Queen Mary
University of London

Truncations of random orthogonal matrices

Khoruzhenko, BA; Sommers, HJ; Zyczkowski, K

"The final publication is available at

http://journals.aps.org/pre/abstract/10.1103/PhysRevE.82.040106"

For additional information about this publication click this link.
http://qmro.qmul.ac.uk/xmlui/handle/123456789/13476

Truncations of Random Orthogonal Matrices

Boris A. Khoruzhenko,[1] Hans-Jürgen Sommers,[2] and Karol Życzkowski[3,4]

[1]*Queen Mary University of London, School of Mathematical Sciences, London E1 4NS, UK*
[2]*Fakultät für Physik, Universität Duisburg-Essen, 47048 Duisburg, Germany*
[3]*Smoluchowski Institute of Physics, Jagiellonian University, ul. Reymonta 4, 30-059 Kraków, Poland*
[4]*Centrum Fizyki Teoretycznej, Polska Akademia Nauk, Al. Lotników 32/44, 02-668 Warszawa, Poland*
(Dated: 12 August, 2010)

Statistical properties of non–symmetric real random matrices of size M, obtained as truncations of random orthogonal $N \times N$ matrices are investigated. We derive an exact formula for the density of eigenvalues which consists of two components: finite fraction of eigenvalues are real, while the remaining part of the spectrum is located inside the unit disk symmetrically with respect to the real axis. In the case of strong non–orthogonality, $M/N =$const, the behavior typical to real Ginibre ensemble is found. In the case $M = N - L$ with fixed L, a universal distribution of resonance widths is recovered.

Random unitary matrices [1] are used in numerous physical applications including chaotic scattering, conductance in mesoscopic systems [2] or statistical properties of periodically driven quantum systems [3]. In several applications one needs to restrict this class to real, orthogonal matrices, which are assumed to be distributed uniformly with respect to the Haar measure on the orthogonal group. This is the case while describing quasiparticle excitations in metals and superconductors [4, 5], or quantum maps performed on real quantum states [6]. In applications one is often led to consider square truncations of random unitary matrices of large matrix dimension N. These matrices have been used to describe quantum systems with absorbing boundaries [7] and found applications in various physical problems including optical and semiconductor superlattices [8], problems of quantum conductance [9], distribution of resonances for open quantum maps [10–13]. The operation of truncation does not preserve unitarity and as a consequence the eigenvalues move inside the unit disk in the complex plane, see, e.g., [3, 14], and the emerging eigenvalue statistics appear to be universal [10].

The aim of this work is to present a comprehensive study of truncations of *random orthogonal matrices.* We consider an ensemble of $N \times N$ real orthogonal matrices Q with flat matrix distribution on the orthogonal group. The goal of this paper is to compute analytically the full distribution of eigenvalues of the top $M \times M$ sub-matrix A of $Q = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$ for any values of M and N[15]. This then gives the eigenvalue correlation functions in a closed form and we investigate two asymptotic regimes of direct physical importance. For simplicity we consider M even, however, our results can be generalized to odd dimensions as well.

The orthogonality condition for the first M columns of Q, $A^T A + C^T C = 1$, defines a manifold in the space of real $N \times M$ matrices which can be identified with the flag manifold $O(N)/O(L)$ where $L = N - M$. Correspondingly, the density of distribution of A, C is

$$P(A, C) = \frac{\text{vol } O(L)}{\text{vol } O(N)} \, \delta(A^T A + C^T C - 1),$$

where $\text{vol } O(N) = 2^N \prod_{j=1}^{N} \pi^{j/2}/\Gamma(j/2)$ is the volume of the orthogonal group $O(N)$ [6]. The density function of A is $P(A, C)$ integrated over C. For $L \geq M$ the C-integral can be performed by scaling $C \rightarrow C\sqrt{1 - A^T A}$. The corresponding Jacobian is $|\det(1 - A^T A)|^{L/2}$ and from the δ-function one gets another factor $|\det(1 - A^T A)|^{-(M+1)/2}$. The remaining integral yields the volume of $O(L)/O(L - M)$, leading to[9]

$$P(A) = \frac{\text{vol } O(L) \, \text{vol } O(L)}{\text{vol } O(N) \, \text{vol } O(L-M)} \, \det(1 - A^T A)^{\frac{L-M-1}{2}} \quad (1)$$

If $L < M$ the probability density function of A contains singular (δ-function) terms due to finite mass of the boundary of the matrix ball $A^T A \leq 1$ which supports the distribution of A.

By setting $M = 1$ in Eq. (1) one recovers the distribution $P(A) \propto (1 - A^2)^{\frac{N-2}{2}}$ of a single entry of Q. Obviously, for N large $P(A)$ can be approximated by a Gaussian. This is also true for matrix blocks if the size of truncation $M \ll N$. Eigenvalue correlations in such a Gaussian regime are now accessible following the recent progress [16–20] made for the real Ginibre ensemble [21]. In this paper we investigate two non-Gaussian regimes: (i) strong non-orthogonality $M, L \propto N$; and (ii) weak non-orthogonality $M \propto N$ and $L \ll N$. In the former we again recover Ginibre eigenvalue correlations while the latter yields eigenvalue statistics from a different universality class [10].

We shall obtain the joint distribution of eigenvalues of A directly from $P(A, C)$ reducing the general case to $M = 2$. Thus, it is instructive to consider first $M = 2$. In this case, one can bring A by a rotation to the form $A = \begin{pmatrix} \lambda_1 & \delta \\ -\delta & \lambda_2 \end{pmatrix}$ with real $\lambda_{1,2}$ and δ from which the eigenvalues z_i of A can be recovered via the relations $\det A = \lambda_1 \lambda_2 + \delta^2$ and $\text{Tr } A = \lambda_1 + \lambda_2$ [18]. The matrix A can have either two real eigenvalues in which case we choose $z_1 > z_2$ or two complex conjugate eigenvalues in which case we choose $\text{Im } z_1 > \text{Im } z_2$. The joint distribution $d\mu(z_1, z_2)$ for $L \geq 2$ in each of these two cases can be obtained from Eq. (1) by making use of $\det(1 - A^T A) = (1 - z_1^2)(1 - z_2^2) - 4\delta^2$ and integrating out δ subject to the constraints $4\delta^2 \leq (1 - z_1^2)(1 - z_2^2)$ and $4\delta^2 \geq -(z_1 - z_2)^2$ arising due to the chosen parameterization of A, see [18], and the positivity of $1 - A^T A$. The resulting

expression can be presented in a form that encodes both cases (real and complex):

$$d\mu(z_1, z_2) = (z_1 - z_2)\, f(z_1)f(z_2)\, dz_1 \wedge dz_2 \qquad (2)$$

with $(z = x + iy)$

$$f^2(z) = \frac{L(L-1)}{2\pi}\, |1 - z^2|^{L-2} \int_{\frac{2|y|}{|1-z^2|}}^{1} du\, (1 - u^2)^{\frac{L-3}{2}}. \quad (3)$$

If the eigenvalues of A are real then the integral in (3) yields a Beta function. By integrating $d\mu(z_1, z_2)$ over the triangle $-1 < z_2 < z_1 < 1$ one gets the probability $p_{2,L}^R$ for A to have two real eigenvalues: $p_{2,L}^R = \mathrm{B}(L/2 + 1/2, L)/\mathrm{B}(L/2, L + 1/2)$. In the limit $L \to \infty$ this converges to $1/\sqrt{2}$ which, as expected, is the value for the real Ginibre ensemble.

One recovers $d\mu(z_1, z_2)$ for $L = 1$ by letting $L \to 1$ in (3). In this limit $f^2(z) = (2\pi|1 - z^2|)^{-1}$. As a simple check one can evaluate the probability $p_{2,L=1}^C$ for A to have two complex eigenvalues. In this case $z_1 = z_2^*$ so that $(z_1 - z_2)dz_1 \wedge dz_2 = 4y dx dy$ in (2). Integrating $d\mu(z_1, z_2)$ over the upper half of the unit disk $|z| < 1$ one gets $p_{2,1}^C = 1 - 2/\pi = 1 - p_{2,1}^R$, in agreement with the above result for $p_{2,L}^R$.

We now proceed with the general case of M even. By an orthogonal transformation $\mathcal{O}$ drawn from $O(M)/O(2)^{M/2}$ one can bring A to upper triangular form $A = \mathcal{O}(\Lambda + \Delta)\mathcal{O}^T$, where Λ is block diagonal with 2×2 blocks Λ_I in the diagonal and Δ is block triangular with 2×2 blocks Δ_{IJ} ($I < J$) above the diagonal. Integrating $P(A, C)$ over $\mathcal{O}$, Δ and C we obtain the distribution of Λ,

$$P(\Lambda) = (4\pi)^{-M/2}\, c_{M,L}\, \Delta'(\Lambda) \times \qquad (4)$$
$$\int d\Delta \int dC\, \delta((\Lambda^T + \Delta^T)(\Lambda + \Delta) + C^T C - 1)$$

with $c_{M,L} = \mathrm{vol} O(L) \mathrm{vol} O(M)/\mathrm{vol} O(L+M)$. Here $\Delta'(\Lambda) = \Delta(\Lambda)/\prod_I \Delta(\Lambda_I)$ is the Jacobian of the transformation from A to $\mathcal{O}$, Δ and Λ, see [23], and $\Delta(\Lambda) = \prod_{i<j}(z_i - z_j)$ is the Vandermonde determinant of the eigenvalues z_j of Λ. Thus Δ' does not count the pairs $i < j$ from the diagonal blocks.

Writing C as $C = [C_1, \ldots, C_{M/2}]$ in terms of $L \times 2$ subblocks C_I, we have for $I < J$ the equations

$$\Lambda_I^T \Delta_{IJ} + \sum_{K<I} \Delta_{KI}\Delta_{KJ} = -C_I^T C_J \qquad (5)$$

due to the δ-function in (4), which start with $\Lambda_1^T \Delta_{1J} = -C_1^T C_J$ for $J \geq 2$, $\Lambda_2^T \Delta_{2J} + \Delta_{12}^T \Delta_{2J} = -C_2^T C_J$ for $J \geq 3$, etc. Because of the triangular structure of Δ these equations can iteratively be solved for Δ_{IJ}. Thus integration over Δ_{IJ} can be performed with a Jacobian $\prod_{I<J} |\det \Lambda_I|^{-2}$. There remains integration over C_J:

$$\int \prod_{J=1}^{M/2} [dC_J\, \delta(\Lambda_J^T \Lambda_J + \sum_{K<I} \Delta_{KI}\Delta_{KJ} + C_J^T C_J - 1)].$$

From Eqs. (5) one sees that Δ_{IJ} for $I < J$ is linear in C_J and depends otherwise on C_K for $K < J$. Thus we have an

integral $\int \prod_{J=1}^{M/2} dC_J\, \delta(\Lambda_J^T \Lambda_J + C_J^T X_J C_J - 1)$ where X_J is a linear operator acting on C_J per matrix multiplication. The scaling $C_J \to (1/\sqrt{X_J})C_J$ results in a Jacobian that cancels the Jacobian $\prod_{I<J} |\det \Lambda_I|^{-2}$ from the previous integration over Δ_{IJ} and we end up with

$$\frac{P(\Lambda)}{c_{M,L}} = \frac{\Delta'(\Lambda)}{(4\pi)^{M/2}} \int \prod_{J=1}^{M/2} dC_J\, \delta(\Lambda_J^T \Lambda_J + C_J^T C_J - 1).$$

where Λ_J's are 2×2 and C_J's are $L \times 2$. Now we proceed calculating $P(\Lambda)$ using our result for $M = 2$. Performing integration over C_J yields

$$\frac{P(\Lambda)}{c_{M,L}} = \frac{\Delta'(\Lambda)}{(4\pi)^{M/2}} \prod_{J=1}^{M/2} \frac{\mathrm{vol}\, O(L)}{\mathrm{vol}\, O(L-2)}\, \det(1 - \Lambda_J^T \Lambda_J)^{\frac{L-3}{2}},$$

and we obtain the joint (full) distribution of eigenvalues of A:

$$d\mu(z_1, \ldots, z_M) = \mathcal{Z}_{M,L}^{-1} \prod_{1 \leq i < j \leq M} (z_i - z_j) \prod_{j=1}^{M} f(z_j)\, dz_j \quad (6)$$

with $\mathcal{Z}_{M,L}^{-1} = [L!/(2\pi)^L]^{M/2} c_{M,L}$, $dz_j = dx_j + idy_j$ and ordering $z_1 > z_2 > \ldots > z_M$ if all eigenvalues are real, $z_1 = z_2^*$, $\mathrm{Im}\, z_1 > 0$, $z_3 > z_4 > \ldots > z_M$ if two eigenvalues are complex conjugate and so on. Eqs. (6), (3) represent our first main result.

Eq. (6) is of the same form as the corresponding one in the real Ginibre ensemble [18] but with a different weight function $f(z)$. Following [18] one can obtain the eigenvalue correlation densities in terms of a skew-symmetric kernel

$$\mathcal{K}(z_1, z_2) = \sum_{k=1}^{M} \sum_{l=1}^{M} (\mathcal{A}^{-1})_{kl} z_1^{k-1} z_2^{l-1}$$

where $\mathcal{A}_{kl} = \int \int d^2 z_1 d^2 z_2\, \mathcal{F}(z_1, z_2)\, z_1^{k-1} z_2^{l-1}$ with $d^2 z = dx dy$ and

$$\mathcal{F}(z_1, z_2) = f(z_1)f(z_2) \times$$
$$[2i\, \delta^{(2)}(z_1 - z_2^*)\, \mathrm{sgn}(y_1) + \delta(y_1)\delta(y_2)\, \mathrm{sgn}(x_2 - x_1)],$$

with $\delta^{(2)}(z_1 - z_2^*) = \delta(x_1 - x_2)\delta(y_1 + y_2)$. For example, the one-point density is given by an integral

$$R_1(z) = \int d^2 z_2\, \mathcal{F}(z, z_2)\mathcal{K}(z_2, z). \qquad (7)$$

Higher order densities R_n are given by a Pfaffian involving $\mathcal{F}$ and $\mathcal{K}$ [18]. Expressing the kernel $\mathcal{K}$ in terms of skew-orthogonal polynomials with respect to weight function $\mathcal{F}$ leads to an alternative Pfaffian representation for R_n [16, 17].

Evaluating the kernel $\mathcal{K}$ (or, equivalently [24], the corresponding skew-orthogonal polynomials) in a closed form is an important step on the way to obtaining eigenvalue statistics [17]. In our case the kernel can be found by exploiting its relation to averages of the characteristic polynomials [18, 25] of truncations of orthogonal matrices of smaller dimension:

$$\mathcal{K}(z_1, z_2) = (z_1 - z_2) \langle \det(z_1 - \tilde{A})(z_2 - \tilde{A}) \rangle_{\tilde{A}}$$

where the matrices $\tilde{A}$ are square truncations of size $M - 2$ of random orthogonal matrices of size $N - 2$. Due to the invariance of the distribution of $\tilde{A}$ the above average is effectively an average over the eigenvalues of $\tilde{A}^T\tilde{A}$:

$$\langle \det(z_1 - \tilde{A})(z_2 - \tilde{A})\rangle_{\tilde{A}} = \sum_{m=0}^{n}(z_1 z_2)^m \frac{\langle \epsilon_{n-m}(\tilde{A}^T\tilde{A})\rangle_{\tilde{A}}}{\epsilon_{n-m}(1)},$$

where $n = M - 2$ is the size of $\tilde{A}$ and $\epsilon_m(X)$ is the m-th elementary symmetric function in eigenvalues of X. The integral $\langle \epsilon_m(\tilde{A}^T\tilde{A})\rangle_{\tilde{A}}$ can be reduced [26] to Selberg's integral [29], yielding the kernel in a closed form,

$$\mathcal{K}(z_1, z_2) = (z_1 - z_2)\sum_{m=0}^{M-2} \frac{(L+m)!}{L!\,m!}(z_1 z_2)^m, \quad (8)$$

which is our second main result. The truncated binomial series on the rhs can be expressed [23] in terms of an incomplete beta function which comes in handy for asymptotic analysis of eigenvalue statistics in the limit of large M.

The one-point eigenvalue density (7) is composed of two parts: $R_1(z) = \rho_C(z) + \delta(y)\rho_R(x)$, where $\rho_C(z) = 2f(z)^2|\mathcal{K}(z, z^*)|$ is the density of complex eigenvalues and

$$\rho_R(x) = \int_{-1}^{1} dx_2\, \text{sgn}(x_2 - x)\mathcal{K}(x_2, x)f(x_2)f(x) \quad (9)$$

is the density of real eigenvalues of truncated orthogonal matrices, with the normalization $\int_{|z|<1} R_1(z)\, d^2 z = M$.

We shall first look at real eigenvalues of A. Integrating $\rho_R(x)$ over $-1 < x < 1$ one obtains the expected total number N_R of real eigenvalues,

$$N_R = 1 + \frac{L}{2}\int_0^1 \frac{ds}{s^{L+1}}\, I_{s^2}\left(L/2, 1/2\right)\, I_{\frac{2s}{1+s}}\left(L+1, M-1\right),$$

where $I_x(a, b) = \int_0^x t^{a-1}(1-t)^{b-1}dt/B(a, b)$ is an incomplete Beta function. For large matrix dimensions ($M \gg 1$) N_R, in the leading order, is described by a simpler expression:

$$N_R \simeq \frac{2\,\text{artanh}\,\sqrt{M/N}}{B(L/2, 1/2)}.$$

Depending on L this leads to different scaling laws for N_R. In the limit of strong non-orthogonality when both $M, L \propto N$, the number of real eigenvalues $\propto \sqrt{M}$ which is characteristic of the real Ginibre ensemble [27]. On the other hand, in the limit of weak non-orthogonality when $L \ll M$, N_R grows logarithmically with M, $N_R \simeq (\log M)/B(L/2, 1/2)$.

Eq. (9) can be transformed to a form,

$$\rho_R(x) = \frac{1}{B(L/2, 1/2)}\frac{I_{1-x^2}(L+1, M-1)}{1 - x^2} +$$
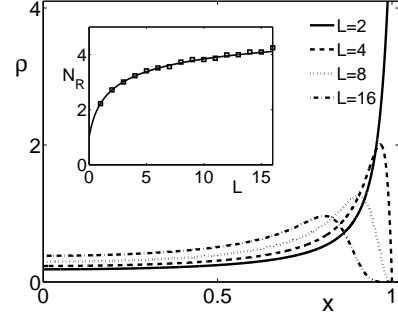$$\frac{(1-x^2)^{\frac{L-2}{2}}|x|^{M-1}}{B(M/2, L/2)}I_{x^2}((M-1)/2, (L+2)/2),$$

FIG. 1. Rescaled density of real eigenvalues $\rho = \frac{\rho_R}{N_R}$ for $M = 32$ and $L = 2, 4, 8, 16$. The inset shows the number of real eigenvalues N_R for $M = 32$ as a function of L obtained from samples of 10^3 realizations (squares) and analytical prediction (solid line)

revealing the Artanh Law of the distribution of real eigenvalues for large matrix dimensions M ($\mu = M/N$):

$$\rho_R(x) \simeq \frac{1}{B(L/2, 1/2)}\frac{1}{1 - x^2} \quad -\sqrt{\mu} < x < \sqrt{\mu}.$$

In the limit of strong non-orthogonality $\mu < 1$ and $\rho_R(x)$ is supported strictly inside the interval $(-1, 1)$, vanishing at a Gaussian rate at the boundaries $x = \pm\sqrt{\mu}$. In the limit of weak non-orthogonality $\mu \to 1$ and $\rho_R(x)$ becomes singular at the boundaries $x = \pm 1$ of the real eigenvalue support. Correspondingly, real eigenvalues accumulate close to $x = \pm 1$. Setting $x = 1 - u/M$ and taking the limit $M \to \infty$ one obtains the density profile near the accumulation point $x = 1$: $\rho_R(x)/M \simeq p(u)$ with

$$p(u) = \frac{u^{\frac{L}{2}-1}e^{-u}}{2\Gamma(L/2)}(1 - u^{\frac{L}{2}+1}\gamma^*(L/2+1, u)) +$$
$$\frac{(2u)^L}{B(L/2, 1/2)}\gamma^*(L+1, 2u),$$

where $\gamma^*(n, x) = (x^n\Gamma(n))^{-1}\int_0^x t^{n-1}e^{-t}dt$ is an incomplete Gamma function. For small u, $p(u) \simeq (2\Gamma(\frac{L}{2}))^{-1}u^{\frac{L}{2}-1}$, showing a transition, as L increases, in nature of the boundary points $x = \pm 1$ from 'attractive' ($L = 1$) to 'repulsive' ($L \geq 4$), see Fig. 1. On the other hand, for large u, $p(u) \simeq (B(L/2, 1/2)u)^{-1}$, exhibiting a heavy tail which manifest itself in the $\log M$ asymptotics for N_R.

Now we shall look at complex eigenvalues. Their density $\rho_C(z)$, visualized in Fig. 2, vanishes on the real line. In the limit of strong non-orthogonality, close to the real line the complex density is described by the same scaling law

$$\rho_C(z) \simeq \rho_R(x)^2 h(y\rho_R(x)), \quad h(y) = 4\pi|y|e^{4\pi y^2}\text{erfc}(\sqrt{4\pi}|y|),$$

as in the real Ginibre ensemble [28], while away from the real axis $\rho_C(z)$ is the same as for truncations of random *unitary* matrices [22], $\rho_C(z) \simeq \frac{L}{\pi}\frac{1}{(1-|z|^2)^2}\Theta(\frac{M}{N} - |z|^2)$. In the limit of weak non-orthogonality away from the real axis

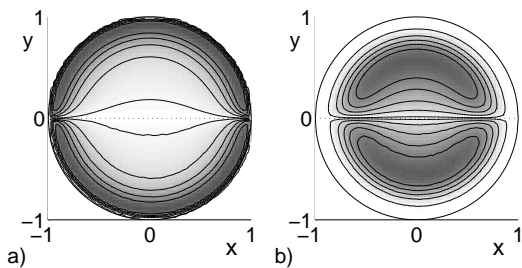$$\rho_C(z) \simeq \nu^2 h(4\pi\nu(1 - |z|)), \quad h(u) = 4\pi L u^{L-1}\gamma^*(L+1, u),$$

FIG. 2. Contour plot of the normalized density of complex eigenvalues $\rho_C(z)/(1-N_R)$ of truncations of random orthogonal matrices for a) $N=12$ and $M=10$ (weak non-orthogonality) and b) $N=12$ and $M=6$ (strong nonorthogonality) with the larger values coded darker. The density plotted in the upper half of the complex planes represents analytical results, while the density in the lower part is obtained numerically for a sample of 0.5×10^7 random orthogonal matrices.

where $\nu = \frac{M}{2\pi}$ is the density of distribution of eigenvalues of random orthogonal matrices along the unit circle. Again, the limiting density profile does not depend on the angle and is the same as for truncations of random unitary matrices [22].

Eigenvalue correlations for truncations of random orthogonal matrices can also be obtained in a closed form. In the limit of strong non-orthogonality, after appropriate rescaling the eigenvalue correlations become identical to those in the real Ginibre ensemble. For example, at the origin, $f^2(z/\sqrt{L}) \simeq \sqrt{L/(8\pi)}e^{y^2-x^2}\,\mathrm{erfc}(\sqrt{2}|y|)$ and $K(z_1/\sqrt{L}, z_2/\sqrt{L}) \simeq \frac{1}{\sqrt{L}}(z_1-z_2)e^{z_1 z_2}$ which is what one gets in the real Ginibre ensemble [17–19]. The limit of weak non-orthogonality away from the real axis the eigenvalue correlations for truncations of random orthogonal matrices are exactly the same as those found for truncated random unitary matrices [22, 23]. New correlation laws arise in the vicinity of $x = \pm 1$.

TABLE I. Ensembles of random nonhermitian matrices

matrices	complex	real	$\mu = M/N$
Haar measure	$U(N)$	$O(N)$	$\mu = 1$
truncations of matrices	unitary	orthogonal	$1 > \mu > 0$
Ginibre ensemble	complex	real	$\mu \to 0$

In conclusion, we have found the full probability distribution of eigenvalues of truncated random orthogonal matrices and obtained the eigenvalue density and higher order correlation functions in a closed form. This work completes our understanding of ensembles of non-hermitian random matrices which are derived from random unitary and orthogonal matrices. As shown in Table I truncations of random orthog-

onal (unitary) matrices form an ensemble which interpolates between matrices distributed according to the Haar measure on the orthogonal (unitary) group [29] and the real (complex) Ginibre ensemble [21].

Note added. After submitting our work we learned about a very recent preprint of Forrester [30] extending our study

[1] F. J. Dyson, J. Math. Phys. **3**, 140 (1962).
[2] C. W. J. Beenakker, Rev. Mod. Phys. **69**, 731 (1997).
[3] F. Haake, *Quantum Signatures of Chaos, 3 ed.* (Springer, 2010).
[4] A. Altland and M. R. Zirnbauer, Phys. Rev. B **55**, 1142 (1997).
[5] J.P. Dahlhaus, B. Béri and C.W.J. Beenakker, Phys. Rev. **B 82**, 014536 (2010).
[6] I. Bengtsson and K. Życzkowski, *Geometry of Quantum States* (Cambridge University Press, Cambridge, 2006).
[7] G. Casati, G. Maspero, and D. L. Shepelyansky, Phys. Rev. Lett. **82**, 524 (1999).
[8] M. Glück, A. R. Kolovsky, and H. J. Korsch, Phys. Rep. **366**, 103 (2002).
[9] P. J. Forrester, J. Phys. A **39**, 6861 (2006).
[10] Y. V. Fyodorov and H.-J. Sommers, J. Phys. A **36**, 3303 (2003).
[11] H. Schomerus and P. Jacquod, J. Phys. A **38**, 10663 (2005).
[12] S. Nonnenmacher and M. Zworski, Comm. Math. Phys. **269**, 311 (2007).
[13] J. M. Pedrosa, G. G. Carlo, D. A. Wisniacki, and L. Ermann, Phys. Rev. E **79**, 016215 (2009).
[14] P. J. Forrester and M. Krishnapur, J. Phys. A **42**, 385204 (2009).
[15] This distribution depends on the size M of the truncation but not on its position in view of the invariance of the Haar measure.
[16] E. Kanzieper and G. Akemann, Phys. Rev. Lett. **95**, 230201 (2005).
[17] P. J. Forrester and T. Nagao, Phys. Rev. Lett. **99**, 050603 (2007); J. Phys. **A41** 375003 (2008).
[18] H.-J. Sommers, J. Phys. A **40**, F671 (2007); H.-J. Sommers and W. Wieczorek, J. Phys. A **41**, 405003 (2008).
[19] A. Borodin and C.D. Sinclair, Commun. Math. Phys. **291** 177 (2009).
[20] P.J. Forrester and A. Mays, arXiv:0910.2531.
[21] J. Ginibre, J. Math. Phys. **6**, 440 (1965).
[22] K. Życzkowski and H.-J. Sommers, J. Phys. A **33**, 2045 (2000).
[23] B. A. Khoruzhenko and H.-J. Sommers, arXiv:0911.5645.
[24] G. Akemann, M. Kieburg, and M.J. Phillips, J. Phys. **A43**, 375207 (2010).
[25] G. Akemann, M.J. Phillips and H.-J. Sommers, J.Phys. **A42**, 012001 (2009).
[26] Y.V. Fyodorov and B.A. Khoruzhenko, J.Phys. **A40**, 669 (2007).
[27] A. Edelman, E. Kostlan, and M. Shub, J. Am. Math. Soc., **7**, 247 (1994).
[28] A. Edelman, J. Multivariate Anal. **60**, 203 (1997).
[29] M. Mehta, *Random Matrices, 3rd ed.* (Academic Press, 2004).
[30] P.J. Forrester, arXiv:1009.3066.