

Reaction-diffusion wave fronts: Multigeneration biological species under climate changeDaniel Campos,^{1,*} Joaquim Fort,^{2,†} and Josep Enric Llebot^{1,‡}¹*Grup de Física Estadística, Departament de Física, Universitat Autònoma de Barcelona, E-08193 Bellaterra, Spain*²*Departament de Física, Universitat de Girona, Campus de Montilivi, 17071 Girona, Catalonia, Spain*

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A generalization of reaction-diffusion models to multigeneration biological species is presented. It is based on more complex random walks than those in previous approaches. The new model is developed analytically up to infinite order. Our predictions for the speed agree to experimental data for several butterfly species better than existing models. The predicted dependence for the speed on the number of generations per year allows us to explain the change in speed observed for a specific invasion.

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I. INTRODUCTION

Many attempts have been made in order to describe biological migrations and colonizations by physical methods [1–4]. A possible approach to these problems is based on reaction-diffusion equations [4]. Extension of such equations by considering time-delayed processes has focused the attention of many physicists in recent years [5–7]. In this paper, we expand this approach to allow the study of multigeneration species, i.e., to explain range expansions for species that have several generations per year (gen/yr), separated by different delay (or resting) times. The problem we want to solve is fundamentally different from that of a waiting time distribution function considered by other authors [8,9]. In their models, particles or individuals may “jump” after a rest time τ_1 with probability p_1 , after a rest τ_2 with probability p_2 , etc., and this happens *at any instant of time*. In contrast, in the case we shall introduce below, there is a seasonal, non-overlapping succession of resting times.

Delayed diffusion-reaction models can be derived from random-walk movements [5]. Every particle or individual is supposed to move at successive steps (with time of travel t), separated by a time of rest τ [Fig. 1(a)]. Here we allow for more complex situations, by introducing the possibility of secondary steps with different travel and rest times $t_1, \tau_1, t_2, \tau_2, \dots$ [Fig. 1(b)].

In Sec. II we derive our model. In Sec. III, we apply it using typical dispersion and reproduction data for British butterflies, some of which present several gen/yr [10]. Such species have been observed to expand their ranges northwards in the past years, and biologists have pointed out climate change as one main reason [11–14]. We use our equation to predict the typical rates of spread, and compare them to experimental data and previous models. There is good agreement between theory and observations. We argue that this shows (i) the convenience of analytical models such as the one presented and (ii) that climatic change, on its own, does not explain the observed speed.

*Email address: daniel.campos@uab.es†Email address: joaquim.fort@udg.es‡Email address: enric.llebot@uab.es**II. MODEL**

Let $n(\vec{r}, t)$ be the density of particles or individuals in a two-dimensional space. If we assume that the particles jump in the way depicted in Fig. 1(b), we may write for the change in particle number in a differential of area dA during a secondary step of duration $T_i \equiv t_i + \tau_i$,

$$[n(\vec{r}, t + T_i) - n(\vec{r}, t)]dA = [n(\vec{r}, t + T_i) - n(\vec{r}, t)]_D dA + [n(\vec{r}, t + T_i) - n(\vec{r}, t)]_R dA, \quad (1)$$

where the first term in the right is due to diffusion, and the second one to net reproduction (if dealing with a biological species). Following Einstein’s approach [15], we write the diffusive term as the number of particles reaching the area dA minus those leaving it during T_i ,

$$[n(\vec{r}, t + T_i) - n(\vec{r}, t)]_D dA = \Phi - n(\vec{r}, t)dA, \quad (2)$$

where

$$\Phi = dA \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} n(x + \Delta x, y + \Delta y, t) \times \varphi_i(\Delta x, \Delta y) d\Delta x d\Delta y, \quad (3)$$

with $\varphi_i(\Delta x, \Delta y)dA$ the fraction of particles which have jumped from an area differential centered at $(x + \Delta x, y + \Delta y)$ at time t into another area differential centered at (x, y) in $t + T_i$. We also assume, as in [15], that all dispersion kernels are isotropic, $\varphi_i(\Delta x, \Delta y) = \varphi_i(\Delta)$, where $\Delta \equiv \sqrt{\Delta x^2 + \Delta y^2}$. Equations (1) and (2) can be approximated by Taylor series if the experimental data on the range expansion span along large enough times ($t \gg \sum_{i=1}^N T_i$) and distances ($x \gg \Delta x, y \gg \Delta y$). Following the same approach as in Ref. [16] we arrive at the expression up to k th order

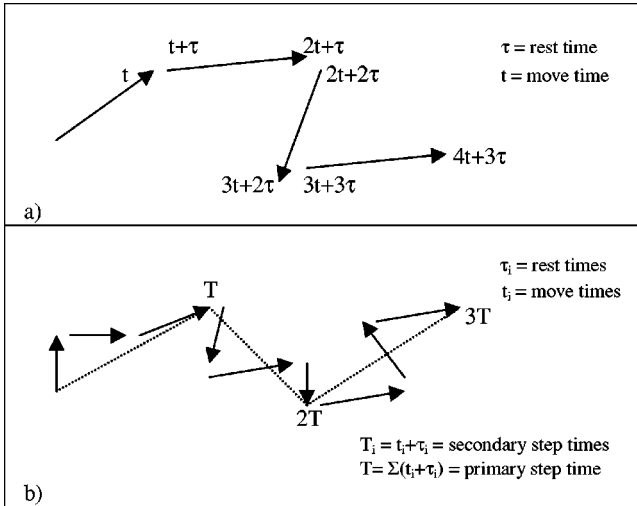


FIG. 1. Two kinds of random walks for the trajectory of a particle (or individual). (a) Classical diffusion, where the travel time (t) alternates with the rest time (τ) [4]. (b) The more general case in which there is a periodic distribution of travel (t_i) and rest (τ_i) times.

$$\sum_{j=1}^k \frac{T_i^j}{j!} \frac{\partial^j n}{\partial t^j} = \frac{2D_i T_i}{2!} \left(\frac{\partial^2 n}{\partial x^2} + \frac{\partial^2 n}{\partial y^2} \right) + \frac{(2D_i T_i)^2}{4!} \left(\frac{\partial^4 n}{\partial x^4} + 6 \frac{\partial^4 n}{\partial x^2 \partial y^2} + \frac{\partial^4 n}{\partial y^4} \right) + \dots + a \sum_{j=1}^k \frac{T_i^j}{j!} \frac{\partial^{j-1} n}{\partial t^{j-1}}, \quad (4)$$

where a represents the initial growth rate (i.e., $[\partial n / \partial t]_R \approx an$ for $n \approx 0$ [16]) and D_i is the diffusion coefficient of the i th substep,

$$D_i \equiv \frac{1}{4T_i} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \varphi_i(\Delta) (\Delta x^2 + \Delta y^2) d\Delta x d\Delta y. \quad (5)$$

As usual in this kind of analysis [4,16], we assume the existence of wave front solutions, by using in Eq. (4) solutions with the form $n \approx \exp(\lambda(x - vt))$, with $\lambda < 0$. Next, to analyze the process during a time interval encompassing N secondary steps [i.e., for a time interval $\sum_{i=1}^N T_i = \sum_{i=1}^N (\tau_i + t_i)$], we just have to write the expression resulting from Eq. (4) for the time interval $(t, t + T_1)$, for the interval $(t + T_1, t + T_1 + T_2)$, etc., and add up all these equations. This yields

$$\sum_{i=1}^N \sum_{j=1}^k \frac{(-T_i v \lambda)^j}{j!} = \sum_{i=1}^N \sum_{j=1}^k \frac{(2D_i T_i \lambda^2)^j}{(2j)!} + a \sum_{i=1}^N \sum_{j=1}^k \frac{T_i^j (-v \lambda)^{j-1}}{j!}, \quad (6)$$

which in the limit $k \rightarrow \infty$ acquires the form

$$\sum_{i=1}^N (e^{-v T_i \lambda} - 1) \left(1 + \frac{a}{v \lambda} \right) = \sum_{i=1}^N \cosh(\lambda \sqrt{2D_i T_i}) - N. \quad (7)$$

The wave front speed v can be found numerically from this equation in the usual way: for given parameter values, the speed is the minimum value of v such that a solution $\lambda < 0$ exists [16].

In order to show the generality of Eq. (7), let us take the limit $T_i \rightarrow 0$ (weakly delayed systems). From Eq. (7) one finds, up to first and second order in T_i , respectively,

$$D^* \lambda^2 + v^{(1)} \lambda + a = 0, \quad (8)$$

$$-v^{(2)} \lambda + \frac{v^{(2)2} T^*}{2} \lambda^2 = D^* \lambda^2 + a \left(1 - \frac{v^{(2)} T^*}{2} \lambda \right), \quad (9)$$

with

$$D^* = \frac{\sum_{i=1}^N D_i T_i}{\sum_{i=1}^N T_i}, \quad T^* = \frac{\sum_{i=1}^N T_i^2}{\sum_{i=1}^N T_i}. \quad (10)$$

In these special cases, using the fact that λ must be real, we reach the expressions for speed

$$v^{(1)} = 2\sqrt{aD^*}, \quad (11)$$

$$v^{(2)} = \frac{2\sqrt{aD^*}}{1 + \frac{aT^*}{2}}, \quad (12)$$

which correspond to the wave front speed for the well-known parabolic (i.e., nondelayed) [17] and hyperbolic (i.e., weakly delayed) [4,18] approximations, respectively, except that here T^* and D^* appear (instead of T and D in Ref. [4]). This is due to our assumption of secondary steps. The classical, single-substep case [4] is recovered for $N=1$. Equation (7) is more general than Eq. (12), because it applies even to strongly delayed diffusion. Moreover, the new diffusive behavior due to i secondary steps [Fig. 1(b)] is considered in the general expression (7) and taken into account even for the parabolic and hyperbolic approximations by means of the new parameters T^* and D^* .

III. APPLICATIONS

(1) *Invasion speeds.* One of the most direct applications of our new model, and indeed its original motivation, is the spread of biological species with N generations per year. Let us see whether our model is able to predict the magnitude of the observed speeds.

Many butterfly species have up to 4 (or sometimes even more) summer generations, whereas they do not appear in colder seasons. The periods during which individuals are able to disperse, known as flight times, are typically $t_1 \approx t_2 \approx \dots \approx 30$ days, separated by $\tau_1 \approx \tau_2 \approx \tau_3 \approx 10$ days [10], ex-

cept for the last annual generation. For this N th generation, because of the annual behavior ($T \equiv \sum_{i=1}^N T_i = 1$ yr), we have $\tau_N = T - t_1 - \tau_1 - t_2 - \dots$. Thus, if $N=4$ gen/yr, $\tau_4 \approx 215$ days (choosing similar values would not change considerably our final results).

The diffusion coefficient D_i (5) can be rewritten, according to a well-known result of basic diffusion theory [19], as

$$D_i = \frac{1}{4T_i} \frac{\langle \Delta^2 \rangle_{obs}}{t_{obs}} t_i, \quad (13)$$

where the subscript *obs* means that these values correspond to direct observations. The values of $\langle \Delta^2 \rangle_{obs}$ and t_{obs} are determined in ecology from mark-recapture experiments, which yield $2 < D_i < 31$ km²/yr for butterflies [20,21], depending on the species considered and the ecological conditions. Such a range may seem rather wide, but in fact, in biological systems the range of diffusion coefficients observed spans many orders of magnitude (from 10^{-5} to 10^4 km²/yr [22]). Finally, typical values of a between 0.6 and 1 yr^{-1} follow from the population field data as a function of time in Ref. [23].

We may now compare the predictions to experimental observations. Nonmigratory butterfly expansions have been measured for the last thirty years in Great Britain and the US [23]. Their observed spread rates are in the range 1.3–8.6 km/yr [24]. Introducing the values above for t_i, T_i, a, D_i into our new Eq. (7), we obtain the prediction (for 4 gen/yr) $1.9 < v < 8.4$ km/yr, which agrees quite well with the observed range above.

The usefulness of assuming several generations can be seen by comparing with the results of a model [16] based in the same Eq. (7), but with $N=1$ gen/yr. In that model we take $t = \sum_{i=1}^4 t_i \approx 120$ days, so the annual flight time is the same in both cases. The range predicted by the previously known model is $3.3 < v < 14.2$, which is less consistent with the experimental data [24]. This is also observed in Fig. 2, where our new model (solid lines) is more consistent with the observed values of a and D_i than the model with $N=1$ gen/yr (dotted lines).

The inset in Fig. 2 shows the convenience of assuming different waiting times from variations in v/v_0 (v_0 is the speed value for 1 gen/yr) as a function of N . Triangles correspond to the case with identical waiting times $T_1 = T_2 = \dots = T_i$ (that is, $T^* = T_i$ and $D^* = D_i$). Squares correspond to experimental values for t_i, τ_i above. The differences between both models are qualitatively important, especially for low N , simply because in our case (squares) generations are concentrated in a specific season (summer). When N is high enough, generations become uniformly distributed throughout the year, and both models lead to similar results, as they should. Of course, for other biological species, differences between both models could be important up to much higher values of N .

(2) *Changes of speed.* One of the very few butterfly species whose spreads have been measured in detail is *Pieris Rapae*. Andow *et al.* [25] analyzed its southward invasion across the US in the 19th century. They mentioned that the number of gen/yr for this species is higher in Missouri (6–7

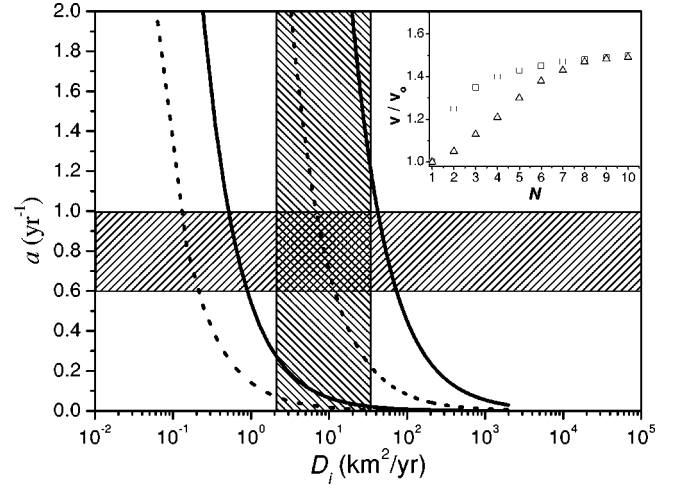


FIG. 2. Hatched regions define the range of independently observed values for the parameters a and D_i . The solid lines correspond to the new multigeneration model (4 gen/yr) with $v = 1.3$ and 8.6 km/yr, which is the experimental range, $T_1 = T_2 = T_3 = 40$ days, and $T_4 = 245$ days. Dotted lines are for the model with $N = 1$ gen/yr [16] and flight time $t = \sum_{i=1}^4 t_i$. Inset: Speed ratio v/v_0 vs number of generations N for the multigeneration model (squares, $T_1 = T_2 = \dots = T_{N-1} = 40$ days, $T_N = 1 \text{ yr} - \sum_{i=1}^{N-1} T_i$ days) and assuming all T_i are equal (triangles, $T_1 = T_2 = \dots = 1 \text{ yr}/N$).

gen/yr) than in more Northern locations (typically 3 gen/yr), all generations with similar flight times ($t_1 \approx t_2 \approx \dots$). They even argued that this could explain the increase in the front speed observed during this expansion. This gives us a unique opportunity to check whether the tendency of the speed on the number of generations agrees with our model or not. We checked the original data from Scudder [26] before *P. Rapae* reached Missouri ($N_{BM} \approx 3$ gen/yr) and after passing Missouri ($N_{AM} \approx 6-7$ gen/yr) and found that the observed ratio of both speeds was $v_{AM}/v_{BM} = 1.5 \pm 0.2$. We will compute the predicted speed ratio, instead of both speeds, because *P. Rapae* is an extremely migrating species so its diffusion coefficient is very difficult to measure and the value of a for this species is also unknown. Equation (7) does not allow us to perform speed ratios, and we use the approximation (11). If we recall the annual behavior ($\sum_{i=1}^N T_i \equiv T = 1$ yr) and assume, as mentioned, the same flight times t_i for all generations i , Eqs. (10) and (13) yield $D^* = N < \Delta^2_{obs} > t_i / (4t_{obs} \cdot 1 \text{ yr})$. Then,

$$\frac{v_{AM}}{v_{BM}} = \sqrt{\frac{N_{AM}}{N_{BM}}}. \quad (14)$$

Finally, we predict from Eq. (14) $v_{AM}/v_{BM} = 1.6$, which is very similar to the experimental value above. Additional agreement and more detailed comparisons to observations are expected to be achieved in the future, as new data for specific species will be obtained.

IV. DISCUSSION

In this work we have extended reaction-diffusion models to consider, up to infinite order, a more general diffusive

pattern than previous authors. The new random walk is characterized by a sequential, periodic succession of move and rest times [Fig. 1(b)]. The need for this more general behavior comes from biological applications, namely from the fact that some species present more than one generation per year.

The situation considered can be regarded as a rapidly time-varying value for the rest time. In principle, this could lead to a rapidly changing wave front speed. But the experimental data are not accurate enough to show such an effect. This is why we have presented an averaged description, leading to a constant speed.

We stress that we have presented a specific application for illustration purposes, and because of the interest that phenological response have as an indicator of climate change. Specifically, nonmigratory butterflies have been intensively studied with the aim to predict the effects of climate, weather and habitat variations [11,14,27,28]. Numerical methods have recently been used to describe the polewards expansions of butterflies and other biological species [29,13]. These approaches are able to predict whether the species can be expected to be present in a given locality, but such a detailed, microscopic method assumes in fact that butterfly range shifts can be as high as allowed by the speed of climate change, which has been criticized [30]. A very recent paper

[14] summarizes a list of species which are expanding their ranges in response to climate change. The speeds listed there have very different values. Thus, it is unlikely that these expansions can be explained just by the speed of climatic isotherms [11,31].

This argument makes us think that analytical approaches to these processes, as proposed here, could be useful. The results obtained by our model are quite satisfactory: predicted (1.9–8.4 km/yr) and observed (1.3–8.6 km/yr) ranges for butterflies are very similar (see Fig. 2). Hence, we consider that our model may be applied to the study of other multigeneration species, as well as to new applications which are in accordance with the diffusive behavior in Fig. 1(b).

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- [1] K. Sznajd-Weron and A. Pekalski, *Physica A* **294**, 424 (2001).
 [2] I. Mróz, A. Pekalski, and K. Sznajd-Weron, *Phys. Rev. Lett.* **76**, 3025 (1996).
 [3] F. Manzo and L. Peliti, *J. Phys. A* **27**, 7079 (1994).
 [4] J. Fort and V. Méndez, *Phys. Rev. Lett.* **82**, 867 (1999).
 [5] V. Méndez and J. Camacho, *Phys. Rev. E* **55**, 6476 (1997).
 [6] G. Abramson, A.R. Bishop, and V.M. Kenkre, *Phys. Rev. E* **64**, 066615 (2001).
 [7] S. Fedotov, *Phys. Rev. Lett.* **86**, 926 (2001).
 [8] V. Méndez, Ph.D. thesis, UAB, 1998.
 [9] S. Fedotov and Y. Okuda, UMIST report (unpublished).
 [10] E. Pollard and T. J. Yeats, *Monitoring Butterflies for Ecology and Conservation* (Chapman & Hall, London, 1993).
 [11] C. Parmesan *et al.*, *Nature (London)* **399**, 579 (1999).
 [12] S.L. Pimm, *Nature (London)* **411**, 531 (2001).
 [13] M.S. Warren *et al.*, *Nature (London)* **414**, 65 (2001).
 [14] G.R. Walther *et al.*, *Nature (London)* **416**, 389 (2002).
 [15] A. Einstein, *Investigations on the Theory of the Brownian Movement* (Dover, New York, 1956).
 [16] J. Fort and V. Méndez, *Phys. Rev. E* **60**, 5894 (1999).
 [17] R.A. Fisher, *Ann. Eugenics* **7**, 355 (1937).
 [18] V. Méndez, J. Fort, and J. Farjas, *Phys. Rev. E* **60**, 5231 (1999).
 [19] J. Crank, *The Mathematics of Diffusion* (Oxford University Press, London, 1956).
 [20] R.E. Jones, N. Gilbert, M. Guppy, and V. Nealis, *J. Anim. Ecol.* **49**, 629 (1980).
 [21] P.M. Kareiva, *Oecologia* **57**, 322 (1983).
 [22] E.D. Grosholz, *Ecology* **77**, 1680 (1996).
 [23] J. L. Banwell and T. J. Crawford, "Butterfly Monitoring Scheme. Key indicators for British wildlife," Report PECD 7/2/84, University of York, 1992. All data used are available at <http://bms.ceh.ac.uk/>
 [24] For 35 nonmigratory butterfly species, the shifts to the north have been by 35–200 km [11] over 27 yr (Ref. [14], Table II), which yields speeds from 1.3 to 7.4 km/yr. This is consistent with Ref. [27], which reports 1.4 to 8.6 km/yr for 12 species, and with the value 2.9 km/yr for *Pyronia Tithonus* (Ref. [10], p. 125).
 [25] D. A. Andow, P. Kareiva, S. Levin, and A. Okubo, in *Evolution of Insect Pests: The Pattern of Invasions*, edited by K. C. Kim (Wiley, New York, 1993).
 [26] S.H. Scudder, *Mem. Boston Soc. Nat. Hist.* **4**, 53 (1887).
 [27] J. Asher *et al.*, *The Millennium Atlas of Butterflies in Britain and Ireland* (Oxford University Press, Oxford, 2001).
 [28] R. L. H. Dennis, *Butterflies and Climate Change* (Manchester University Press, Manchester, 1993).
 [29] J.K. Hill, C.D. Thomas, and B. Huntley, *Proc. R. Soc. London, Ser. B* **266**, 1197 (1999).
 [30] A.T. Peterson *et al.*, *Nature (London)* **416**, 626 (2002).
 [31] M. Beniston *et al.*, in *The Regional Impacts of Climate Change, Intergovernmental Panel on Climate Change Working Group II*, edited by R. T. Watson, M. C. Zinyowera, and R. H. Moss (Cambridge University Press, Cambridge, England, 1998), pp. 149–185.