Non-perturbative renormalization of the static quark theory in a large volume

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We report on progress to renormalize non-pertubatively the static heavy quark theory on the lattice. In particular, we present first results for position-space renormalization scheme for heavy-light bilinears. We test our approach on RBC's $16^3 \times 32$ lattice ensemble with $m_\pi \approx 420$ MeV, Iwasaki gauge action and domain wall light fermions.

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1. Introduction

Studies of properties of heavy quarks and heavy mesons are an important part of lattice simulations of Quantum Chromodynamics since it can be argued that they are especially sensitive to the physics beyond the Standard Model. Such calculations are however difficult to perform in a straigthforward manner because of a large hierarchy of scales, namely $m_B/m_\pi \approx 40$. Consequently it is useful to employ an effective description of heavy quarks, which in the case of the present contribution, is formalized as the Heavy Quark Effective Theory. More precisely we consider the leading order of HQET and we propose a new non-perturbative renormalization scheme, which allows for a fully non-perturbative renormalization of the action and of any composite operators of interest. Our scheme is based on correlation functions in coordinate space and we discuss its practical implementation for the static Lagrangian as well as for the temporal component of the heavy-light axial and vector currents. In the following we introduce very briefly the static quark theory (section 2) and present details of our scheme (section 3). Present numerical results for the ratio of renormalization constants $Z_V^{\text{stat}}/Z_A^{\text{stat}}$, heavy quark mass and the renormalization factors themselves are discussed in sections 4.2-4.5. We conclude in Section 5.

2. HQET: an effective field theory for QCD

HQET provides an effective description of QCD processes with initial and final states containing a single heavy quark. The high momentum components of the massive quark field are integrated out and their contribution is summarized in the HQET parameters, whereas the low momentum components are present as a new two-component effective field ψ_h. The quark mass dependence is explicitly removed and hence all masses computed within HQET must be shifted by an overall energy offset. At leading order, in the so-called 'static quark theory', one assumes that the heavy quark is infinitely heavy and the Lagrangian reduces to

$$\mathcal{L}_{\text{stat}} = \bar{\psi}_h D_0 \psi_h. \tag{2.1}$$

A peculiarity of this theory is that D_0 has the same quantum numbers as the lower dimensional mass operator and thus those two operators mix under renormalization with a linearly divergent coefficient. Note that Wilson fermions suffer from the same linearly divergent additive mass renormalization. In the latter case the non-perturbative renormalization is usually done in a hadronic scheme, where the mass of an appropriate meson is tuned to its experimental value.

It was argued that renormalization of the HQET Lagrangian must be done non-perturbatively otherwise uncancelled divergent terms can combine with lattice artefacts giving finite, non-vanishing contributions [1]. A practical prescription of such renormalization for HQET at any order in $1/m$ was worked out by the ALPHA collaboration. It consists in renormalizing HQET and matching it to QCD in a small volume using SF boundary conditions [2]. Then the HQET parameters are evolved non-perturbatively using step scaling techniques to a large volume where hadronic matrix elements can be evaluated. At the leading order of HQET one is allowed to use perturbation theory [3, 4], however at the cost of large systematic errors. In this work we discuss a non-perturbative renormalization scheme which avoids using correlation functions with SF boundary conditions.

3. Description of the method

Our primary object of interest is the correlator of two heavy-light currents

$$C_\Gamma(t) = \langle \bar{\psi}_h(t,\vec{0}) \Gamma \psi(t,\vec{0}) \bar{\psi}(0,\vec{0}) \Gamma \psi_h(0,\vec{0}) \rangle. \tag{3.1}$$

It depends only on the time separation, since the heavy quark does not propagate in space. The correlator Eq.(3.1) contains all the information we need. The rate of decrease of $C_\Gamma(t)$ encodes the energy shift, and hence the large distance behavior of $C_\Gamma(t)$, where the slope corresponds to the static energy of the heavy-light meson, can be used to define a renormalization condition for the static quark mass. The slope denoted by $\delta m(t^*)$ can be estimated in the following way [5]

$$\delta m(t^*) = -\log\left(\frac{C_\Gamma(t^*+1)}{C_\Gamma(t^*)}\right). \tag{3.2}$$

Hence, on each configuration we extract the heavy-light meson mass and define the mass-renormalized correlator $\tilde{C}_\Gamma(t)$

$$\tilde{C}_\Gamma(t,t^*) = e^{-\delta_m(t^*)(t+1)} C_\Gamma(t). \tag{3.3}$$

An additional term $e^{-\delta_m(t^*)}$ was explicitly added in Eq.(3.3). Such factor is commonly included in studies using off-shell renormalization conditions and reabsorbed into the heavy quark wave function renormalization (see for example [4, 3]). We follow this convention. Note, however, that our scheme is an on-shell scheme and is defined using gauge-invariant quantities and hence the wave function renormalization factor does not appear explicitly. In other words our set of renormalization conditions is consistent with or without that factor. Note also that a renormalization condition Eq.(3.3) imposed at large t^* would be inaccessible in perturbation theory and hence would prohibit the translation to the usually employed continuum $\overline{\text{MS}}$ scheme. If one works in this non-perturbative regime of t^* a solution would be to perform a non-perturbative step scaling to a higher scale, where contact to perturbation theory can be made. Otherwise, if the lattice spacing is fine enough one can choose t^* small so that perturbation theory is reliable at that scale.

Subsequently, we use the correlator $\tilde{C}_\Gamma(t,t^*)$ to impose the renormalization condition which fixes the renormalization constant of the heavy-light current

$$\left(Z_\Gamma^X(t_0)\right)^2 \tilde{C}_\Gamma(t_0,t^*) = C_\Gamma^{\text{lattice tree-level}}(t_0). \tag{3.4}$$

We stress that one can considerably suppress discretization effects by using the lattice free correlator on the right hand side of this equation. We now describe how the above conditions work in practice.

4. Results

4.1 Ensemble details

In this feasibility study we used one ensemble generated by the RBC collaboration with three flavors of domain-wall light fermions on Iwasaki gauge action [6]. The lattice was $16^3 \times 32$ with a lattice spacing $a = 0.11\text{fm} = 1.73\text{GeV}^{-1}$. The pion mass was estimated on this ensemble to be $m_\pi \approx 420$ MeV. The following results were obtained using 20 configurations separated by 200

MDU so that autocorrelations in the HMC time are negligible. The analysis was done with two sets of propagators: point-to-all propagators and stochastic random Z_2 wall propagators. The correlation functions were averaged over the spatial volume at the sink. We used three definitions of the lattice static theory: the original Eichten-Hill proposal [7], which will be denoted in the following by EH, and two smeared versions proposed in [8] denoted by HYP1 and HYP2. We observed (see figure 1) that in general it is advantageous to use stochastic wall sources instead of point sources because the statistical error is significantly smaller. Also, the smeared actions behaved better at larger distances than the Eichten-Hill action. Besides the unitary quark mass $am = 0.01$ we used two other values: one heavier and one lighter by a factor 2.

4.2 Mass renormalization

We use mass renormalization condition as explained in Eq.(3.3) for $t^* = 8$ where the statistical noise is under control. Note that according to [4] the effective mass plateau starts at $t = 10$. As will be shown in section 4.5 the lattice spacing of our ensemble is too coarse to reliably make contact with perturbation theory and therefore in this feasibility study we take the perspective that a step scaling calculation will be needed to translate our results to the $\overline{\text{MS}}$ scheme. Hence, at this stage of our investigation we decided to renormalize the static quark mass with t^* in this non-perturbative regime. We checked that the estimated mean values of $\delta m(t^*)$ agree with the results published in Ref.[4] for the HYP1 and HYP2 static actions. The mass renormalization condition was imposed on the $C_{\gamma_5\gamma_0}(t)$ correlator at the unitary mass $am = 0.01$ configuration per configuration and used for all remaining correlators.

4.3　Z_A^{stat} and Z_V^{stat} in position space scheme

The currents' renormalization constants Z_A^{stat} and Z_V^{stat} were obtained by imposing the condition (3.4) on the correlation function $\tilde{C}_\Gamma(t)$. The values computed at $am = 0.005, 0.01$ and $am = 0.02$ were extrapolated linearly to the chiral limit. The results are shown on the right panel of figure 1.
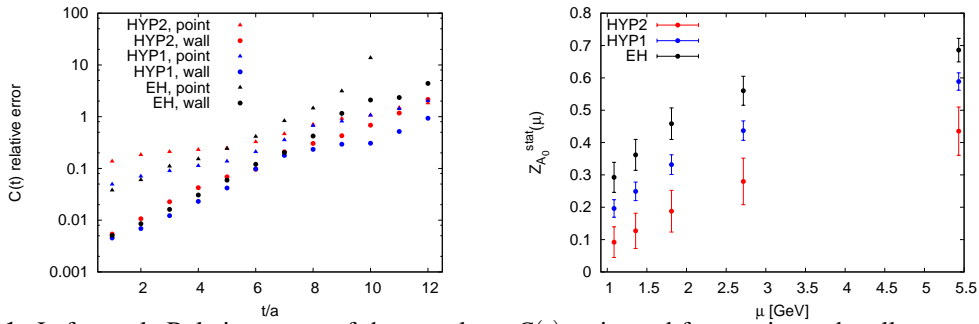

Figure 1: Left panel: Relative errors of the correlator $C(t)$ estimated from point and wall sources. Right panel: $Z_{A_0}^{\text{stat}}$ in the chiral limit. Different data sets correspond to different static actions. The scale μ corresponds to $\mu = \pi/t_0$.

4.4 $Z_A^{\text{stat}}/Z_V^{\text{stat}}$ as a check of precision

In order to test the entire setup we computed the ratio of the renormalization constants

$$R(t_0/a, am) = Z_A^{\text{stat}}(t_0/a, am)/Z_V^{\text{stat}}(t_0/a, am), \qquad (4.1)$$

which is independent of the mass renormalization. $R(t_0/a, am)$ is expected to be equal to unity in the high energy regime where effects of chiral symmetry breaking become negligible and therefore provides a check for our scheme. By using a chirally symmetric discretization of the Dirac operator one can bring this regime down so that it can be approached on current lattices. In order to check the above statement we performed a combined, correlated fit to all our data using the following fit ansatz

$$R(t_0/a, am) = R + \alpha/(t_0/a)^2 + \beta(t_0/a)^2 + \gamma am, \qquad (4.2)$$

where the term proportional to β is supposed to describe non-perturbative low-energy effects and the term proportional to α parametrizes the discretization errors. The left panel of figure 2 shows
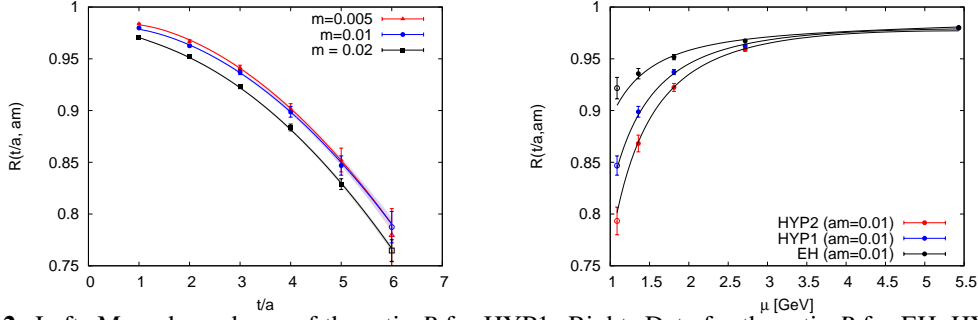

Figure 2: Left: Mass dependence of the ratio R for HYP1. Right: Data for the ratio R for EH, HYP1 and HYP2 action for am=0.01.

the dependence of R on the quark mass. We see that R approaches unity as am, the explicit chiral symmetry breaking term, is decreased. On the right panel we show the comparison of R for different static actions. For the improved actions, HYP1 and HYP2, we find the extrapolated value to be equal to unity with a good precision. The calculation with the original Eichten-Hill proposal is one standard deviation away, but it is known to suffer from large discretization effects. Results for the fit parameters are collected in table 1. The first error is statistical and the second is an estimate of the systematic error obtained by the maximal difference between fit results when changing the fitting range, i.e. including or excluding points with $t/a = 5, 6$. We conclude that our setup is working as expected and that for this ratio cutoff effects are under control even at small distances of few lattice spacings.

4.5 Z_A^{stat} and Z_V^{stat} in $\overline{\text{MS}}$

As another check of our results we translated our renormalization factors to the $\overline{\text{MS}}$ scheme and evolved perturbatively to a common scale of 3 GeV. The continuum perturbative approximation for C_Γ can be found in [9]. We use a one-loop conversion factor between continuum HQET position space scheme and continuum HQET $\overline{\text{MS}}$ scheme and two-loop running. The summary of our final

	R	α	β	γ
E-H	0.9875(9)(97)	0.0041(7)(422)	-0.0030(2)(5)	-0.79(4)(40)
HYP1	0.9995(14)(96)	-0.0033(12)(50)	-0.0057(2)(3)	-1.16(9)(37)
HYP2	1.0013(27)(38)	-0.0075(32)(40)	-0.0076(3)(2)	-0.92(12)(10)

Table 1: Results for the fit parameters from Eq. (4.2). The first error is statistical and the second is systematic.

		HYP1		HYP2	
scale [GeV]	α_s	$Z_{A_0}^{\mathrm{stat},\overline{\mathrm{MS}}}(3\mathrm{GeV})$	$Z_{A_0}^{\mathrm{stat},\overline{\mathrm{MS}},\mathrm{RGI}}$	$Z_{A_0}^{\mathrm{stat},\overline{\mathrm{MS}}}(3\mathrm{GeV})$	$Z_{A_0}^{\mathrm{stat},\overline{\mathrm{MS}},\mathrm{RGI}}$
5.44	0.1967	0.702(32)	0.523(24)	0.818(46)	0.609(32)
2.72	0.2563	0.589(40)	0.439(30)	0.755(60)	0.562(45)
1.81	0.3166	0.498(45)	0.371(34)	0.688(73)	0.513(54)
1.36	0.3864	0.419(48)	0.312(36)	0.607(80)	0.452(60)

Table 2: Numerical values of renormalization constants in $\overline{\mathrm{MS}}$. Data is ploted on figure 3.

numbers is given in table 2. In the third column we list the values of the renormalization constant in the $\overline{\mathrm{MS}}$ scheme at 3 GeV, whereas in the last column we translate them to the RGI scheme by cancelling the running with a two-loop perturbative approximation. Both, the results at 3GeV and the RGI values, should be scale independent, whereas our results show a 15% difference between subsequent scales at which we evaluated our renormalization constant. Taking into account that the size of the coupling constant, i.e. the lattice spacing of our ensemble, is rather large and we use only a two-loop approximation of the running at the scale as low as ≈ 1.5 GeV, we conclude that the size of the discrepancies agrees with our expectations. This behavior should be milder on a finer lattice.
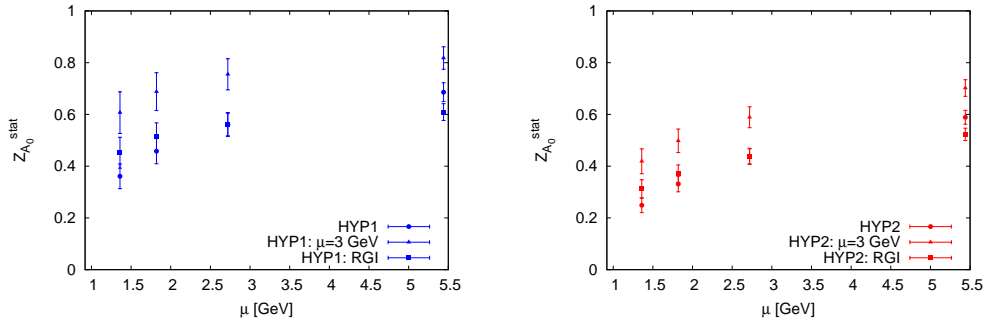

Figure 3: Comparison of $Z_{A_0}^{\mathrm{stat}}$, $Z_{A_0}^{\mathrm{stat},\overline{\mathrm{MS}}}(3\mathrm{GeV})$ and $Z_{A_0}^{\mathrm{stat},\overline{\mathrm{MS}},\mathrm{RGI}}$.

5. Discussion and conclusions

In this contribution we proposed a non-perturbative renormalization scheme for the static

quark theory which can be used on large volume ensembles. The main ingredients of our proposal are renormalization conditions formulated in position space, reduction of cut-off effects through a tree-level improvement and the use of stochastic wall sources to decrease the statistical error of heavy-light correlators. The renormalization conditions are formulated using a single correlation function, which assures that they are gauge invariant and on-shell. Hence, the wave function renormalization constants are not needed. The main drawback of our renormalization scheme is that it suffers from the same window problem as the well-known RI-MOM scheme. This short-comming can be aleviated by using finer lattices or additional ensembles which then allow to run non-perturbatively the renormalization constants through the step scaling functions to higher scales where the perturbative matching to $\overline{\mathrm{MS}}$ could be done more reliably.

In the future, we will provide a companion RI-MOM study and utilize a step-scaling setup.

Acknowledgments

References

[1] C. Sachrajda, Nucl.Phys.Proc.Suppl. 185 (2008) 62-67

[2] J. Heitger et al., JHEP 0402 (2004) 022, hep-lat/0310035
B. Blossier et al., JHEP 1006 (2010) 002, arXiv:1001.4783
B. Blossier et al., JHEP 1209 (2012) 132, arXiv:1203.6516

[3] Y. Aoki, T. Ishikawa, T. Izubuchi, C. Lehner, A. Soni, Phys. Rev. D 91 (2015) 114505, arXiv:1406.6192

[4] T. Ishikawa, Y. Aoki, J. Flynn, T. Izubuchi, O. Loktik, JHEP 1105 (2011) 040, arXiv: 1101.1072

[5] V. Gimenez, G. Martinelli, C. Sachrajda, M. Telavi, unpublished notes

[6] N. H. Christ et al., Phys. Rev. Lett. 105 (2010) 241601, arXiv:1002.2999
T. Blum et al., Phys. Rev. D 84, (2011) 114503, arXiv:1106.2714

[7] E. Eichten, B. Hill, Phys. Lett., B234, 511
E. Eichten, B. Hill, Phys. Lett., B240, 193
E. Eichten, B. Hill, Phys. Lett., B243, 427

[8] A. Hasenfratz, F. Knechtli, Phys. Rev., D64, 034504
M. Della Morte, A. Shindler, R. Sommer, JHEP 0508 (2005) 051, hep-lat/0506008

[9] E. Bagan, P. Ball, W.M. Braun, H.G. Dosch, Phys. Lett. B 278 (1992) 457