1994].

161 Experimentally obtained hydraulic conductivity measurements were made for comparison with

162 the calculated values.

163

164 *3.2. Measurement of the soil water release characteristic*

165 In order to investigate the effect of drying on soil water distribution and soil structural changes

166 in 3D, a custom-built tension table was designed to hold the soil sample at a given matric

167 potential, whilst undergoing CT scanning. A small vacuum chamber was constructed

168 (Supplementary Figure 1), that contained a porous ceramic plate (Soil Moisture Corp, Santa

169 Barbara, CA, U.S.A) on top of which a field soil core was placed, with kaolin clay at the base

170 to ensure a good contact. The sample size was kept small to ensure high resolution scanning

171 based on the sample size:resolution trade off that limits CT studies [*Wildenschild et al.*, 2002].

172 The porous ceramic was first submerged in de-aired deionized water and a vacuum applied to

173 ensure no air bubbles remain trapped within the ceramic. The plastic chamber had an O-ring

174 seal at the base and a flange, which was screwed together to ensure an air-tight fit. All

175 components of the chamber were made from plastic to avoid possible image artefacts, which

176 could result from using high X-ray attenuating construction materials such as metal. A 0387

177 Millipore vacuum pump (Merck Millipore, MA, USA) was attached to the chamber and the

178 soil columns were initially saturated from below with deionized water before being placed

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181 under a vacuum of -5 kPa (50.8 mmHg), -10 kPa (76.2 mmHg), -20 kPa (152.4 mmHg), -40

182 kPa (304.8 mmHg), -60 kPa (457.2 mmHg) and -75 kPa (533.4 mmHg). It was not possible to

183 achieve pressures in excess of -75 kPa in the system described. The vacuum pump was enabled

184 for 120 min then the valve sealed to retain the vacuum inside the chamber. At each matric

185 potential the soil core inside the chamber was scanned (further details in next section). After

186 each scan the soil core was removed from the chamber and weighed to calculate gravimetric

187 water content.

188

189 Using a combination of sand tables, pressure plates and pressure membrane apparatus, the

190 WRC were obtained for both soil types. Using conventional methods the van Genuchten model

191 [*van Genuchten*, 1980] was fitted to the data. The sand table was prepared by filling the water

192 reservoir and raising it above the base height of the table to ensure full saturation of the sand

193 table with no air bubbles. Both undisturbed field cores and sieved samples (packed to a bulk

194 density of 1.2 g cm^{-3}) were placed flat on the sand table. To obtain the mid-range points on the

195 WRC, a pressure plate Model 1600 Pressure Plate Extractor (Soil Moisture Corp, Santa

196 Barbara, CA, U.S.A) was used. Samples were placed on the plate and the samples were

197 weighed frequently until equilibrating at a series of matric potentials. To obtain measurements

198 at lower matric potentials, a pressure membrane apparatus was used. For the sieved samples

199 the required mass of soil was carefully placed into the metal vessels of the pressure membrane

200 apparatus and saturated with air-free water. The field core samples were directly collected in

201 the metal cores and placed onto the apparatus. The collection tubes were weighed frequently

202 and once equilibrated the system was adjusted to higher pressures. After the final measurement,

203 soil samples were oven dried at 105 °C for 24 hr then weighed.

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209 The measured volumetric water content of the field structured soil cores (θ), accounted for the

210 likely presence of stones and their overall influence on the WRC [*Gardner*, 1965; *Reinhart*,

211 1961]. The volumetric water content of the soil is written as θ which is measured in $\frac{cm^3(H_2O)}{cm^3(Soil)}$.

212 To calculate this a correction was applied to the measured volumetric water content of the

213 sample, θ_m, which is measured as $\frac{cm^3(H_2O)}{cm^3(Total)}$. This stone correction was calculated using

$$\theta = \theta_m \left(1 + \frac{V_s}{V_f}\right), \tag{1}$$

214 where V_f is the fine soil volume measured in $cm^3(soil)$ and V_s is the volume of stones which

215 is measured in $cm^3(stones)$ and calculated by subtracting the soil volume from the total

216 sample volume.

217

218 *3.3. X-ray Computed Tomography and analysis*

219 Four replicates from each soil type of the field cores were scanned at the seven matric potentials

220 giving a total of 56 scans for the field structured cores. As they would not be used for individual

221 pore characterisation analysis only three replicates were scanned for the sieved soils leading to

222 42 additional scans and a total overall number of 98 scans. The field of view for each scan

223 included the entire sample and each scanned sample created a dataset approximately 25

224 gigabytes in size, which includes all the associated image stacks for analysis, therefore 2.4

225 terabytes of data was collected. X-ray CT scanning was performed using a Phoenix Nanotom

226 180NF (GE Sensing & Inspection Technologies GmbH, Wunstorf, Germany). The scanner

227 consisted of a 180 kV nanofocus X-ray tube fitted with a diamond transmission target and a 5-

228 megapixel (2316 x 2316 pixels, 50 x 50 μm pixel size) flat panel detector (Hamamatsu

229 Photonics KK, Shizuoka, Japan). A maximum X-ray energy of 100 kV and 140 μA was used

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to scan each soil core. A total of 1440 projection images were acquired over a 360° rotation. Each projection was the average of 3 images acquired with a detector exposure time of 1 s. The resulting isotropic voxel edge length was 10.17 μm and total scan time was 105 minutes per core. Although much faster scan times are possible it was necessary in this instance to use a longer scan time to acquire the highest quality images to aid with the phase separation. Two small aluminium and copper reference objects (< 1 mm^2) were attached to the side of the soil core to assist with image calibration and alignment during image analysis. Reconstruction of the projection images to produce 3D volumetric data sets was performed using the software datos|rec (GE Sensing & Inspection Technologies GmbH, Wunstorf, Germany).

The reconstructed CT volumes were visualised and quantified using VG StudioMAX® 2.2 (Volume Graphics GmbH, Heidelberg, Germany). Air, soil and water phases of the scanned volume were segmented separately using a threshold technique based on the greyscale value of each voxel using a calibration tool within VG StudioMax v2.2. This works by selecting specific areas of the scanned volume based on greyscale values, which are a result of (X-ray attenuation) density differences within the sample. Image segmentation is the classification of voxels within a CT volume that share common grayscale values and thus X-ray attenuation. The calibration tool allows the user to sample the greyscale value of a selection of voxels that correspond to the background (*e.g.* the phase/s not considered, this is usually the background air) and then the process is repeated for voxels that correspond to the material of interest (*e.g.* water, air or soil). Based on the greyscale range we segmented all voxels above a 50% mean value between the background and material of interest and define them as a particular phase (i.e. a region of interest). In a two phase sample air is usually used as the background greyscale value to calibrate against. However, as we required the segmentation of three phases, two

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257   vessels, one containing soil pore water and the other finely sieved soil (< 100 μm) for either

258   the sand or clay soil were securely fitted to the inside of the chamber for each scan. The soil

259   pore water and finely sieved soil remained separate from the soil core, but within the imaging

260   field of view, which was important as the soil core sample can change in overall greyscale

261   range over the course of the experiment due to localised drying and X-ray filament aging. Using

262   this approach these separate vessels were used as reference objects during image analysis. It

263   was important that reference vessels were included in every scan as CT scanning is subject to

264   minor variations in greyscales between scans as the system filament ages. A schematic

265   representation of the water segmentation method is presented in Supplementary Figure 2. The

266   first stage was to create a region of interest (ROI) that included all phases except the air filled

267   pore space. This was done by selecting the air space around the sample as the background,

268   ensuring all voxels except those of air would be included as the material of interest and a 3D

269   ROI was then created and labelled 'Water and soil ROI' *e.g.* air not included. The process was

270   then repeated, but utilising the soil inside the reference vessel as a reference value for the solid

271   material. It was used as a reference as it contained limited water and air, hence this 3D ROI

272   was labelled 'Soil'. Subtraction of the 'Soil ROI' from the 'Water and soil ROI' resulted in a

273   ROI with voxels of grey scale values attributed to 'soil water only', *e.g.* the range of voxels

274   remaining after the soil and air ROIs have been subtracted. The volume of the resulting 'Soil

275   Water' ROI was validated against the water reference object and traditional methods of

276   determining volumetric water content (weighing). Using this approach we assume all organic

277   matter is classified as 'soil'. To create an ROI that included solely the air filled pore space, the

278   ROI that included all material except the air was inverted using the 'Invert ROI' function in

279   VG StudioMax software. The volume analyser tool in VG StudioMax was subsequently used

280   to quantify the total volume of the air and water filled pores. Reconstructed volumes for each

281   matric potential were aligned in VG StudioMax using the metal reference objects on the outside

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295 of the sample container. After segmentation of the soil core volumes, a cylindrical ROI shape

296 template was used to subtract the surrounding column and air space from the volume to remove

297 the cylinder of soil. This template was used to ensure the same volume was used for each

298 sample before analysis. No significant evidence of shrinkage was observed across the range of

299 low matric potentials considered here. We note that when three phases (air water and soil) are

300 present in an image erroneous films of the intermediate phase can often be segmented out

301 between the high and low density phases. In this case erroneous films of water could be

302 introduced through segmentation. However, as thin water films only contribute a total flow

303 proportional to the third power of the film thickness these are not a concern in this study.

304

305 Image stacks of the extracted volumes were exported and subsequently analysed for individual

306 water filled and air filled pore characteristics for the field structured soil only using ImageJ

307 v1.42 (http://rsbweb.nih.gov/ij/) [*Ferreira and Rasband*, 2011]. For 2D analysis objects less

308 than two pixels (twice the resolution) in diameter (0.02 mm) and for 3D analysis objects less

309 than two voxels in each direction ($8 \times 10^{-6} \text{ mm}^3$) were regarded as potential noise as a

310 precaution [*Wildenschild et al.*, 2005] and subsequently excluded from the analysis. The BoneJ

311 plugin algorithm [*Doube et al.*, 2010] (http://bonej.org/) in ImageJ software was used to

312 measure discrete individual 3D water filled and air filled pore characteristics namely volume,

313 surface area and thickness (diameter). A voxel is classed as connected to another voxel if it at

314 least touches corner to corner. ImageJ was used to measure the 2D pore shape characteristic

315 'circularity', which is a measure of an object's similarity to a circle. A value of 1 describes a

316 perfect circle and as the value decreases, the object becomes increasingly elongated. The

317 circularity was determined using

318
$$C = \frac{4\pi(A_p)}{(P_p)^2}, \tag{2}$$

12

327   where A_p is the pore area and P_p is the pore perimeter. The water-filled pore image stack was

328   skeletonised and from this the 3D connectivity (the sum of all the thin (1 voxel) pathways that

329   still preserve the connected topology and original shape of the object) of the water-filled pore

330   volume was also calculated.

331

332   In order for the geometries of the water-filled pores to be modelled, surface mesh files (.stl)

333   were required, which were generated in VG StudioMax v2.2. After segmentation of the soil

334   water phase, a cube shaped ROI template was imported to create identical cubes for the surface

335   mesh generation. Each sample was subsampled, from random initial coordinates, with 6 cubes

336   comprising side lengths of 3.8 mm. The same coordinates were used for different matric

337   potentials of the same sample. Each individual surface mesh file took between $0.5 - 2$ hours to

338   generate depending on the complexity of the surface.

339

340   *3.4. Modelling*

341   The hydraulic conductivity of the soil was calculated using the method of homogenization

342   [*Pavliotis and Stuart*, 2008]. This method allows an average Darcy's law to be defined which

343   is applicable to the soil column and is parameterised by the pore scale geometry of the soil.

344   The key assumption used was that the hydraulic properties of the soil can be accurately

345   captured by studying a small subsample of the soil [*Fowler*, 1997], this is often referred to as

346   the Representative Elementary Volume (REV). In this paper we were particularly interested in

347   determining the minimum size REV for which the calculated hydraulic conductivity converges

348   to the macroscopic hydraulic conductivity for soil samples of the order of a few millimetres.

361 This method allowed us to theoretically determine the required sample size for CT analysis of

362 water flow in soil at this scale.

363

364 We note that the averaged hydraulic conductivity and WRC calculated at this scale are not

365 globally applicable. Rather they describe the average flow rate in soil of this type assuming no

366 large scale features such as cracks or larger voids are encountered. In theory these features

367 could be included by deriving an averaged Darcy's law for each soil type, Stokes' equations

368 for large void spaces and approximate Darcy's laws for fractures equation [*Hornung*, 1997].

369 These systems could, in principle, be upscaled to derive averaged Darcy's laws on a much

370 larger scale. However, in this paper our focus is in obtaining an estimate of the hydraulic

371 properties of soils based on their pore structure. Hence, we neglect these larger scale features

372 but emphasise that they could be included through an additional level of upscaling.

373

374 As we are able to segment the air and water separately from the CT images, the fluid dynamics

375 can be greatly simplified. Rather than focusing on the moving interface between each phase,

376 we consider the relatively slow, flow of water about a fixed interface. The resulting equations

377 are introduced in the appendix. We further simplify the equations by assuming that the non-

378 wetting phase, in this case air, is stationary. This assumption is valid assuming that the air

379 phase is disconnected. If this is not the case then the movement of the air effectively lubricates

380 the movement of water resulting in an increase in the hydraulic conductivity. This approach is

381 valid assuming that the pressure gradients are sufficiently low, such that the interface remains

382 fixed, and that the non-wetting phase is not connected, hence, the trapped non-wetting phase

383 has zero average velocity.

14

392

Strictly speaking the theory of homogenization requires the soil structure to be periodic. Clearly for real soil samples this is not the case. This is overcome by enforcing an apparent image based periodicity either by translation or reflection of the CT image (Figure 1). In this paper reflection was chosen as it simplifies the resulting calculations and does not introduce discontinuities in the soil structure. The error induced by enforcing periodicity is that the geometry considered numerically is now fully periodic rather than quazi-periodic and does not truly represent the imaged soil structure. To overcome this, different size REVs were taken from the segmented .stl files, see Figure 1. Specifically, the 6 cubes which were segmented from each scan were of the same volume $V_m = 54.9$ mm^3, and were assumed to be sufficiently large that the soil properties would have converged. The REVs sampled from the six cubes were of volume $V = V_m/(2^j)$, where j is a positive integer in the range 0 to 10 such that the smallest volume we consider is 0.053 mm^3 and the largest is V_m. As j is decreased and, hence, the size of the REV is increased, the relative size of the errors induced by the reflection decreases. Similarly as the REV size increases, the hydraulic properties of the subsample will, in principle, converge to the hydraulic properties of the soil.

408

A rigorous analysis of equations governing fluid flow through soil is given in [*Daly and Roose*, 2014]. Here we provide a brief description of the theory in the applied context with further details given in the appendix. We consider a sample of soil which is large with respect to the REV size. Specifically if the REV is a small cube with side length L_y and the large sample we are interested in has characteristic length L_x then we require that the ratio of these two lengths is small, *i.e.*, $\epsilon = L_y/L_x \ll 1$.   For typical pore sizes observed in soil, the viscous forces dominate [*Fowler*, 1997]. Hence, we may consider the Stokes limit of the Navier-Stokes

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421 equations where all inertial terms are neglected. Mathematically the macroscopic hydraulic

422 conductivity, which is valid for the whole soil column, is obtained in two steps. First, it is

423 shown, see appendix for details, that pressure variations across the REV of size ϵ will induce

424 a water velocity, also of size ϵ. Secondly, a set of equations are derived which allow the

425 pressure driven fluid velocity to be determined based on the soil (Figure 1). Finally, the average

426 velocity over the REV is used to determine a Darcy's law which is independent of ϵ. This value

427 is valid for the bulk soil and describes fluid driven by an external pressure gradient, see [*Daly*

428 *and Roose*, 2014]

$$\boldsymbol{u} = -\mathcal{K}(\boldsymbol{\nabla} p_0 - \rho g \hat{\boldsymbol{e}}_z), \tag{3}$$

429 where ρ is the fluid density ($\rho = 10^3 \text{kg m}^{-3}$ in the case of water), $g = 9.8 \text{ m s}^{-2}$ is the

430 acceleration due to gravity, p_0 is the applied pressure, $\boldsymbol{u}$ is the volume averaged water velocity

431 and $\mathcal{K}$ is the relative permeability (in the general case a tensor) which has components defined

432 as

$$\mathcal{K}_{jk} = \frac{L_y^2}{\mu} \int_{\Omega_w} \hat{\boldsymbol{e}}_j \cdot \boldsymbol{v}_k^w \, dy. \tag{4}$$

433 Here $\hat{\boldsymbol{e}}_j$, for $j = x, y, z$ is a unit vector in the j-th direction and $\boldsymbol{v}_k^w$ is the local velocity.

434 Assuming that the air velocity is slower than the water velocity then $\boldsymbol{v}_k^w$ satisfies the cell

435 problem

$$\boldsymbol{\nabla} \cdot \sigma_k^w - \boldsymbol{\nabla} \pi_k^w = \hat{\boldsymbol{e}}_k, \qquad\qquad \boldsymbol{x} \in \Omega_w, \tag{5a}$$

$$\boldsymbol{\nabla} \cdot \boldsymbol{v}_k^w = 0, \qquad\qquad \boldsymbol{x} \in \Omega_w, \tag{5b}$$

$$\boldsymbol{v}_k^w = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_s, \tag{5c}$$

$$\boldsymbol{v}_k^w = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_{aw}, \tag{5d}$$

16

440   where $\sigma_k^w = (\nabla v_k^w) + (\nabla v_k^w)^T$, π_k^w is the local pressure correction due to the microscale

441   geometry, the superscript w denotes the water phase, Ω_w, Γ_s and Γ_{aw} denote the water domain,

442   the soil boundary and the air-water interface respectively. Physically this can be thought of as

443   calculating the fluid velocity subject to a unit pressure gradient. As the equations are linear

444   Darcy's law follows by multiplying the result by the desired pressure gradient.

445

446   Equations (5) were solved numerically on each subsample obtained from the CT images using

447   OpenFOAM, an open source Computational Fluid Dynamics toolbox running on IRIDIS, the

448   High Performance Computing Facility at the University of Southampton. The hydraulic

449   conductivity is then calculated as the average water velocity due to gravity. Results were

450   obtained for each soil sample using the method illustrated in Figure 1. The 6 cube surface

451   meshes, generated from each soil core, were repeatedly sampled to obtain 3D REVs of

452   increasing size. The result is a set of hydraulic conductivity calculations which we expect to

453   converge to the true hydraulic conductivity of the soil, at each point along the WRC, as the

454   sub-volume size is increased.

455

456   *3.5. Statistical analysis*

457   The results obtained directly from the CT images were analysed by general analysis of

458   variance (ANOVA) containing soil type and matric potential and all possible interactions

459   as explanatory variables using Genstat 15.1 (VSN International, UK).

460

461   **4. Results & Discussion**

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462

4.1 Hydraulic properties

464

The clay soil contained an average stone volume of 0.2 cm^3 whereas the average stone volume for the sand soil was 5.5 cm^3 (total soil volume was 55 cm^3). Therefore although field structured soil was able to retain water for longer, as shown by the curves from the conventional method (Figure 2), the influence of stones was not a consideration and the greater water holding capacity is attributed to the more complex pore network in the field structured soil.

470

There are clear differences between sieved soil and field structured soil (Figures 2 and 3). For the field structured soil the imaging technique shows a good agreement with the conventional methods with an average error of $\leq$ 5% across the range of matric potentials considered. The sieved soil responded more strongly to pressure change in comparison to the field structured soil. However, the imaged sieved soil data did not show this trend, hence, there was quite a large error of 9% at 0 kPa and as high as 65% at -75 kPa. This suggests that, for the field structured samples, the segmentation procedure was successful at identifying the majority of the water filled pores (WFPs).

479

It is somewhat surprising that the more homogeneous sieved soil is less well represented through the imaging technique than the field structured soil. However, we attribute this to the pore size distribution in the two soils. The field structured soil has a wider range of pore sizes than the sieved soil which contains macropores defined by the largest grain size. As the matric potential becomes increasingly negative, the water drains from smaller and smaller pores. In

18

495   the case of the field structured soil this results in a gradual decrease in volumetric water content

496   with matric potential which the imaging technique can follow assuming the pore size is larger

497   than the image resolution. In the case of the sieved soil there is a reduced range of pore sizes

498   and the water quickly drains from the larger pores and remains trapped in a large number of

499   smaller pores. The result is that the volumetric water content drops quickly as the matric

500   potential is reduced and the imaging technique, which cannot resolve the smallest pores, fails

501   to capture this. We also note that the differences observed in the WRC between the two soil

502   types is smaller when measured using the imaging technique than using conventional methods.

503   This is again attributed to the finite resolution of the measurements and is analogous to the

504   increasing error observed as the saturation is decreased.

505

506

507   The hydraulic conductivity, which corresponds to the water release curves shown in Figure 2,

508   has been calculated by solving equations (5) for increasing subsample size. As the subsample

509   size was increased, the hydraulic conductivity converged to a fixed value which we interpret

510   as the macroscopic hydraulic conductivity (Figure 4). The sample size at which this occurred

511   was smaller for lower matric potentials. This is attributed to the increased air-filled porosity

512   which decreases the contact area between the water and the soil. Hence, the effect of the

513   heterogeneous soil boundary is decreased and the overall homogeneity of the sample is

514   increased. We note that, as convergence is seen to occur at smaller sample sizes for lower

515   saturation, fewer simulations were run at these saturation values. This occurs at approximately

516   ¼ of the sample volume for all matric potentials for both soils. We note that although the REV

517   volume is approximately the same for all soil types considered, the properties obtained from

518   this analysis are different for different soil types. This method therefore enabled the required

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521 sample size for CT analysis of water flow in these samples to be determined, which was a cube

522 volume of 11 mm³. This sample was large enough for convergence of the calculated hydraulic

523 conductivity, meaning that the sample size contained an adequate volume of soil to capture the

524 averaged hydraulic properties on this scale. Therefore considering the averaging method used

525 in this paper, the hydraulic conductivity is globally valid. This means that the representative

526 volume of 11 mm³ is an appropriate representative volume element for the global hydraulic

527 conductivity of the soil samples used in this study. However, there may be significant features,

528 such as cracks or heterogeneities in the soil, which are not captured in the 11 mm³ volume.

529 Hence, care must be taken in applying these results on the field scale and an additional level of

530 upscaling may be required to capture the properties of these features.

531

532 The method described in Section 3.4 and the appendix allowed us to also obtain the hydraulic

533 conductivity in unsaturated soils as well as in *x*, *y* and *z* directions. It was expected the flow

534 would be dominant in the *z* direction due to gravity, however (possibly due to the size of the

535 samples) we did not observe this. The calculated hydraulic conductivities are shown as a

536 function of matric potential in Figure 5. The calculated hydraulic conductivities and the

537 corresponding WRC have been fitted to the van Genuchten model for the WRC and the

538 unsaturated hydraulic conductivity [*van Genuchten*, 1980] using a non-linear least squares

539 method. We note that we have only done this fitting for the imaged data as we do not have

540 hydraulic conductivity measurements for the full range of matrix potentials using conventional

541 measurements. The volumetric water content θ is given by

$$\theta = (\theta_s - \theta_r)\left(\frac{1}{1+(\alpha h)^n}\right)^m + \theta_r, \qquad (6)$$

20

550 where θ_s and θ_r are the saturated and residual volumetric water content respectively, h is the

551 is the matric head, $m = 1 - 1/n$ and n and α are the van Genuchten parameters. The

552 corresponding hydraulic conductivity is given by $K = K_{sat} k_r^{vg}$. Here K_{sat} is the saturated

553 hydraulic conductivity and the relative hydraulic conductivity is given by

$$k_r^{vg} = \frac{\{1 - (\alpha h)^{n-1}[1 + (\alpha h)^n]^{-m}\}^2}{[1 + (\alpha h)^n]^{m/2}}. \qquad (7)$$

554 We take θ_r to be negligible and fit the remaining parameters to the imaged data. The average

555 saturated hydraulic conductivity and the fitted parameters: $K_{sat}, \theta_s, \alpha$ and n are shown in Table

556 1. The fitted hydraulic conductivity and WRC, obtained using these parameters, are shown in

557 Figure 5.

558 **Table 1. Fitted van Genuchten parameters**

Soil	Measured K_{sat} [$cm\ s^{-1}$]	Calculated K_{sat} [$cm\ s^{-1}$]	Saturated volumetric water content θ_s	α [cm^{-1}]	n
Sand	0.0013	0.0014	0.418	0.153	1.71
Clay	0.0004	0.0010	0.423	0.148	1.66

559

560 The hydraulic conductivity decreases as the matric potential decreases from 0 kPa to -75kPa

561 (P<0.001) and is slightly larger in the sand soil compared to the clay soil. The van Genuchten

562 model fits the imaged data well with a slight discrepancy in the WRC at 0 kPa (RMSE=0.016

563 for the clay soil and RMSE=0.022 for the clay soil). The experimentally measured hydraulic

564 conductivity agreed well with the calculated value for the sand soil. However, the data

565 compares less favourably in the case of the clay soil where the experimentally measured

566 hydraulic conductivity is less than half that of the calculated value.

567

The differences observed for the clay soil between the experimentally measured hydraulic conductivity and the calculated values are attributed to a combination of image resolution and segmentation limitations. Typically clay soils have much smaller pores than sand soils, due to the differences in particle size and as such will have a large number of pores which are comparable with or below the imaging resolution used in this study. Further complications include partial volume effects, which can cause errors in voxel classification, if the boundary is between two phases with wide ranging attenuation values, such as air and rock [*Ketcham and Carlson*, 2001]. This therefore could lead to an incorrect classification of voxels resulting in either an overestimate or an underestimate of the pore space and pore size distribution due to the effect it has on objects of differing sizes, which may explain why we have an overestimation at the wet end, and underestimation at the dry end. As the calculated hydraulic conductivity is large with respect to the measured value we expect that the pore space is being over estimated in this case.

4.2 Pore characteristics

In order to obtain a better understanding of the pore scale processes which contribute to the macroscopic hydraulic properties we consider the distribution of air and water within the soil matrix on the pore scale. In order to quantify the air and water content we define Air Filled Pores (AFPs) and Water Filled Pores (WFPs) as single connected regions of air or water respectively. We also define the pore space as the union of all the AFPs and WFPs. Throughout this paper when discussing connected regions of water and air we exclusively use this terminology in keeping with the soil science literature. In addition we refer to individual pores within the soil to define simply connected pathways between two distinct points within the pore space. Typically, the pore space contains a single large WFP which contains over 50% of

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596   the water within the pore space and a large number of much smaller AFPs and WFPs. The

597   connected WFPs are the main contributor to both the WRC and the hydraulic conductivity

598   calculations. However, insight may be gained into the wetting and drying behaviour of the

599   soils by considering the properties of the AFPs. It is possible to determine the interfacial area

600   between water and air phases [*Costanza-Robinson et al.*, 2008]. However, as previous studies

601   have made the calculations based on porous media other than soil, this is not comparable to the

602   work undertaken in this study. In addition the majority of the previous work in soil [*Falconer*

603   *et al.*, 2012; *García-Marco et al.*, 2014; *Tucker*, 2014], refer directly to WFPs and AFPs as

604   such it is more meaningful in the context of the literature to consider the WFP and AFP surface

605   area.

606

607 At 0 kPa we have a residual AFP volume of approximately 6% of the total pore space, see

608 supplementary Figure 3. This volume corresponds to trapped air pockets which, due to the high

609 tortuosity of the soil structure, remain in the pore space at saturation. As the matric potential

610 is made more negative, the total WFP volume decreases and the corresponding AFP volume

611 increases (Figure 6; P<0.001). In addition to the change in WFP and AFP volume, the total soil

612 volume is seen to increase by up to 10%. Whilst this may, in part, be attributed to segmentation

613 issues and partial volume effects, it is also likely that changes in soil structure are responsible

614 for this apparent increase. The total AFP volume and the corresponding AFP surface area did

615 not vary significantly across the different soil types (sand and clay). However, the average

616 volume and the corresponding surface area of an individual AFP averaged across all matric

617 potentials shows significant variation across soil types (P<0.001) with average AFP volume of

618 0.00115 mm^3 in the clay soil compared with 0.00163 mm^3 in the sand soil and with average

619 AFP surface area of 0.0109 mm^2 in the clay soil and 0.0182 mm^2 in the sand soil. This greater

620 surface area corresponds to a much greater number of AFPs at -75 kPa in the clay soil (104

23

pores/mm³) than the sand soil (55 pores/ mm³), (P<0.05). The average pore thickness

decreased significantly with decreasing matric potential (Figure 7; P<0.001), from an average

of 0.030 mm at 0 kPa to 0.026 mm at -75 kPa. There was no significant difference between soil

types. At 0 kPa the average pore thickness for the clay soil was 0.027 mm compared to 0.033

mm for sand. At -75 kPa the average pore thickness for the clay and the sand soils was just

0.026 mm.

To highlight the differences in the soil types we classify the AFPs as air filled macro (> 100

μm), air filled meso (30 - 100 μm) and air filled micropores (21 – 30 μm), see supplementary

Figure 4. This classification is based on the pore diameter, *i.e.* the maximum sphere diameter

which can fit inside the pore. The air filled micropore range is limited by the achievable

resolution based on our estimation of image noise (> 2 voxels). As the matric potential is

decreased, the total number of AFPs classified into each category changes with an overall

increase in the number of air filled macropores seen in the case of the sand soil. In the sand

soil there is a linear increase in the percentage of air filled macropores with soil drainage, this

was expected due to the high number of macropores throughout the sand structure, which is

typically more homogenous compared to the clay. The clay soil had a less regular drainage

pattern than the sand soil. There is a linear increase in the air filled micropore percentage as

the clay dried. However, the percentage of air filled macropores is quite irregular and can be

attributed to the occurrence of crack formation which is more likely in clay soil. The average

pore circularity of the AFPs, equation (2), also increased as matric potential decreased. At -75

kPa the AFP space was comprised of pores with an average circularity of 0.78 compared to a

circularity value of 0.72 at 0 kPa (P<0.001). Our results also show that at drier matric potentials

the shape of the air and water filled pores become increasingly rounded possibly linked to air

bubble-style pore formation. This measurement enables us to quantify the influence of pore

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shape on soil drainage. However we note this is only a single measure and further descriptors that account for the shape of the soil porous architecture in 3D are required to advance understanding in this area. The 3D connectivity of the WFPs decreased with decreasing matric potential (P<0.05), pore connectivity was significantly greater in the clay soil at all matric potentials (P<0.01).   The differences in the AFP structure in the two soil types are indicative of the soil structure. The clay soil contains a large number of micro-pores which corresponds to the large number of air filled micropores trapped within the soil structure. In contrast the sand soil contains a larger number of meso- and macro-pores causing the air to be contained in relatively few, larger volumes. It is highly likely that the clay soil contained many pores that were too small to be imaged at the achievable resolution in this study.

5. Conclusions

In this paper we have used a combination of CT imaging and mathematical modelling to derive the WRC and unsaturated hydraulic conductivity in contrasting soil textures (*i.e.* clay and sand) and structures (i.*e.* sieved and field structured soil). The ability to non-destructively scan the soil cores and segment the water and air filled pores from the soil matrix combined with the mathematical modelling have provided a unique insight into the WRC and hydraulic properties of contrasting soils. A comprehensive study of the hydraulic properties of different soil types based on a quantitative visual assessment of the soil structure has been carried out across a range of different saturation values. The results predict the increase in hydraulic conductivity with saturation and highlight the structural basis for the differences between sand and clay soils.

699   The data obtained via CT imaging allows us to visualise the soil structure in 3D and, hence, to

700   measure the effects of pore geometry on the WRC. Using the scanned geometry combined with

701   detailed mathematical modelling, the unsaturated hydraulic conductivity and, hence, hydraulic

702   conductivities were obtained at the precise measured matric potentials. These values are not

703   easily measured in the laboratory through experimentation alone. Using the measured saturated

704   hydraulic conductivity and typical soil parameters we have compared the computed values

705   against the van Genuchten model. The simulated values compared well with the measured

706   saturated hydraulic conductivity values. The calculated hydraulic properties of the clay soil

707   against the van Genuchten model compare well. However, in the case of the sand soil, the

708   calculated values are lower than those obtained using the van Genuchten model.   This

709   highlights the potential structural differences between two soils of a particular textural class

710   and the potential role played by the micro-porous structure of the aggregates which are

711   currently below the imaging resolution of CT for this sample size.

712

713   The combination of CT imaging and modelling provides a new level of understanding of the

714   role that structure plays in the averaged properties of soil. The WRC has been obtained via CT

715   imaging and provides a detailed geometric representation of how the water is retained in pores

716   of different sizes. A downside to the imaging approach is the time investment required to

717   collect the data. However, this is comparable to the conventional laboratory method as the

718   majority of time is spent allowing the sample to equilibrate. However, this detailed information

719   on the pore-scale architecture represents a significant development in our ability to study soil

720   systems which will continue to advance pore scale modelling by providing three dimensional

721   geometries against which the numerical models can be tested [19].

722

26

737 **Table A1. Values of parameters used in numerical calculations**

Constant	Value	Units	Description
ρ^w	10^3	kg m^{-3}	Density of water
ρ^a	1.2	kg m^{-3}	Density of air
μ^w	10^{-3}	kg m^{-1} s^{-}	Viscosity of water
μ^a	20×10^{-6}	kg m^{-1} s^{-}	Viscosity of air
$\tilde{g}$	9.8	m s^{-2}	Acceleration due to gravity
r_y	100×10^{-6}	m	Typical pore radius
L_y	10^{-3}	m	Typical size of microscale geometry
L_x	10^{-1}	m	Typical size of macroscale geometry
$\tilde{\gamma}$	7.28×10^{-2}	kg s^{-2}	Surface tension
$\tilde{\kappa}$	10^4	m^{-1}	Curvature
ϵ	10^{-2}		Ratio of geometry sizes
κ	2.97×10^2		Scaled surface tension-curvature product
$\lvert \boldsymbol{g} \rvert$	1		Scaled gravitational force
δ_u	0.16		Typical velocity ratio
δ_p	8.33		Typical pressure ratio

738

739 *Appendix*

740 The starting point is the non-dimensional Stokes equations using the scaling given in [*Daly and*
741 *Roose*, 2014]. The dimensional variables, denoted with a $\tilde{\ }$, are scaled using

$$\boldsymbol{v}^w = \frac{4\mu^w}{\rho^w \tilde{g} r_y^2}\,\widetilde{\boldsymbol{v}}^w, \qquad\qquad p^w = \left(\frac{L_x}{L_y}\right)\frac{4L_y}{\rho^w \tilde{g} r_y^2}\,\tilde{p}^w,$$

$$\boldsymbol{v}^a = \left(\frac{L_x}{L_y}\right)\frac{4\mu^a}{\rho^a \tilde{g} r_y^2}\,\widetilde{\boldsymbol{v}}^a, \qquad\qquad p^a = \left(\frac{L_x}{L_y}\right)^2 \frac{4L_y}{\rho^a \tilde{g} r_y^2}\,\tilde{p}^a,$$

$$t = \left(\frac{L_x}{L_y}\right)\frac{\rho^w r_y^2 \tilde{g}}{4\mu^w L_y}\,\tilde{t},$$

742 where $\boldsymbol{v}^w$ and $\boldsymbol{v}^a$ are the water and air velocities, p^w and p^a are the water and air pressures, t
743 is the dimensionless time and the remaining parameters are listed in table A1 and the non-
744 dimensional constants are

$$\kappa = \left(\frac{L_y}{L_x}\right)\frac{4L_y\,\tilde{\gamma}}{\rho^w r_y^2\,\tilde{g}}\,\tilde{\kappa}, \qquad \boldsymbol{g} = \frac{\hat{\boldsymbol{e}}_x L_y^3}{r_y^2 L_x}, \qquad \delta_u = \left(\frac{L_y}{L_x}\right)\frac{\rho^w \mu^a}{\rho^a \mu^w}, \qquad \delta_p = \left(\frac{L_y}{L_x}\right)\frac{\rho^w}{\rho^a},$$

745 Typical values for the non-dimensional numbers are given in table A1. The resulting non-
746 dimensional Stokes equations are:

$$\epsilon \boldsymbol{\nabla} \cdot \sigma^w - \boldsymbol{\nabla} p^w = \epsilon \boldsymbol{g} \qquad\qquad\qquad \boldsymbol{x} \in \Omega_w, \quad \text{(A1a)}$$

$$\epsilon \boldsymbol{\nabla} \cdot \sigma^a - \boldsymbol{\nabla} p^a = \epsilon^2 \boldsymbol{g} \qquad\qquad\qquad \boldsymbol{x} \in \Omega_a, \quad \text{(A1b)}$$

$$\epsilon \phi \frac{\partial s^w}{\partial t} + \boldsymbol{\nabla} \cdot \boldsymbol{v}^w = 0, \qquad\qquad\qquad \boldsymbol{x} \in \Omega_w, \quad \text{(A1c)}$$

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$$\epsilon\phi\frac{\partial S^a}{\partial t} + \delta_u \boldsymbol{\nabla} \cdot \boldsymbol{v}^a = 0, \qquad\qquad \boldsymbol{x} \in \Omega_a, \quad \text{(A1d)}$$

$$\boldsymbol{v}^w = 0, \qquad \boldsymbol{v}^a = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_s, \quad \text{(A1e)}$$

$$\hat{\boldsymbol{n}} \cdot \boldsymbol{v}^w = 0, \qquad \hat{\boldsymbol{n}} \cdot \boldsymbol{v}^a = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_{aw}, \quad \text{(A1f)}$$

$$\hat{\boldsymbol{\tau}} \cdot [\boldsymbol{v}^w - \delta_u \boldsymbol{v}^a] = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_{aw}, \quad \text{(A1g)}$$

$$\hat{\boldsymbol{n}} \cdot \left[(\epsilon\sigma^w - Ip^w) - \delta_p(\epsilon\sigma^a - Ip^a)\right] = \hat{\boldsymbol{n}}\kappa, \qquad\qquad \boldsymbol{x} \in \Gamma_{aw}, \quad \text{(A1h)}$$

747   combined with periodic boundary conditions in the x y and z directions. To simplify the
748   problem we notice that $\delta_u^{-1} < 1$ and $\delta_p < 1$ and, therefore, they can be neglected to first
749   approximation. The result is that the air and water velocities decouple we need only consider
750   the flow of water through the soil. The air velocity could in principle be calculated as a
751   perturbation of order δ_u which in turn would create a water velocity perturbation of order
752   δ_u/δ_p. However, we expect that this perturbation will be small with respect to the uncertainty
753   in the segmentation and the temperature dependent variations of the water viscosity. The
754   simplified equations are:

$$\epsilon\nabla^2\boldsymbol{v}^w - \boldsymbol{\nabla}p^w = \epsilon\boldsymbol{g} \qquad\qquad \boldsymbol{x} \in \Omega_w, \quad \text{(A2a)}$$

$$\boldsymbol{\nabla} \cdot \boldsymbol{v}^w = 0, \qquad\qquad \boldsymbol{x} \in \Omega_w, \quad \text{(A2e)}$$

$$\boldsymbol{v}^w = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_s, \quad \text{(A2f)}$$

$$\boldsymbol{v}^w = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_{aw}, \quad \text{(A2g)}$$

755   We expand the gradient operator $\boldsymbol{\nabla} = \epsilon\boldsymbol{\nabla}_x + \boldsymbol{\nabla}_y$ and look for a power series solution to
756   equations (A4) of the form

$$\boldsymbol{v}^w = \boldsymbol{v}_0^w + \epsilon\boldsymbol{v}_1^w + O(\epsilon^2) \qquad\qquad \text{(A3a)}$$
$$p^w = p_0^w + \epsilon p_1^w + O(\epsilon^2) \qquad\qquad \text{(A3b)}$$

757   Substituting equations (A3) into equations (A2) and retaining terms $O(\epsilon^0)$ we obtain $\boldsymbol{\nabla}_y p_0^w = $
758   0, *i.e.*, the largest component of the pressure is constant over the microscale. Expanding to
759   $O(\epsilon^1)$ we effectively convert the leading order pressure drop into a body force and obtain

760
$$\boldsymbol{v}_0^w = \sum_{k=1}^{3} \boldsymbol{v}_k \partial_{x_k} p_0^w$$

761   Where $\boldsymbol{v}_k$ satisfy the cell problem

$$\nabla_y^2 \boldsymbol{v}_k - \boldsymbol{\nabla}_y \pi_k^w = \hat{\boldsymbol{e}}_k, \qquad\qquad \boldsymbol{x} \in \Omega_w, \quad \text{(A4a)}$$

$$\boldsymbol{\nabla}_y \cdot \boldsymbol{v}_k = 0, \qquad\qquad \boldsymbol{x} \in \Omega_w, \quad \text{(A4b)}$$

$$\boldsymbol{v}_k = 0, \qquad\qquad \boldsymbol{x} \in \Gamma_s, \quad \text{(A4c)}$$

$$\boldsymbol{v}_k = 0, \qquad\qquad\qquad\qquad x \in \Gamma_{aw}, \quad \text{(A4d)}$$

$$\boldsymbol{v}_k, \pi_k^w \text{ periodic with period 1} \qquad\qquad\qquad\qquad \text{(A4e)}$$

762 Averaging the velocity over the unit cell we obtain the non-dimensional form of Darcy's law
763 for the average velocity

$$\boldsymbol{u}_0^w = \int_\Omega \boldsymbol{v}_k \otimes \hat{\boldsymbol{e}}_k \, dy \, (\nabla_x p_0^w + \boldsymbol{g})$$

765 which in dimensional form is equation (3) and (4) in the main paper.

766 As discussed in the main text we impose periodicity via reflection of the geometry about the x,
767 y and z axis. The result is that the velocity is, for the k-th cell problem, the pressure correction
768 π_k^w is even odd in the direction x_k and even in the direction x_j for $j \neq k$. Similarly we find
769 that the k-th component of the velocity, $\hat{\boldsymbol{e}}_k \cdot \boldsymbol{v}_k$ is even in all directions. The remaining velocity
770 components, $\hat{\boldsymbol{e}}_j \cdot \boldsymbol{v}_k$ for $j \neq k$, are odd in the directions x_k and x_j, however, they are even in
771 the direction x_p for $p \neq k$ and $p \neq j$. Specifically the resulting boundary equations and
772 boundary conditions are:

$$\nabla_y^2 \boldsymbol{v}_k - \boldsymbol{\nabla}_y \pi_k^w = \hat{\boldsymbol{e}}_k, \qquad\qquad\qquad x \in \Omega_w, \quad \text{(A5a)}$$

$$\boldsymbol{\nabla}_y \cdot \boldsymbol{v}_k = 0, \qquad\qquad\qquad x \in \Omega_w, \quad \text{(A5b)}$$

$$\boldsymbol{v}_k = 0, \qquad\qquad\qquad x \in \Gamma_s, \quad \text{(A5c)}$$

$$\boldsymbol{v}_k = 0, \qquad\qquad\qquad x \in \Gamma_{aw}, \quad \text{(A5d)}$$

$$\pi_k^w = 0, \ \frac{\partial}{\partial x_k}(\hat{\boldsymbol{e}}_k \cdot \boldsymbol{v}_k) = 0, \ \hat{\boldsymbol{e}}_j \cdot \boldsymbol{v}_k = 0, j \neq k \qquad x \in \partial x_k \quad \text{(A5e)}$$

$$x \in \partial x_j \quad \text{(A5f)}$$

$$\frac{\partial}{\partial x_j}\pi_k^w = 0, \ \frac{\partial}{\partial x_p}(\hat{\boldsymbol{e}}_p \cdot \boldsymbol{v}_k) = 0, \ \hat{\boldsymbol{e}}_j \cdot \boldsymbol{v}_k = 0, p \neq k, , p \neq j$$

773 where ∂x_k is the boundary located at $x_k = 0$ and $x_k = 1/2$, ∂x_j is the union of the boundaries
774 located at $x_j = 0$ and $x_j = 1/2$ for $j \neq k$.

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778 **7. References**
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897 **8. Figures**

898 **Figure 1.** Outline of steps used in the subsampling, meshing and solution of cell

899 problem. Figures (a) and (b) show the subsampling of a 1.2 mm side length cube of the

900 segmented .stl file. Figure (c) shows the mesh generation. Figures (d) and (e) show the

901 generation of a truly periodic geometry through translation (d) and reflection (e), the original

902 volume mesh is shown in a lighter shade. Figure (f) shows the numerical solution for the local

903 velocity magnitude in the subsampled soil.

904

905 **Figure 2.** Water release characteristic of the sand and clay sieved (a) and field structured (b)

906 soils fitted to the van Genuchten (VG) curve. Error bars associated with histograms show one

907 standard deviation. Note the different scaling on the y axis for figures (a) and (b).

908

909 **Figure 3.** Greyscale images showing examples of field structured (a) & (b) and sieved (c) &

910 (d) clay (a) & (c) sand soil (b) & (d) at -75 kPa. Scale bar = 0.25 mm.

911

912 **Figure 4.** Hydraulic conductivity of field structured sand and clay soils averaged over all

913 directions for decreasing j corresponding to increasing volume. The dots show the result from

914 each of the cubic volumes, the lines show the average. The corresponding saturated hydraulic

915 conductivity measured for the sand and clay soils are 1.3×10^{-3} cm s^{-1} and $0.4 \times$

916 10^{-3} cm s^{-1} respectively.

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919 **Figure 5.** Hydraulic conductivity of <u>field structured </u>sand and clay soils as a function of

920 matric potential.

921 **Figure 6.** 3D core section of a sand (a) and clay (b) soil sample at the specific matric potentials

922 from saturation (left) to drier (right). Segmented phases are coloured brown (soil), blue (water

923 filled pores) and black (air filled pores). Scale bar = 5 mm.

924

925 **Figure 7.** 3D pore thickness heat maps for a representative sand (a-c) and clay (e-f) sample at

926 0 (a, d), -20 (b, e) and -75 (c, f) kPa. (g) total air filled pore volume (bars) and average individual

927 pore volumes (lines) for sand and clay soil. Error bars associated with histograms show one

928 standard deviation.

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942 **Figure 1.**

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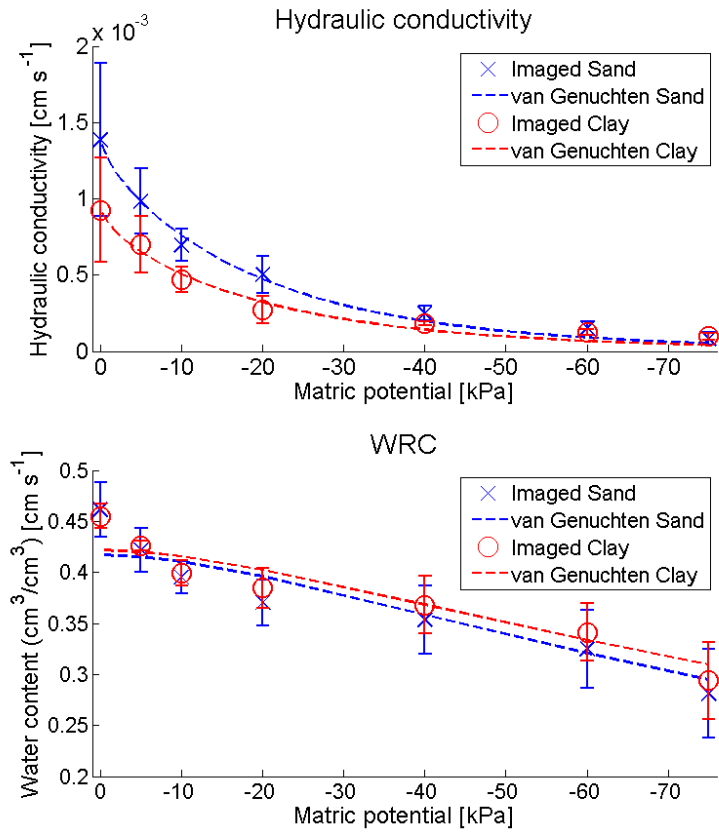
971   **Figure 4.**

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