

PLASMA AND MAGNETIC FIELD CHARACTERISTICS OF SOLAR CORONAL MASS EJECTIONS IN RELATION TO GEOMAGNETIC STORM INTENSITY AND VARIABILITY

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ABSTRACT

The largest geomagnetic storms of solar cycle 24 so far occurred on 2015 March 17 and June 22 with D_{st} minima of -223 and -195 nT, respectively. Both of the geomagnetic storms show a multi-step development. We examine the plasma and magnetic field characteristics of the driving coronal mass ejections (CMEs) in connection with the development of the geomagnetic storms. A particular effort is to reconstruct the in situ structure using a Grad-Shafranov technique and compare the reconstruction results with solar observations, which gives a larger spatial perspective of the source conditions than one-dimensional in situ measurements. Key results are obtained concerning how the plasma and magnetic field characteristics of CMEs control the geomagnetic storm intensity and variability: (1) a sheath-ejecta-ejecta mechanism and a sheath-sheath-ejecta scenario are proposed for the multi-step development of the 2015 March 17 and June 22 geomagnetic storms, respectively; (2) two contrasting cases of how the CME flux-rope characteristics generate intense geomagnetic storms are found, which indicates that a southward flux-rope orientation is not a necessity for a strong geomagnetic storm; and (3) the unexpected 2015 March 17 intense geomagnetic storm resulted from the interaction between two successive CMEs plus the compression by a high-speed stream from behind, which is essentially the “perfect storm” scenario proposed by Liu et al. (i.e., a combination of circumstances results in an event of unusual magnitude), so the “perfect storm” scenario may not be as rare as the phrase implies.

Key words: shock waves – solar–terrestrial relations – solar wind – Sun: coronal mass ejections (CMEs)

1. INTRODUCTION

A topic of increasing interest to space weather is how the plasma and magnetic field characteristics of coronal mass ejections (CMEs) result in geomagnetic storm activity, in particular those intense events. The southward magnetic field and speed of CMEs at the Earth have received the most attention, because their cross product, the dawn-dusk electric field, controls the rate of the solar wind energy coupling to the terrestrial magnetosphere (Dungey 1961). However, it is still not clear how the ejecta speed and southward magnetic field work together to achieve a sustained, enhanced dawn-dusk electric field and how they lead to the variability of geomagnetic storms.

The southward magnetic field is often found within the ejecta reaching the Earth in the form of an interplanetary CME with a preceding shock. This usually leads to a classic geomagnetic storm sequence: a sudden commencement generated by the shock, a main decrease phase caused by the ejecta’s southward magnetic field, and then a recovery phase. In addition to the driver gas, the sheath region between the shock and ICME can also be geo-effective (e.g., Tsurutani et al. 1988; Liu et al. 2008b) as both the sheath speed and southward magnetic field are amplified by shock compression. The sheath-ejecta scenario has been invoked to explain the two-step development of geomagnetic storms (Kamide et al. 1998). Complex ejecta resulting from interactions between CMEs (Burlaga et al. 2001, 2002) can be very geo-effective owing to their prolonged durations (e.g., Farrugia & Berdichevsky 2004; Zhang et al. 2007; Lugaz & Farrugia 2014; Mishra et al. 2015). They could also cause two-step geomagnetic storms (Farrugia et al. 2006; Liu et al. 2014b).

A special case of complex ejecta is the interaction of a preceding ejecta with an overtaking shock (e.g., Harrison et al. 2012; Liu et al. 2012, 2014b; Möstl et al. 2012; Webb et al. 2013). The shock enhances the pre-existing southward magnetic field inside the ejecta, an idea for increased geo-effectiveness dating back several decades (Burlaga 1991; Vandas et al. 1997). A statistical analysis indicates that 19 out of 49 shocks propagating inside ICMEs are associated with an intense geomagnetic storm (minimum $D_{st} < -100$ nT; Lugaz et al. 2015). A recent study combining remote-sensing and in situ observations suggests a “perfect storm” scenario for the generation of an extreme space weather event (Liu et al. 2014a): preconditioning of the upstream solar wind by an earlier CME plus in-transit interaction between two later, closely launched CMEs, in order to have an exceptionally high solar wind speed and unusually strong ejecta magnetic fields at 1 AU. This, again, emphasizes the crucial importance of CME–CME interactions for space weather.

On 2015 March 17 and June 22 the Earth underwent an intense geomagnetic storm with the minimum D_{st} of -223 and -195 nT, respectively. These are the largest geomagnetic storms of solar cycle 24 so far. They occurred in the declining phase of the solar cycle, a phenomenon that is not uncommon (Gopalswamy et al. 2005; Kilpua et al. 2015). We provide a timely analysis of the driving CMEs, in an attempt to identify the plasma and magnetic field characteristics controlling the geomagnetic storm intensity and variability. Another focus of this letter is to test the “perfect storm” scenario proposed by Liu et al. (2014a): whether it is a rare coincidence or if it happens more frequently than what the phrase suggests. We examine the solar wind signatures and their connections with the

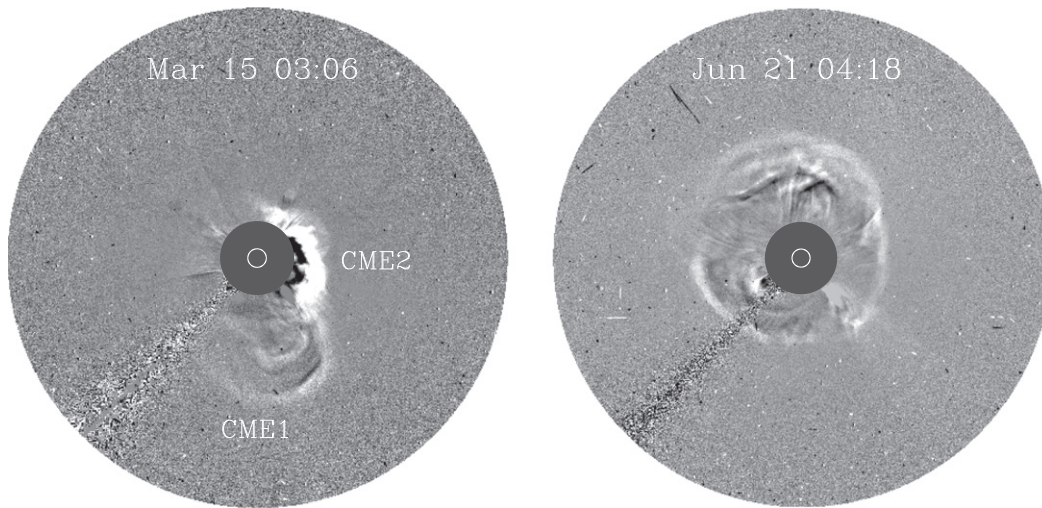


Figure 1. Difference images of the 2015 March 15 (left) and June 21 (right) CMEs from LASCO C3 aboard *SOHO*. Note the interaction between the March 15 CME (CME2) and a preceding one that occurred on March 14 (CME1).

development of the geomagnetic storms, complemented with the modeling of the D_{st} index using two empirical formulae based on the solar wind measurements (Burton et al. 1975; O’Brien & McPherron 2000). We also use a Grad–Shafranov (GS) technique (Hau & Sonnerup 1999; Hu & Sonnerup 2002), which has been validated by well separated multi-spacecraft measurements (Liu et al. 2008a; Möstl et al. 2009), to reconstruct the in situ ICME structure. The GS method can give a cross section as well as flux-ropes orientation without prescribing the geometry. In conjunction with solar observations it provides a larger spatial perspective of ICMEs than one-dimensional in situ measurements (Liu et al. 2010). These efforts are key to understanding how the plasma and magnetic field characteristics of CMEs are connected with the intensity of geomagnetic storms as well as their variability.

2. THE 2015 MARCH EVENT

Tracing back to the Sun, the drivers of the 2015 March 17 geomagnetic storm were two interacting CMEs on March 15 (Figure 1, left). The second CME (CME2) had a maximum speed of about 1100 km s^{-1} and was associated with a long-duration C9.1 flare from AR 12297 (S22°W25°) that peaked at 02:13 UT on March 15. The first CME (CME1) occurred on March 14 and had a speed of about 350 km s^{-1} . It was likely associated with a C2.6 flare from the same active region (S21°W20°) that peaked around 11:55 UT on March 14. A first impression from the coronagraph images is that CME1 was largely propagating southward while CME2 had a major component heading west. One may expect that the Earth would encounter the flank of the ejecta, so this would not raise the alarm for a major geomagnetic storm. Another fact that also makes the occurrence of an intense geomagnetic storm surprising is that the associated flares are relatively weak. In this sense, the occurrence of the 2015 March 17 intense geomagnetic storm is similar to the formation of the 2012 July 23 super solar storm that impacted STEREO A (Liu et al. 2014a). Without white-light observations from STEREO accurate CME kinematics cannot be obtained.

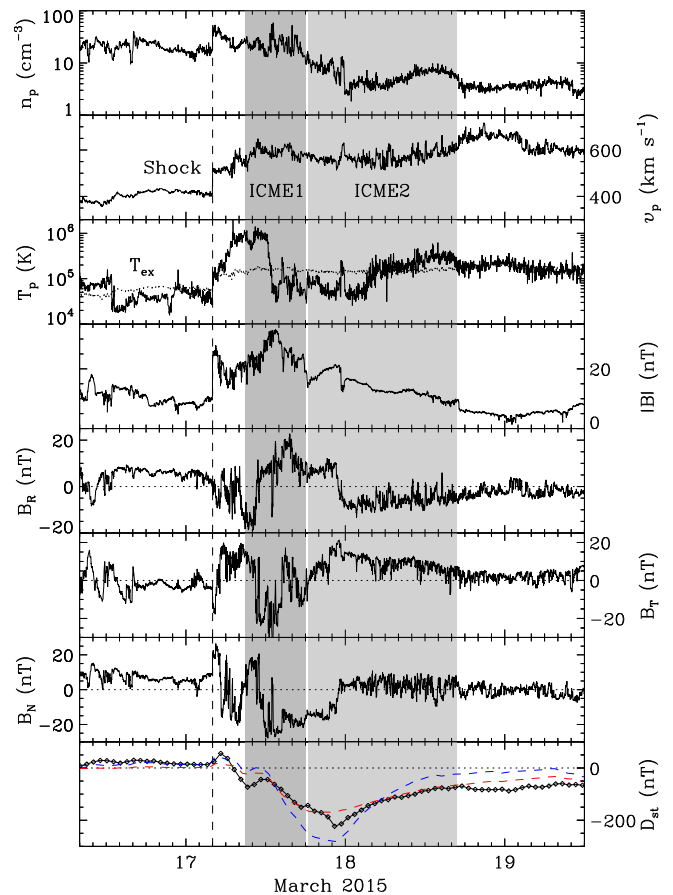


Figure 2. Solar wind measurements at *Wind* and associated D_{st} index for the 2015 March 17 event. From top to bottom, the panels show the proton density, bulk speed, proton temperature, magnetic field strength and components, and D_{st} index, respectively. The shaded regions indicate two ICME intervals. The vertical dashed line marks the associated shock. The dotted curve in the third panel denotes the expected proton temperature calculated from the observed speed (Lopez 1987). The red and blue curves in the bottom panel represent D_{st} values estimated using the formulae of O’Brien & McPherron (2000) and Burton et al. (1975), respectively.

Figure 2 shows the in situ signatures observed at *Wind*. A shock passed *Wind* at 04:01 UT on March 17 and caused the sudden commencement of the geomagnetic storm. It is difficult to unambiguously connect the in situ signatures with the coronagraph images without wide-angle imaging observations from STEREO. However, application of an empirical model (Gopalswamy et al. 2000) with the CME speed of 1100 km s^{-1} gives a predicted arrival time of 23:59 UT on March 16 at *Wind* (0.99 AU from the Sun), which is only 4 hr earlier than the observed shock arrival. As can be seen from the figure, two ICMEs (or flux ropes) are identified. Our interpretation is different from those of Kataoka et al. (2015) and Gopalswamy et al. (2015) who identify a single, shorter ICME interval from the data (although with different durations). The reason that we believe there are two ICMEs is as follows. First, there are multiple rotations in the magnetic field components whose polarities change twice in the shaded data intervals. Obviously these features cannot be explained by a single flux rope. Second, our ICME intervals are an outcome of the GS reconstruction, which is sensitive to the chosen boundaries. Despite the magnetic field fluctuations, both of the ICMEs can be reconstructed fairly well (see description below). Third, the interpretation of two ICMEs is consistent with what the solar observations indicate (see Figure 1 and discussions below). One may argue for a single ICME interval (say, from 12:58 UT on March 17 to 03:22 UT on March 18) based on the depressed proton temperature, a signature often used to identify ICMEs. A reasonable GS reconstruction, however, cannot be obtained for the interval and its variations. Given the presence of the compression by a high-speed stream from behind and CME–CME interactions, the criterion of a low proton temperature for identifying ICMEs may not be valid in the current case.

Within the ICME intervals the maximum magnetic field strength is about 33 nT while the southward component reaches -25 nT . These are not small magnetic fields considering only the flanks of the CMEs are encountered (see Figure 1). Interactions between these two ICMEs may have inhibited their expansion, a mechanism to create strong ejecta magnetic fields as we have seen from the 2012 July 23 event (Liu et al. 2014a). Also note the high-speed stream compressing ICME2 from behind, which may help maintain a strong ejecta magnetic field and a relatively high speed as well. The D_{st} profile indicates a two-step geomagnetic storm sequence with a global minimum of -223 nT . The first dip is produced by the southward magnetic field component in the sheath region behind the shock, while the second one results from the southward fields within the two ICMEs that last about 12 hr. Given the presence of a preceding shock and two interacting ICMEs, the two-step development of the geomagnetic storm can be classified as a sheath-ejecta-ejecta scenario. The modeled D_{st} index using the O’Brien & McPherron (2000) formula (minimum -170 nT) generally agrees with actual D_{st} measurements but underestimates the global minimum. The Burton et al. (1975) scheme gives a deeper global minimum (-280 nT) and a shallower recovery phase than measured.

Further information on how the plasma and magnetic field characteristics control the geomagnetic storm activity is obtained from the GS reconstruction, as shown in Figure 3. The reconstructions give a right-handed flux rope structure for both of the ICMEs, as can be judged from the transverse fields along the spacecraft trajectory together with the axial direction. Their axis orientations are almost opposite to each other: an elevation angle of about 33° and azimuthal angle of about 256°

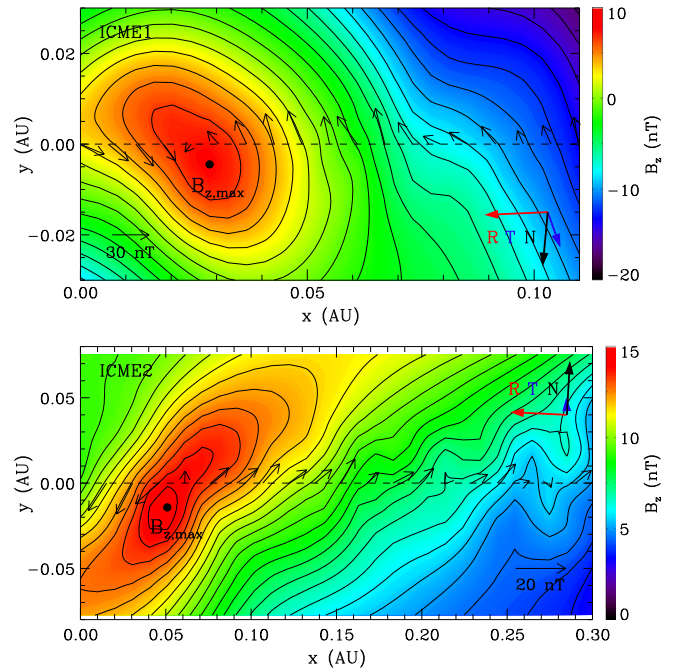


Figure 3. Reconstructed cross sections of ICME1 (upper) and ICME2 (lower). Black contours show the distribution of the vector potential, and the color shading indicates the value of the axial magnetic field. The location of the maximum axial field is indicated by the black dot. The dashed line marks the trajectory of the *Wind* spacecraft. The thin black arrows denote the direction and magnitude of the observed magnetic fields projected onto the cross section, and the thick colored arrows show the projected RTN directions.

(in RTN coordinates) for ICME1, and an elevation angle of about -18° and azimuthal angle of about 92° for ICME2. These low inclinations are consistent with the slightly tilted neutral line (N. Gopalswamy 2015, private communication) and filament channel (M. Temmer 2015, private communication) associated with the active region. ICME1 has a larger elevation angle (33°), which may help explain the encounter in spite of the largely southward propagation direction of CME1 (Figure 1). The impact of ICME2 may be accounted for by its lower elevation angle (-18°), although CME2 had a major section propagating westward. In addition, the angular span of an ICME can be 60° or even larger (Bothmer & Schwenn 1998; Richardson et al. 2002), so it is not surprising that both of the ICMEs hit the Earth. It is also likely that the CME–CME interactions, the following high-speed stream and even the surrounding coronal magnetic field structures may have changed the propagation directions of both CMEs (e.g., Gopalswamy et al. 2009; Lugaz et al. 2012; Zuccarello et al. 2012; Liu et al. 2014a; Kay et al. 2015; Möstl et al. 2015).

Given those relatively low inclination angles, the geomagnetic storm was mainly caused by the azimuthal magnetic field components of the flux ropes rather the axial components. A particularly interesting feature shown by the cross section of ICME1 is that the azimuthal field component (with a maximum value of about 30 nT) is much larger than the axial component (maximum value of only 8 nT). Therefore, if the flux rope were vertically orientated with the axis pointing southward the geomagnetic storm would be much weaker since the southward field would be less. This is contrary to the common belief that a southward pointing flux rope favors a strong geomagnetic storm. ICME2 has comparable axial and azimuthal magnetic field components. Another prominent feature is the vortices

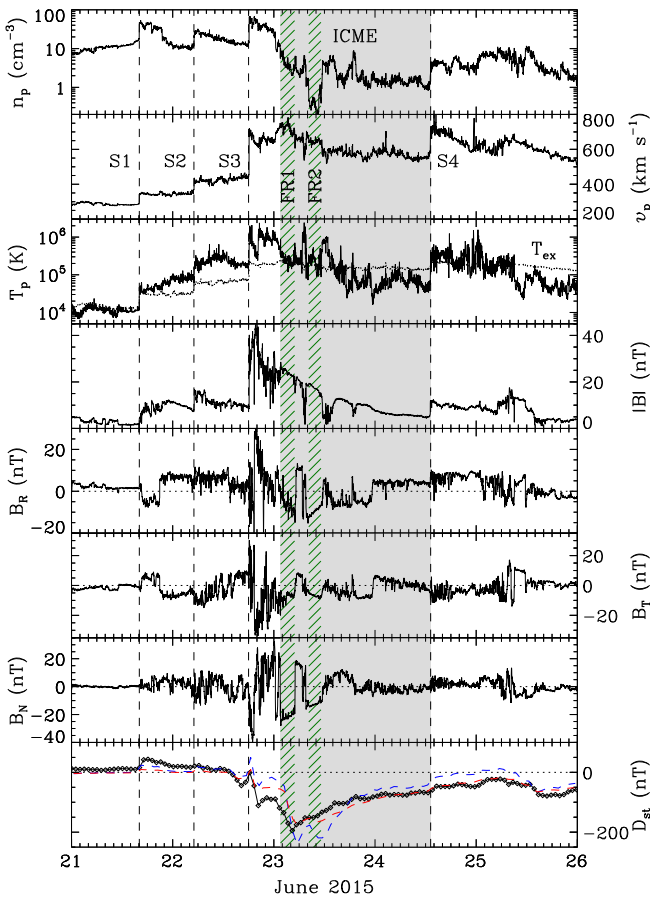


Figure 4. Solar wind measurements at *Wind* and associated D_{st} index for the 2015 June 22 event. Similar to Figure 2. The shaded region shows the overall ejecta interval, while the hatched areas indicate two small flux ropes identified within the ICME. Three shocks are observed ahead of the ejecta. The last shock (S4) was driven by the CME that occurred at the Sun on 2015 June 22 and was overtaking the ICME at 1 AU.

visible in the cross section of ICME2, reminiscent of the Kelvin–Helmholtz instability. This is probably a result of the fast stream interacting with ICME2 from behind (see Figure 2).

3. THE 2015 JUNE EVENT

In 2015 June AR 12371 exhibited elevated activity, somewhat similar to AR 11429 in 2012 March (Liu et al. 2013, 2014c; Wang et al. 2014; Sun et al. 2015). The active region produced a CME of about 1200 km s^{-1} associated with an M3.0 flare from $\text{N}13^\circ\text{E}45^\circ$ that peaked at 17:36 UT on June 18, a CME of about 1300 km s^{-1} associated with an M2.0 flare from $\text{N}12^\circ\text{E}13^\circ$ peaking at 01:42 UT on June 21, a CME of about 1000 km s^{-1} associated with an M6.5 flare from $\text{N}13^\circ\text{W}05^\circ$ peaking at 18:23 UT on June 22, and another one of about 1700 km s^{-1} associated with an M7.9 flare from $\text{N}10^\circ\text{W}42^\circ$ peaking at 08:16 UT on June 25. All these CMEs impacted the Earth. The June 21 CME appeared as a single halo event in the coronagraph images (see Figure 1, right), and a near head-on collision with the Earth was expected.

The corresponding in situ signatures at *Wind* are displayed in Figure 4. A cluster of shocks passed *Wind* at 16:05 UT on June 21, 05:02 UT and 18:00 UT on June 22, and 13:12 UT on June 24, respectively. The ICME boundaries are determined from the magnetic field in conjunction with the proton temperature and density. The first shock (S1) seemed driven by the June 18

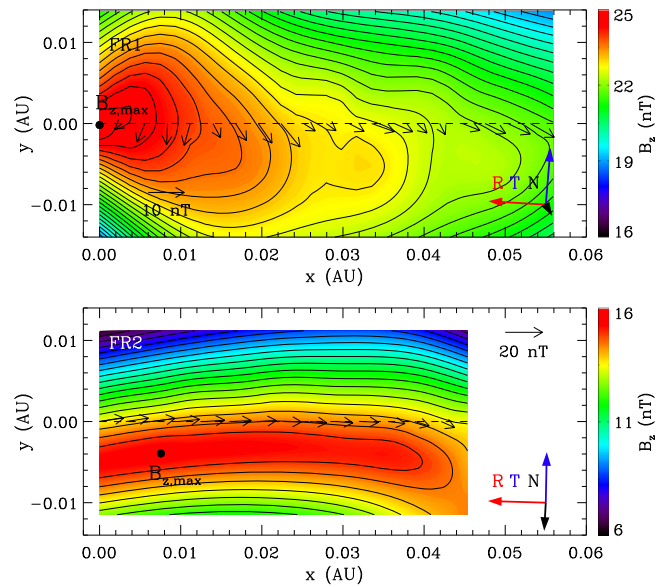


Figure 5. Reconstructed cross sections of FR1 (upper) and FR2 (lower) identified inside the 2015 June 22 ICME. Similar to Figure 3. These small flux ropes are both left-handed and have similar axis orientations.

CME, and the second one (S2) was likely associated with a CME from June 19. No driver signatures are observed at *Wind* for these two shocks, presumably owing to the largely eastward and southward propagation directions of the June 18 and 19 CMEs respectively (not shown here). The ICME and its preceding shock (S3) were produced by the June 21 CME; again, application of the empirical model (Gopalswamy et al. 2000) with the CME speed of 1300 km s^{-1} yields a predicted arrival time of 17:02 UT on June 22 at *Wind* (1.02 AU from the Sun), which is only 1 hr earlier than the observed S3 arrival. The fourth shock (S4) that was overtaking the ICME at 1 AU was associated with the June 22 CME. A series of dips in the magnetic field strength are observed inside the ICME, suggestive of the presence of current sheets. This signature is possibly due to the heliospheric current sheet cutting through the ejecta, which may lead to a chain of small flux ropes within the ICME (see below).

The D_{st} profile shows a multi-step geomagnetic storm with a global minimum of -195 nT . The first dip is produced by the fluctuating southward field component upstream of S3 (likely owing to amplification by the two preceding shocks), the second one by the southward field in the sheath downstream of S3 (further enhanced by S3), and the major dip by the southward field in the first hatched interval inside the ejecta. The southward field in the second hatched interval only creates a flattening of the D_{st} value, perhaps because of the extremely low density. Note that the solar wind density upstream of the ejecta is significantly enhanced by the three preceding shocks. This high density may feed the plasma sheet of the magnetosphere, which in turn helps intensify the ring current (Farrugia et al. 2006; Lavraud et al. 2006). Given the presence of more than one preceding shocks and a single ICME, the multi-step development of the geomagnetic storm can be classified as a sheath-sheath-ejecta scenario. Again, the modeled D_{st} index using the O’Brien & McPherron (2000) formula (minimum -174 nT) underestimates the global minimum but reproduces the recovery phase fairly well, whereas application of the Burton et al. (1975) scheme gives a larger global minimum (-241 nT) than measured.

Figure 5 shows the reconstructed cross sections of two small flux ropes identified inside the ICME. Here we call them small flux ropes rather than ICMEs, in order to distinguish from the 2015 March 17 case. Specifically, the June 22 event is a single ejecta instead of multiple ICMEs. These small flux ropes may have formed from the interaction between the CME and the heliospheric current sheet. The reconstructions yield a left-handed structure for both of the flux ropes and similar axis orientations: an elevation angle of about -77° and azimuthal angle of about 236° for the first one (FR1), and an elevation angle of about -61° and azimuthal angle of about 272° for the second one (FR2). The same chirality and similar axis orientations support the interpretation of a single ICME. Both flux ropes have strong axial magnetic field components compared with the azimuthal ones. Also note the largely southward orientation of the flux ropes. It is the strong axial field component on top of the largely southward flux-rope orientation, in conjunction with the relatively high solar wind speed, that may have resulted in the intense geomagnetic storm.

4. CONCLUSIONS

We have examined the sources of the 2015 March 17 and June 22 intense geomagnetic storms, the largest ones of solar cycle 24 up to the time of this writing. Key findings are obtained on how the plasma and magnetic field characteristics of CMEs control the geomagnetic storm intensity and variability.

1. A sheath-ejecta-ejecta mechanism and a sheath-sheath-ejecta scenario are proposed for the development of multi-step geomagnetic storms, based on the observed and reconstructed solar wind structures associated with the CMEs. The 2015 March 17 geomagnetic storm shows a two-step development, which is produced by the southward magnetic field components behind the preceding shock and those within two interacting CMEs. This falls into the sheath-ejecta-ejecta category. The 2015 June 22 geomagnetic storm exhibits a multi-step development, which is caused by the southward fields due to amplification by a series of preceding shocks and those within a single ejecta. This is classified as a sheath-sheath-ejecta scenario. The multiple preceding shocks and sheaths may precondition the magnetosphere for the growth of an intense geomagnetic storm.
2. We find two contrasting cases of how the CME flux-rope characteristics generate intense geomagnetic storms. Our GS reconstruction of the ejecta responsible for the 2015 June 22 geomagnetic storm indicates that the geomagnetic storm resulted from the largely southward flux-rope orientation with a strong axial magnetic field component. However, for the 2015 March 17 geomagnetic storm the GS reconstruction reveals a much larger azimuthal field component than the axial component. The intense geomagnetic storm occurred despite low flux-rope inclinations. A southward flux-rope orientation is thus not a necessity for a strong geomagnetic storm to occur.
3. The “perfect storm” scenario proposed by Liu et al. (2014a) may not be as rare as the phrase implies. The 2015 March 17 intense geomagnetic storm occurred in spite of the relatively weak solar flares and an encounter with the CME flank. What makes it an intense

geomagnetic storm is the interaction between two successive CMEs plus the compression by a high-speed stream from behind, which helps maintain strong ejecta magnetic fields and a relatively high speed. This is essentially the “perfect storm” scenario—a combination of circumstances results in an event of unusual magnitude, although the 2015 March 17 event is not “super” in the same sense as the 2012 July 23 solar storm. Note that there are many combinations of circumstances that can occur to make an event more geo-effective, including pileup of events, pre-event rarefactions and field line stretching, shock enhancement of southward fields, and following high-speed streams causing compressions. This “perfect storm” scenario now seems useful and necessary to worry about because complex events with these combinations are common.

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