# The R136 star cluster dissected with Hubble Space Telescope/STIS. I. Far-ultraviolet spectroscopic census and the origin of He iI $\lambda 1640$ in young star clusters 

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#### Abstract

We introduce a Hubble Space Telescope (HST)/Space Telescope Imaging Spectrograph (STIS) stellar census of R136a, the central ionizing star cluster of 30 Doradus. We present low resolution far-ultraviolet STIS spectroscopy of R136 using 17 contiguous $52 \operatorname{arcsec} \times 0.2 \operatorname{arcsec}$ slits which together provide complete coverage of the central 0.85 parsec ( 3.4 arcsec ). We provide spectral types of 90 per cent of the 57 sources brighter than $m_{\text {F555W }}=16.0 \mathrm{mag}$ within a radius of 0.5 parsec of R136a1, plus 8 additional nearby sources including R136b (O4 If/WN8). We measure wind velocities for 52 early-type stars from C iv $\lambda \lambda 1548-51$, including $16 \mathrm{O} 2-3$ stars. For the first time, we spectroscopically classify all Weigelt and Baier members of R136a, which comprise three WN5 stars (a1-a3), two O supergiants (a5-a6) and three early O dwarfs (a4, a7, a8). A complete Hertzsprung-Russell diagram for the most massive O stars in R136 is provided, from which we obtain a cluster age of $1.5_{-0.7}^{+0.3} \mathrm{Myr}$. In addition, we discuss the integrated ultraviolet spectrum of R136, and highlight the central role played by the most luminous stars in producing the prominent $\mathrm{He}_{\text {II }} \lambda 1640$ emission line. This emission is totally dominated by very massive stars with initial masses above $\sim 100 \mathrm{M}_{\odot}$. The presence of strong He $_{\text {II }} \lambda 1640$ emission in the integrated light of very young star clusters (e.g. A1 in NGC 3125) favours an initial mass function extending well beyond a conventional upper limit of $100 \mathrm{M}_{\odot}$. We include montages of ultraviolet spectroscopy for Large Magellanic Cloud O stars in the appendix. Future studies in this series will focus on optical STIS medium resolution observations.


Key words: stars: early-type - stars: massive - stars: winds, outflows - Magellanic Clouds galaxies: star clusters: individual: R136a.

## 1 INTRODUCTION

The formation of massive stars remains an unsolved problem in astrophysics. This is especially challenging since very massive stars (VMS, $\geq 100 \mathrm{M}_{\odot}$ ) are usually located within young massive star clusters ( $\geq 10^{4} \mathrm{M}_{\odot}$ ) with small core radii ( $r_{\mathrm{c}} \sim 0.1-0.2 \mathrm{pc}$; Portegies Zwart, McMillan \& Gieles 2010). The short lifetime of such stars ( $2-3 \mathrm{Myr}$; Yusof et al. 2013) restricts suitable Milky Way star clusters to only a handful, including the Arches (Martins et al. 2008) and NGC 3603 (Schnurr et al. 2008), both of which are

[^0]highly reddened. Therefore, observations of young star clusters at ultraviolet (UV) wavelengths - corresponding to the peak output of hot, luminous stars - are largely restricted to extragalactic cases.
R136 was included in the census of optically bright Magellanic Cloud stars by Feast, Thackeray \& Wesselink (1960), and subsequently resolved into three components, a, b and c (Feitzinger et al. 1980), with component a subsequently shown to be multiple (Weigelt \& Baier 1985). R136a is the closest example of an extragalactic young massive star cluster (Hunter, Shaya \& Holtzman 1995). It is the central cluster of the NGC 2070 complex within the 30 Doradus star-forming region in the Large Magellanic Cloud (LMC; Massey \& Hunter 1998; Doran et al. 2013), and is relatively lightly reddened, facilitating study at UV wavelengths (de Marchi
et al. 1993). However, the 0.1 pc core radius of R136a (Portegies Zwart et al. 2010; Selman \& Melnick 2013) subtends only 0.41 arcsec at the 50 kpc distance to the LMC (Pietrzyński et al. 2013), which led to the cluster at one time being mistaken for a $\geq 1000$ $\mathrm{M}_{\odot}$ star (Cassinelli, Mathis \& Savage 1981; Feitzinger, Hanuschik \& Schmidt-Kaler 1983), although Walborn (1973) had earlier predicted that R136a would have a structure analogous to the core of NGC 3603. Consequently, observations at high spatial resolution are required to resolve individual stars within the core of R136a, requiring either Hubble Space Telescope (HST) or large ground-based telescopes with adaptive optics (VLT/MAD; Campbell et al. 2010).

The star formation history of NGC 2070 has recently been investigated by Cignoni et al. (2015) using Hubble Tarantula Treasury Project imaging, revealing an upturn in star formation $\sim 7 \mathrm{Myr}$ ago, peaking 2-3 Myr ago, and a total stellar mass of $9 \times 10^{4} \mathrm{M}_{\odot}$ assembled within 20 pc of R136. The mass function for intermediateand low-mass stars in R136 have been obtained from HST imaging (Hunter et al. 1995; de Marchi et al. 2011). However, the colour/reddening degeneracy of hot massive stars at UV and optical wavelengths necessitates spectroscopy for robust physical properties, and in turn, masses, measuring present-day mass functions (though see Maíz Apellániz et al. 2014). To date, Koornneef \& Mathis (1981) and Vacca et al. (1995) have presented UV scans across NGC 2070 with the International Ultraviolet Explorer (IUE), while UV spectroscopy of individual stars of R136a has been limited to the brightest members with $H S T /$ GHRS (de Koter, Heap \& Hubeny 1997, 1998).

Meanwhile, Massey \& Hunter (1998) have obtained optical spectroscopy for dozens of stars within NGC 2070 using HST/FOS, which revealed large numbers of very early O stars, indicating a young cluster age ( $<1-2 \mathrm{Myr}$ ). More recently Schnurr et al. (2009) used VLT/SINFONI to obtain spatially resolved $K$-band spectroscopy of the brightest sources, namely R136a1, a2, a3 and c, which are WN-type emission-line stars. Analysis of the VLT and $H S T$ spectroscopy plus VLT/MAD imaging led Crowther et al. (2010) to conclude that such stars possess very high (initial) stellar masses of $160-320 \mathrm{M}_{\odot}$, results largely supported by Bestenlehner et al. (2011) and Hainich et al. (2014).

The existence of such high-mass stars in R136 challenges the previously accepted $\sim 150 \mathrm{M}_{\odot}$ upper stellar mass limit (Figer 2005), unless they are products of stellar mergers (Banerjee, Kroupa \& Oh 2012; Schneider et al. 2014b). If such stars are indeed merger remnants one would expect younger apparent ages for this with respect to lower mass stars, and differences in the present day mass function with respect to a Salpeter slope (Schneider et al. 2014b; Schneider et al. 2015). Massey \& Hunter (1998) suggested that the mass function of R136 is normal, although their study only included a subset of the brightest stars in the vicinity of R136a.

A census of the hot, luminous stars in R136a with Space Telescope Imaging Spectrograph (STIS) is the focus of the present series of papers. We discuss the suitability of UV spectroscopy for spectral classification of O stars, based in part upon spectroscopic data sets in the usual blue visual wavelength range. Wind velocities for O stars can uniquely be obtained from P Cygni profiles in the far-UV (e.g. Groenewegen, Lamers \& Pauldrach 1989; Prinja, Barlow \& Howarth 1990). Comparisons of empirical (reduced) wind momenta ${ }^{1}$ with theoretical predictions usually rely on approximate $v_{\infty}$ calibrations (e.g. Lamers, Snow \& Lindholm 1995). Indeed,

[^1]high-quality far-UV spectroscopy has been restricted to relatively few LMC O stars to date (Prinja \& Crowther 1998; Walborn et al. 1995; Massey et al. 2005). One byproduct of this study is a calibration of wind velocities for (mostly) early O dwarfs at the $\sim 0.5 \mathrm{Z} \odot$ metallicity of the LMC.

In the first study in this series, we introduce our new STIS observations of R136a in Section 2. We present and analyse UV STIS/Multi-Anode Microchannel Array (MAMA) spectroscopy of individual stars in Section 3, including approximate spectral types and estimates of terminal wind velocities. An UV spectral atlas of LMC O stars is compiled using literature spectra, in a few instances complemented with new STIS/Charged Couple Device (CCD) spectroscopic data sets to provide templates of missing spectral types. The current data set also allows the contributions from individual stars to its integrated UV spectrum to be quantified. This is discussed in Section 4, of relevance for other young massive star clusters in the local universe (e.g. Chandar, Leitherer \& Tremonti 2004; Leitherer et al. 2011) and the integrated rest-frame UV spectrum of highredshift star-forming galaxies (e.g. Shapley et al. 2003). Finally, brief conclusions are drawn in Section 5.

## 2 OBSERVATIONS

Here we present our new long-slit HST STIS observations of R136, and use of archival WFC3/UVIS imaging discussed by de Marchi et al. (2011).

### 2.1 UV spectroscopy

We have obtained UV STIS/G140L spectroscopy using the far-UV MAMA detector and a $52 \operatorname{arcsec} \times 0.2$ arcsec slit at 17 contiguous pointings in R136, identified by their position within the cluster, namely SE9, SE8, ...SE1, NW1, ...NW7, NW8 (Fig. 1). The spectral coverage of the G140L grating is $\lambda \lambda 1150-1700 \AA$, with a pixel scale of $0.6 \AA \mathrm{pixel}^{-1}$, and notional point source resolving power of 1250 at $\lambda=1500 \AA$. Since the MAMA plate scale is 0.024 arcsec pixel ${ }^{-1}$, a slit length of $\sim 25$ arcsec is sampled in all cases. All MAMA exposures were obtained in two visits, 4 and 5 from GO 12465, totalling 6 orbits, as set out in Table 1. All observations were obtained at a fixed position angle $64^{\circ}$ ( E of N ), to ensure that R136a1 and R136a2 lie within the same slit position (NW1) as shown in Fig. 1. For each visit, the nearby isolated star Melnick 34 was acquired using the CCD and F28X50LP aperture, followed by a peakup with the CCD and $52 \operatorname{arcsec} \times 0.1$ arcsec aperture, followed by an offset to individual pointings, with two exposures acquired per pointing, the second spatially offset along the slit by 11 pix ( 0.27 arcsec), with integration times shown in Table 1.

MAMA exposures were combined and extracted using the specialized stistools package (Bostroem \& Proffitt 2011). No automated wavelength correction was made for sources not centred in individual slits, while standard calstis slit loss corrections were applied. Therefore, absolute wavelength scales may be incorrect by up to $\pm 2$ pixel ( $0.56 \AA$ pixel $^{-1}$ ) while absolute far-UV fluxes will be underestimated by up to a factor of 2 in extreme cases.

Extractions were also independently made using multispec (Maíz Apellániz 2005; Knigge et al. 2008) which is particularly suited to the crowded region at the geometric centre of the cluster. Fig. 2 shows 2D spectral data sets for particularly crowded slits NW1 (a1, a2 and a5) and NW3 (a4 and a8), including collapsed spectral images between $\lambda \lambda 1475-1525$. multispec extractions relied upon empirical point spread functions for G140L/MAMA spectroscopy employing the position of individual sources in slits from short


Figure 1. HST/STIS slits ( $52 \operatorname{arcsec} \times 0.2$ arcsec) superimposed upon an F336W WFC3/UVIS image of R136 (de Marchi et al. 2011), oriented with north up and east to left, together with a circle of radius $4.1 \mathrm{arcsec}(1 \mathrm{parsec})$ and identification of the acquisition star Melnick 34. The active slit length for MAMA observations is the central 25 arcsec. The zoom highlights the central region, including identification of individual slits and the integrated R136a cluster ( 2.05 arcsec radius circle, centred upon R136a1 equivalent to 0.5 parsec at the distance of the LMC).

Table 1. Log of HST/STIS spectroscopic observations of R136 from Cycles 19 (GO 12465) and 20 (GO 13052).

| GO <br> /visit | Orbits | Target | Detector | Grating <br> /position | PA <br> $\circ$ | $T_{\exp }$ <br> (sec) | Date |
| :--- | :--- | :--- | :--- | :--- | ---: | :---: | ---: |
| $12465 / 1$ | 4 | SE9... NW1 | CCD | G430M/3936 | 64 | $2 \times(440-570)$ | 2012 April 6 |
| $12465 / 2$ | 4 | SE9...NW1 | CCD | G430M/4451 | 64 | $2 \times(440-570)$ | 2012 April 4 |
| $12465 / 3$ | 4 | SE9...NW1 | CCD | G430M/4706 | 64 | $2 \times(440-570)$ | 2012 April 6 |
| $12465 / 4$ | 4 | SE9...NW3 | MAMA | G140L | 64 | $2 \times(400-470)$ | 2012 April 7 |
| $12465 / 5$ | 2 | NW4...NW8 | MAMA | G140L | 64 | $2 \times(470-520)$ | 2012 April 8 |
| $12465 / 6$ | 5 | SE9...NW8 | CCD | G750M/6581 | 244 | $2 \times(350-370)$ | 2012 October 23 |
| $12465 / 7$ | 3 | NW2...NW8 | CCD | G430M/3936 | 244 | $2 \times 540$ | 2012 October 20 |
| $12465 / 8$ | 3 | NW2...NW8 | CCD | G430M/4451 | 244 | $2 \times 540$ | 2012 October 20 |
| $12465 / 9$ | 3 | NW2...NW8 | CCD | G430M/4706 | 244 | $2 \times 540$ | 2012 October 20 |
| $13052 / 1$ | 4 | SE9...NW1 | CCD | G430M/4194 | 244 | $2 \times(440-570)$ | 2012 October 21 |
| $13052 / 2$ | 3 | NW2...NW8 | CCD | G430M/4194 | 64 | $2 \times 540$ | 2013 April 3 |

exposure HST WFC3/UVIS F336W imaging of R136 from 2009 October (for details, see de Marchi et al. 2011). The latter provided wavelength corrections if sources were not centred in individual slits, but a standard slit loss correction is again adopted.

For relatively isolated sources, calstis and multispec extractions agree extremely well, while the latter approach usually performed better in situations where sources were closely spaced within individual slits. For consistency, spectroscopic data sets presented here were extracted with multispec, except for R136 H70, 141, 149 for which mULTISPEC extractions were problematic.

### 2.2 Optical spectroscopy

We have obtained optical STIS spectroscopy using the CCD detector and identical $52 \operatorname{arcsec} \times 0.2 \operatorname{arcsec}$ slit at identical positions and position angle (or $\mathrm{PA}+180^{\circ}$ ) to the UV spectroscopy, although the coarser plate scale of $0.05 \mathrm{arcsec}_{\text {pixel }}{ }^{-1}$ provided complete spatial
coverage of the slit length. In total, two exposures at four grating positions with the G 430 M grating provided complete blue coverage $\lambda \lambda 3793-4849 \AA$, with a pixel scale of $0.28 \AA$ pixel $^{-1}$, and point source resolving power of 7700 at $\lambda=4400 \AA$. A single grating position with the G750M grating provided $\lambda \lambda 6482-7054 \AA$ spectroscopy, with a pixel scale of $0.56 \AA$ pixel $^{-1}$, and point source resolving power of 6000 at $\mathrm{H} \alpha$. For all CCD observations, the second exposure was offset spatially along the slit by 7 pixel ( 0.35 arcsec ). Three of the four G430M observations were taken over the course of a few days, with the fourth (G430M/4194) obtained approximately 6 months later to enable to search for low-period binaries. All G750M exposures were taken at a single epoch.
In total, CCD exposures were split between six visits during Cycle 19 (GO 12465, 26 orbits), plus two visits during Cycle 20 (GO 13052, 7 orbits), also set out in Table 1. mUltispec and Calstis extractions of CCD observations will be the focus of the next paper in this series (Caballero-Nieves et al., in preparation).


Figure 2. (Top): 2D spectral data set for NW3 (left) and NW1 (right) slits, oriented such that the $\lambda \lambda 1150-1700$ dispersion axis runs from top to bottom (horizontal line near top is geocoronal Ly $\alpha$ emission, occulting bar lies to the right of each image) and the ( $25 \mathrm{arcsec}=6.2 \mathrm{pc}$ ) spatial axis is oriented from left (SW) to right (NE). Prominent (stellar plus interstellar) absorption features include Ly $\alpha$ (top) and C iv $\lambda 1548-51$ (bottom); (centre): collapsed spectral image for NW3 (left) and NW1 (right) between $\lambda \lambda 1475-1525$; (bottom) $2.5 \operatorname{arcsec}(0.62 \mathrm{pc})$ zoom of collapsed spectral image for NW3 (left) and NW1 (right) indicating the dense cores, resolving R136a1 and a2, and R136a4 and a8, respectively, together with multispec fits (faint grey line).

LMC stars that have been observed in the UV and for which reliable optical classifications are available are fairly scarce. Our sample allows 10 stars with robust optical classifications spanning the full O2-O9.5 subtype sequence, which are included here as LMC templates (see Section 3).

### 2.3 Stellar census of R136

Hunter et al. (1995) have used $H S T /$ WFPC2 to identify the visually brightest sources in R136, while deeper HST/WFC3 images have been analysed by de Marchi et al. (2011) to derive multicolour photometry. We compare WFPC2 and WFC3/UVIS photometry in F336W, F555W and F814W filters for all sources in common within $10 \operatorname{arcsec}(2.4 \mathrm{pc})$ of R136a1 in Fig. 3. Overall, the earlier WFPC2 results are confirmed by WFC3, albeit with modest systematic offsets of $m_{\text {F336W }}($ WFPC2 $)-m_{\text {F336W }}(\mathrm{WFC} 3)=+0.15 \pm$ $0.20 \mathrm{mag}, m_{\text {F555W }}(\mathrm{WFPC} 2)-m_{\text {F555W }}(\mathrm{WFC} 3)=+0.17 \pm 0.21 \mathrm{mag}$, $m_{\mathrm{F} 14 \mathrm{~W}}(\mathrm{WFPC} 2)-m_{\mathrm{F} 14 \mathrm{~W}}(\mathrm{WFC} 3)=+0.18 \pm 0.22$ mag. In sources
for which de Marchi et al. (2011) photometry is unavailable, we have substituted Hunter et al. (1995) results to which these offsets have been applied. Overall, agreement between early WFPC2 and new WFC3 photometry is generally very good, with a few notable exceptions - WFC3 photometry from de Marchi et al. (2011) is significantly brighter for R136a1, a2 and a6.

In Table 2, we provide a catalogue of all 55 sources within $2.05 \operatorname{arcsec}(0.5 \mathrm{pc})$ of R136al brighter than $m_{\mathrm{F} 555 \mathrm{~W}}=16.0 \mathrm{mag}$, while an additional 8 bright sources observed with STIS/MAMA are presented in Table 3. Additional sources brighter than $m_{\mathrm{F} 555 \mathrm{~W}}=$ 17.0 mag within $4.1 \operatorname{arcsec}(1 \mathrm{pc})$ are listed in Appendix D (Table D1). There is evidence for a non-standard extinction law in 30 Doradus, and significant variation thereof (Maíz Apellániz et al. 2014). Specifically for R136, Doran et al. (2013) favoured $R_{\mathrm{V}} \sim 4.2 \mathrm{mag}$ for individual stars based on a combination of optical (WFC3/UVIS) and near-IR (VLT/MAD, Campbell et al. 2010) photometry. We have applied chorizos to 41 O stars with WFC3 photometry and approximate spectral types, and obtained good fits in 34 cases,


Figure 3. Comparison between WFPC2 (Hunter et al. 1995) and WFC3/UVIS photometry (de Marchi et al. 2011) for all sources in common within 10 arcsec of R136a1 (O stars: triangles, WN stars: squares) for F336W (top), F555W (middle) and F814W (bottom). Agreement is generally excellent aside from R136a1, a2 and a6 (see Section 2.3).
from which average values of $R_{5495}=4.1 \pm 0.5$ and $A_{\mathrm{F555W}}=$ $1.72 \pm 0.25 \mathrm{mag}$ is obtained, in good agreement with the Tatton et al. (2013) near-IR study, for which $A_{V}=1.75$ mag was inferred for R136. Here, we adopt tailored extinctions of WN5 stars from Crowther et al. (2010), while for the O stars we adopt a reddening law of $R_{5495}=4.1$ and individual extinctions based upon temperature dependent F336W-F555W intrinsic colours from multicolour chorizos (Maíz Apellániz 2004):
$\left(m_{\text {F336W }}-m_{\text {F555W }}\right)_{0}=-1.475-0.155\left(T_{\text {eff }} / 10000\right) \mathrm{mag}$,
plus $A_{\mathrm{F} 555 \mathrm{~W}}=R_{5495} E_{\mathrm{F} 336 \mathrm{~W}-\mathrm{F} 555 \mathrm{~W}} / 2.28$ from the Seaton (1979)/Howarth (1983) LMC extinction law, plus an LMC distance modulus of 18.49 mag (Pietrzyński et al. 2013). For stars without UV derived spectral types, we adopt O6 V ( $\left.T_{\text {eff }}=39900 \mathrm{~K}\right)$ for stars brighter than $m_{\text {F555w }}=16.0 \mathrm{mag}$, and $09.5 \mathrm{~V}\left(T_{\text {eff }}=32900 \mathrm{~K}\right)$ for fainter stars. We adopt $A_{\text {F555W }}=1.72 \mathrm{mag}$ for all sources without WFC3 photometry.
Severe crowding results in significant differences in photometry between Hunter et al. (1995) and de Marchi et al. (2011) for some of the brightest sources in the core of R136a. HST/WFC3 photometry implies F336W/F555W/F814W flux ratios of 0.88/0.95/0.84 for a2/a1, while $K$-band VLT/SINFONI observations favour a significantly lower ratio of 0.76 (Crowther et al. 2010). From our STIS/MAMA spectroscopy, the $1500 \AA$ flux ratio inferred from the multispec extraction is lower still ( 0.56 ), while Gaussian fits to spatial profiles in our 2D UV NW1 data set (lower left panel in Fig. 2) favour 0.62. We adopt this flux ratio, so $m_{\mathrm{F} 555 \mathrm{~W}}^{a 1}-m_{\mathrm{F} 555 \mathrm{~W}}^{a 2}=$ 0.52 under the reasonable assumption that their interstellar extinctions are equivalent. From comparison between the UV fluxes of a1 and a2 and relatively isolated sources in the NW1 slit, we support WFC3/F555W photometry of a1, resulting in a revised a2 apparent magnitude of $m_{\text {F555W }}=12.80 \mathrm{mag}$ in Table 2. This leads to a downward revision of 0.18 dex in stellar luminosity, and corresponding 20 per cent lower mass estimate.

In addition, WFC3 photometry of a6 from de Marchi et al. (2011) is significantly brighter than Hunter et al. (1995). From our MAMA SE1 data set, we obtain a $1500 \AA$ Alux ratio of 0.70 for a6/a3, in perfect agreement with $m_{\mathrm{F} 555 \mathrm{~W}}^{a 3}-m_{\mathrm{F} 555 \mathrm{~W}}^{a 6}=-0.38$ from WFC3, so adopt photometry from de Marchi et al. Finally, de Marchi et al. (2011) provide a single photometric measurement for $24+a 8$, which agrees satisfactorily with the combination of these sources in Hunter et al. (1995). Again, we utilize UV data sets (lower right panel in Fig. 2) to verify their relative fluxes. From our multispec extraction, we obtain a $1500 \AA \AA$ flux ratio of 0.66 for a8/a4, in good agreement with 0.65 from Gaussian fits to their spatial profiles in the NW3 slit. We obtain $m_{\mathrm{F} 555 \mathrm{~W}}^{a 4}-m_{\mathrm{F} 555 \mathrm{~W}}^{a 8}=0.46$, again assuming identical interstellar extinctions for these stars. From the combined WFC3 photometry of $m_{\mathrm{F} 555 \mathrm{~W}}^{a 4+a 8}=13.41 \pm 0.01$, we obtain revised a4 and a8 apparent magnitudes of $m_{\text {F555W }}=13.91$ and 14.42 mag (Table 2).

## 3 SPATIALLY RESOLVED FAR-UV SPECTROSCOPY

We have classified individual R136 stars, based upon a comparison between their UV morphology and LMC template stars with reliable optical classifications. As noted above, we have incorporated 10 reference $\mathrm{O} 2-09.5$ stars based on optical classifications provided in Caballero-Nieves et al. (in preparation) which are indicated in Tables 2 and 3. Our approach therefore complements the UV classification system of LMC OB stars by Smith Neubig \& Bruhweiler (1999) using IUE SWP/LORES spectroscopy. We

Table 2. HST/WFC3 photometry (de Marchi et al. 2011) for sources ( $m_{\mathrm{F} 555 \mathrm{~W}} \leq 16.0 \mathrm{mag}$ ). within a projected distance of 2.05 arcsec ( 0.5 parsec ) from R136a1. In case WFC3 photometry is unavailable, WFPC2 photometry from Hunter et al. (1995) is used, albeit offset in F555W by -0.17 mag (indicated in italic font, see Section 2.3). Identifications are from Weigelt \& Baier (1985, WB85) or Hunter et al. (1995, HSH95), while spectral types are obtained from optical spectroscopy. $A_{\mathrm{F} 555 \mathrm{~W}}$ is either obtained via the observed $m_{\mathrm{F} 336 \mathrm{~W}}-m_{\mathrm{F} 555 \mathrm{~W}}$ colour plus $R_{5495}=4.1$, or $A_{\mathrm{F} 555 \mathrm{~W}}=1.72$ mag is adopted (values for WN5 stars are from Crowther et al. 2010).

| WB85 | HSH95 | Spectral type | Ref | $r$ (arcsec) | $m_{\text {F555W }}$ <br> (mag) | $\begin{gathered} m_{\text {F336W-F438W }} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} m_{\text {F438W-F555W }} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} m_{\text {F555W-F814W }} \\ (\mathrm{mag}) \end{gathered}$ | $A_{\text {F555W }}$ (mag) | $M_{\text {F555W }}$ (mag) | Slit |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| a1 | 3 | WN5h | b | 0.00 | $12.28 \pm 0.01$ | $-1.25 \pm 0.01$ | $0.17 \pm 0.01$ | $0.49 \pm 0.01$ | 1.88 | -8.09 | NW1 |
| a2 | 5 | WN5h | b | 0.08 | $12.80 \pm 0.01^{\ddagger}$ | $-1.23 \pm 0.01$ | $0.23 \pm 0.01$ | $0.37 \pm 0.01$ | 1.83 | $-7.52^{\ddagger}$ | NW1 |
| a5 | 20 | O2 If* | e | 0.28 | $13.71 \pm 0.01$ | $-1.28 \pm 0.02$ | $0.06 \pm 0.01$ | $0.48 \pm 0.01$ | 1.87 | -6.65 | NW1 |
| a7 | 24 | O3 III(f*) | a | 0.36 | $13.97 \pm 0.01$ | $-1.28 \pm 0.02$ | $-0.04 \pm 0.02$ | $0.30 \pm 0.02$ | 1.58 | -6.10 | NW3 |
| a8 | 27 |  |  | 0.39 | $14.42 \pm 0.01^{\ddagger}$ | $-1.30 \pm 0.01$ | $0.12 \pm 0.01$ | $0.27 \pm 0.01$ | 1.98 | $-6.05^{\ddagger}$ | NW3 |
| a4 | 21 |  |  | 0.43 | $13.96 \pm 0.01^{\ddagger}$ |  |  |  |  | $-6.51{ }^{\ddagger}$ |  |
|  | 186 |  |  | 0.43 | $15.63 \pm 0.02$ | $-1.23 \pm 0.03$ | $0.10 \pm 0.03$ | $0.35 \pm 0.04$ | 1.75 | -4.62 | NW1 |
|  | 86 |  |  | 0.45 | $14.73 \pm 0.01$ | $-1.42 \pm 0.02$ | $0.41 \pm 0.02$ |  | 2.15 | -5.91 | NW1 |
| a3 | 6 | WN5h | b | 0.49 | $12.97 \pm 0.01$ | $-1.31 \pm 0.01$ | $0.16 \pm 0.01$ | $0.55 \pm 0.01$ | 1.87 | -7.39 | SE1 |
|  | 66 | O 3 V | a | 0.49 | $15.06 \pm 0.02$ | $-1.29 \pm 0.02$ | $0.17 \pm 0.02$ | $0.29 \pm 0.03$ | 1.95 | -5.38 | SE2 |
|  | 119 |  |  | 0.51 | 15.75 |  |  |  | 1.72 | -4.46 | NW3 |
|  | 58 | O3 III(f*) | a | 0.59 | $14.80 \pm 0.01$ | $-1.30 \pm 0.02$ | $0.11 \pm 0.02$ | $0.34 \pm 0.02$ | 1.96 | -5.65 | NW1 |
|  | 30 | 07 V | a | 0.62 | $14.21 \pm 0.01$ | $-1.31 \pm 0.01$ | $0.13 \pm 0.01$ | $0.29 \pm 0.01$ | 1.61 | -5.89 | NW4 |
|  | 70 | O 5 Vz | f | 0.62 | $14.96 \pm 0.01$ | $-1.27 \pm 0.02$ | $0.13 \pm 0.02$ | $0.37 \pm 0.03$ | 1.79 | -5.32 | NW2 |
|  | 89 |  |  | 0.66 | $14.76 \pm 0.01$ | $-1.11 \pm 0.02$ | $0.34 \pm 0.02$ | $0.20 \pm 0.02$ | 2.52 | -6.25 | SE2 |
|  | 62 |  |  | 0.66 | $14.91 \pm 0.01$ | $-1.36 \pm 0.02$ | $0.11 \pm 0.02$ | $0.31 \pm 0.02$ | 1.85 | -5.43 | SE3 |
| a6 | 19 |  |  | 0.73 | $13.35 \pm 0.01$ | $-1.26 \pm 0.01$ | $0.07 \pm 0.01$ | $0.46 \pm 0.01$ | 1.82 | -6.96 | SE1 |
|  | 50 |  |  | 0.74 | $14.65 \pm 0.01$ | $-1.25 \pm 0.01$ | $0.11 \pm 0.01$ | $0.37 \pm 0.02$ | 2.05 | -5.89 | SE3 |
|  | 90 | O5 V:: | a | 0.77 | $15.48 \pm 0.02$ | $-1.36 \pm 0.03$ | $0.12 \pm 0.03$ | $0.34 \pm 0.04$ | 1.67 | -4.68 | NW3 |
|  | 149 |  |  | 0.78 | $15.44 \pm 0.02$ | $-1.28 \pm 0.03$ | $0.17 \pm 0.03$ | $0.26 \pm 0.04$ | 1.96 | -5.01 | SE2 |
|  | 141 |  |  | 0.78 | $15.82 \pm 0.03$ | $-1.35 \pm 0.03$ | $0.09 \pm 0.03$ | $0.29 \pm 0.05$ | 1.55 | -4.22 | SE4 |
|  | 80 | O8 V | f | 0.87 | $15.17 \pm 0.02$ | $-1.30 \pm 0.02$ | $0.11 \pm 0.02$ | $0.25 \pm 0.03$ | 1.53 | -4.85 | SE3 |
|  | 35 | O3 V | f | 0.88 | $14.32 \pm 0.01$ | $-1.43 \pm 0.01$ | $0.11 \pm 0.01$ |  | 1.64 | -5.81 | NW5 |
|  | 78 | O4: V | f | 0.97 | $15.26 \pm 0.02$ | $-1.33 \pm 0.02$ | $0.11 \pm 0.02$ | $0.29 \pm 0.03$ | 1.70 | -4.93 | NW3 |
|  | 73 | O9 V | a | 1.01 | $15.13 \pm 0.02$ | $-1.12 \pm 0.02$ | $0.11 \pm 0.02$ |  | 1.78 | -5.14 | SE4 |
|  | 92 | O3 V | a | 1.05 | $15.46 \pm 0.02$ | $-1.36 \pm 0.02$ | $0.10 \pm 0.03$ | $0.30 \pm 0.04$ | 1.52 | -4.55 | NW6 |
|  | 143 |  |  | 1.05 | $15.94 \pm 0.03$ | $-1.30 \pm 0.04$ | $0.19 \pm 0.04$ | $0.36 \pm 0.06$ | 1.70 | -4.25 | SE2 |
|  | 112 | O8.5 III(f) | a | 1.06 | $15.74 \pm 0.03$ | $-1.25 \pm 0.04$ | $0.12 \pm 0.04$ | $0.29 \pm 0.05$ | 1.65 | -4.40 | NW2/3 |
|  | 135 |  |  | 1.08 | $15.90 \pm 0.03$ | $-1.20 \pm 0.04$ | $0.13 \pm 0.04$ |  | 1.67 | -4.26 | NW5 |
|  | 69 | O3-6 V | a | 1.09 | 15.05 | $-1.32 \pm 0.04$ |  |  | 1.72 | -5.16 | SE5 |
|  | 52 | O 3 V | a | 1.13 | $14.72 \pm 0.01$ | $-1.34 \pm 0.01$ | $0.09 \pm 0.01$ | $0.22 \pm 0.02$ | 1.71 | -5.48 | SE5 |
|  | 48 | O2-3 III(f*) | f | 1.22 | $14.75 \pm 0.01$ | $-1.26 \pm 0.01$ | $0.10 \pm 0.01$ | $0.46 \pm 0.02$ | 2.01 | -5.75 | SE2 |
|  | 77 | O5.5 V+O5.5 V | d | 1.29 | $15.21 \pm 0.01$ | $-1.31 \pm 0.02$ | $0.04 \pm 0.02$ | $0.30 \pm 0.02$ | 2.87 | -6.15 | SE6 |
|  | 94 | O3 V | a | 1.31 | $15.57 \pm 0.02$ | $-1.31 \pm 0.03$ | $0.10 \pm 0.03$ | $0.28 \pm 0.04$ | 1.69 | -4.61 | SE6 |
|  | 115 |  |  | 1.41 | $15.76 \pm 0.02$ | $-1.32 \pm 0.03$ | $0.12 \pm 0.03$ | $0.28 \pm 0.04$ | 1.44 | -4.17 | NW8 |
|  | 132 |  |  | 1.48 | $15.86 \pm 0.03$ | $-1.35 \pm 0.03$ | $0.13 \pm 0.03$ | $0.22 \pm 0.04$ | 1.40 | -4.03 | NW8 |
|  | 36 | O2 If* | f | 1.51 | $14.41 \pm 0.01$ | $-1.26 \pm 0.01$ | $0.11 \pm 0.01$ | $0.48 \pm 0.01$ | 1.89 | -5.97 | SE3 |
|  | 173 |  |  | 1.53 | $15.98 \pm 0.03$ | $-1.32 \pm 0.04$ | $0.19 \pm 0.04$ | $0.25 \pm 0.05$ | 1.56 | -4.07 | NW8 |
|  | 75 |  |  | 1.56 | $15.08 \pm 0.01$ | $-1.32 \pm 0.02$ | $0.10 \pm 0.02$ | $0.18 \pm 0.02$ | 1.70 | -5.11 | NW8 |
|  | 114 |  |  | 1.57 | $15.68 \pm 0.02$ | $-1.35 \pm 0.03$ | $0.10 \pm 0.03$ | $0.21 \pm 0.04$ | 1.44 | -4.25 | SE7 |
|  | 108 |  |  | 1.59 | $15.44 \pm 0.02$ | $-1.35 \pm 0.02$ | $0.18 \pm 0.02$ | $0.25 \pm 0.03$ | 1.60 | -4.65 | NW8 |
|  | 31 |  |  | 1.64 | $14.12 \pm 0.01$ | $-1.31 \pm 0.01$ | $0.10 \pm 0.01$ | $0.32 \pm 0.01$ | 1.92 | -6.29 | SE8 |
|  | 49 | O3 V | a | 1.67 | $14.75 \pm 0.01$ | $-1.41 \pm 0.01$ | $0.15 \pm 0.01$ | $0.27 \pm 0.02$ | 1.60 | -5.34 | NW8 |
|  | 207 |  |  | 1.72 | $15.65 \pm 0.02$ |  | $0.09 \pm 0.02$ | $0.47 \pm 0.03$ | 1.72 | -4.56 | SE2 |
|  | 46 |  |  | 1.73 | $14.56 \pm 0.01$ | $-1.25 \pm 0.01$ | $0.21 \pm 0.01$ | $0.49 \pm 0.01$ | 2.14 | -6.07 | SE2 |
|  | 47 | O3 III(f*) | a | 1.75 | $14.72 \pm 0.01$ | $-1.22 \pm 0.01$ | $0.15 \pm 0.01$ | $0.45 \pm 0.01$ | 2.10 | -5.87 | NW3 |
|  | 40 | O3 V | a | 1.80 | $14.56 \pm 0.01$ | $-1.34 \pm 0.01$ | $0.08 \pm 0.01$ | $0.29 \pm 0.01$ | 1.83 | -5.76 | NW8 |
|  | 118 |  |  | 1.81 | 15.74 |  |  |  | 1.72 | -4.49 | SE2 |
|  | 116 |  |  | 1.81 | $15.79 \pm 0.02$ | $-1.32 \pm 0.03$ | $0.10 \pm 0.03$ |  | 1.51 | -4.21 | SE6 |
|  | 42 | $\mathrm{O} 3 \mathrm{~V}+\mathrm{O} 3 \mathrm{~V}$ | d | 1.83 | $14.72 \pm 0.01$ | $-1.36 \pm 0.01$ | $-0.03 \pm 0.01$ | $0.44 \pm 0.01$ | 1.45 | -5.22 | NW8 |
|  | 55 | O3 V | a | 1.84 | $14.74 \pm 0.01$ | $-1.31 \pm 0.01$ | $0.09 \pm 0.01$ | $0.31 \pm 0.01$ | 1.91 | -5.66 | SE9 |
|  | 39 | $\mathrm{O} 3 \mathrm{~V}+\mathrm{O} 5.5 \mathrm{~V}$ | d | 1.95 | $14.50 \pm 0.01$ | $-1.32 \pm 0.01$ | $0.08 \pm 0.01$ | $0.21 \pm 0.01$ | 1.76 | -5.75 | - |
|  | 137 |  |  | 1.97 | $15.97 \pm 0.02$ | $-1.33 \pm 0.03$ | $0.07 \pm 0.03$ | $0.23 \pm 0.04$ | 1.51 | -4.03 | , |
|  | 71 | O3-6 V | a | 2.03 | $15.16 \pm 0.01$ | $-1.35 \pm 0.01$ | $0.11 \pm 0.01$ |  | 1.67 | -5.00 | SE9 |
|  | 121 | O9.5 V | f | 2.05 | $15.85 \pm 0.02$ | $-1.24 \pm 0.03$ | $0.12 \pm 0.03$ |  | 1.58 | -4.22 | SE8 |

Notes. (a): Massey \& Hunter (1998); (b): Crowther \& Dessart (1998); (c): Walborn et al. (2002a); (d): Massey, Penny \& Vukovich (2002); (e): Crowther \& Walborn (2011); (f): Caballero-Nieves et al. (in preparation).
$\ddagger$ : Absolute magnitudes have been adjusted for R136a2, a4 and a8 (see Section 2.3).

Table 3. HST/WFC3 photometry (de Marchi et al. 2011) for additional STIS/MAMA sources ( $m_{\text {F555W }} \leq=16.0$ mag) beyond a projected distance of $2.05 \operatorname{arcsec}(0.5 \mathrm{parsec})$ from R136a1. In case WFC3 photometry is unavailable, WFPC2 photometry from Hunter et al. (1995) is used, albeit offset in F555W by -0.17 mag (indicated in italic font, see Section 2.3). Identifications are from Weigelt \& Baier (1985, WB85) or Hunter et al. (1995, HSH95), while spectral types are obtained from optical spectroscopy. $A_{\mathrm{F} 555 \mathrm{~W}}$ is either obtained via the observed $m_{\mathrm{F} 336 \mathrm{~W}}-m_{\mathrm{F} 555 \mathrm{~W}}$ colour plus $R_{5495}=4.1$, or $A_{\mathrm{F} 555 \mathrm{~W}}=1.72 \mathrm{mag}$ is adopted (values for WN5 stars are from Crowther et al. 2010).

| WB85 | HSH95 | Spectral type | Ref | $\begin{gathered} r \\ (\operatorname{arcsec}) \end{gathered}$ | $m_{\text {F555W }}$ <br> (mag) | $\begin{gathered} m_{\mathrm{F} 336 \mathrm{~W}-\mathrm{F} 438 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} m_{\text {F438W-F555W }} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{aligned} & m_{\text {F555W-F814W }} \\ & (\mathrm{mag}) \end{aligned}$ | $A_{\text {F555W }}$ (mag) | $M_{\text {F555W }}$ (mag) | Slit |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| b | 9 | O4 If/WN8 | b | 2.12 | $13.24 \pm 0.01$ | $-1.14 \pm 0.01$ | $0.18 \pm 0.01$ | $0.57 \pm 0.01$ | 2.10 | -7.31 | SE8 |
|  | 65 |  |  | 2.20 | 15.01 |  |  |  | 1.72 | -5.20 | SE7 |
|  | 134 | 07 Vz | b | 2.28 | $15.97 \pm 0.01$ | $-1.36 \pm 0.02$ | $0.10 \pm 0.02$ | $0.23 \pm 0.03$ | 1.57 | -3.97 | SE6 |
|  | 64 | O7 V((f)) | a | 2.44 | $15.03 \pm 0.01$ | $-1.19 \pm 0.01$ | $0.14 \pm 0.01$ | $0.38 \pm 0.01$ | 1.96 | -5.42 | NW7 |
|  | 45 | O3 V | a | 2.45 | $14.65 \pm 0.01$ | $-1.32 \pm 0.01$ | $0.15 \pm 0.01$ | $0.43 \pm 0.01$ | 1.77 | -5.61 | SE7 |
|  | 123 | O6 V | b | 2.47 | $15.91 \pm 0.01$ | $-1.32 \pm 0.02$ | $0.10 \pm 0.02$ | $0.09 \pm 0.03$ | 1.58 | -4.16 | SE3 |
|  | 68 |  |  | 4.61 | $15.15 \pm 0.01$ | $-1.24 \pm 0.01$ | $0.18 \pm 0.01$ | $0.46 \pm 0.01$ | 1.88 | -5.22 | SE4 |
|  | 102 |  |  | 4.83 | $15.70 \pm 0.01$ | $-1.18 \pm 0.02$ | $0.26 \pm 0.01$ | $0.66 \pm 0.01$ | 2.44 | -5.24 | NW1 |

Notes. (a): Massey \& Hunter (1998); (b): Caballero-Nieves et al. (in preparation).


Figure 4. Comparison of the background (diamonds) to total (pluses) farultraviolet counts in the central $4 \operatorname{arcsec}(\sim 1 \mathrm{pc})$ of each slit. This background represents a combination of unresolved cluster members ( 20 per cent) plus extended emission from bright stars in adjacent slits ( 15 per cent).
adopt existing (optical) subtypes for the three WN5 stars in R136a (Crowther \& Dessart 1998).

In addition to the point sources, a significant fraction of the far-UV spectrum of the central parsec of R136 arises from the intra-cluster background, on the basis of the multispec fits to the 100 brightest cluster members. This is illustrated in Fig. 4, and represents a combination of unresolved cluster members plus extended emission from bright cluster stars originating in adjacent slits. In Fig. 5, we present the far-UV spectra from the sum of selected background regions of SE3 (1 arcsec E of R136a1) and SE4 (0.7 arcsec SE of R136a1), reminiscent of R136 H70 (O5 Vz), except for extended $\mathrm{He}_{\text {II }}$ emission originating from either R136al or a2. $\sim 20$ per cent of the integrated far-UV continuum arises from the diffuse background, with bright stars in adjacent slits contributing the remaining 15 per cent of the total.

### 3.1 UV morphology of $\mathbf{O}$ stars in the LMC

Walborn, Nichols-Bohlin \& Panek (1985) have previously presented an atlas of O stars based on IUE/HIRES spectroscopy. The overwhelming majority of these stars are located in the Milky Way. More recently, Walborn et al. $(1995,2002 \mathrm{c}, 2004)$ have presented HST and FUSE UV spectroscopy of Magellanic Cloud O stars, although


Figure 5. Intra-cluster far-UV spectra of selected inner regions from SE3 and SE4 slits, in close proximity ( $0.7-1$ arcsec, solid) to R136a1 together with R136 H70 (O5 Vz, dotted) for reference. Extended He ii $\lambda 1640$ emission from R136a1 is apparent in the extracted background spectra.
subtypes and luminosity classes are poorly sampled. Since then, a number of HST programmes have obtained UV spectroscopy for a range of LMC O stars (e.g. Massey et al. 2004, 2005, 2009). In Appendix A, we present a montage of UV spectroscopic sequences for dwarfs, giants and supergiants (Figs A1-A4), with luminosity sequences shown in Figs A5-A11. A log of observations of these data sets is presented in Table A1.

There is a clear variation in UV morphology amongst $O$ stars from early to late subtypes and from dwarf to supergiant, although we wish to highlight wind line profile variation at an individual subtype/luminosity class. By way of example, we present C iv $\lambda \lambda 1548-$ 51 P Cygni line profiles amongst six O4-5.5 dwarfs in the LMC in Fig. 6. Both LH81/W28-5 and R136 H78 possess saturated P Cygni (black) absorption troughs, whereas other stars possess shallow extended absorption troughs which may be flat (e.g. $\mathrm{Sk}-70^{\circ} 60$, LH 58-496, Sk-70 ${ }^{\circ} 69$ ) or deeper at high velocities (R136 H70). Within this sample, P Cygni emission components are relatively weak, with the exception of $\mathrm{Sk}-70^{\circ} 69$ for which C IV emission is very prominent.

This variation indicates (at least) a third parameter affecting the wind profiles beyond ionization and luminosity class. While further investigation is required, it may well be age, i.e. evolutionary distance from the zero-age main sequence, since some of these stars are in very young regions. While LH 81/W28-5 does have a strong


Figure 6. Variation in spectral morphology of C iv 1548-51 amongst O45.5 dwarf stars in the LMC. The O4 dwarf LH81/W28-5 (Massey et al. 2004) is shown for reference in each case (dotted lines).

He $_{\text {II }} \lambda 4686$ absorption line (Massey et al. 2004, their fig. 17), it is substantially weaker than $\mathrm{He}_{\text {II }} \lambda 4541$, indicating a higher luminosity and/or stronger wind than for typical class V . It is well known that mass-loss rates of O stars at a given temperature and luminosity span a wide range of values (Hillier et al. 2003). Binaries and/or binary evolution may also be relevant contributors to this diversity.

### 3.2 UV morphology of stars in R136

In Appendix B we present a series of montages of UV spectroscopy of stars in R136 (Figs B1-B9), together with UV templates where appropriate.

In Fig. B1, we present HST/STIS spectroscopy of the three luminous WN stars R136a1, a2 and a3 (Crowther et al. 2010). These stars have previously been observed in the UV with HST/GHRS (de Koter et al. 1997), although a1 and a2 were blended in the GHRS 0.22 arcsec entrance aperture. Our long-slit STIS spectroscopy was obtained at a position angle of $64^{\circ}$ and extracted with multispec to disentangle a1 from a2. Nevertheless, the UV morphologies of all WN stars are extremely similar to one another, as is the case optically (HST/FOS, de Koter et al. 1997) and in the near-IR (VLT/SINFONI, Schnurr et al. 2008). We include an IUE/LORES spectrum of R146 (Brey 88, WN5) for comparison.

In Fig. B2, we present $H S T /$ STIS spectroscopy of two Of/WN stars, Melnick 37 and $\mathrm{Sk}-67^{\circ}$ 22, together with a5 and b. R136a5 has previously been observed in the UV by HST/GHRS and classified as O3f/WN (de Koter et al. 1994). From optical HST/FOS spectroscopy, R136a5 has been assigned O3 If/WN (Massey \& Hunter 1998) or O2 If* (Crowther \& Walborn 2011). The closest

LMC analogue to R136a5 in the UV is Sk-67 22 (O2 If/WN5). R136b has also previously been observed with HST/GHRS (GO 6018/Heap) while Massey \& Hunter (1998) assigned O4 If ${ }^{+}$from HST/FOS spectroscopy (see also Crowther \& Walborn 2011). From STIS/CCD spectroscopy, the ratio of $\mathrm{N}_{\text {IV }} \lambda 4058 / \mathrm{N}_{\text {III }} \lambda \lambda 4634,41$ is characteristic of O4 and WN8 subtypes (rather than O3.5 and WN7) so O4 If/WN8 is newly inferred here. Не II $\lambda 1640$ and Si IV $\lambda \lambda 1394-1402$ P Cygni profiles are much stronger in R136b than for conventional LMC O4 supergiants (e.g. $\mathrm{Sk}-65^{\circ} 47$, Fig. A8).

HST/STIS spectroscopy of O2-3 (super)giants in R136 is presented in Fig. B3, including the LMC O2-3 III and O2 If reference stars R136 H48 and H36, from our new STIS/CCD data set. Massey \& Hunter (1998) previously assigned O3 III for H48 while Walborn et al. (2002a) assigned O2-3 If for H36. The O2 supergiant classification of R136a6 is new here, with R136 H46 and H47 assigned O2 III-If, owing to their morphological similarity to the candidate runaway star VFTS 16 (Evans et al. 2010). Another example of a O2-3 III star is observed beyond the central parsec, namely R136 H102 located 4.8 arcsec to the NE.

Fig. B4 presents STIS/MAMA spectroscopy of eight O2-3 V stars in R136 together with the templates BI $237(\mathrm{O} 2 \mathrm{~V})$ and R136 H35 (O3 V) based on our new STIS/CCD data sets. All stars exhibit prominent $\mathrm{Ov} \lambda 1371$ absorption, P Cygni C iv $\lambda \lambda 1548-51$, $\mathrm{N} v \lambda \lambda 1238-42$ and negligible/weak absorption at He II $\lambda 1640$. Of these, six stars have previously been spectroscopically observed with HST/FOS by Massey \& Hunter (1998), who assigned O3 V or O3 III classifications. R136a4, a8, H31, H50 and H62 are newly classified as $\mathrm{O} 2-3 \mathrm{~V}$ in this study.
$H S T / S T I S$ spectroscopy of seven UV-classified O3-4 dwarfs in R136 is presented in Fig. B5 together with the reference stars R136 H35 (O3 V) and H78 (O4 V) from our STIS/CCD spectroscopy. All O3-4 dwarfs possess weak $\mathrm{O} v \lambda 1371$ absorption, plus prominent, albeit unsaturated C iv $\lambda \lambda 1548-51 \mathrm{P}$ Cygni profiles (Fig. A6). Of these, four have previously been spectroscopically observed by Massey \& Hunter (1998), who assigned O3 V, except for R136a7 for which O3 III was proposed from HST/FOS spectroscopy. R136 H86 and H149 are newly classified here as O3-4 V from STIS/MAMA data sets, although the latter, centrally located in the SE2 slit, suffers from significant Не ІІ $\lambda 1640$ contamination by R136a3 (SE1).

Fig. B6 presents STIS/MAMA spectroscopy for dwarf O4-5 stars in R136 together with new templates R136 H78 (O4 V) and H70 (O5 V) from new STIS/CCD observations. We favour O4 V for R136 H65, H71, H75, H89 and H90 from their morphological similarity to H 78 , while $\mathrm{O} 4-5 \mathrm{~V}$ is preferred for $\mathrm{H} 45, \mathrm{H} 49, \mathrm{H} 94$ and H 119. Classifications are in good agreement with Massey \& Hunter (1998) who proposed O 3 V for $\mathrm{H} 45, \mathrm{H} 49$ and H 94 , O5 V for H 90 and $\mathrm{O} 3-$ 6 V for H71. Fig. B7 extends this sequence to O5 V and O6 V subtypes, exploiting a STIS/CCD subtype of H123 (O6 V). We propose O 5 V subtypes for H 64 and H 69 , 06 V subtypes for H 30 , H77 and H92, and O5-6 V for H141, again in reasonable agreement with Massey \& Hunter (1998) who adopted O7 V for H30 and H64, O3-6 V for H69, O3 V for H92. Massey et al. (2002) obtained O5.5 V+O5.5 V for H77. The UV morphology of R136 H68, located 4.6 arcsec from R136a1, is also consistent with O6 V.

Fig. B8 presents STIS spectroscopy for O7-8 dwarfs in R136, exploiting new optical subtypes of O7 V and O8 V for R136 H134 and H 80 , respectively. All stars are newly classified except for H112, for which Massey \& Hunter (1998) proposed O8.5 III. A very weak wind signature of C Iv $\lambda \lambda 1548-51$ is observed, via either very shallow P Cygni absorption or a clear blueshifted absorption. We assign $\mathrm{O} 9+\mathrm{V}$ subtypes for the final group of stars, presented in Fig. B9, for which solely photospheric C iv $\lambda \lambda 1548-51$ absorption is

Table 4. STIS/MAMA-derived properties of R136 stars within a projected distance of $2.05 \operatorname{arcsec}(0.5 \mathrm{pc})$, brighter than $\left.m_{\mathrm{F} 555 \mathrm{~W}}=16.0 \mathrm{mag}\right)$, sorted by projected distance from R136a1. References to optical (literature) spectral types are provided in Table 2. $F_{1500}$ fluxes are averages for the range $1500 \pm$ $25 \AA$ (estimated values are provided in parentheses). Aside from the WN5 stars (Crowther et al. 2010), temperatures and luminosities are from spectral type calibrations (Doran et al. 2013), while inferred current masses $M$ and ages $\tau$ for O stars are from bonnsai (Schneider et al. 2014a) using the Bonn evolutionary models (Brott et al. 2011; Köhler et al. 2015). Velocities include $v_{\text {edge }}, v_{\text {black }}$ and $v_{\infty}$ (colons indicate less secure values).

| HSH95 <br> (WB85) | Optical Sp type | Ultraviolet <br> Sp type | $r(\operatorname{arcsec})$ | $\begin{array}{r} 10^{14} F_{1500} \\ \left(\mathrm{erg} \mathrm{~s}^{-1} \mathrm{~cm}^{-2} \AA\right) \end{array}$ | $M_{\text {F555W }}$ <br> (mag) | $\begin{gathered} T_{\text {eff }} \\ (\mathrm{kK}) \end{gathered}$ | $\log L / L_{\odot}$ | $\begin{gathered} M_{\text {current }} \\ \left(\mathrm{M}_{\odot}\right) \end{gathered}$ | $\begin{aligned} & \tau \\ & (\mathrm{Myr}) \end{aligned}$ | $\begin{array}{r} v_{\text {edge }} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{array}$ | $\begin{gathered} v_{\text {black }} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} v_{\infty} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | Fig. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 3 (a1) | WN5h | WN5 | 0.00 | $22.9 \pm 1.4$ | -8.09 | $53 \pm 3$ | $6.94 \pm 0.09$ | $315_{-50}^{+60}$ | $0.0_{-0.0}^{+0.3}$ | 3270 : | 2600 | 2600 | B1 |
| 5 (a2) | WN5h | WN5 | 0.08 | $13.1 \pm 1.2$ | -7.52 | $53 \pm 3$ | $6.63 \pm 0.09$ | $195_{-30}^{+35}$ | $0.3_{-0.3}^{+0.4}$ | 3440 | 2425 | 2425 | B1 |
| 20 (a5) | O2 If* | O2 If/WN5 | 0.28 | $6.9 \pm 0.3$ | -6.65 | $50_{-5}^{+4}$ | $6.32_{-0.15}^{+0.16}$ | $101_{-26}^{+28}$ | $0.8_{-0.7}^{+0.4}$ | 3615 | 3045 | 3045 | B2 |
| 24 (a7) | O3 III(f*) | O3-4 V | 0.36 | $4.4 \pm 0.4$ | -6.10 | $46_{-3}^{+4}$ | $5.99_{-0.08}^{+0.10}$ | $69_{-10}^{+12}$ | $0.5_{-0.3}^{+0.1}$ | 3385 | - | 2710: | B5 |
| 27 (a8) |  | O2-3 V | 0.39 | $5.2 \pm 0.3$ | -6.05 | $51 \pm 6$ | $6.28{ }_{-0.15}^{+0.13}$ | $96_{-22}^{+27}$ | $0.8_{-0.8}^{+0.5}$ | 3550 | 2980 | 2980 | B4 |
| 21 (a4) |  | O2-3 V | 0.43 | $7.8 \pm 0.3$ | -6.51 | $51 \pm 6$ | $6.46_{-0.15}^{+0.13}$ | $124_{-31}^{+37}$ | $0.3_{-0.3}^{+0.2}$ | 3275 | 2475: | 2475: | B4 |
| 186 |  | - | 0.43 | (1.3) | -4.62 | - | - | - | - | - | - | - | - |
| 86 |  | O3-4 V | 0.45 | $2.1 \pm 0.2$ | -5.91 | $46_{-3}^{+4}$ | $5.91{ }_{-0.08}^{+0.10}$ | $63_{-8}^{+10}$ | $1.5{ }_{-0.9}^{+0.5}$ | 3275: | 2475 | 2475 | B5 |
| 6 (a3) | WN5h | WN5 | 0.49 | $16.3 \pm 2.1$ | -7.39 | $53 \pm 3$ | $6.58 \pm 0.09$ | $180_{-30}^{+30}$ | $0.3_{-0.3}^{+0.4}$ | 2950 | 2400 | 2400 | B1 |
| 66 | O3 V | O3-4 V | 0.49 | $2.2 \pm 0.2$ | $-5.38$ | $46_{-3}^{+4}$ | $5.70_{-0.08}^{+0.10}$ | $51 \pm 7$ | $1.4{ }_{-1.1}^{+0.6}$ | 3045 | 2590 | 2590 | B5 |
| 119 |  | O4-5 V | 0.51 | $2.7 \pm 0.2$ | -4.46 | $43 \pm 3$ | $5.24{ }_{-0.09}^{+0.08}$ | $32_{-3}^{+4}$ | $0.8_{-0.8}^{+1.4}$ | 3075 | - | 2460: | B6 |
| 58 | O3 III(f*) | O2-3 V | 0.59 | $2.7 \pm 0.3$ | -5.65 | $51 \pm 6$ | $5.93{ }_{-0.15}^{+0.13}$ | $63_{-12}^{+16}$ | $0.6_{-0.6}^{+0.9}$ | - | 2980 | 2980 | B4 |
| 30 | 07 V | 06 V | 0.62 | $5.7 \pm 0.3$ | -5.89 | $38 \pm 2$ | $5.67{ }_{-0.06}^{+0.05}$ | $42 \pm 3$ | $2.9{ }_{-0.3}^{+0.3}$ | 3110: | - | 2490: | B7 |
| 70 | O 5 Vz | O5 V | 0.62 | $2.6 \pm 0.2$ | -5.32 | $42 \pm 2$ | $5.57{ }_{-0.08}^{+0.07}$ | $41_{-4}^{+5}$ | $2.3{ }_{-0.7}^{+0.5}$ | 3340: | - | 2670: | B6 |
| 89 |  | O4 V | 0.66 | $1.3 \pm 0.1$ | -6.25 | $44 \pm 2.5$ | $5.99_{-0.07}^{+0.08}$ | $65_{-9}^{+10}$ | $1.7 \pm 0.4$ | 2425: | - | 1940: | B6 |
| 62 |  | O2-3 V | 0.66 | $3.3 \pm 0.2$ | -5.43 | $51 \pm 6$ | $5.84{ }_{-0.15}^{+0.13}$ | $58_{-12}^{+13}$ | $0.4_{-0.4}^{+1.2}$ | 3340 | 2770 | 2770 | B4 |
| 19 (a6) |  | O2 If | 0.73 | $11.6 \pm 0.4$ | -6.96 | $46_{-2}^{+4}$ | $6.52_{-0.05}^{+0.10}$ | $150_{-40}^{+36}$ | $0.9_{-0.6}^{+0.5}$ | 3440 | 2650 | 2650 | B3 |
| 50 |  | O2-3 V | 0.74 | $3.4 \pm 0.2$ | -5.89 | $51 \pm 6$ | $6.02_{-0.13}^{+0.15}$ | $70_{-14}^{+17}$ | $0.8_{-0.8}^{+0.7}$ | 3390 | 2620 | 2620 | B4 |
| 90 | O5 V | O 4 V | 0.77 | $1.7 \pm 0.1$ | -4.68 | $44 \pm 2.5$ | $5.36{ }_{-0.07}^{+0.08}$ | $36 \pm 4$ | $1.0_{-1.0}^{+0.8}$ | - | 2475: | 2475: | B6 |
| 149 |  | O3-4 V:: | 0.78 | $1.0 \pm 0.1$ | -5.01 | - | - | - | - | 2900 | 2210 | 2210 | B5 |
| 141 |  | O5-6 V | 0.78 | $1.0 \pm 0.1$ | -4.22 | $41 \pm 3$ | $5.09 \pm 0.12$ | $27_{-3}^{+4}$ | $1.0_{-1.0}^{+1.5}$ | - | - | - | B7 |
| 80 | O8 V | O8 V | 0.87 | $2.3 \pm 0.1$ | -4.85 | $36 \pm 2$ | $5.20_{-0.10}^{+0.08}$ | $26_{-2}^{+3}$ | $4.0_{-0.9}^{+0.8}$ | 2070: | - | 1655: | B8 |
| 35 | O3 V | O3 V | 0.88 | $5.5 \pm 0.2$ | -5.81 | $48 \pm 3$ | $5.92{ }_{-0.09}^{+0.08}$ | $66_{-9}^{+10}$ | $1.2_{-0.8}^{+0.5}$ | 3440 | 2770 | 2770 | B4 |
| 78 | O4: V | O4: V | 0.97 | $1.8 \pm 0.2$ | -4.93 | $44 \pm 2.5$ | $5.47{ }_{-0.07}^{+0.08}$ | $40_{-4}^{+5}$ | $1.5{ }_{-1.1}^{+0.7}$ | 3000 | 2375: | 2375: | B6 |
| 73 | O9 V | O9+ V | 1.01 | $2.0 \pm 0.1$ | -5.14 | $33 \pm 2$ | $5.21_{-0.10}^{+0.09}$ | $24 \pm 2$ | $4.8 \pm 0.7$ | - | - | - | B9 |
| 92 | O3 V | 06 V | 1.05 | $2.2 \pm 0.2$ | -4.55 | $40 \pm 2$ | $5.20_{-0.08}^{+0.07}$ | $30 \pm 3$ | $2.3_{-1.4}^{+0.9}$ | - | 2080: | 2080: | B7 |
| 143 |  | O7-8 V | 1.05 | $0.8 \pm 0.1$ | -4.25 | $39 \pm 3$ | $4.99_{-0.14}^{+0.12}$ | $24 \pm 3$ | $1.6_{-1.6}^{+1.4}$ | 1850: | - | 1480: | B8 |
| 112 | O8.5 III(f) | O7+V | 1.06 | $1.0 \pm 0.1$ | -4.40 | $36 \pm 4$ | $5.01_{-0.11}^{+0.10}$ | $23 \pm 3$ | $3.8{ }_{-2.5}^{+1.6}$ | - | - | - | B8 |
| 135 |  | O9+ V | 1.08 | $1.0 \pm 0.1$ | -4.26 | $33 \pm 2$ | $4.86_{-0.10}^{+0.09}$ | $20 \pm 2$ | $5.3_{-1.7}^{+1.3}$ | - | - | - | B9 |
| 69 | O3-6 V | O5 V | 1.09 | $2.2 \pm 0.2$ | -5.16 | $42 \pm 2$ | $5.49_{-0.08}^{+0.07}$ | $38 \pm 4$ | $2.2_{-0.9}^{+0.6}$ | 3225 | - | 2580 | B7 |
| 52 | 03 V | O3-4 V | 1.13 | $4.0 \pm 0.3$ | -5.48 | $46_{-3}^{+4}$ | $5.74{ }_{-0.08}^{+0.10}$ | $53_{-7}^{+8}$ | $0.3 \pm 0.2$ | 3260 | 2820 | 2820 | B5 |
| 48 | O2-3 III( $\mathrm{f}^{*}$ ) | O2-3 III | 1.22 | $2.7 \pm 0.2$ | -5.75 | $51 \pm 6$ | $5.97{ }_{-0.15}^{+0.13}$ | $66_{-12}^{+16}$ | $0.7_{-0.7}^{+0.8}$ | 3615 | 3045 | 3045 | B3 |
| 77 | O5.5 V+O5.5 V | O6 V | 1.29 | $2.4 \pm 0.2$ | -6.15 | - |  | - | - | 1920: | 1510: | 1510: | B7 |
| 94 | O3 V | O4-5 V | 1.31 | $1.5 \pm 0.1$ | -4.61 | $43 \pm 3$ | $5.31{ }_{-0.09}^{+0.08}$ | $34 \pm 4$ | $1.3{ }_{-1.2}^{+0.9}$ | 3110: | - | 2490: | B6 |
| 115 |  | O9+V | 1.41 | $1.4 \pm 0.1$ | -4.17 | $33 \pm 2$ | $4.82_{-0.10}^{+0.09}$ | $19_{-1}^{+2}$ | $5.3{ }_{-1.9}^{+1.4}$ | - | - | - | B9 |
| 132 |  | O9+V | 1.48 | $1.1 \pm 0.1$ | -4.03 | $33 \pm 2$ | $4.76{ }_{-0.10}^{+0.09}$ | $18_{-1}^{+2}$ | $6.2_{-0.5}^{+1.6}$ | - | - | - | B9 |
| 36 | O2 If* | O2 If* | 1.51 | $3.6 \pm 0.2$ | -5.97 | $46_{-2}^{+4}$ | $5.94{ }_{-0.05}^{+0.10}$ | $66_{-9}^{+10}$ | $1.5_{-0.9}^{+0.4}$ | 4625: | 3500 | 3500 | B3 |
| 173 |  | O9+ V | 1.53 | $0.7 \pm 0.1$ | -4.07 | $33 \pm 2$ | $4.78{ }_{-0.10}^{+0.09}$ | $19 \pm 1$ | $6.6{ }_{-0.8}^{+1.4}$ | - | - | - | B9 |
| 75 |  | O4 V | 1.56 | $2.2 \pm 0.2$ | -5.11 | $44 \pm 2.5$ | $5.54{ }_{-0.07}^{+0.08}$ | $42 \pm 5$ | $1.7_{-1.1}^{+0.7}$ | 3000 | 2550 | 2550 | B6 |
| 114 |  | O7-8 V | 1.57 | $1.6 \pm 0.1$ | -4.25 | $37 \pm 3$ | $4.99_{-0.14}^{+0.12}$ | $23 \pm 3$ | $3.1_{-2.2}^{+1.4}$ | 2215: | - | 1770: | B8 |
| 108 |  | O7-8 V | 1.59 | $1.5 \pm 0.1$ | -4.65 | $37 \pm 3$ | $5.15{ }_{-0.14}^{+0.12}$ | $25_{-3}^{+4}$ | $3.6{ }_{-1.9}^{+1.2}$ | 1300: | - | 1040: | B8 |
| 31 |  | O2-3 V | 1.64 | $6.0 \pm 0.3$ | -6.29 | $51 \pm 6$ | $6.19_{-0.15}^{+0.13}$ | $87_{-20}^{+21}$ | $0.8_{-0.8}^{+0.5}$ | 3500 | 2815: | 2815: | B4 |
| 49 | O3 V | O4-5 V | 1.67 | $0.9 \pm 0.1$ | -5.34 | $43 \pm 3$ | $5.60_{-0.09}^{+0.08}$ | $43_{-5}^{+6}$ | $2.0_{-1.0}^{+0.7}$ | 3780 | 2980: | 2980: | B6 |
| 207 |  | - | 1.72 | $0.3 \pm 0.1$ | -4.56 | - | - | - | - | 4560: | - | 3650: | - |
| 46 |  | O2 III-If | 1.73 | $2.3 \pm 0.2$ | -6.07 | $48 \pm 4$ | $6.02{ }_{-0.09}^{+0.12}$ | $74_{-11}^{+14}$ | $1.2{ }_{-0.9}^{+0.4}$ | 4450: | 3440 | 3440 | B3 |
| 47 | O3 III(f*) | O2 III-If | 1.75 | $2.6 \pm 0.3$ | -5.87 | $48 \pm 4$ | $5.95{ }_{-0.13}^{+0.14}$ | $65_{-11}^{+14}$ | $1.1{ }_{-0.9}^{+0.5}$ | 3615 | 3045 | 3045 | B3 |
| 40 | O3 V | O2-3 V | 1.80 | $4.5 \pm 0.2$ | -5.76 | $51 \pm 6$ | $5.97{ }_{-0.15}^{+0.13}$ | $66_{-12}^{+16}$ | $0.7_{-0.7}^{+0.8}$ | 3435 | 2750 | 2750 | B4 |

Table 4 - continued

| HSH95 <br> (WB85) | Optical Sp type | Ultraviolet Sp type | $r(\operatorname{arcsec})$ | $\begin{array}{r} 10^{14} F_{1500} \\ \left(\mathrm{erg} \mathrm{~s}^{-1} \mathrm{~cm}^{-2} \AA\right) \end{array}$ | $\begin{gathered} M_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} T_{\text {eff }} \\ (\mathrm{kK}) \end{gathered}$ | $\log L / L_{\odot}$ | $\begin{aligned} & M_{\text {current }} \\ & \left(\mathrm{M}_{\odot}\right) \end{aligned}$ | $\tau$ <br> (Myr) | $\begin{array}{r} v_{\text {edge }} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{array}$ | $\begin{gathered} v_{\text {black }} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} v_{\infty} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | Fig. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 118 |  | O7-8 V | 1.81 | $1.1 \pm 0.2$ | -4.49 | $39 \pm 3$ | $5.07{ }_{-0.14}^{+0.12}$ | $25_{-3}^{+4}$ | $2.2_{-1.8}^{+1.2}$ | 1460: | - | 1170: | B8 |
| 116 |  | O7-8 V | 1.81 | $1.3 \pm 0.1$ | -4.21 | $37 \pm 3$ | $4.97{ }_{-0.14}^{+0.12}$ | $23 \pm 3$ | $3.0_{-2.2}^{+1.4}$ | 1200: | - | 960: | B8 |
| 42 | O3 V+O3 V | O3-4 V | 1.83 | $3.5 \pm 0.2$ | -5.22 | $46_{-3}^{+4}$ | $5.64{ }_{-0.08}^{+0.10}$ | $48_{-6}^{+7}$ | $0.3_{-0.3}^{+0.2}$ | 3045 | 2245 | 2245 | B5 |
| 55 | O3 V | O2-3 V | 1.84 | $3.6 \pm 0.2$ | -5.66 | $51 \pm 6$ | $5.93_{-0.15}^{+0.13}$ | $63_{-12}^{+15}$ | $0.6_{-0.6}^{+0.9}$ | 3390 | 2880 | 2880 | B4 |
| 39 | $\mathrm{O} 3 \mathrm{~V}+\mathrm{O} 5.5 \mathrm{~V}$ | - | 1.95 | (3.7) | -5.75 | - | , | - | , | - | - | - | - |
| 137 |  | - | 1.97 | (1.0) | -4.03 | - | . | - | - | - | - | - | - |
| 71 | O3-6 V | O4 V | 2.03 | $2.4 \pm 0.2$ | -5.00 | $44 \pm 2.5$ | $5.49_{-0.07}^{+0.08}$ | $40_{-4}^{+5}$ | $1.6{ }_{-1.1}^{+0.7}$ | 3030 | 2475 | 2475 | B6 |
| 121 | 09.5 V | O9.5 V | 2.05 | $1.2 \pm 0.1$ | -4.22 | $33 \pm 1.5$ | $4.84{ }^{ \pm} 0.07$ | $19_{-1}^{+2}$ | $5.3_{-1.3}^{+1.1}$ | - | - | - | B9 |

Table 5. STIS/MAMA-derived properties of R136 stars beyond a projected distance of $2.05 \operatorname{arcsec}(0.5 \mathrm{pc})$ from R136a1, brighter than $\left.m_{\mathrm{F} 555 \mathrm{~W}}=16.0 \mathrm{mag}\right)$. References to optical (literature) spectral types are provided in Table 3. $F_{1500}$ fluxes are averages for the range $1500 \pm 25 \AA$. Aside from the WN5 stars (Crowther et al. 2010), temperatures and luminosities are from spectral type calibrations (Doran et al. 2013), while inferred current masses $M$ and ages $\tau$ for O stars are from bonnsai (Schneider et al. 2014a) using the Bonn evolutionary models (Brott et al. 2011; Köhler et al. 2015). Velocities include $v_{\text {edge }}$, $v_{\text {black }}$ and $v_{\infty}$ (colons indicate less secure values).

| HSH95 <br> (WB85) | Optical Sp type | Ultraviolet Sp type | $r(\operatorname{arcsec})$ | $\begin{array}{r} 10^{14} F_{1500} \\ \left(\mathrm{erg} \mathrm{~s}^{-1} \mathrm{~cm}^{-2} \AA\right) \end{array}$ | $\begin{gathered} M_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} T_{\text {eff }} \\ (\mathrm{kK}) \end{gathered}$ | $\begin{aligned} & \log L / L \odot \\ & \left(\mathrm{~L}_{\odot}\right) \end{aligned}$ | $\begin{aligned} & M \\ & \left(\mathrm{M}_{\odot}\right) \end{aligned}$ | $\begin{aligned} & \tau \\ & (\mathrm{Myr} \end{aligned}$ | $\begin{array}{r} v_{\text {edge }} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{array}$ | $\begin{gathered} v_{\text {black }} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} v_{\infty} \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | Fig. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 9 (b) | O4 If/WN8 | O4 If/WN8 | 2.12 | $7.5 \pm 0.6$ | -7.31 | $41 \pm 3$ | $6.30_{-0.12}^{+0.11}$ | $93 \pm 19$ | $1.5 \pm 0.3$ | 2080 | 1400 | 1400 | B2 |
| 65 |  | O4 V | 2.20 | $1.9 \pm 0.2$ | -5.20 | $44 \pm 2.5$ | $5.56{ }_{-0.07}^{+0.08}$ | $43 \pm 5$ | $1.7_{-1.0}^{+0.6}$ | 3110 | 2540 | 2540 | B6 |
| 134 | 07 Vz | O7 V | 2.28 | $1.2 \pm 0.1$ | -3.97 | $38 \pm 2$ | $4.91{ }_{-0.06}^{+0.05}$ | $23 \pm 2$ | $1.8_{-1.6}^{+1.3}$ | 1465 | - | 1170: | B8 |
| 64 | $07 \mathrm{~V}($ (f) ) | O5 V | 2.44 | $1.7 \pm 0.1$ | -5.42 | $42 \pm 2$ | $5.60_{-0.08}^{+0.07}$ | $42_{-4}^{+5}$ | $2.33_{-0.6}^{+0.5}$ | 2215: | - | 1770: | B7 |
| 45 | O3 V | O4-5 V | 2.45 | $3.0 \pm 0.2$ | -5.61 | $43 \pm 3$ | $5.76{ }_{-0.09}^{+0.08}$ | $50_{-6}^{+7}$ | $2.0_{-0.7}^{+0.6}$ | 3275: | - | 2620: | B6 |
| 123 | O6 V | O6 V | 2.47 | $1.1 \pm 0.1$ | -4.16 | $40 \pm 2$ | $5.04{ }_{-0.08}^{+0.07}$ | $27 \pm 2$ | $1.22_{-1.2}^{+1.1}$ | 2020 | - | 1615: | B7 |
| 68 |  | O6 V | 4.61 | $1.6 \pm 0.1$ | -5.22 | - | - | - | - | 2390 | - | 1910: | B7 |
| 102 |  | O2-3 III | 4.83 | $0.6 \pm 0.1$ | -5.24 | - | - | - | - | 2865: | 2540 | 2540 | B3 |

observed, in common with H 121 for which O 9.5 V is obtained from our STIS/CCD spectroscopy. Low S/N becomes increasingly problematic for these stars, whose far-UV fluxes are typically $F_{1500} \sim$ $10^{-14} \mathrm{erg} \mathrm{s}^{-1} \mathrm{~cm}^{-2} \AA^{-1}$ or lower.

Of the remaining stars within 2.05 arcsec of R136a1 that are listed in Table 4 two lie beyond our HST/STIS slit spectroscopy, namely R136 H39 and H137, of which H39 has been observed previously by Massey et al. (2002) and classified O3 V+O5.5 V. Properties of 8 bright sources beyond the central region of R136a are included in Table 5. Finally, H186 lies close to a much brighter source (H86), within its slit, so we are unable to isolate its UV spectrum.

### 3.3 Wind velocities for R136 stars

The presence of strong, saturated P Cygni profiles amongst the O stars of R136a allows us to measure their wind velocities. To date, wind velocities for large numbers of Galactic OB stars have been obtained from IUE SWP/HIRES spectroscopy (Prinja et al. 1990; Howarth et al. 1997), whereas wind velocities for O stars in the Magellanic Clouds have been relatively scarce (e.g. Prinja \& Crowther 1998; Massey et al. 2004, 2005).

Here we measure the maximum blueward extend of saturated P Cygni absorption lines (usually $\mathrm{C}_{\text {IV }} \lambda \lambda 1548-51$ ), $v_{\text {black }}$, and use it as a probe of terminal wind velocity, i.e. $v_{\infty}=v_{\text {black }}$ (Prinja et al. 1990). The velocity at which the violet absorption meets the continuum, the edge velocity $v_{\text {edge }}$, is usually larger than $v_{\text {black }}$ as a result of turbulent motions in the outflow. Such stochastic motions cause some of the material to move at speeds beyond the local flow
speed, hence, will cause absorption beyond the terminal speed if the line is fully saturated. If P Cygni profiles are not fully saturated, $v_{\infty}$ can be approximated as a function of $v_{\text {edge }}$. In such cases, the terminal velocity could be underestimated because the ion may no longer be present in the outer wind.

We measure $v_{\text {black }}$ and/or $v_{\text {edge }}$ with respect to (LMC) interstellar absorption lines C iv $\lambda \lambda 1548-51$ and/or Si II $\lambda 1527$ (e.g. Howarth \& Phillips 1986). The systemic radial velocity of R136 is $\sim 268 \mathrm{~km} \mathrm{~s}^{-1}$ (Hénault-Brunet et al. 2012). Individual measurements should be reliable to $\pm 100 \mathrm{~km} \mathrm{~s}^{-1}$, so $v_{\infty}$ should be accurate to $\pm 200 \mathrm{~km} \mathrm{~s}^{-1}$.

For this sample, Civ $\lambda \lambda 1548-51$ exhibits one of three spectral morphologies. A significant subset, typically $\mathrm{O} 2-4$ stars and supergiants, exhibit saturated P Cygni profiles, from which $v_{\text {black }}$ (and $v_{\text {edge }}$ ) can be measured, as shown for R136 H36 in the top panel of Fig. 7. A further subset possess strong P Cygni profiles, albeit without saturated absorption features, from which $v_{\text {edge }}$ can be measured, typical of O5-6 dwarfs, as shown for R136 H70 in the central panel of Fig. 7. O7-8 dwarfs tend not to exhibit prominent P Cygni profiles, but some show blue absorption that are sufficiently distinct from interstellar C iv, as shown in the lower panel of Fig. 7 for H80.

Individual results are presented in Tables 4 and 5. The highest terminal wind velocities are $v_{\infty} \sim 3500 \mathrm{~km} \mathrm{~s}^{-1}$ for the O 2 (super)giants H 36 and H 46 , while late O dwarfs often possess $v_{\text {edge }} \leq$ $1500 \mathrm{~km} \mathrm{~s}^{-1}$, or no wind measurements are possible. Fig. 8 shows edge and black velocities for 28 R136 stars, from which we obtain $v_{\text {black }} / v_{\text {edge }}=0.8 \pm 0.05$, which is adopted to estimate wind velocities for stars in which solely edge velocities are measured. For reference, Prinja \& Crowther (1998) obtained $v_{\text {black }} / v_{\text {edge }} \sim 0.7$ for


Figure 7. Representative HST/STIS C iv $\lambda \lambda 1548-51$ profiles for stars with saturated P Cygni profiles (top), unsaturated P Cygni profiles (centre) and non-P Cygni profiles (bottom). Profiles are shown in velocity space in the LMC rest frame of interstellar C iv $1548.2 \AA$ in all cases.

LMC stars and $v_{\text {black }} / v_{\text {edge }} \sim 0.9$ for 6 R136 stars (in common with this sample).

In Table 6, average values of wind velocities in R136 are provided for O2-3, O3-4, O4-5, O5-6, O7-8 dwarfs, plus early O giants and supergiants, including cases for which wind velocities are based upon $v_{\infty}=(0.8 \pm 0.05) v_{\text {edge }}$. Amongst dwarfs, the trend towards lower wind velocities for later subtypes is apparent. The standard deviation at O5-6 is unusually large since individual estimates span a wide range from $1510 \mathrm{~km} \mathrm{~s}^{-1}$ (H77, O6 V from UV morphology) to $2580 \mathrm{~km} \mathrm{~s}^{-1}$ (H69, O5 V from UV morphology). Statistics of wind velocities for 'field' O stars elsewhere in the LMC are included, drawn from Walborn et al. (1995), de Koter et al. (1998), Prinja \& Crowther (1998), Massa et al. (2003), Massey et al. (2004, 2005, 2009), Evans et al. (2010) and Bestenlehner et al. (2014), with individual measurements from the literature and comparisons with present results provided in Table C1 in Appendix C.

The wind velocity of $1400 \mathrm{~km} \mathrm{~s}^{-1}$ for R136b (O4 If/WN8) is extremely low, even with respect to O2-3.5 If/WN stars in the LMC (Table C1). In contrast, O2-3 supergiants within R136 possess higher wind velocities than dwarfs, which is the reverse of O2-3 stars elsewhere in the LMC, albeit based on low-number statistics. For completeness, we obtain an average wind terminal velocity of


Figure 8. Comparison between edge and black velocities (in $\mathrm{Mm} \mathrm{s}^{-1}$ ) for R136 WN (stars), Of/WN (squares), (super)giants (triangles) and dwarfs (circles), together with the best-fitting $v_{\text {black }}=(0.8 \pm 0.05) v_{\text {edge }}$ relationship for those stars without saturated black absorption troughs (solid line).

Table 6. Terminal wind velocities $\left(\mathrm{km} \mathrm{s}^{-1}\right)$ and standard deviations ( $\sigma$ ) for O stars in R136 (this study) and other O stars in the LMC (Walborn et al. (1995); Massa et al. (2003); Massey et al. (2004, 2005, 2009), Evans et al. (2010); Bestenlehner et al. (2014). N is the sample size for each category.

| Subtype | Dwarf |  |  | Giants |  |  | Supergiants |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $v_{\infty}$ | $\sigma$ | N | $v_{\infty}$ | $\sigma$ | N | $v_{\infty}$ | $\sigma$ | N |
| - R136 (HST/STIS) - |  |  |  |  |  |  |  |  |  |
| O2-3 | 2780 | 160 | 9 | 3020 | 370 | 4 | 3065 | 425 | 3 |
| O3-4 | 2525 | 235 | 8 | - | - | - | 1400 | - | 1 |
| O4-5 | 2475 | 245 | 11 | - | - | - | - | - | - |
| O5-6 | 2095 | 415 | 8 | - | - | - | - | - | - |
| O7-8 | 1320 | 315 | 7 | - | - | - | - | - | - |
| - LMC field (HST FOS/STIS/COS; FUSE) - |  |  |  |  |  |  |  |  |  |
| O2 | 3290 | 155 | 2 | 3310 | 180 | 7 | 2600 | 345 | 4 |
| O3-3.5 | 2950 | 375 | 4 | 2880 | 450 | 2 | 2000 | - | 1 |
| O4-6 | 2510 | 205 | 5 | 2360 | 245 | 3 | 1900 | 130 | 5 |
| O6.5-8 | 1950 | - | 1 | 2675 | 1025 | 4 | 1950 | 70 | 2 |
| 08.5-9.7 | - | - | - | - | - | - | 1375 | 375 | 6 |

$2475 \pm 110 \mathrm{~km} \mathrm{~s}^{-1}$ for the three WN5 stars within R136a. Crowther et al. (2010) obtained $v_{\infty}=2200-2600 \mathrm{~km} \mathrm{~s}^{-1}$ for these stars from their UV to near-IR spectroscopic analysis.

Our new data set effectively double the sample of LMC early O stars with measured wind velocities. It is important to note that a range of wind velocities have been obtained for individual stars. In part this arises from the methodology, either direct measurement as applied here or line fitting based on the Sobolev with Exact Integration method (SEI; Groenewegen \& Lamers 1989). Still, different authors applying the same approach to the same data set have obtained significantly different wind velocities. By way of example, de Koter et al. (1998) obtained $v_{\infty}=3000 \mathrm{~km} \mathrm{~s}^{-1}$ from HST GHRS spectroscopy of R136a5, while Massey et al. (2004) obtained $3400 \mathrm{~km} \mathrm{~s}^{-1}$ from the same data set; we find $3045 \pm 200 \mathrm{~km} \mathrm{~s}^{-1}$


Figure 9. Wind velocities (in $\mathrm{km} \mathrm{s}^{-1}$ ) versus estimated stellar temperature [Of/WN stars: squares, (super)giants: triangles, dwarfs: circles] based on $v_{\text {black }}$ (filled) or $v_{\text {edge }}$ (open) assuming the temperature scale of Doran et al. (2013). Average velocities versus stellar temperature for Galactic O stars from Prinja et al. (1990) are indicated as dotted lines (supergiants), dashed lined (giants) and solid lines (dwarfs), incorporating the updated classification scheme of Walborn et al. (2002a) and temperature scale of Martins et al. (2005) plus $T_{\text {eff }}(\mathrm{O} 2 \mathrm{I})=46 \mathrm{kK}$.
from STIS spectroscopy. Mindful of concerns about systematic differences between alternative approaches, an SEI analysis of H36 was carried out (Prinja, private communication) revealing $v_{\infty}=$ $3500 \mathrm{~km} \mathrm{~s}^{-1}, v_{\text {turb }} / v_{\infty}=0.15$ for a $\beta=1$ velocity law, in agreement with the $v_{\text {black }}$ result. This does not imply that all such measurements will be entirely consistent, although it does suggest that there is no inherent systematic offset.

Despite the low-number statistics, wind velocities of O2-4 dwarfs in R136 are somewhat lower than elsewhere in the LMC. By way of example, the average wind velocity of 8 R136 O3-4 dwarfs is $2525 \mathrm{~km} \mathrm{~s}^{-1}$, versus $2900 \mathrm{~km} \mathrm{~s}^{-1}$ for five counterparts elsewhere in the LMC (and $3040 \mathrm{~km} \mathrm{~s}^{-1}$ for Galactic O3-4 dwarfs). The reason for this is not clear, if not simply due to small sample sizes. If physical, it is likely not the result of a lower metallicity in 30 Doradus ( $v_{\infty} \propto Z^{0.13}$; Leitherer, Robert \& Drissen 1992) relative to the rest of the LMC, although youth may be a factor, recalling the discussion linked to Fig. 6.

Wind velocities for individual R136 stars are presented as a function of stellar temperature in Fig. 9, based on UV-derived subtypes and the temperature calibration of LMC O stars of Doran et al. (2013), in which results from $v_{\text {black }}$ (filled symbols) and $v_{\text {edge }}$ (open symbols) are distinguished. We have included average UV-derived wind velocities for Galactic O stars from Prinja et al. (1990) based on the same methodology as used here - updated for O2-4 stars following Walborn et al. (2002a), and adopting the Martins, Schaerer \& Hillier (2005) Galactic observational temperature calibration adapted for LMC O stars by Doran et al. (2013).

To date, comparisons of empirical (reduced) wind momenta of O stars with theoretical predictions either adopt wind velocities or are based on calibrations. Pre-empting our stellar mass estimates from Section 3.4, we are able to compare terminal wind velocities


Figure 10. Hertzsprung-Russell diagram [WN: stars; Of/WN stars: squares; (super)giants: triangles; dwarfs: circles] brighter than $m_{\mathrm{F} 555 \mathrm{~W}}=$ 16.0 mag within 0.6 pc from R136a1, based on either optically-derived (filled) or UV-derived (open) spectral types using the temperature calibrations from Doran et al. (2013). LMC isochrones for slowly rotating stars $\left(v_{\text {eq }}=100 \mathrm{~km} \mathrm{~s}^{-1}\right)$ at ages of $0,1,2,3$ and 4 Myr are overlaid (solid), together with individual tracks for $25,40,60,100$ and $200 \mathrm{M}_{\odot}$ stars (dashed), from Brott et al. (2011) and Köhler et al. (2015). Representative uncertainties are indicated in the lower left. A dotted line marking $m_{\text {F555W }}=16.0 \mathrm{mag}$ for an adopted $A_{\mathrm{F} 555 \mathrm{~W}}=1.72 \mathrm{mag}$ is indicated, emphasizing that our census is sensitive to all luminous stars in R136a.
with escape velocities. Lamers et al. (1995) obtained $v_{\infty} / v_{\text {esc }}=$ $2.65 \pm 0.2$ for 16 Galactic O stars earlier than O7. From 35 O2-6 stars in R136, we find $v_{\infty} / v_{\text {esc }}=2.5 \pm 0.4$, in agreement with the Galactic sample. Therefore, although R136 early O stars possess $\sim 15$ per cent lower wind velocities than Galactic counterparts, the ratio $v_{\infty} / v_{\text {esc }}$ is unchanged. Still, considerable variation is observed for other O stars in metal-poor galaxies (Garcia et al. 2014).

### 3.4 Physical parameters for R136 stars

De Koter et al. (1998) provided UV-derived parameters of the brightest dozen or so members in R136, while Massey \& Hunter (1998) provided a more complete Hertzsprung-Russell diagram from HST/FOS spectroscopy. Since then, Doran et al. (2013) have provided updated physical parameters for stars within R136 using contemporary calibrations for O stars in the LMC, based on the Galactic scale from Martins et al. (2005).

We defer more robust R136 O-type classifications until future papers in this series, but already our UV morphological classification allows us to present a preliminary, near complete HertzsprungRussell diagram for luminous stars within a radius of 2.5 arcsec ( 0.6 parsec) of R136a1. Temperatures and bolometric corrections of O stars are adopted from tables 5-7 of Doran et al. (2013), while properties of Of/WN stars are adopted from their Appendix A. HST/WFC3 photometry (de Marchi et al. 2011), except in particularly crowded regions (recall a1/a2, a4/a8), and interstellar extinctions provide absolute magnitudes as set out in Section 2.3.

In Fig. 10, we present the first spectroscopically derived Hertzsprung-Russell diagram for stars brighter than
$m_{\mathrm{F} 555 \mathrm{~W}}=16.0$ mag within 0.6 pc from R136a1. A dotted line corresponding to $m_{\mathrm{F} 555 \mathrm{~W}}=16.0 \mathrm{mag}$ for an adopted $A_{\mathrm{F} 555 \mathrm{~W}}=1.72 \mathrm{mag}$ is indicated, highlighting the region where our sample is complete. A total of 25 stars ( 26 including R136c), mostly early O dwarfs, are more massive than $\sim 50 \mathrm{M}_{\odot}$. The remaining stars are mid- to late O dwarfs with initial masses above $\sim 20 \mathrm{M}_{\odot}$.

### 3.5 Ages and masses of individual stars in R136

LMC isochrones for rotating stars with $v_{\mathrm{eq}}=100 \mathrm{~km} \mathrm{~s}^{-1}$ from Brott et al. (2011) and Köhler et al. (2015) are included, together with tracks for non-rotating $25,40,60,100$ and $200 \mathrm{M} \odot$. These tracks were selected as representative since they closely match the average rotation rate of O stars in 30 Doradus (Ramírez-Agudelo et al. 2013), though do not differ significantly from non-rotating models. It is apparent that, R136b aside, all stars with initial masses above $100 \mathrm{M}_{\odot}$ lie between the $0-1 \mathrm{Myr}$ isochrones, whereas stars with masses between 40 and $100 \mathrm{M}_{\odot}$ span $0-2.5 \mathrm{Myr}$, and those below $40 \mathrm{M}_{\odot}$ typically span $0-4 \mathrm{Myr}$, although uncertainties in physical parameters do not exclude ages below $\sim 2 \mathrm{Myr}$, with the exception of R136 H73 (O9+ V) for which an age of $\sim 5 \mathrm{Myr}$ is favoured.

For quantitative mass and age determinations, we have exploited bonNSAI ${ }^{2}$ (Schneider et al. 2014a) which implements a Bayesian method of quantifying ages and masses from comparison between UV-estimated temperatures and luminosities and evolutionary models, involving two key priors: a Salpeter mass function and an adopted rotation distribution of O stars from VLT Flames Tarantula Survey (VFTS, see Ramírez-Agudelo et al. 2013). Individual age and mass estimates are presented in Tables 4 and 5, and include results for the WN5h stars solely based upon their positions in the HR diagram, i.e. for consistency initially neglecting their He-abundances.
As suggested by Fig. 10, application of the Bonn evolutionary models (via bonNsai) to the three WN5h stars implies very young ages of $\leq 0.3 \mathrm{Myr}$ in all cases. Low initial rotation rates - favoured by observations of O stars in 30 Dor (Ramírez-Agudelo et al. 2013) - will lead to a redward evolution on the main sequence in evolutionary models, so young ages are required to explain the blue location of the WN5h stars. The next most luminous stars are extreme O supergiants, R136a5 (O2 If/WN5), a6 (O2 If) and b (O4 If/WN8), each with current masses of $100-150 \mathrm{M}_{\odot}$, plus O2-3 V stars R136a4 and a8. Once again, young ages of $0.6 \pm 0.3 \mathrm{Myr}$ are inferred for these stars, with the exception of R136b whose lower temperature suggests an age of $\sim 1.5 \mathrm{Myr}$. Schnurr et al. (2009) have previously argued that R136b, offset by a projected distance of 0.5 pc is more evolved than other core members of R136 (though see appendix of Hainich et al. 2014).
Consequently, the HR diagram positions of the WN5h stars and the most luminous O stars in R136a favours very young ages. However, it should be emphasized that envelope inflation is implemented in the Bonn models, which is most pronounced for the most massive models ( $\geq 50 \mathrm{M}_{\odot}$ ), since these lie closest to the Eddington limit (Sanyal et al. 2015). Inflation moves isochrones to lower temperatures, such that younger ages are inferred for stars with models accounting for inflation, compared with those that do not.

Turning to the lower mass stars, all $\mathrm{O} 2-6$ stars lie between the $0-3 \mathrm{Myr}$ isochrones in Fig. 10, although all stars do not lie along a

[^2]single isochrone. This is unsurprising given the approximate physical parameters of individual stars, initial rotation rate distributions, and uncertain binarity. For an extreme mass ratio system, the primary would completely dominate the spectrum of the system, so its location in the HR diagram would not be significantly displaced. In contrast, for a mass ratio of unity, individual stars would lie 0.3 dex lower in the HR diagram than the composite system, so the age would be overestimated. The effect would be modest for the highest mass systems at very young ages, but would become more significant for lower mass systems. Indeed, later O stars possess higher age estimates ( $3-6 \mathrm{Myr}$ ), although the inferred properties for such stars are least secure of the entire sample.

Our STIS/CCD data sets will help identify close binaries in R136 (Caballero-Nieves et al., in prep), while new HST Fine Guidance Sensor (FGS) observations - currently being analysed (GO 13477, PI, S.E. de Mink) - seek to reveal binaries with separations down to $0.02 \operatorname{arcsec}(1000 \mathrm{AU})$, following on from Cycle 3 FGS observations of R136 (Lattanzi et al. 1994).
Individual age estimates follow solely from the position of stars in the HR diagram. We are in the process of analysing STIS/CCD data sets of R136a members, from which rotational velocities and helium abundances can be obtained. For the moment, He abundances for the WN5h stars are available (Crowther et al. 2010). From their luminosities and He-rich atmospheres, we can infer minimum ages since the surface helium mass fraction sets a lower limit to the core helium mass fraction. The minimum ages of R136a1-3 are $1-1.3 \mathrm{Myr}$, in reasonable agreement with results for the O stars. In order to simultaneously reproduce the high surface temperatures of the WN5hs stars, Bonn evolutionary models favour rapid initial rotation ( $\sim 400 \mathrm{~km} \mathrm{~s}^{-1}$ ), and ages of $1.1-1.6 \mathrm{Myr}$, as shown in Table 7. These results compare favourably with Crowther et al. (2010) who applied Geneva evolutionary models for VMS (Yusof et al. 2013), also requiring rapid rotation ( $v_{\text {rot }} / v_{\text {crit }}=0.4$ or $v_{\text {rot }} \sim$ $400 \mathrm{~km} \mathrm{~s}^{-1}$ ) and ages of $\sim 1.5 \mathrm{Myr}$ (see Table 7). Details of the Bonn and Geneva codes differ significantly regarding the consideration of mixing, core overshooting and the treatment of envelope inflation. Both employ the metallicity-dependent Vink, de Koter \& Lamers (2001) prescription for mass-loss on the main sequence, so estimates of $M_{\text {init }}$ from current luminosities ought to be consistent, even if there are uncertainties in their validity for VMS (see Crowther et al. 2010; Gräfener \& Vink 2013)

### 3.6 R136 cluster age and mass

Early estimates of the age of R136 were 3-4 Myr on the basis that the WN stars were classical He-burning stars. Subsequently, a lower age of $<1-2 \mathrm{Myr}$ was obtained once it was recognized that the WN stars are core H -burning stars, supported by a rich early O-type population (de Koter et al. 1998, Crowther \& Dessart 1998). The most extensive study of massive stars to date has been by Massey \& Hunter (1998) who inferred a very young age of $\leq 1 \mathrm{Myr}$ using the Vacca, Garmany \& Schull (1996) O star temperature scale, or $\leq 2$ Myr from the Chlebowski \& Garmany (1991) scale, although the majority of their sample was drawn from the R136 region, rather than R136a itself.

Here, we have combined the age estimates of 52 O stars in R136a (excluding WN5h stars) from bonnsai to produce a probability density distribution, presented in Fig. 11, from which an age of $1.5_{-0.7}^{+0.3} \mathrm{Myr}$ is obtained. The $\pm 1 \sigma$ estimate represents the standard deviation in each bin of the age distribution computed with a bootstrapping method. As discussed above, lower ages are obtained

Table 7. Comparison of mass and age estimates for the WN5h stars in R136a from Geneva (Crowther et al. 2010; Yusof et al. 2013) and Bonn (Köhler et al. 2015) based on physical and chemical abundances obtained by Crowther et al. (2010), incorporating the adjustment to the absolute magnitude for R136a2, as discussed in the present study.

| Star | $T_{*}$ <br> $(\mathrm{kK})$ | $\log L$ <br> $(L \odot)$ | $X_{\mathrm{H}}$ <br> (per cent) | Code | $M_{\text {init }}$ <br> $\left(\mathrm{M}_{\odot}\right)$ | $v_{\text {init }}$ <br> $\left(\mathrm{km} \mathrm{s}^{-1}\right)$ | $M_{\text {current }}$ <br> $\left(\mathrm{M}_{\odot}\right)$ | $\tau$ <br> $(\mathrm{Myr})$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| R136a1 | $53 \pm 3$ | $6.94 \pm 0.09$ | - | Bonn | $325_{-45}^{+55}$ | $100_{-60}^{+180}$ | $315_{-50}^{+60}$ | $0.0_{-0.0}^{+0.3}$ |
|  | $53 \pm 3$ | $6.94 \pm 0.09$ | $40 \pm 10$ | Bonn | $315_{-20}^{+50}$ | $440_{-85}^{+20}$ | $280_{-30}^{+35}$ | $0.8 \pm 0.2$ |
|  | $53 \pm 3$ | $6.94 \pm 0.09$ | $40 \pm 5$ | Geneva | $320_{-40}^{+100}$ | 400 | $265_{-35}^{+80}$ | $1.4_{-0.1}^{+0.2}$ |
| R136a2 | $53 \pm 3$ | $6.63 \pm 0.09$ | - | Bonn | $195_{-30}^{+35}$ | $100_{-55}^{+325}$ | $190_{-35}^{+35}$ | $0.3_{-0.3}^{+0.4}$ |
|  | $53 \pm 3$ | $6.63 \pm 0.09$ | $35 \pm 5$ | Bonn | $160_{-20}^{+25}$ | $380_{-20}^{+85}$ | $130_{-20}^{+20}$ | $1.6 \pm 0.2$ |
|  | $53 \pm 3$ | $6.63 \pm 0.09$ | $35 \pm 5$ | Geneva | $180_{-30}^{+35}$ | 400 | $150_{-25}^{+30}$ | $1.7 \pm 0.1$ |
| R136a3 | $53 \pm 3$ | $6.58 \pm 0.09$ | - | Bonn | $180 \pm 30$ | $100_{-55}^{+330}$ | $175_{-35}^{+35}$ | $0.3_{-0.3}^{+0.4}$ |
|  | $53 \pm 3$ | $6.58 \pm 0.09$ | $40 \pm 5$ | Bonn | $155_{-20}^{+25}$ | $370_{-30}^{+80}$ | $130_{-15}^{+25}$ | $1.5 \pm 0.2$ |
|  | $53 \pm 3$ | $6.58 \pm 0.09$ | $40 \pm 5$ | Geneva | $165 \pm 30$ | 400 | $135_{-20}^{+25}$ | $1.7 \pm 0.1$ |



Figure 11. Probability density distribution of the apparent ages of UVderived spectral types of individual O stars (excluding R136a1, a2, a3) based on BONNSAI (Schneider et al. 2014a) and main-sequence evolutionary models (Brott et al. 2011; Köhler et al. 2015), favouring an age of $1.5_{-0.7}^{+0.3} \mathrm{Myr}$. The shaded area is a bootstrapped $\pm 1 \sigma$ estimate of the age distribution.
for the highest luminosity stars, although the treatment of envelope inflation for the most massive stars ( $>50 \mathrm{M}_{\odot}$ ) remains uncertain.

Discrepant age estimates for massive stars in other young massive star clusters (e.g. Arches and Quintuplet in Galactic Centre) have been obtained from their Wolf-Rayet and O star contents, in the sense that the luminous WR stars possess apparent ages which are younger than O stars by 1-2 Myr (Martins et al. 2008; Liermann, Hamann \& Oskinova 2012). Schneider et al. (2014b) have interpreted this as rejuvenation of the most massive stars through stellar mergers. Banerjee et al. (2012) have suggested the very luminous WN5h stars in R136 are post-merger systems. If this were the case, one might anticipate younger (apparent) age estimates for these rejuvenated stars with respect to the massive O star population.

R136 differs from the Quintuplet and Arches clusters in the sense that the luminous O stars also appear young, yet the He-enrichment of WN5h stars favours ages of $\sim 1.5 \mathrm{Myr}$ and unusually rapid rotation, the latter predicted by stellar mergers (de Mink et al. 2013). According to Schneider et al. (2014b), the probability that the most massive star in a $5 \times 10^{4} \mathrm{M}_{\odot}$ star cluster is a binary product rises from 50 per cent at 1 Myr to 100 per cent at 3 Myr . However,
only the most massive system is anticipated to have merged after 1.5 Myr (four are anticipated after 3 Myr ) for a Salpeter-type mass function in the core. Therefore, the youth of R136a does not support a scenario in which all the VMS are merger products. Indeed, as we discuss below, the mass within the central parsec is significantly lower than $5 \times 10^{4} \mathrm{M}_{\odot}$ further strengthening this conclusion. A more definitive answer to this question is deferred until the analysis of STIS/CCD data sets, but the apparent youth of the most luminous stars likely arises from the treatment of inflation in single star models than the physics of binary mergers.

A variety of mass estimates of R136 have been obtained. Hunter et al. (1995) infer a mass of $2.2 \times 10^{4} \mathrm{M}_{\odot}$ for stars with $>2.8 \mathrm{M}_{\odot}$ out to a radius of 5 pc from HST/WFPC2, suggesting a total mass of $\sim 5 \times 10^{4} \mathrm{M}_{\odot}$, while Andersen et al. (2009) derived a mass of $5 \times 10^{4} \mathrm{M}_{\odot}$ down to $2.1 \mathrm{M}_{\odot}$ out to a radius of 7 pc from HST/NICMOS. The actual mass of the R136 cluster may be somewhat lower, as it may suffer contamination from the surrounding halo of NGC 2070, for which Cignoni et al. (2015) infer a global mass of $9 \times 10^{4} \mathrm{M}_{\odot}$. Indeed, the two components contribute equally to the surface brightness at a radius of $\sim 5 \mathrm{pc}$ from R136a1 according to Mackey \& Gilmore (2003).

For a total cluster mass of $5 \times 10^{4} \mathrm{M}_{\odot}$, an initial mass function following Kroupa (2001) with an upper mass limit of $300 \mathrm{M}_{\odot}$, we would anticipate $10^{4} / 1250=40$ stars with initial masses in excess of $50 \mathrm{M} \odot$. Since only 25 stars with such high masses are observed within the central 0.5 pc radius, the actual R136a cluster mass may be somewhat lower. Of course, some massive stars may have been ejected from the cluster (e.g. VFTS 16; Evans et al. 2010). Fujii \& Portegies Zwart (2011) suggest a few per cent of stars with $\sim 50 \mathrm{M}_{\odot}$ would be dynamically ejected from dense clusters of this mass. If we were to consider a radius of 5 pc , analogous to the region surveyed by Hunter et al. (1995), there are another ~30 stars with masses in excess of $50 \mathrm{M}_{\odot}$ (from the census of Doran et al. 2013), supporting their high enclosed mass for the R136 'region'.

Sabbi et al. (2012) have proposed that R136 forms part of an ongoing merging binary cluster system within NGC 2070, on the basis of a clump of older stars located $\sim 5 \mathrm{pc}$ to its NE (see also Cignoni et al. 2015). If this clump were to reflect stars in the process of dynamically accreting, one might expect older stars to have already merged with R136 itself (see Gieles 2013). Indeed, some of the apparently older O stars in R136 (e.g. R136 H73) may arise from this clump, although the determination of physical parameters for these late-type O stars awaits detailed spectroscopic analysis of

Table 8. Census of the most luminous $\left(\log \left(L / L_{\odot}\right) \geq 6.2\right)$ objects in 30 Doradus, sorted by projected distance from R136a1, drawn from the present study, Crowther et al. (2010), Doran et al. (2013) or Bestenlehner et al. (2014). Initial mass estimates from bonnsai (based upon estimated temperatures and luminosities) are limited to apparently main-sequence stars.

| Star, Alias | r (pc) | Subtype | $\log L / L_{\odot}$ | $M / \mathrm{M}_{\odot}$ | Ref |
| :---: | :---: | :---: | :---: | :---: | :---: |
| R136a1, H3 | 0.00 | WN5h | $6.94 \pm 0.09$ | $325_{-45}^{+55}$ | a |
| R136a2, H5 | 0.02 | WN5h | $6.63 \pm 0.09$ | $195_{-30}^{+35}$ | a, e |
| R136a5, H20 | 0.07 | O2 If/WN5 | $6.32_{-0.16}^{+0.15}$ | $106_{-24}^{+30}$ | e |
| R136a8, H27 | 0.10 | O2-3 V | $6.28_{-0.15}^{+0.13}$ | $102_{-22}^{+26}$ | e |
| R136a4, H21 | 0.11 | O2-3 V | $6.46_{-0.15}^{+0.13}$ | $132_{-29}^{+37}$ | e |
| R136a3, H6 | 0.12 | WN5h | $6.58 \pm 0.09$ | $180 \pm 30$ | a, e |
| R136a6, H19 | 0.18 | O2 If | $6.52_{-0.05}^{+0.10}$ | $158_{-28}^{+35}$ | e |
| R136b, H9 | 0.52 | O4 If/WN8 | $6.30_{-0.12}^{+0.11}$ | $100_{-17}^{+22}$ | e |
| R136c, H10 | 0.83 | WN5h | $6.75 \pm 0.11$ | $230_{-45}^{+50}$ | a |
| Mk 42, H2 | 1.96 | O2 If | $6.56 \pm 0.1$ | $170_{-30}^{+40}$ | d |
| Mk 34, H8 | 2.61 | WN5h | $6.85 \pm 0.1$ | $275_{-45}^{+55}$ | c |
| R134, BAT99-100 | 2.86 | WN6(h) | 6.2 | - | d |
| Mk 37Wa, H11 | 2.86 | O4 If | $6.34 \pm 0.1$ | $110 \pm 20$ | d |
| Mk 37Wb, H44 | 2.94 | O2 If/WN5 | $6.21 \pm 0.1$ | $100 \pm 15$ | d |
| Mk 39, VFTS 482 | 3.01 | O2.5 If/WN6 | $6.4 \pm 0.1$ | $125_{-20}^{+25}$ | d |
| Mk 37, H14 | 3.02 | O3.5 If/WN7 | $6.48 \pm 0.1$ | $140_{-25}^{+40}$ | d |
| Mk 35, VFTS 545 | 3.07 | O2 If/WN5 | $6.3 \pm 0.1$ | $110 \pm 20$ | d |
| Mk 30, VFTS 542 | 4.15 | O2 If/WN5 | $6.16 \pm 0.1$ | $90 \pm 15$ | d |
| Mk 25, VFTS 506 | 10.9 | ON2 V((f*)) | $6.24 \pm 0.1$ | $100_{-15}^{+20}$ | d |
| R140b, BAT99-103 | 12.0 | WN5(h)+O | 6.4 | - | d |
| Mk 51, VFTS 457 | 12.6 | O3.5 If/WN7 | $6.2 \pm 0.1$ | $90 \pm 15$ | d |
| R139, VFTS 527 | 16.2 | O6.5 Iafc | 6.3 | $>78$ | b |
|  |  | O6 Iaf |  | $>66$ | b |
| R145, VFTS 695 | 19.6 | WN6h+ | 6.5 | - | d |
| R135, VFTS 402 | 21.6 | WN7h+OB | 6.2 | - | c |
| VFTS 682 | 29.8 | WN5h | $6.51 \pm 0.1$ | $155_{-25}^{+30}$ | d |
| R144, BAT99-118 | 62.3 | WN5-6 | 6.4 | - | c |
|  |  | WN6-7 |  | - | c |
| R147, VFTS 758 | 71.7 | WN5h | $6.36 \pm 0.1$ | $125_{-20}^{+25}$ | d |
| R146, VFTS 617 | 84.7 | WN5ha | $6.29 \pm 0.1$ | $115 \pm 20$ | d |
| VFTS 16 | 120: | O2 III-If | $6.23 \pm 0.1$ | $105_{-15}^{+20}$ | d |

Notes. (a): Crowther et al. (2010); (b): Taylor et al. (2011); (c): Doran et al. (2013); (d): Bestenlehner et al. (2014); (e): This work.
our optical STIS/CCD data sets. The evolutionary status of some stars in the NE clump (e.g. Mk 33Sb, WC5) implies an apparent age of 3-4 Myr, providing these are coeval (Melnick 1985; Walborn \& Blades 1997). The age of this region may be higher, should the highest luminosity stars be rejuvenated binary products.

### 3.7 Census of very luminous stars in 30 Doradus

Doran et al. (2013) have undertaken a study of the massive-star content of 30 Doradus, together with the contribution of R136 to the ionizing and mechanical feedback of the entire star-forming region, albeit reliant upon an incomplete census of R136 itself. We are now able to reassess the total number of very luminous stars in 30 Doradus, updated to include tailored parameters of massive stars from Bestenlehner et al. (2014) based on VFTS observations (Evans et al. 2011). Table 8 lists 29 stars within 30 Doradus which are more luminous than $\log \left(L / L_{\odot}\right) \approx 6.2$, corresponding to initial masses of $\sim 100 \mathrm{M}_{\odot}$ or more for zero age main-sequence stars. Only 9 ( 31 per cent) of the most luminous stars in 30 Doradus lie within R136 itself!

A similar number (10) of very luminous stars are located in the halo around R136, at projected distances of $1.5-5 \mathrm{pc}^{3}$ The other (11) high-luminosity stars in 30 Doradus lie $\geq 10 \mathrm{pc}$ away from R136a1, approximately half within NGC 2070 and the remainder elsewhere. The majority of these are luminous WN-type stars (see also Bestenlehner et al. 2014; Hainich et al. 2014).

The overwhelming majority of the VMS in 30 Doradus are early O supergiants, transition Of/WN stars or luminous WN stars. Crowther et al. (2010) proposed that VMS may possess Osupergiant spectral appearances on the main sequence, owing to their proximity to the (classical) Eddington limit. However, R136 $\mathrm{a} 4, \mathrm{a} 8$ and Mk 25 indicate that very luminous $\mathrm{O} 2-3$ dwarf stars do exist despite having classical Eddington parameters of $\Gamma_{\mathrm{e}} \sim 0.5$, based upon bonnsai-derived masses.

## 4 INTEGRATED FAR-UV SPECTROSCOPY

UV spectroscopy of the central region of NGC 2070 was first obtained with IUE (e.g. Koornneef \& Mathis 1981; Vacca et al. 1995), while the advent of $H S T / G H R S$-permitted high-resolution spectroscopy (Walborn et al. 1992). The integrated R136 spectrum revealed prominent He $_{\text {II }} \lambda 1640$ emission, together with strong Nv $\lambda \lambda 1238-42$ and C iv $\lambda \lambda 1548-51$ P Cygni features, with stellar Si Iv $\lambda \lambda 1393-1402$ absent.

Fig. 12 presents the integrated far-UV spectrum of R136 considering all sources up to 0.5 pc from R136a1. Its appearance closely reflects earlier descriptions, $\mathrm{He}_{\text {II }} \lambda 1640$ emission is relatively strong ( $W_{\lambda} \sim 4.5 \AA$ ) and broad (FWHM $\sim 10 \AA$ or $1800 \mathrm{~km} \mathrm{~s}^{-1}$ ), so the integrated spectral morphology is closer to an Of/WN supergiant than a genuine WN star - recall fig. 4 of Walborn et al. (1992) which compares R136a to Melnick 42 (O2 If). A comparison between the $2 \operatorname{arcsec} \times 2 \operatorname{arcsec}$ aperture GHRS/G140L observation of R136a from 1991 April (GO 1188, PI D.C. Ebbetts) and an equivalent region of our STIS data set reveals excellent agreement, and is morphologically identical to the integrated spectrum shown in Fig. 12 aside from a 30 per cent lower continuum flux.

Fig. 12 includes the co-added spectrum of all 70 stars for which $F_{1500} \geq 5 \times 10^{-15} \mathrm{erg} \mathrm{s}^{-1} \mathrm{~cm}^{2} \AA^{-1}$, which collectively contribute $\sim 77$ per cent of the integrated continuum flux, with the difference resulting from UV faint late O-type stars, plus the unresolved early $B$ population. The cumulative spectrum of the seven most luminous members of R136a - each with inferred initial masses in excess of $100 \mathrm{M}_{\odot}$ - comprising three WN stars (a1-a3), one Of/WN star (a5), an O2 supergiant (a6) plus two O2-3 dwarfs (a4, a8) - is also included in Fig. 12. Collectively, these most closely resemble a weak lined mid-WN integrated star. It is apparent that the $\mathrm{He}_{\text {II }} \lambda 1640$ emission line flux, and one third of the continuum flux originates from these few stars, the former dominated by R136a1-a3. The remaining 63 far-UV bright stars cumulatively resemble an early

[^3]

Figure 12. Integrated HST/STIS spectrum of all sources within 0.5 parsec of R136a1 (solid, all stars), the composite spectrum of all 70 bright stars with $F_{1500} \geq 5 \times 10^{-15} \mathrm{erg} \mathrm{s}^{-1} \mathrm{~cm}^{-2} \AA^{-1}$ (dot-dashed), comprising seven very massive stars (VMS, dotted), and the remaining 63 far-UV bright stars (dashed). He II $\lambda 1640$ emission in R136 is totally dominated by VMS. The difference between the 'stars' and 'all' arises from the contribution from UV faint O late-type stars and unresolved early B stars.
non-supergiant O star, with prominent $\mathrm{O} v \lambda 1371$, and P Cygni $\mathrm{N} v$ $\lambda \lambda 1238-42$, C iv $\lambda \lambda 1548-51$.

It is apparent that prominent $\mathrm{He}_{\text {II }} \lambda 1640$ emission in R136 arises from the presence of VMS (initial masses $\geq 100 \mathrm{M}_{\odot}$ ). Usually, the upper mass limit assumed in population synthesis models is $100 \mathrm{M}_{\odot}$, so a synthetic model for R136 would fail to reproduce its $\mathrm{He}_{\text {II }} \lambda 1640$ emission. Indeed, if this feature is observed in the integrated light of sufficiently young (unresolved) high-mass star clusters, one might legitimately conclude that the stellar mass function extends well beyond $100 \mathrm{M}_{\odot}$.

At very young ages, early O stars will dominate the UV appearance of massive clusters, so strong $\mathrm{O} v \lambda 1371$ absorption will be present. The absence of mid- to late O supergiants will ensure that P Cygni Si iv $\lambda \lambda 1393-1402$ is absent. At later ages (3-5 Myr) Ov will weaken while $\mathrm{Si}_{\text {IV }}$ will strengthen, plus $\mathrm{He}_{\text {II }} \lambda 1640$ from classic WR stars will appear. Therefore, the simultaneous presence of $\mathrm{O} v$ absorption, $\mathrm{He}_{\text {II }}$ emission together with the absence of P Cygni Si iv would suggest both a very young age and the presence of VMS. To date, the most promising example of such a star cluster is A1 within NGC 3125 (Chandar et al. 2004; Hadfield \& Crowther 2006; Wofford et al. 2014).

In addition, since our data sets enable spatially resolved spectroscopy of the constituents of R136, we are able to quantify the contribution of individual stars to the integrated UV continuum. five stars brighter than $m_{\mathrm{F} 555 \mathrm{~W}}=17.0 \mathrm{mag}$ located within $2.05 \operatorname{arcsec}$ of R136a1 lie beyond our STIS slit spectroscopy, so we have estimated their fluxes by adopting an average $F_{1500} / F_{\text {F555W }}$ ratio from cluster members; they are presented in parentheses in Tables 4 and B1. We estimate $\sim 2$ per cent of the integrated UV flux originates from these stars, primarily R136 H39 (O3 V+O5.5 V, Massey et al. 2002).

Fig. 13 presents the cumulative $1500 \AA$ Å flux of R136a to a radius of 0.5 pc , versus a ranked list of the 70 far-UV brightest stars in this region. The three brightest sources alone (R136a1-a3) contribute 20 per cent of the far-UV continuum flux, while the 18 brightest


Figure 13. Cumulative $F_{1500}$ flux for the brightest 70 far-UV sources within $2.05 \operatorname{arcsec}(0.5 \mathrm{pc})$ of R136a1, including estimates for stars beyond our slit spectroscopy. The three WN stars R136a1-a3 contribute 20 per cent of the total, while collectively the 18 brightest far-UV sources provide 50 per cent of the integrated continuum flux.
far-UV sources provide 50 per cent of the integrated far-UV flux. Walborn et al. (1992) cite an higher (pre-COSTAR) estimate of 40 per cent for R136a1-a3 from GHRS acquisition images, according to D. Ebbets and E. Malumuth.

## 5 SUMMARY

We have introduced a census of the visually brightest members of the R136 star cluster based on new long-slit HST/STIS spectroscopy. We have obtained UV, blue visual and $\mathrm{H} \alpha$ spectroscopic data sets using 17 contiguous $52 \operatorname{arcsec} \times 0.2$ arcsec slits sampling the core of R136. In this study, we reveal the following.
(i) We classify 51 of the 57 stars brighter than $m_{\text {F555W }}=16.0 \mathrm{mag}$ within a radius of 0.5 pc of R136al based on their UV spectroscopic morphology. We confirm previous results of Massey \& Hunter (1998) that the bulk of the visually brightest members of R136 are O2-3 stars. For the first time, we classify all Weigelt \& Baier (1985) stars within R136a, which comprise three WN5 stars (a1-a3), two O supergiants (a5-a6) and three early O dwarfs (a4, a7-a8). Eight additional sources are spectrally classified within the central ~pc, including R136b (O4 If/WN8).
(ii) We obtain wind terminal velocities from measurements of $v_{\text {black }}$ (or $v_{\text {edge }}$ ) from P Cygni C iv $\lambda \lambda 1548-51$ profiles, from which we obtain average $\mathrm{O} 2-3$ wind velocities of 2750 and $3120 \mathrm{~km} \mathrm{~s}^{-1}$ for dwarfs and (super)giants, respectively. Comparisons with early O dwarfs elsewhere in the LMC and Milky Way indicate somewhat lower wind velocities in R136.
(iii) We estimate physical parameters for the visually brightest members of R136 based on the temperature calibration of LMC O-type stars from Doran et al. (2013) and exploit bonnsai (Schneider et al. 2014a) to estimate masses and ages based on rotating evolutionary models appropriate for the LMC (Brott et al. 2011; Köhler et al. 2015). Three WN stars (R136a1-a3), two Of/WN stars (R136a5 and R136b), newly identified O2 If R136a6 and O2-3 V stars R136a4 and a8 possess initial masses $\geq 100 \mathrm{M}_{\odot}$.
(iv) We obtain a cluster age of $1.5_{-0.7}^{+0.3} \mathrm{Myr}$ from the O star content, which is consistent with that estimated for the luminous WN5h stars by Crowther et al. (2010), although apparent ages of these and other luminous O stars is $0.6 \pm 0.3 \mathrm{Myr}$, likely to arise from the treatment of envelope inflation in evolutionary models (Sanyal et al. 2015), rather than rejuvenation following binary mergers. The mass and youth of R136a indicate that no more than the most massive system is potentially a binary product (Schneider et al. 2014b). More definitive conclusions await analysis of STIS/CCD spectroscopy, permitting rotation rates and He abundances of O stars to be obtained. We identify 25 stars more massive than $50 \mathrm{M}_{\odot}$ within 0.5 pc of R136a1 indicating a lower enclosed mass than previous estimates, although reasonable consistency is achieved with respect to Hunter et al. (1995) and Andersen et al. (2009) in the R136 'region' from the total number of high-mass stars within a 5 pc radius.
(v) Incorporating results from our earlier VFTS-based census (Doran et al. 2013; Bestenlehner et al. 2014) we estimate a total of 29 stars more luminous than $\log (L / \mathrm{L} \odot)=6.2$ in the entire 30 Doradus nebula, of which 26 possess O2-4 If, Of/WN or WN subtypes. Only 9 stars are located within R136 itself, with an additional 9 stars within the extended halo, $1.5-5 \mathrm{pc}$ away from R136a1 (Melnick 1985 ) and another 11 stars at much greater distance ( $10-120 \mathrm{pc}$ ).
(vi) We have also considered the integrated UV spectrum of R136a in the context of its constituent members. The prominent Не II $\lambda 1640$ emission feature arises primarily from the most massive stars. Specifically, close to 100 per cent of the Не II $\lambda 1640$ flux and 32 per cent of the far-UV continuum originates in the seven stars for which we estimate (initial) masses above $100 \mathrm{M}_{\odot}$, the former dominated by from R136a1-a3. Usually, such stars are neglected in conventional population synthesis models. Prominent He it $\lambda 1640$ emission in the integrated spectrum of suitably young star clusters prominent $\mathrm{O} \vee \lambda 1371$ absorption without strong Si iv $\lambda \lambda 1393$, 1402P

Cygni emission - would favour a mass function that extends well above this value. Indeed, extremely strong $\lambda 1640$ emission in cluster A1 within NGC 3125 favours extremely massive stars therein (Chandar et al. 2004; Hadfield \& Crowther 2006; Wofford et al. 2014).

In future papers in this series, we will present the corresponding blue, visual and $\mathrm{H} \alpha$ STIS/CCD spectroscopic data sets. These will be used to refine spectral classifications, obtain more robust physical and wind properties, and search for close binary systems. These will enable a refined Hertzsprung-Russell diagram for R136a. The multi-epoch nature of the CCD data sets will enable an attempt at identifying close binaries (see Cottaar \& Hénault-Brunet 2014), together new HST FGS observations of R136a (S.E. de Mink, P.I.).

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## APPENDIX A: UV ATLAS OF LMC O STARS

In this appendix, we present an UV atlas of LMC O stars, primarily obtained with the FOS or STIS instruments of HST. A log of observations is presented in Table A1, with morphological sequences for O dwarfs, giants and supergiants presented in Figs A1-A4. Luminosity sequences for $\mathrm{O} 2-\mathrm{O} 9+$ stars are presented in Figs A5-A11. More extensive IUE SWP/LORES atlases for LMC OB stars are presented by Smith Neubig \& Bruhweiler (1999).

Table A1. Log of ultraviolet spectroscopic observations used in the LMC O star montages.

| Star | Alias | Sp type | Ref | Tel | Inst | HST GO/PI | Ref | Fig. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| BI 237 |  | O2 V | W02b | HST | STIS/G140L | 9412/Massey | M04 | A1, A5 |
| R136 H35 |  | O3 V | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A1, A6 |
| LH 81/W28-23 |  | O3.5 V ((f+)) | M05 | HST | STIS/G140L | 9412/Massey | M05 | A1, A7 |
| LH 81/W28-5 |  | O4 V((f+)) | M04 | HST | STIS/G140L | 8633/Massey | M04 | 6 |
| R136 H78 |  | O4: V | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A1, A8 |
| Sk-70 ${ }^{\circ} 6$ |  | O4-5 V((f)) | M09 | HST | FOS/G130H | 5444/Robert | PC98 | 6 |
| LH 58-496 | LH58-10a | O5 V(f) | M05 | HST | STIS/G140L | 9412/Massey | M05 | 6 |
| R136 H70 |  | 05 Vz | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A1, A2 |
| Sk-70 ${ }^{\circ} 69$ |  | O5.5 V((f)) | M09 | HST | FOS/G130H | 2233/Kudritzki | W95 | 6 |
| R136 H123 |  | O6 V | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A2, A9 |
| R136 H134 |  | 07 Vz | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A2, A10 |
| R136 H80 |  | O8 V | CN16 | HST | STIS/G140L | 12465/Crowther | Ths study | A2, A10 |
| R136 H121 |  | 09.5 V | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A2, A11 |
| HDE 269810 | Sk-67 211 | O2 III(f) | W02b | HST | FOS/G130H | 4110/Kudritzki | W95 | A3, A5 |
| R136 H48 |  | O2-3 III | CN16 | HST | STIS/G140L | 12465/Crowther | This work | A3, A6 |
| LH90 ST2-22 |  | O3.5 III(f+) | M05 | HST | STIS/G140L | 9412/Massey | M05 | A3, A7 |
| HDE 269676 | Sk-71 ${ }^{\circ} 45$ | O4-5 III | W77 | IUE | SWP/HIRES | Nandy | G80 | A3, A8 |
| Sk-66 ${ }^{\circ} 100$ |  | O6 II(f) | W95 | HST | FOS/G130H | 2233/Kudritzki | W95 | A3, A9 |
| Sk-67 ${ }^{\circ} 101$ |  | O8 II(f) | R03 | HST | STIS/E140M | 7299/Bomans | M03 | A3, A10 |
| R136 H36 |  | O2 If | CN16 | HST | STIS/G140L | 12465/Crowther | This study | A4, A5 |
| Sk-65 ${ }^{\circ} 47$ | LH 43-18 | O4 I(n)f+p | W10 | HST | STIS/G140L | 9412/Massey | M05 | A4, A8 |
| Sk-67 ${ }^{\circ} 111$ |  | O6 Ia(n)fpv | W02a | HST | STIS/E140M | 12218/Massa |  | A4, A9 |
| HDE 270952 | $\mathrm{Sk}-65^{\circ} 22$ | O6 Iaf+ | W77 | IUE | SWP/HIRES |  | C 02 | A9 |
| Sk-69 50 |  | O7(n)(f)p | W10 | HST | STIS/E140M | 12218/Massa |  | A10 |
| HDE 269702 | Sk-67 ${ }^{\circ} 168$ | O8 I(f)p | W10 | HST | STIS/E140M | 12218/Massa |  | A4, A10 |
| Sk-69 ${ }^{\circ} 124$ |  | O9 Ib | C86 | HST | FOS/G130H | 5444/Robert | PC98 | A11 |
| Sk-66 ${ }^{\circ} 169$ |  | O9.7 Ia+ | F88 | IUE | SWP/HIRES | Patriarchi | C02 | A4, A11 |
| Sk-68 ${ }^{\circ} 41$ |  | B0.5 Ia | F88 | HST | FOS/G130H | 4110/Kudritzki | W95 | A11 |
| Sk-67 ${ }^{\circ} 22$ |  | O2 If/WN5 | CW11 | HST | STIS/G140L | 9412/Massey | M05 | A5 |
| Mk 39 | R136 H7 | O2.5 If/WN6 | CW11 | HST | STIS/G140L | 9412/Massey | M05 | A6 |
| Mk 37 | R136 H14 | O3.5 If/WN7 | CW11 | HST | STIS/G140L | 9412/Massey | M05 | A7 |

Notes. C86: Conti, Garmany \& Massey 1986; C02: Crowther et al. 2002; CW11: Crowther \& Walborn 2011; CN16: Caballero-Nieves et al. (in preparation); G80: Gondhalekar et al. 1980; M03: Massa et al. 2003; M04: Massey et al. 2004; M05: Massey et al. 2005; M09: Massey et al. 2009; PC98: Prinja \& Crowther 1998; R03: Robert et al. 2003; W77: Walborn 1977; W95: Walborn et al. 1995; W02a: Walborn et al. 2002c; W02b: Walborn et al. 2002a; W10: Walborn et al. 2010.


Figure A1. Ultraviolet morphological progression of $\mathrm{O} 2-5$ dwarfs, showing strong, broad $\mathrm{O} v \lambda 1371$ absorption at $\mathrm{O} 2-3$, together with strong P Cygni Nv $\lambda \lambda 1238-42$ and C iv $\lambda \lambda 1548-51 . \mathrm{O} v$ weakens at O 3.5 and disappears at O 4 , while C iv P Cygni weakens at O 5 .


Figure A2. Ultraviolet morphological progression of O5-9.5 dwarfs, showing prominent broad, unsaturated P Cygni C iv $\lambda \lambda 1548-51$ at O5 V, which weakens in both emission strength and absorption width at $\mathrm{O} 6-7 \mathrm{~V}$ and becomes photospheric at $\mathrm{O} 9.5 \mathrm{~V} . \mathrm{N} v \lambda \lambda 1238-42$ is apparent at O 5 V but also weakens at $\mathrm{O} 6-8 \mathrm{~V}$ and is undetected at O 9.5 V .


Figure A3. Ultraviolet morphological progression of $\mathrm{O} 2-8$ giants, showing strong, broad $\mathrm{O} v \lambda 1371$ absorption at $\mathrm{O} 2-3$, with prominent $\mathrm{N} v \lambda \lambda 1238-42$ and C iv $\lambda \lambda 1548-51$ P Cygni profiles until O6 giants, and unsaturated thereafter.


Figure A4. Ultraviolet morphological progression of $\mathrm{O} 2-9.7$ supergiants, showing prominent $\mathrm{O} v \lambda 1371$ absorption at O 2 , strong, saturated C iv $\lambda \lambda 1548-51$ P Cygni profiles until O8 supergiants, Si iv $\lambda \lambda 1393-1402 \mathrm{P}$ Cygni appearing weakly at O 4 and saturating for O 6 and later subtypes. He II $\lambda 1640$ emission is prominent between O 2 and 6 , but thereafter blended with iron forest features.


Figure A5. Montage of O 2 luminosity sequence. Strong/broad $\mathrm{O} v \lambda 1371$ absorption is common to all O2 subtypes. C iv $\lambda \lambda 1548-51$ and $\mathrm{Nv} \lambda \lambda 1238-42$ are prominent P Cygni profiles at all luminosity classes, although the former is unsaturated in dwarfs/giants. Note the reversal in the He if $\lambda 1640$ line from absorption to strong P Cygni emission from dwarf to Of/WN.


Figure A6. Montage of $\mathrm{O} 2-3$ luminosity sequence. $\mathrm{O} v \lambda 1371$ absorption is common to all luminosity classes. Note the reversal in the $\mathrm{He}_{\text {II }} \lambda 1640$ line from absorption to strong P Cygni emission from dwarf to Of/WN.


Figure A7. Montage of O3.5 luminosity sequence. Ov $\lambda 1371$ absorption weakly is present, with Siiv $\lambda \lambda 1393-1402$ P Cygni appearing in Of/WN stars together with prominent $\mathrm{He}_{\text {II }} \lambda 1640 \mathrm{P}$ Cygni emission.


Figure A8. Montage of O4-5 luminosity class. $\mathrm{O} v \lambda 1371$ absorption is absent, although P Cygni $\mathrm{N} v \lambda \lambda 1238-42$ and C iv $\lambda \lambda 1548-51$ remain prominent for all classes, albeit unsaturated in dwarfs/giants. P Cygni Si iv $\lambda \lambda 1393-1402$ and He iI $\lambda 1640$ emission are seen in O4 supergiants.


Figure A9. Montage of O6 luminosity sequence. C iv $\lambda \lambda 1548-51$ and $N v \lambda \lambda 1238-42$ exhibit a weak P Cygni profile for dwarfs, in contrast to strong P Cygni profiles at class II-I, with Si iv $\lambda \lambda 1393-1402$ also exhibiting saturated P Cygni profiles for supergiants. Narrow He ir 1640 emission is also observed in supergiants.


Figure A10. Montage of O7-8 luminosity sequence. C Iv $\lambda \lambda 1548-51$ exhibits a narrow, weak $P$ Cygni profile for dwarfs, with $N v \lambda \lambda 1238-42$ barely present. Both lines display a prominent P Cygni profile at Class II (unsaturated for Nv v), while Si Iv $\lambda \lambda 1393-1402$ exhibit saturated P Cygni profiles amongst supergiants.


Figure A11. Montage of O9-B0. 5 luminosity sequence. C iv $\lambda \lambda 1548-51$ is photospheric in dwarfs, $\mathrm{N} v \lambda \lambda 1238-42$ absent, while C iv and $\operatorname{Si}$ iv $\lambda \lambda 1393-1402$ are saturated P Cygni profiles in luminous supergiants. $\mathrm{C}_{\text {III }} \lambda 1247$ is also prominent in luminous supergiants.

APPENDIX B: ULTRAVIOLET ATLAS OF EARLY-TYPE STARS IN R136


Figure B1. HST/STIS ultraviolet spectroscopy of WN stars in R136, together with an LMC template WN5 star R146 (Brey 88=BAT99-117, IUE SWP/LORES).


Figure B2. HST/STIS ultraviolet spectroscopy of Of/WN stars in R136, together with LMC templates Sk-67²2 and Melnick 37 (see Table A1).


Figure B3. HST/STIS ultraviolet spectroscopy of O2-3 (super)giants in R136, together with LMC templates R136 H48 (O2-3 III) and H36 (O2 If), see Table A1.


Figure B4. HST/STIS ultraviolet spectroscopy of O2-3 dwarfs in R136, together with templates BI 237 (O2 V) and R136 H35 (O3 V), see Table A1.


Figure B5. HST/STIS ultraviolet spectroscopy of O3-4 dwarfs in R136, together with LMC templates R136 H35 (O3 V) and H78 (O4 V); see Table A1.


Figure B6. HST/STIS ultraviolet spectroscopy of O4-5 dwarfs in R136, together with LMC templates R136 H78 (O4 V) and H70 (O5 Vz); see Table A1.

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Figure B7. HST/STIS ultraviolet spectroscopy of O5-6 dwarfs in R136, together with LMC templates R136 H70 (O5 Vz) and H123 (O6 V); see Table A1.


Figure B8. HST/STIS ultraviolet spectroscopy of O7-8 dwarfs in R136, together with LMC templates R136 H134 (O7 V) and H80 (O8 V); see Table A1.


Figure B9. HST/STIS ultraviolet spectroscopy of O9+ dwarfs in R136, together with LMC template R136 H121 (O9.5 V); see Table A1.

## APPENDIX C: WIND VELOCITIES OF LMC O-TYPE STARS

Table C1. Comparison of terminal wind velocities ( $\mathrm{km} \mathrm{s}^{-1}$ ) of LMC O stars in the present study (indicated with a colon if via $v_{\text {edge }}$ ) with literature values from far-UV spectroscopy (direct measurement or Sobolov with Exact Integral SEI approach).

| Star | Alias | Sp type | $\begin{aligned} & v_{\infty} \\ & \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{aligned}$ | Data set/diagnostic (this study) | $\begin{aligned} & v_{\infty} \\ & \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{aligned}$ | Data set/diagnostic (literature) | Reference |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| BI 237 |  | O2 V |  |  | 3400 | HST STIS/C iv | f |
| BI 253 | VFTS 72 | O2 V |  |  | 3180 | HST STIS/C iv | f |
| R136 H40 |  | O2-3 V | 2750 | HST STIS/C iv | 3400, 3000, 3400 | HST GHRS/C Iv | b, c, e |
| R136 H50 |  | O2-3 V | 2620 | HST STIS/C iv | 3125 | HST GHRS/C iv | b |
| R136 H55 |  | O2-3 V | 2880 | HST STIS/C iv | 3150, 2955, 3250 | HST GHRS/C iv | b, c, e |
| R136 H58 |  | O2-3 V | 2980 | HST STIS/C iv | 3050, 2840 | HST GHRS/C iv | b, c |
| LH 101/W3-14 |  | O3 V |  |  | 3100 | HST STIS/C iv | f |
| R136 H33 |  | O3 V |  |  | 3250 | HST STIS/C iv | f |
| LH 101/W3-24 |  | O3 V |  |  | 2400 | HST STIS/C iv | e |
| R136 H39 |  | O3 V+O5.5 V |  |  | 2900 | HST GHRS/C iv | b |
| LH 81/W28-23 |  | O3.5 V((f+)) |  |  | 3050 | HST STIS/C iv | f |
| R136a7 | R136 H24 | O3-4 V | 2710: | HST STIS/C iv | 2900, 3100 | HST GHRS/C iv | b, e |
| R136 H42 |  | O3-4 V | 2245 | HST STIS/C iv | 2850 | HST GHRS/C iv | b |
| LH 81/W28-5 |  | O4 V((f+)) |  |  | 2700 | HST STIS/C iv | e |
| Sk-70 60 |  | O4-5 V((f)) |  |  | 2300 | FUSE/SEI | d |
| LMC 054383 |  | O4-5 V |  |  | 2380 | HST STIS/C iv | g |
| LH 58-496 |  | O5 V(f) |  |  | 2400 | HST STIS/C iv | f |
| Sk-70 69 |  | O5.5 V((f)) |  |  | 2750 | FUSE/SEI | d |
| BI 229 |  | O7 V-III |  |  | 1950 | FUSE/SEI | d |
| Sk-67* 191 |  | O8 V |  |  | 1950 | FUSE/SEI | d |
| HDE 269810 | Sk-67 211 | O2 III(f) |  |  | 3600 | FUSE/SEI | d |
| Sk-66 ${ }^{\circ} 172$ |  | O2 III |  |  | 3100 | FUSE/SEI | d |
| Sk-70 ${ }^{\circ} 91$ | LH 62-1 | O2 III |  |  | 3150 | FUSE/SEI | d |
| Sk-68 137 |  | O2 III |  |  | 3400 | HST FOS/SEI | a |
| LH 64-16 |  | ON2 III |  |  | 3250 | HST STIS/C iv | g |
| VFTS 16 |  | O2 III-If |  |  | 3250 | HST COS/SEI | h |
| R136 H46 |  | O2 III-If | 3440 | HST STIS/C iv | 3550, 3355 | HST GHRS/C iv | b, c |
| R136 H47 |  | O2 III-If | 3045 | HST STIS/C iv | 3625, 3305, 3500 | HST GHRS/C iv | b, c, e |
| R136 H18 | Melnick 33S | O3 III |  |  | 3200 | $H S T /$ STIS/C iv | f |
| LH 90 ST2-22 |  | O3.5 III(f+) |  |  | 2560 | HST STIS/C iv | f |
| Sk-67 69 |  | O4 III |  |  | 2500 | FUSE/SEI | d |
| HDE 269676 | Sk-71 ${ }^{\circ} 45$ | O4-5 III |  |  | 2500 | FUSE/SEI | d |
| Sk-66 ${ }^{\circ} 100$ |  | O6 II(f) |  |  | 2075 | FUSE/SEI | d |
| BI 272 |  | O7: III-II |  |  | 3400 | FUSE/SEI | d |
| Sk-67 101 | LH 54-21 | O8 II(f) |  |  | 2300 | FUSE/SEI | d |
| BI 173 |  | O8 II |  |  | 2850 | FUSE/SEI | d |
| LH 101/W3-19 |  | O2 If |  |  | 2850 | HST STIS/C iv | f |
| Melnick 42 | R136 H2 | O2 If |  |  | 2800 | HST STIS/C iv | , |
| R136 H36 |  | O2 If | 3500 | HST STIS/C iv | 3640, 3700, 3750 | HST GHRS/C iv | b, c, e |
| Sk-65 ${ }^{\circ} 47$ | LH 43-18 | O4 I(n)f+p |  |  | 2100 | HST STIS/C iv | f |
| HDE 269698 | Sk-67 166 | O4 If + |  |  | 1900 | HST FOS/SEI | a |
| Sk-67* 167 | LH 76-21 | O4 Inf+ |  |  | 2150 | HST FOS/SEI | a |
| Sk-67 111 | LH 60-53 | O6 Ia(n)fp var |  |  | 2000 | FUSE/SEI | d |
| HDE 270952 | Sk-65 22 | O6 Iaf+ |  |  | 1350 | FUSE/SEI | d |
| HDE 270145 | Sk-70 115 | O6.5 Iaf |  |  | 2200 | FUSE/SEI | d |
| BI 170 |  | 09.5 Ib |  |  | 1700 | FUSE/SEI | d |
| Sk-69 124 |  | 09 Ib |  |  | 1600 | FUSE/SEI | d |
| Sk-65 ${ }^{\circ} 21$ |  | O9.7 Iab |  |  | 1700 | FUSE/SEI | d |
| Sk-66 ${ }^{\circ} 169$ |  | O9.7 Ia+ |  |  | 800 | FUSE/SEI | d |
| HDE 269896 | Sk-68 ${ }^{\circ} 135$ | ON9.7 Ia+ |  |  | 1050 | FUSE/SEI | d |
| HDE 268605 | Sk-67 05 | 09.7 Ib |  |  | 1400 | FUSE/SEI | d |
| Sk-67 ${ }^{\circ} 22$ |  | O2 If/WN5 |  |  | 2650 | HST STIS/C iv | f |
| R136a5 | R136 H20 | O2 If/WN5 | 3045 | HST STIS/C iv | 3000, 3400 | HST/GHRS C iv | b, e |
| Melnick 39 | R136 H7 | O2.5 If/WN6 |  |  | 2100 | HST STIS/C iv+Si iv | f |
| Melnick 37 | R136 H14 | O3.5 If/WN7 |  |  | 2000 | HST STIS/C iv+Si iv | f |
| LH 90 Brey 58 | AB4 | O3.5 If/WN7 |  |  | 1900 | HST STIS/Si iv | f |

Notes. (a) Walborn et al. (1995); (b) de Koter et al. (1998); (c) Prinja \& Crowther (1998); (d) Massa et al. (2003); (e) Massey et al. (2004); (f) Massey et al. (2005); (g) Massey et al. (2009); (h) Evans et al. (2010); (i) Bestenlehner et al. (2014).

## APPENDIX D: ADDITIONAL R136 SOURCES

Table D1. HST/WFC3 photometry (de Marchi et al. 2011) for additional R136 sources brighter than $m_{\mathrm{F} 555 \mathrm{~W}}=17.0$ mag within a projected distance of $4.1 \operatorname{arcsec}(1 \mathrm{parsec})$ from R136a1. Identifications are from Weigelt \& Baier (1985, WB85) or Hunter et al. (1995, HSH95). AF555W is calculated from observed colours for an adopted reddening law of $R_{5495}=4.1$, and $\mathrm{DM}=18.49$ (Pietrzyński et al. 2013). F555W-band photometry shown in italics is from Hunter et al. (1995) offset by -0.17 mag (see Section 2.3). R136 H351 is included since it is one of the 70 far-UV brightest sources in R136a1.

| HSH95 <br> (WB85) | Sp <br> type | Ref | $\begin{gathered} r \\ (\operatorname{arcsec}) \end{gathered}$ | $\begin{gathered} m_{\text {F555W }} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} m_{\mathrm{F} 336 \mathrm{~W}}-m_{\mathrm{F} 438 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} m_{\mathrm{F} 438 \mathrm{~W}}-m_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} m_{\mathrm{F} 555 \mathrm{~W}}-m_{\mathrm{F} 814 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{gathered} A_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{array}{r} M_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{array}$ | $\begin{gathered} 10^{14} F_{1500} \\ \left(\mathrm{erg} \mathrm{~s}^{-1} \mathrm{~cm}^{-2} \AA\right) \end{gathered}$ | Slit |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 379 |  |  | 0.41 | $16.51 \pm 0.05$ | $-0.97 \pm 0.06$ | $0.01 \pm 0.06$ | $0.60 \pm 0.08$ | 1.87 | -3.85 | $0.5 \pm 0.1$ | SE2 |
| 235 |  |  | 0.50 | 16.63 |  |  |  | 1.72 | -3.58 | $1.2 \pm 0.1$ | SE2 |
| 210 |  |  | 0.54 | 16.47 |  |  |  | 1.72 | -3.74 | $1.2 \pm 0.1$ | NW3 |
| 378 |  |  | 0.56 | $16.06 \pm 0.03$ |  |  |  | 1.72 | -4.15 | $1.0 \pm 0.1$ | SE2 |
| 306 |  |  | 0.70 | $16.42 \pm 0.05$ |  |  | $0.33 \pm 0.08$ | 1.81 | -3.88 | $0.5 \pm 0.1$ | NW4 |
| 196 |  |  | 0.74 | 16.36 |  |  |  | 1.72 | -3.85 |  | NW5 |
| 231 |  |  | 0.88 | $16.66 \pm 0.06$ | $-1.23 \pm 0.08$ | $0.11 \pm 0.08$ | $0.34 \pm 0.11$ | 1.58 | -3.41 | $0.4 \pm 0.1$ | NW2 |
| 305 |  |  | 0.91 | 16.96 |  |  |  | 1.72 | -3.25 |  | NW4 |
| 291 |  |  | 0.94 | $16.67 \pm 0.05$ | $-1.20 \pm 0.07$ | $0.16 \pm 0.07$ | $0.42 \pm 0.08$ | 1.72 | -3.54 | $0.4 \pm 0.1$ | SE4 |
| 259 |  |  | 1.00 | 16.75 |  |  |  | 1.72 | -3.46 | $0.4 \pm 0.1$ | NW5 |
| 217 |  |  | 1.01 | $16.52 \pm 0.04$ | $-1.36 \pm 0.06$ | $0.18 \pm 0.06$ |  | 1.48 | -3.45 | $0.5 \pm 0.1$ | NW1 |
| 274 |  |  | 1.06 | $16.87 \pm 0.07$ |  |  | $0.29 \pm 0.12$ | 1.72 | -3.34 | $0.4 \pm 0.1$ | NW4 |
| 187 |  |  | 1.07 | $16.16 \pm 0.03$ | $-1.31 \pm 0.04$ | $0.11 \pm 0.05$ | $0.27 \pm 0.05$ | 1.44 | -3.77 | $1.0 \pm 0.1$ | SE4 |
| 203 |  |  | 1.21 | $16.33 \pm 0.04$ | $-1.27 \pm 0.06$ | $0.17 \pm 0.06$ | $0.32 \pm 0.07$ | 1.61 | -3.77 | $0.6 \pm 0.1$ | NW3 |
| 233 |  |  | 1.24 | $16.76 \pm 0.05$ | $-1.25 \pm 0.08$ | $0.15 \pm 0.08$ | $0.32 \pm 0.09$ | 1.61 | -3.34 | $0.3 \pm 0.1$ | SE2 |
| 293 |  |  | 1.26 | $16.99 \pm 0.07$ | $-1.03 \pm 0.11$ | $0.17 \pm 0.10$ | $0.30 \pm 0.11$ | 2.05 | -3.55 | $0.3 \pm 0.1$ | NW7 |
| 159 |  |  | 1.34 | $16.14 \pm 0.03$ | $-1.24 \pm 0.04$ | $0.10 \pm 0.04$ | $0.33 \pm 0.04$ | 1.65 | -4.00 | $0.7 \pm 0.1$ | SE6 |
| 176 |  |  | 1.35 | $16.14 \pm 0.03$ | $-1.36 \pm 0.04$ | $0.13 \pm 0.04$ | $0.22 \pm 0.05$ | 1.36 | -3.71 | $0.9 \pm 0.1$ | NW7 |
| 179 |  |  | 1.40 | $16.23 \pm 0.03$ | $-1.31 \pm 0.05$ | $0.13 \pm 0.05$ |  | 1.47 | -3.73 | $0.9 \pm 0.1$ | NW6 |
| 254 |  |  | 1.43 | $16.87 \pm 0.05$ | $-1.25 \pm 0.08$ | $0.10 \pm 0.08$ | $0.32 \pm 0.09$ | 1.51 | -3.13 | $0.5 \pm 0.1$ | NW3 |
| 302 |  |  | 1.48 | $16.98 \pm 0.06$ | $-1.28 \pm 0.09$ | $0.10 \pm 0.09$ | $0.18 \pm 0.12$ | 1.47 | -2.98 | $0.9 \pm 0.1$ | NW8 |
| 223 |  |  | 1.53 | $16.61 \pm 0.04$ | $-1.30 \pm 0.06$ | $0.15 \pm 0.06$ | $0.26 \pm 0.07$ | 1.51 | -3.39 | $0.6 \pm 0.1$ | NW2 |
| 285 |  |  | 1.58 | $16.82 \pm 0.05$ | $-1.23 \pm 0.08$ | $0.17 \pm 0.07$ | $0.53 \pm 0.08$ | 1.72 | -3.39 | $0.4 \pm 0.1$ | NW2 |
| 288 |  |  | 1.62 | $16.80 \pm 0.04$ | $-1.18 \pm 0.06$ | $0.13 \pm 0.06$ | $0.29 \pm 0.07$ | 1.69 | -3.38 |  | - |
| 281 |  |  | 1.63 | $16.71 \pm 0.04$ | $-1.21 \pm 0.06$ | $0.14 \pm 0.05$ | $0.39 \pm 0.06$ | 1.65 | -3.43 | $0.5 \pm 0.1$ | SE1 |
| 139 |  |  | 1.65 | $16.02 \pm 0.02$ |  |  | $0.22 \pm 0.04$ | 1.72 | -4.19 | $1.1 \pm 0.1$ | SE3 |
| 162 |  |  | 1.72 | 16.06 |  |  |  | 1.72 | -4.15 | $0.8 \pm 0.1$ | SE2 |
| 386 |  |  | 1.75 | $16.87 \pm 0.05$ | $-0.97 \pm 0.09$ | $0.12 \pm 0.07$ | $0.36 \pm 0.08$ | 2.07 | -3.69 | $0.2 \pm 0.1$ | NW4 |
| 311 |  |  | 1.78 | $16.87 \pm 0.05$ | $-0.97 \pm 0.08$ | $0.12 \pm 0.07$ | $0.36 \pm 0.08$ | 2.07 | -3.69 | $0.3 \pm 0.1$ | NW4 |
| 211 |  |  | 1.78 | $16.52 \pm 0.03$ | $-1.31 \pm 0.04$ | $0.07 \pm 0.04$ | $0.22 \pm 0.06$ | 1.72 | -3.69 |  | - |
| 181 |  |  | 1.87 | 16.17 |  |  |  | 1.72 | -4.04 | $0.7 \pm 0.1$ | NW1 |
| 240 |  |  | 2.11 | 16.66 |  |  |  | 1.72 | -3.55 |  | - |
| 237 |  |  | 2.18 | $16.66 \pm 0.03$ | $-1.04 \pm 0.04$ | $0.09 \pm 0.04$ | $0.30 \pm 0.05$ | 1.90 | -3.73 |  | - |
| 416 |  |  | 2.20 | $16.95 \pm 0.04$ | $-0.90 \pm 0.06$ | $0.11 \pm 0.05$ | $0.44 \pm 0.06$ | 2.18 | -3.72 |  | - |
| 213 |  |  | 2.23 | $16.68 \pm 0.03$ | $-1.10 \pm 0.05$ | $0.12 \pm 0.05$ | $0.36 \pm 0.05$ | 1.83 | -3.64 |  | - |
| 253 |  |  | 2.30 | $16.89 \pm 0.04$ | $-1.27 \pm 0.05$ | $0.14 \pm 0.05$ | $0.27 \pm 0.06$ | 1.56 | -3.16 |  | - |
| 392 |  |  | 2.33 | $16.73 \pm 0.03$ | $-1.11 \pm 0.07$ | $0.07 \pm 0.07$ | $0.46 \pm 0.06$ | 1.72 | -3.48 | $0.3 \pm 0.1$ | SE8 |
| 205 |  |  | 2.37 | $16.37 \pm 0.02$ | $-1.22 \pm 0.03$ | $0.19 \pm 0.03$ | $0.46 \pm 0.03$ | 1.75 | -3.87 |  | - |
| 258 |  |  | 2.39 | $16.74 \pm 0.03$ |  | $0.17 \pm 0.05$ | $0.44 \pm 0.05$ | 1.72 | -3.47 |  | - |
| 189 |  |  | 2.42 | $16.31 \pm 0.02$ | $-1.23 \pm 0.03$ | $0.18 \pm 0.03$ | $0.32 \pm 0.03$ | 1.72 | -3.89 | $1.0 \pm 0.2$ | SE8 |
| 97 |  |  | 2.42 | $15.46 \pm 0.01$ | $-1.31 \pm 0.02$ | $0.18 \pm 0.02$ | $0.34 \pm 0.02$ | 1.76 | -4.79 |  | - |
| 87 |  |  | 2.47 | $15.25 \pm 0.01$ | $-1.32 \pm 0.01$ | $0.13 \pm 0.01$ | $0.32 \pm 0.02$ | 1.65 | -4.89 |  | - |
| 145 |  |  | 2.47 | $16.10 \pm 0.01$ | $-1.30 \pm 0.02$ | $0.10 \pm 0.02$ | $0.28 \pm 0.03$ | 1.42 | -3.81 | $0.8 \pm 0.1$ | NW2 |
| 547 |  |  | 2.51 | $16.75 \pm 0.03$ |  |  | $0.64 \pm 0.04$ | 1.72 | -3.46 |  | - |
| 255 |  |  | 2.55 | $16.77 \pm 0.03$ | $-1.17 \pm 0.04$ | $0.18 \pm 0.04$ | $0.48 \pm 0.04$ | 1.81 | -3.53 |  | - |
| 267 |  |  | 2.60 | $16.90 \pm 0.02$ | $-1.19 \pm 0.03$ | $0.12 \pm 0.03$ | $0.27 \pm 0.04$ | 1.68 | -3.27 |  | - |
| 206 |  |  | 2.63 | $16.53 \pm 0.02$ |  | $0.16 \pm 0.03$ | $0.23 \pm 0.04$ | 1.72 | -3.68 |  | - |
| 127 |  |  | 2.67 | $15.94 \pm 0.02$ | $-1.33 \pm 0.02$ | $0.11 \pm 0.02$ | $0.30 \pm 0.03$ | 1.59 | -4.15 |  | - |
| 177 |  |  | 2.80 | $16.26 \pm 0.02$ | $-1.18 \pm 0.02$ | $0.26 \pm 0.02$ | $0.52 \pm 0.03$ | 1.94 | -4.17 | $0.4 \pm 0.1$ | NW5 |
| 33 | O3 V | b | 2.80 | $14.36 \pm 0.01$ | $-1.36 \pm 0.01$ | $0.06 \pm 0.01$ | $0.28 \pm 0.01$ | 1.45 | -5.58 |  | - |
| 251 |  |  | 2.81 | $16.83 \pm 0.02$ |  | $0.13 \pm 0.03$ |  | 1.72 | -3.38 |  | - |
| 184 |  |  | 2.82 | $16.40 \pm 0.02$ | $-1.32 \pm 0.02$ | $0.17 \pm 0.02$ | $0.29 \pm 0.03$ | 1.52 | -3.61 |  | - |
| 294 |  |  | 2.86 | $17.00 \pm 0.03$ | $-1.24 \pm 0.04$ | $0.17 \pm 0.04$ | $0.46 \pm 0.04$ | 1.67 | -3.16 |  | - |
| 236 |  |  | 2.90 | $16.73 \pm 0.02$ | $-1.24 \pm 0.03$ | $0.13 \pm 0.03$ | $0.35 \pm 0.04$ | 1.59 | -3.35 |  | - |
| 241 |  |  | 2.92 | $16.80 \pm 0.02$ |  |  | $0.17 \pm 0.04$ | 1.72 | -3.41 |  | - |

Table D1 - continued

| HSH95 <br> (WB85) | $\begin{gathered} \mathrm{Sp} \\ \text { type } \end{gathered}$ | Ref | $r$ <br> (arcsec) | $m_{\text {F555W }}$ <br> (mag) | $\begin{gathered} m_{\mathrm{F} 336 \mathrm{~W}}-m_{\mathrm{F} 438 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $m_{\mathrm{F} 438 \mathrm{~W}}-m_{\mathrm{F} 555 \mathrm{~W}}$ (mag) | $\begin{gathered} m_{\mathrm{F} 555 \mathrm{~W}}-m_{\mathrm{F} 814 \mathrm{~W}} \\ (\mathrm{mag}) \end{gathered}$ | $\begin{array}{r} A_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{array}$ | $\begin{array}{r} M_{\mathrm{F} 555 \mathrm{~W}} \\ (\mathrm{mag}) \end{array}$ | $\begin{gathered} 10^{14} F_{1500} \\ \left(\mathrm{erg} \mathrm{~s}^{-1} \mathrm{~cm}^{-2} \AA\right) \end{gathered}$ | Slit |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 272 |  |  | 2.93 | $16.95 \pm 0.02$ | $-1.08 \pm 0.04$ | $0.21 \pm 0.04$ | $0.54 \pm 0.04$ | 2.03 | -3.57 |  | - |
| 170 |  |  | 2.97 | $16.25 \pm 0.01$ | $-1.23 \pm 0.02$ | $0.07 \pm 0.02$ | $0.33 \pm 0.03$ | 1.50 | -3.74 |  | - |
| 216 |  |  | 2.99 | $16.60 \pm 0.01$ | $-1.23 \pm 0.02$ | $0.10 \pm 0.02$ | $0.27 \pm 0.03$ | 1.55 | -3.44 |  | SE7 |
| 222 |  |  | 3.06 | $16.71 \pm 0.02$ | $-1.15 \pm 0.03$ | $0.22 \pm 0.03$ | $0.55 \pm 0.03$ | 1.92 | -3.70 | 0.1: | SE8 |
| 250 |  |  | 3.09 | $16.95 \pm 0.02$ | $-1.09 \pm 0.02$ | $0.22 \pm 0.03$ | $0.46 \pm 0.03$ | 2.04 | -3.58 | $0.2 \pm 0.1$ | NW1 |
| 277 |  |  | 3.22 | $16.96 \pm 0.02$ | $-1.13 \pm 0.03$ | $0.25 \pm 0.03$ | $0.62 \pm 0.03$ | 2.01 | -3.54 | $0.2 \pm 0.1$ | SE6 |
| 317 |  |  | 3.27 | $16.25 \pm 0.02$ | $-1.50 \pm 0.02$ | $0.63 \pm 0.02$ |  | 2.03 | -4.27 |  | - |
| 120 |  |  | 3.29 | $15.88 \pm 0.01$ | $-1.23 \pm 0.02$ | $0.05 \pm 0.02$ | $0.32 \pm 0.02$ | 1.65 | -4.26 | $1.1 \pm 0.1$ | SE4 |
| 10 (c) | WN5h | a | 3.44 | $13.43 \pm 0.01$ | $-1.20 \pm 0.02$ | $0.31 \pm 0.01$ | $0.82 \pm 0.01$ | 2.33 | -7.39 |  | - |
| 128 |  |  | 3.52 | $16.00 \pm 0.01$ | $-1.28 \pm 0.01$ | $0.07 \pm 0.02$ | $0.28 \pm 0.02$ | 1.41 | -3.90 |  | - |
| 136 |  |  | 3.64 | $15.96 \pm 0.01$ | $-1.15 \pm 0.02$ | $0.23 \pm 0.02$ | $0.62 \pm 0.02$ | 1.94 | -4.47 |  | - |
| 150 |  |  | 3.74 | $16.11 \pm 0.01$ | $-1.30 \pm 0.02$ | $0.16 \pm 0.02$ | $0.32 \pm 0.02$ | 1.54 | -3.92 |  | - |
| 165 |  |  | 3.81 | $16.16 \pm 0.01$ | $-1.27 \pm 0.02$ | $0.18 \pm 0.02$ | $0.35 \pm 0.02$ | 1.63 | -3.98 |  | - |
| 57 |  |  | 3.87 | $14.80 \pm 0.01$ | $-1.16 \pm 0.01$ | $0.26 \pm 0.01$ | $0.66 \pm 0.01$ | 2.17 | -5.86 |  | - |
| 197 |  |  | 4.00 | $16.38 \pm 0.01$ | $-1.17 \pm 0.02$ | $0.13 \pm 0.02$ | $0.45 \pm 0.02$ | 1.72 | -3.83 |  | - |
| 183 |  |  | 4.07 | $16.33 \pm 0.01$ | $-1.26 \pm 0.02$ | $0.24 \pm 0.02$ | $0.49 \pm 0.02$ | 1.76 | -3.92 |  | - |
| 140 |  |  | 4.09 | $15.99 \pm 0.01$ | $-1.29 \pm 0.02$ | $0.11 \pm 0.02$ | $0.26 \pm 0.02$ | 1.47 | -3.97 |  | - |
| 246 |  |  | 4.11 | $16.75 \pm 0.01$ | $-1.36 \pm 0.02$ | $0.12 \pm 0.02$ | $0.15 \pm 0.02$ | 1.36 | -3.10 |  | - |
| 351 |  |  | 0.55 | 17.18 |  |  |  | 1.72 | -3.03 | $0.5 \pm 0.1$ | SE2 |

Notes. (a) Crowther \& Dessart (1998); (b) Massey et al. (2005).
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[^1]:    ${ }^{1}$ Reduced wind momenta $\dot{M} v_{\infty} R^{1 / 2}$, where $\dot{M}$ is the mass-loss rate, $v_{\infty}$ is the terminal wind velocity and $R$ is the stellar radius.

[^2]:    ${ }^{2}$ BONNSAI is available at http://www.astro.uni-bonn.de/stars/bonnsai

[^3]:    ${ }^{3}$ We have excluded VFTS 512 (O2 V) from the Doran et al. (2013) statistics since the luminosity obtained by Bestenlehner et al. (2014) is $\log \left(L / \mathrm{L}_{\odot}\right)=$ 6.0. Source 1222 from Parker (1993) has also been excluded owing to its O9 V spectral type from Walborn, Maíz Apellániz \& Barbá (2002b), as has VFTS 1014 (Parker 863) despite Bestenhehner et al. obtaining $\log \left(L / L_{\odot}\right)=$ 6.2 since this ARGUS spectrum is likely composite ( $\mathrm{O} 3 \mathrm{~V}+$ mid-/late O ). Similarly, we favour the lower luminosities $\left.\left(\log L / L_{\odot}\right)=6.0-6.1\right)$ inferred by Doran et al. (2013) for R140a2 (WN5+O), Mk 49 (WN6) and BAT9996 (WN8) over those by Hainich et al. (2014), who obtained significantly higher values by assuming the WR dominated the continuum flux (R140a2) or on the basis of high interstellar extinctions (Mk 49, BAT99-96).

