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RESEARCH PAPER

ERROR ESTIMATES FOR APPROXIMATIONS OF DISTRIBUTED ORDER TIME FRACTIONAL DIFFUSION WITH NONSMOOTH DATA

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Abstract

In this work, we consider the numerical solution of a distributed order subdiffusion model, arising in the modeling of ultra-slow diffusion processes. We develop a space semidiscrete scheme based on the Galerkin finite element method, and establish error estimates optimal with respect to data regularity in $L^2(\Omega)$ and $H^1(\Omega)$ norms for both smooth and nonsmooth initial data. Further, we propose two fully discrete schemes, based on the Laplace transform and convolution quadrature generated by the backward Euler method, respectively, and provide optimal $L^2(\Omega)$ error estimates, which exhibits exponential convergence and first-order convergence in time, respectively. Extensive numerical experiments are provided to verify the error estimates for both smooth and nonsmooth initial data.

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Key Words and Phrases: distributed order, fractional diffusion, Galerkin finite element method, fully discrete scheme, error estimates

1. Introduction

We consider the following initial boundary value problem for a distributed order time fractional diffusion equation for $u(x, t)$:

$$\begin{aligned} D_t^{[\mu]}u - \Delta u &= 0 \quad \text{in } \Omega \times (0, T], \\ u &= 0 \quad \text{on } \partial\Omega \times (0, T], \\ u(0) &= v \quad \text{in } \Omega, \end{aligned} \tag{1.1}$$

where Ω is a bounded convex polygonal domain in \mathbb{R}^d ($d = 1, 2, 3$) with a boundary $\partial\Omega$, v is a given function on Ω , and $T > 0$ is a fixed value. Here, $D_t^{[\mu]}u$ denotes the distributed order fractional derivative of u in time t (with respect to μ) defined by

$$D_t^{[\mu]}u(t) = \int_0^1 \partial_t^\alpha u(t) \mu(\alpha) d\alpha,$$

where $\partial_t^\alpha u$, $0 < \alpha < 1$, is the Caputo derivative of order α in t , defined by [16, p. 91]

$$\partial_t^\alpha u(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t (t-s)^{-\alpha} u'(s) ds,$$

where $\Gamma(\cdot)$ denotes Euler's Gamma function. In this paper we consider a weight function $\mu \in C[0, 1]$ with $0 \leq \mu < 1$ and $\mu(0)\mu(1) > 0$.

In the last few decades, fractional calculus has been extensively studied and successfully employed to model subdiffusion, in which the mean squared variance grows slower than that in a Gaussian process. The subdiffusion model involves a Caputo derivative $\partial_t^{\alpha_0}u$ of order $\alpha_0 \in (0, 1)$ in t

$$\partial_t^{\alpha_0}u - \Delta u = f \quad \text{in } \Omega \times (0, T] \tag{1.2}$$

describes subdiffusion processes. The model (1.2) is recovered from (1.1) with a singular weight $\mu(\alpha) = \delta(\alpha - \alpha_0)$, the Dirac delta function at α_0 . Physically, subdiffusion can be characterized by a unique diffusion exponent known as the Hurst exponent [3]. In practice, the physical process may not possess a unique Hurst exponent, and the model (1.1) provides a flexible framework for describing a host of continuous and nonstationary signals [3, 4, 29]. It is often employed to describe ultraslow diffusion, where the mean squared variance grows only logarithmically with time.

In recent years, the theoretical study of problem (1.1) has attracted some attention. Kochubei [17] made some early contributions to the rigorous analysis of the model (1.1), by constructing fundamental solutions and establishing their positivity and subordination property. Mainardi et al. [22] studied the existence of a solution, asymptotic behavior, and positivity etc. Meerschaert and Scheffler [25] gave a stochastic model for ultraslow diffusion; see also [31]. Luchko [21] showed a weak maximum principle for the problem. Li et al [18] established a sharp asymptotic behavior of the solution for $t \rightarrow 0$ and $t \rightarrow \infty$, in the case of continuous density μ with $\mu(1) > 0$; see also [1] for further discussions.

The solution to (1.1) is rarely available in closed form, which necessitates the development of efficient numerical schemes. Despite the extensive studies on (1.2), there are only very few studies [6, 15, 26, 9, 7] on (1.1). Diethelm and Ford [6] developed a numerical scheme for distributed order fractional ODEs. It approximates $D_t^{[\mu]}u(t)$ by quadrature. This technique was also employed to solve nonlinear distributed-order fractional ODEs in [15], but without any analysis. Recently, Morgado and Rebelo [26] developed an implicit finite difference method for (1.1) with a Lipschitz nonlinear source term. The scheme approximates $D_t^{[\mu]}u(t)$ by quadrature together with the backward finite difference in time, and the second-order finite difference in space. Its stability, and an error estimate $O(h^2 + \tau + (\delta\alpha)^2)$ (with h, τ and $\delta\alpha$ being the mesh size, time step size and step size for quadrature rule, respectively) were established by assuming that the solution u is C^2 in time and C^4 in space and $\mu(\alpha)$ is sufficiently regular. See also [9] a high-order difference scheme with $v \equiv 0$, and also [7]. The regularity requirement is restrictive. The development of numerical schemes for (1.1) with nonsmooth data and their rigorous analysis have not been carried out, despite its importance, e.g., in inverse and optimal control problems [14].

In this work, we develop a Galerkin finite element method (FEM) for problem (1.1). It is based on the finite element space X_h of continuous piecewise linear functions over a family of shape regular quasi-uniform partitions $\{\mathcal{T}_h\}_{0 < h < 1}$ of Ω into d -simplexes, where h is the mesh size. Then the space semidiscrete Galerkin FEM is given by: find $u_h(t) \in X_h$ such that

$$(D_t^{[\mu]}u_h, \chi) + a(u_h, \chi) = 0, \quad \forall \chi \in X_h, \quad T \geq t > 0, \quad (1.3)$$

with $u_h(0) = v_h$, where (\cdot, \cdot) denotes the $L^2(\Omega)$ -inner product, $a(u, w) = (\nabla u, \nabla w)$ for $u, w \in H_0^1(\Omega)$, and $v_h \in X_h$ is an approximation of v . Further, we develop two fully discrete schemes based on the Laplace transform and convolution quadrature generated by the backward Euler method.

Our main contributions are as follows. First, in Theorem 2.1, we establish sharp regularity estimates for problem (1.1). Second, in Theorems 3.1 and 3.2, we derive optimal error estimates for the semidiscrete scheme (1.3). Third, we develop a fully discrete scheme based on the Laplace transform, using a contour representation of the semidiscrete solution with a hyperbolic contour and trapezoidal quadrature. We show its exponential convergence for a fixed time t , cf. Theorem 4.1. Last, we develop a second fully discrete scheme based on convolution quadrature, generated by the backward Euler method, and in Theorem 5.3, establish a first-order convergence. All these error estimates are nearly optimal and expressed in terms of the data regularity directly.

The model (1.1) is closely related to parabolic equations with a positive type memory term, for which there are many studies on numerical schemes

based on convolution quadrature [19, 20, 5] and Laplace transform [27, 28, 32]. These interesting works have inspired the current work. However, they do not cover (1.1), due to the general kernel function involved. Instead, we shall opt for a direct strategy by bounding the kernel function.

The rest of the paper is organized as follows. In Section 2, we recall the solution theory of the model (1.1). In Section 3, we develop a space semidiscrete Galerkin scheme. Two fully discrete schemes are given in Sections 4 and 5. Finally, to verify the theory, we present in Section 6 some numerical experiments. Throughout, the notation c , with or without a subscript, denotes a generic constant, which may differ at different occurrences, but it is always independent of the mesh size h , the number N of quadrature points, and time step size τ . Further, we denote by $\|\cdot\|$ the $L^2(\Omega)$ norm.

2. Solution Theory

In this part, we discuss the solution theory of problem (1.1). We denote by $\widehat{\cdot}$ the Laplace transform. Next we denote by A the operator $-\Delta$ with a homogeneous Dirichlet boundary condition with a domain $D(A) = H_0^1(\Omega) \cap H^2(\Omega)$. It is known that the operator A generates a bounded analytic semigroup of angle $\pi/2$, i.e., for any $\theta \in (\pi/2, \pi)$ [10, p. 321, Prop. C.4.2]

$$\|(zI + A)^{-1}\| \leq |\Im(z)|^{-1} \leq |z \sin(\theta)|^{-1} \quad \forall z \in \Sigma'_\theta, \quad (2.1)$$

where Σ'_θ is a sector with the origin excluded, i.e., $\Sigma_\theta = \{z \in \mathbb{C} : |\arg(z)| < \theta\}$, $\Sigma'_\theta = \Sigma_\theta \setminus \{0\}$. Then by (1.1) and the relation $\widehat{\partial_t^\alpha u} = z^\alpha \widehat{u} - z^{\alpha-1} u(0)$

$$zw(z)\widehat{u}(z) + A\widehat{u}(z) = w(z)v,$$

with $w(z) = \int_0^1 z^{\alpha-1} \mu(\alpha) d\alpha$. Hence, $u(t)$ is represented by

$$u(t) = S(t)v := \frac{1}{2\pi i} \int_{\Gamma_{\theta,\delta}} e^{zt} H(z)v dz, \quad (2.2)$$

where $H(z) = (zw(z)I + A)^{-1}w(z)$, and $\Gamma_{\theta,\delta} = \{z \in \mathbb{C} : |z| = \delta, |\arg z| \leq \theta\} \cup \{z \in \mathbb{C} : z = \rho e^{\pm i\theta}, \rho \geq \delta\}$.

Now we give a few properties of $w(z)$. The first is the sector-preserving property, since $z^\alpha \in \Sigma_\theta$ for any $\alpha \in (0, 1)$ and $z \in \Sigma_\theta$.

LEMMA 2.1. *Let $\theta \in (\pi/2, \pi)$. Then $zw(z) \in \Sigma_\theta$ for all $z \in \Sigma_\theta$.*

The second result is an upper bound on $w(z)$.

LEMMA 2.2. *Let $\mu \in C[0, 1]$ with $\mu \geq 0$. Then there holds*

$$|w(z)| \leq \|\mu\|_{C[0,1]} (|z| - 1) / (|z| \ln |z|).$$

The third result gives a lower bound on $zw(z)$.

LEMMA 2.3. *Let $\theta \in (\pi/2, \pi)$ with $\mu(0)\mu(1) > 0$. Then there exists a constant $c > 0$ dependent only on θ and μ such that for any $z \in \Sigma'_\theta$*

$$|zw(z)| \geq c(|z| - 1)/\ln |z|, \quad (2.3)$$

$$|z|w(|z|) \geq |zw(z)| \geq c|z|w(|z|). \quad (2.4)$$

P r o o f. Let $z = re^{i\varphi}$. Using $\mu(1) > 0$ and $\mu \in C[0, 1]$, we can find a small $\epsilon_1 > 0$ such that $\min_{\alpha \in [1-\epsilon_1, 1]} \mu(\alpha) \geq \delta_1 > 0$. Thus for $r \geq 1$, there holds $\int_0^1 r^\alpha \mu(\alpha) d\alpha \geq \int_{1-\epsilon_1}^1 r^\alpha \mu(\alpha) d\alpha \geq \delta_1 \int_{1-\epsilon_1}^1 r^\alpha d\alpha \geq \epsilon_1 \delta_1 \int_0^1 r^\alpha d\alpha$. Similarly, we may find a small $\epsilon_2 > 0$ such that $\min_{\alpha \in [0, \epsilon_2]} \mu(\alpha) \geq \delta_2 > 0$ and then for all $r < 1$, there holds $\int_0^1 r^\alpha \mu(\alpha) d\alpha \geq \epsilon_2 \delta_2 \int_0^1 r^\alpha d\alpha$. Hence for $\varphi \in (\theta - \pi, \pi - \theta)$, we get for $c_1 = \min(\epsilon_1 \delta_1, \epsilon_2 \delta_2)$

$$|zw(z)| \geq \Re(zw(z)) \geq \cos(\pi - \theta) \int_0^1 r^\alpha \mu(\alpha) d\alpha \geq c_1 \cos(\pi - \theta) \int_0^1 r^\alpha d\alpha.$$

It suffices to consider $\varphi \in [\pi - \theta, \theta]$. Since $\mu(0) > 0$, there exists a small $\epsilon_0 > 0$ such that $\min_{\alpha \in [0, \epsilon_0]} \cos(\alpha\pi)\mu(\alpha) = \delta_0 > 0$, and thus

$$\begin{aligned} \Re(zw(z)) &\geq \int_0^{\epsilon_0} r^\alpha \cos(\alpha\varphi)\mu(\alpha) d\alpha - \int_{\epsilon_0}^1 r^\alpha |\cos(\alpha\varphi)|\mu(\alpha) d\alpha \\ &\geq \delta_0 \int_0^{\epsilon_0} r^\alpha d\alpha - \|\mu\|_{C[0,1]} \int_{\epsilon_0}^1 r^\alpha d\alpha \\ &\geq -(\ln r)^{-1} [\delta_0 - r^{\epsilon_0}(\delta_0 + \|\mu\|_{C[0,1]})]. \end{aligned}$$

Consequently,

$$|zw(z)| \geq \Re(zw(z)) \geq \frac{\delta_0}{2} \int_0^1 r^\alpha d\alpha \quad \forall r \leq r_0 =: (\delta_0 / (2(\delta_0 + \|\mu\|_{C[0,1]})))^{1/\epsilon_0}.$$

Similarly, (2.3) holds for $r \geq r_0$ and $\varphi \in [\pi - \theta, \theta]$, thereby showing (2.3). The estimate (2.4) follows from $\|\mu\|_{C[0,1]} \int_0^1 r^\alpha d\alpha \geq \int_0^1 r^\alpha \mu(\alpha) d\alpha = |z|w(|z|)$ and the inequality $|zw(z)| \leq |z|w(|z|)$. \square

Now we give the stability of problem (1.1).

THEOREM 2.1. *Let $\mu \in C[0, 1]$ with $\mu \geq 0$ and $\mu(0)\mu(1) > 0$. Then the solution u to problem (1.1) satisfies for $t \in (0, T]$ and $\nu = 0, 1$:*

$$\|A^\nu S^{(m)}(t)v\| \leq c_T t^{-m-\nu} \ell_1(t)^\nu \|v\|, \quad v \in L^2(\Omega), m \geq 0, \quad (2.5)$$

$$\|A^\nu S^{(m)}(t)v\| \leq ct^{-m+1-\nu} \ell_2(t)^{1-\nu} \|Av\|, \quad v \in D(A), \nu + m \geq 1, \quad (2.6)$$

where $\ell_1(t) = (\ln(2T/t))^{-1}$, $\ell_2(t) = \ln(\max(t^{-1}, 2))$ and $c_T > 0$ may depend on d, Ω, μ, M, m and T .

P r o o f. The existence and uniqueness of a weak solution was shown in [18], and it suffices to show (2.5) and (2.6). First, by (2.1), we have $\|H(z)\| = \|(zw(z)I + A)^{-1}\| |w(z)| \leq M/|z|$ for all $z \in \Sigma'_\theta$. Let $t > 0$, $\theta \in (\pi/2, \pi)$, $\delta > 0$. We choose $\delta = t^{-1}$ and denote for short $\Gamma = \Gamma_{\theta, \delta}$. First we derive (2.5) for $\nu = 0$ and $m \geq 0$. By (2.2), we deduce

$$\begin{aligned} \|S^{(m)}(t)\| &= \left\| \frac{1}{2\pi i} \int_{\Gamma} z^m e^{zt} H(z) dz \right\| \leq c \int_{\Gamma} |z|^m e^{\Re(z)t} \|H(z)\| |dz| \\ &\leq c \left(\int_{t^{-1}}^{\infty} r^{m-1} e^{rt \cos \theta} dr + \int_{-\theta}^{\theta} e^{\cos \psi} t^{-m} d\psi \right) \leq ct^{-m}. \end{aligned}$$

Next we prove (2.5) for $\nu = 1$ and $m \geq 0$, by taking $\delta = 2T/t$ in Γ . By applying A to (2.2) and differentiating with respect to t we arrive at

$$AS^{(m)}(t) = \frac{1}{2\pi i} \int_{\Gamma} z^m e^{zt} AH(z) dz. \quad (2.7)$$

Since $AH(z) = A(zw(z)I + A)^{-1}w(z) = (I - zH(z))w(z)$, Lemma 2.2 yields $\|AH(z)\| \leq c|w(z)| \leq c(|z| - 1)/(|z| \ln |z|)$ for all $z \in \Sigma'_\theta$. Hence (2.7) gives

$$\begin{aligned} \|AS^{(m)}(t)\| &\leq c \int_{\Gamma} |z|^m \frac{|z| - 1}{|z| \ln |z|} e^{\Re(z)t} |dz| \leq c \int_{\delta}^{\infty} r^{m-1} \frac{r-1}{\ln r} e^{rt \cos \theta} dr \\ &\quad + c_T t^{-m} \frac{\delta - 1}{\ln \delta} \int_{-\theta}^{\theta} e^{2T \cos \psi} d\psi =: I + II. \end{aligned}$$

Since $\delta \geq 2$, we can bound the first term I by

$$I \leq c \int_{\delta}^{\infty} r^m (\ln r)^{-1} e^{rt \cos \theta} dr \leq c \ell_1(t) \int_{\delta}^{\infty} r^m e^{rt \cos \theta} dr \leq c_T t^{-m-1} \ell_1(t).$$

The second term II is bounded by $II = c_T t^{-m} (\delta - 1)/\ln \delta \leq c_T t^{-m-1} \ell_1(t)$. This shows (2.5). To prove (2.6) with $\nu = 0$, we choose $\delta = t^{-1}$. Then,

$$S^{(m)}(t)v = \frac{1}{2\pi i} \int_{\Gamma} z^m e^{zt} H(z)v dz = \frac{1}{2\pi i} \int_{\Gamma} z^{m-1} e^{zt} zA^{-1}H(z)Av dz.$$

Since $zA^{-1}H(z) = A^{-1} - (zw(z)I + A)^{-1}$ and $\int_{\Gamma} z^{m-1} e^{zt} dz = 0$ for $m \geq 1$, we have

$$\begin{aligned} S^{(m)}(t)v &= \frac{1}{2\pi i} \int_{\Gamma} z^{m-1} e^{zt} v dz - \frac{1}{2\pi i} \int_{\Gamma} z^{m-1} e^{zt} (zw(z)I + A)^{-1} dz Av \\ &= -\frac{1}{2\pi i} \int_{\Gamma} z^{m-1} e^{zt} (zw(z)I + A)^{-1} dz Av. \end{aligned}$$

By (2.1) and Lemma 2.3 we obtain $\|(zw(z)I + A)^{-1}\| \leq M|zw(z)|^{-1} \leq c(|z| - 1)^{-1} \ln |z|$. This and the monotonicity of $f(x) = -\ln(x)/(1-x)$ on the positive real axis \mathbb{R}^+ yield

$$\begin{aligned}
 \|S^{(m)}(t)v\| &\leq c \left(\int_{\Gamma} |z|^{m-1} e^{\Re(z)t} \|(zw(z)I + A)^{-1}\| |dz| \right) \|Av\| \\
 &\leq c \left(\int_{t^{-1}}^{\infty} e^{rt \cos \theta} r^{m-1} \frac{\ln r}{r-1} dr + t^{-m} \frac{\ln t^{-1}}{t^{-1}-1} \int_{-\theta}^{\theta} e^{\cos \psi} d\psi \right) \|Av\| \\
 &\leq ct^{-m+1} \ln(t^{-1})(1-t)^{-1} \|Av\|.
 \end{aligned}$$

For $t^{-1} \geq 2$, $\ln(t^{-1})/(1-t) \leq 2 \ln(t^{-1})$, and for $t^{-1} < 2$, $\ln(t^{-1})/(1-t) = \ln(t)/(t-1) \leq 2 \ln(2)$. Thus $\|S^{(m)}(t)v\| \leq ct^{-m+1} \ell_2(t) \|Av\|$. Last, (2.6) with $\nu = 1$ is equivalent to (2.5) with $\nu = 0$ and v by Av . \square

3. Semidiscrete Discretization by Galerkin FEM

Now we discuss the space semidiscrete scheme (1.3) based on the Galerkin FEM. On the space X_h , we define the $L^2(\Omega)$ projection $P_h : L^2(\Omega) \rightarrow X_h$ and the Ritz projection $R_h : H_0^1(\Omega) \rightarrow X_h$, respectively, by

$$\begin{aligned}
 (P_h \varphi, \chi) &= (\varphi, \chi) \quad \forall \chi \in X_h, \\
 (\nabla R_h \varphi, \nabla \chi) &= (\nabla \varphi, \nabla \chi) \quad \forall \chi \in X_h.
 \end{aligned}$$

With the discrete Laplacian $\Delta_h : X_h \rightarrow X_h$ defined by $-(\Delta_h \varphi, \chi) = (\nabla \varphi, \nabla \chi)$ for all $\varphi, \chi \in X_h$, and $A_h = -\Delta_h$, (1.3) can be rewritten as

$$D_t^{[\mu]} u_h(t) + A_h u_h(t) = 0, \quad t > 0 \quad (3.1)$$

with $u_h(0) = v_h \in X_h$. For the error analysis, we employ an operator trick [8]. To this end, we first represent u_h by

$$u_h(t) = S_h(t)v_h := \frac{1}{2\pi i} \int_{\Gamma_{\theta, \delta}} e^{zt} (zw(z)I + A_h)^{-1} w(z) v_h dz. \quad (3.2)$$

Now we introduce the error function $e(t) := u(t) - u_h(t)$, which, by (2.2) and (3.2), is given by

$$e(t) = \frac{1}{2\pi i} \int_{\Gamma_{\theta, \delta}} e^{zt} w(z) (\widehat{\varphi}(z) - \widehat{\varphi}_h(z)) dz, \quad (3.3)$$

with $\widehat{\varphi}(z) = (zw(z)I + A)^{-1}v$ and $\widehat{\varphi}_h(z) = (zw(z)I + A_h)^{-1}P_h v$. The next lemma shows a bound on $\widehat{\varphi}_h - \widehat{\varphi}$. It follows from Lemma 2.1, similar to [2, Lemmas 3.3 and 3.4], and hence the proof is omitted.

LEMMA 3.1. *Let $v \in L^2(\Omega)$, $z \in \Sigma_{\theta}$ with $\theta \in (\pi/2, \pi)$, $\widehat{\varphi}(z) = (zw(z)I + A)^{-1}v$ and $\widehat{\varphi}_h(z) = (zw(z)I + A_h)^{-1}P_h v$. Then there holds*

$$\|\widehat{\varphi}(z) - \widehat{\varphi}_h(z)\| + h \|\nabla(\widehat{\varphi}(z) - \widehat{\varphi}_h(z))\| \leq ch^2 \|v\|.$$

Now we can state an error estimate for nonsmooth data $v \in L^2(\Omega)$.

THEOREM 3.1. *Let u and u_h be the solutions of problem (1.1) and (3.1) with $v \in L^2(\Omega)$ and $v_h = P_h v$, respectively. Then for $t > 0$ and $\ell_1(t) = \ln(2T/t)^{-1}$, there holds*

$$\|u(t) - u_h(t)\| + h\|\nabla(u(t) - u_h(t))\| \leq c_T h^2 t^{-1} \ell_1(t) \|v\|.$$

P r o o f. By (3.3), with $\delta = 2T/t$ in $\Gamma_{\theta,\delta}$, and Lemmas 3.1 and 2.2,

$$\|\nabla e(t)\| \leq ch \int_{\delta}^{\infty} e^{rt \cos \theta} \frac{r-1}{r \ln r} dr \|v\| + ch \int_{-\theta}^{\theta} e^{2T \cos \psi} \frac{\delta-1}{\ln \delta} d\psi \|v\| := I + II.$$

Now the first term I can be bounded by

$$I \leq ch \int_{\delta}^{\infty} \frac{e^{rt \cos \theta}}{\ln r} dr \|v\| \leq \frac{ch}{\ln \delta} \int_{\delta}^{\infty} e^{rt \cos \theta} dr \|v\| \leq c_T h t^{-1} \ell_1(t) \|v\|,$$

and the second term II by

$$II \leq \frac{c_T h}{t \ln \delta} \int_{-\theta}^{\theta} e^{2T \cos \psi} d\psi \|v\| \leq c_T h t^{-1} \ell_1(t) \|v\|.$$

Thus the bound on $\|\nabla e(t)\|$ follows. The L^2 estimate is similar. \square

Next we turn to the case of smooth data $Av \in L^2(\Omega)$.

THEOREM 3.2. *Let u and u_h be the solutions of problem (1.1) and (3.1) with $v \in D(A)$ and $v_h = R_h v$, respectively. Then for $t > 0$,*

$$\|u(t) - u_h(t)\| + h\|\nabla(u(t) - u_h(t))\| \leq ch^2 \|Av\|.$$

P r o o f. Like before, we take $\theta \in (\pi/2, \pi)$ and $\delta = 1/t$ in the contour $\Gamma_{\theta,\delta}$. Then the error $e_h(t) = u(t) - u_h(t)$ can be represented by

$$e_h(t) = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\delta}} e^{zt} w(z) ((zw(z)I + A)^{-1} - (zw(z)I + A_h)^{-1} R_h) v dz.$$

Using $w(z)(zw(z)I + A)^{-1} = z^{-1}I - z^{-1}(zw(z)I + A)^{-1}A$, we deduce

$$e_h(t) = \frac{1}{2\pi i} \left(\int_{\Gamma_{\theta,\delta}} e^{zt} z^{-1} (\widehat{\varphi}_h(z) - \widehat{\varphi}(z)) dz + \int_{\Gamma_{\theta,\delta}} e^{zt} z^{-1} (v - R_h v) dz \right),$$

where $\widehat{\varphi}(z) = (zw(z)I + A)^{-1}Av$ and $\widehat{\varphi}_h(z) = (zw(z)I + A_h)^{-1}A_h R_h v$. Then Lemma 3.1 and the identity $A_h R_h = P_h A$ give

$$\|\widehat{\varphi}(z) - \widehat{\varphi}_h(z)\| + h\|\nabla(\widehat{\varphi}(z) - \widehat{\varphi}_h(z))\| \leq ch^2 \|Av\|.$$

Now it follows that

$$\|e_h(t)\| \leq ch^2 \|Av\| \left(\int_{t^{-1}}^{\infty} e^{rt \cos \theta} r^{-1} dr + \int_{-\theta}^{\theta} e^{\cos \psi} d\psi \right) \leq ch^2 \|Av\|,$$

which gives the L^2 -estimate. The H^1 estimate is analogous. \square

REMARK 3.1. The estimate for $v \in L^2(\Omega)$ deteriorates like $t^{-1}\ell_1(t)$ as $t \rightarrow 0^+$, which agrees with Theorem 2.1. The factor $t^{-1}\ell_1(t)$ is different from that for subdiffusion [12] and multi-term counterpart [11]. For $Av \in L^2(\Omega)$, the estimate is uniform in time.

4. Fully Discrete Scheme I: Laplace Transform

The first fully discrete scheme is based on Laplace transform, by applying a quadrature rule to the representation (3.2). We follow [24, 27, 28, 32] and deform the contour $\Gamma_{\theta,\delta}$ to be a curve with a parametric representation

$$z(\xi) := \lambda(1 + \sin(i\xi - \psi)), \quad (4.1)$$

with $\lambda > 0$, $\psi \in (0, \pi/2)$ and $\xi \in \mathbb{R}$. The optimal choices of λ and ψ will be given in Lemma 4.4. This deformation is valid since it does not transverse the poles of $H(z)v = (zw(z) + A_h)^{-1}w(z)v$, cf., Lemmas 2.1 and 4.3. With $z = x + iy$, the contour (4.1) is the left branch of the hyperbola

$$((x - \lambda)/(\lambda \sin \psi))^2 - (y/(\lambda \cos \psi))^2 = 1,$$

which intersects the real axis at $x = \lambda(1 \pm \sin \psi)$ and has asymptotes $y = \pm(\lambda - x) \cot \psi$. Now we can represent $u_h(t)$ by

$$u_h(t) = \int_{-\infty}^{\infty} \widehat{g}(\xi, t) d\xi \quad (4.2)$$

with the integrand $\widehat{g}(\xi, t)$ being defined by

$$\widehat{g}(\xi, t) = (2\pi i)^{-1} e^{z(\xi)t} (z(\xi)w(z(\xi))I + A_h)^{-1} w(z(\xi))z'(\xi)v_h. \quad (4.3)$$

Now we describe the quadrature approximation. By setting $z_j = z(\xi_j)$ and $z'_j := z'(\xi_j)$ with $\xi_j = jk$ and k being the step size, we have

$$U_h(t) = \frac{k}{2\pi i} \sum_{j=-\infty}^{\infty} e^{z_j t} \widehat{\phi}_j z'_j,$$

and the truncated quadrature approximation

$$U_{N,h}(t) = \frac{k}{2\pi i} \sum_{j=-N}^N e^{z_j t} \widehat{\phi}_j z'_j, \quad (4.4)$$

with $\widehat{\phi}_j = (z_j w(z_j)I + A_h)^{-1} w(z_j)v_h$. To compute $U_{N,h}(t)$, we need to solve only $N+1$ elliptic problems, instead of $2N+1$, by exploiting conjugacy: $z_{-j} = \overline{z_j}$, $w(z_{-j}) = \overline{w(z_j)}$, $\widehat{\phi}_{-j} = \overline{\widehat{\phi}_j}$, $j = 1, \dots, N$. Indeed, since $z'_j = z'(\xi_j) = i\lambda \cos(i\xi_j - \psi)$, with $\zeta_j = \lambda \cos(i\xi_j - \psi)$, (4.4) is reduced to

$$U_{N,h}(t) = \frac{k}{2\pi} e^{z_0 t} \widehat{\phi}_0 \zeta_0 + \frac{k}{\pi} \sum_{j=1}^N \Re\{e^{z_j t} \widehat{\phi}_j \zeta_j\}.$$

Hence we solve the following complex-valued elliptic problems

$$(z_j w(z_j)I + A_h) \widehat{\phi}_j = w(z_j)v_h, \quad j = 0, \dots, N.$$

Next we define a strip by $\mathcal{S}_{a,b} = \{p = \xi + i\eta : \xi \in \mathbb{R}, \eta \in (-b, a)\}$. The following lemma gives the quadrature error [23] [32, Theorem 2.1].

LEMMA 4.1. *Let g be an analytic function in a strip $\mathcal{S}_{a,b}$ for some $a, b > 0$, and I and I_k , for $k > 0$, be defined by $I = \int_{-\infty}^{\infty} g(x) dx$ and $I_k = k \sum_{j=-\infty}^{\infty} g(jk)$, respectively. Furthermore, assume that $g(z) \rightarrow 0$ uniformly as $|z| \rightarrow \infty$ in $\mathcal{S}_{a,b}$, and that there exist $M_+, M_- > 0$ such that*

$$\lim_{r \rightarrow a^-} \int_{-\infty}^{\infty} |g(x + ir)| dx \leq M_+ \quad \text{and} \quad \lim_{s \rightarrow b^-} \int_{-\infty}^{\infty} |g(x - is)| dx \leq M_-.$$

Then with $E^+ = M_+/(e^{2a\pi/k} - 1)$ and $E^- = M_-/(e^{2b\pi/k} - 1)$, there holds

$$|I - I_k| \leq E^+ + E^-.$$

The next lemma gives one crucial estimate on the map $z(p)$ over the strip $\mathcal{S}_{a,b}$. Even though the hyperbolic contour (4.1) has been extensively used, the estimate on the map $z(p)$ below seems to be new.

LEMMA 4.2. *Let $p = \xi + i\eta$ with $\xi, \eta \in \mathbb{R}$. Then with $a = \pi/2 - \psi - \epsilon$ and $b = \psi - \epsilon$, for small $\epsilon > 0$, there holds*

$$z(p) \in \Sigma_{\pi-\psi} \quad \text{and} \quad |z'(p)||z(p)|^{-1} \leq c\epsilon^{-1} \quad \forall p \in \overline{\mathcal{S}}_{a,0}, \quad (4.5)$$

$$z(p) \in \Sigma_{\pi-\epsilon} \quad \text{and} \quad |z'(p)||z(p)|^{-1} \leq c \quad \forall p \in \overline{\mathcal{S}}_{0,b}. \quad (4.6)$$

P r o o f. For $p = \xi + i\eta$ with $\xi, \eta \in \mathbb{R}$, then the image $z(p)$ in the parameterization (4.1) is given by $z(p) = \lambda(1 - \sin(\psi + \eta) \cosh(\xi)) + i\lambda \cos(\psi + \eta) \sinh(\xi)$, and its derivative $z'(p)$ is given by $z'(p) = \lambda \cosh \xi \cos(\psi + \eta) - i \sinh \xi \sin(\psi + \eta)$. By writing $z = x + iy$, it can be expressed as the left branch of the hyperbola $((x - \lambda)/(\lambda \sin(\psi + \eta)))^2 - (y/(\lambda \cos(\psi + \eta)))^2 = 1$. It intersects the real axis at $x = \lambda(1 - \sin(\psi + \eta))$ and has the asymptotes $y = \pm(x - \lambda) \cot(\psi + \eta)$. Next we show (4.5) and (4.6). First, for $p \in \overline{\mathcal{S}}_{a,0}$, i.e., $\eta \in [0, a]$, $z(p)$ lies in the sector $\Sigma_{\pi-\psi}$. Since $\varphi := \eta + \psi \in (\psi, \pi/2 - \epsilon)$, and $\sin(\pi/2 - \epsilon) \sim 1 - \epsilon^2/2 \leq 1 - \epsilon^2/3$ for small ϵ , we have for all $\xi \in \mathbb{R}$

$$\begin{aligned} \left| \frac{z'(p)}{z(p)} \right|^2 &= \left| \frac{\cos(\varphi) \cosh(\xi) - i \sin(\varphi) \sinh(\xi)}{(1 - \cosh(\xi) \sin(\varphi)) + i \sinh(\xi) \cos(\varphi)} \right|^2 \\ &= \frac{\cosh^2(\xi) - \sin^2(\varphi)}{(\cosh(\xi) - \sin(\varphi))^2} = \frac{\cosh(\xi) + \sin(\varphi)}{\cosh(\xi) - \sin(\varphi)} \leq \frac{1 + \sin(\varphi)}{1 - \sin(\varphi)} \\ &\leq \frac{2}{1 - (1 - \epsilon^2/3)} = 6\epsilon^{-2}. \end{aligned}$$

Hence (4.5) holds true. For $p \in \overline{\mathcal{S}}_{0,b}$, i.e., $\eta \in [-b, 0]$, $z(p)$ lies in $\Sigma_{\pi-(\eta+\psi)} \subset \Sigma_{\pi-\epsilon}$. Further, since $\varphi := \eta + \psi \in (\epsilon, \psi)$, for all $\xi \in \mathbb{R}$

$$|z'(p)|^2|z(p)|^{-2} \leq (1 + \sin(\varphi))/(1 - \sin(\varphi)) \leq (1 + \sin(\psi))/(1 - \sin(\psi)).$$

Then the desired result (4.6) follows directly. \square

The next result gives analyticity of and an estimate on $\widehat{g}(\xi, t)$ on $\mathcal{S}_{a,b}$.

LEMMA 4.3. *Let $p = \xi + i\eta$ with $\xi, \eta \in \mathbb{R}$, a and b be defined as in Lemma 4.2, and $\widehat{g}(p, t)$ be defined by (4.3). Then $\widehat{g}(p, t)$ is analytic on the strip $\mathcal{S}_{a,b}$, and there holds*

$$\|\widehat{g}(p, t)\| \leq c\epsilon^{-1}e^{\lambda(1-\sin(\psi+\eta))\cosh(\xi)t}\|v_h\| \quad \forall p \in \mathcal{S}_{a,b}. \quad (4.7)$$

P r o o f. For $p = \xi + i\eta$ with $\xi, \eta \in \mathbb{R}$, $z(p)$ in (4.1) is given by

$$z(p) = \lambda(1 - \sin(\psi + \eta))\cosh(\xi) + i\lambda\cos(\psi + \eta)\sinh(\xi).$$

By Lemmas 4.2 and 2.1, $z(p)w(z(p)) \in \Sigma_{\pi-\epsilon}$, and thus the function

$$\widehat{g}(p, t) = (2\pi i)^{-1}e^{z(p)t} (z(p)w(z(p))I + A_h)^{-1}w(z(p)) z'(p)v_h$$

is analytic in $\mathcal{S}_{a,b}$. It suffices to show (4.7). For $p \in \overline{\mathcal{S}}_{0,b}$, by (4.6), $z(p) \in \Sigma_{\pi-\epsilon}$. By Lemma 2.1, $z(p)w(z(p)) \in \Sigma_{\pi-\epsilon}$. By (2.1), for small $\epsilon > 0$

$$\|(zI + A_h)^{-1}\| \leq c/|\Im(z)| \leq c/|z\sin(\pi - \epsilon)| \leq c/(|z|\epsilon) \quad \forall z \in \Sigma'_{\pi-\epsilon}.$$

For any $p \in \overline{\mathcal{S}}_{0,b}$, $\Re(z(p)) = \lambda(1 - \sin(\psi + \eta))\cosh(\xi)$. By Lemma 2.1,

$$\begin{aligned} \|\widehat{g}(p, t)\| &\leq ce^{\Re(z(p))t}|z'(p)w(z(p))| \|(z(p)w(z(p)) + A)^{-1}\| \|v_h\| \\ &\leq c\epsilon^{-1}e^{\lambda(1-\sin(\psi+\eta))\cosh(\xi)t}|z'(p)||z(p)|^{-1}\|v_h\|. \end{aligned}$$

This and (4.6) yield (4.7). The case $p \in \overline{\mathcal{S}}_{a,0}$ is direct: (4.6) and Lemma 2.1 imply $z(p)w(z(p)) \in \Sigma_\theta$. Then (4.7) follows from (4.5) and (2.1). \square

Now we can give an error estimate on the approximation $U_{N,h}$.

LEMMA 4.4. *Let $u_h(t)$ and $U_{N,h}(t)$ be defined in (4.2) and (4.4), respectively, and the contour be parametrically represented by (4.1). Then with the choice $k = c_0/N$ and $\lambda = c_1N/t$, there holds*

$$\|u_h(t) - U_{N,h}(t)\| \leq ce^{-c'N}\|v\|. \quad (4.8)$$

P r o o f. Let $u_h - U_{N,h} = (u_h - U_h) + (U_h - U_{N,h}) =: E_q + E_t$. Let $a = \pi/2 - \psi - \epsilon$ and $b = \psi - \epsilon$. For $p = \xi + ia$, $zw(z)$ lies in Σ_θ for some

$\theta \in (\pi/2, \pi)$. Since $\cosh \xi \geq 1 + \xi^2/2$ and $1 - \sin(\pi/2 - \epsilon) \leq \epsilon$ for small $\epsilon > 0$, the choice $\lambda = c_1 N/t$ and Lemma 4.3 yield

$$\begin{aligned} \left\| \int_{-\infty}^{\infty} |\widehat{g}(\xi + ia)| d\xi \right\| &\leq c\epsilon^{-1} \int_0^{\infty} e^{c_1 N(1 - \sin(\pi/2 - \epsilon) \cosh(\xi))} d\xi \|v_h\| \\ &\leq c\epsilon^{-1} e^{c_1 N\epsilon} \int_0^{\infty} e^{-c_1 N \sin(\pi/2 - \epsilon) \xi^2/2} d\xi \|v_h\| \\ &\leq c\epsilon^{-1} N^{-1/2} e^{c_1 N\epsilon} \|v_h\|. \end{aligned}$$

By Lemma 4.1, for $k = c_0/N$ we have

$$\|E_q^+\| \leq c\epsilon^{-1} N^{-\frac{1}{2}} e^{-(2\pi(\frac{\pi}{2} - \psi - \epsilon)/c_0 - \epsilon c_1)N} \|v_h\|.$$

Next we bound the error due to the lower half. For the choice $p = \xi - ib$, $\lambda = c_1 N/t$ and by the inequality $\cosh \xi \geq 1 + \xi^2/2$, we deduce

$$\begin{aligned} \left\| \int_{-\infty}^{\infty} |\widehat{g}(\xi - ib)| d\xi \right\| &\leq c\epsilon^{-1} \int_0^{\infty} e^{c_1 N(1 - \sin(\epsilon) \cosh(\xi))} d\xi \|v_h\| \\ &\leq c\epsilon^{-1} e^{c_1 N(1 - \sin(\epsilon))} \int_0^{\infty} e^{-c_1 N \sin(\epsilon) \xi^2/2} d\xi \|v_h\| \\ &\leq c\epsilon^{-3/2} N^{-1/2} e^{c_1 N(1 - \epsilon)} \|v_h\|. \end{aligned}$$

Then for the choice $k = c_0/N$, Lemma 4.1 yields the following estimate $\|E_q^-\| \leq c\epsilon^{-3/2} N^{-1/2} e^{-(2\pi(\psi - \epsilon)/c_0 - c_1(1 - \epsilon))N} \|v_h\|$. Further, since $\cosh(\xi) \geq \cosh(c_0) + \sinh(c_0)(\xi - c_0)$ for $\xi \geq c_0$, the error $\|E_t\|$ is bounded by

$$\begin{aligned} \|E_t\| &\leq c\epsilon^{-1} \int_{c_0}^{\infty} e^{c_1 N(1 - \sin(\psi) \cosh(\xi))} d\xi \|v_h\| \\ &\leq c\epsilon^{-1} e^{c_1 N(1 - \sin(\psi) \cosh(c_0))} \int_{c_0}^{\infty} e^{-c_1 N \sin(\psi) \sinh(c_0)(\xi - c_0)} d\xi \|v_h\| \\ &= c/(c_1 \sin(\psi) \sinh(c_0)\epsilon) N^{-1} e^{c_1 N[1 - \sin(\psi) \cosh(c_0)]} \|v_h\|. \end{aligned}$$

By disregarding ϵ terms, balancing the exponentials in $\|E_q^+\|$, $\|E_q^-\|$ and $\|E_t\|$, we get $2\pi(\pi/2 - \psi)/c_0 = 2\pi\psi/c_0 - c_1 = -c_1(1 - \sin(\psi) \cosh(c_0))$. Next we express c_0 and c_1 using ψ : $c_0 = \cosh^{-1}(2\pi\psi/((4\pi\psi - \pi^2) \sin \psi))$ and $c_1 = (4\pi\psi - \pi^2)/\cosh^{-1}(2\pi\psi/((4\pi\psi - \pi^2) \sin \psi))$. Finally we minimize the ratio $B(\psi) = c_1 - 2\pi\psi/c_0$ with respect to ψ , which achieves the minimum at $\psi = 1.1721$ and hence, $c_0 = 1.0818$, $c_1 = 4.4920$ and $B(\psi) = -2.32$, which are identical to the values in [32]. Then collecting the balanced asymptotic bound and $\|E_q^+\|$, $\|E_q^-\|$ and $\|E_t\|$, and choosing $\epsilon = 1/N$ yield

$$\begin{aligned} \|u_h(t) - U_{N,h}(t)\| &\leq c\epsilon^{-1} (N^{-\frac{1}{2}} + \epsilon^{-\frac{1}{2}} N^{-\frac{1}{2}} + N^{-1}) e^{-[2.32 - (\frac{2\pi}{c_0} + c_1)\epsilon]N} \|v_h\| \\ &\leq c\epsilon^{(-2.32 + \frac{\ln N}{N})N} \|v_h\|. \end{aligned}$$

Since $x^{-1} \ln x \leq 1/e$ for $x \geq 1$, the L^2 -stability of P_h yields (4.8). \square

Last, we give error estimates for the fully discrete scheme (4.4). It follows from Theorems 3.1 and 3.2, and Lemma 4.4.

THEOREM 4.1. *Let $u(t)$ be the solution of problem (1.1), and $U_{N,h}(t)$ be the approximation (4.4), with the parameters chosen as in Lemma 4.4. Then with $\ell_1(t) = (\ln 2T/t)^{-1}$, the following statements hold.*

(a) *If $Av \in L^2(\Omega)$ and $v_h = R_h v$, then*

$$\|u(t) - U_{N,h}(t)\| \leq c(e^{-c'N} + h^2)\|Av\|.$$

(b) *If $v \in L^2(\Omega)$ and $v_h = P_h v$, then*

$$\|u(t) - U_{N,h}\| \leq c_T(e^{-c'N} + h^2 t^{-1} \ell_1(t))\|v\|.$$

5. Fully Discrete Scheme II: Convolution Quadrature

Now we develop a second fully discrete scheme based on convolution quadrature. We divide the interval $[0, T]$ into a uniform grid with a time step size $\tau = T/N$, $N \in \mathbb{N}$, with $t_n = n\tau$, $n = 0, \dots, N$. Following [19, 5], we consider the convolution quadrature generated by the backward Euler method. The weights $\{b_j\}_{j=0}^\infty$ are given by [19] $\tilde{\omega}(\xi) = \sum_{j=0}^\infty b_j \xi^j = (\frac{1-\xi}{\tau})w(\frac{1-\xi}{\tau})$. Then the convolution quadrature approximation is given by: for $n = 1, 2, \dots, N$: $Q_n(\varphi) = \sum_{j=0}^n b_{n-j} \varphi(t_j)$. Then we get a fully discrete scheme for (1.1): with $U_h^0 = v_h$, for $n = 1, 2, \dots, N$

$$Q_n(U_h - v_h) + A_h U_h^n = 0. \quad (5.1)$$

We denote the generating function $\tilde{\beta}$ of $\{\beta_j\}_{j=0}^\infty$ by $\tilde{\beta}(\xi) = \sum_{j=0}^\infty \beta_j \xi^j$.

Now we analyze (5.1), following [20]. First we split the error into

$$e^n = u(t_n) - U_h^n = (u(t_n) - u_h(t_n)) + (u_h(t_n) - U_h^n).$$

By Theorems 3.1 and 3.2, it suffices to bound $\|u_h(t_n) - U_h^n\|$. We write $u_h(t_n) - U_h^n = y_h(t) - Y_h^n$, where $y_h(t) = u_h(t) - v_h$ and $Y_h^n = U_h^n - v_h$.

First, we derive representations of y_h and Y_h .

LEMMA 5.1. *Let*

$$K(z) = -z^{-1}(zw(z)I + A_h)^{-1}A_h \quad \text{and} \quad \chi(z) = (1 - e^{-z\tau})/\tau. \quad (5.2)$$

Then y_h and Y_h^n can be represented by

$$y_h(t) = \frac{1}{2\pi i} \int_{\Gamma_{\theta,\delta}} e^{zt} K(z) v_h dz \quad \text{and} \quad Y_h^n = \frac{1}{2\pi i} \int_{\Gamma_\tau} e^{zt_{n-1}} K(\chi(z)) v_h dz,$$

respectively, with the contour $\Gamma_\tau = \{z \in \Gamma_{\theta,\delta} : |\Im(z)| \leq \pi/\tau\}$.

P r o o f. By definition, y_h satisfies $D_t^{[\mu]}y_h + A_h y_h = -A_h v_h$, with $y_h(0) = 0$. The Laplace transform gives $zw(z)\widehat{y}_h(z) + A_h \widehat{y}_h(z) = -z^{-1}A_h v_h$. Hence, $\widehat{y}_h(z) = K(z)v_h$, with $K(z) = -z^{-1}(zw(z)I + A_h)^{-1}A_h$, and thus follows the representation of $y_h(t)$. Likewise, Y_h^n satisfies

$$Q_n(Y_h) + AY_h^n = -A_h v_h \quad \text{with} \quad Y_h^0 = 0.$$

Multiplying both sides by ξ^n and summing in n from 1 to ∞ yield

$$\sum_{n=1}^{\infty} Q_n(Y_h)\xi^n + A_h \widetilde{Y}_h(\xi) = -\xi/(1-\xi)A_h v_h.$$

Using the condition $Y_h^0 = 0$, we have

$$\sum_{n=1}^{\infty} Q_n(Y_h)\xi^n = \sum_{n=0}^{\infty} \sum_{j=0}^n (b_{n-j}\xi^{n-j})(Y_h^j \xi^j) = ((1-\xi)/\tau)w((1-\xi)/\tau)\widetilde{Y}_h(\xi).$$

Thus, simple calculation shows $\widetilde{Y}_h(\xi) = (\xi/\tau)K((1-\xi)/\tau)v_h$, and it is analytic at $\xi = 0$. Then Cauchy theorem implies that for small ϱ

$$Y_h^n = \frac{1}{2\tau\pi i} \int_{|\xi|=\varrho} \xi^{-n} K((1-\xi)/\tau)v_h d\xi.$$

Now, by changing variable $\xi = e^{-z\tau}$, we obtain

$$Y_h^n = \frac{1}{2\pi i} \int_{\Gamma_0} e^{zt_{n-1}} K((1-e^{-z\tau})/\tau)v_h dz,$$

where the contour $\Gamma_0 = \{z = -\ln(\varrho)/\tau + iy : |y| \leq \pi/\tau\}$ is oriented counterclockwise. The desired representation follows by deforming the contour Γ_0 to $\Gamma_\tau = \{z \in \Gamma_{\theta,\delta} : |\Im(z)| \leq \pi/\tau\}$. \square

By Lemma 5.1, we have $y_h(t_n) - Y_h^n = I + II$, where

$$\begin{aligned} I &= \frac{1}{2\pi i} \int_{\Gamma_{\theta,\delta} \setminus \Gamma_\tau} e^{zt_n} K(z)v_h dz, \\ II &= \frac{1}{2\pi i} \int_{\Gamma_\tau} e^{zt_n} (K(z) - e^{-z\tau} K(\chi(z)))v_h dz. \end{aligned} \tag{5.3}$$

Since $|e^{-z\tau}|$ is uniformly bounded on Γ_τ , we have

$$\begin{aligned} \|K(z) - e^{-z\tau} K(\chi(z))\| &\leq |e^{-z\tau}| \|K(z) - K(\chi(z))\| + |1 - e^{-z\tau}| \|K(z)\| \\ &\leq c \|K(z) - K(\chi(z))\| + c\tau |z| \|K(z)\| \\ &\leq c \|K(z) - K(\chi(z))\| + c\tau, \end{aligned}$$

where the last line, using (2.1), follows from the inequality

$$\|K(z)\| = |z|^{-1} \| -I + zw(z)(zw(z) + A_h)^{-1} \| \leq c|z|^{-1}.$$

Hence, it remains to bound the term $\|K(z) - K(\chi(z))\|$. First we recall a bound on $\chi(z) = \tau^{-1}(1 - e^{-z\tau})$ [13, Lemma 3.1].

LEMMA 5.2. Let $\chi(z) = \tau^{-1}(1 - e^{-z\tau})$. Then for all $z \in \Gamma_\tau$, there hold

$$|\chi(z) - z| \leq c|z|^2\tau \quad \text{and} \quad c_1|z| \leq |\chi(z)| \leq c_2|z|,$$

and $\chi(z)$ lies in a sector $\Sigma_{\theta'}$ for some $\theta' \in (\pi/2, \pi)$.

Next we give one estimate on the approximation $\chi(z)w(\chi(z))$ to $zw(z)$.

LEMMA 5.3. For $z \in \Gamma_\tau$, there holds

$$|\chi(z)w(\chi(z)) - zw(z)| \leq c\tau|z|^2w(|z|).$$

P r o o f. By the intermediate value theorem, for $z \in \Gamma_\tau$, we have

$$|\chi(z)^\alpha - z^\alpha| = \alpha \left| \int_z^{\chi(z)} s^{\alpha-1} ds \right| \leq \alpha |\chi(z) - z| \max_{\eta \in [0,1]} |z_\eta|^{\alpha-1},$$

where $z_\eta = \eta\chi(z) + (1 - \eta)z$ with $\eta \in [0, 1]$. Next we claim $|z_\eta|^{-1} \leq c|z|^{-1}$ for $\eta \in [0, 1]$. We split Γ_τ into $\Gamma_\tau = \Gamma_\tau^+ \cup \Gamma_\tau^c \cup \Gamma_\tau^-$, with Γ_τ^\pm being the rays in the upper and lower half plane, and Γ_τ^c is the circular arc. For $z \in \Gamma_\tau^c$, Taylor expansion gives $z_\eta = z(1 + \eta \sum_{j=1}^{\infty} (-1)^j \frac{z^j \tau^j}{(j+1)!})$. Since $|z\tau| \leq 1$ for $z \in \Gamma_\tau^c$, $|z_\eta|^{-1} \leq c|z|^{-1}$ for $z \in \Gamma_\tau^c$. It remains to consider $z \in \Gamma_\tau^+$. For $z = re^{i(\pi-\theta)}$ with $r\tau \in (\delta, \pi/\sin\theta)$ we have $\chi(z) = \frac{1}{\tau}(1 - e^{r\tau \cos\theta} e^{-ir\tau \sin\theta})$, and since $r\tau \sin\theta \leq \pi$, $\Im(\chi(z)) \geq 0$. Then Lemma 5.2 yields $|z_\eta| > \min(|z|, |\chi(z)|) \cos \frac{\theta}{2} \geq c|z|$. This shows the desired claim. Hence, by Lemma 5.2, for $z \in \Gamma_\tau$ there holds $|\int_0^1 (\chi(z)^\alpha - z^\alpha) \mu(\alpha) d\alpha| \leq \int_0^1 |\chi(z)^\alpha - z^\alpha| \mu(\alpha) d\alpha \leq c\tau|z| \int_0^1 |z|^\alpha \mu(\alpha) d\alpha = c\tau|z|^2w(|z|)$. \square

Next we give an error estimate on the approximation $K(\chi(z))$ to $K(z)$.

LEMMA 5.4. For $z \in \Gamma_\tau$, there holds $\|K(z) - K(\chi(z))\| \leq c\tau$.

P r o o f. Let $B(z) = zK(z)$. Simple computation shows

$$\begin{aligned} B(z) - B(\chi(z)) &= zw(z)((zw(z)I + A_h)^{-1} - (\chi(z)w(\chi(z))I + A_h)^{-1}) \\ &\quad + (zw(z) - \chi(z)w(\chi(z))) (\chi(z)w(\chi(z))I + A_h)^{-1} := I + II. \end{aligned}$$

First, by Lemmas 2.3 and 5.2, there holds $|\chi(z)w(\chi(z))| \geq c|\chi(z)|w(|\chi(z)|) \geq c|z|w(|z|)$. Further, by Lemma 2.1, (2.1) and Lemma 2.3, we have

$$\|(zw(z)I + A_h)^{-1}\| \leq c|zw(z)|^{-1} \leq c(|z|w(|z|))^{-1}. \quad (5.4)$$

Likewise, in view of Lemmas 5.2 and 2.1 and (2.1), we have

$$\|(\chi(z)w(\chi(z))I + A_h)^{-1}\| \leq c|\chi(z)w(\chi(z))|^{-1} \leq c(|z|w(|z|))^{-1}. \quad (5.5)$$

Now, by the identity

$$\begin{aligned} & (zw(z)I + A_h)^{-1} - (\chi(z)w(\chi(z))I + A_h)^{-1} \\ &= (zw(z) - \chi(z)w(\chi(z))) (zw(z)I + A_h)^{-1} (\chi(z)w(\chi(z))I + A_h)^{-1}, \end{aligned}$$

Lemma 2.3, (5.4) and (5.5), the first term I can be bounded by

$$\|I\| \leq c\tau|z|^3 w(|z|)^2 \|(zw(z)I + A_h)^{-1}\| \|(\chi(z)w(\chi(z))I + A_h)^{-1}\| \leq c\tau|z|.$$

Likewise, by Lemma 5.3 and (5.5)

$$\begin{aligned} \|II\| &\leq |zw(z) - \chi(z)w(\chi(z))| \|(\chi(z)w(\chi(z))I + A_h)^{-1}\| \\ &\leq c\tau|z|^2 w(|z|) |zw(z)|^{-1} \leq c\tau|z|. \end{aligned}$$

Hence, $\|B(z) - B(\chi(z))\| \leq c\tau|z|$. Last, by Lemma 5.2 and $\|B(z)\| \leq c$,

$$\begin{aligned} \|K(z) - K(\chi(z))\| &\leq |z^{-1} - \chi(z)^{-1}| \|B(z)\| + |z|^{-1} \|B(z) - B(\chi(z))\| \\ &\leq c|z - \chi(z)| |z|^{-2} + c\tau \leq c\tau. \end{aligned}$$

□

Now we can state an error estimate for nonsmooth data $v \in L^2(\Omega)$.

THEOREM 5.1. *Let u_h and U_h^n be the solutions of problems (3.1) and (5.1) with $v \in L^2(\Omega)$, $v_h = P_h v$ and $f \equiv 0$, respectively. Then there holds*

$$\|u_h(t_n) - U_h^n\| \leq c\tau t_n^{-1} \|v\|.$$

P r o o f. It suffices to bound the terms I and II in (5.3). By choosing $\delta = t_n^{-1}$ in $\Gamma_{\delta, \theta}$ and (2.1), we bound I by

$$\|I\| \leq c \int_{\frac{\pi}{\tau \sin \theta}}^{\infty} e^{rt_n \cos \theta} r^{-1} dr \|v_h\| \leq c\tau \|v_h\| \int_0^{\infty} e^{rt_n \cos \theta} dr \leq c\tau t_n^{-1} \|v_h\|.$$

Using Lemma 5.4, we deduce

$$\|II\| \leq c\tau \|v_h\| \left(\int_{t_n^{-1}}^{\frac{\pi}{\tau \sin \theta}} e^{rt_n \cos \theta} dr + \int_{-\theta}^{\theta} e^{\cos \psi} t_n^{-1} d\psi \right) \leq ct_n^{-1} \tau \|v_h\|.$$

Hence, $\|y_h(t_n) - Y_h^n\| \leq c\tau t_n^{-1} \|v_h\|$, and the estimate follows from $U_h^n - u_h(t_n) = Y_h^n - y_h(t_n)$ and the L^2 stability of P_h . □

Next we turn to smooth data $Av \in L^2(\Omega)$.

LEMMA 5.5. *Let $K^s(z) = -z^{-1}(zw(z)I + A_h)^{-1}$. Then for any $z \in \Gamma_\tau$,*

$$\|K^s(z) - K^s(\chi(z))\| \leq c\tau(|z| - 1)^{-1} \ln |z|.$$

P r o o f. Let $B^s(z) = -(zw(z)I + A_h)^{-1}$. Then by the identity $B^s(z) - B^s(\chi(z)) = (\chi(z)w(\chi(z)) - zw(z))(zw(z)I + A_h)^{-1}(\chi(z)w(\chi(z))I + A_h)^{-1}$, Lemma 5.3, and (5.4) and (5.5), we deduce $\|B^s(z) - B^s(\chi(z))\| \leq c\tau|w(z)|^{-1}$. Appealing to (5.4) again yields $\|B^s(z)\| \leq c|zw(z)|^{-1}$, and thus

$$\begin{aligned} \|K^s(z) - K^s(\chi(z))\| &\leq |z^{-1} - \chi(z)^{-1}| \|B^s(z)\| + |\chi(z)|^{-1} \|B^s(z) - B^s(\chi(z))\| \\ &\leq c|z - \chi(z)||z|^{-3}|w(z)|^{-1} + c\tau|zw(z)|^{-1} \leq c\tau|zw(z)|^{-1}. \end{aligned}$$

Then the desired result follows from Lemma 2.3. \square

THEOREM 5.2. *Let u_h and U_h^n be the solutions of problems (3.1) and (5.1) with $Av \in L^2(\Omega)$ and $v_h = R_h v$, respectively. Then with $\ell_2(t) = \ln(\max(t^{-1}, 2))$,*

$$\|u_h(t_n) - U_h^n\| \leq c\tau\ell_2(t)\|Av\|.$$

P r o o f. Let $K^s(z) = -z^{-1}(zw(z)I + A_h)^{-1}$. Then

$$\begin{aligned} y_h(t_n) - Y_h^n &= \frac{1}{2\pi i} \int_{\Gamma_{\theta, \delta} \setminus \Gamma_\tau} e^{zt_n} K^s(z) A_h v_h dz \\ &\quad + \frac{1}{2\pi i} \int_{\Gamma_\tau} e^{zt_n} (K^s(z) - e^{-z\tau} K^s(\chi(z))) A_h v_h dz := I + II. \end{aligned}$$

By Lemma 5.5, $\|K^s(z) - e^{-z\tau} K^s(\chi(z))\| \leq c\tau \ln|z|/(|z| - 1)$ for $z \in \Gamma_\tau$. With $\delta = t_n^{-1}$, the monotonicity of $f(x) = (1 - x)^{-1} \ln x$ yields

$$\begin{aligned} \|II\| &\leq c\tau \|A_h v_h\| \left(\int_{t_n^{-1}}^{\frac{\pi}{\tau \sin \theta}} e^{rt_n \cos \theta} \frac{\ln r}{r - 1} dr + \int_{-\theta}^{\theta} e^{\cos \psi} \frac{\ln(t_n^{-1})}{1 - t_n} d\psi \right) \\ &\leq c(1 - t_n)^{-1} \ln(t_n^{-1}) \tau \|A_h v_h\|. \end{aligned}$$

By (2.1), for $z \in \Gamma_{\theta, \delta}$, $\|K^s(z)\| \leq c|z|^{-1}|zw(z)|^{-1}$. Now Lemma 2.3 yields

$$\begin{aligned} \|I\| &\leq c \|A_h v_h\| \int_{\frac{\pi}{\tau \sin \theta}}^{\infty} e^{rt_n \cos \theta} r^{-2} |w(r)|^{-1} dr \\ &\leq c\tau \|A_h v_h\| \int_{t_n^{-1}}^{\infty} e^{rt_n \cos \theta} \frac{\ln r}{r - 1} dr \leq c \frac{\ln(t_n^{-1})}{1 - t_n} \tau \|A_h v_h\|. \end{aligned}$$

Now the identity $A_h R_h = P_h A$ completes the proof. \square

The next theorem gives error estimates for the scheme (5.1), which follow from Theorems 3.1, 3.2, 5.1 and 5.2.

THEOREM 5.3. *Let u and U_h^n be the solutions of problems (1.1) and (5.1) with $U_h^0 = v_h$, respectively. Then the following statements hold.*

(a) If $Av \in L^2(\Omega)$ and $v_h = R_h v$, then with $\ell_1(t) = \ln(2T/t)^{-1}$

$$\|u(t_n) - U_h^n\| \leq c(\tau \ell_2(t_n) + h^2) \|Av\|.$$

(b) If $v \in L^2(\Omega)$ and $v_h = P_h v$, then with $\ell_2(t) = \ln(\max(t^{-1}, 2))$

$$\|u(t_n) - U_h^n\| \leq c_T (\tau + h^2 \ell_1(t_n)) t_n^{-1} \|v\|.$$

6. Numerical Experiments and Discussions

Now we present numerical results to verify the convergence theory. Throughout, let the domain $\Omega = (0, 1)$ and consider

- (a) $v(x) = \sin(2\pi x) \in H^2(\Omega) \cap H_0^1(\Omega)$;
- (b) $v = \chi_{(0,1/2)} \in H^{1/2-\epsilon}(\Omega)$ with $\epsilon \in (0, 1/2)$, and χ_S the characteristic function of a set S ;
- (c) $v(x) = x^{-1/4} \in H^{1/4-\epsilon}(\Omega)$ with $\epsilon \in (0, 1/4)$.

We measure the temporal error by the L^2 errors $\|u(t_n) - U_{N,h}(t_n)\|/\|v\|$ or $\|u(t_n) - U_h^n\|/\|v\|$, and the spatial error by $\|u(t) - u_h(t)\|/\|v\|$ and $\|\nabla(u(t) - u_h(t))\|/\|v\|$. Throughout we divide the domain Ω into M equally spaced subintervals with a mesh size $h = 1/M$. Since the exact solution $u(t)$ is not unavailable, we compute a reference solution using a finer mesh.

6.1. Numerical results for the semidiscrete scheme. First we examine the convergence behavior of (1.3). Here we fix $N = 10$ in the Laplace transform approach such that the error due to time discretization is negligible. The numerical results are given in Table 1. In the table, **rate** denotes the empirical rates when the mesh size h halves, and the numbers in the bracket denote the theoretical rates. For all cases, the L^2 and H^1 norms of the error exhibit $O(h^2)$ and $O(h)$ convergence, respectively, agreeing with Theorems 3.1 and 3.2. The scheme is robust since the convergence rates hold for both smooth and nonsmooth data. The error increases as $t \rightarrow 0$, due to the weak solution singularity around $t \rightarrow 0$, cf. Theorem 2.1.

6.2. Numerical results for the fully discrete scheme I. Next we illustrate the convergence of the scheme (4.4). To make the spatial discretization error negligible, we fix h at $h = 10^{-5}$. The numerical results are summarized in Tables 2 and 3 for $\mu(\alpha) = (\alpha - 1/2)^2$ and $\mu(\alpha) = \chi_{[1/2,1]}(\alpha)$, respectively. The notation r denotes the exponent in the estimate $\|u_{N,h}^n - u(t_n)\| \leq C e^{-rN}$. The results indicate an exponential convergence with respect to N , with a rate $e^{-2.15N}$ and $e^{-2.14N}$ for $\mu(\alpha) = (\alpha - 1/2)^2$ and $\mu(\alpha) = \chi_{[1/2,1]}(\alpha)$, respectively, which agree with Theorem 4.1. Note that even though $\mu(\alpha) = \chi_{[1/2,1]}(\alpha)$ does not satisfy the condition $\mu(0)\mu(1) > 0$, the empirical rates still agree well with the theoretical one, which calls for further study. The convergence rate is independent of t , and thus the

case	t	M	10	20	40	80	160	320	rate
(a)	1	L^2	2.79e-5	7.02e-6	1.76e-6	4.39e-7	1.09e-7	2.70e-8	2.00 (2.00)
		H^1	8.84e-4	4.44e-4	2.22e-4	1.11e-4	5.23e-5	2.36e-5	1.05 (1.00)
	10^{-3}	L^2	6.38e-3	1.61e-3	4.03e-4	1.01e-4	2.51e-5	6.21e-6	2.00 (2.00)
		H^1	1.41e-1	7.04e-2	3.53e-2	1.76e-2	3.75e-3	1.65e-3	1.07 (1.00)
(b)	1	L^2	3.97e-5	9.94e-6	2.48e-6	6.21e-7	1.55e-7	3.87e-8	2.00 (2.00)
		H^1	1.26e-3	6.29e-4	3.15e-4	1.55e-4	7.63e-5	3.68e-5	1.01 (1.00)
	10^{-3}	L^2	6.34e-3	1.59e-3	3.96e-4	9.92e-5	2.48e-5	6.18e-6	2.00 (2.00)
		H^1	1.73e-1	8.65e-2	4.32e-2	2.14e-2	1.04e-2	5.06e-3	1.02 (1.00)
(c)	1	L^2	3.82e-5	9.67e-6	2.44e-6	6.12e-7	1.53e-7	3.79e-8	2.00 (2.00)
		H^1	1.21e-3	6.13e-4	3.09e-4	1.55e-4	7.33e-5	3.33e-5	1.05 (1.00)
	10^{-3}	L^2	3.48e-3	8.76e-4	2.20e-4	5.49e-5	1.37e-5	3.36e-6	2.00 (2.00)
		H^1	1.49e-1	7.45e-2	3.73e-2	1.86e-2	8.76e-3	3.97e-3	1.07 (1.00)

TABLE 1. Numerical results for the semidiscrete scheme (1.3) with $\mu(\alpha) = (\alpha - 1/2)^2$.

case	$t \setminus N$	3	5	7	9	11	13	r
(a)	1	1.33e-6	1.49e-8	1.26e-10	2.20e-12	3.54e-14	8.24e-17	2.35
	10^{-2}	4.78e-6	7.36e-7	2.77e-9	5.45e-11	4.88e-13	2.23e-14	1.92
	10^{-3}	8.30e-5	8.78e-7	3.81e-9	7.55e-11	6.43e-13	1.23e-14	2.26
(b)	1	3.34e-6	3.56e-8	2.85e-10	5.76e-12	8.68e-14	1.25e-15	2.17
	10^{-2}	1.24e-5	8.29e-7	2.31e-9	6.09e-11	4.78e-13	2.18e-14	2.02
	10^{-3}	6.99e-5	1.73e-6	1.09e-8	5.38e-11	1.17e-12	1.59e-14	2.22
(c)	1	8.04e-6	9.05e-8	6.80e-10	1.39e-11	2.08e-13	3.02e-15	2.17
	10^{-2}	3.01e-5	1.71e-6	3.85e-9	1.26e-10	9.22e-13	4.21e-14	2.04
	10^{-3}	1.16e-4	4.09e-6	2.65e-8	6.65e-11	2.75e-12	3.49e-14	2.19

TABLE 2. The L^2 errors for (a)-(c) with $h = 10^{-5}$ and $\mu(\alpha) = (\alpha - 1/2)^2$, by the Laplace transform method.

scheme is robust. Interestingly, the smoothness of the initial data does not affect much the time discretization errors, even for small time, cf. Table 4.

One salient feature of the fully discrete scheme I is that it allows computing the solution at large time. This allows one to examine the asymptotic behavior of the solution as the time $t \rightarrow \infty$; see Table 5 and Fig. 1. In particular, one clearly observes the logarithmic decay of the solution [18, Theorem 2.1]; see also Fig. 1. This numerically verifies the ultraslow decay asymptotics for distributed order diffusion process.

case	$t \setminus N$	3	5	7	9	11	13	r
(a)	1	4.54e-6	2.30e-7	1.63e-9	1.69e-11	2.36e-13	8.46e-15	2.02
	10^{-2}	6.21e-5	1.65e-6	3.71e-9	1.07e-10	7.00e-13	2.58e-14	2.16
	10^{-3}	8.02e-4	3.61e-6	1.66e-8	4.17e-10	3.10e-12	6.73e-15	2.55
(b)	1	4.78e-6	4.74e-7	2.43e-9	3.44e-11	3.49e-13	1.87e-14	1.94
	10^{-2}	1.03e-4	1.13e-6	3.58e-9	8.78e-11	5.04e-13	1.93e-14	2.24
	10^{-3}	5.12e-4	4.79e-6	4.95e-8	5.23e-10	5.15e-12	5.58e-14	2.29
(c)	1	4.79e-6	5.61e-7	2.75e-9	4.07e-11	3.94e-13	2.23e-14	1.92
	10^{-2}	1.18e-4	6.08e-7	3.37e-9	7.22e-11	2.84e-13	8.94e-14	2.10
	10^{-3}	1.09e-4	5.24e-6	6.02e-8	5.62e-10	5.95e-12	1.02e-13	2.07

TABLE 3. The L^2 errors for (a)-(c) with $h = 10^{-5}$ and $\mu(\alpha) = \chi_{[1/2,1]}(\alpha)$, by the Laplace transform method.

case \ t	10^{-4}	10^{-5}	10^{-6}	10^{-7}	10^{-8}	10^{-9}
(b)	7.05e-6	9.39e-6	1.58e-5	1.75e-5	1.81e-5	1.82e-5
(c)	6.39e-6	1.17e-5	1.53e-5	1.68e-5	1.75e-5	1.79e-5

TABLE 4. The L^2 errors for (b) and (c) with $h = 10^{-5}$, $\mu(\alpha) = (\alpha - 1/2)^2$ and $N = 5$ at time $t = 10^{-k}$, $k = 4, \dots, 9$, by the Laplace transform method.

case \ k	6	8	10	12	14	16	18	rate
(a)	3.33e-4	2.70e-4	2.26e-4	1.95e-4	1.71e-4	1.52e-4	1.37e-4	$1/k$
(c)	1.06e-3	8.54e-4	7.17e-4	6.17e-4	5.41e-4	4.82e-4	4.34e-4	$1/k$

TABLE 5. The L^2 norm of the solution for (a) and (c) with $h = 10^{-5}$, $\mu(\alpha) = (\alpha - 1/2)^2$ and $N = 10$ at time $t = 10^k$, $k = 6, 8, \dots, 18$, by the Laplace transform method.

6.3. Numerical results for the fully discrete scheme II. Last we verify the convergence of the fully discrete scheme II, i.e., (5.1). The results are shown in Tables 6 and 7 for $\mu(\alpha) = (\alpha - 1/2)^2$ and $\mu(\alpha) = \chi_{[1/2,1]}(\alpha)$, respectively. An $O(\tau)$ convergence is always observed, cf. Theorem 5.3.

To examine more closely its convergence behavior, we consider $t = 10^{-k}$, $k = 4, \dots, 9$, and at each time t , divide the interval $[0, t]$ into $N = 10$ subintervals. The scheme works well for the smooth initial data, however, it works poorly for the singular initial data, cf. Table 8. This behavior is predicted by Theorem 5.1: the error is dominated by the factor τ/t for L^2 initial data. In Fig. 2, we plot the ratio $\|U_h^1 - u(\tau)\|/\tau$ versus $\ln \tau$ for

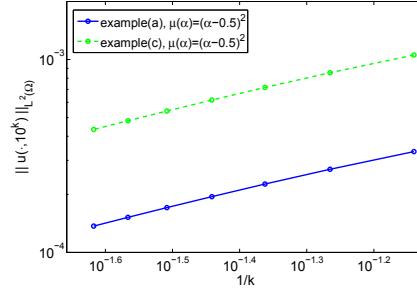


FIGURE 1. The L^2 norm of the solution for (a) and (c) at $t = 10^k$, $k = 6, 8, \dots, 18$, by the Laplace transform method.

case	$t \setminus N$	10	20	40	80	160	320	rate
(a)	1	1.82e-5	8.78e-6	4.31e-6	2.12e-6	1.01e-6	4.74e-7	1.05 (1.00)
	10^{-2}	8.64e-4	3.91e-4	1.88e-4	9.20e-5	4.55e-5	2.26e-5	1.05 (1.00)
	10^{-3}	2.17e-2	1.10e-2	5.51e-3	2.76e-3	1.38e-3	6.92e-4	0.99 (1.00)
(b)	1	4.81e-5	2.32e-5	1.14e-5	5.60e-6	2.67e-6	1.26e-6	1.05 (1.00)
	10^{-2}	8.11e-3	3.87e-3	1.88e-3	9.29e-4	4.61e-4	2.30e-4	1.03 (1.00)
	10^{-3}	1.48e-2	7.46e-3	3.74e-3	1.88e-3	9.39e-4	4.70e-4	1.00 (1.00)
(c)	1	5.81e-5	2.81e-5	1.38e-5	6.76e-6	3.23e-6	1.52e-6	1.05 (1.00)
	10^{-2}	1.01e-2	4.80e-3	2.34e-3	1.15e-3	5.72e-4	2.85e-4	1.03 (1.00)
	10^{-3}	7.35e-3	3.66e-3	1.82e-3	9.11e-4	4.55e-4	2.27e-4	1.00 (1.00)

TABLE 6. The L^2 errors for (a)-(c) with $h = 10^{-4}$ and $\mu(\alpha) = (\alpha - 1/2)^2$, by convolution quadrature.

case	$t \setminus N$	10	20	40	80	160	320	rate
(a)	1	2.20e-4	1.06e-4	5.20e-5	2.58e-5	1.28e-5	6.40e-6	1.02 (1.00)
	10^{-2}	1.76e-2	8.81e-3	4.40e-3	2.20e-3	1.10e-3	5.49e-4	1.00 (1.00)
	10^{-3}	3.92e-3	1.98e-3	9.95e-4	4.99e-4	2.50e-4	1.25e-4	0.99 (1.00)
(b)	1	6.52e-4	3.11e-4	1.52e-4	7.53e-5	3.74e-5	1.87e-5	1.03 (1.00)
	10^{-2}	1.25e-2	6.26e-3	3.13e-3	1.56e-3	7.82e-4	3.91e-4	1.00 (1.00)
	10^{-3}	5.76e-3	2.88e-3	1.44e-3	7.18e-4	3.59e-4	1.79e-4	1.00 (1.00)
(c)	1	7.92e-4	3.78e-4	1.85e-4	9.14e-5	4.54e-5	2.27e-5	1.03 (1.00)
	10^{-2}	7.40e-3	3.71e-3	1.86e-3	9.28e-4	4.64e-4	2.32e-4	1.00 (1.00)
	10^{-3}	6.10e-3	3.06e-3	1.53e-3	7.65e-4	3.83e-4	1.91e-4	1.00 (1.00)

TABLE 7. The L^2 errors for (a)-(c) with $h = 10^{-4}$ and $\mu(\alpha) = \chi_{[1/2,1]}(\alpha)$, by convolution quadrature.

smooth initial data. Theorem 5.2 predicts $\|U_h^1 - u(\tau)\| \leq c\tau \ln \tau^{-1}$. The log factor $\ell_2(t)$ in Theorem 5.2 is confirmed by Fig. 2, and thus it is sharp.

case \ t	10^{-4}	10^{-5}	10^{-6}	10^{-7}	10^{-8}	10^{-9}	rate
(a)	2.42e-3	1.03e-4	7.87e-6	7.59e-7	7.58e-8	7.44e-9	1.01 (1.00)
(c)	7.44e-3	5.67e-3	4.30e-3	3.27e-3	2.49e-3	1.88e-3	0.12 (0.12)

TABLE 8. The L^2 errors with $h = 10^{-5}$ and $N = 10$, at $t = 10^{-k}$, $k = 4, \dots, 9$, by convolution quadrature.

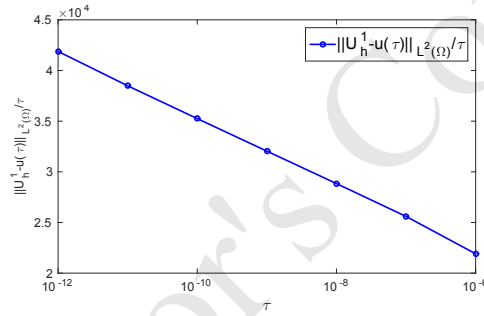


FIGURE 2. The L^2 errors for (a) at $t_1 = \tau = 10^{-k}$, $k = 5, \dots, 12$, by convolution quadrature.

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References

- [1] E. Bazhlekova, Completely monotone functions and some classes of fractional evolution equations. *Integral Transforms Spec. Funct.* **26**, No 9 (2015), 737–752; DOI:10.1080/10652469.2015.1039224.
- [2] E. Bazhlekova, B. Jin, R. Lazarov, Z. Zhou, An analysis of the Rayleigh-Stokes problem for the generalized second grade fluid. *Numer. Math.* **131**, No 1 (2015), 1–31; DOI:10.1007/s00211-014-0685-2.
- [3] A. V. Chechkin, R. Gorenflo, I. M. Sokolov, Retarding subdiffusion and accelerating superdiffusion governed by distributed-order fractional diffusion equations. *Phys. Rev. E* **66** (2002), 046129; DOI:10.1103/PhysRevE.66.046129

- [4] A. V. Chechkin, R. Gorenflo, I. M. Sokolov, V. Y. Gonchar, Distributed order time fractional diffusion equation. *Fract. Calc. Appl. Anal.* **6**, No 3 (2003), 259–279.
- [5] E. Cuesta, C. Lubich, C. Palencia, Convolution quadrature time discretization of fractional diffusion-wave equations. *Math. Comp.* **75**, No 254 (2006), 673–696; DOI: 10.1090/S0025-5718-06-01788-1.
- [6] K. Diethelm, N. J. Ford, Numerical analysis for distributed-order differential equations. *J. Comput. Appl. Math.* **225**, No 1 (2009), 96–104; DOI: 10.1016/j.cam.2008.07.018.
- [7] N. J. Ford, M. L. Morgado, M. Rebelo, An implicit finite difference approximation for the solution of the diffusion equation with distributed order in time. *Electron. Trans. Numer. Anal.* **44** (2015), 289–305.
- [8] H. Fujita, T. Suzuki, Evolution problems. In: *Handbook of Numerical Analysis*, Vol. II, North-Holland, Amsterdam (1991), 789–928.
- [9] G. Gao, H. Sun, Z. Sun, Some high-order difference schemes for the distributed-order differential equations. *J. Comput. Phys.* **298**, No 1 (2015), 337–359; DOI: 10.1016/j.jcp.2015.05.047.
- [10] M. Hasse, *The Functional Calculus for Sectorial Operators*. Birkhäuser, Basel (2006).
- [11] B. Jin, R. Lazarov, Y. Liu, Z. Zhou, The Galerkin finite element method for a multi-term time-fractional diffusion equation. *J. Comput. Phys.* **281** (2015), 825–843; DOI: 10.1016/j.jcp.2014.10.051.
- [12] B. Jin, R. Lazarov, Z. Zhou, Error estimates for a semidiscrete finite element method for fractional order parabolic equations. *SIAM J. Numer. Anal.* **51**, No 1 (2013), 445–466; DOI:10.1137/120873984
- [13] B. Jin, R. Lazarov, Z. Zhou, An analysis of the L1 scheme for the subdiffusion equation with nonsmooth data. *IMA J. Numer. Anal.* **36**, No 1 (2016), 197–221; DOI: 10.1093/imanum/dru063.
- [14] B. Jin, W. Rundell, A tutorial on inverse problems for anomalous diffusion processes. *Inverse Problems* **31**, No 3 (2015), 035003; DOI: 10.1088/0266-5611/31/3/035003.
- [15] J. T. Katsikadelis. Numerical solution of distributed order fractional differential equations. *J. Comput. Phys.* **259** (2014), 11–22; DOI: 10.1016/j.jcp.2013.11.013.
- [16] A. Kilbas, H. Srivastava, J. Trujillo, *Theory and Applications of Fractional Differential Equations*. Elsevier, Amsterdam (2006).
- [17] A. N. Kochubei, Distributed order calculus and equations of ultra-slow diffusion. *J. Math. Anal. Appl.* **340**, No 1 (2008), 252–281; DOI: 10.1016/j.jmaa.2007.08.024.

- [18] Z. Li, Y. Luchko, M. Yamamoto, Asymptotic estimates of solutions to initial-boundary-value problems for distributed order time-fractional diffusion equations. *Fract. Calc. Appl. Anal.* **17**, No 4 (2014), 1114–1136; DOI: 10.2478/s13540-014-0217-x; <http://www.degruyter.com/view/j/fca.2014.17.issue-4/issue-files/fca.2014.17.issue-4.xml>.
- [19] C. Lubich, Convolution quadrature and discretized operational calculus, I. *Numer. Math.* **52**, No 2 (1988), 129–145; DOI: 10.1007/BF01398686.
- [20] C. Lubich, I. H. Sloan, V. Thomée, Nonsmooth data error estimates for approximations of an evolution equation with a positive-type memory term. *Math. Comp.* **65**, No 213 (1996), 1–17; DOI: 10.1090/S0025-5718-96-00677-1.
- [21] Y. Luchko, Boundary value problems for the generalized time fractional diffusion equation of distributed order. *Fract. Calc. Appl. Anal.*, **12**, No 4 (2009), 409–422; at <http://www.math.bas.bg/~fcaa>.
- [22] F. Mainardi, A. Mura, G. Pagnini, R. Gorenflo, Time-fractional diffusion of distributed order. *J. Vibr. Control* **14**, No 9–10 (2008), 1267–1290; DOI: 10.1177/1077546307087452.
- [23] E. Martensen, Zur numerischen Auswertung uneigenlicher Integrale. *Z. Angew. Math. Mech.* **48** (1968), T83–T85.
- [24] W. McLean, V. Thomée, Numerical solution via Laplace transforms of a fractional order evolution equation. *J. Integral Equations Appl.* **22**, No 1 (2010), 57–94; DOI: 10.1216/JIE-2010-22-1-57.
- [25] M. M. Meerschaert, H.-P. Scheffler, Stochastic model for ultraslow diffusion. *Stochastic Process. Appl.* **116**, No 9 (2006), 1215–1235; DOI: 10.1016/j.spa.2006.01.006.
- [26] M. L. Morgado, M. Rebelo, Numerical approximation of distributed order reaction diffusion equations. *J. Comput. Appl. Math.* **275** (2015), 216–227; DOI: 10.1016/j.cam.2014.07.029.
- [27] D. Sheen, I. Sloan, V. Thomée, A parallel method for time-discretization of parabolic problems based on contour integral representation and quadrature. *Math. Comp.* **69**, No 229 (2000), 177–195; DOI: 10.1090/S0025-5718-99-01098-4.
- [28] D. Sheen, I. Sloan, V. Thomée, A parallel method for time discretization of parabolic equations based on Laplace transformation and quadrature. *IMA J. Numer. Anal.* **23**, No 2 (2003), 269–299; DOI: 10.1093/imanum/23.2.269.
- [29] I. M. Sokolov, A. V. Chechkin, J. Klafter, Distributed-order fractional kinetics. *Acta Phys. Polon. B* **35**, No 4 (2004), 1323–1341.

- [30] V. Thomée, *Galerkin Finite Element Methods for Parabolic Problems*. Springer-Verlag, Berlin (2006).
- [31] S. Umarov, Continuous time random walk models associated with distributed order diffusion equations. *Fract. Calc. Appl. Anal.* **18**, No 3 (2015), 821–837; DOI: 10.1515/fca-2015-0049; <http://www.degruyter.com/view/j/fca.2015.18.issue-3/issue-files/fca.2015.18.issue-3.xml>.
- [32] J. A. C. Weideman, L. N. Trefethen, Parabolic and hyperbolic contours for computing the Bromwich integral. *Math. Comp.* **76**, No 259 (2007), 1341–1356; DOI: 10.1090/S0025-5718-07-01945-X.

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