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Abstract

We report current work to study transfer of population and orientation in collisions of NaK molecules with argon and potassium atoms using polarization labeling (PL) and laser-induced fluorescence (LIF) spectroscopy. In the PL experiment, a circularly polarized pump laser excites a specific NaK A¹ Σ^+ ($v'=16, J') \leftarrow$ $X^{1}\Sigma^{+}(v''=0, J' \pm 1)$ transition, creating an orientation (non-uniform $m_{I'}$ level distribution) in both levels. The linearly polarized probe laser is scanned over various $3^{1}\Pi(v, J'\pm 1) \leftarrow A^{1}\Sigma^{+}(v''=16, J')$ transitions. The probe laser passes through a crossed linear polarizer before detection, and signal is recorded if the probe laser polarization has been modified by the vapor (which occurs when it comes into resonance with an oriented level). Using both spectroscopic methods, analysis of weak collisional satellite lines adjacent to these directly populated lines, as a function of argon buffer gas pressure and cell temperature, allows us to discern separately the effects collisions with argon atoms and potassium atoms have on the population and orientation of the molecule. In addition, code has been written which provides a theoretical analysis of the process, through a solution of the density matrix equations of motion for the system.

Background



- We use Optical-Optical Double Resonance (OODR) to study the NaK molecule in both Polarization Labeling and Laser-Induced Fluorescence spectroscopy
- Quantitative study of collisional population and orientation transfer Separately study the effects of
- collisions with Argon vs. Potassium

Polarization Spectroscopy

Due to asymmetric m_{J} distribution created by circular polarized pump beam $(\Delta m = +1$ for right circular polarization), vapor is birefringent for any probe that shares upper, lower, or both levels of pump transition

Note: linearly polarized probe can be thought of as equal components of left and right circular polarization. ($\Delta m = -1$ and +1)



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-.J+1 *J*-1 Polarization signal depends on the difference of absorption between left and right circular polarization components of the linearly polarized probe beam

Transfer of Orientation (Polarization Spectroscopy)

 F_{-J}^{L} F_{-J+1}^{L} F_{-J}^{R} F_{-J+1}^{R} F_{J-1}^{L} F_{J}^{L} F_{J-1}^{R} F_{J}^{R}

Can model this using: Signal $\propto \sum$ Right Absorptions – \sum Left Absorptions

$$= \sum_{m=-J'}^{+J'} n_m F_m^R - \sum_{m=-J'}^{+J'} n_m F_m^L$$
$$= \sum_{m=-J'}^{+J'} n_m \left(F_m^R - F_m^L \right) = f(J, J') \sum_{m=-J'}^{+J'} mn_m = f(J, J') \langle m \rangle \sum_{m=-J'}^{+J'} n_m$$

In this expression, *F* factors are relative absorption coefficients:

$$\left|\left<\beta_{1},J_{1},m_{1}\left|\vec{r}\right|\beta_{2},J_{2},m_{2}\right>\right|^{2}=F_{J_{1},J_{2}}^{m_{1},m_{2}}\left|\left<\beta_{1}\left|\left|r\right|\right|\beta_{2}\right>\right|^{2}$$

Evaluation of difference in *F* factors yields

$$F_m^R - F_m^L = f(J, J')m$$

Thus the ratio of polarization spectroscopy signal strengths for the collisional and direct line are roughly equal to the ratio of orientation multiplied by the population of the levels. We can use a rate equation model to solve for the ratio of the intensities of the collisional and direct line: $\Delta J \Omega$ $1 \Delta J \Omega$

$$R_p = \frac{\langle m_2 \rangle n_2}{\langle m_1 \rangle n_1} = \frac{\stackrel{k_{\mathrm{Ar}}^{-3,0}}{\Gamma} n_{\mathrm{Ar}} + \stackrel{k_{\mathrm{K}}^{-3,0}}{\Gamma} n_{\mathrm{Ar}}}{1 + \stackrel{g_{\mathrm{Ar}}}{\Gamma} n_{\mathrm{Ar}} + \frac{g_{\mathrm{K}}}{\Gamma} n_{\mathrm{K}}}$$

Here $k^{\Delta J,O}$ is the rate of transfer of orientation from level 1 to the level characterized by ΔJ , and g is the total collisional rate of destruction of orientation. Separating out the population and orientation dependence yields 1 Q

 $f_{Ar,K}$ represents the fraction of orientation destroyed in the initial collision with either Ar or K, and g' is the rate of destruction of orientation by subsequent collisions that do not change J. g' is small compared to k^Q , so we approximate $g'_{Ar} = \frac{1}{2} \left(f_{Ar}^{-1} k_{Ar}^{-1} + f_{Ar}^{+1} k_{Ar}^{+1} \right)$

 $f = \frac{1}{2} \left(f_{\rm K}^{-1} k_{\rm K}^{-1} + f_{\rm K}^{+1} k_{\rm K}^{+1} \right)$

A fit of intensity ratios obtained from fluorescence and polarization spectroscopy, using our empirical formulas R_f and R_p , as functions of Ar or K number density has been carried out, for $\Delta J = \pm 1, \pm 2, \pm 3, \pm 4$

Fitting Results

The parameters k_{Ar} , k_{K} , f_{Ar} , and f_{K} give us information on the population transferred, as well as the orientation destroyed, in the initial collision from the directly populated intermediate level to the collisional level.



 $k_{\rm Ar}$ values are much larger for even numbered collisional transitions than for odd numbered transitions

 $k_{\rm K}$ values do not seem to show such a propensity f values show that collisions with potassium atoms are

more likely to destroy angular momentum orientation than collisions with argon atoms

 $k_{\rm Ar}^{\rm Q} = 1.16 \pm .07 \times 10^{-9} \ {\rm cm}^3 \ {\rm s}^{-1}$ $k_{\rm K}^{\rm Q} = 5.3 \pm .6 \times 10^{-9} \ {\rm cm}^3 \ {\rm s}^{-1}$





Density Matrix Model

To completely describe a system of quantum mechanical energy levels in a statistical manner, it is necessary to solve the density matrix equations of motion

$$\dot{\rho}_{nm} = -\frac{i}{\hbar} \sum_{k} \left(H_{nk} \rho_{km} - \rho_{nk} H_{km} \right) + relaxation \ terms$$

A diagonal element ρ_{nn} represents the population of sublevel n, while an off-diagonal element ρ_{nm} represents the laser-induced coherence between sublevels n and m. A computer code has been written which solves the density matrix equations of motion for either three coupled angular momentum states for the direct probe model, or four states for the collisional model. +1 +2



Model Includes

- Laser coherence terms
- Fluorescence into and out of various levels
- Transit relaxation
- Additional "dump" levels representing othe
- ground/intermediate levels Velocity group and laser detuning loops



Collisional model is equipped to incorporate any theoretical model for collisional transfer of population and orientation linking states "2" and "c"



- incorporate into density matrix
- Investigate ground state collisional lines
- Thus far we have only investigated excited state collisional lines
- *v*-changing collisions have been observed in ground state. We plan to investigate these in greater detail
- At Temple U., we have observed ground state collisional lines for very large ΔJ (up to $\Delta J = 58$) in polarization spectroscopy of Rb₂
- How can orientation be preserved in collisions where *J* changes by a very large amount? Effects on orientation/population by other collsion partners (He, Xe, Kr) Propensity toward $\Delta J = \pm 2, \pm 4$ not observed in NaCs
- Is this due to NaCs being "more heteronuclear" than NaK?