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A roadmap for PERC cell efficiency towards 22%, focused on technology-related constraints

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Abstract

Presently, the crystalline silicon (c-Si) photovoltaic (PV) industry is switching from standard cells to PERC cells to increase cell efficiency from about 18% to about 20%. This paper gives a roadmap for increasing PERC cell efficiency further towards 22%. Which equipment and which process conditions are feasible to go beyond 20% efficiency? To help answer this as generally as possible, we conduct state-of-the-art modelling in which we sweep the inputs that represent major *technology-related constraints*, such as diffusion depth, metal finger width and height, alignment tolerances, etc. (these are assigned to the x- and y-axes of our graphs). We then predict the *optimum device parameters resulting from these restrictions* (shown as contour lines). There are many different ways to achieve 22%. Our modelling predicts, for example, that 60 μ m wide screen-printed metal fingers are sufficiently narrow if the alignment tolerance (width of the n⁺⁺ region) is below 90 μ m. The rear may be contacted with 30 μ m wide openings of the Al₂O₃/SiN_x stack and with local J_{0,BSF} values as high as 900 fA/cm². If these requirements cannot be met, they may be compensated by improvements in other device parts. Regardless of this, the wafer material requires a SRH lifetime of at least 1 ms at excess carrier densities near 10¹⁴ cm⁻³.

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1. Evolutionary or revolutionary development?

In the past, the c-Si PV industry went from 14% to 18% efficiency on an evolutionary path, where only single fabrication processes were optimised or altered. It is presently going from about 18% towards 20% by switching from standard cells to PERC cells. The switch cannot be achieved by an evolutionary approach, because both the front and the rear must be improved at the same time to make PERC design effective. Therefore, PERC cells need some new fabrication equipment. We show in this paper, however, that going beyond 20% can be achieved on an

evolutionary path again. It is therefore the main purpose of this abstract to give quantitative foundations for choosing the new PERC production equipment so it stays useful in the near future for achieving 22% efficiency. For our predictions, we use state-of-the-art SENTAURUS [1] device modelling, combined with the EDNA and GRID calculators [2], and SPICE models.

2. Reaching 22% efficiency on a standard path

From the Shockley diode equation, it follows that to reach 22% efficiency, the total saturation current J_0 must be no more than 200 fA/cm². This conclusion arises by assuming $J_{sc} = 40 \text{ mA/cm}^2$, an ideality factor m = 1, and a total $R_s = 0.4 \ \Omega \cdot \text{cm}^2$. (A lower J_0 is required if m > 1 or if R_s is larger.) Since the Shockley equation serves only as a rough guide, we do not discuss its assumptions.

Having set an upper limit to the J_0 of a 22% solar cell, we next set a target J_0 for the three major components of the solar cell: the emitter J_{0e} , the base $J_{0,base}$, and the rear $J_{0,rear}$. We can then examine how those targets might be met by today's equipment—or perhaps tomorrow's equipment.

In the following example, the target J_0 of the three components is chosen to be $J_{0e} \approx 80$, $J_{0,\text{base}} \approx 50$, and $J_{0,\text{rear}} \approx 70$ fA/cm², but there is freedom in dividing the components.

2.1. Constraints on the emitter

The J_{0e} of today's emitters is about 150-200 fA/cm². In our quest for $J_{0e} = 80$ fA/cm² and a 22% solar cell, it must be substationally improved. The J_{0e} depends not only on the emitter's dopant profile, but also on the amount of phosphorus precipitates [3] introduced during fabrication, and on the quality of the surface passivation [4].

Because it is difficult to reach 22% cell efficiency with a homogeneous emitter and conventional screen-printing, we choose a selective emitters. J_{0e} consists of three parts: the metal contacts $J_{0,met}$, the passivated n^{++} part $J_{0,n++}$, and the passivated n^{+} part $J_{0,n++}$, which sum to give J_{0e} respective to their area fraction:

$$J_{0e} = \frac{w_{met}}{p_{\rm f}} J_{0,\rm met} + \frac{w_{\rm n++}}{p_{\rm f}} J_{0,\rm n++} + \frac{p_{\rm f} - w_{\rm met} - w_{\rm n++}}{p_{\rm f}} J_{0,\rm n+}$$
(1)

where p_f is the front finger pitch, w_{mel} and w_{n++} are the width of the contact or n^{++} region, respectively. In the following, we assume that the emitter is fabricated with a POCl₃ diffusion, and we use the experimental values from a recent study [5] shown in Fig. 1(a). Using these values, we optimise Eq. (1) and cross-check with SENTAURUS.



Fig.1. Experimentally achieved J_{0e} by POCl₃ diffusions [5] in dependence of ρ_{em} .

For the simulation of selective emitters, we examine how the $J_{\theta e}$ depends on two technological constraints. Firsly, the alignment tolerance of the metal fingers w_{tol} (in units of μm) is varied; this tolerance implies that the n⁺⁺ region must be wider than the metal contact by w_{tol} . And secondly, the width of the metal contacts w_f is varied. Fig. 2 presents the results, giving a contour plot of J_{0e} and ρ_{n+} as as a function of w_{tol} (x-axis) and w_f (y-axis). It was simulated for $p_f = 1.1 \text{ mm}$, $J_{0,n+} = 60 \text{ fA/cm}^2$ and $\rho_{n+} \approx 120 \Omega/\text{sq}$.

In Fig. 2, we see that as w_{tol} decreases, the optimum ρ_{n++} decreases and a lower J_{0e} is attained. We find that, for example, the total $J_{0e} = 80$ fA/cm² (or 90 fA/cm²) can indeed be achieved with 60 µm wide metal contacts, if the alignment tolerance is maximally $w_{tol} \approx 50$ µm (or 90 µm). In this case, $\rho_{n++} \approx 40$ Ω/sq (or ≈ 50 Ω/sq) is optimum. Note, how sensitively the w_{tol} impacts on J_{0e} . However, an alignment tolerance down to 30 µm seems to be feasible [6].



Fig. 2. Simulation of a selective emitter with $J_{n+} = 60 \text{ fA/cm}^2$. The minimally achievable total J_{0e} (solid lines) is plotted with its corresponding optimum ρ_{n+} (dashed lines), in dependence of the alignment tolerance of the front metal finger, and the width of the front contact (the technological restrictions). A front finger distance of $p_f = 1.1 \text{ mm}$ is assumed (from Fig. 3).

2.2. Optimization of the metal grid

Optimizing the finger pitch is the next task. It affects the resistive losses in the emitter and in the front metallization, as well as J_{sc} , and is therefore a highly non-linear optimization task. The series resistance contribution (in units of Ωcm^2) of the selective emitter can be calculated as

$$R_{s.em} = \frac{1}{3l_{\rm f}(2n_{\rm f}2n_{\rm bb})L^2} \left[\rho_{\rm n++}L_{\rm n++}L^2 \left(1 + \frac{L_{\rm n+}}{L} + \frac{L_{\rm n+}^2}{L^2} \right) + \rho_{\rm n+}L_{\rm n+}L_{\rm n+}^2 \right] A_{\rm cell}$$
(2)

where l_f is the finger length, $n_f (n_{bb})$ is the number of fingers (busbars) in the cell, L is half of the finger separation (which is the longest way electrons travel laterally in the emitter), A_{cell} the cell area, and there is $L_{n+} = w_{n+}/2$, $L_{n++} = w_{n++}/2$, and $L = p_f - w_{met}$.

We now determine how the efficiency of the solar cell depends on technological constraints placed by the screenprinting technology. The results are plotted in Fig. 3(a), which plots efficiency contours against the finger width w_f (x-axis) and the finger height (y-axis). In these simulations, series resistance contribution of the metal grid is calculated with the freeware GRID [7] for a standard resistivity of screen-printed Ag fingers of $4.5 \times 10^{-6} \Omega$ cm; the shading is assumed to be 50% of the finger coverage and 90% of the busbar coverage, as was measured in a module [8]. For each simulation, the optimal finger pitch is selected, as plotted by dashed lines.

Fig. 3(a) shows that with 60 μ m wide and about 25 μ m high fingers, 22% is possible if the front finger pitch p_f

is about 1.1 mm. Narrower fingers (and accordingly smaller p_f) are only necessary if going above 22% cell efficiency, and instead, multi-wire busbars would be similarly effective.



Fig. 3. (a) Predicted optimum front finger pitch p_f (blue dashed lines) with the resulting maximally achievable cell efficiency (red solid lines) in dependence of the front finger width and front finger height (the technological restrictions). (b) Simulated demands of the rear metallisation, BSF and passivation: the optimum rear finger pitch (blue dashed lines) with the resulting maximally achievable cell efficiency (red lines). Note that these predictions depend on the emitter quality.

2.3. Optimization of the rear

To achieve 22% cell efficiency, the losses in the local BSFs and at the rear surface should combine to give a $J_{0,rear}$ of no more than 70 fA/cm² (as stated in the beginning). Fig. 3(b) shows how this may be achieved assuming 30 µm wide rear contact openings. For example, a $J_{0,rear}$ of 70 fA/cm² and therefore an efficiency of 22.0% arises when $J_{0,BSF} \approx 800$ fA/cm² and $S_{eff} \approx 10$ cm/s, where the latter has been experimenally demonstrated for Al₂O₃/SiN_x passivation. In these simulations, the rear finger pitch that yields the maximised efficiency is selected—this is plotted as the dashed lines in Fig.3(b). In the example case, the optimum rear finger distance then is 850 µm, but it may be chosen as wide as 1.2 mm without compromising on cell efficiency (not shown here).

We note that with $S_{\text{eff}} = 100 \text{ cm/s}$, only 21.5% efficiency is achievable, unless the emitter or the front metallization are made better than in the example of Fig. 2.

2.4. Wafer quality requirements

In all of our simulations, we assume a SRH lifetime τ in the wafer of 1 - 2 ms. With lower τ , the equivalent $J_{0,\text{base}}$ exceeds 40 fA/cm². This lifetime is presently achievable in boron-doped Cz materials by applying deactivation procedures [9]. It has been also achieved in Ga-doped or magnetic Cz materials.

3. Conclusions

To reach 22% cell efficiency, 60 μ m wide and 25 μ m high front metal fingers with a pitch of about 1.1 mm are sufficient, if an emitter saturation current density $J_{0e} = 80$ fA/cm² (including contacts) is realized, e.g. by a selective emitter with 60 μ m wide n⁺⁺ regions (with $\rho_{n++} \approx 40 \Omega/sq$), and with an n⁺ region that has $J_{0,n+} = 60$ fA/cm², and $\rho_{n+} \approx 120 \Omega/sq$. At the rear, $S_{eff} < 30$ cm/s is favourable, as otherwise the front must be improved to compensate. Such a rear structure may be realised by a passivation with S = 10 cm/s, about 30 μ m wide openings with a pitch near 1 mm, and local Al-BSFs with $J_{0,BSF} < 700$ fA/cm². The main issue is the bulk lifetime τ : it must reach about 2 ms at the

prevailing excess carrier densities at MPP, which are near 1×10^{13} cm⁻³ (and at open-circuit conditions about 2×10^{15} cm⁻³) in 1 Ω cm material.

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