

## THE CURIOUS CASE OF HD 41248. A PAIR OF STATIC SIGNALS BURIED BEHIND RED NOISE

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## ABSTRACT

Gaining a better understanding of the effects of stellar-induced radial velocity noise is critical for the future of exoplanet studies since the discovery of the lowest-mass planets using this method will require us to go below the intrinsic stellar noise limit. An interesting test case in this respect is that of the southern solar analog HD 41248. The radial velocity time series of this star has been proposed to contain either a pair of signals with periods of around 18 and 25 days, which could be due to a pair of resonant super-Earths, or a single and varying 25 day signal that could arise due to a complex interplay between differential rotation and modulated activity. In this work, we build up more evidence for the former scenario, showing that the signals are still clearly significant, even after more than 10 yr of observations, and they likely do not change in period, amplitude, or phase as a function of time, the hallmarks of static Doppler signals. We show that over the last two observing seasons, this star was more intrinsically active and the noise reddened, highlighting why better noise models are needed to find the lowest amplitude signals, in particular, models that consider noise correlations. This analysis shows that there is still sufficient evidence for the existence of two super-Earths on the edge of, or locked into, a 7:5 mean motion resonance orbiting HD 41248.

*Key words:* planetary systems – stars: fundamental parameters – stars: individual (HD41248) – stars: rotation

*Online-only material:* color figures, machine-readable table

## 1. INTRODUCTION

The discovery of low-mass planets in the super-Earth regime using the radial velocity method is at the forefront of modern exoplanet science as it pushes the boundaries of what is possible using current technology (Pepe et al. 2011; Tuomi et al. 2014b; Anglada-Escudé et al. 2013). However, the Doppler signals imposed on the host stars of such orbiting bodies can also fight for dominance with signals induced in the data by rotationally modulated activity features like star spots (see Boisse et al. 2011).

A radial velocity analysis of the star HD 166435 by Queloz et al. (2001) found a repeating short-period signal of less than four days, suggesting the presence of a planetary companion to the star. After photometric follow-up, they found a period matching the period of the radial velocity signal, indicating the star was actually active and the signal they had detected was due to rotationally modulated star spots. This was confirmed when they found that the coherence time of the radial velocity signal was only  $\sim 30$  days and correlations were found with the bisector inverse slope (BIS), meaning it was not a static signal as expected of a genuine Doppler velocity profile.

GJ581 provides another example of false-positive radial velocity signals where a possible candidate planet (GJ581 d) was reported in Udry et al. (2007) with a period of 82 days, later shown to be the one-year alias of the real planet candidate period of 67 days (Mayor et al. 2009). The existence of the habitable zone super-Earth GJ581 g (Vogt et al. 2010) has also been disputed (Tuomi 2012; Baluev 2013; Hatzes 2013a), later countered by Vogt et al. (2012), as has the existence of the Earth-mass planet reported to be orbiting Alpha Cen B (Dumusque et al. 2012; Hatzes 2013b). Clearly, the detection of low-mass planets approaching the intrinsic noise level of the star and instrument combination is fraught with difficulty.

Jenkins et al. (2013b) announced the discovery of a pair of planetary candidates orbiting the star HD 41248 in, or close to, a 7:5 mean motion resonance (MMR) configuration. Both signals reported in their work were statistically significant, even when considering correlations between the activity indicators and the radial velocities. However, although the time baseline was long, around 7.5 yr, they only had a total of 62 Doppler velocities, yet the MMR configuration (period ratio of  $1.400 \pm 0.002$ ) seemed to favor a planetary hypothesis as such a pair of periods so close to a 7:5 integer ratio seems difficult to attribute to the star. The metal-poor nature of HD 41248 ( $[\text{Fe}/\text{H}] = -0.43$  dex) also agrees with the emerging notion that metal-poor stars have a higher fraction of the lowest-mass planets (Jenkins et al. 2013a).

Recently, Santos et al. (2014) have claimed that the longer-period signal in the HD 41248 radial velocity data is likely due to rotationally modulated magnetic activity; after adding more than 160 new velocities and when subtracting off that signal, there is no remaining evidence for the 18 day signal. With this in mind, we decided to reanalyze all data for HD 41248 and test whether the pair of signals still remain in the new data from Santos et al. Moreover, we discuss whether these signals could still be interpreted as being due to a pair of planets.

## 2. HD41248 STATISTICAL MODEL

We modeled the HARPS radial velocities of HD 42148 by adopting the analysis techniques and the statistical model applied in Tuomi et al. (2014a). This model contains Keplerian signals, a linear trend, a moving average component with exponential smoothing, and linear correlations with activity indices, namely, BIS, FWHM, and chromospheric activity  $S$  index. According to Tuomi et al., such a model can filter out activity-related variations in radial velocities and even suppress the velocity variations caused by the co-rotation of star spots on the stellar surface below the detection threshold, enabling

**Table 1**  
HARPS Time Series Data for HD 41248

BJD	RV (ms <sup>-1</sup> )	S Index (dex)	BIS (ms <sup>-1</sup> )	FWHM (ms <sup>-1</sup> )
2452943.8528426	3526.59 ± 2.59	0.169	35.93	6721.78
2452989.7102293	3519.14 ± 4.06	0.170	27.40	6719.01
2452998.6898180	3526.43 ± 5.43	0.179	33.53	6701.21
2453007.6786518	3526.63 ± 2.53	0.162	28.61	6718.20
2453787.6079555	3522.44 ± 2.76	0.162	31.31	6718.54
2454055.8375443	3523.18 ± 2.06	0.168	23.95	6714.52
2454789.7207967	3522.99 ± 0.82	0.171	27.43	6722.19
2454790.6943362	3519.49 ± 0.90	0.170	30.83	6724.20
2454791.7055725	3522.47 ± 0.83	0.171	29.54	6720.60
2454792.7042506	3522.29 ± 0.80	0.172	28.09	6728.65

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

the detection of low-amplitude variations of planetary origin, as witnessed on CoRoT-7 (Tuomi et al.). We write the statistical model as

$$m_{i,l} = \gamma_l + \dot{\gamma} t_i + f_k(t_i) + \epsilon_{i,l} + \sum_{j=1}^q c_{j,l} \xi_{j,i,l} + \sum_{j=1}^p \phi_{j,l} \exp\left\{\frac{t_i - j - t_i}{\tau_l}\right\} \epsilon_{i,l}, \quad (1)$$

where  $m_{i,l}$  is the measurement made at time  $t_i$  and the index  $l$  denotes that it corresponds to an independent  $l$ th data set, parameter  $\gamma_l$  is the reference velocity, function  $f_k$  denotes the superposition of  $k$  Keplerian signals,  $\epsilon_{i,l}$  is a Gaussian white noise with zero mean and a variance of  $\sigma_i^2 + \sigma_l^2$ , where  $\sigma_i$  is the estimated instrument uncertainty corresponding to the radial velocity measurement  $m_{i,l}$  and  $\sigma_l$  quantifies the excess white noise in the  $l$ th data set, parameters  $c_{j,l}$  describe the linear correlations with the activity indices  $\xi_{j,i,l}$ , for  $j = 1, \dots, q$ , and parameters  $\phi_{j,l}$  quantify the moving average components,  $j = 1, \dots, p$ , with exponential smoothing in a timescale of  $\tau_l$ . In practice, we apply a first-order moving average model (MA(1)) as we believe it is a sufficiently accurate description of this data, parameterized by setting the moving average components ( $p$ ) in Equation (1) equal to unity.

The prior probability densities have to be defined in order to use the techniques of Tuomi et al. (2014a) relying on the Bayes' rule of conditional probabilities. We define these densities according to Tuomi & Anglada-Escudé (2013) by choosing uninformative and uniform densities for all but two model parameters, namely, the eccentricities ( $e$ ) and excess jitter ( $\sigma_j$ ). These are set such that low eccentricities and low jitters are preferred but that higher values are not ruled out a priori (Tuomi et al. 2014a).

As the noise caused by inhomogeneities of the stellar surface and activity cannot be expected to be time-invariant over the baseline of the observations of over 10 yr, we model the velocities already analyzed in Jenkins et al. (2013b) and the new ones obtained during the last 2 yr as independent data sets. In this way, we can account for the possibility that the noise properties have changed over the data baseline and the potential effects that the data sampling, which is more dense during the last two years, has on the parameters of the noise model. Finally, we also split the last two observing seasons up into two subsets of data, and although this is detrimental to the information content,

this was done to directly compare our results with those recently published in Santos et al. (2014).

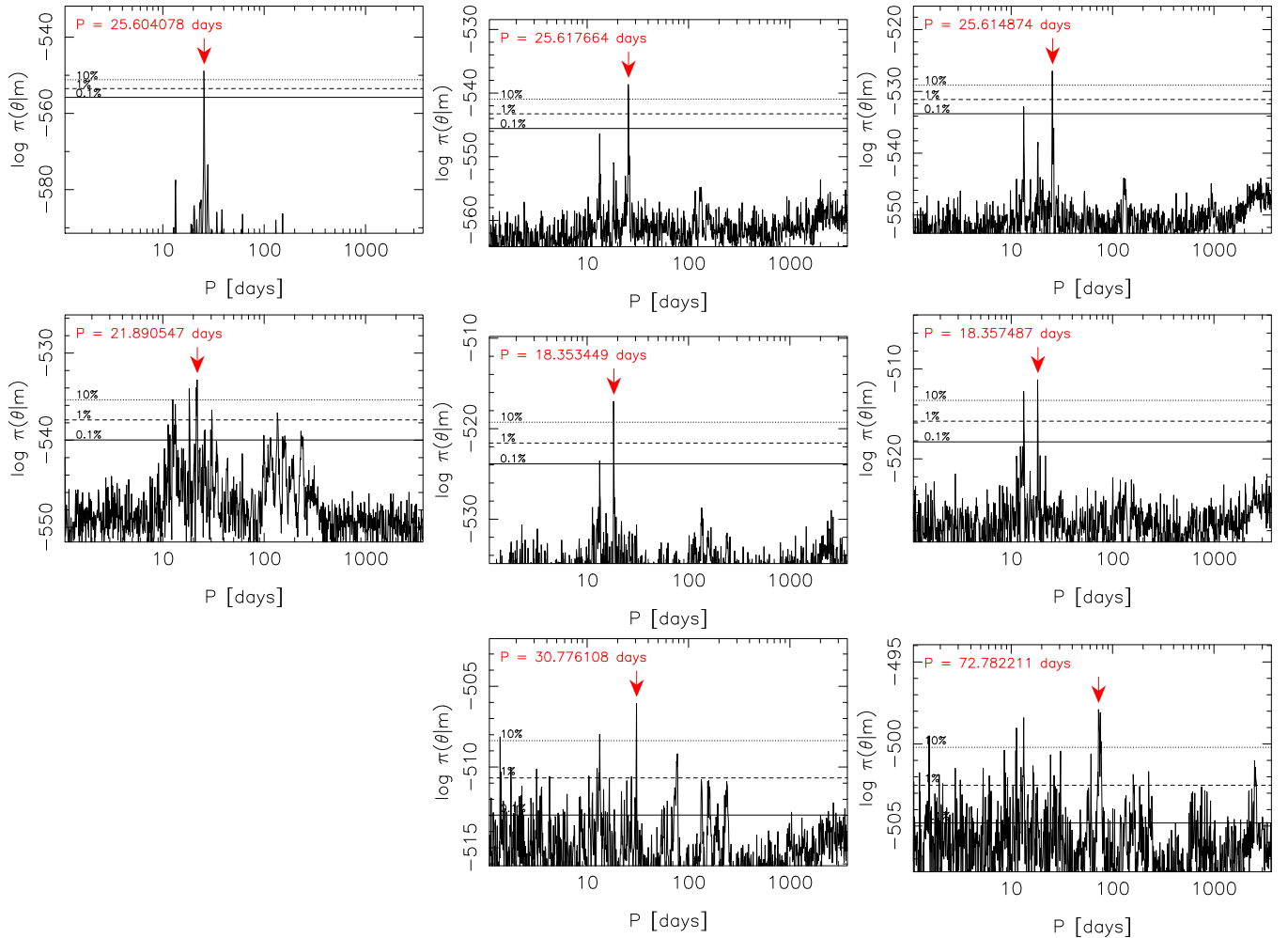
### 3. HD41248 REANALYSIS

We applied our statistical model outlined above to the full data set of radial velocities for HD 41248, combining the previously published data in Jenkins et al. (2013b) with the newly published data in Santos et al. (2014), giving rise to a total time series of 223 HARPS (Mayor et al. 2003) velocities.<sup>4</sup> We applied both tests with and without the linear activity correlation terms and also compared to the white-noise model applied in Santos et al. Table 1 contains the measured radial velocities, BIS, and FWHM values from the HARPS CCF, and the chromospheric  $S$  indices that were measured following the procedures in Jenkins et al. (2006, 2008, 2011).

In the top panels of Figure 1, we show the posterior probability densities as functions of period for tempered Markov chain samplings employing three different  $k = 1$  signal (Keplerian) models. We can see that there appears to be three regions in the period space where the Markov chains identified considerable maxima for the pair of one-component moving average (MA(1)) models (middle and right columns) and only one region for the white-noise model (left column). The most significant of these, i.e., the global maximum, was found to be at 25.6 days in all models. The existence of the local maxima means we are likely to find other significant periodic signals in the data.

The panels in the middle row in Figure 1 show what we find when applying the  $k = 2$  signal models, the first one being at a period of 25.6 days. This time the most significant maximum is found to be at 18.35 days when dealing with the red noise and also including the correlation terms, in excellent agreement with the pair of signals published previously in Jenkins et al. (2013b). Therefore, we confirm that there are two significant frequencies in the extended time series data for HD 41248. The log-Bayesian evidences for these model comparisons can be found in Table 2 listed as Full Data Set. It can be seen that the white-noise model search (left panel) did not find a unique second periodicity since it is masked by the correlated noise, variability related to activity, and increased jitter. Santos et al. (2014) could not confirm this second signal in their data, and it is likely that this was due to inadequate noise modeling as they assumed a white-noise model, fixed the excess white noise to an value of

<sup>4</sup> The data were obtained from the European Southern Observatory archive under request number JENKINS-110394.



**Figure 1.** Estimated posterior probability densities based on tempered Markov chain samplings as functions of signal period. From top to bottom we show the  $k = 1, 2,$  and 3 planet models. The maximum a posteriori estimates identified by the chains are highlighted on the plots, as are the 0.1%, 1%, and 10% equi-probability thresholds with respect to the maxima. The left column shows the pure white-noise model, the middle column is for an MA1 red-noise model without activity correlation terms included, and the right column is for an MA(1) red-noise model that includes the activity correlations. Note that there is no  $k = 3$  planet model for the white-noise model analyses (left column).

(A color version of this figure is available in the online journal.)

**Table 2**

Log-Bayes Factors  $\ln B_{1,0}$  and  $\ln B_{2,1}$ , i.e., in Favor of One Against Zero and in Favor of Two Against One-Keplerian Models, Given Various Divisions of the Data and Different Models

Data/Model	1	2	Periods (day)
Old Data	18.4	8.8	25, 18
New Data	13.8	...	25
Full Data	21.8	15.2	25, 18
Full Data with Activity Correlations	22.8	16.7	25, 18

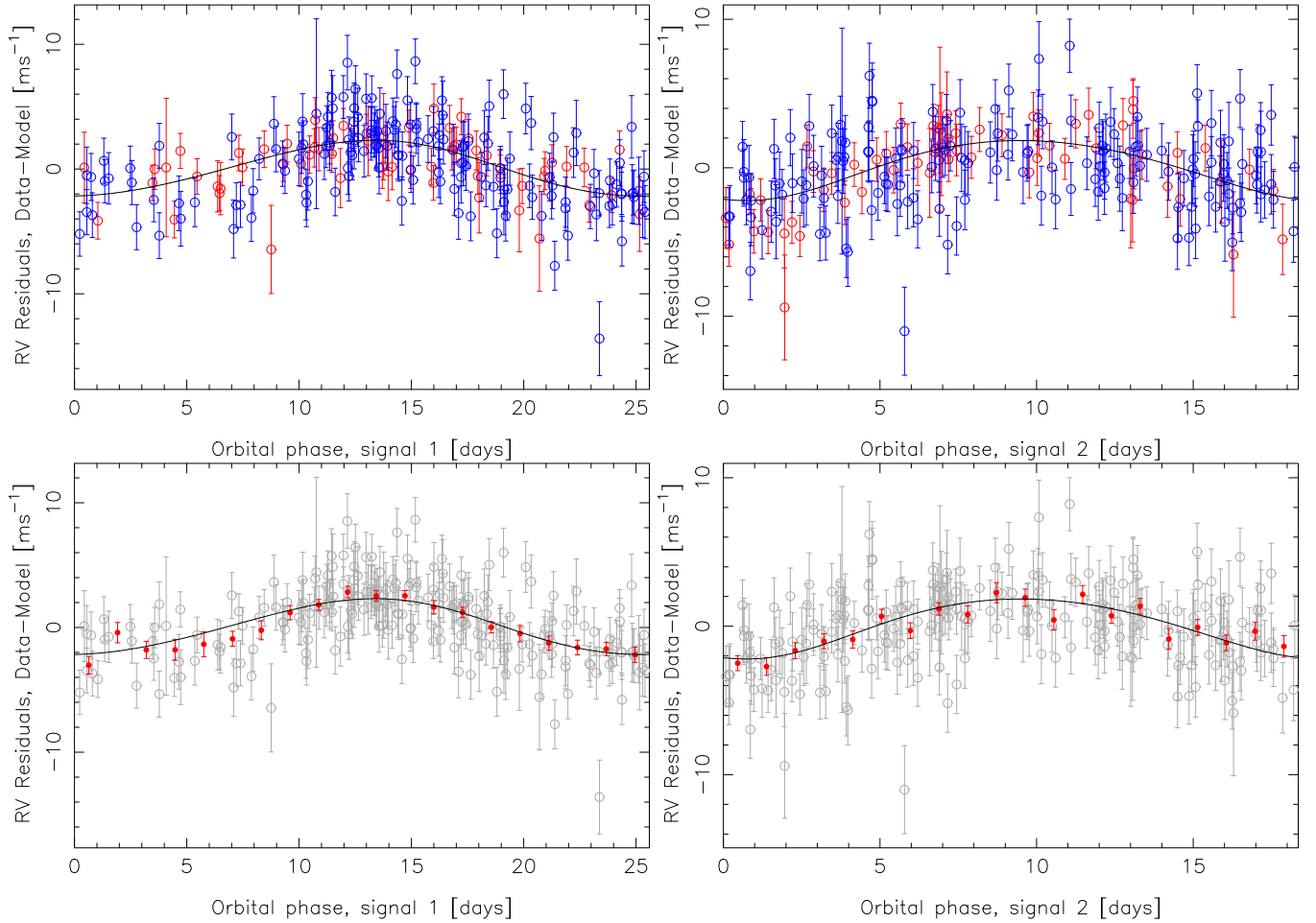
**Notes.** The last column denotes the periods of the significantly detected signals. The models contain a moving average component. Two alternative two-Keplerian solutions are shown for the full data set.

$0.7 \text{ ms}^{-1}$ , and relied on analyses of model residuals that cause severe biases to the obtained solutions (Tuomi 2012; Tuomi & Anglada-Escudé 2013). We also note that the parameter density widths for both signals decrease significantly by including this new data, a feature not expected from a quasi-static source. Furthermore, the linear trend applied to the data in Santos et al.

(2014) is not significant, agreeing well with zero within any reasonable confidence level (Table 2).

When employing the  $k = 2$  models, a second significant peak around 13 days was found to cross the 10% probability threshold in the analysis with activity correlation terms included, indicating a third signal could be present in the combined data. The probability thresholds are scaled to the global maximum of the posterior probability space, whereby the maximum is set to 100% and the thresholds mark 10%, 1%, and 0.1% of this maximum. Santos et al. (2014) also detected this signal and attributed it to the first harmonic of the 25 day signal; however, our parameter densities suggest otherwise as the distribution did not overlap with one-half of the period of the 25 day signal. In any case, we then applied the  $k = 3$  models to test if this was indeed the case, and we show the posterior diagrams in the bottom panels of Figure 1. No additional signals that were unique and passed our signal detection criteria were found in the analysis. We did not employ the  $k = 3$  model to the white-noise model as we could not constrain any secondary signal under that assumption.

Given that no 13 day signal is found in the full data when including the noise correlations, and since the signal in the



**Figure 2.** Phase-folded radial velocities for both signals detected in the HD 41248 radial velocities as a function of time. The red points are for old low-cadence data, and the blue points are for new high-cadence data in the top panels. The lower panels show the same data but with filled red points representing binned velocities to highlight the significance of the signals.

(A color version of this figure is available in the online journal.)

later data that we detect does not pass our planetary signal selection criteria, which are (1) the model including the signal must be 10,000 times more probable than the previous model, i.e.,  $P(\mathcal{M}_{k+1}) \gg P(\mathcal{M}_k)$  and (2) the signal must not vary in time in period, phase, and amplitude over the baseline of the observations (see Tuomi et al. 2014b), it cannot be considered a static Doppler signal. This result is a reproduction of the same result found for the CoRoT-7 radial velocity time series (Tuomi et al. 2014a) where the rotation period of that star, known from the previously measured CoRoT photometry, did not correspond to a genuine velocity signal. This strongly indicates that the 13 day signal reflects a quasi-static nature that changes as a function of time and could be the actual rotational period of HD 41248 or a mix of the signal from the rotational period and differential rotation. The final system parameters are listed in Table 3, and the log-Bayesian evidence ratios for these tests are shown in Table 2. The phase-folded signals are shown in Figure 2.

### 3.1. ASAS Photometry

A large part of the problem in the signal characterization for HD 41248 surrounds the star’s rotational period. Therefore, in order to see if we could pin down the rotational period, we searched the latest version of the ASAS photometric catalog (Pojmanski 1997) to see what useful data there is for HD 41248

**Table 3**  
Solutions for HD 41248

Parameter	HD41248 <i>b</i>	HD41248 <i>c</i>
$P$ (day)	25.595 [25.551, 25.652]	18.361 [18.337, 18.392]
$K$ ( $\text{ms}^{-1}$ )	2.30 [1.39, 3.21]	1.95 [0.99, 2.83]
$e$	0.09 [0, 0.26]	0.10 [0, 0.28]
$\omega$ (rad)	0.3 [0, $2\pi$ ]	3.3 [0, $2\pi$ ]
$M_0$ (rad)	3.7 [0, $2\pi$ ]	5.6 [0, $2\pi$ ]
$a$ (AU)	0.166 [0.148, 0.180]	0.132 [0.118, 0.146]
$m_p \sin i$ ( $M_\oplus$ )	9.8 [5.9, 14.6]	7.6 [3.6, 11.6]
	Old data	New data
$\tau$ (day)	16.8 [0, 100]	1.4 [0, 100]
$\phi$	0.17 [−0.47, 0.87]	0.35 [0.10, 0.73]
$\sigma_j$ ( $\text{ms}^{-1}$ )	1.10 [0.77, 2.12]	2.25 [1.74, 2.81]
$c_{\text{BIS}}$	−0.05 [−0.29, 0.19]	−0.11 [−0.28, 0.07]
$c_{\text{FWHM}}$	0.03 [−0.09, 0.14]	0.13 [0.02, 0.26]
$c_S$ ( $\text{ms}^{-1} \text{dex}^{-1}$ )	91 [−135, 343]	67 [−93, 226]
$\dot{\gamma}$ ( $\text{ms}^{-1} \text{yr}^{-1}$ )	0.04 [−0.30, 0.37]	

and if we could locate a plausible rotational period. We obtained *V*-band photometry for this star, and after weeding out strong outliers (beyond  $5\sigma$ ) and selecting only the best data, those classed as “A” in the ASAS photometric grading, we were left



with a total of 420–650 photometric points processed through five ASAS *V*-band apertures.

Periodogram analyses of each of the aperture data revealed some significant peaks. A long-period peak is found beyond 1000 days, likely attributed to the magnetic cycle of the star. Another cluster of peaks emerge around 200–300 days, peaks that also appear in the radial velocities under the assumption of pure white noise (see the left column in Figure 1). We found neither a statistically significant peak that matched the signals we see in the radial velocity data nor a period that could plausibly relate to the stellar rotation period.

### 3.2. Activity Periods

In addition to analyzing the photometry, we also tested the activity indices by running both periodogram analyses and posterior samplings to constrain any frequencies in these indicators that would show activity cycles that could be the source of these radial velocity variations. Santos et al. (2014) show possible correlations between the 25 day radial velocity period and similar periods in the  $\log R'_{\text{HK}}$  and the CCF FWHM measurements.

First, we performed a periodogram analysis on the chromospheric activity indices after removing  $3\sigma$  outliers from the sample that, when included, only served to add noise. This analysis revealed a statistically significant frequency with a period of 27.62 days, close to the radial velocity period for the primary signal in the data. We then performed the same analysis on the FWHM values, using the same data set, and we found another statistically significant periodogram peak that matched the activity index peak, having a period of 27.93 days. The period of this FWHM changed to 25.31 days when we included all the FWHM data and subtracted off a linear trend from the time series, closely matching the radial velocity signal. The question to answer is whether these signals are related to the radial velocity variations.

In order to answer this question, we then ran MCMC samplings under our Bayesian approach to search for the signals independently and to constrain the significance of the signals and their possible extent in period space to see if they overlap with the radial velocity periods. We found the global probability maximum in the cleaned  $\log R'_{\text{HK}}$  indices to be located at 60 days with our samplings, closely followed by a 27.7 day maximum that matched the periodogram analysis. However, neither of these signals could be detected in the data according to the signal detection criteria because their periods and amplitudes could not be constrained from above and below. This means that we cannot rule out the possibility that these maxima are in fact statistical flukes caused by the combination of random noise and data sampling coupled with correlations.

We then chose to perform the same analysis on the FWHM measurements and found the strongest frequency to be at  $\sim 800$  days, after considering the linear trend. A cluster of probability maxima was found between 20 and 30 days in this analysis. However, none of these maxima corresponded to a genuine signal because they did not satisfy our signal detection criteria. Our interpretation was that the reason for the 25 day peak in the periodogram analysis was simply due to making the assumption that the noise is distributed in a Gaussian fashion, which does not appear to be the case as we find a significant red-noise component in both the  $S$  indices and the FWHM measurements. Therefore, the 25 day peak in the activity indicators is heavily model-dependent. This is in stark

contrast to the radial velocity signal at 25 days, which is found irrespective of the assumed noise model.

### 3.3. Signal Coherence

As discussed in the [Introduction](#), the signal found in the radial velocity time series of HD 166435 only had a coherence time of around 30 days. Combining this with the bisector correlations Queloz et al. (2001) ruled out the existence of this planetary candidate. In order to test the probability that the proposed planet candidates HD 41248*b* and *c* are real Doppler signals and independent from any activity correlations, we split the data up into two independent sets as a function of time. The first set was the original data published in Jenkins et al. (2013b) but analyzed using our current statistical model (data set 1), and the second set (data set 2) is the new data that is around 2.5 times larger than the first data and was added in the analysis from Santos et al. (2014). The full time baseline of data covers more than 10 yr, where data set 1 spans over 8 yr and data set 2 covers 2 yr, but at much higher cadence.

We analyzed both sets independently and in combination and recovered two signals each time with high statistical significance. The phase, period, and amplitude of the detected signals were in agreement with those published in Jenkins et al. (2013b). Table 2 shows the results for the sample for these analyses and includes results for the full data set with and without correlations with the activity indices.

The fact we see no change in the properties of the 25 day signal between the full data and the two subsets is remarkable because the star itself does change with time. We found that the intrinsic noise, parameterized with the standard deviation of the excess white noise  $\sigma_J$ , increased between the old and new data sets by roughly a factor of two. For the old data, the jitter was found to be  $1.3 \text{ ms}^{-1}$ , but for the new data, the jitter increased to  $2.6 \text{ ms}^{-1}$ . Furthermore, the data are more correlated in the second data set, with the correlation parameter ( $\phi_{1,1}$ ) changing from being consistent with zero for the old data to being significantly clustered around a value of 0.6 in the new data, meaning the noise becomes redder in the new data. This increase in stellar noise and correlation parameter means that the second signal at 18 days is no longer detected in the second set of data and it is likely the reason Santos et al. (2014) struggled to locate the 18 day signal as they did not account for these differences in the noise. This is also likely the reason they could not find a circular solution for the 25 day signal. We found a circular solution by both including the eccentricity prior or assuming a flat prior; however, Zakamska et al. (2011) show that there is a bias toward higher eccentricities in radial velocity surveys, a bias that our eccentricity prior helps alleviate. We also note that correlations between the activity indicators and the velocities became significant in the new data, whereas these correlations were not significant in the old data set (see Jenkins et al. 2013b).

Although the 18 day signal cannot be independently detected in the new data, likely due to the fact that the star has become more active and therefore disabling the detection of this weak signal behind the increased jitter and red noise, it is still well supported by the new data because its significance increases considerably when comparing between the old data and the full data (see Table 2). The significance of the  $k = 2$  model increases by a factor of 2700 between the old data and the full data, where it is  $1.8 \times 10^7$  times more probable than the  $k = 1$  model in the full data. This means that this signal together with the 25 day one retain their properties throughout the data baseline and cannot

be shown to be dependent on the changes in the stellar activity. This result also shows that serendipitously observing the first epoch of data when the star was intrinsically more inactive and the noise was whiter allowed both signals to be confirmed in that data set. There is the sampling cadence to consider here also since in the first epoch of data, the sampling density was lower than the later two observing seasons and it could also be that this decreased cadence helped to suppress the effects of the correlated noise.

We then split the data further into three groups so that we could examine the data sets presented in Santos et al. (2014) directly. We first consider the data published in Jenkins et al. (2013b) as one set as before (First Data), but then we analyze the following two observing seasons individually, where the middle set in time runs from JD 55904.8–56414.5 and the later set runs from JD 56521.9–56632.7. Santos et al. claim that the 25 day signal amplitude evolves with time across these three observing epochs, a strong argument in favor of an activity-induced signal; however, we could not confirm this to be the case. What we find is a complex evolution of the properties of the data due to changes in the star and sampling. First off, we did not find any significant evidence for a change in the amplitude of the 25 day signal in the first and middle parts of the data; the amplitudes are in strong statistical agreement. We could not confirm this result on the later data since we could not constrain any signal at all in this set, and below we explain why.

As discussed previously, it appears that HD 41248 became more intrinsically active throughout the time series of this data, in agreement with findings in Santos et al. (2014), however, not in a linear fashion. The jitter noise increased in the middle part of the data significantly, going from  $1.3 \text{ ms}^{-1}$  in the first part of the data, up to  $2.4 \text{ ms}^{-1}$  in the middle part, and then dropped again to  $2.1 \text{ ms}^{-1}$  for the later data, although this drop is currently not statistically significant. Added to this, the red-noise components also change with time, going from being consistent with zero in the first and middle parts of the data to a value of 0.6 in the later data, showing an increased red noise in the later data as well.

The activity correlations were also found to evolve throughout these three epochs of data. Again, these correlations were consistent with zero in the first part of the data for all three indicators,  $S$  index, FWHM, and BIS, but apart from the  $S$  index, whose correlations with the velocities do not appear to evolve with time, the other two indicators do. The correlations with FWHM ( $c_{\text{FWHM}}$ ) go from 0.13 in the middle data to 0.16 in the later data, both values statistically consistent but also statistically different from zero. The correlations with the BIS go from being negatively correlated in the middle data, with a value of  $c_{\text{BIS}} = -0.22^{+0.21}_{-0.20}$  (uncertainties mark the 99% credibility limits), and then become in agreement with zero again in the later data. Clearly the red noise presents a complex pattern in this time series, and when including all data after JD 55904.8, they are important and must be considered when searching for any low-amplitude signals in this data. In any case, it appears that our model does a reasonable job of describing the noise in the radial velocity time series for HD 41248, similar to the case of CoRoT-7 (Tuomi 2014).

Santos et al. (2014) claim a stable period and phase could be maintained in HD 41248 by an active longitude impacting the radial velocities (Berdyugina & Usoskin 2003; Ivanov 2007), under their hypothesis that the amplitude of the signal varies with time, which we have shown here cannot be concluded to be the case. Yet, Ivanov also show that although such active longitude spot formation zones maintain their form much longer than the

lifetimes of the individual spots, the solar data suggest they are only stable for 15–20 rotational periods. We show that the 25 day radial velocity signal found for HD 41248 has been static for *at least* 5 yr (70 rotational periods assuming the 25 day signal is the rotation period), and likely for the entire 10 yr baseline of data (140 rotational periods).

As a further test, we decided to again split the data up into independent sets, this time based on their chromospheric activity levels, to see if we were sensitive to the impact of the increased magnetic activity and spot formation of HD 41248. We built three almost equally numbered sets, comprising an inactive set ( $\log R'_{\text{HK}} \leq -4.93$ ) of 71 velocities, an intermediately active set ( $-4.93 < \log R'_{\text{HK}} \leq -4.91$ ) of 73 velocities, and an active set ( $\log R'_{\text{HK}} > -4.91$ ) of 79 velocities. We proceeded to search for the 25 day signal in these data sets and found that we could detect this signal in the inactive and intermediately active sets, but could not constrain anything in the active set. This affirms why the 18 day signal cannot be detected in the later data and also calls for a noise model to be scaled as a function of chromospheric activity, an upgrade we plan to include in future versions of our model. This feature also highlights that we are sensitive to activity-related features in our data; therefore, if the 25 day signal was genuinely due to rotationally modulated magnetic activity, we would expect the signal to appear stronger in the active data set than in the inactive data set since the sensitivity to the features causes the signal increase. We might also expect it to be more significant, depending on the structure in the increased jitter noise, which would also be modulated by the rotation. In any case, the signal parameters are invariant between the inactive and moderately active sets, showing that changes in the magnetic activity of the star do not change the period, amplitude, or phase of the signals, arguing against a magnetic origin for these signals.

Finally, we also tested the signals as a function of wavelength using the reddest HARPS orders only (see Anglada-Escudé & Butler (2012) and Tuomi et al. (2013) for details) and found no dependence of the signal properties or significances on wavelength. This indicates that neither of the signals show evidence for a dependence on wavelength, at least across the wavelength domain offered by HARPS. This would again argue against the origin of these signals being from magnetic activity cycles modulated by rotation.

#### 4. SUMMARY

We have shown that the radial velocity time series for the star HD 41248, covering nearly 10 yr of observation, clearly supports the existence of two signals with close to circular morphologies once red-noise components are considered. This analysis provides additional evidence that the pair of signals detected in this data could be due to a pair of planets in, or very near to, the 7:5 MMR with periods of 18.361 and 25.595 days and a period ratio ( $P_c/P_b$ ) of  $1.394 \pm 0.005$  at 99% confidence level, in excellent agreement with the results published in Jenkins et al. (2013b). Such resonances are known to be a byproduct of the planet formation and evolution process (Baruteau & Papaloizou 2013) in the early history of a star's life. It seems difficult to give rise to signals so close to such a period ratio simply by rotationally modulated activity in the presence of differential rotation, except in the most unique circumstances.

By analyzing the signals as a function of time in an independent fashion, we were able to obtain evidence for their static nature over the full baseline of observations. This analysis also

allowed us to show that the star got intrinsically more active within the period of the most recent data, with the jitter noise taking on a value twice that reported in Jenkins et al. (2013b). The noise also got significantly redder and the linear correlations with the velocities and FWHM increased such that they became statistically significant over the last two observing seasons. Splitting the data further revealed a complex pattern of evolving red-noise and activity correlations, both of which would serve to mask weak signals under a white-noise assumption. We conclude that noise correlations must be taken into account when attempting to search for periodic signals that are at the noise level of the star/instrument combination.

We analyzed the radial velocities with and without activity indicator correlations and found that both signals are supported by the old and new HARPS velocity data. Moreover, the significances of the signals increase when including the new data and when considering the activity indicator correlations. This is characteristic behavior of a pair of static Doppler signals. Including the linear correlation terms and red-noise correlations also results in removing spurious peaks from the white-noise model search, peaks that appear in the activity indicators. We also show that the signals in the activity indicators are highly model-dependent, only peaking at 25 days when Gaussian noise is assumed, whereas the 25 day signal in the radial velocity measurements is found no matter what the assumption of the noise is.

Further tests revealed that we are sensitive to changes in the magnetic activity, and we found that we could detect the 25 day signal in the radial velocities of the star when it was in its most inactive and moderately active states. No signals were detected when the star was in its most active state, contrary to what would be expected if these signals were due to magnetic activity since we might expect these signals to be strongest when the star is in an active state as the source of the signals should give rise to a stronger signal. We also note that the signals are independent of wavelength in the band covered by HARPS, a further argument against the magnetic activity cycle theory.

Additional confirmation that these signals could represent a pair of resonant planets may have to wait until future instruments operating in the near-infrared come online. Instruments such as HPF (Ramsey et al. 2008; Mahadevan et al. 2012) or CARMENES (Quirrenbach et al. 2012) could search for a change in the period or amplitudes of these signals as a function of wavelength over a much wider waveband than that offered by HARPS, which would attribute them to rotationally modulated spots on the surface of the star. Future direct imaging systems, such as the previously proposed TPF or Darwin missions (see Léger 2000), could be another way to confirm the existence of

these planets, yet this type of mission is a long way off in the future and the distance of 52 pc to HD 41248 makes this a real challenge. In any case, more high-cadence velocity observations over the coming years might be able to shed some light on the nature of these signals, either by searching for variations in the periods, amplitudes, and phases or by confirming the nature of these signals with more high-quality data. HD 41248, therefore, represents a very interesting target to monitor radial velocity signals buried within evolving red-noise and activity correlations.

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