

AlGaInP red-emitting light emitting diode under extremely high pulsed pumping

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ABSTRACT

Efficiency of commercial 620 nm AlGaInP Golden Dragon-cased high-power LEDs has been studied under extremely high pump current density up to 4.5 kA/cm² and pulse duration from microsecond down to sub-nanosecond range. To understand the nature of LED efficiency decrease with current, pulse width variation is used. Analysis of the pulse-duration dependence of the LED efficiency and emission spectrum suggests the active region overheating to be the major factor controlling the LED efficiency reduction at CW and sub-microsecond pumping. The overheating can be effectively avoided by the use of sub-nanosecond current pulses. A direct correlation between the onset of the efficiency decrease and LED overheating is demonstrated.

Keywords: AlGaInP LEDs, external quantum efficiency, self-heating, efficiency droop

1. INTRODUCTION

Light emitting diodes (LEDs) based on AlGaInN and AlGaInP materials systems cover the whole visible spectrum from 400 to 700 nm. While the AlGaInN based LEDs are most efficient in the blue-green spectral region, the long-wavelength region of the visible spectrum around yellow to red is benefitted by the efficient emission from the AlGaInP LEDs. The LEDs utilizing (Al_xGa_{1-x})_{0.5}In_{0.5}P quaternary alloys are normally lattice matched to their GaAs substrate of growth. These materials can provide light emission at the wavelengths from 560 nm to 650 nm by varying the aluminium (Al) content until the band gap becomes indirect around 555 nm ($x \sim 0.53-0.56$)¹. The application area for these devices ranges from indicator lamps to outdoor and automotive lighting². Though specific applications influence requirements on the device parameters, LEDs with higher efficiency over broad range of operating conditions are always beneficial and desirable.

It has been commonly observed that external quantum efficiency (EQE) of both AlGaInP and AlGaInN emitters, decreases with the forward current at high current densities³. This decrease of the efficiency limits LEDs to achieve their utmost performance and is being investigated intensively. In case of AlGaInN LEDs, the proposed dominant mechanisms responsible for the efficiency decrease (droop) are: (a) carrier leakage either due to asymmetric junction or polarization or both^{4,5}, (b) saturated radiative recombination⁶, (c) Auger recombination^{7,8}, (d) delocalization of carriers⁹ and (e) dislocation-induced non-radiative recombination.¹⁰ Observed similarities in EQE vs. current behaviour of AlGaInP and AlGaInN LEDs points to a possible common mechanism leading to the efficiency decrease for both material systems¹¹. However, for AlGaInP devices, polarization-induced electron leakage, increased non-radiative recombination at dislocations, and carrier delocalization are unlikely to dominate because of the absence of strong piezoelectric fields, the

use of heterostructures lattice-matched to GaAs substrate, and much weaker tendency of the alloy to carrier localization due to composition fluctuation, respectively.

At low current densities, carrier leakage due to asymmetric p-n junction has been suggested to contribute towards efficiency decrease in AlGaInP LEDs^{12,13}. Comprehensive studies of internal quantum efficiency (IQE) of amber/red AlGaInP LEDs^{14,15} have attributed the leakage to insufficiently high barriers separating quantum wells (QWs) in the LED active regions. They found the leakage to depend strongly on temperature with the activation energy nearly equal to the difference between the electron energy levels in the QWs and the conduction band bottom in the barriers. From this point of view, pushing up the electron quasi-Fermi level due to increase of the electron concentration in the QWs and device self-heating should produce similar effects on the efficiency reduction. It should be noted that elevated operating temperature may lower the LED efficiency by itself without enhancing the electron leakage. Indeed, the radiative recombination constant generally decreases with temperature, whereas the non-radiative carrier lifetime becomes shorter, i.e. both factors reduce IQE of the active region material. In order to distinguish between the roles of the MQW barrier heights controlling the electron leakage and elevated temperature by itself in operation of amber/red LEDs, special experiments are required, including LED examination under short-pulse current injection and rather high current densities. To our knowledge, such experiments have not yet been reported before.

The effect of device overheating and efficiency thermal rollover in LEDs can be mitigated by a packaged heat sink and/or by driving the device with short pulses with the duty cycle $\leq 1\%$ ¹⁶. In this paper, we will compare the behaviour of EQE in commercial 620 nm red high-power AlGaInP LED under continuous wave (CW) and microsecond-to-sub-nanosecond current-pulse excitation in the width modulation (PWM) regime up to the current densities of $\sim 4.5 \text{ kA/cm}^2$. Detailed characterization of the device under CW and PWM pumping is important for making conclusion on the nature of the efficiency reduction with current in phosphide LEDs.

2. DEVICE STRUCTURE

The $1 \text{ mm} \times 1 \text{ mm}$ high power red LEDs based on AlGaInP (Fig.1) used in this study have multiple quantum well (MQW) active regions emitting at 620 nm at room temperature. $\text{Al}_{0.5}\text{In}_{0.5}\text{P}$ layers are used for confining the carriers with an AlGaAs window and p-GaAs contact layer. The structure was grown on GaAs substrate. OSRAM OS's standard Thinfilm process has been applied for processing the devices and their packaging in Golden Dragon cases.

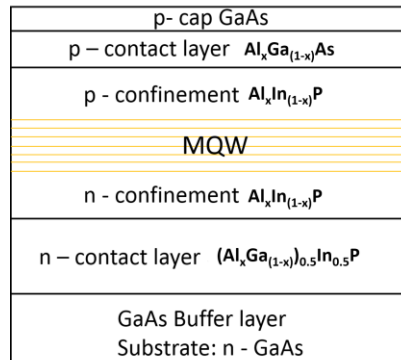


Figure 1: Schematic layer structure of high power red LED.

3. EXPERIMENTAL DETAILS

The optical parameters of the device under test were measured in both continuous wave (CW) and pulsed regime. A “CDS-600” spectrometer from Labsphere Co, Ltd. was used to collect light output along with the “LightMtrx” software to record the radiometric and photometric quantities. A “Keithley 2400” source was used to drive the LED under CW regime. Under pulsed regime, for currents up to 2 A, an “Agilent 8114A” pulse generator was used capable of producing pulses as short

as 200 ns. For the current range of 2-23 A and sub-microsecond pulse width, we utilized IXYS Colorado DEI generator with external triggering. For the sub-nanosecond high current (20-45 A) pulses, a specially designed high power hand-made generator was used. The pulse width and duty cycle were monitored on the oscilloscope using a current probe. To measure the sub-nanosecond electroluminescence (EL) spectra an ultrafast photo detector of 30 GHz bandwidth (BW) was used together with an Agilent 86100A 50 GHz oscilloscope. And, Si-based photo detector with the 14 ns rise time was used for sub-microsecond measurement. The time-resolved response was monitored and recorded on the oscilloscope with BW of 1GHz. The duty cycle was chosen in such a way, as to avoid the thermal impact of the previous pulses on the current one. Also, a closed-cycle Helium cryostat “Janis-CCS450” was used to perform temperature-dependent electroluminescence (TDEL) measurements.

4. RESULTS AND DISCUSSION

4.1 Spectral behaviour

Figure 2a demonstrates EL spectral dependence on pulsed current input. The pulsed input is varied over a broad current range from 100 mA to 5 A at a constant pulse duration of 100 μ s and 1% duty cycle. With increasing current, the emission peak shifts towards longer wavelengths, i.e. a redshift is observed¹⁷. This behaviour is similar to the behaviour observed under CW regime with the redshift starting at currents \geq 10 mA and greatly enhanced at currents $>$ 100 mA.

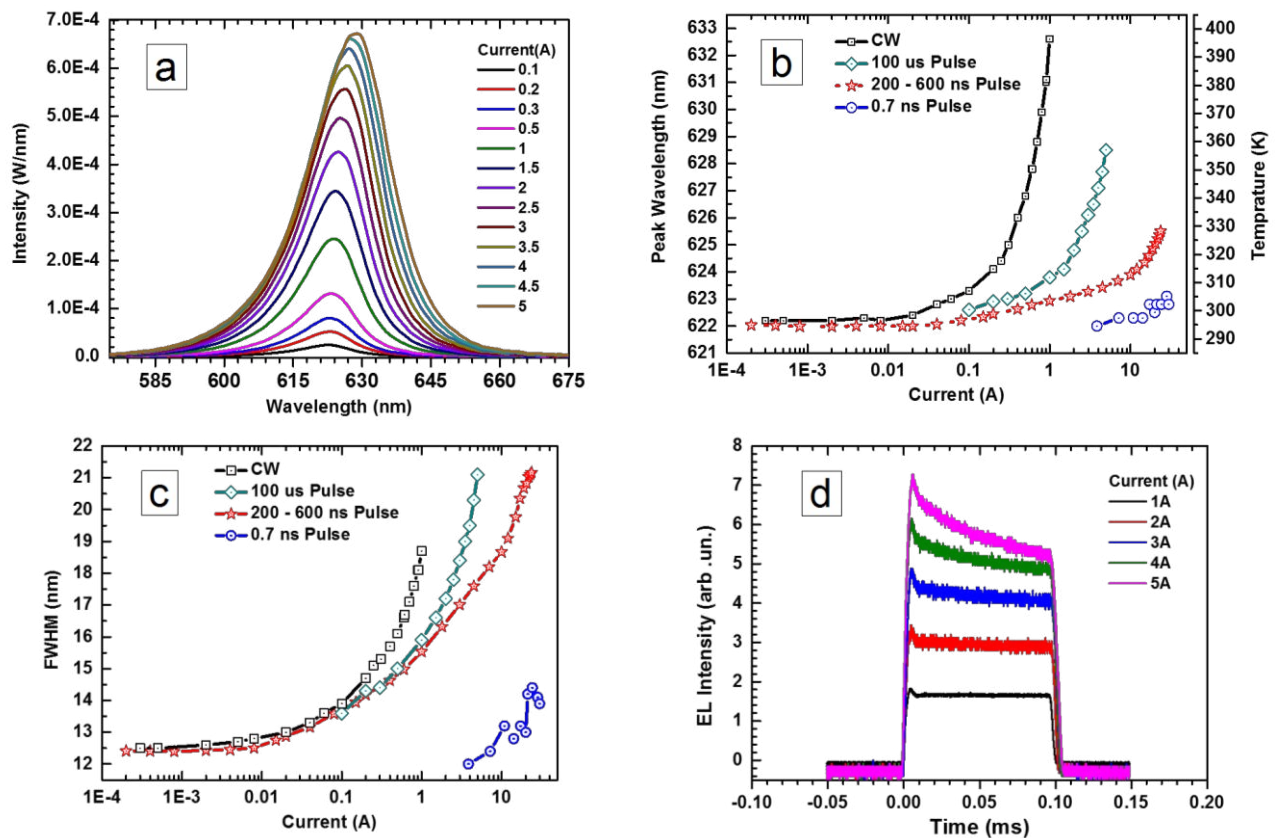


Figure 2: EL emission spectra of a red LED pumped by 100 μ s pulses with 1% duty cycle in the current range from 100 mA to 1 A (a). Evolution of peak wavelength/corresponding temperature variation (b) and FWHM of the emission spectrum (c) with current at CW (squares) and 100 μ s (rhombus), 200-600 ns (stars), and 0.7 ns pulsed excitation (circles). A short-term, i.e. during-the-pulse EL intensity reduction observed at 100 μ s pulsed excitation for current magnitudes from 1A to 5A.

Under 100 μs pulse regime at driving currents up to 1 A, a peak wavelength shift of ~ 1 nm with a ~ 2.5 nm increase in the spectrum full width half maximum (FWHM) is observed, as shown in Fig.2b,c, suggesting insignificant heating of the active region. However, at further increase in current up to 5 A, a red shift of ~ 6 nm (Fig.2b) and a spectral broadening of ~ 28.5 nm (Fig.2c) present a strong evidence for elevated active region temperature.

The asymmetric shapes and the slope variation of the short-wavelength edge of the emission spectrum with current shown in Fig.2a, can be attributed to increasing junction temperature¹³. To further investigate into behaviour of the peak wavelength and spectral width, the EL intensity evolution during the microsecond current pulses was monitored via photodiode response under the same operating conditions, as shown in Fig.2d. To determine EQE, in this case, radiant flux was collected for time much longer than the pulse period and then normalised to the duty cycle. A decrease in the EL intensity during the pulse for more than 2 A of the current magnitude is evident

After the above observations with microsecond pulses, spectral redshift under sub-microsecond (~ 100 -600 ns) and sub-nanosecond (0.7 ns) pulsed pumping in the extended range of currents up to 23 A was investigated.

Also, as shown in figure 2b and 2c respectively a redshift on ~ 4 nm in peak wavelength and ~ 10 nm of FWHM broadening is observed. Both the onset of decrease in efficiency and peak wavelength redshift starts in the current range of 600-700 mA. To understand temperature dependent shift of peak wavelength TDEL measurements over a wide temperature range from 13 to 450 K under CW regime at various currents were done. The experimental dependence (peak wavelength vs. temperature) was best described by a liner function with a coefficient of about 7.5 K/nm for temperatures above 150 K. Thus suggesting self-heating with change in temperature ΔT was about 30 K corresponding to 4 nm red shift in peak wavelength.

It is seen from Fig.2b that LED pumping by sub-nanosecond (~ 0.7 ns) current pulses with the 0.1% duty cycle provides the redshift of ~ 1 nm only. This suggests that the active region of the LED has negligible overheating under these pumping conditions.

4.2 Efficiency vs. current analysis

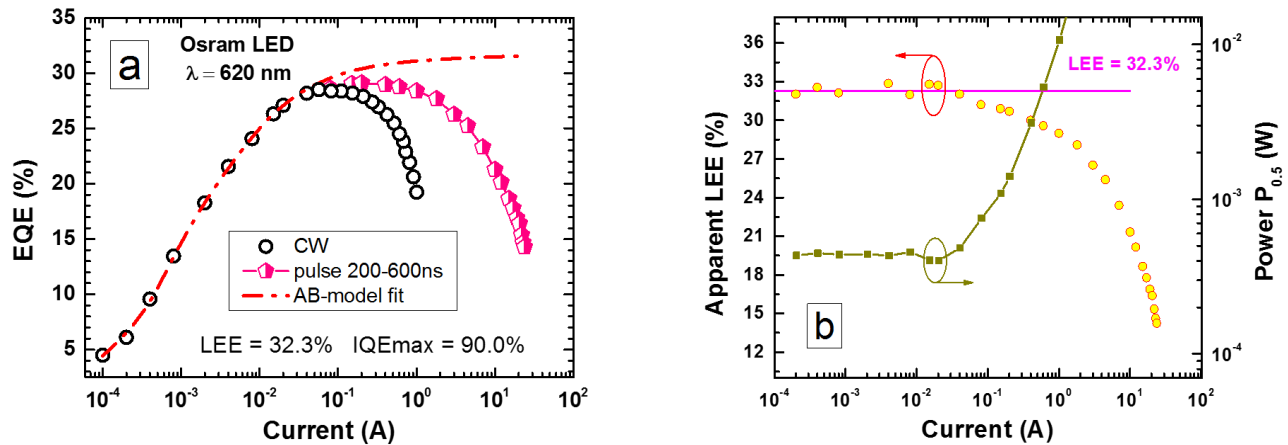


Figure 3: (a) External quantum efficiency under sub-microsecond pulses (half-filled pentagons) and CW regime (black open circles). Efficiency decrease is present for sub-microsecond pulses for current > 1 A after it starts to saturate beyond 100 mA of current. EQE is fitted with AB model (reduced ABC model) neglecting Auger coefficient, (b) apparent LEE and $P_{0.5}$ vs. current evaluated from experimental EQE data.

In order to estimate CW light extraction efficiency (LEE) and IQE of blue LEDs from the measured EQE(I) dependence, a special procedure based on the ABC recombination model has been suggested elsewhere¹⁸. Direct application of the procedure to red LEDs fails because of (i) small contribution of Auger recombination to the total recombination rate in the range of the current densities of interest and (ii) interfering the recombination processes with the electron leakage producing the efficiency decrease at high currents. Therefore, we suggest here an alternative procedure based on the reduced ABC-model called, hereafter, the AB-model. According to the model, only the Shockley-Read-Hall non-radiative recombination with the recombination constant A and radiative recombination with the recombination constant B are

accounted for, assuming the 100% -injection efficiency at low currents. Then EQE (η_e) of an LED structure can be estimated from the equation

$$\eta_e = \frac{\eta_{ext} P_{out}^{1/2}}{P_{0.5}^{1/2} + P_{out}^{1/2}} \quad (1)$$

where $P_{out} = \hbar\omega\eta_{ext}V_rBn^2$ is the output optical power of the LED with the active region volume V_r and non-equilibrium electron/hole concentration n ; $\hbar\omega$ is the mean energy of the emitted photons and η_{ext} is LEE. In Eq.(1), $P_{0.5} = \hbar\omega\eta_{ext}V_rA^2B^{-1}$ is the optical power corresponding to IQE of 50%. This parameter is generally unknown and cannot be directly evaluated from the measured EQE(I) dependence. Nevertheless, it can be excluded by regarding any pair of experimental points from the dependence, giving

$$\eta_{ext} = \frac{P_2^{1/2} - P_1^{1/2}}{P_2^{1/2}/\eta_2 - P_1^{1/2}/\eta_1}, \quad P_{0.5}^{1/2} = \frac{\eta_2 - \eta_1}{\eta_1/P_1^{1/2} - \eta_2/P_2^{1/2}}, \quad (2)$$

where $\eta_j = \eta_e(P_j)$, $j = 1, 2$ and P_j being the output optical power of the LED measured at two different driving currents. So, regarding a measured point as the reference one we can obtain η_{ext} and $P_{0.5}$ as a function of current similar to that shown in Fig.3b. One can see from the figure that both parameters calculated by Eqs.(2) remain to be nearly constant and equal to $\eta_{ext} = 32.3\%$ and $P_{0.5} = 0.45$ mW at the currents less than ~ 50 -60 mA. This value of η_{ext} is associated with the real LEE. At higher currents, apparent η_{ext} and $P_{0.5}$ estimated by Eqs.(2) start to vary, which is the evidence for IQE degradation caused by the device overheating or/and electron leakage.

Figure 3a shows the EQE(I) under CW and sub-microsecond excitation measured in an integrating sphere. Under CW regime of operation, the EQE rises steadily at lower current densities before it reaches its maximum at ~ 60 mA. A decrease in efficiency is observed when pumped at currents ≥ 150 mA. At the current of 700 mA, EQE drops by 15% of its maximum value. This sharp decrease in the efficiency can be explained by device self-heating. Unlike CW regime, the maximum efficiency is approached at the current of ~ 200 -250 mA at the sub-microsecond pulsed excitation, and the 15% efficiency drop is observed at ~ 8 -9 A.

The dash line on the Fig.3a shows EQE(I) predicted by the reduced AB-model, assuming isothermal operation conditions and no electron leakage to occur. Comparison of its behavior with the measured CW EQE(I) dependences enables estimation of the maximum IQE obtained at the currents of ~ 0.2 -0.5 A, which is of $\sim 90.0\%$ and compares well with the state-of-the-art. The low EQE maximum of $\sim 29\%$ is a result of a relatively low LEE = 32.3% specific for the devices under investigation that do not feature the typical lens cover. The AB-model fits very well to the experimental data of the low current CW branch, neglecting Auger recombination.

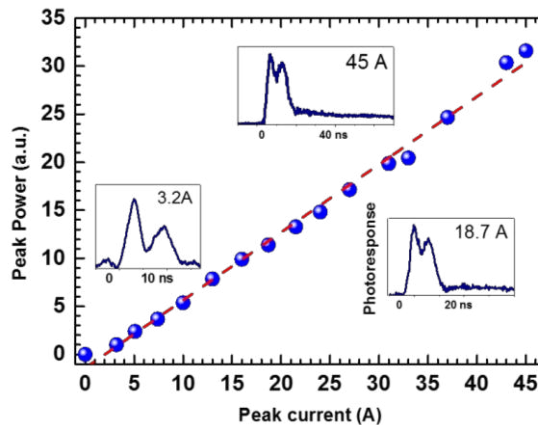


Figure 4: Peak power behavior under wide range of 0.7 ns peak current up to 45 A with 1% duty cycle, where blue circles represent the data and linear fitting is depicted by the dashed red line. In inset the photo-response pulse at 5, 20 and 45 A as observed on the oscilloscope is shown.

In order to exclude the effect of self-heating in our experiments, we studied the photo-response of LEDs under a wide range of 0.7 ns pulsed current pumping up to 45 A with 0.1% duty cycle. Figure 4 shows the typical dependence of the maximum peak output power of LED on the peak current of the pumping pulse. The insets to Fig.4 show the time-traces of the LED output recorded by 30 GHz BW photodetector and 50 GHz oscilloscope. The observed second peak of smaller amplitude of the LED output can be explained by parasitic oscillations. The experimental data (blue circles) is well expressed by a linear fitting function (dashed red line). The observed linear dependence would imply EQE not to decrease at high currents. This conclusion should be regarded, however with some care. Preliminary simulations of the time-dependent EL of the LED has shown that duration of the optical response is remarkably longer than that of the sub-nanosecond current pulse. This requires reconsidering the usual way of IQE and, hence, EQE determination, which will be done elsewhere. Nevertheless, the general conclusion could be done as following: the EQE does not decrease even at high currents, that means that increase in the carrier concentration in the active region without overheating does not result in the enhanced electron leakage and, hence, in the efficiency reduction.

5. CONCLUSION

To understand the nature of efficiency decrease at extremely high pump currents, the Golden Dragon-cased high power red LEDs were investigated under CW and pulsed regimes. Under CW operation the observed redshift of peak emission wavelength by ~ 10 nm and spectral broadening of ~ 4 nm for current range of 100 mA to 1 A is present due to device overheating. This is manifested on the EQE behavior as 10 % decrease in efficiency. Further, as expected, a low LEE due to the absence of silicone lens, of 32.3% and maximum IQE of 90% was estimated using AB model based on ABC recombination model excluding the Auger coefficient. Pulses with duty cycle of 1% and wide range of durations from 100 μ s to sub-microsecond with increasing extended current range were applied to mitigate the effect of self-heating. For 100 μ s second pulses the observed photodiode response showed a decrease in intensity during the pulse, indicative of device self-heating at this pulse duration. EQE behavior similar to CW regime is observed for sub-microsecond (100 – 600 ns) pulses. However, the onset of efficiency decrease is shifted to higher current of ≥ 600 mA along with the redshift of peak wavelength. This shift of 4 nm corresponds to 30 K temperature increase. Thus a direct correlation between onset of the efficiency decrease and LED overheating is observed. To exclude the effect of self-heating sub-nanosecond pulses of 0.7 ns with duty cycle of 0.1% and up to 45 A of peak current were applied. A linearly dependence of LED peak output power on increasing peak current is observed indicating towards lowering or no decrease in EQE under this regime even at such high current pulsed pumping along with a red shift of ≤ 1 nm. Nevertheless, it is observed that self-heating is an important factor contributing to efficiency decrease in these devices under all pumping regimes except for sub-nanosecond (0.7 ns) pulses. Since electron overflow is also temperature-dependent, we still cannot distinguish between the thermal effect on the competition between radiative and non-radiative recombination and electron overflow. The sub-nanosecond pumping would clarify this issue but this requires, at least, another way for EQE determination, as the duration of the optical response seems to be remarkably longer than the current pulse duration.

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