

EFFECT OF IMPURITIES ON CO₂ STREAM PROPERTIES

by

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ABSTRACT

CO₂ obtained by capture process (such as post combustion, pre combustion and oxy-fuel combustion) is not 100% pure and may contain impurities such as H₂, Ar, CO, H₂S and water. The presence of such impurities in CO₂ stream can lead to challenging flow assurance and processing issues. The gaseous CO₂-rich stream is generally compressed to be transported as liquid in order to avoid two-phase flow and increase the density of the system. One aim of this work is to evaluate the effect of impurities on the physical properties of CO₂ such as density, viscosity, speed of sound and on the phase behaviour of such systems. Speed of sound and isothermal compressibility of CO₂/impurities mixtures were measured at condition above the saturation curve and temperature from 268.15 to 301.15 K. A new volume correction was implemented to the Peng-Robinson equation of state in order to minimise the error associated with the isothermal compressibility prediction. Moreover, density and viscosity are two of the most important properties in transport properties. Therefore, the effect of impurities on density and viscosity were experimentally and theoretically investigated in liquid CO₂ and liquid CO₂/impurities systems. The viscosity measurements were performed using in-house capillary tube apparatus in the range from 280 to 343.15 K and pressure up to 40 MPa. Two viscosity models, LBC and Pedersen, were modified in order to predict the viscosity of both pure and impure CO₂. The density measurements were carried-out using an Anton Paar densitometer in both liquid and supercritical regions from 283.15 to 423.15 K. In order to improve the accuracy of EOSs in density of CO/impurities systems, a new modification was developed based on mixing the volume obtained from EOSs (SRK, PR and VPT) and the volume obtained from CO₂-MBWR.

The presence of water may result in ice and/or gas hydrate formation and cause blockage of pipelines. Several measurements were also conducted to evaluate the hydrate stability zone of pure and rich CO₂ systems in free water. A thermodynamic model based on the VPT EOS was adopted to predict the hydrate phase of the systems. In addition, few saturation measurements of synthetic alkane mixture plus pure or impure CO₂ were performed at 344.3 K in order to investigate the effect of impurities on the saturation pressure of CO₂/alkane system. IFT and swelling factor properties on CO₂/n-decane mixture were investigated at 310.95 K from ambient to near the minimum miscibility pressure of the mixture. The experiments were extended to cover the presence of impurities on the properties at the same range of pressure. Minimum

miscibility pressure of the systems was estimated by both Vanishing Interfacial Tension method and multiple-mixing-cell calculation.

DEDICATION

This work is dedicated to my parents and my wife Maimoona.

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In the Name of Allah, the Most Gracious, the most Merciful

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TABLE OF CONTENTS

TABLE OF CONTENTS	i
LIST OF FIGURES	iv
LIST OF TABLES	viii
LIST OF MAIN SYMBOLS	xi
Chapter 1: INTRODUCTION	1
Chapter 2: VLE OF CO₂/IMPURITIES+ SATURATION PRESSURE AND SPEED OF SOUND OF CO₂/ALKANE MIXTURES	9
2.1 Introduction	9
2.2 VLE of CO₂ and CO₂/Impurities System	10
<i>2.2.1 Critical Behaviour of CO₂ Rich Fluid</i>	<i>21</i>
2.3 Bubble Point and Speed of Sound Measurements of CO₂ and Alkanes System	23
2.4 Conclusions	32
Chapter 3: SPEED OF SOUND AND ISOTHERMAL COMPRESSIBILITY OF CO₂-RICH SYSTEMS	41
3.1 Introduction	41
3.2 Experimental Section	41
<i>3.2.1 Experimental Apparatus</i>	<i>41</i>
<i>3.2.2 Materials</i>	<i>42</i>
<i>3.2.3 Procedures</i>	<i>43</i>
3.3 Modelling Section	43
3.4 Results	46
<i>3.4.1 Speed of Sound</i>	<i>46</i>
<i>3.4.2 Isothermal Compressibility</i>	<i>56</i>
3.5 Conclusions	60
Chapter 4: DENSITY OF CARBON DIOXIDE RICH SYSTEMS	63
4.1 Introduction	63

4.2	Experimental Methods and Equipment Materials	64
4.2.1	<i>Experimental Equipment.....</i>	66
4.2.2	<i>Experimental Calibration and Procedure.....</i>	67
4.3	Modelling	68
4.4	Results & Discussion.....	71
4.4.1	<i>Density of Liquid CO₂/impurities System.....</i>	71
4.4.2	<i>Density of Supercritical CO₂/Impurities System.....</i>	80
4.5	Conclusions	83
Chapter 5: HYDRATE STABILITY OF CARBON DIOXIDE RICH SYSTEMS.....		87
5.1	Introduction.....	87
5.2	Experimental Section.....	90
5.2.1	<i>Materials.....</i>	90
5.2.2	<i>Apparatus.....</i>	90
5.2.3	<i>Procedures.....</i>	91
5.3	Modelling Section.....	92
5.4	Results & Discussion.....	93
5.5	Conclusions.....	101
Chapter 6: VISCOSITY OF CARBON DIOXIDE RICH SYSTEMS.....		112
6.1	Introduction.....	112
6.2	Experimental Methods and Equipment.....	114
6.2.1	<i>Materials.....</i>	114
6.2.2	<i>Experimental Equipment.....</i>	115
6.2.3	<i>Experimental Procedure.....</i>	117
6.3	Modelling	118
6.3.1	<i>CO₂-Pedersen Model.....</i>	118
6.3.2	<i>Lohrenz–Bray–Clark (LBC) model.....</i>	123
6.4	Results and Discussion.....	125
6.4.1	<i>Experimental Results.....</i>	125
6.4.2	<i>Modelling Results.....</i>	140
6.5	Conclusions.....	143

Chapter 7: IFT, SWELLING FACTOR AND MMP OF CO₂/N-DECANE SYSTEM	148
7.1 Introduction	148
7.2 Experimental Section	150
7.2.1 <i>Materials</i>	150
7.2.2 <i>Equipment and procedures</i>	151
7.3 Results and Discussion	153
7.3.1 <i>IFT</i>	153
7.3.2 <i>Swelling factor</i>	157
7.3.3 <i>MMP</i>	159
7.3.3.1 <i>Vanishing Interfacial Tension Approach (VIT)</i>	159
7.3.3.2 <i>Multiple Mixing Cell Algorithm for predicting the MMP</i>	161
7.4 Conclusions	168
Chapter 8: CONCLUSIONS AND RECOMMENDATIONS FOR FUTURE WORK.....	174
8.1 Introduction	174
8.2 Conclusions from Literature Survey	176
8.3 Experimental Apparatus	177
8.4 Modelling Results	178
8.5 The Effect of Impurities on CO₂ Properties & Future Work Recommendations	179

LIST OF FIGURES

Figure 1.1 Various operations involved in the CO ₂ transport	3
Figure 2.1 AAD Deviations from PR EOS on CO ₂ /N ₂ mixture at 270 K, data from [23]	14
Figure 2.2 AAD Deviations from SRK EOS on CO ₂ /N ₂ mixture at 270 K, data from [23]	15
Figure 2.3 AAD Deviations from VPT EOS on CO ₂ /N ₂ mixture at 270 K, data from [23]	15
Figure 2.4 (a-p) AAD of PR (▲Δ), VPT (●○) and SRK (■□) EOSs on the calculated VLE properties of binary CO ₂ /mixtures; empty symbols represent AAD on vapour fraction of CO ₂ (Y _{CO₂}) and full symbols represent AAD on saturation pressure (Sources are listed in Table 2.1).	19
Figure 2.5 (a-h) prediction results on VLE of CO ₂ mixtures by PR EOS	20
Figure 2.6 (a-f) critical properties of CO ₂ /CH ₄ (a and b), CO ₂ /H ₂ S (c and d), CO ₂ /N ₂ (e and f); SRK EOS is used for modelling	23
Figure 2.7 Critical lines of the binary CO ₂ mixtures. T _{c_m} and P _{c_m} are the critical temperature and pressure of mixtures, respectively. T _{c_{CO₂}} and P _{c_{CO₂}} are the critical temperature and pressure of CO ₂ respectively	23
Figure 2.8 (a-c) Plots showing an example of bubble point determination of pure CO ₂ at 344.3 from plot of change in cell pressure versus (a) wave travelling time, (b) piston displacement and (c) speed of sound	25
Figure 2.9 Saturation pressure measurements of CO ₂ and n-decane system at 344.3 K	26
Figure 2.10 Effect of 5 mole% impurities on the saturation pressure of CO ₂ /alkane mixture	27
Figure 2.11 The speed of sound of CO ₂ /alkane system at 344.3 K, data are correlated using Equation 2.27	31

Figure 2.12 (a-b) Effect of impurities on speed of sound property of CO ₂ /alkane system at 344.3 K, the SoS of CO ₂ /alkane mixture is correlated using Equation 2.27 at the same condition of temperature and mixture composition. _____	31
Figure 3.1 Schematic diagram of the ultrasonic set-up _____	42
Figure 3.2 Liquid and saturation liquid density of CO ₂ _____	46
Figure 3.3 Measurements in speed of sound of CO ₂ _____	52
Figure 3.4 Effect of CO ₂ impurities on speed of sound property at 268.15 K _____	54
Figure 3.5 Effect of CO ₂ impurities on speed of sound property at 301.15 K _____	54
Figure 3.6 Change in speed of sound (defined as $SoS_{pureCO_2} - SoS_{impureCO_2}$) as a function of pressure and temperature in CO ₂ /CH ₄ mixture _____	55
Figure 3.7 Change in speed of sound (defined as $SoS_{pureCO_2} - SoS_{impureCO_2}$) as a function of pressure and temperature in MIXb (CO ₂ -CH ₄ -H ₂ -N ₂) _____	55
Figure 3.8 Isothermal compressibility in pure CO ₂ (measurements & modelling), error bar 5% _____	57
Figure 3.9 Isothermal compressibility in MIXb (CO ₂ +Ar+CO) (measurements & modelling), error bar 5% _____	58
Figure 3.10 Effect of three selected impurities systems on CO ₂ stream at 268.15 K _____	59
Figure 3.11 Effect of three selected impurities systems on CO ₂ stream at 301.15 K _____	59
Figure 4.1 Schematic diagram of the compressed liquid densimeter apparatus. PP, Ruska motorized HP pump; PS, pressure sensor; SC, sample cylinder; VP, vacuum pump system; TB, thermostatic bath; EU, mPDS 2000V3 unit; NV, nanovolt meter; V, valves _____	66
Figure 4.2 Deviation between PR EOS and the data provided in [23] for pure CO ₂ _____	69
Figure 4.3 Deviation between VPT EOS and the data provided in [23] for pure CO ₂ _____	69

Figure 4.4 Density of CO ₂ /CH ₄ at 283.15 K	78
Figure 4.5 Density of CO ₂ /CH ₄ at 301.15 K	78
Figure 4.6 Density reduction in liquid CO ₂ in the presence of impurities at 283.15 K	79
Figure 4.7 Density reduction in supercritical CO ₂ in the presence of impurities at 323.15 K	83
Figure 5.1 Schematic of hydrate mixed autoclave rig	91
Figure 5.2 Hydrate phase equilibria of the binary CO ₂ + water mixtures	94
Figure 5.3 Hydrate phase equilibria of the ternary CO ₂ + N ₂ + water mixtures	97
Figure 5.4 Hydrate phase equilibria of the ternary CO ₂ + H ₂ + water mixtures	98
Figure 5.5 Hydrate phase equilibria of the ternary CO ₂ + CH ₄ + water mixtures	98
Figure 5.6 Hydrate phase equilibria of the ternary CO ₂ + CO + water mixtures	99
Figure 5.7 Hydrate phase equilibria of the ternary CO ₂ + Ar + water mixtures	99
Figure 5.8 Hydrate phase equilibria of the ternary CO ₂ + O ₂ + water mixtures	100
Figure 5.9 Effects of impurities on the hydrate phase equilibria of CO ₂	100
Figure 6.1 schematic view of the viscosity experiments setup	116
Figure 6.2 Our viscosity data of pure CO ₂ together with some literature data at three isothermal conditions	137
Figure 6.3 The effect of Ar and H ₂ on CO ₂ viscosity at 280 K	137
Figure 6.4 The effect of Ar and H ₂ on CO ₂ viscosity at 288.15 K	138
Figure 6.5 The effect of Ar and H ₂ on CO ₂ viscosity at 300.15 K	138
Figure 6.6 The effect of Ar and H ₂ on CO ₂ viscosity at 308.15 K	139
Figure 6.7 The effect of Ar and H ₂ on CO ₂ viscosity at 323.15 K	139

Figure 6.8 The effect of Ar and H ₂ on CO ₂ viscosity at 343.15 K _____	140
Figure 6.9 Deviation of the CO ₂ -Pedersen model from the experimental results for pure CO ₂ _____	142
Figure 6.10 Deviation of the LBC model from the experimental results for pure CO ₂ ____	142
Figure 7.1 Schematic of capillary rise technique used for CO ₂ /n-decane system _____	152
Figure 7.2 Images of CO ₂ /n-decane at 310.95 K, shown that the capillary rise reduces as the system approaches the MMP by increasing the injection gas pressure _____	153
Figure 7.3 Experimental data for CO ₂ /n-decane system at 310.95 K, literature data [31]	154
Figure 7.4 The cell volume Calibration in terms of the vertical distance (height) _____	158
Figure 7.5 Swelling factor measurements in CO ₂ and CO ₂ /impurities systems with n-decane at 310.95 K _____	159
Figure 7.6 Determination of VIT Miscibility in n-Decane-CO ₂ System and n-Decane impure CO ₂ System at 310.95 K _____	160
Figure 7.7 The Effect of impurities on MMP of CO ₂ /n-decane system at 310.95 K ____	161
Figure 7.8 Multiple-mixing cell mechanism (constant volume cells, moving-excess-oil option) proposed in [25] _____	164
Figure 7.9 Recovery factor vs. Pressure for CO ₂ and CO ₂ + impurities with n-decane at 310.95 K _____	165
Figure 7.10 Recovery Factor (RF) vs. pressure supposed, using 50 cells. MMP predicted (1): using the range of RF (0.1-0.95) and MMP predicted (2): using the range of RF (0.7:0.99), for CO ₂ +nC ₁₀ at 310.95 K _____	166

LIST OF TABLES

Table 1.1 Impurities expected in CO ₂ stream from capture processes [1-2]	2
Table 2.1 Literature data for CO ₂ /mixtures	13
Table 2.2 BIPs for different binary CO ₂ /mixtures	14
Table 2.3 AAD of EOSs on VLE data of binary CO ₂ mixtures	16
Table 2.4 Alkane mixture compositions	24
Table 2.5 SoS measurements of liquid alkane mixture in pure and impure CO ₂	28
Table 3.1 Synthetic gas composition of CO ₂ /impurities mixtures	43
Table 3.2 Speed of sound and isothermal compressibility measurements of CO ₂ and CO ₂ /impurities system	47
Table 4.1 The Composition of mixtures used for the experiments	65
Table 4.2 Mixture composition (MIX1)	65
Table 4.3 The interaction parameters of CO ₂ /binary system used in VPT EOS	70
Table 4.4 Experimental and modelling data for liquid CO ₂ and liquid CO ₂ /impurities system	72
Table 4.5 Experimental and modelling data for supercritical CO ₂ and supercritical CO ₂ /impurities system	80
Table 5.1 Literature data for hydrate equilibria of CO ₂ and CO ₂ mixtures	89
Table 5.2 List of impurities composition used in the experiments	90
Table 5.3 VLE and BP data for binary CO ₂ /impurities system	92
Table 5.4 Equilibrium hydrate dissociation pressure for the CO ₂ /impurities mixtures	95

Table 6.1 Source of viscosity data for CO ₂ /impurities system _____	113
Table 6.2 The Composition of mixtures used for the experiments _____	115
Table 6.3 The mixture composition (MIX1) _____	115
Table 6.4 Values of coefficients, a _i , in Equation 6.6 for CO ₂ _____	120
Table 6.5 Values of coefficients, d _{ij} , in Equation 6.10 _____	120
Table 6.6 The original Pedersen mixture tuning parameters _____	122
Table 6.7 Parameters in the original LBC Viscosity Correlation _____	123
Table 6.8 Experimental and modelling results of the 95% CO ₂ + 5% CH ₄ binary system	126
Table 6.9 Experimental and modelling results of the 95% CO ₂ + 5% Ar binary system _	128
Table 6.10 Experimental and modelling results of the 95% CO ₂ + 5% CO binary system	129
Table 6.11 Experimental and modelling results for the 95% CO ₂ + 5% O ₂ binary system	131
Table 6.12 Experimental and modelling results of the 95% CO ₂ + 5% H ₂ binary system	133
Table 6.13 Experimental and modelling results for MIX1 _____	135
Table 6.14 Average Absolute Deviation (AAD) for the different studied systems _____	141
Table 7.1 Injection gas compositions _____	150
Table 7.2 Mixture compositions MIX1 and MIX2 _____	151
Table 7.3 Equilibrated fluid densities, capillary rise heights and IFTs Measured in n-decane/CO ₂ System and n-decane/impure CO ₂ Systems at 310.95 K _____	155
Table 7.4 Pure components properties _____	162
Table 7.5 Binary Interaction Parameters (BIPs) used in PR EOS _____	162

Table 7.6 Comparison of experimental and predicted MMP for the studies systems in
this work _____167

LIST OF MAIN SYMBOLS

x_{co_2}	Mole fraction of CO ₂ in the mixture
τ	Period of oscillation
θ	Equilibrium contact angle in degrees
<i>AAD</i>	Average Absolute Deviation
<i>BWRS</i>	Benedict-Webb-Rubin-Starling EOS
<i>CCS</i>	Carbon Capture and Storage
<i>EOR</i>	Enhanced Oil Recovery
<i>EOS</i>	Equation of State
g	Acceleration due to gravity
H	Hydrate
h	Capillary rise
<i>IFT</i>	Interfacial Tension
k_{ij}	Binary interaction parameter
K_T	Isothermal compressibility
<i>LBC</i>	Lohrenz–Bray–Clark viscosity model
L_{CO_2}	Liquid CO ₂
M	Molecular weight

<i>MIX</i>	Mixture
<i>MMP</i>	Minimum Miscibility Pressure
<i>MOC</i>	Method of the Characteristics
<i>NDD</i>	Non density dependent mixing rules
<i>OF</i>	Objective function
<i>P</i>	Pressure
<i>P_c</i>	Critical pressure
<i>PR</i>	Peng-Robinson EOS
<i>P_{sat}</i>	Saturation pressure
<i>PT</i>	Patel-Teja EOS
<i>PV</i>	Pore Volume
<i>R</i>	Universal gas constant
<i>RF</i>	Recovery Factor
<i>SH</i>	Structure H
<i>SI</i>	Structure I
<i>SII</i>	Structure II
<i>SoS</i>	Speed of sound
<i>SRK</i>	Soave-Redlich-kwong EOS
<i>T</i>	Temperature

T_c	Critical temperature
V	Volume
VIT	Vanishing Interfacial Tension Approach
VLE	Vapour Liquid Equilibrium
VPT	Valderrama-Patel-Teja EOS
Y_{CO_2}	Vapour phase composition of CO ₂
z	Piston displacement
Z_{oil}	Oil height in the cell
α	Alpha function in EoS
η	Viscosity
ρ	Density

CHAPTER 1: INTRODUCTION

It is now widely accepted that anthropogenic CO₂ emissions from the burning of fossil fuels are largely responsible for the rapid rise in global temperatures recorded over the past century. Worldwide concerns over the threat of global warming have encouraged industrialised countries into working together to reduce carbon emissions, with specific targets being laid out in the 1997 Kyoto protocol agreement. To meet these goals, nations must increase investment in ‘clean’, renewable sources of energy, and develop solutions for reducing CO₂ (and other greenhouse gases) emissions from existing and new fossil fuel usage. With regard to carbon emissions, the most promising solution currently on the table is CO₂ capture and subsurface storage, primarily by injection into ageing oil reservoirs (where it can be used to improve oil recovery) or deep saline aquifers. A particularly promising fossil fuel technology, which can help addressing the above issues, is Integrated Gasification Combined Cycle (IGCC) reformation (e.g. White et al., 2000). In IGCC reformation, hydrocarbons (gas, coal, oil) are broken down to yield hydrogen and CO₂. The produced hydrogen can be used for electricity generation and/or as a clean fuel (e.g. to power motor vehicles), while the CO₂ is separated for subsurface disposal.

There are three capture processes: post-combustion, pre-combustion and oxy-fuel combustion. In a post-combustion, a solvent (mainly amines) process is required to capture CO₂ in flue gas of power plants. The process involves high energy to regenerate the solvent.

The pre-combustion process means that CO₂ is separated before the combustion process. The fuel is treated prior the combustion to produce CO₂ and H₂. The H₂ is used for generating power and the CO₂ is captured and stored. In oxy-fuel combustion process, oxygen is captured and removed from air to around 90% purity. Then O₂ is used to burn fuel instead of air to produce power. This process results in an exhaust stream consists of mainly pure CO₂ typically around 85 to 99%.

However CO₂ coming from the above capture processes is generally not pure and can contain impurities such as N₂, H₂O, H₂S, CH₄, CO and water. Scenarii of excepted impurities on CO₂ are reported in [Table 1.1](#) from [1], [2]. The type and amount of

impurities introduced in CO₂ depends on the type of capture process and the fuels beside the type of solvent used. The presence and type of other components may differ considerably between post-combustion, pre-combustion and oxy-fuel capture processes [3]. From Table 1.1, the non-condensable components N₂, Ar, O₂, CO and N₂ plus water are common impurities in all capture processes. The impurities fraction ranges from 15 volume% high to 0.05 volume% low. The removal of impurities is possible with further gas cleaning. However, CO₂ purification process is expensive, thus the overall cost will increase. The volume fraction of non-condensable gases such as N₂, H₂, CH₄, O₂ and Ar should be within 4% and H₂ should be removed as much as possible for its energy use [4].

Table 1.1 Impurities expected in CO₂ stream from capture processes [1]

Components	Pre-combustion		Post-combustion			Oxyfuel		
	Selexol*	Rectisol*	Post-1*	Post-2*	Post-3*	Oxy-1*	Oxy-2*	Oxy-3*
CO ₂ (Volume%)	97.95	99.7	99.93	99.92	99.81	85	98	99.94
N ₂ (Volume%)	0.9	0.21	0.045**	0.045**	0.09	5.8	0.71	0.01
O ₂ (Volume%)	-	-	0.015	0.015	0.03	4.7	0.67	0.01
Ar(Volume%)	0.03	0.15	-	-	-	4.47	0.59	0.01
H ₂ O(ppm)	600	10	100	100	600	100	100	100
CO(ppm)	400	400	10	10	20	50	50	50
H ₂ (ppm)	10000	20	-	-	-	-	-	-
CH ₄ (ppm)	100	100						
SO ₂ (ppm)	-	-	10	10	20	50	50	50
H ₂ S+COS(ppm)	100	20	-	-	-	-	-	-

*More details about these processes see [2]

**Concentration of Ar+N₂

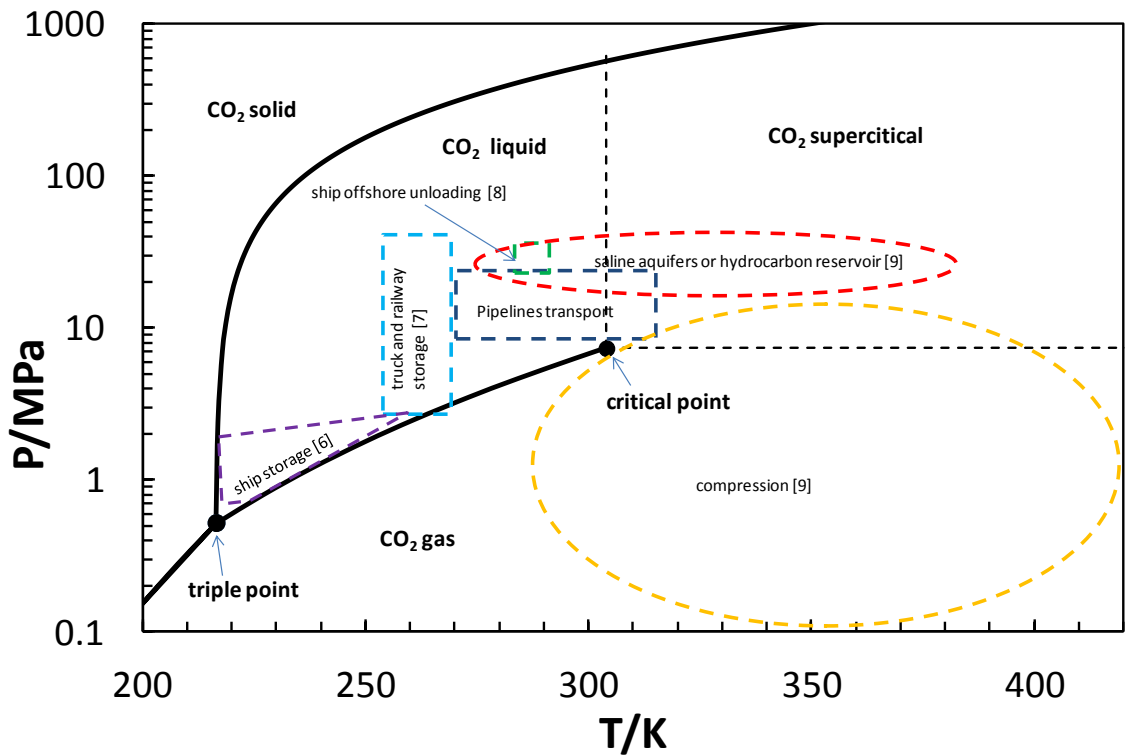


Figure 1.1 Various operations involved in the CO₂ transport

The intermediate step between CO₂ storage (or EOR) and capture is the transport. The most effective CO₂ transport options are either by pipeline or ship due to high bulk volume required for CCS. There are other CO₂ transport options such as railway or truck. Figure 1.1 shows the phase diagram of CO₂. The critical point of CO₂ [9] is (304.25 K and 7.3 MPa) and its triple point is (216.65 K and 0.51 MPa). CO₂ is compressed for transportation by ship, at a pressure of 0.6 to 1.7 MPa and cooled down to near the triple point. CO₂ is loaded to the ship tank at operating pressure and temperature while the unloading condition takes place at higher pressure and temperature (see example of offshore unloading Figure 1.1).

CO₂ is transported in long distance pipelines as a dense, gas or supercritical fluid or in the sub-cooled liquid state. The transport as a dense fluid seems to be the most economical among the others [5]. It is possible to transport CO₂ in gas phase; however this means lower density and high pressure drop per unit length i.e. lower capacity and higher costs [10]. The operating temperature of the pipelines is controlled by the underground soil temperature. CO₂ dense phase is the phase above the critical pressure of CO₂, 7.3 MPa. Above the critical pressure and temperature, CO₂ becomes supercritical. In this phase, CO₂ has high density as a liquid and low viscosity as gas.

However, in cold and moderate sites where the soil temperature is lower than 304.3 K, CO₂ is transported as liquid dense phase. The pipelines pressure for liquid dense phase should be kept above the critical pressure to avoid two phase flow. Geological and ocean storages are the two main possible methods to store the captured CO₂ [11]. Depleted oil and gas fields and deep saline formations are suitable for the geological storage of CO₂. Under high pressure (higher than 15 MPa), CO₂ is pumped into the depleted fields. The properties of CO₂, such as density and viscosity change as CO₂ is injected further deep underground. Knowing the properties will help to identify the capacity of the compressors required to compress CO₂ and hence minimizing the costs. As mentioned earlier, CO₂ can be used for EOR where CO₂ is injected into oil reservoirs to improve the flow of the remaining oil. The displacements range from viscosity reduction, oil swelling near the miscible pressure to completely miscible above the Minimum Miscible Pressure (MMP). The displacements depend on many factors under the same pressure and temperature such as VLE of CO₂ and oil and composition of both oil and CO₂ where impurities can strangely affect those factors.

The presence of the impurities combined with potentially long distance transportation of CO₂ could lead to challenging engineering and flow assurance issues. The presence of water may result in corrosion, ice and/or gas hydrate formation and pipeline blockage, so the fluid system should meet certain dehydration requirements. Furthermore, the gaseous CO₂ rich stream is generally compressed to be transported as liquid or dense-phase state in order to avoid two-phase flow and increase the density of the fluid system. The presence of the above impurities will also change the physical properties of the stream, i.e., the system's bubble point pressure and viscosity, hence affecting the compression requirement. Furthermore, in the majority of feasibility studies of CO₂ injection, the effects of these impurities have been overlooked or completely ignored.

The aim of this work is to investigate the effect of impurities such as N₂, O₂, H₂, CH₄, SO₂, CO, H₂S and Ar on the phase behaviour, speed of sound, density, viscosity, and isothermal compressibility. The work is also extended to investigate the effect of impurities on the properties of CO₂/oil mixture. The properties include IFT, MMP, swelling factor and saturation pressure.

In [Chapter 2](#) Vapour-Liquid Equilibria of CO₂ and CO₂/impurities systems have been evaluated using 3 cubic equations (SRK, PR and VPT) combined with the classical

mixing rules. To match the experimental data, the three equations have been tuned based on the VLE or saturation pressure of the binary CO₂/impurities mixtures. The critical point curves of the binary mixtures have been also estimated with the equation. In addition, the saturation pressure and speed of sound of a synthetic alkane mixture with CO₂ + impurity binary mixtures were measured at constant temperature, 344.3 K. The effect of four impurities, Ar, H₂, CO and CH₄, on the properties has been investigated.

Accurate measurements and prediction of speed of sound in CO₂-rich systems can have many applications, e.g., determining some thermodynamic properties, 4-D seismic, or monitoring compositional changes. In [Chapter 3](#), speed of sound and isothermal compressibility in pure liquid CO₂ and CO₂ + impurities have been measured under pressure up to 50 MPa and temperature ranging from 268.15 to 301.15 K. The tested impurities were N₂, CH₂, CO, O₂, H₂ and Ar. Measurements have also been conducted on two multi-components CO₂-rich mixtures, i.e., CO₂-CH₄-H₂-N₂ and CO₂-Ar-CO. A new generalized correction based on the PR-model for specific volume of liquid CO₂ has been proposed. The model is tested with the measured isothermal compressibility for both CO₂ and CO₂/impurities system.

The aim of [Chapter 4](#) is to evaluate the risk of hydrate formation in a rich carbon dioxide stream. The three phases HLV, HLL and HLLV equilibria of the ternary CO₂ + (N₂ or CH₄ or O₂ or Ar or H₂ or CO) + water systems have been determined at a constant composition (4.56% N₂, 7.09% H₂, 5.85% CH₄, 5.87% CO, 5.03% Ar, and 5.34% O₂ in dry basis). Measurements have been performed using a standard constant-volume equilibrium step-heating technique. A thermodynamic approach has been employed to model the phase equilibria. The thermodynamic model was used to predict the hydrate dissociation conditions of CO₂ and CO₂ rich streams in the presence of free water.

Effect of impurities on CO₂ density has been investigated and explained in [Chapter 5](#). The density of CO₂ in the presence of several gaseous impurities was measured over a range of temperature from 283.15 K to 423.15 K and pressures up to 50 MPa in both liquid and supercritical phases. The measurements have been carried-out with an Anton Parr densitometer. Within 4 to 5 mole% of various impurities (i.e. CH₄, N₂, O₂, H₂, Ar and CO) in rich-CO₂ binary system have been tested in this study. The findings show a

significant reduction in the CO₂ + impurities systems density in presence of the above impurities. A new volume correction added to the EOSs was implemented to model the experimental data in both liquid and supercritical phases.

Viscosity is a key transport property for transport pipeline systems as well as sub-surface and process systems. In [Chapter 6](#), the viscosity of CO₂ systems with impurities, such as methane, nitrogen, argon, carbon monoxide, oxygen and hydrogen, were measured using capillary viscosity measurement technique in an in-house designed and constructed set-up at pressures up to 5.5 MPa and temperatures of 280, 288.15, 300.15, 308.15, 323.15 and 343.15 K. The results show that all of the tested impurities caused a reduction of viscosity compared to pure CO₂. In addition, two models were modified in order to predict and estimate the viscosity of CO₂ systems. First, the predictive model of Pedersen was modified substituting CO₂ as a reference fluid to the original reference fluid. Second, the LBC correlation was used after tuning to match the obtained experimental data.

In [Chapter 7](#):, the effect of impurities (N₂, CH₄, CO, O₂, H₂ and Ar) of CO₂/n-Decane system on the IFT (Interfacial Tension), the swelling factor and the MMP (Minimum Miscibility Pressure) have been investigated at 310.95 K. The impact of impurities in the injection gas was studied, four binary mixtures were investigated: CO₂-CH₄, CO₂-N₂, CO₂-O₂, CO₂-H₂ and two multi-components mixtures. A capillary rise apparatus was used to measure the IFTs and the swelling factors at the same pressure and temperature conditions and the VIT (Vanishing Interfacial Tension) approach was used to estimate the MMPs of the systems containing impurities. The experimental MMPs were compared with the predictions made by an in-house software developed based on the multiple-mixing cell approach. The above properties are compared between systems: CO₂/n-decane in the presence of impurities and the pure CO₂.

[Chapter 8](#) presents the conclusions of this study and recommendations for future work.

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CHAPTER 2: VLE OF CO₂/IMPURITIES+SATURATION

PRESSURE AND SPEED OF SOUND OF CO₂/ALKANE

MIXTURES

2.1 Introduction

CO₂ is transported under different pressure/temperature conditions, ranging from the triple point to near the critical point. Impurities can have high impacts on the thermodynamic properties of CO₂-rich systems that can further affect the design, transport cost and operation. In Enhanced Oil Recovery (EOR) process, laboratory studies showed that impurities can affect solubility, PVT properties and Minimum Miscible Pressure (MMP) behaviours. In addition, impurities affect thermodynamics and transport properties of CO₂ stream such as vapour liquid equilibrium (VLE), viscosity, density and hydrate stability. The need for a reliable model to predict the vapour liquid equilibrium is essential for CO₂ transport. Reliable VLE and physical properties data on CO₂ rich system are essential to design the necessary purification processes, compression and transportation of CO₂ for carbon capture and storage (CCS).

Accurate VLE models are important because the operation region for CO₂ should be within dense or supercritical area in which non-accurate prediction can lead to high risk or inefficient transport. In addition, two phase flow in a pipeline should be avoided due to the potential for damage to the pipelines [1]. The experimental data of CO₂ impurities on the phase behaviour properties are limited in the range of transport conditions and in some cases scattered. Therefore, the need of an accurate theoretical modelling approach to overcome this gap becomes potentially recommended in both industries and research studies.

Many papers are available about predicting pure CO₂ and impure CO₂ VLE using equation of states [2- 8]. For the design of CO₂ pipelines, no consensus has been achieved in which equations of state should be used for rich CO₂ mixtures. The EOSs used by authors ranged from specialized to complex equations such as Span and Wagner and Benedict-Webb-Rubin-Starling (BWRS) for pure CO₂, to general and

simplest ones such as Peng-Robinson (PR), Soave-Redlich-kwong (SRK) and Patel-Teja EOSs (PT). The term “complex” in equations of state comes from the parameters number and the structure of the equations. For example Span and Wagner EOS used around 51 terms (many of which include logarithms and exponentials) and BWRS with a minimum of 8 parameters (while the classic EOSs such as PR and SRK contain only two parameters). These equations may give better estimation of phase properties, however because of their complexity, Helmholtz and Gibbs free energy derivatives are more complicated to evaluate. CO₂ in the presence of impurities classified as mixtures, in which general EOSs interaction parameters are required. The EOSs should be tuned using experimental data to evaluate accuracy level achieved from their calculations. However, the non-ideal behaviour of mixture approaching the critical region is a challenge for the equations. In this work, three equations of state are used to predict the phase behaviour of CO₂/impurities system, PR, SRK and VPT (Valderrama-Patel-Teja equation of state) models. In addition, SRK is used to predict the critical locus curve of the binary CO₂/impurities mixture. The impurities that are covered in the work H₂, N₂, CH₄, O₂, Ar, CO, H₂S and SO₂.

Vapour liquid equilibrium (VLE), volumetric properties, and solubility data of alkanes in supercritical carbon dioxide are important information in the petroleum process design area. VLE and volumetric properties in those binary and ternary mixtures have been studied both computationally and experimentally extensively in the last decade [9-14]. There is no data available for CO₂ containing impurities in the literature. In this work, a mixture of alkane (n-dodecane, n-tridecane and n-hexadecane) plus pure or impure CO₂ were prepared. The CO₂ gas mixture was prepared with 5 mole% of impurities (H₂, Ar, CH₄ and CO). The saturation pressure of the mixtures was measured at 344.3 K. The corresponding speed of sound property of the mixtures in the presence of impurities was measured above the mixture saturation pressure. The measurements are compared with those obtained for pure CO₂.

2.2 VLE of CO₂ and CO₂/Impurities System

The study of the phase behaviour of CO₂/impurities is essentially important for the operation and the design of all the processes of CO₂ transport (ship carriers, loading, unloading) and capture (purification/liquefaction). Beside the experimental work, reliable empirical equations of state can play as a solution key in terms of cost and time. In this

study, in-house code is used to predict Vapour-Liquid Equilibrium of CO₂ and CO₂ binary systems by using Peng-Robinson (PR), Soave-Redlich-Kwong (SRK) and Valderrama-Patel-Teja (VPT) EOSs. Details of these equations are given below.

Peng-Robinson EOS:

$$P = \frac{RT}{V-b} - \frac{a}{V(V+b)+b(V-b)} \quad (2.1)$$

$$a = a_c(1+m(1-\sqrt{T_r}))^2 \quad (2.2)$$

$$a_c = \frac{0.45724R^2T_c^2}{P_c} \quad (2.3)$$

$$m = 0.37464 + 1.54226\omega - 0.26992\omega^2 \quad (2.4)$$

$$b = \frac{0.07780RT_c}{P_c} \quad (2.5)$$

Soave-Redlich-Kwong EOS:

$$P = \frac{RT}{V-b} - \frac{a}{V(V+b)} \quad (2.6)$$

$$a = a_c(1+m(1-\sqrt{T_r}))^2 \quad (2.7)$$

$$a_c = \frac{0.42748R^2T_c^2}{P_c} \quad (2.8)$$

$$m = 0.480 + 1.574\omega - 0.176\omega^2 \quad (2.9)$$

$$b = \frac{0.08664RT_c}{P_c} \quad (2.10)$$

Valderrama-Patel-Teja EOS:

$$P = \frac{RT}{V-b} - \frac{a}{V^2 + V(b+c) - bc} \quad (2.11)$$

$$a = \Omega_a \frac{R^2T_c^2}{P_c} \left((1+m(1-T_r^{0.5}))^2 \right) \quad (2.12)$$

$$b = \Omega_b \frac{RT_c}{P_c} \quad (2.13)$$

$$c = \Omega_c \frac{RT_c}{P_c} \quad (2.14)$$

$$\Omega_a = 0.66121 - 0.76105Z_c \quad (2.15)$$

$$\Omega_b = 0.02207 - 0.20868Z_c \quad (2.16)$$

$$\Omega_c = 0.57765 - 1.87080Z_c \quad (2.17)$$

$$m = 0.462823 + 3.5830\omega Z_c + 8.117(\omega Z_c)^2 \quad (2.18)$$

Where P , V , T , R are pressure, the molar volume, the temperature and the universal gas constant, respectively. The parameter a is a measure of the attractive forces between the molecules and the parameter b is a measure of the size of the molecules (intrinsic volume). ω and Z_c are acentric factor and critical compressibility factor, respectively. The above equations are developed for pure components. The van der Waals conventional random mixing rules are used to determine, b , and, c , parameters for multi-components systems:

$$b = \sum_{i=1}^N x_i b_i \quad (2.19)$$

$$c = \sum_{i=1}^N x_i c_i \quad (2.20)$$

Non Density Dependent (NDD) mixing rule was used in this work to calculate the parameter, a . The attraction parameter, a , is divided into two parts, classical and asymmetric, a^c , and, a^a , as following:

$$a = a^c + a^a \quad (2.21)$$

$$a^c = \sum_{i=1}^N \sum_{j=1}^N x_i x_j (a_i a_j)^{0.5} (1 - k_{ij}) \quad (2.22)$$

$$a^a = \sum_p x_p^2 \sum_i x_i a_{pi} l_{pi} \quad (2.23)$$

$$a_{pi} = (a_p a_i)^{0.5} \quad (2.24)$$

$$l_{pi} = l_{pi}^0 - l_{pi}^1 (T - T_0) \quad (2.25)$$

Where k_{ij} is the binary interaction parameter and l_{pi} is the binary interaction parameter between polar component and other components. l_{pi} is a function of temperature. It is calculated using [Equation 2.25](#). l_{pi}^0 and l_{pi}^1 are the interaction parameters while T_0 is the freezing point in K. However, the asymmetric, a^a , is not calculated for CO₂ or CO₂ mixtures since CO₂ is classified as a non-polar component.

An extensive literature review was carried-out to obtain VLE data of the binary CO₂/mixtures. Eight binary CO₂/impurities mixtures are included: H₂, N₂, CH₄, O₂, Ar, CO, H₂S and SO₂. The experimental data, listed in [Table 2.1](#), were used for tuning the EOSs. Binary interaction parameters (BIPs) were tuned by minimizing either the experimental bubble point pressures or Vapour-Liquid Equilibrium data. VLE and bubble point data have been gathered for the various binary CO₂/impurities systems. The binary interaction parameters between CO₂ and the impurities were adjusted using the references mentioned in [Table 2.1](#) through a Simplex algorithm using the Objective Function, OF, displayed in [Equation 2.26](#).

$$OF = \left\{ \begin{array}{l} \frac{1}{N} \sum_{i=1}^N \left(\frac{P_{Exp}^{Bubble} - P_{Cal}^{Bubble}}{P_{Exp}^{Bubble}} \right) \\ \frac{1}{N} \sum_{i=1}^N \left[\left(\frac{y_{Exp} - y_{Cal}}{y_{Exp}} \right)^2 + \left(\frac{x_{Exp} - x_{Cal}}{x_{Exp}} \right)^2 \right] \end{array} \right\} \quad (2.26)$$

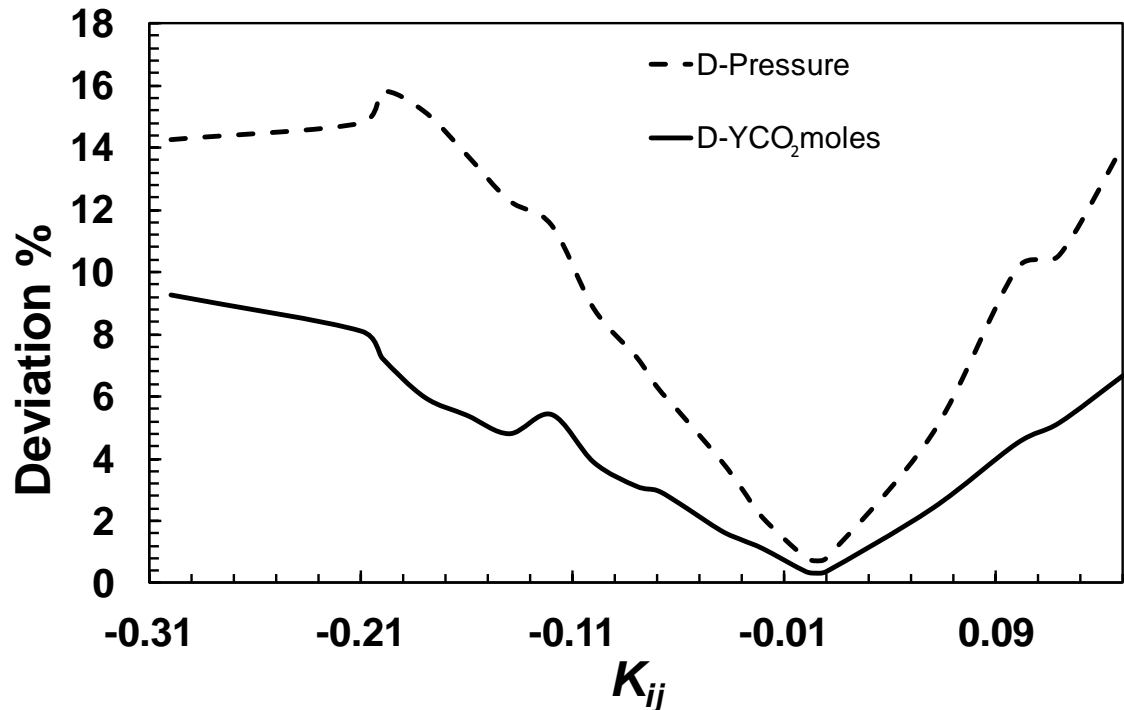
The BIPs obtained for the binary mixtures are summarised in [Table 2.2](#). According to the mixing rule used in EOSs, the BIP should be temperature independent and its value is between -1 to 1. The default value is zero for SRK and PR and 1 for VPT (i.e. $(1 - k_{ij}) = 1$). The change in BIPs can lead to high deviation, accordingly the need to choose the appropriate values. [Figure 2.1](#) to [Figure 2.3](#) show the average deviations in the saturation pressure (P_{sat}), vapour phase composition of CO₂ (Y_{CO_2}) and ($P_{sat} + Y_{CO_2}$) as a function of k_{ij} in CO₂/N₂ mixture at 270 K. The plots show that the EOS BIPs for the mixture are very close to their default values (zero). The BIPs of N₂ regressed using the objective function for wide range of temperatures as listed in [Table 2.2](#). The BIPs result for CO₂/N₂ using the three EOSs shown in [Figure 2.1](#) to [Figure 2.3](#) at 270 K are close to the general BIP for the same mixture. Our results are also matched the BIPs results generated by Li and Yan (2006) [7].

Table 2.1 Literature data for the CO₂ + impurities systems

Mixtures	Source
CO ₂ /CO	[15-17]
CO ₂ /N ₂	[18-28]
CO ₂ /O ₂	[18,29]
CO ₂ /Ar	[30-31]
CO ₂ /CH ₄	[21],[25], [32-47]
CO ₂ /SO ₂	[48-50]
CO ₂ /H ₂	[39],[51,53]
CO ₂ /H ₂ S	[54-56]

Table 2.2 BIPs for different binary CO₂/mixtures

	CO	N ₂	O ₂	Ar	H ₂	CH ₄	SO ₂	H ₂ S
VPT	0.005	-0.014	0.111	0.129	-0.387	0.099	0.02	0.082
PR	-0.042	-0.048	0.091	0.113	0.0202	0.092	0.02	0.090
SRK	-0.022	-0.046	0.106	0.123	-0.142	0.100	0.02	0.096

Figure 2.1 AAD Deviations from PR EOS on CO₂/N₂ mixture at 270 K, data from [23]

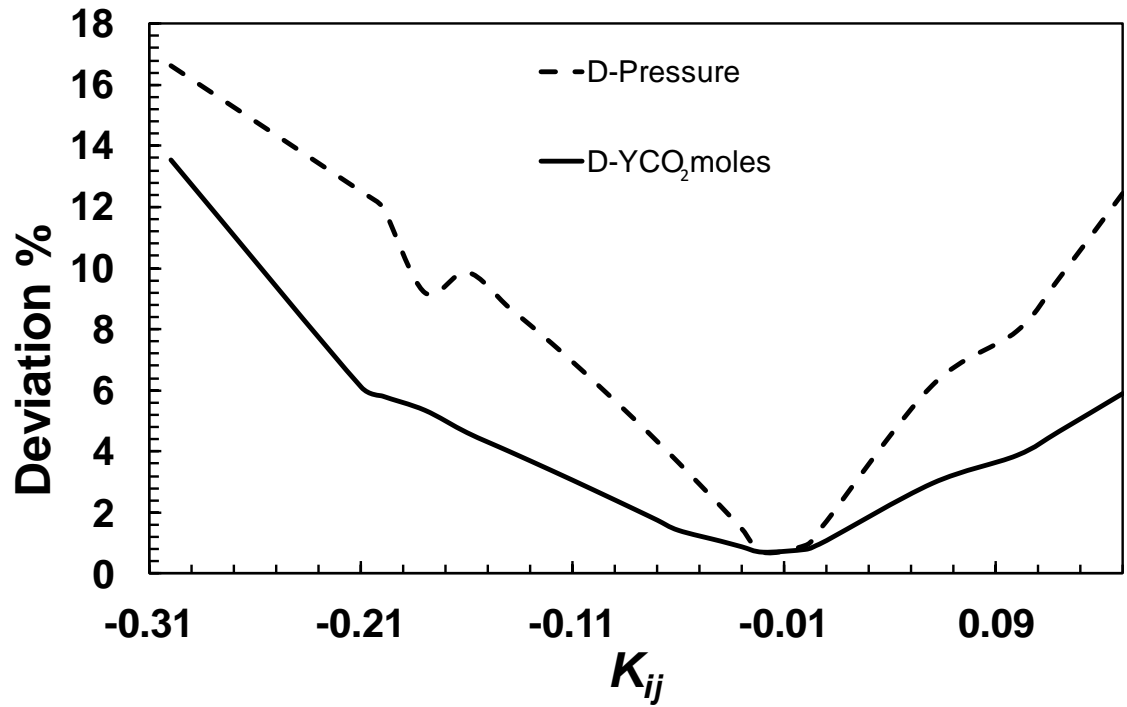


Figure 2.2 AAD Deviations from SRK EOS on CO₂/N₂ mixture at 270 K, data from [23]

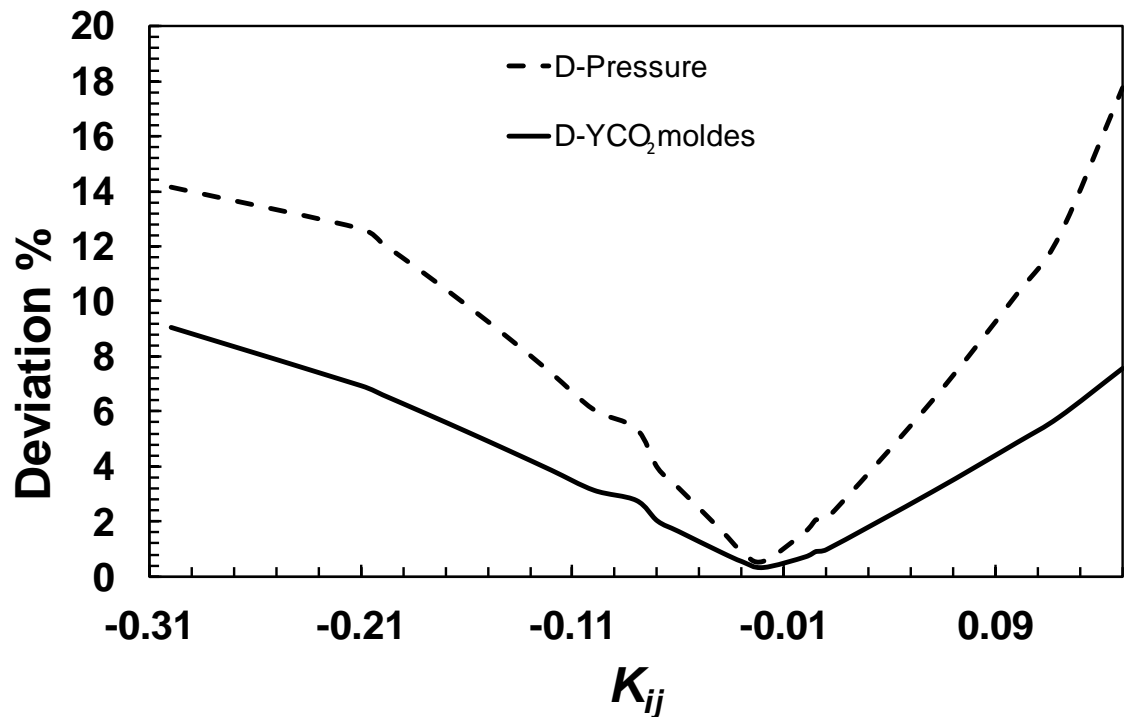


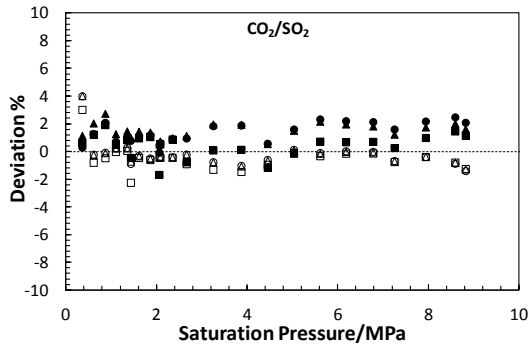
Figure 2.3 AAD Deviations from VPT EOS on CO₂/N₂ mixture at 270 K, data from [23]

Table 2.3 AAD of EOSs on VLE data of binary CO₂ mixtures

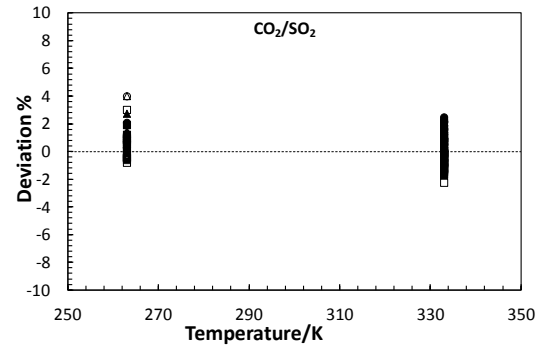
Model	AAD %	O ₂	CH ₄	H ₂ S	H ₂	N ₂	CO	SO ₂	Ar
SRK	P _{sat}	4.68	1.87	1.28	10.34	3.79	5.34	0.82	4.49
	Y _{CO2}	3.99	2.69	4.19	11.48	2.41	4.12	0.76	4.02
PR	P _{sat}	4.76	1.70	2.25	10.37	3.25	6.44	1.48	4.87
	Y _{CO2}	3.17	2.32	5.09	10.37	1.99	2.92	0.60	2.76
VPT	P _{sat}	4.64	1.70	1.73	9.92	3.46	5.03	1.39	4.62
	Y _{CO2}	3.37	2.28	4.35	9.50	2.22	3.79	0.60	3.06

The experimental data (Table 2.1) of the binary CO₂ systems are used to tune the three EOSs models at different conditions. From Figure 2.4 (a-p), the models show a good agreement with experimental data. However, they are weaker near the critical point for most of these binary mixtures. The models were tested in a wide range of pressure and temperature conditions. Table 2.3 lists the AADs of the EOSs on the saturated pressure (P_{sat}) and vapour compositions (Y_{CO2}) for the systems studied here. Figure 2.4 (a-p) shows the calculated AAD as function of saturated pressure and temperature for the eight CO₂/impurities systems. From the tables and figures we can observe that the PR, SRK and VPT equations predict bubble pressure with about the same accuracy. The overall AAD on the saturation pressure for the eight binary systems are 4.08, 4.39 and 4.06% by using SRK, PR and VPT EOSs respectively. The overall AAD on vapour composition are 4.21, 3.65 and 3.64% by using SRK, PR and VPT EOSs respectively. This indicates that VPT is slightly superior to SRK and PR. The higher deviations found in the pressure predictions of the CO₂/impurities systems are due to higher errors near the critical region of the mixtures. The predictions in the vapour phase compositions are more accurate than the saturation pressure predictions in most of the tested impurities such as O₂, N₂, CO, Ar and SO₂. In terms of the impurities, the highest deviation is found for the binary CO₂/H₂ system. This is a more or less general behaviour of cubic equations of state when applied to H₂ containing mixtures. This is mainly because of the low critical temperature of H₂ compare to the other impurities and CO₂. The AAD on the P_{sat} for the CO₂/H₂ mixture are 10.34, 10.37 and 9.92% by using SRK, PR and VPT respectively and the deviation can reach up to 55%. If the composition range for the mixture is restricted to a mixture up to 10% CO₂, the overall AAD on the P_{sat} is significantly reduced to 2.80, 2.93 and 2.52% for SRK, PR and VPT, respectively. The

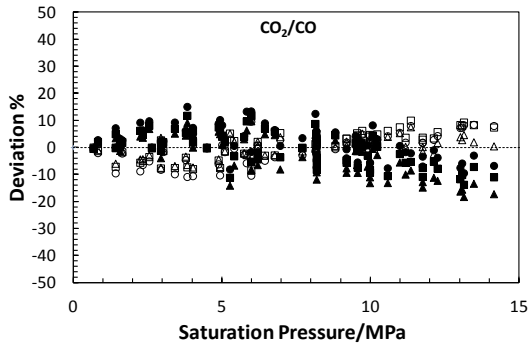
effect of modelling the previous impurities by using PR on the phase envelope of CO₂ is shown in [Figure 2.5 \(a-h\)](#) which also includes the saturation pressure for both pure carbon dioxide and 100% of the contaminants. The envelope size of the mixture is important in CO₂ transport and pipelines design and the two phase region should be avoided as possible. It is obvious from the figures that the envelope size depends on both the type of impurity and the composition of the impurity in the mixture. Each square dot along the critical locus curve represents 10% mole of impurity. It is also very clear that the size of phase envelope is increased as the impurity concentration is increased. Both mixture saturation and vapour lines of the tested mixtures (except for SO₂ and H₂S) tend to increase compared with the saturated line of pure CO₂. The figures indicate that the saturation pressures with CO₂ are increased greatly when the CO₂ rich stream contains H₂, O₂, Ar, CO, CH₄ and N₂. In the other hand, adding SO₂ and H₂S reduces the saturation pressures. The effect of the tested impurities can be classified into four categories. The first includes H₂ which has the highest impact on VLE property. The second category is the impurities group (CO, Ar, O₂ and N₂). The third and fourth categories are CH₄ and the impurities group (SO₂ and H₂S) respectively. For instance at 290 K, 5% mole fraction of impurity, H₂ can raise the saturation pressure of pure CO₂ to 80%, the group (CO, Ar, O₂ and N₂) to 33 to 35% and CH₄ to 22%. However, at the same temperature and composition, the presence of SO₂ and H₂S reduce the saturation pressure to 7% and 2%, respectively.



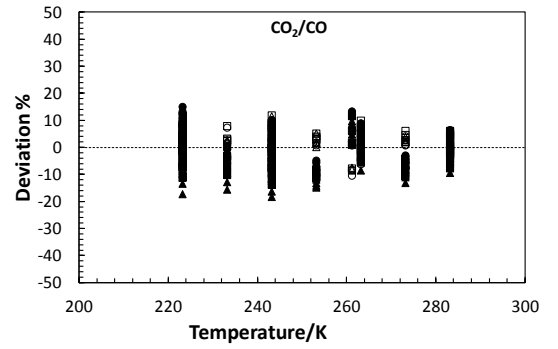
a)



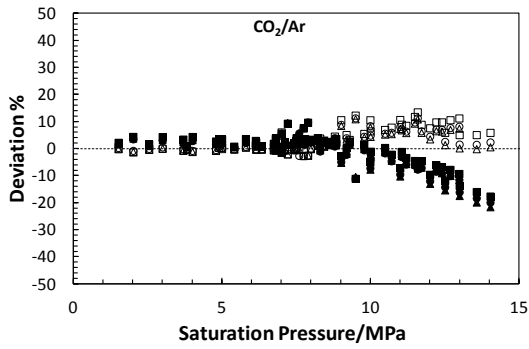
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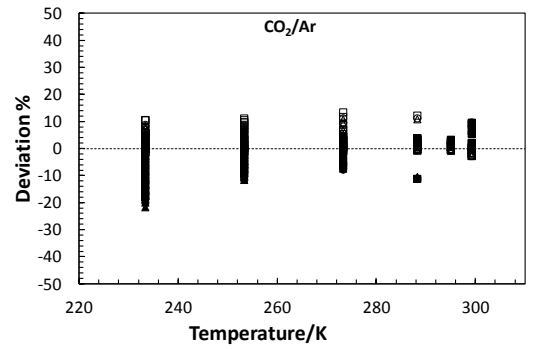
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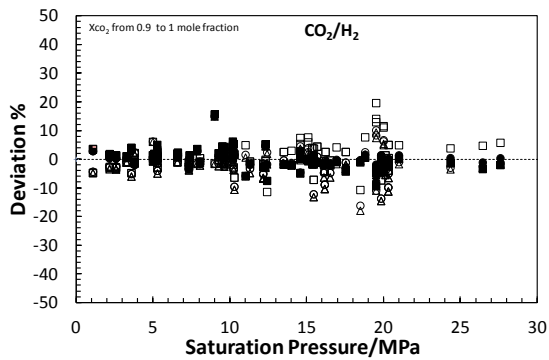
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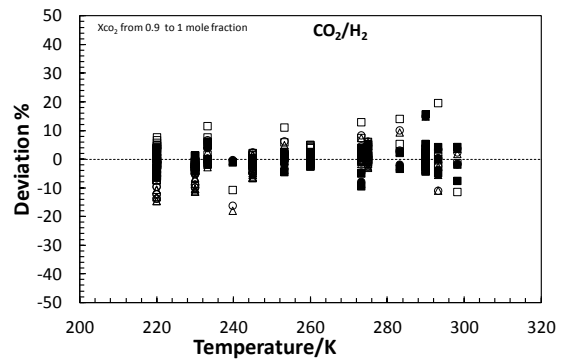
e)



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g)



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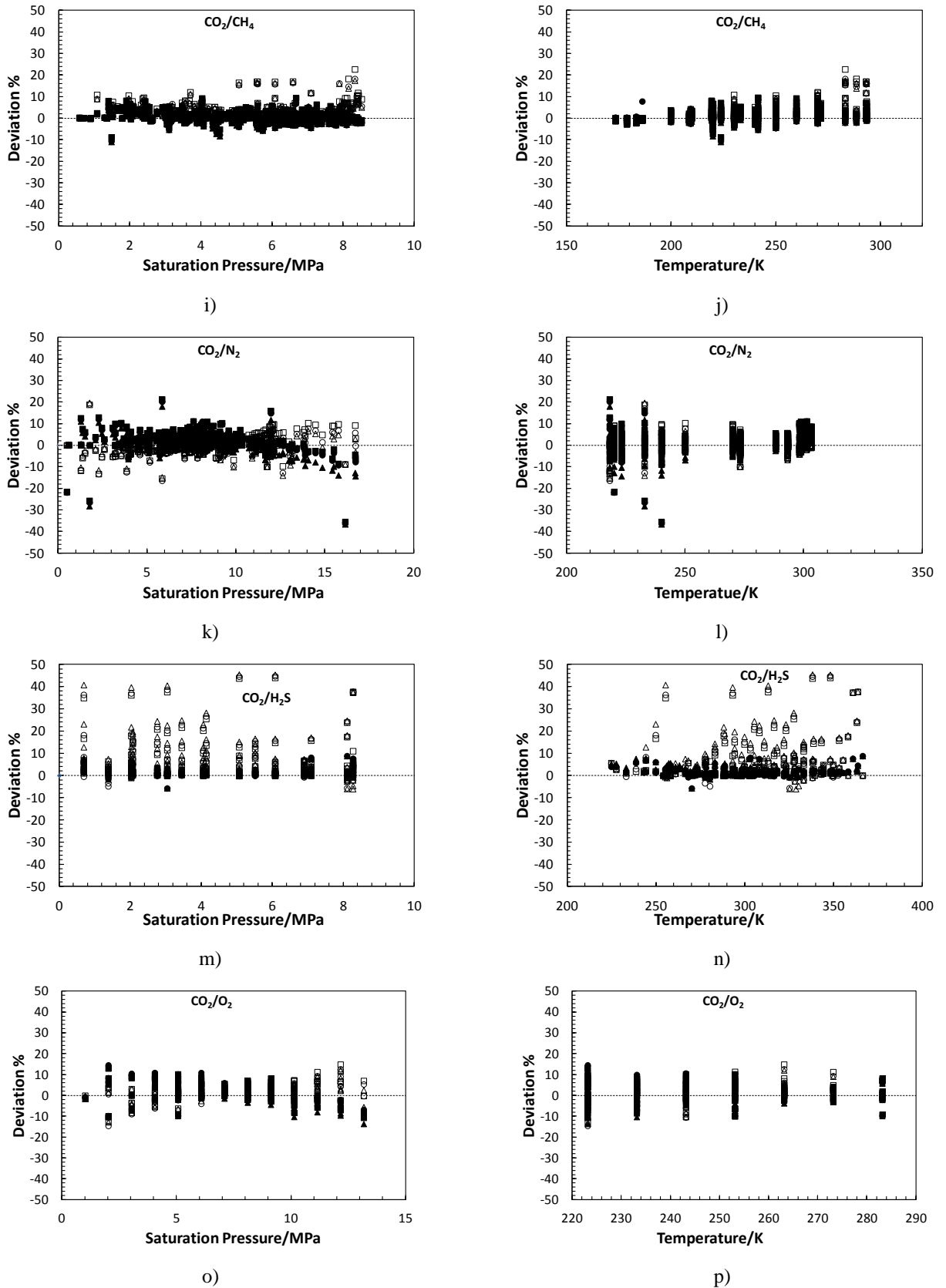


Figure 2.4 (a-p) the deviation of PR ($\blacktriangle\triangle$), VPT ($\bullet\circ$) and SRK ($\blacksquare\square$) EOSs on the calculated VLE properties of binary CO_2 /mixtures; empty symbols represent AAD on vapour fraction of CO_2 (Y_{CO_2}) and full symbols represent AAD on saturation pressure (Sources are listed in Table 2.1)

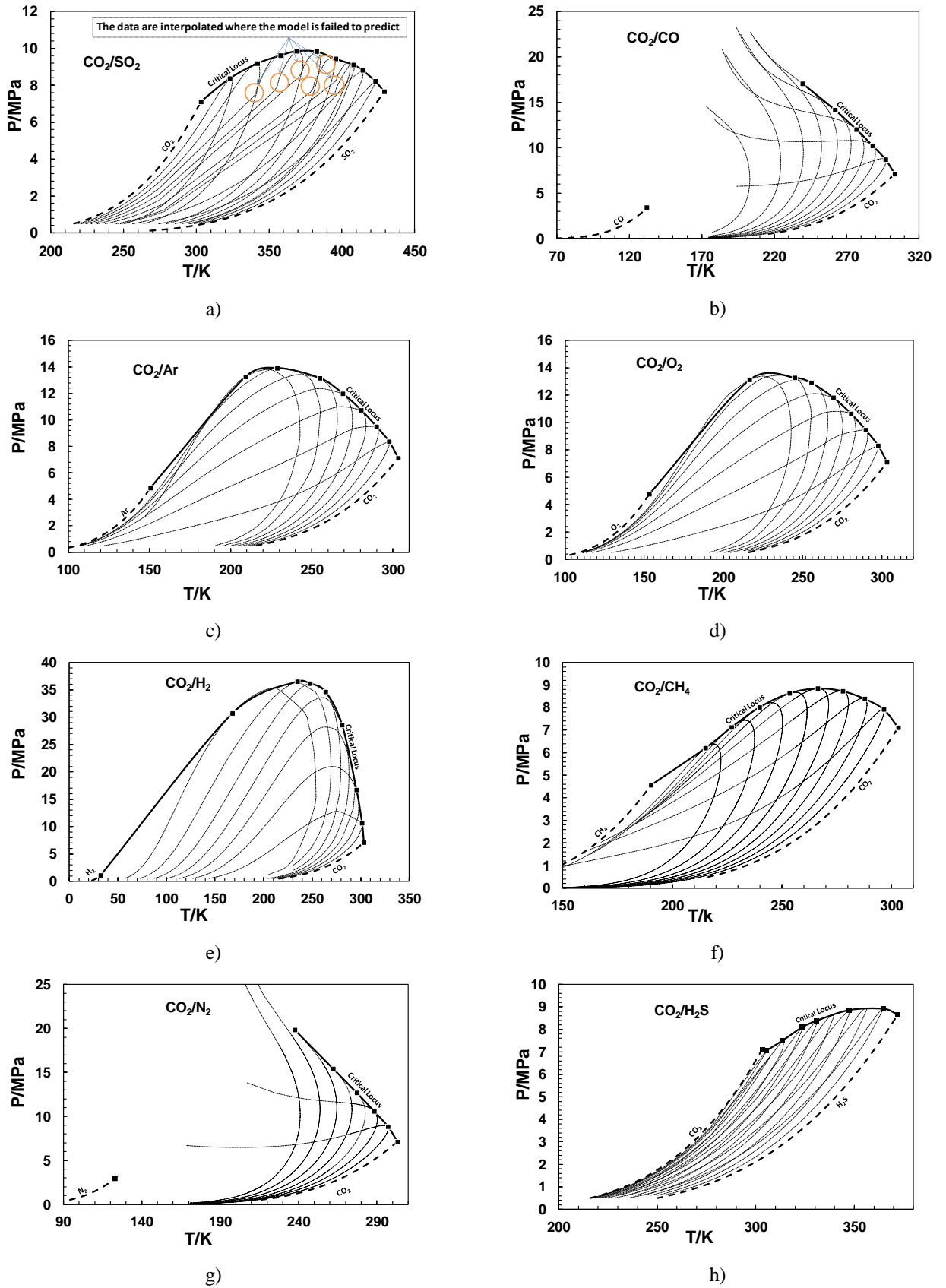


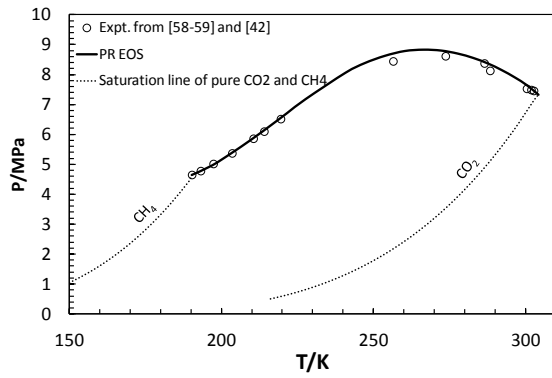
Figure 2.5 (a-h) prediction results on VLE of CO₂ mixtures by PR EOS

2.2.1 Critical Behaviour of CO₂ Rich Fluid

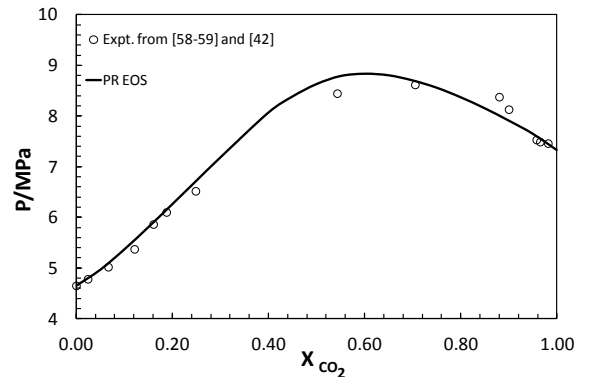
For CO₂ to be transported in pipelines conditions (near or above supercritical phase), knowledge of the critical properties of the stream are essential. For economical and engineering reasons, it is better to have a high critical temperature and a low critical pressure; high critical temperature for reducing the cooling capacity (lower insulation cost and therefore lower maintenance cost) and low critical pressure for transportation at lower pressure [57]. Few critical points data were found in the literature for CO₂/impurities systems, such as CO₂/N₂, CO₂/CH₄ and CO₂/H₂S and no critical data for the other impurities systems concern in this work. Figure 2.6 (a-f) shows the experimental and modelling critical locus curve for three binary mixtures, CO₂/N₂, CO₂/CH₄ and CO₂/H₂S. The data were calculated with the PR EOS. The experimental critical point data were not used for adjustment of the model but nevertheless reproduces the experimental results equally well as shown in the figures. For the binary CO₂/CH₄ (Figure 2.6-a and Figure 2.6-b), the model can predict well the critical pressure and temperature in high concentration CO₂ (up to 8% CO₂) and high concentration CH₄ (up to 20% CH₄). However, away from the above composition, the model cannot predict well the critical temperature or the critical pressure. The reason is probably that the source is from 1954 [42] and more reliable experimental data are required to test the reliability of the model. The critical curve of the binary CO₂/H₂S is shown in Figure 2.6-c and Figure 2.6-d. From Figure 2.6-d, it is very clear that the experimental data of the critical temperature in rich CO₂ are in a good agreement with the model. However, the model is over-predicting compared to the measured values of the critical pressure. It is usually necessary to have more than one set of experimental data points from different sources in order to evaluate the model successfully. Figure 2.6-e and Figure 2.6-f show also plots for the critical properties of the system CO₂/N₂ system. The modelling results for the critical points in the system are in agreement with the measured values at high CO₂ mole fraction (above 80%). Then the model starts to deviate as more N₂ is added to the mixture.

The eight binary mixtures critical lines are shown in Figure 2.7 in the temperature range 243.3 K ($\frac{T_{Cm}}{T_{C_{CO_2}}} \approx 0.8$) to 365 K ($\frac{T_{Cm}}{T_{C_{CO_2}}} \approx 1.2$). Each dot on the critical locus curves represents 10 mole% of impurity in CO₂. In terms of the critical pressure, it is clear that the binary mixtures except for H₂S tend to raise the critical

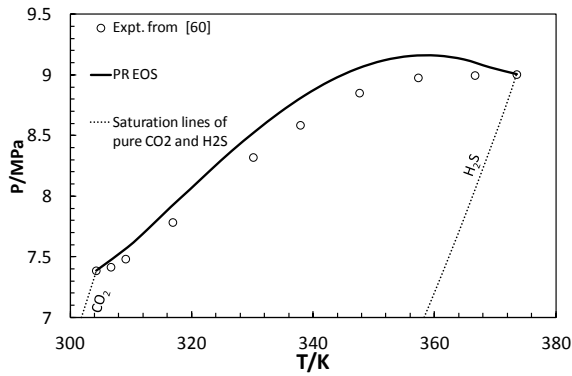
pressure of CO₂. For instance at 10 mole% of impurities, H₂ has the highest impact ($P_{Cm}/P_{CCO_2} \approx 1.4$) while H₂S has the lowest ($P_{Cm}/P_{CCO_2} \approx 0.99$). The other impurities (CO, Ar, O₂, N₂, SO₂ and CH₄) impact varies in the range $1.1 \leq P_{Cm}/P_{CCO_2} \leq 1.2$. In the other hand, the presence of H₂S or SO₂ increases the critical temperatures of CO₂ by 0.7 % to 7% respectively while the other impurities (CO, Ar, O₂, N₂, and CH₄) decreases the critical temperature by an average of 2% and H₂ to less than 1%.



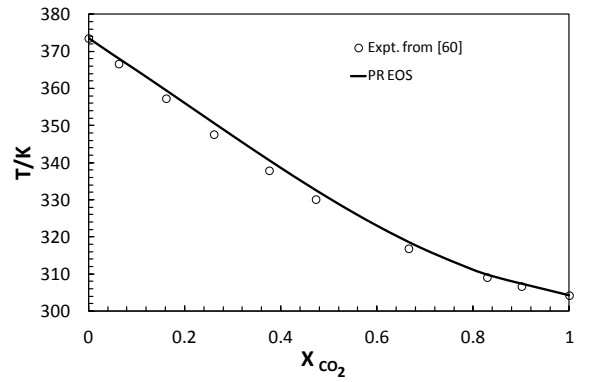
a)



b)



c)



d)

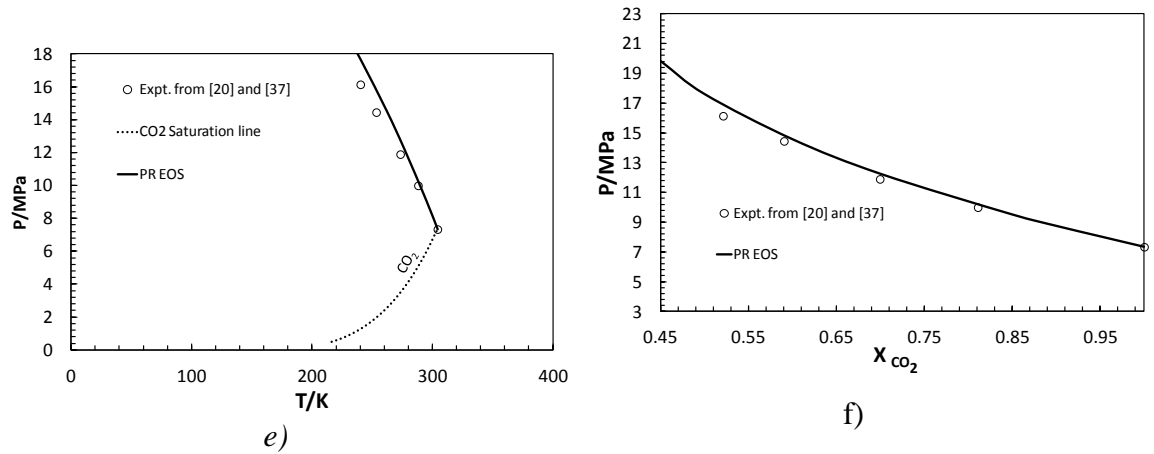


Figure 2.6 (a-f) critical properties of CO_2/CH_4 (a and b), CO_2/H_2S (c and d), CO_2/N_2 (e and f); PR EOS is used for modelling

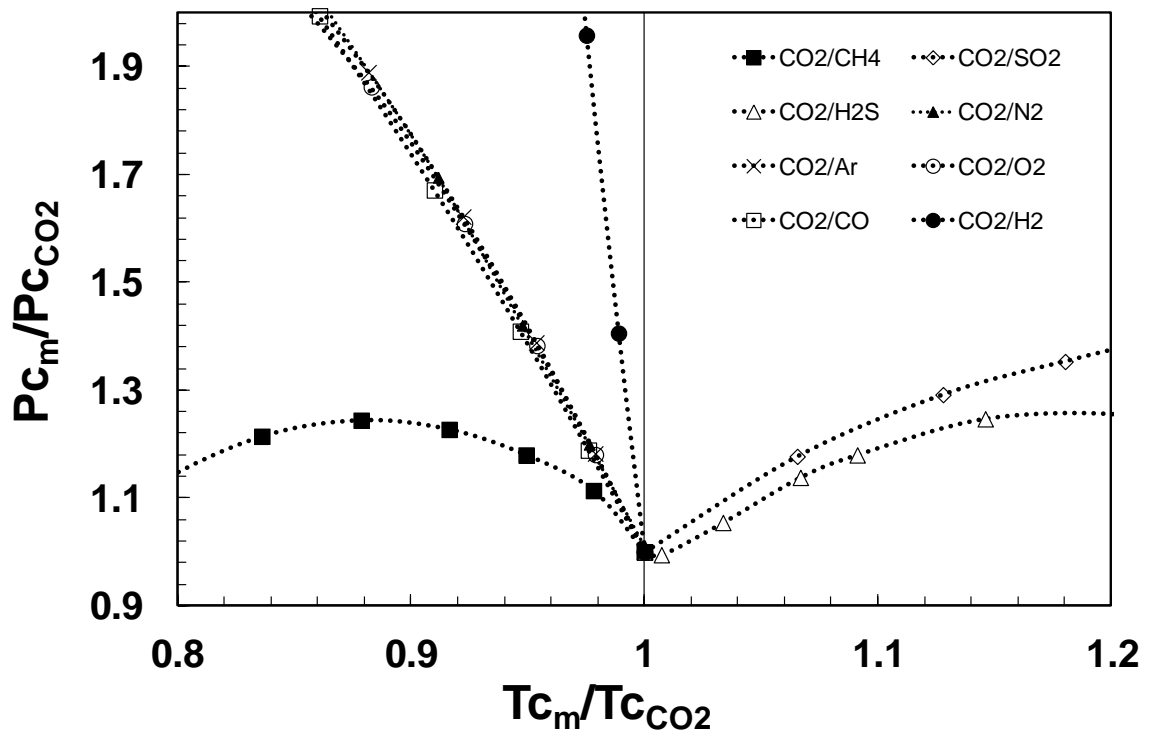


Figure 2.7 Critical lines of the binary CO_2 mixtures. $T_{c,m}$ and $P_{c,m}$ are the critical temperature and pressure of mixtures, respectively. T_{c,CO_2} and P_{c,CO_2} are the critical temperature and pressure of CO_2 respectively

2.3 Bubble Point and Speed of Sound Measurements of CO_2 and Alkanes System

The efficiency of oil displacement by CO_2 is related to the phase equilibrium since the process involves intimate contact of gases and liquids. Most studies on the phase behaviour of CO_2 with hydrocarbons have been limited to mixtures of pure CO_2 with hydrocarbons. These studies were usually made to obtain binary interaction parameters.

One of the aims of this study is to provide few experimental data on how impurities such as N₂, CH₄, CO, Ar etc. affect the phase behaviour of CO₂/crude oil mixtures.

Materials: n-decane (C₁₀H₂₂), n-tridecane (C₁₃H₂₈) and n-hexadecane (C₁₆H₃₄) of purity (weight) > 99 % were purchased from Fisher Scientific Co. and used without further purification. Carbon dioxide with a minimum purity of 99.9% (mole) was used in this work. Compositional details of the alkane oil used is given in [Table 2.4](#).

Table 2.4 Alkane mixture compositions

C. number	MW	Weight % ($\pm 0.03\%$)
C10	142.28	52.9
C13	184.37	30.2
C16	226.45	17.9
Total		100

The equipment used to measure the speed of sound was specially designed for measuring acoustic properties of aqueous solutions. A schematic diagram of the ultrasonic set-up can be found in [Chapter 3](#). All the mixtures were prepared gravimetrically using a digital balance (Sartorius) with a precision of (10⁻²g). The maximum uncertainty of the sample preparation in mole fraction was ± 0.002 . The system temperature is controlled by circulating coolant from a cryostat which is capable of maintaining the cell temperature stability to within better than 0.05 °C. The piston displacement is measured by a digital piston displacement indicator attached to one end of the cell (accuracy 10⁻³ mm). The maximum length of the cell is 10 cm. The digital indicator measures the in-displacement of the piston into the cell. The liquid components of the sample were charged into the cell before the gas to avoid the possibility of backflush. The amounts of liquid and gas charged into the recombination cell were measured gravimetrically by weighing the containers before and after charging. In order to investigate the effects of CO₂ impurities on alkanes-CO₂ mixture, a mixture of three components were prepared n-decane (C₁₀H₂₂), n-tridecane (C₁₃H₂₈) and n-hexadecane (C₁₆H₃₄). The experiment was run with a constant weight composition of alkane mixture, given in [Table 2.4](#). The total mixture composition was then recalculated after CO₂ or CO₂ with impurities injection. In this part of study, five tests

were carried-out in order to understand the effect of Ar, H₂, CO and CH₄ on CO₂ saturation pressure at 344.3 K. All tests were performed with 5 mole% of impurities.

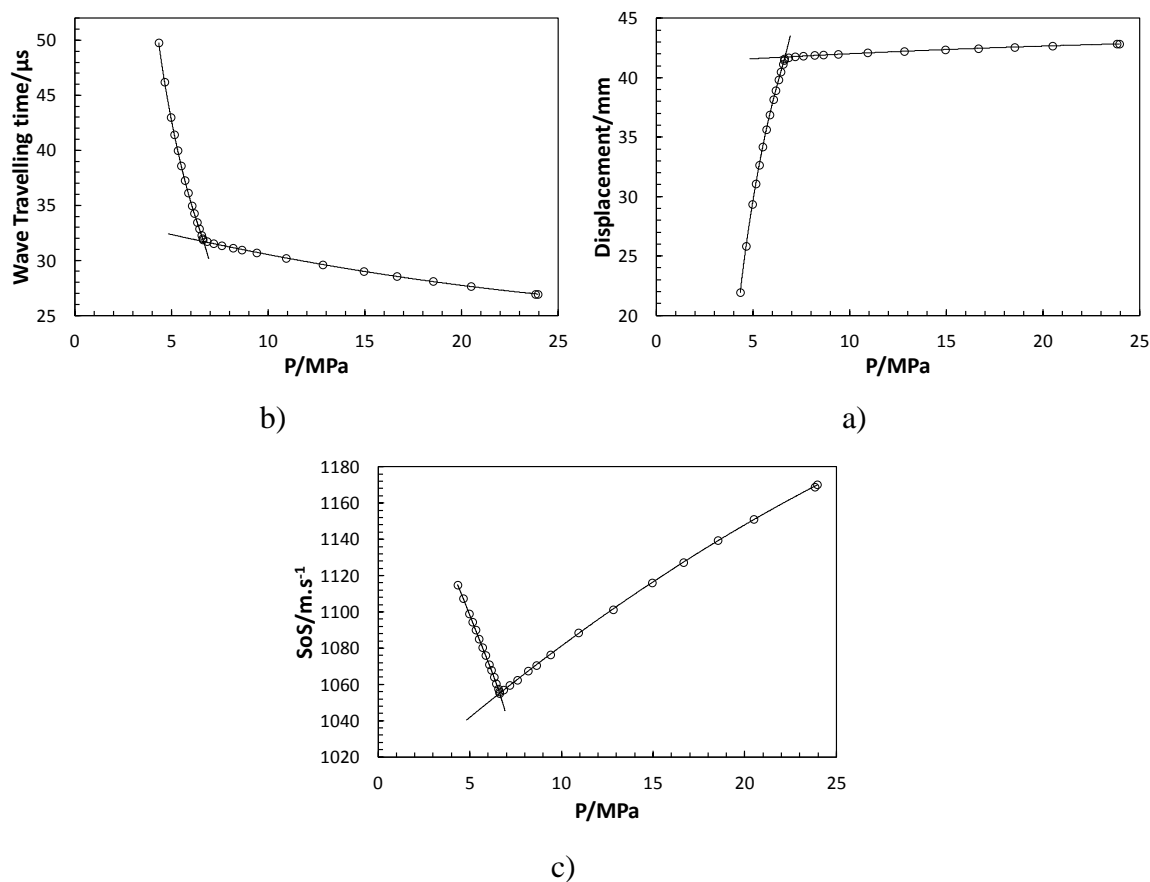


Figure 2.8 (a-c) Plots showing an example of bubble point determination of pure CO₂ at 344.3 K from plot of change in cell pressure versus (a) wave travelling time, (b) piston displacement and (c) speed of sound

A typical test to determine the bubble point of pure CO₂ (and impure CO₂) with /alkane mixture is as follows: the cell was charged with the test sample and set to the desired temperature for the measurement. The sample volume was then reduced by pumping a hydraulic fluid into the cell (behind the piston), at the opposite end to the sample. By this means the sample pressure was increased such that the sample was at a pressure significantly higher than the expected bubble point pressure. The cell was then rocked to mix the contents and ensure equilibrium. The sample volume was then increased step-wise by withdrawing liquid out of the cell. The piston displacement reading from the indicator is then taken at the corresponding pressure. At each step, mixing was continued until equilibrium was achieved, indicated by a constant pressure. The

stabilized equilibrium pressures and change in the displacement were plotted and the bubble point was indicated by a sharp change in the pressure. In addition, the bubble point was determined and confirmed from the break-point of the pressure versus travelling time and speed of sound plots as shown in Figure 2.8 (a-c). The plots show the bubble point estimation for pure CO₂ at 20.14 weight% in the alkane system. The calculated bubble point pressures of the mixture are 6.65, 6.66 and 6.63 MPa by using travelling distance, travelling wave time and speed of sound plots respectively. The values obtained from travelling wave time and the travelling distance are very close because they are based on the volume change. From Figure 2.8 (a-c), the speed of sound of the liquid mixture of CO₂/alkane starts to decrease as pressure decreases until the system reaches the bubble point pressure. This is mainly due to the reduction of the system bulk modulus. As the pressure decreases, below the bubble point, the dissolved CO₂ begins to escape from the alkane/CO₂ mixture. This dramatic phenomenon occurs because the compressibility decreases as CO₂ concentration decreases.

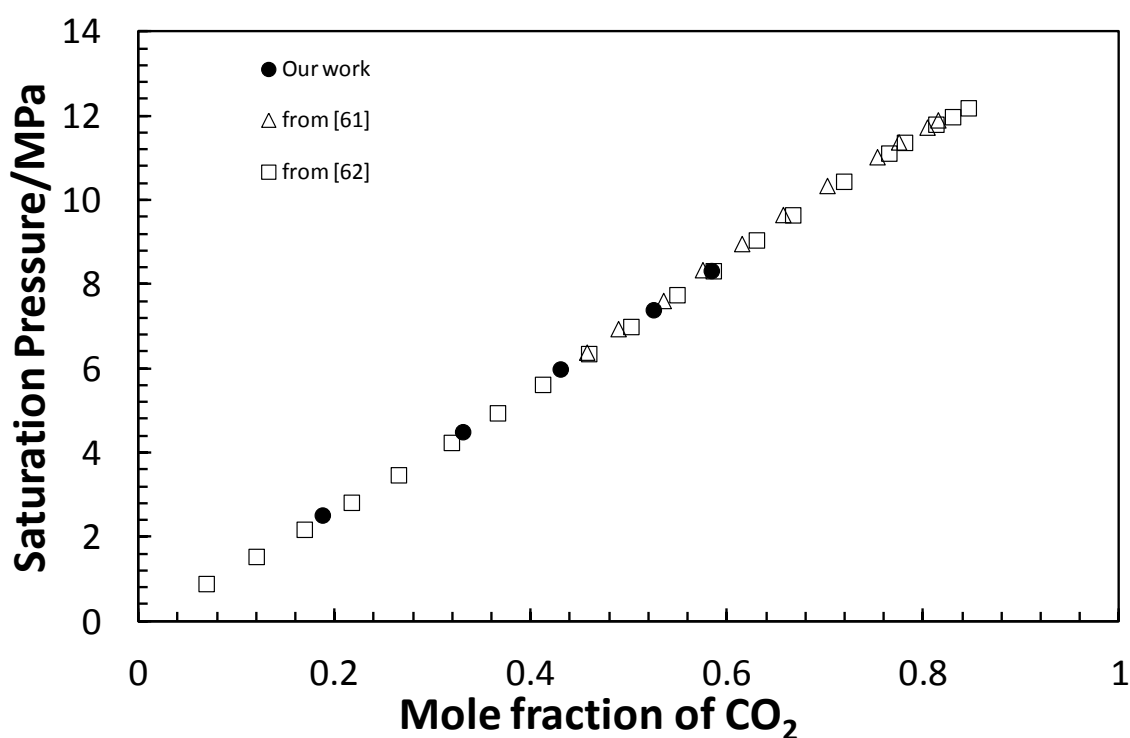


Figure 2.9 Saturation pressure measurements of CO₂ and n-decane system at 344.3 K

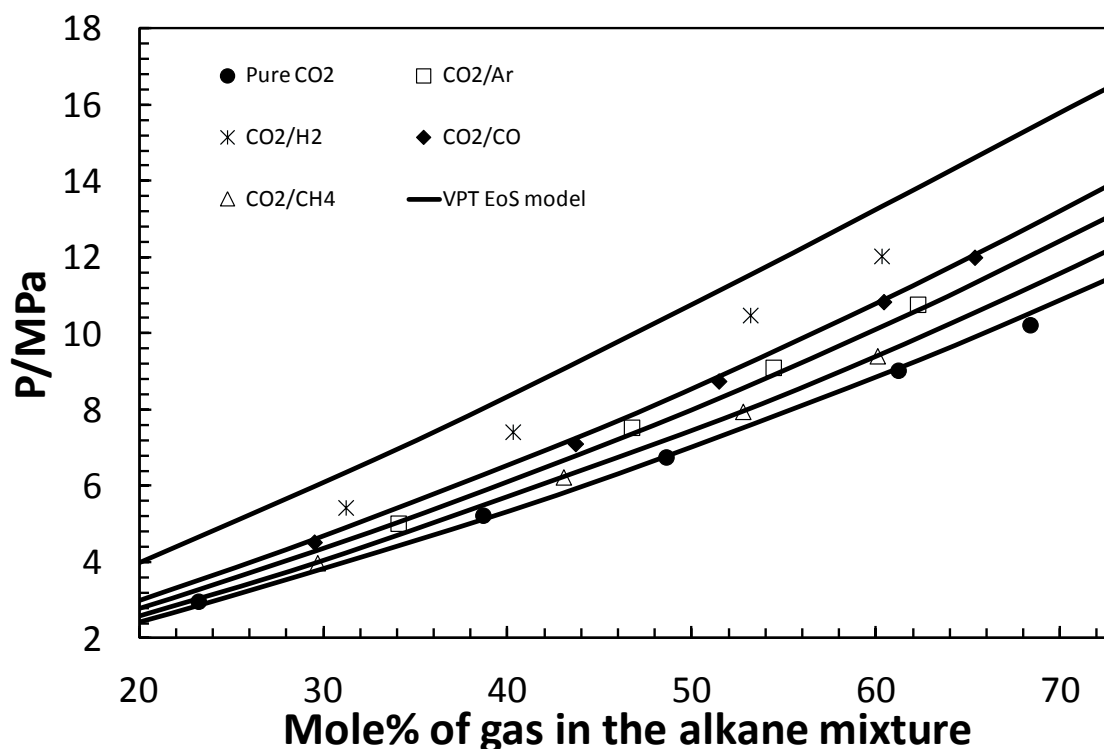


Figure 2.10 Effect of 5 mole% impurities on the saturation pressure of CO₂/alkane mixture

The experiment was first carried-out with CO₂ and n-decane system at 344.3 K for validation. Our results were in a good agreement with those performed by Nagarajan and Robinson [61] and Shaver et al. [62] as shown in Figure 2.9. Figure 2.10 shows the effect of 5 mole% impurity on the bubble point measurements of CO₂/alkane mixture. The increase in bubble pressure that is caused by additional mole percentage of gas is significantly greater in the presence of impurity than in pure CO₂. The change is as high as 1.2 MPa and 5.3 MPa at 30 mole% and 60 mole% respectively in the presence of hydrogen, while methane has the lowest impact to the CO₂/alkane saturation pressure (the saturation pressure of the mixture at 60 mole% CO₂/CH₄ gas is 2.4 MPa higher than in pure CO₂). The presence of argon and carbon monoxide in CO₂ stream can increase the saturation pressure of CO₂/alkane mixture to 4.1 MPa and 3.4 MPa. The VPT EOS is used to predict the experimental data of CO₂/alkane (from 7.46 weight% to 36.57 weight% CO₂) and CO₂/alkane in the presence of impurities (from 9.83 weight% to 30.46 weight%) as shown in Figure 2.10.

Usually, equations of state (EOSs) are used to predict vapour liquid equilibrium of such mixtures and acceptable accuracy can be found with the VPT model as shown in Figure

2.10. To obtain better results with this EOS, binary interaction parameters, k_{ij} , are required. These parameters in such mixtures are temperature dependent and they are usually calculated from isothermal experimental VLE data [63]. The interaction parameters used in the VPT EOS for the binary CO₂/n-decane, CO₂/n-tridecane and CO₂/n-hexadecane are 0.104, 0.098 and 0.095 respectively. The other binary alkane/impurities systems are set to their default values. The deviation on the saturation pressure is less than 3% in all the mixture of alkane containing pure and impure CO₂ except for the one containing hydrogen where it may exceed 10%. This is mainly due the low critical temperature of hydrogen compare to other impurities.

Table 2.5 SoS measurements of liquid alkane mixture in pure and impure CO₂

Pure CO ₂ in alkane mixture									
7.46		14.39		20.14		29.61		36.57	
Weight%		weight %		weight %		weight %		weight %	
P/	SoS/	P/	SoS/	P/	SoS/	P/	SoS/	P/	SoS/
MPa	m.s ⁻¹	MPa	m.s ⁻¹	MPa	m.s ⁻¹	MPa	m.s ⁻¹	MPa	m.s ⁻¹
23.02	1298	23.39	1229	23.95	1170	25.46	1089	28.33	1036
20.99	1291	21.30	1220	23.82	1169	22.86	1072	26.28	1022
20.09	1287	18.99	1208	20.49	1151	21.43	1062	23.32	999
18.68	1283	16.61	1195	18.53	1139	19.94	1051	21.78	987
16.50	1273	14.14	1181	16.66	1127	17.97	1036	19.11	965
14.19	1263	11.86	1167	14.94	1116	15.79	1020	17.11	949
12.28	1253	9.93	1155	12.82	1101	13.83	1005	14.85	930
9.15	1235	7.87	1140	10.92	1089	11.83	989	12.64	912
7.12	1223	6.01	1126	9.40	1076	9.97	974	11.80	906
5.88	1215	5.42	1122	8.63	1071	9.19	968	11.08	900
4.43	1205			8.18	1067				
3.72	1200			7.59	1062				
3.32	1198			7.17	1060				
				6.83	1057				
CO/CO ₂ in alkane mixture									
9.89		16.89		21.73		28.55		33.07	
weight %		weight %		weight %		weight %		weight %	

P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹
27.40	1274	39.76	1271	40.93	1234	39.29	1170	39.17	1134
22.14	1249	34.43	1248	35.61	1208	33.14	1136	30.30	1081
17.27	1224	28.88	1221	29.50	1176	27.74	1103	22.11	1022
13.57	1203	23.18	1190	24.51	1147	22.51	1067	16.35	971
8.19	1170	16.33	1148	18.72	1110	17.48	1028	11.84	923
5.83	1154	10.43	1107	14.30	1078	13.15	990		
		7.79	1087	10.43	1047	10.81	968		
Ar/CO ₂ in alkane mixture									
12.07		18.89		24.07		30.46			
weight %		weight %		weight %		weight %			
P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹
37.54	1294	41.33	1253	39.26	1197	36.68	1129		
30.47	1263	34.05	1219	29.57	1134	32.10	1103		
24.91	1237	27.49	1186	20.93	1092	25.04	1056		
17.25	1196	20.74	1147	15.25	1051	16.45	988		
13.01	1172	15.39	1113	10.60	1012	11.60	941		
7.16	1134	10.57	1078						
4.98	1119	7.44	1054						
H ₂ /CO ₂ in alkane mixture									
10.33		14.64		22.39		27.84			
weight %		weight %		weight %		weight %			
P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹	P/ MPa	SoS/ m.s ⁻¹
31.04	1284	28.20	1237	34.26	1188	34.18	1138		
27.81	1270	23.05	1206	28.26	1157	30.29	1116		
22.78	1247	19.32	1185	23.58	1128	24.71	1079		
16.66	1214	15.75	1162	19.22	1098	19.91	1043		
10.50	1178	10.15	1125	13.97	1059	14.34	995		
6.79	1154	7.28	1105	10.35	1031	11.79	976		
5.43	1144								

CH ₄ /CO ₂ in alkane mixture							
9.83		16.55		22.4		27.98	
weight %		weight %		weight %		weight %	
P/	SoS/	P/	SoS/	P/	SoS/	P/	SoS/
MPa	m.s ⁻¹	MPa	m.s ⁻¹	MPa	m.s ⁻¹	MPa	m.s ⁻¹
32.70	1300	36.32	1254	39.57	1221	44.62	1195
26.74	1273	27.98	1213	34.07	1193	36.70	1155
19.90	1238	20.66	1173	28.02	1160	28.79	1109
16.10	1219	14.65	1136	19.97	1109	22.05	1063
10.20	1185	9.74	1100	14.00	1066	15.43	1010
6.52	1161	6.64	1076	11.66	1047	11.36	972
4.59	1147			8.35	1019		
4.22	1144						

The injection of gas in a liquid reduces the sound speed in the liquid [64-67]. The speed of sound in an air water mixture can fall to 20 m.s⁻¹ while the speed of pure air and pure water are 1440 m.s⁻¹ and 340 m.s⁻¹ respectively [67]. This occurs because the mixture has the compressibility of gas and the density of a liquid [68]. The results of speed of sound under different pressures and mixture compositions are listed in Table 2.5. The results are given for CO₂/alkane mixture plus CO₂/alkane mixture containing impurities (CO, Ar, H₂ and CH₄). The mole percentage of impurities in the binary CO₂/impurities system is 5 mole%. The measurements were performed for a liquid single phase above the saturation line. The speed of sound measurements of CO₂/alkane are plotted in Figure 2.11. At a fixed mixture composition, the speed of sound increases as pressure increases. It is also clear that by increasing the concentration of CO₂ in the mixture, the SoS decreases. In order to investigate the effect of impurities on the SoS, a mathematical correlation was developed to match the data, taking the following form:

$$SoS_{co_2/alkane} = \frac{1282 - 1285x + 25.51P}{1 + 0.128x + 0.0154P} \quad (2.27)$$

P is the pressure in MPa, x is the weight fraction (0 to 1) and the calculated SoS is given in m.s⁻¹. The AAD and maximum deviation from the correlation are 0.17% and 0.53% respectively. The calculated reduction in the speed of sound is around 10 m.s⁻¹ per

adding 1 weight% of CO₂ at the same pressure and temperature condition. Figure 2.12-a and Figure 2.12-b show the effect of two impurities (Ar and H₂) on the speed of sound of the CO₂/alkane mixture. The presence of impurities on the CO₂ stream reduces slightly the speed of sound of the mixture. The effect is found in the range from 1 to 3% reduction. The reduction on the speed of sound increases as the pressure increases. The reduction on the property is not significantly affected by the type of impurities used.

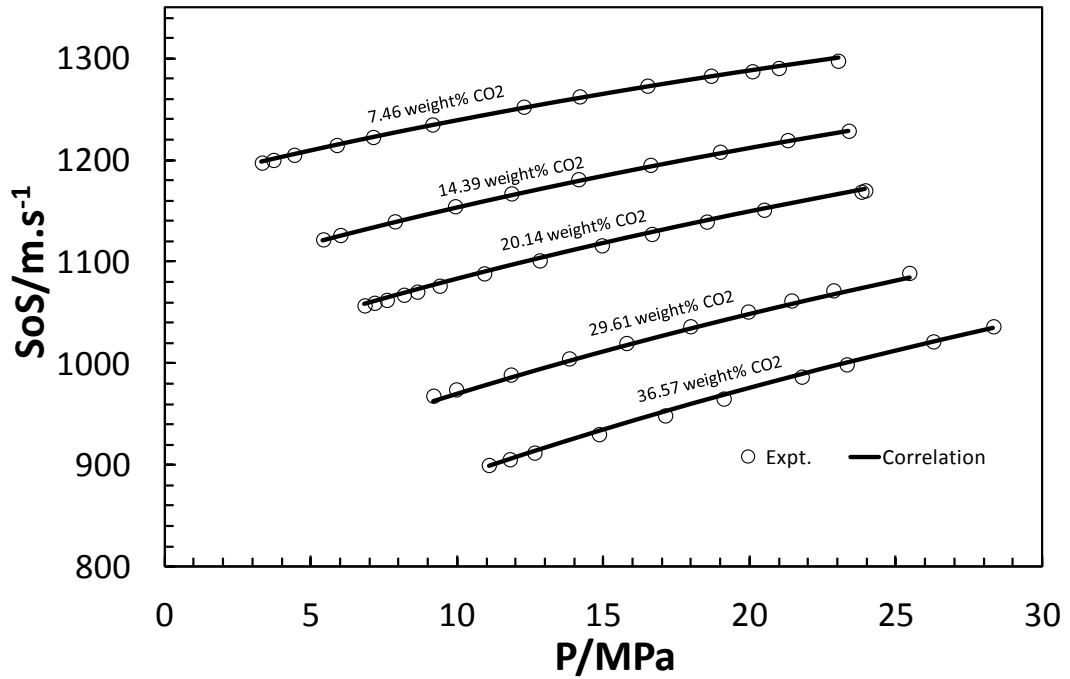


Figure 2.11 The speed of sound of CO₂/alkane system at 344.3 K, data are correlated using Equation 2.27

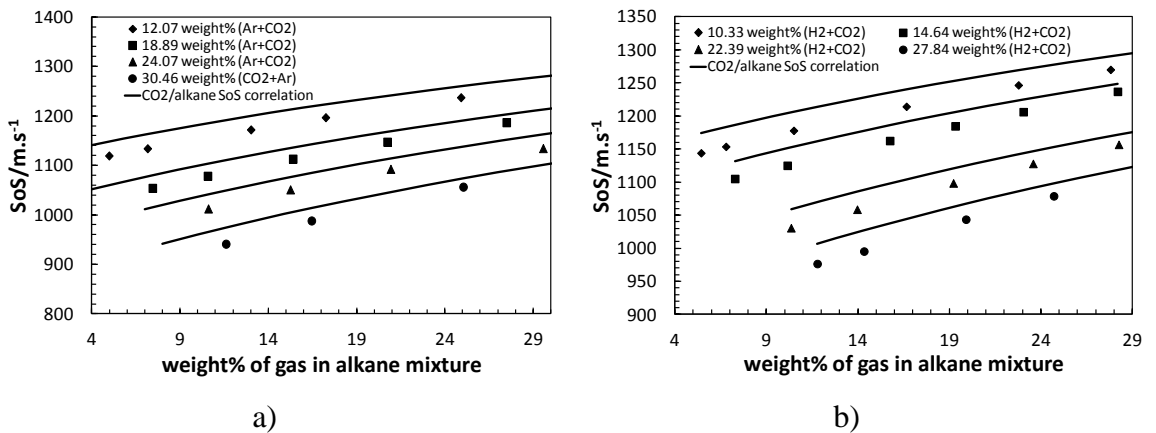


Figure 2.12 (a-b) Effect of impurities on speed of sound property of CO₂/alkane system at 344.3 K, the SoS of CO₂/alkane mixture is correlated using Equation 2.27 at the same condition of temperature and mixture composition.

2.4 Conclusions

Extensive experimental data on VLE of CO₂/impurities system were collected in order to tune the three EOSs (PR, VPT, and SRK). The effects of impurities of CO₂ systems on the VLE and critical curve properties were investigated for eight binary mixtures (impurities: H₂, CH₄, CO, Ar, O₂, N₂, SO₂ and H₂S). The three EOSs predict the saturation pressure of binary mixtures with similar accuracy (4.06% to 4.39%) in all the tested impurities system. The average deviation is as high as 10% using the three EOSs in the binary CO₂/H₂ Because of the low critical temperature of hydrogen component. The prediction of VLE on the CO₂/CH₄ binary system has the lowest deviation, i.e., less than 2.8%. The presence of impurities in CO₂ rich system tends to raise the critical pressure of CO₂ except for H₂S. The impurities H₂, CH₄, CO, Ar, O₂ and N₂ have a negative impact by increasing the CO₂ critical temperature while SO₂ and H₂S have a positive impact.

Few saturation pressure points were experimentally measured for CO₂/alkane mixture with pure and impure CO₂ at 344.3 K. The effect of four impurities were investigated, Ar, H₂, CO and CH₄. The increase on the saturation pressure of the mixture is ranged from the highest to the lowest impact as the following order: H₂, CO, Ar, and CH₄. In addition, the effects of impurities on the speed of sound of the previous mixtures were studied above the saturation line. Two key findings can be highlighted: 1) the impurities reduce approximately the speed of sound equally at the same conditions of pressure and temperature; 2) The speed of sound decreases as the amount of gas (pure or impure CO₂) increases.

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CHAPTER 3: SPEED OF SOUND AND ISOTHERMAL COMPRESSIBILITY OF CO₂-RICH SYSTEMS

3.1 Introduction

The speed of sound in CO₂ and CO₂-rich systems is an important physical property. It is required for the design of pulsation dampeners (natural frequency of the dampener is very much dependent on gas sound velocity) and/or the calculation of the velocities in throat in supersonic nozzles [1]. The sound speed is used to locate hydrate plugs and other obstructions in gas pipelines [2], and in testing pipelines for leaks [3]. It is required for the study of vibrations in high speed, high pressure gas compressors [4]. It also can have many applications, e.g., determining some thermodynamic properties, 4-D seismic, or monitoring compositional changes. The speed of sound is an important parameter in a depressurization case, because it determines how fast the pressure drop will propagate through the pipe [5]. Speed of sound is important for modelling and it is used to predict other fluid properties such as density and heat capacity ratio. Density and other thermodynamic properties that are predicted from the speed of sound data of liquids at high pressure can be found in literatures [6-10].

In this work, speed of sound and isothermal compressibility of CO₂ and CO₂/impurities system are investigated experimentally in the liquid region from 268.15 to 301.15 K and pressure up to 45 MPa. The isothermal compressibility of CO₂ and impure CO₂ is predicted by a proposed volume shift correction based on the general Peng Robinson model.

3.2 Experimental Section

3.2.1 *Experimental Apparatus*

The equipment used to measure the speed of sound was specially designed for measuring acoustic properties of high pressure fluids. It has a maximum volume of 280 ml and a length of 10.05 cm which can be adjusted by a movable piston. The system temperature is controlled by circulating coolant from a cryostat which is capable to

maintaining the cell temperature stability to within better than 0.05 K. The equilibrium cell is well isolated from the outside temperature. The system pressure is regulated by a pump connected to the non-sample side of moveable piston. The piston displacement is measured by a digital piston displacement indicator attached to appropriate end of the cell (accuracy higher than 10^{-2} cm). The sample is introduced into the cell directly from a sampling cylinder connected to a cell valve as shown in the schematic diagram of the ultrasonic set-up, see Figure 3.1. An ultrasonic pulser/receiver is used to transmit ultrasonic signal and receive it after passing through the fluid sample in the cell. After pre-amplification, the pulser/receiver sends the signal to a digital oscilloscope for analyses of the received signals. The total distance travelled by the ultrasound pulse was calibrated by measuring the time of flight in deionised water. The speed of sound in water was taken from [11]. The distance calibration was confirmed by testing at different temperatures (278.15 K, 288.15 K and 298.15 K).

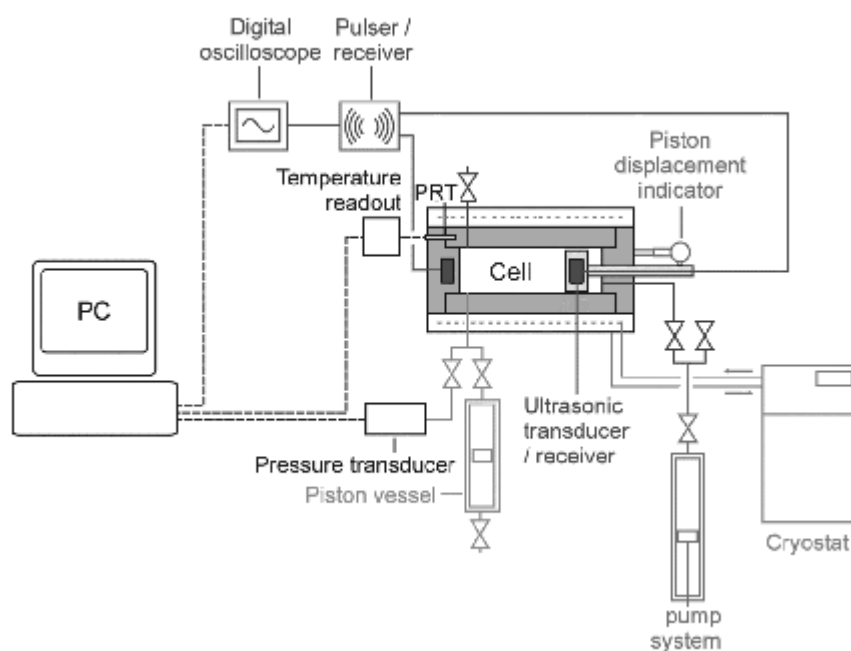


Figure 3.1 Schematic diagram of the ultrasonic set-up

3.2.2 Materials

The following compounds were used to make the different synthetic mixtures studied in this work:

- Carbon Dioxide and methane supplied by Air Products, Research Grade

- Nitrogen, argon, oxygen, hydrogen and carbon monoxide supplied by BOC, Research Grade

Binary and multi-component mixtures of CO₂/impurities system are listed in [Table 3.1](#). The mixtures were prepared gravimetrically (uncertainty ± 0.3 mole %)

Table 3.1 Synthetic gas composition of CO₂/impurities mixtures

Mixture	CO ₂	N ₂	CH ₄	CO	O ₂	Ar	H ₂
CO ₂ +N ₂		4.44	---	---	---	---	---
CO ₂ +CH ₄		---	4.61	---	---	---	---
CO ₂ +CO		---	---	4.31	---	---	---
CO ₂ +O ₂	Balance	---	---		6.52	---	---
CO ₂ +Ar		---	---	---	---	4.46	---
CO ₂ +H ₂		---	---	---	---	---	4.53
CO ₂ +Ar+CO (MIXa)		---	---	2.29	---	1.16	---
CO ₂ +CH ₄ +H ₂ +N ₂ (MIXb)		1.78	2.1	---	---	---	0.8

3.2.3 Procedures

The equilibrium cell and its loading lines were vacuumed prior to introducing CO₂ or CO₂ mixtures. The prepared sample was pressurised to 40 MPa in a sampling cylinder to ensure that the mixture is loaded into the cell as a single liquid phase. Then the system was pressurised by injecting hydraulic fluid into the non-sample side of the piston. The sample is given time to reach equilibrium. The equilibrium is confirmed by observing constant reading in both pressure and waveform signal. The mean wave travelling time and the corresponding displacement taken from the digital piston displacement indicator are recorded in a spreadsheet. The calibration equation obtained from the speed of sound in water is used for calculating the speed of sound in liquid CO₂ and CO₂ mixtures. The accuracy of the measurement is estimated within ± 1 m/s.

3.3 Modelling Section

The Peng Robison-EoS is classified as a two-parameter cubic equation and has become one of the models of choice in the process modelling. The PR-EoS was chosen to model

the volumetric behaviour of CO₂ and CO₂/impurities mixtures with the quadratic van der Waals mixing rules and take the following form:

$$P = \frac{RT}{(V-b)} - \frac{a}{[V(V+b)+b(V-b)]} \quad (3.1)$$

Where,

$$a = a(T_c)\alpha(T_r, \omega) \quad (3.2)$$

$$\alpha(T_r, \omega) = [1 + m(1 - T_r^{0.5})]^2 \quad (3.3)$$

$$m = 0.37464 + 1.54226\omega - 0.26992\omega^2 \quad (3.4)$$

$$a(T_c) = 0.45724 \frac{R^2 T_c^2}{P_c} \quad (3.5)$$

$$b = 0.0778 \frac{RT_c}{P_c} \quad (3.6)$$

In general an equation of state is developed first for pure compounds, and then extended to mixtures through the use of mixing rules for combining the parameters of pure compounds. For mixtures, the conventional van der Waals mixing rules were used in this work:

$$a_m = \sum_i^n \sum_j^n x_i x_j a_{ij} \quad (3.7)$$

$$b_m = \sum_i^n x_i b_i \quad (3.8)$$

$$a_{ij} = (1 - k_{ij}) \sqrt{a_i a_j} \quad (3.9)$$

Where, a_i , and, b_i , are calculated from [Equations 3.2](#) and [3.6](#) using the critical pressure, P_c , critical temperature, T_c , and acentric factor, ω , for each component. In [Equation 3.9](#), k_{ij} is the binary interaction parameter.

The Mathias-Copeman [\[12\]](#) alpha function with three adjustable parameters is used to improve vapour pressure calculation of pure CO₂

$$\text{if } T < T_c \quad \alpha(T_r) = \left[1 + m_1(1 - \sqrt{T_r}) + m_2(1 - \sqrt{T_r})^2 + m_3(1 - \sqrt{T_r})^3 \right]^2 \quad (3.10)$$

$$\text{if } T \geq T_c \quad \alpha(T_r) = \left[1 + m_1(1 - \sqrt{T_r}) \right]^2 \quad (3.11)$$

Where m_1 , m_2 and m_3 are the three EoS-dependent adjustable parameters, the parameters were tuned for PR-EoS. Their values for carbon dioxide are 0.7138, -0.4422 and 2.4364, respectively.

Significant deviations can be observed (later in this chapter) in the prediction of liquid density compared to the experimental data. The liquid density of carbon dioxide predicted by the Peng Robinson (PR) equation of state is often off by 10% or more at temperature and pressure encountered in most of transport conditions. In this work, a volume shift is added to the PR model in order to improve liquid density of both pure and impure CO₂. The correction is defined as follow:

$$V_{corrected} = V_{EoS} + VC \quad (3.12)$$

where V_{EoS} is the volume estimated from the PR equation of state and VC is the volume-shift parameter, which in our calculation is defined as:

$$VC = s + \frac{a_c T_r}{b_c + T_r(304.2 + c_c P_r)} \quad (3.13)$$

Where $a_c = -0.001254$, $b_c = -291.6$, $c_c = -0.0655$ and $s = -3.004$

The modified volume correction is applied for CO₂ and CO₂/impurities system (The effect of adding the above correction to the other thermodynamic properties is not tested). The binary interaction parameters (k_{ij}) of CO₂ and its impurities can be found in [Chapter 2](#). The molecular weight of the mixture is defined by:

$$M_{Mix} = M_{CO_2} x_{CO_2} + M_{impurity} (1 - x_{CO_2}) \quad (3.14)$$

Where, M_{CO_2} and $M_{impurity}$ are the molecular weight of CO₂ and the impurity respectively. x_{CO_2} is the mole fraction of CO₂ in the mixture.

Span and Wagner [13] have the latest review of the thermodynamic properties of CO₂. Their work covered a wide range of pressure and temperature including the critical area based on large amount of experimental data. They generated highly accurate tables based on a complex equation of state. The tabulated data of their work is used in this work in order to compare our modified model of CO₂ liquid density. Figure 3.2 shows the liquid density of CO₂ from 228 through 301.15 K and pressure up to 50 MPa. Our modified PR-EoS model matches the results tabulated by Span and Wagner (1996) and the average deviation is less than 0.1%.

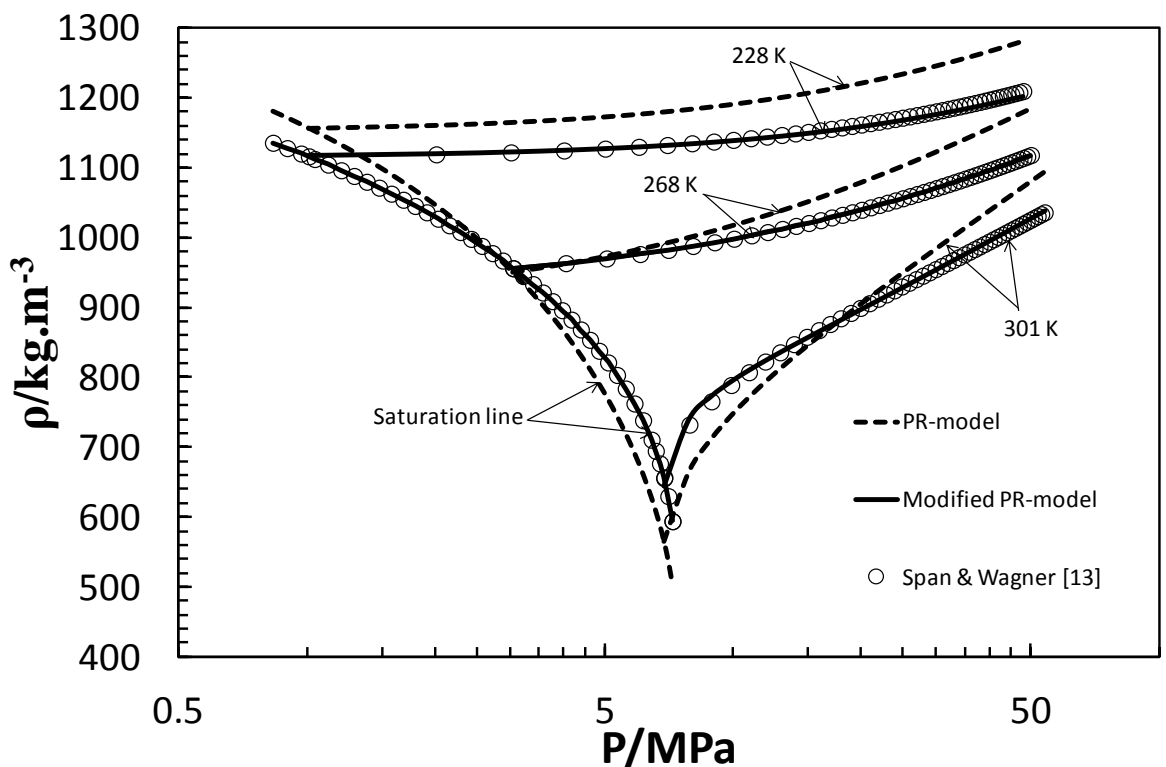


Figure 3.2 Liquid and saturation liquid density of CO₂

3.4 Results

3.4.1 Speed of Sound

Six isothermal measurements of speed of sound were carried-out from 268.15 K to 301.15 K and pressure up to 45 MPa for pure CO₂ and the CO₂/impurities system. Speed of sound in liquid CO₂ as a function of pressure and temperature is shown in Figure 3.3 and listed in Table 3.2. The speed of sound of liquid CO₂ increases smoothly and monotonously at a constant temperature with increasing pressure, while at a similar

pressure away from the saturation pressure, it decreases slightly with increasing temperature. The reduction in the sound property becomes significant as it approaches the saturation line. The variation of the speed of sound with pressure is high compared to its temperature dependence.

Table 3.2 Speed of sound and isothermal compressibility measurements of CO₂ and CO₂/impurities system

Pure CO ₂								
268.15 K			278.15 K			283.15 K		
P /MPa	K _T /MPa ⁻¹	SoS /m.s ⁻¹	P /MPa	K _T /MPa ⁻¹	SoS /m.s ⁻¹	P /MPa	K _T /MPa ⁻¹	SoS /m.s ⁻¹
3.63	0.00721	594	6.89	0.00898	557	7.34	0.01100	519
7.37	0.00585	649	8.11	0.00817	577	8.34	0.01004	538
10.08	0.00509	682	10.34	0.00698	610	10.24	0.00854	570
14.00	0.00427	723	13.85	0.00569	654	13.41	0.00679	614
18.00	0.00369	760	17.40	0.00483	693	17.36	0.00545	661
21.93	0.00330	792	20.93	0.00421	726	20.71	0.00471	694
24.94	0.00306	814	24.10	0.00377	753	24.28	0.00413	727
28.35	0.00284	838	27.84	0.00336	782	27.25	0.00373	751
31.63	0.00264	861	30.82	0.00311	804	30.82	0.00337	779
35.36	0.00243	883	34.43	0.00290	829	34.56	0.00311	805
41.13	0.00216	916	37.45	0.00274	848	38.47	0.00286	831
			41.14	0.00240	870	41.22	0.00255	848
288.15 K			293.15 K			301.15 K		
6.93	0.01042	459	8.00	0.01837	434	9.43	0.02592	391
10.23	0.00926	528	10.41	0.01357	491	11.14	0.01952	438
14.35	0.00787	592	13.89	0.00941	551	14.51	0.01198	505
17.46	0.00690	630	16.89	0.00748	593	17.18	0.00917	546
20.64	0.00598	665	21.08	0.00597	641	21.74	0.00708	604
23.99	0.00513	697	24.08	0.00522	671	25.08	0.00606	640
28.23	0.00425	734	27.50	0.00450	703	27.43	0.00537	662
31.33	0.00376	757	31.23	0.00393	733	31.60	0.00442	699
34.65	0.00339	782	34.27	0.00368	757	34.63	0.00418	723
38.05	0.00318	805	37.96	0.00347	783	37.89	0.00406	748
41.64	0.00317	828	41.96	0.00250	809	41.38	0.00366	772
CO ₂ +N ₂ (4.44 % N ₂)								
268.15 K			278.15 K			283.15 K		
9.45	0.00678	618	9.49	0.00954	534	10.51	0.01134	512
11.14	0.00626	640	11.08	0.00859	561	12.23	0.00977	542
13.10	0.00571	663	13.75	0.00722	600	14.20	0.00834	572
16.75	0.00481	702	16.91	0.00598	640	17.44	0.00669	615

20.30	0.00413	735	20.51	0.00501	679	20.73	0.00562	653
24.52	0.00355	771	24.20	0.00436	714	24.31	0.00485	688
27.87	0.00323	797	27.51	0.00393	742	27.83	0.00428	720
31.39	0.00300	822	31.15	0.00352	771	31.16	0.00383	747
35.06	0.00276	847	34.86	0.00318	798	34.72	0.00345	774
38.19	0.00249	867	38.58	0.00312	823	38.30	0.00330	799
288.15 K			293.15 K			301.15 K		
13.29	0.01093	523	13.71	0.01282	495	15.03	0.01520	464
15.80	0.00875	562	15.76	0.01052	529	16.57	0.01303	491
18.31	0.00735	595	18.08	0.00868	562	18.62	0.01080	523
20.99	0.00634	627	21.17	0.00708	602	20.98	0.00896	555
23.92	0.00554	658	24.24	0.00605	635	24.12	0.00733	593
27.35	0.00480	691	27.92	0.00518	671	27.96	0.00606	634
30.97	0.00420	722	30.79	0.00465	697	30.95	0.00535	662
34.82	0.00378	752	34.00	0.00420	723	35.20	0.00456	698
37.51	0.00355	772	38.07	0.00383	749	38.45	0.00412	723
39.78	0.00323	787	41.64	0.00350	779	41.52	0.00382	745
CO ₂ +CH ₄ (4.61 % CH ₄)								
268.15 K			278.15 K			283.15 K		
8.87	0.00681	614	7.43	0.01232	503	9.58	0.01251	502
11.09	0.00613	644	8.54	0.01099	526	10.99	0.01068	529
13.86	0.00531	678	10.54	0.00909	562	13.94	0.00837	576
17.20	0.00449	713	13.89	0.00702	610	16.95	0.00711	617
20.75	0.00388	746	17.68	0.00568	657	21.09	0.00584	663
24.41	0.00349	776	21.33	0.00486	695	24.71	0.00476	700
27.72	0.00325	802	24.43	0.00433	723	27.91	0.00402	728
31.29	0.00301	828	28.30	0.00380	756	31.21	0.00369	753
34.82	0.00273	852	31.41	0.00347	780	34.94	0.00366	782
38.05	0.00253	872	35.27	0.00316	808	37.37	0.00324	799
288.15 K			293.15 K			301.15 K		
8.47	0.01957	429	9.26	0.02546	397	10.36	0.03579	354
10.77	0.01418	483	10.74	0.01968	439	12.27	0.02454	411
13.82	0.01007	538	12.96	0.01389	488	14.81	0.01571	467
16.95	0.00792	584	18.12	0.00837	570	17.98	0.01087	521
20.89	0.00639	632	20.76	0.00726	603	20.37	0.00923	555
24.46	0.00534	670	23.95	0.00619	639	24.26	0.00736	602
27.87	0.00458	701	27.74	0.00506	676	27.55	0.00595	637
31.22	0.00419	730	31.05	0.00451	705	30.81	0.00521	667
34.40	0.00399	755	34.60	0.00431	734	34.56	0.00515	700
37.82	0.00300	780	37.87	0.00320	759	38.22	0.00252	728
CO ₂ +H ₂ (4.53 %)								
268.15 K			278.15 K			283.15 K		

10.54	0.01004	620	10.77	0.01148	542	9.71	0.01851	492
13.02	0.00816	652	13.96	0.00850	592	12.73	0.01233	542
15.16	0.00687	677	17.29	0.00682	634	15.03	0.00962	574
17.52	0.00579	701	20.88	0.00575	673	17.58	0.00783	609
20.58	0.00485	725	24.48	0.00499	708	20.66	0.00656	646
25.39	0.00407	767	28.11	0.00437	739	24.11	0.00561	681
28.41	0.00376	790	31.33	0.00398	765	27.96	0.00479	716
31.38	0.00347	812	34.74	0.00369	790	31.35	0.00433	744
34.66	0.00317	834	38.33	0.00319	814	35.16	0.00401	773
38.07	0.00314	856				38.27	0.00335	795
288.15 K			293.15 K			301.15 K		
9.78	0.02310	458	10.07	0.03112	418	13.34	0.02310	420
13.39	0.01339	521	13.88	0.01590	493	15.86	0.01667	472
17.23	0.00906	576	17.71	0.01023	553	18.21	0.01289	512
20.99	0.00731	622	21.28	0.00830	599	21.35	0.00997	556
24.55	0.00621	660	24.57	0.00708	635	24.31	0.00838	591
28.17	0.00527	694	28.25	0.00583	671	27.63	0.00713	627
31.48	0.00472	722	31.54	0.00520	700	31.02	0.00611	659
35.25	0.00442	751	34.96	0.00495	728	34.32	0.00539	688
38.38	0.00352	774	38.22	0.00336	752	37.76	0.00492	715
						40.75	0.00435	737
CO ₂ +O ₂ (6.52 %)								
268.15 K			278.15 K			283.15 K		
10.42	0.00665	599	11.46	0.00920	538	8.93	0.01809	444
13.18	0.00611	635	12.18	0.00883	550	11.54	0.01281	500
17.57	0.00502	686	14.55	0.00769	584	14.22	0.00951	546
21.09	0.00427	720	17.15	0.00658	618	17.26	0.00744	589
24.28	0.00378	749	20.36	0.00551	655	20.84	0.00613	632
28.27	0.00339	780	23.94	0.00470	691	24.10	0.00534	666
31.36	0.00316	803	28.01	0.00417	726	27.83	0.00456	700
34.54	0.00294	825	31.48	0.00388	754	31.16	0.00400	728
37.65	0.00272	845	34.82	0.00358	779	35.03	0.00361	758
40.84	0.00257	865	37.97	0.00318	801	38.35	0.00342	781
			40.52	0.00273	817	41.00	0.00303	799
288.15 K			293.15 K			301.15 K		
10.64	0.01809	438	12.92	0.01587	454	12.80	0.02219	391
13.87	0.01206	505	14.73	0.01301	489	13.98	0.01969	419
17.06	0.00885	555	17.26	0.01020	530	17.63	0.01354	488
20.79	0.00694	603	21.84	0.00734	590	21.39	0.00961	543
24.01	0.00596	638	24.28	0.00644	618	24.12	0.00793	577
28.33	0.00495	680	27.68	0.00552	652	27.35	0.00674	612
31.21	0.00441	705	31.09	0.00483	683	31.15	0.00580	648
34.67	0.00399	732	34.54	0.00432	712	34.59	0.00503	678

38.15	0.00371	758	38.16	0.00389	739	37.91	0.00442	705
40.78	0.00326	777	40.82	0.00344	758	40.83	0.00444	727
CO ₂ +Ar (4.46 %)								
268.15 K			278.15 K			283.15 K		
9.31	0.00636	607	10.26	0.00816	547	11.66	0.01064	531
11.15	0.00596	635	12.49	0.00745	581	14.38	0.00814	573
13.69	0.00536	665	16.35	0.00632	632	17.51	0.00661	613
17.15	0.00460	703	18.35	0.00579	654	20.95	0.00563	652
20.73	0.00400	736	21.70	0.00498	688	24.29	0.00491	685
23.86	0.00362	763	24.66	0.00438	715	28.01	0.00423	717
28.01	0.00325	795	28.23	0.00380	745	31.44	0.00381	745
31.57	0.00296	820	31.33	0.00344	769	34.86	0.00355	771
35.01	0.00272	843	35.31	0.00318	798	38.32	0.00295	795
38.32	0.00269	863	38.58	0.00315	819			
288.15 K			293.15 K			301.15 K		
12.38	0.01512	505	10.60	0.01778	421	11.75	0.02614	382
15.02	0.01205	548	15.04	0.01249	512	14.64	0.01756	451
16.80	0.01025	577	18.20	0.00925	559	17.06	0.01328	494
19.94	0.00767	612	21.15	0.00720	597	19.66	0.01045	534
22.39	0.00621	639	24.59	0.00592	634	22.11	0.00868	565
25.11	0.00510	667	28.08	0.00533	668	24.79	0.00725	596
28.13	0.00442	694	31.21	0.00492	695	28.08	0.00598	630
31.18	0.00413	720	34.56	0.00443	722	31.12	0.00531	658
34.91	0.00399	749	37.85	0.00440	747	34.82	0.00494	689
38.25	0.00364	773				38.44	0.00349	718
41.75	0.00253	796						
CO ₂ +CO (4.31 %)								
268.15 K			278.15 K			283.15 K		
11.53	0.00641	622	11.98	0.00905	553	9.86	0.01469	475
14.38	0.00562	658	14.87	0.00754	597	14.47	0.00914	557
18.28	0.00481	700	17.75	0.00626	632	17.15	0.00745	595
20.27	0.00447	718	20.79	0.00531	666	20.91	0.00604	640
24.03	0.00389	752	24.04	0.00474	699	24.12	0.00524	673
27.65	0.00343	781	27.81	0.00437	731	27.90	0.00452	708
31.19	0.00311	808	30.72	0.00406	755	31.18	0.00404	735
34.59	0.00292	831	34.75	0.00342	785	34.73	0.00369	762
37.92	0.00274	852	38.18	0.00294	809	37.96	0.00342	786
40.77	0.00243	870	41.11	0.00323	829	40.91	0.00294	806
288.15 K			293.15 K			301.15 K		
10.31	0.01809	439	9.71	0.02832	376	10.37	0.04742	315
13.99	0.01159	515	13.52	0.01558	471	12.49	0.03008	386
16.89	0.00883	560	16.84	0.01039	529	15.08	0.01768	446
20.38	0.00700	605	20.50	0.00790	581	17.44	0.01245	489

24.10	0.00583	646	23.99	0.00664	620	20.80	0.00979	539
27.93	0.00494	683	27.91	0.00548	660	24.08	0.00847	582
31.46	0.00431	714	31.21	0.00470	691	27.55	0.00668	618
34.94	0.00393	741	34.95	0.00426	722	31.21	0.00507	654
38.27	0.00364	766	37.99	0.00403	744	34.72	0.00495	684
40.93	0.00310	784	40.74	0.00314	764	38.20	0.00511	717
						40.98	0.00154	733

CO ₂ +CH ₄ +N ₂ +H ₂ , MIXb,(2.10% CH ₄ , 1.78 % N ₂ , 0.8 % H ₂)								
268.15 K			278.15 K			283.15 K		
7.57	0.01055	564	8.10	0.01411	488	7.54	0.02091	415
10.39	0.00718	609	10.67	0.01030	535	10.53	0.01335	491
14.02	0.00529	656	13.78	0.00771	585	13.89	0.00901	551
17.15	0.00465	691	17.02	0.00626	627	17.54	0.00691	602
20.96	0.00415	728	20.71	0.00523	668	21.24	0.00578	645
24.00	0.00373	754	24.24	0.00448	702	24.57	0.00496	679
28.13	0.00319	787	27.85	0.00394	734	27.97	0.00425	710
31.26	0.00297	810	31.48	0.00370	763	31.41	0.00383	739
35.25	0.00286	837	34.63	0.00355	786	34.99	0.00361	766
38.36	0.00238	857	38.26	0.00254	812	38.47	0.00276	791
288.15 K			293.15 K			301.15 K		
7.89	0.03562	374	8.55	0.03868	330	10.76	0.03822	342
10.38	0.02028	446	11.40	0.02119	427	14.11	0.02045	429
13.62	0.01091	512	14.09	0.01303	485	17.15	0.01281	487
17.15	0.00796	567	17.17	0.00939	536	20.67	0.00933	540
20.88	0.00697	614	21.64	0.00747	595	24.15	0.00775	583
24.30	0.00576	651	24.29	0.00639	625	27.57	0.00638	620
28.12	0.00438	688	27.70	0.00507	659	31.44	0.00513	657
31.41	0.00409	716	31.23	0.00455	691	34.89	0.00478	687
35.05	0.00442	745	34.73	0.00463	720	38.18	0.00453	713
38.70	0.00130	771	38.15	0.00179	746	40.55	0.00288	731

CO ₂ +Ar+CO, MIXa, (1.16% Ar, 2.29% CO)								
268.15 K			278.15 K			283.15 K		
13.19	0.00566	650	12.88	0.00827	580	11.33	0.01060	517
16.60	0.00469	690	16.31	0.00655	625	14.14	0.00845	563
20.24	0.00406	727	18.84	0.00567	654	17.34	0.00679	606
23.95	0.00359	759	21.68	0.00497	683	20.57	0.00567	643
27.92	0.00319	790	24.80	0.00440	712	24.11	0.00485	679
31.43	0.00292	815	28.19	0.00392	741	27.86	0.00428	713
34.99	0.00273	839	31.32	0.00352	765	31.37	0.00392	742
38.55	0.00254	861	34.66	0.00315	790	34.56	0.00366	766
40.99	0.00225	876	38.43	0.00292	816	38.25	0.00323	793
			40.80	0.00301	831	41.06	0.00255	811
288.15 K			293.15 K			301.15 K		

10.25	0.01617	453	10.35	0.02073	410	14.28	0.01800	440
13.65	0.01084	520	13.95	0.01307	490	17.13	0.01255	493
17.36	0.00784	576	17.43	0.00917	546	20.70	0.00909	545
20.90	0.00635	619	20.95	0.00716	592	24.73	0.00715	594
24.71	0.00533	659	24.62	0.00595	634	28.39	0.00594	632
28.16	0.00461	691	28.02	0.00513	666	31.52	0.00512	662
31.48	0.00408	719	31.53	0.00445	697	34.97	0.00458	691
35.18	0.00371	748	34.98	0.00399	725	38.27	0.00426	717
38.44	0.00342	772	38.29	0.00369	750	41.22	0.00341	738
40.96	0.00289	789	40.97	0.00327	768			

In order to observe the effect of impurities, the speed of sound of six binary CO₂ mixtures and two multi-components CO₂/mixtures (Table 3.1) were measured in the same range of pressure and temperature. Figure 3.4 and Figure 3.5 show the experimental results of speed of sound property on both pure and impure CO₂ at 268.15 and 301.15 K respectively.

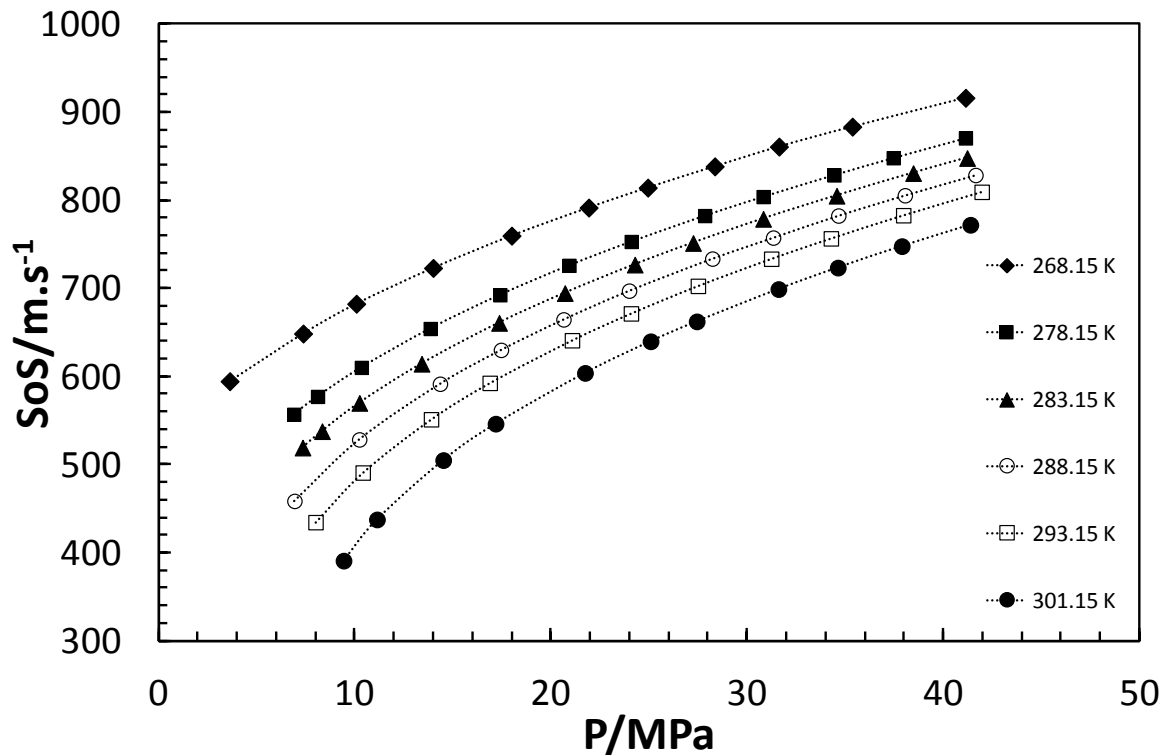


Figure 3.3 Measurements in speed of sound of CO₂

It can be seen that non-condensable impurities such as O₂, Ar, N₂, CO significantly reduce the speed of sound of CO₂ stream. The change is much higher as the temperature increases and it becomes more significant around the critical point of the mixture. The

binary CO₂/O₂ mixture (6.52% O₂) which contains the highest impurity concentration compared to the other investigated system has major effect on the property.

In order to analyse the effect of each impurity at different pressure and temperature conditions, the change in speed of sound (defined as $SoS_{pureCO_2} - SoS_{impureCO_2}$) is plotted at three different temperatures (268.15 , 278.15 and 293.15 K). Two CO₂/impurity systems were selected CO₂/CH₄ and MIXb CO₂+ (CH₄+H₂+N₂). Their effect on the CO₂ speed of sound is plotted in [Figure 3.6](#) and [Figure 3.7](#) respectively. The mole fraction of CO₂ in both systems is around 94.4%. The lines are plotted based on the experimental data (shown in Table 3.2) using cubic spline interpolation. The change is larger at pressures near the saturation line of the mixture and it becomes less significant as pressure increases. The change in speed of sound is strongly dependant on the pressure and the number of impurities present in CO₂ stream. In term of temperature dependency, the change is larger at higher temperatures near the saturation line. As the pressure increases the change reduces to a point where all the isothermal lines intersect. Beyond that point, the change in speed of sound becomes higher at low temperature and less significant at high temperature. This finding was observed in all the other tested CO₂ impurities systems. Speed of sound is a function of three properties: isothermal compressibility (K_T), density (ρ) and heat capacity (γ). The change of both density (see [Chapter 4](#)) and isothermal compressibility (see 3.4.2 in this chapter) of CO₂ in the presence of impurities is found to be consistent as pressure and temperature change. Therefore, more investigation is required to define the effect of heat capacity ratio on impure CO₂ stream and perhaps it can explain the above mentioned behaviour on speed of sound.

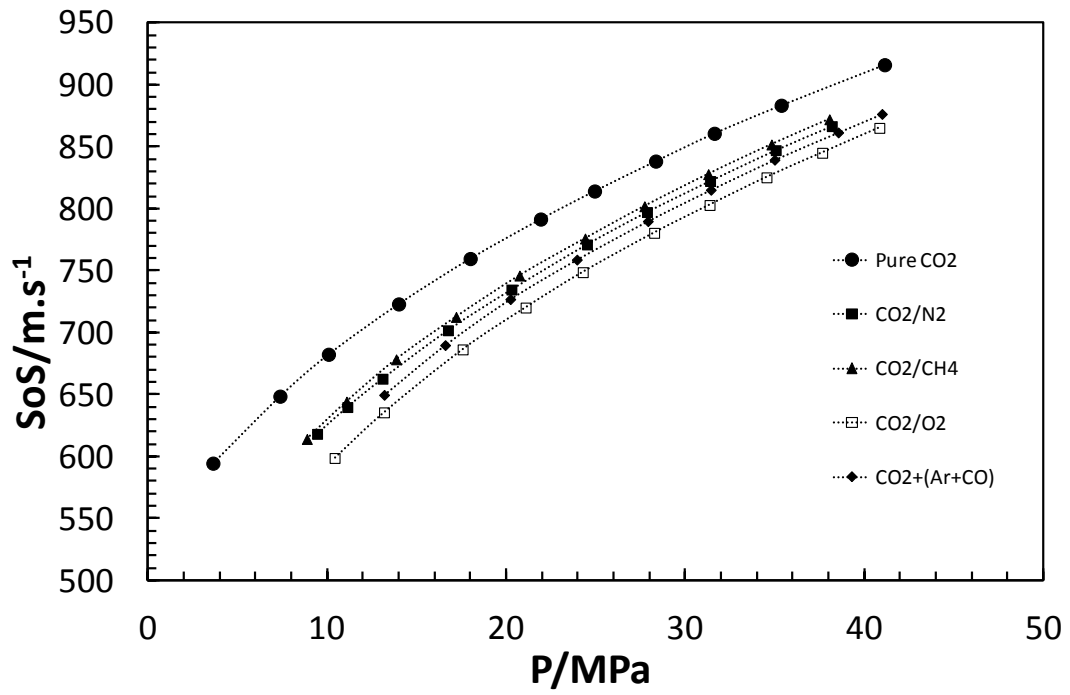


Figure 3.4 Effect of CO₂ impurities on speed of sound property at 268.15 K

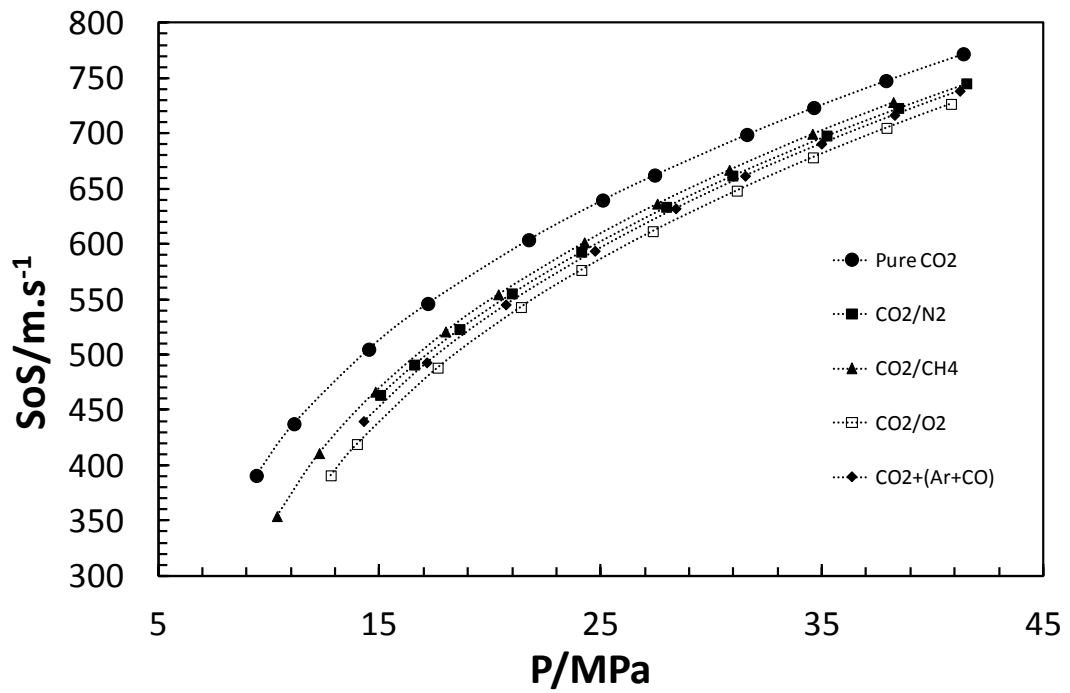


Figure 3.5 Effect of CO₂ impurities on speed of sound property at 301.15 K

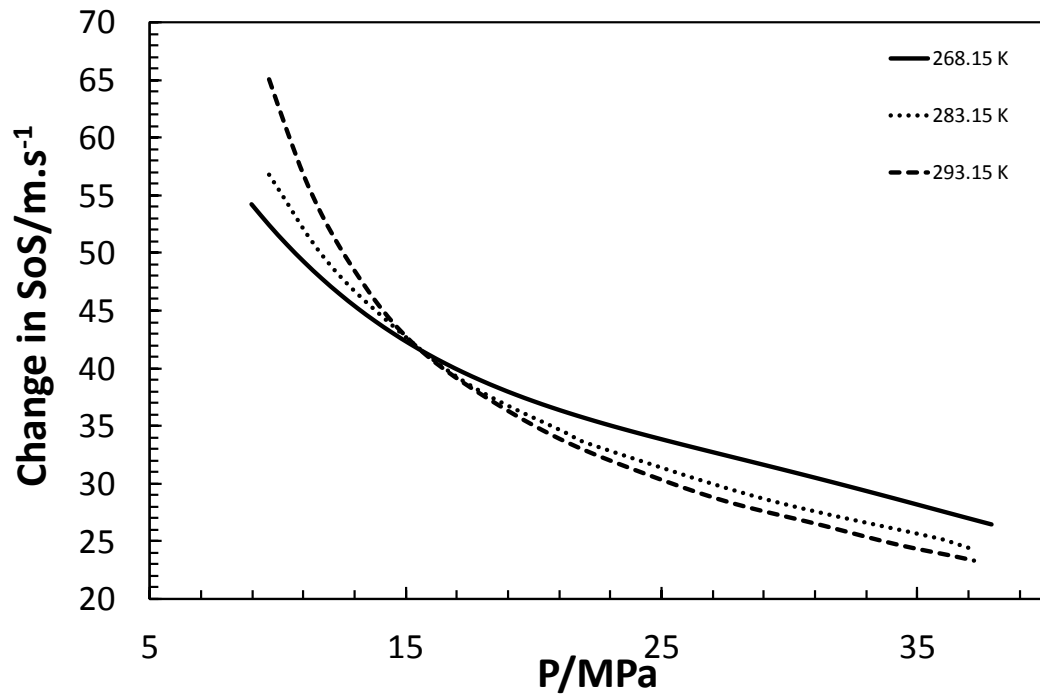


Figure 3.6 Change in speed of sound (defined as $SoS_{pureCO_2} - SoS_{impureCO_2}$) as a function of pressure in CO₂/CH₄ mixture

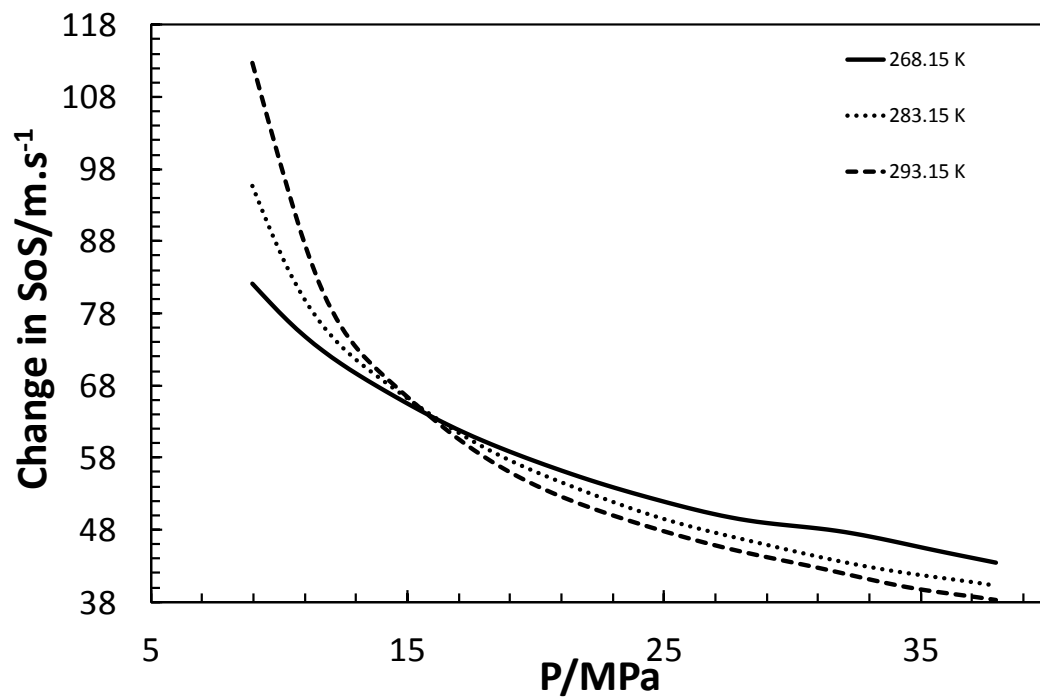


Figure 3.7 Change in speed of sound (defined as $SoS_{pureCO_2} - SoS_{impureCO_2}$) as a function of pressure in MIXb (CO₂-CH₄-H₂-N₂)

3.4.2 Isothermal Compressibility

The isothermal compressibility coefficients are required in several reservoir engineering applications such as transient fluid flow problems and also in the determination of physical properties of crude oils. The isothermal compressibility coefficient (K_T) is defined as the rate of change in volume with pressure per unit volume of fluid at constant temperature. It is also defined as the quantitative expression of the sensitivity of density to pressure so it can be written in either of the following forms:

$$K_T = -\frac{1}{v} \left(\frac{\partial v}{\partial p} \right)_T \quad \text{or} \quad K_T = \rho \left(\frac{\partial \rho}{\partial p} \right)_T \quad (3.15)$$

Where v is molar volume, p is the pressure and ρ is the density

In our study, the effect of the impurities on the isothermal compressibility of CO₂ and CO₂/impurities are experimentally estimated by calculating the effects on the derivative of the displacement with respect to the pressure (Under the assumption that the effective cross-section area is constant in all the tests).

$$K_T = -\frac{1}{z} \left(\frac{\partial z}{\partial p} \right) \quad (3.16)$$

where z is the piston displacement

To calculate the isothermal compressibility (K_T), a cubic-spline method was used to smooth the measured displacement data, and K_T was obtained by differentiate the displacement z data with respect to pressure. It was estimated that the uncertainty of the K_T data was better than 2% (The value is obtained by calculating the error of experimental data to the data obtained from spline interpolation). The isothermal compressibility results of pure and impure CO₂ systems are listed in [Table 3.2](#). [Figure 3.8](#) shows the isothermal compressibility of pure CO₂ at 268.15, 278.15 and 293.15 K. As can be seen from the figure, K_T is sensitive to the pressure as the pressure approaches the critical point of CO₂. On the other hand, K_T is non-sensitive to the pressure when the pressure is far enough from the saturation line. The figure also illustrates that K_T is temperature dependent; its changes are limited as the temperature moves down away from the critical temperature of CO₂ even near the saturation line.

The proposed model in this work is tested to predict the measured isothermal compressibility of pure CO₂. The model predictions are in a good agreement with our experimental data and the average deviation was less than 3% which is comparable to the uncertainty estimated in the experimental measurements. The model was also used to predict the isothermal compressibility of CO₂/impurities systems. An example is given in Figure 3.9. This figure shows the isothermal compressibility of MIXa (CO₂+Ar+CO) at different temperatures. The results demonstrate the capability of the model to predict the compressibility of CO₂ mixtures. The calculated average deviation in these mixtures was reported less than 4%. The PR-model with the proposed shifting volume was not tested for higher impurity concentration in CO₂ mixture.

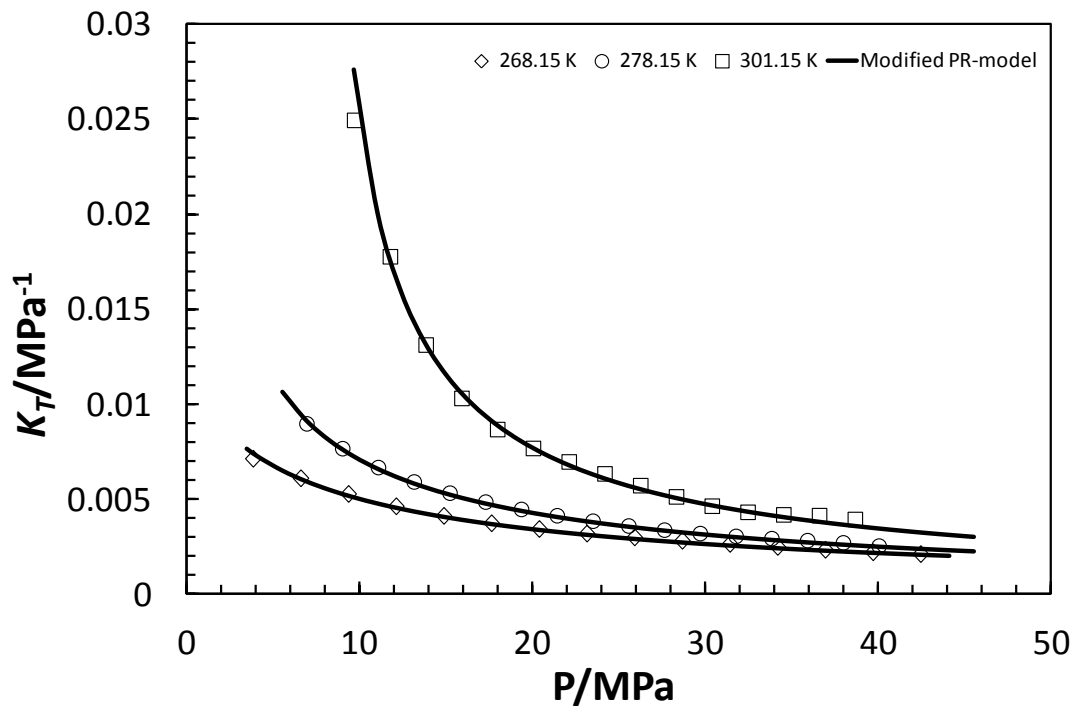


Figure 3.8 Isothermal compressibility in pure CO₂ (measurements & modelling)

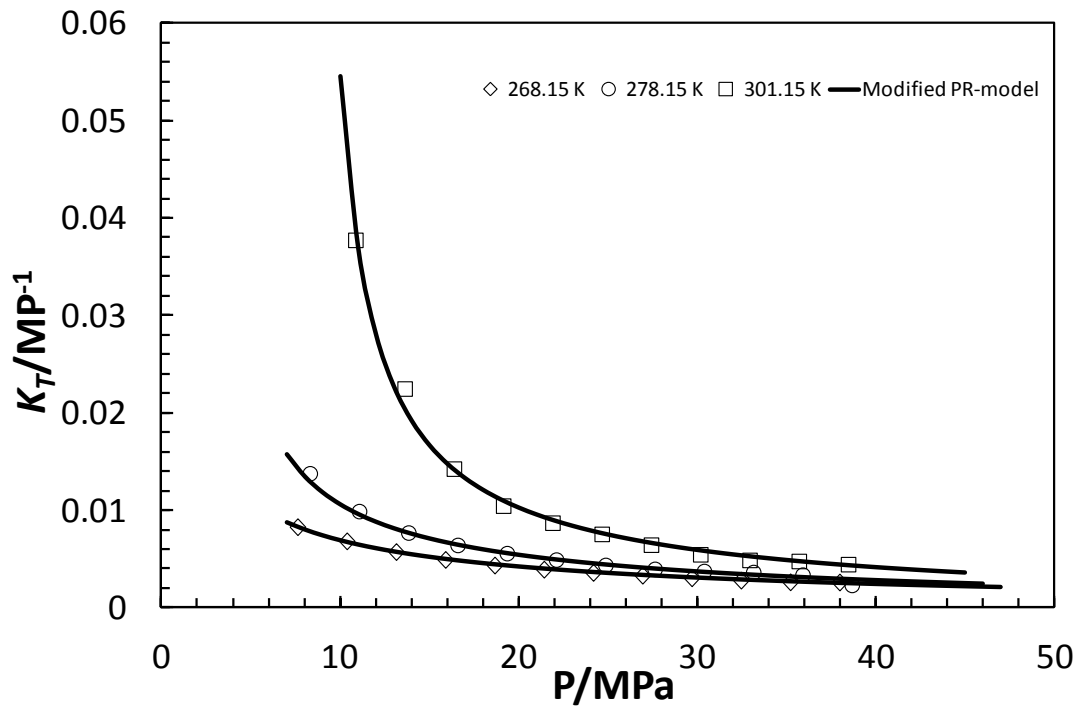


Figure 3.9 Isothermal compressibility in MIXb ($\text{CO}_2+\text{Ar}+\text{CO}$) (measurements & modelling)

Figure 3.10 and Figure 3.11 show the effect of three impurities systems at two different temperatures (268.15 and 301.15 K). In these systems, it is obvious that the presence of the impurities in CO_2 tends to increase the compressibility. The effect is higher at higher temperature and it is very significant as the pressure approaches the saturation line of the mixtures. This variation of K_T with pressure and temperature of the mixtures are similar for all the investigated systems. In general, the effect in the compressibility also depends mainly on two factors: the concentration of the impurity in CO_2 mixture and the type of impurities present in the mixture.

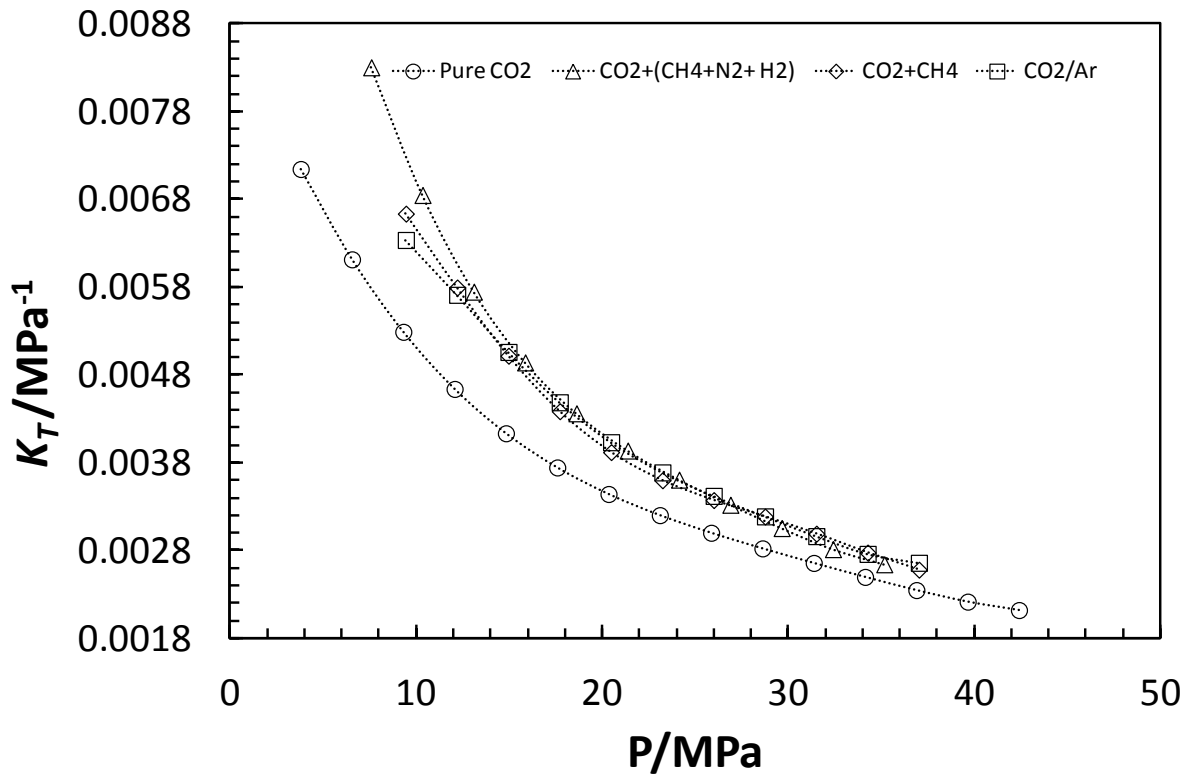


Figure 3.10 Effect of three selected impurities systems on CO₂ stream at 268.15 K

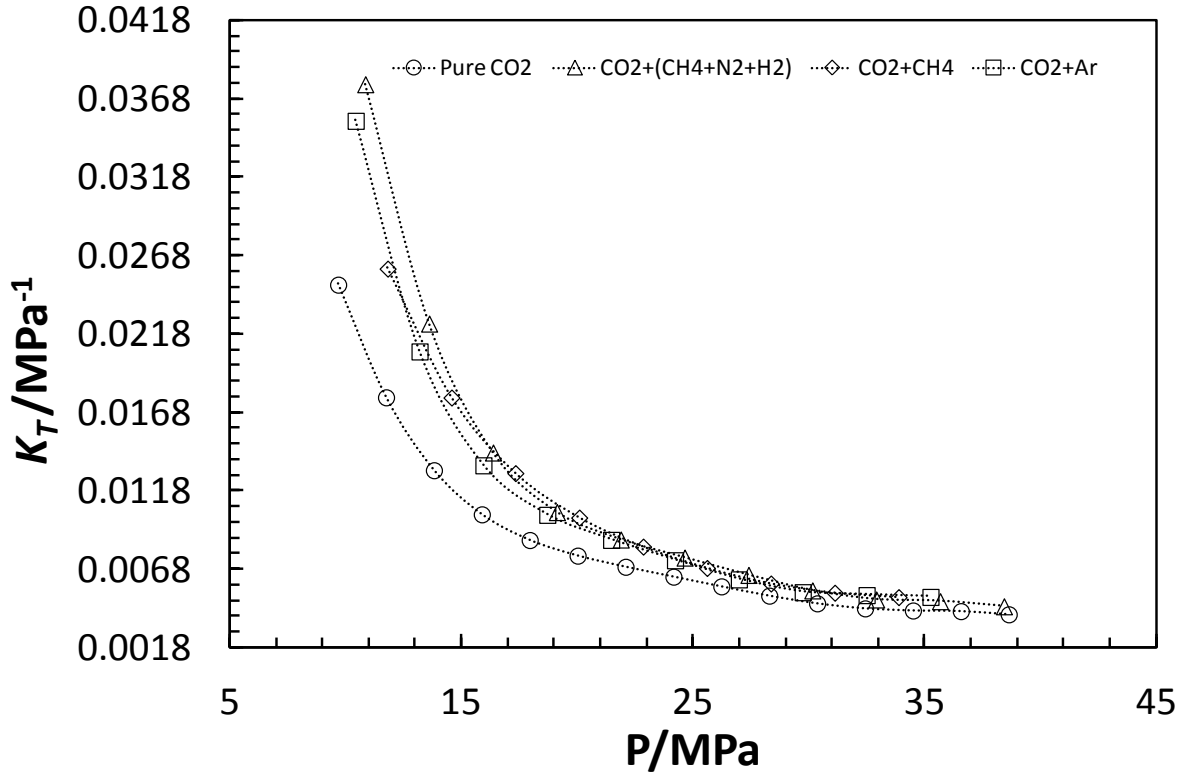


Figure 3.11 Effect of three selected impurities systems on CO₂ stream at 301.15 K

3.5 Conclusions

The presence of impurities tested in this work in CO₂ stream tends to reduce the speed of sound. The reduction depends more on pressure than temperature. The plot of ($SoS_{pureCO_2} - SoS_{impureCO_2}$) versus pressure at different temperatures shows that this change is high at high temperature for a pressure near the saturation line. As the pressure increases, the isothermal lines of the change will intersect at a pressure ≈ 15 MPa. This finding was observed in all the investigated mixtures. Below that pressure, the change is higher at lower temperatures. Conversely, above that pressure the change is higher at higher temperatures. In the other hand, the presence of impurities increases the isothermal compressibility of CO₂ stream. The effect is much higher as the pressure approaches the critical temperature of the mixture.

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CHAPTER 4: DENSITY OF CARBON DIOXIDE RICH SYSTEMS

4.1 Introduction

The density is important when calculating the flow in the pipeline. Knowledge of the inlet temperature, ambient temperature and heat transfer is needed to achieve a correct model of the flow. Maintaining a single phase and avoiding the abrupt pressure drops in CO₂ pipelines are an important issue for engineers where pipelines required boosting stations. The density also (as we can see later in the chapter) is very sensitive near the critical points, .i.e. a small change in both pressure or temperature results in a very large change in the density of CO₂ [1]. Properties of CO₂ stream are the temperature, pressure and density profile along the pipelines therefore, the effect of these parameters translates into recompression distances and compressor power requirements, which in turn would have cost implications [2].

PVT_{xy} for CO₂/impurities system cited in the literature were collected recently by Li et al. [3]. Their conclusion was that the data for CO₂ in the presence of H₂O, CH₄, N₂ and H₂S are available in a wide range of pressure and temperature. However, the other impurities associated with the pipelines transport are still limited. Beside the bibliographic references provided by Li et al. [3], Jiang et al. [4], provided volumetric properties for CO₂/N₂ binary system at 293 K and pressure from 0.60 to 5.18 MPa. Esper et al. [5] measured the density of the binary CO₂/CH₄ system at composition 0.48 CO₂, pressure from 0.08 to 48.3 MPa and temperature range 206 to 320 K. Seitz and Blencoe [6] extended their previous experimental work with the binary CO₂/CH₄ system [7] by covering temperature and pressure up to 673 K and 100 MPa respectively. Hwang et al. [8] reported experimental density measurements for CO₂ + CH₄ mixtures at temperatures from 225 K to 350 K and pressures up to 35 MPa (one isotherm up to 69 MPa). Additional PVT data for the binary mixture have been reported by Nederlandse Gasunie, and Ruhrgas as cited by Jaeschke and Humphreys [9]. Three references were found for the densities of the ternary CO₂/N₂/CH₄ system ([10], [11] and [7]). Bezanehtak et al. [12] contacted VLE and liquid density measurements for CO₂/H₂ binary system at temperatures from 278.15 K to 308.15 K and for pressures in the range 1.5 to 19.3 MPa. The authors compared their VLE results with those obtained

by Tsang and Streett [13], and the deviations for liquid and vapour phases were as high as 7.1% and 9.4%, respectively. Three PVT measurements in CO₂/Ar binary mixture were reported until 1977. Two isothermal data found in literature, at vapour and liquid regions, were measured by Altunin et al. [14] and Sarashina et al. [15] respectively. Kestin et al. [16] (cited from [3].) investigated the volumetric property for the binary mixture ranging from 293.15-303.15 K and pressure up to 2.58 MPa. A single PVT data was reported in the literature [17] for CO₂/CO mixture at a constant composition (0.5732 CO₂ mole%). The measurements were made in the temperature range 323 to 423 K and pressures up to 6.5 MPa. Recently, new supercritical volumetric measurements on three binary mixtures (CO₂/N₂, CO₂/O₂ and CO₂/Ar) were measured by Mantovani et al. [18]. Two different fractions in impurity mixture were prepared for each binary system ranging from 4.15 mole% to 16.94 mole%. The measurements were performed isothermally crossing vapour/supercritical boundary region, temperature range from 303 K to 383 K and in a pressure range from 1 MPa to 20 MPa.

The aim of this work is to investigate the effects of impurities on the CO₂ transport density. Six types of binary CO₂ rich systems were prepared plus a multi-components CO₂-rich mixture. The liquid CO₂ density measurements were carried-out from 283.15 K to 301.15 K and at a pressure up to 50 MPa. In addition, the densities of four binary mixtures of CO₂/Ar, CO₂/CO, CO₂/H₂ and CO₂/O₂, at constant fraction of impurities, were measured in the supercritical region. The measurements were carried-out isothermally at three temperatures, 323.15 K, 373.15 K and 423.15 K and at pressure up to 50 MPa.

4.2 Experimental Methods and Equipment Materials

The following compounds were used to make the different synthetic mixtures studied in this work:

- Carbon Dioxide, Supplied by Air Products, Research Grade
- Methane, Supplied by Air Products, Grade N4.5
- Nitrogen, Supplied by BOC, Research Grade
- Argon, Supplied by BOC, Research Grade
- Carbon Monoxide, Supplied by BOC, Research Grade
- Oxygen, Supplied by BOC, Research Grade

- Hydrogen, Supplied by BOC Gases, \geq %99.995

The systems listed in [Table 4.1](#) were prepared gravimetrically and used to conduct density tests. A multi component mixture, MIX1, (composition given in [Table 4.2](#)) supplied by BOC was also used to carry out the tests.

Table 4.1 The Composition of mixtures used for the experiments

Liquid above the saturation line		
Component	CO ₂ % mole (± 0.3)	Impurity %mole (± 0.3)
Pure CO ₂	100	0.0
CO ₂ +CH ₄	94.1	5.9
CO ₂ +N ₂	95.4	4.6
CO ₂ +Ar	95.0	5.0
CO ₂ +CO	94.1	5.9
CO ₂ +O ₂	95.0	5.0
CO ₂ +H ₂	95.0	5.0
CO ₂ +H ₂	97.4	2.6
MIX1	95.6	See Table 4.2
Supercritical above vapour/supercritical boundary line		
Pure CO ₂	100	0.0
CO ₂ +Ar	95	5.0
CO ₂ +CO	95	5.0
CO ₂ +O ₂	95	5.0
CO ₂ +H ₂	95	5.0
MIX1	95.6	See Table 4.2

Table 4.2 Mixture composition (MIX1)

Mixture	% mole (± 0.05)
Carbon Dioxide	Balance
Methane	0.67
Hydrogen	0.82
Nitrogen	1.41
Carbon Monoxide	0.21
Argon	1.21
Oxygen	0.08
Total	100

4.2.1 Experimental Equipment

The schematic diagram of the experimental system used in this work is shown in [Figure 4.1](#). The compressed liquid and supercritical density data of CO₂ and CO₂/impurities systems were measured with a high temperature and pressure vibrating tube densitometer Anton Paar DMA-HPM which consists of a measuring cell and an interface module. The measuring cell contains a U-shaped Hastelloy C-276 tube that is being excited to vibrate at its characteristic frequency electronically and the interface module generates the period of oscillation and measures the period of oscillation. The DMA-HPM has been connected to a mPDS 2000V3 evaluation unit which indicates the vibration period with seven significant digits.

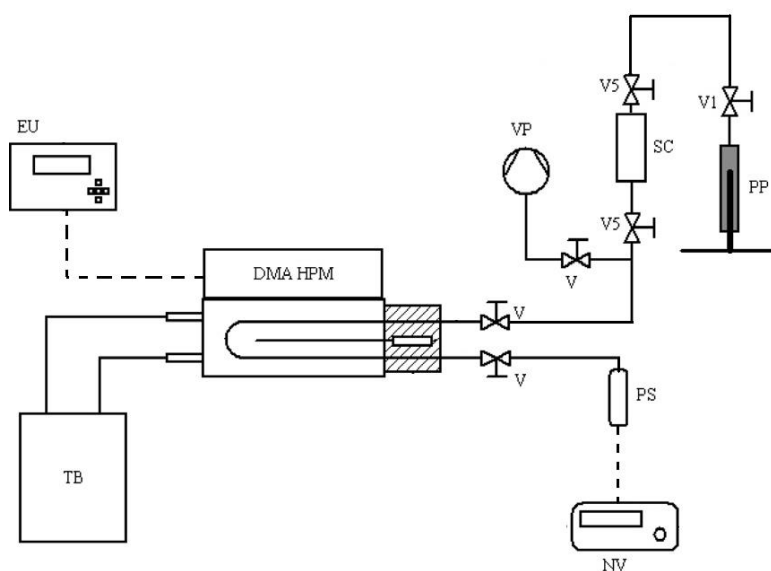


Figure 4.1 Schematic diagram of the compressed liquid densimeter apparatus. PP, Ruska motorized HP pump; PS, pressure sensor; SC, sample cylinder; VP, vacuum pump system; TB, thermostatic bath; EU, mPDS 2000V3 unit; NV, nanovolt meter; V, valves

The vibrating tube is thermally controlled by a circulating coolant coming from an external thermostatic bath for liquid CO₂ mixture samples and by an oven for supercritical CO₂ samples. The measuring cell was insulated from the external environment and a maximum temperature fluctuation of ± 0.01 K could be obtained all along the measurements. The temperature of the vibrating tube cell was measured by a built-in thermometer connected to the mPDS 2000V3 unit. The pressure of the system was applied with a Ruska motorised High Pressure Pump and measured by a pressure transducer (max 70 MPa); the pressure was kept constant to show significant digits in density do not change during the time of the measurement.

4.2.2 *Experimental Calibration and Procedure*

After the entire circuit was purged, vacuum was applied to the system. Then the sample was introduced into the circuit by corresponding valves operation. When the temperature of the vibrating tube was stable, the vibration period of the U-tube was determined from the appropriate initial pressure to the maximum pressure. Then the temperature of the bath was changed and a new isotherm was repeated.

The measurement of density with a vibrating tube densitometer is not absolute, thus, the raw data (period of oscillation) should be further treated to obtain the densities. The relationship between them is [19]

$$\rho(T, P) = A(T, P)\tau^2(T, P) - B(T, P) \quad (4.1)$$

Where $\rho(T, P)$ is the sample density at temperature T and pressure P , $\tau(T, P)$ is the period of oscillation at temperature and pressure, $A(T, P)$ and $B(T, P)$ are the apparatus parameters depending on temperature and pressure, and they must be determined from calibration measurements. In our calibration, CO₂ density is used as a reference substance at two different pressures (at a point just above the saturation point and a point at higher pressure both should be measured at the same temperature). The apparatus parameters were defined as follow:

$$A(T, P_1) = \frac{\rho(T, P_1) - \rho(T, P_2)}{\tau^2(T, P_1) - \tau^2(T, P_2)} \quad (4.2)$$

$$B(T, P_2) = \frac{\tau^2(T, P_2) \cdot \rho(T, P_1) - \tau^2(T, P_1) \cdot \rho(T, P_2)}{\tau^2(T, P_1) - \tau^2(T, P_2)} \quad (4.3)$$

4.3 Modelling

Many studies have been done to compare the abilities of various equations of state. Danesh et al. [20] compared 10 equations of state with classical mixing rules for predicting the phase behaviour and volumetric properties of hydrocarbon fluids. They concluded that the Valderrama modification of the Patel and Teja cubic equation of state [21] was superior to the other tested equations of state, particularly when the EOSs were compared without any use of binary interaction parameters. The other important equation of state is Peng Robinson EOS [22], which is classified as a two-parameter cubic equation and has become one of the models of choice in the process modelling. Span and Wagner [23] have the latest review of the thermodynamic properties of CO₂. Their work covered a wide range of pressure and temperature including the critical area based on large experimental data. They generated highly accurate tables based on a complex equation of state. Therefore, the tabulation data of their work is used to evaluate the absolute average deviation (ADD) for both PR and VPT EOS models. The deviations of the models from the experimental data were drawn for both liquid and supercritical CO₂ from 230 to 390 K and up to 50 MPa as shown in [Figure 4.2](#) and [Figure 4.3](#). The ADDs are 4.4 and 2.7% for PR and VPT EOSs respectively in liquid CO₂. In the supercritical region, the ADD is around 2.52% in both models. From the above figures, the maximum deviation can be as high as 10% in both cases. Hence, the use of those models without further modification for CO₂/impurities systems can definitely increase the deviation even with proper interaction parameters tuning.

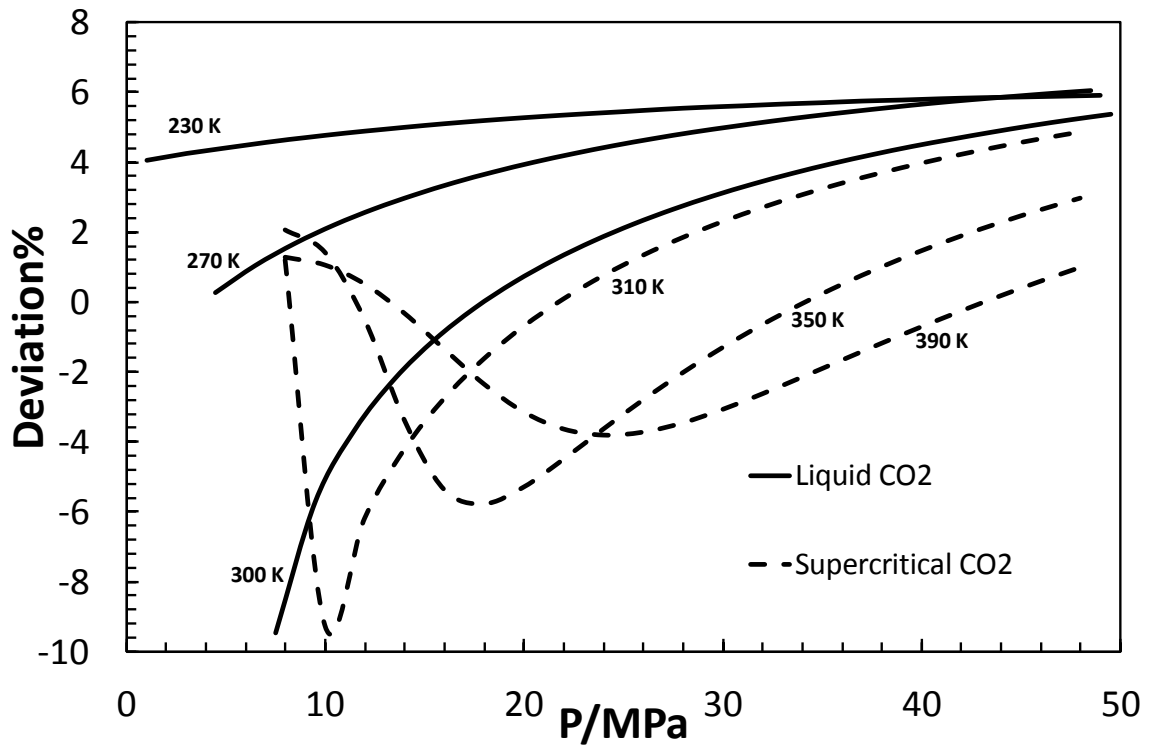


Figure 4.2 Deviation between PR EOS and the data provided in [23] for pure CO₂

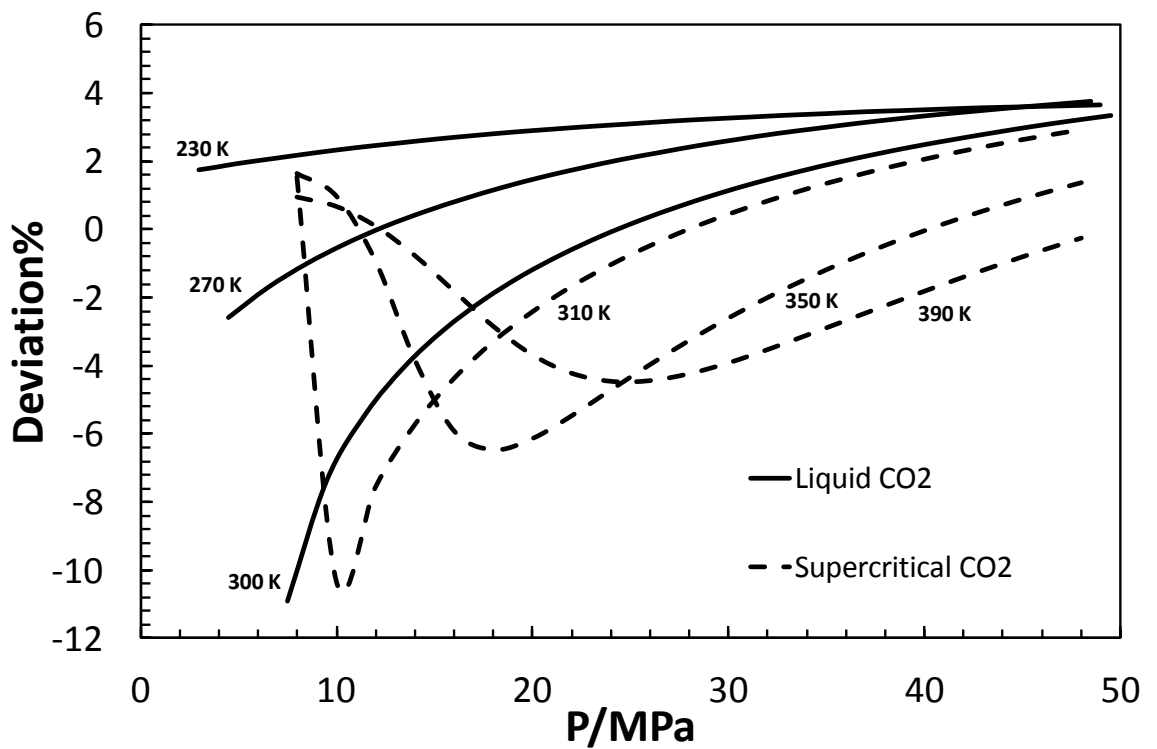


Figure 4.3 Deviation between VPT EOS and the data provided in [23] for pure CO₂

Table 4.3 The interaction parameters of CO₂/binary system used in VPT EOS

Mixture	k_{ij}
CO ₂ /Ar	0.113
CO ₂ /O ₂	0.0909
CO ₂ /H ₂	-0.386
CO ₂ /CO	-0.0415
CO ₂ /N ₂	-0.0607
CO ₂ /CH ₄	0.0915

In this work, the volume for CO₂ or a CO₂-rich mixture calculated by the EOS is corrected using the exact volume of CO₂ at the given T and P.

$$V^{new} = V^{EoS} - V^c \quad (4.4)$$

Where V^{EoS} is the molar volume obtained from VPT EOS, the binary interaction parameters, K_{ij} associated with the equation are tuned with independent experimental data and their values are given in

Table 4.3. The molar volume correction in the Equation 4.4, V^c , is defined as

$$V^c = \sum_i^{NComp} x_i V_i^c \quad (4.5)$$

x_i is the composition of component i in the phase in which the volume is calculated. The volume calculated for impurities V_i^c is set equal to 0 when there is no volume translation is used. For CO₂, V_i^c is defined by

$$V^c_{CO_2} = V^{EoS}_{pureCO_2} - V^{MBWR} \quad (4.6)$$

The carbon dioxide density is computed from the MBWR equation in the form suggested by McCarty [24]

$$P = \sum_{n=1}^9 a_n(T) \rho^n + \sum_{n=10}^{15} a_n(T) \rho^{2n-17} e^{-\gamma \rho^2} \quad (4.7)$$

4.4 Results & Discussion

4.4.1 Density of Liquid CO₂/Impurities System

Liquid densities of CO₂ and six binary CO₂/impurities systems (see [Table 4.1](#)) were measured from (283.15 to 301.15) K, at pressures up to 50 MPa. Measurements were also conducted on a multi-components CO₂-rich mixture, MIX1, shown in [Table 4.2](#). A total of 382 points were obtained in this test, as listed in [Table 4.4](#). Our modified model prediction and its deviation with the experimental data are calculated and given in the table. The AAD of the modified and original models in pure CO₂ are 0.29% and 3.4% while the maximum deviation is 0.87% and 11.4% respectively. The modified model is found to be in a good agreement with our experimental data of impure CO₂ except near the critical point of the mixture. The actual critical temperature of the mixture tends to decrease in the presence of impurities (HydraFlash software). For example, at a composition fraction of 4.78% methane, the critical temperature of CO₂ is reduced from 304.2 K to 300.8 K. Therefore, the tested temperature, 301.15 K is even closer to the critical temperature of the mixture than of pure CO₂. [Figure 4.4](#) and [Figure 4.5](#) show the predictions of the proposed model correction at 283.15 and 301.15 K, respectively. The overall AAD for the tested impurities systems (excluding the binary CO₂/H₂ system) is 0.5%. The modified model deviates up to 5.57% and with AAD of 3.86% in CO₂/H₂ system below the critical temperature of the mixture. This is mainly due to the low molecular weight and critical temperature of H₂ comparing to the other impurities tested. From [Table 4.4](#) some interesting observations can be made. The density of pure CO₂ and the CO₂/impurities mixtures have a stronger dependency on temperature compared to pressure at lower temperatures below their critical temperature. At condition very close to the critical point, a small change in temperature or pressure yields a very large change in the density of CO₂. In transport, this change could lead to a sudden change of phase and fluid velocity.

Table 4.4 Experimental and modelling data for liquid CO₂ and liquid CO₂/impurities system

Pure CO ₂					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
42.95	1059.05	0.0	42.2	1041.95	0.02
39.49	1049.65	0.0	39.3	1033.55	0.02
36.03	1039.65	0.0	36.4	1024.65	0.02
32.58	1028.85	0.0	33.51	1015.15	0.02
29.12	1017.15	0.0	30.61	1004.85	0.03
25.66	1004.35	0.0	27.72	993.88	0.04
22.2	990.14	0.0	24.82	981.87	0.04
18.74	974.19	0.0	21.93	968.66	0.05
15.29	955.88	0.0	19.03	953.9	0.05
11.83	934.06	0.0	16.13	937.11	0.06
8.37	906.63	0.0	13.24	917.46	0.05
			10.34	893.41	0.03
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
44.41	1033.25	0.0	44.82	1010.45	0.04
40.89	1022.75	0.0	41.18	998.73	0.04
37.38	1011.55	0.0	37.54	986.01	0.05
33.86	999.43	0.0	33.9	972.08	0.06
30.34	986.14	0.1	30.25	956.63	0.07
26.83	971.42	0.1	26.61	939.24	0.07
23.31	954.89	0.1	22.97	919.22	0.08
19.8	935.9	0.1	19.33	895.45	0.08
16.28	913.4	0.1	15.69	865.75	0.08
12.76	885.34	0.1	12.05	824.86	0.08
9.25	846.64	0.0	8.41	751.38	0.09
			7.27	698.07	0.13
CO ₂ /CH ₄					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
45.11	1015.73	0.1	44.55	999.77	0.13
41.23	1004.89	0.1	41.51	990.84	0.1
38.1	995.5	0.1	37.91	979.66	0.06
34.38	983.55	0.1	34.53	968.06	0.04
29.57	966.41	0.0	30.57	953.22	0.01

25.21	948.75	0.1	27.37	940.11	0.09
21.89	933.54	0.1	24.04	924.43	0.14
18.28	914.54	0.3	20.59	906.16	0.24
14.81	892.78	0.4	16.97	883.4	0.41
11.38	865.93	0.7	13.98	860.38	0.63
8.91	840.66	1.1	11.91	840.82	0.84
6.64	807.35	1.6	9.32	809.88	1.42
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
41.6	975.03	0.2	32.95	917.08	0.06
38.29	963.93	0.2	44.31	960.81	0.11
34.31	949.56	0.1	40.8	948.59	0.11
31.07	936.55	0.1	37.83	937.49	0.09
27.68	921.47	0.1	34.68	924.67	0.07
23.99	902.91	0.0	29.31	899.79	0.01
20.49	882.37	0.2	23.36	865.75	0.14
17.3	859.85	0.3	18.8	831.41	0.33
13.82	829.5	0.7	14.53	786.72	0.71
11.76	806.12	1.1	12.09	748.58	1.08
9.43	768.44	1.6	10.49	711.91	1.51
8.52	747.64	2.0	9.01	652.38	1.98
7.68	722.4	2.5	8.44	607.33	1.71
CO ₂ /N ₂					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
37.93	1007.12	0.1	44.5	1012.84	0.1
32.47	988.25	0.1	44.05	1011.48	1.6
29.65	977.46	0.1	39.69	997.67	0.1
25.27	959	0.0	36.34	986.1	0.1
21.78	941.49	0.1	33.21	974.42	0.0
17.68	918.1	0.2	28.17	953.48	0.0
14.73	898.61	0.5	24.43	935.32	0.1
14.56	895.9	0.4	20.93	915.93	0.2
10.74	861.27	0.7	17.44	892.88	0.4
7.12	810.01	1.4	13.78	862.37	0.7
			9.96	810.54	0.5
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
44.36	996.98	0.1	45.17	976.41	0.0
40.98	985.87	0.1	41.04	961.61	0.0
38.27	976.3	0.1	37.91	949.35	0.0
34.72	962.81	0.1	34.78	936.03	0.0
33.72	959	0.0	31.15	918.91	0.0

30.95	947.25	0.0	26.21	892.55	0.2
27.72	932.09	0.1	22.22	865.74	0.4
24.28	913.9	0.2	18.34	832.64	0.5
21.02	893.9	0.3	15.09	793.4	0.5
17.42	867.25	0.5	12.14	744.13	1.1
14.57	840.14	0.7	10.34	691.48	1.2
12.25	812.02	1.0			
9.98	772.78	1.4			
CO ₂ /Ar					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
44.33	1051.57	0.0	44.45	1036.96	0.0
41.22	1042.52	0.0	40.92	1026	0.0
38.27	1033.37	0.0	37.74	1015.44	0.0
34.08	1019.35	0.0	33.45	999.6	0.0
28.83	999.63	0.1	28.72	980.73	0.1
24.92	982.98	0.2	24.67	961.69	0.2
21.68	967.32	0.2	20.56	939.53	0.3
17.67	944.95	0.4	17.84	921.84	0.4
14.28	922.11	0.6	14.26	894.11	0.7
10.91	893.38	0.9	10.93	859.96	1.1
7.8	856.93	1.5	8.76	828.81	1.6
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
44.34	1021.1	0.0	48.23	1008.44	0.1
40.62	1008.73	0.0	44.27	995.27	0.1
37.89	999.48	0.0	41.04	983.69	0.1
34.01	984.51	0.1	38.24	972.89	0.1
30.92	971.85	0.1	34.32	956.39	0.0
27.33	955.74	0.3	30.99	940.74	0.0
24.71	941.4	0.2	27.37	921.65	0.1
20.43	915.83	0.4	24.1	901.91	0.2
17.32	893.71	0.7	20.6	876.91	0.4
14.01	864.12	1.1	17.44	848.92	0.6
11.79	838.03	1.5	14.95	820.79	0.9
9.56	802.04	2.2	12.14	777.46	1.5
			9.7	709.31	2.5
CO ₂ /CO					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
47.04	1039.77	0.0	47.86	1027.62	0.0
42.94	1028.4	0.0	44.58	1018.26	0.0
38.14	1013.94	0.0	41.24	1008.15	0.1

33.36	997.9	0.1	38.14	998.12	0.1
28.65	980.27	0.2	34.21	984.13	0.1
24.54	962.69	0.2	30.81	970.99	0.2
20.66	943.7	0.3	27.36	956.24	0.2
17.04	922.87	0.5	23.63	938.51	0.3
13.72	899.91	0.7	20.44	920.47	0.4
10.83	874.68	1.0	17.23	899.7	0.6
8.74	851.51	1.4	13.57	870.16	0.9
			10.63	838.42	1.3
			9.56	825.54	1.7
			9.6	825.23	1.6
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
46.8	1009	0.1	48.55	992.07	0.1
42.49	995.59	0.0	44.56	979.01	0.1
38.12	980.18	0.1	40.76	965.5	0.1
33.74	964.13	0.0	40.75	965.4	0.1
29.19	944.65	0.1	37.97	954.72	0.1
24.23	919.6	0.2	34.21	938.95	0.1
20.91	899.76	0.3	30.63	922.08	0.2
17.21	873.13	0.6	27.35	904.68	0.2
13.71	840.4	0.9	24.17	885.41	0.3
10.7	800.61	1.5	20.49	858.88	0.5
9.27	773.67	2.0	17.35	830.68	0.7
			13.75	786.38	1.2
			10.67	723	1.9
CO ₂ /O ₂					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
9.4	860.57	1.0	9.34	820.34	1.4
16.27	921.14	0.2	13.67	873.35	0.6
19.27	939.75	0.1	18.6	912.54	0.2
21.55	952.05	0.0	20.77	926.35	0.1
25.15	969.81	0.1	24.52	947.14	0
27.56	980.64	0.1	26.59	957.35	0.1
29.7	989.35	0.2	29.17	969.1	0.1
31.93	997.9	0.2	32.87	984.36	0.1
34.79	1008.26	0.2	35.36	993.76	0.2
38.29	1020.02	0.2	39.05	1006.66	0.2
42.53	1032.83	0.3	41.67	1015.22	0.2
46.85	1045.16	0.3	46.05	1028.49	0.3
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%

8.93	761.88	1.9	9.34	623.5	3.9
12.85	832.34	0.9	13.73	770.65	1.1
17.57	879.81	0.4	18.46	834.87	0.7
21.67	909.48	0.2	22.59	871.08	0.7
25.4	931.49	0.0	25.59	896.57	0.2
28.66	948.12	0.0	30.04	915.82	0.9
32.33	964.76	0.1	33.54	934.97	0.7
36.17	980.31	0.1	36.77	952.29	0.4
39.84	993.76	0.1	40.2	967.15	0.3
43.75	1006.91	0.2	44.3	981.52	0.4
47.2	1017.63	0.2	47.79	993.19	0.4
CO ₂ /H ₂ (2.6%)					
T = 283.15 K			T = 288.15 K		
P/MPa	ρ/Kg.m ⁻³	Dev%	P/MPa	ρ/Kg.m ⁻³	Dev%
19.87	930.54	0.6	-	-	-
17.93	919	0.7	-	-	-
16.55	910.05	0.7	-	-	-
15	899.04	0.7	-	-	-
13.79	890.01	0.7	-	-	-
11.88	872.66	0.8	-	-	-
11.03	864.69	0.8	-	-	-
10.34	857.07	0.8	-	-	-
9.65	848.52	0.9	-	-	-
8.96	839.44	0.9	-	-	-
8.83	837.23	1.0	-	-	-
T = 293.15 K			T = 301.15 K		
P/MPa	ρ/Kg.m ⁻³	Dev%	P/MPa	ρ/Kg.m ⁻³	Dev%
17.76	864.65	1.2	17.93	821.2	1.5
16.55	854.11	1.2	16.55	806.24	1.6
15	838.75	1.3	15	786.8	1.7
13.79	825.7	1.4	13.79	768.45	1.9
11.98	800.47	1.7	12.07	734.19	2.4
11.03	785.71	1.8	11.03	706.16	3.1
9.65	754.47	2.7	9.65	649.44	5
9.31	745.49	2.9	9.31	623.95	6.7
9.1	738.84	3.1	9.17	623.54	5.6
9.03	737.26	3.1			
CO ₂ /H ₂ (5%)					
T = 283.15 K			T = 288.15 K		
P/MPa	ρ/Kg.m ⁻³	Dev%	P/MPa	ρ/Kg.m ⁻³	Dev%
11.18	856.75	4.7	12.67	840.12	4.7
15.45	894.91	3.7	14.72	858.96	3.9

17.02	906.04	3.5	18.51	888.71	3.1
19.75	922.93	3.1	20.77	903.75	2.9
23.37	941.91	2.7	23.86	921.63	2.6
26.88	958.09	2.5	27.14	937.81	2.4
30.37	972.68	2.4	31.54	956.5	2.1
34.95	989.58	2.2	36.38	975.18	1.9
38.9	1000.5	1.9	40.4	991.03	2
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
10.27	761.44	6.6	11.26	740.19	12
16.35	845.21	3.8	13.1	776.35	9.2
16.38	845.44	3.8	15.89	814.62	7.1
21.22	899.62	4.7	18.49	842.86	6.3
27.5	934.14	3.9	22.66	874.23	5
35.07	966.59	3.4	25.28	888.7	4.3
42.97	994.11	3.0	27.6	900.8	3.9
49.24	1013.13	2.8	31.78	920.98	3.4
			36.92	945.2	3.2
			40.22	956.85	3
			42.3	964.16	2.8
MIX1					
T = 283.15 K			T = 288.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
10.09	858	1.7	9.94	830.48	3.5
14.56	899.5	1.1	11.49	852.47	3.2
18.99	929.3	0.8	14.78	885.22	2.5
21.55	944.55	0.7	21.27	931.72	1.9
25.42	963.55	0.6	27.32	964.77	1.8
27.8	972.83	0.4	33.81	990.75	1.5
33.04	994.54	0.4	43.86	1026.12	1.5
37.83	1009.76	0.2			
41.89	1022.55	0.2			
T = 293.15 K			T = 301.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
41.99	988.96	0.0	11.79	734.75	1.9
37.31	972.51	0.0	13.42	767.83	1.4
33.32	956.91	0.1	15.93	805.27	1.1
27.35	929.73	0.2	20.5	851.65	0.7
23.9	911.16	0.3	24.46	881.58	0.6
19.75	884.31	0.5	29.31	910.8	0.5
16.64	859.67	0.8	34.25	935.48	0.4
13.32	825.01	1.2	37.92	951.41	0.4
11.12	793.31	1.7			

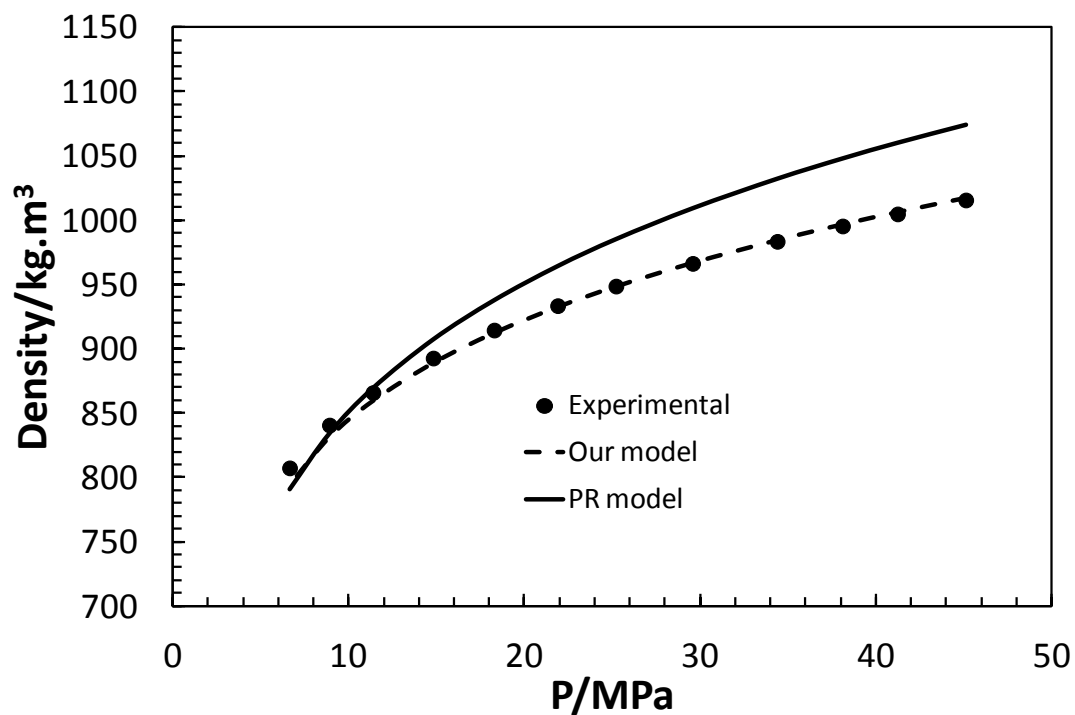


Figure 4.4 Density of CO₂/CH₄ at 283.15 K

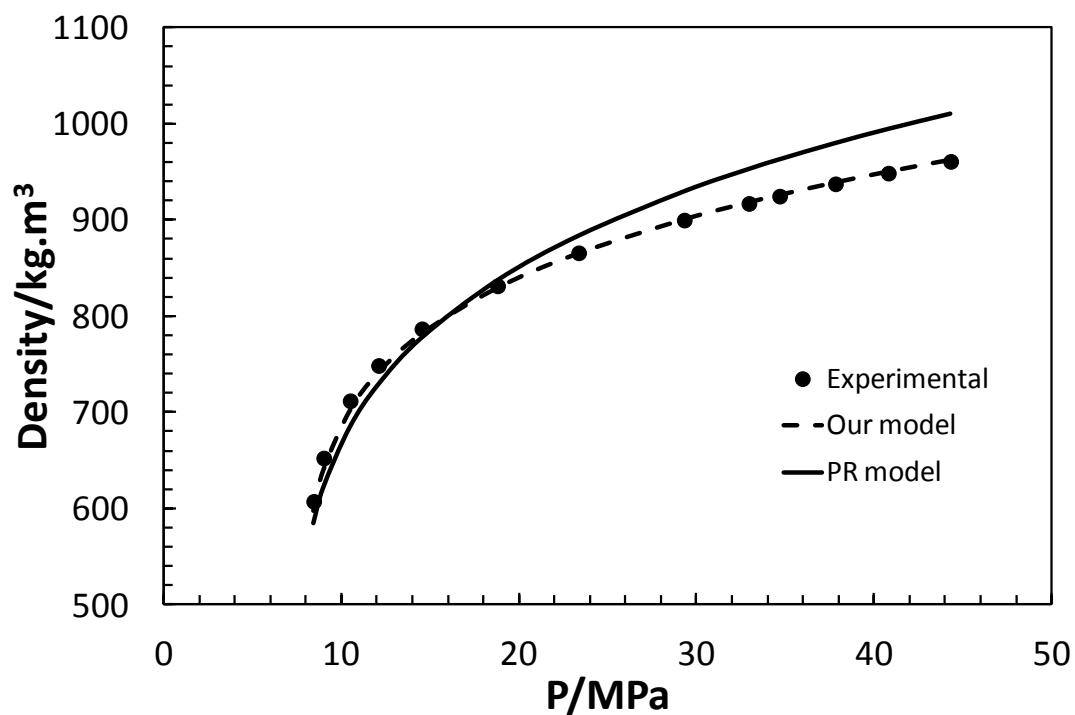


Figure 4.5 Density of CO₂/CH₄ at 301.15 K

Our modified model was used to evaluate the change in CO₂ density in the presence of impurities. An example of the change on the density when impurities are present can be seen in Figure 4.6. The change is plotted as a function of pressure at 283.15 K for four binary mixtures (CO₂/Ar, CO₂/O₂, CO₂/H₂ and CO₂/CH₄). $\Delta\rho$ is calculated by subtracting the density of the corresponding CO₂/impurities from the density of pure CO₂. In general, the impurities reduce CO₂ density over the range of pressure and temperature tested. The change is found to be higher near the critical points of the mixture. The reduction is higher in the presence of methane and it has a minimal effect in the presence of Argon. This is mainly due to the molecular weight ratio of impurity in the mixture to that of pure CO₂, i.e. as $\left(\frac{MW_{CO_2}}{MW_{impurity}}\right)$ tends to 1, the density change becomes insignificant. This rule is obvious in all the tested mixture except for hydrogen binary mixture where the effect of smaller molecular weight is also significant. The reduced density is largely related to the increase volume of the mixture. This can be understood from the fact that non-condensable impurities are less dense than CO₂ and hence take greater volumes.

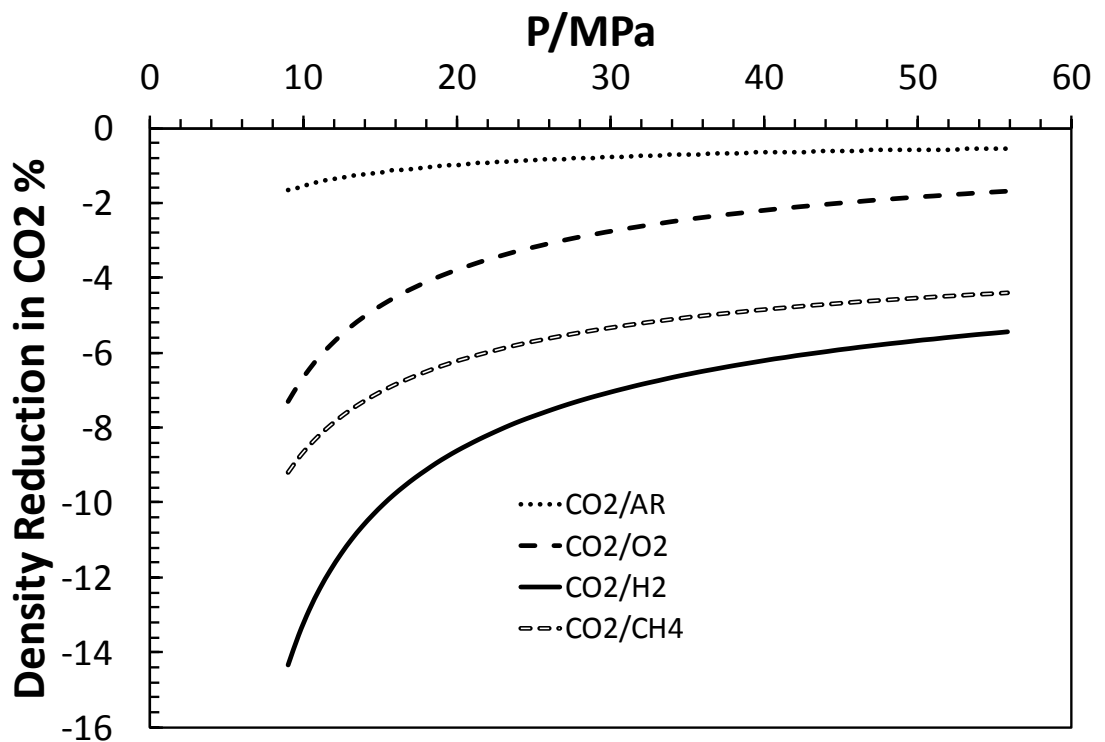


Figure 4.6 Density reduction in liquid CO₂ in the presence of impurities at 283.15 K

4.4.2 Density of Supercritical CO₂/Impurities System

In this part, five systems of CO₂/Ar, CO₂/CO, CO₂/H₂, CO₂/O₂ and MIX1 (composition given in Table 4.1 and Table 4.2) were measured at three different temperatures, 323.15 K, 373.15 K and 423.15 K. The experiment was carried-out at a pressure above the critical pressure of the mixture and up to 50 MPa. The results are listed in Table 4.5 with 160 measurement points. In supercritical region, the density of pure CO₂ is seen to be highly dependent on pressure compare to temperature. Near the critical pressure, a small change in pressure can have a large effect on the density. CO₂ fluids with high density will reduce the pipeline size and the running cost. This is not the case in the presence of impurities. In order to understand this effect, spline interpolation is implemented to the experimental data from 8.3 MPa to 41.3 MPa. Figure 4.7 shows the reduction on CO₂ density of three impurities, Ar, H₂ and O₂ at 323.15 K. A maximum reduction of the CO₂ density at a certain pressure under the given temperature is observed for the CO₂ mixtures. The reduction exceeds 25% in CO₂/H₂, 23% in CO₂/O₂ and 20% in CO₂/Ar mixtures. The maxima occur at a pressure equals to 11 MPa. A critical density change in the CO₂/impurities systems can be found from a pressure above the supercritical boundary to 20 MPa. The overall AAD% of the modified PR model is 1.67 in the tested systems listed in Table 4.5. The deviation increases as the impurities systems approach 11 MPa.

Table 4.5 Experimental and modelling data for supercritical CO₂ and supercritical CO₂/impurities system

Pure CO ₂								
T = 323.15 K			T = 373.15 K			T = 423.15 K		
P/MPa	ρ/Kg.m ⁻³	Dev%	P/MPa	ρ/Kg.m ⁻³	Dev%	P/MPa	ρ/Kg.m ⁻³	Dev%
8.38	240.88	0.06	8.87	161.14	0.01	8.95	128.51	0.04
10.33	423.64	0.44	12.64	260.65	0.15	12.91	195.58	0.50
11.22	522.48	0.36	15.96	362.88	0.07	16.51	262.66	0.04
15.30	704.18	0.73	20.84	502.33	0.26	19.29	314.78	0.08
18.20	756.67	0.68	24.86	586.91	0.29	24.57	409.47	0.04
22.25	806.29	0.48	28.33	641.24	0.18	28.74	475.33	0.02
26.81	845.67	0.41	32.61	692.94	0.03	32.36	524.96	0.04
30.72	872.33	0.37	35.64	722.61	0.07	35.87	566.83	0.07

35.33	898.54	0.33	39.62	755.54	0.17	40.13	610.35	0.13
42.72	933.02	0.26	42.83	778.45	0.22	45.49	656.46	0.20
46.25	947.16	0.23	47.37	806.56	0.27	49.15	683.56	0.26
54.48	976.12	0.16						
58.54	988.78	0.13						
CO ₂ /Ar								
T = 323.15 K			T = 373.15 K			T = 423.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
11.14	408.62	1.23	10.20	204.90	8.93	10.88	151.55	3.88
12.59	520.40	0.42	11.15	227.99	8.05	15.47	228.63	3.23
14.93	630.88	1.29	18.75	435.53	3.75	20.23	314.55	1.91
19.57	735.25	0.92	23.60	549.00	2.99	23.52	367.16	2.72
23.17	782.35	0.61	32.60	683.81	2.73	28.11	441.55	1.75
31.31	852.38	0.20	40.28	757.45	2.79	32.99	504.45	2.08
39.19	898.47	0.02	45.15	795.80	3.09	36.43	542.17	2.28
50.05	945.95	0.15				41.77	596.54	1.78
57.77	973.03	0.20						
CO ₂ /O ₂								
T = 323.15 K			T = 373.15 K			T = 423.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
7.72	187.17	0.50	12.25	230.07	2.49	8.29	111.03	3.50
9.66	284.36	2.53	16.38	340.68	2.76	36.91	539.44	2.91
12.14	468.36	4.19	21.39	468.56	2.80	31.79	481.00	3.07
14.26	586.50	1.55	26.57	567.98	2.26	27.89	428.82	3.15
17.18	674.54	0.70	30.96	629.53	1.86	23.54	363.04	3.11
19.78	723.23	0.54	35.78	681.48	1.57	19.31	292.85	3.05
22.55	760.80	0.54	39.62	715.11	1.43	15.27	223.93	2.99
26.61	802.36	0.58						
29.09	822.63	0.63						
29.25	824.10	0.60						
33.87	855.57	0.67						
38.58	882.17	0.71						
CO ₂ /H ₂								
T = 323.15 K			T = 373.15 K			T = 423.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%

7.92	190.22	3.86	8.94	148.12	0.62	9.30	122.97	1.73
9.33	258.13	4.82	11.49	205.62	0.48	14.82	209.16	1.49
10.59	340.86	4.58	15.64	311.41	0.88	19.04	278.08	1.45
11.78	431.42	5.88	20.26	429.98	0.96	23.16	344.24	1.44
13.92	561.01	7.10	25.52	536.82	1.35	27.50	408.52	1.39
17.64	674.17	5.72	29.83	600.93	1.51	31.83	465.07	1.30
20.81	728.22	4.81	34.22	651.34	1.56	35.70	509.04	1.23
24.92	776.69	4.07	38.61	691.84	1.55	40.66	557.39	1.12
28.85	810.95	3.56	42.79	724.08	1.53	43.00	577.94	1.03
36.55	861.13	2.96	43.15	727.26	1.61			
MIX1								
T = 323.15 K			T = 373.15 K			T = 423.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
9.58	276.08	0.94	8.38	139.44	1.44	9.60	130.90	1.56
12.19	466.66	0.45	8.94	151.66	1.71	13.42	191.99	1.57
15.67	630.18	3.22	13.16	253.60	2.45	17.36	258.28	1.54
19.86	718.93	3.90	18.41	396.44	2.32	21.03	320.01	1.41
23.65	767.89	4.00	23.13	506.27	2.55	26.03	399.28	1.37
27.12	801.45	4.03	27.48	582.01	2.97	33.65	500.95	1.66
39.07	879.55	4.05	31.74	637.98	3.26	37.59	544.07	1.88
			36.25	684.48	3.46			
CO ₂ /CO								
T = 323.15 K			T = 373.15 K			T = 423.15 K		
P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%	P/MPa	$\rho/\text{Kg.m}^{-3}$	Dev%
11.62	419.48	4.42	9.37	161.33	2.08	10.20	138.79	3.63
14.45	587.43	0.23	15.48	312.21	2.47	14.16	202.69	3.26
19.82	721.02	1.30	20.50	441.22	2.58	19.83	298.23	2.96
25.41	789.06	1.12	22.68	494.59	1.21	21.73	330.88	2.58
28.37	814.40	0.94	27.74	583.71	0.65	27.75	425.29	2.02
32.99	846.71	0.73	32.45	644.81	0.28	32.30	485.98	1.75
36.63	867.49	0.55	35.33	675.30	0.07	34.64	515.42	1.27
40.77	888.03	0.39	40.56	720.97	0.12	39.98	572.37	0.73
44.02	902.02	0.24	42.59	736.66	0.23	43.06	602.29	0.25
46.85	913.01	0.10	46.11	760.92	0.33	47.23	638.10	0.25

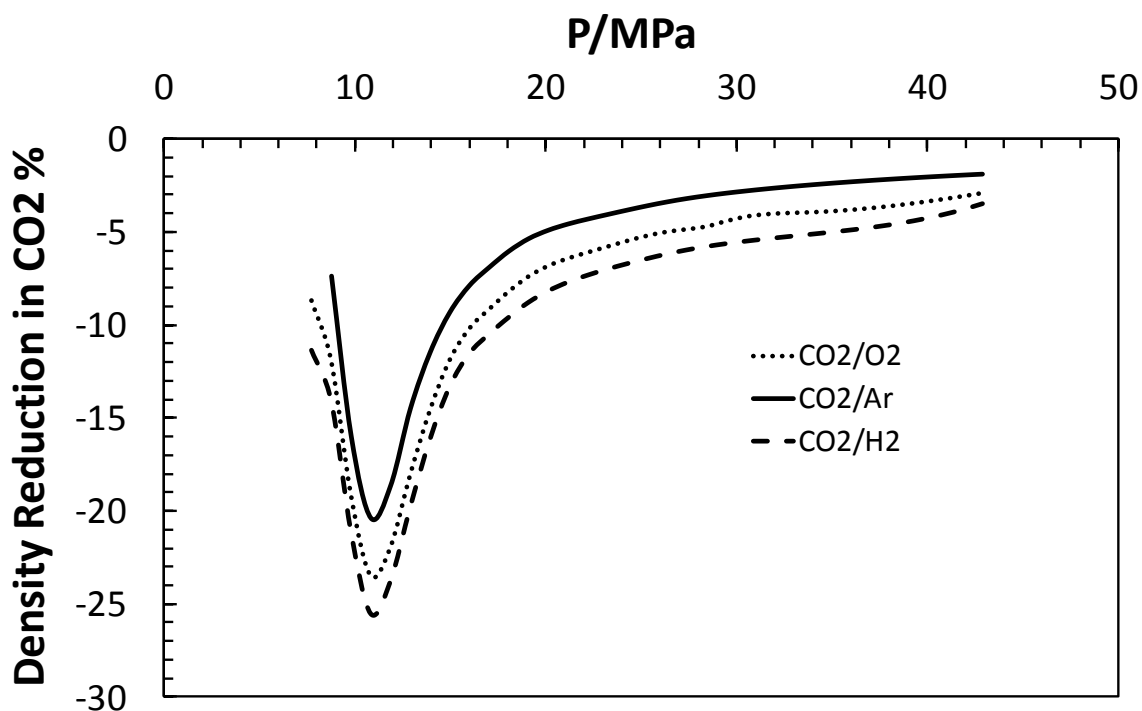


Figure 4.7 Density reduction in supercritical CO_2 in the presence of impurities at 323.15

4.5 Conclusions

The presence of impurities on CO_2 stream can potentially affect the density in both liquid and supercritical regions. Our model density calculations are in a good agreement with the experimental data for both pure and impure CO_2 . The density of pure CO_2 and the CO_2 /impurities mixtures have a stronger dependency on temperature compared to pressure at lower temperatures below the critical temperature. Above the critical temperature, a maximum reduction on the density at a certain pressure under a given temperature was found for each CO_2 /impurity system. Overall, the light molecular weight impurities tend to reduce CO_2 density much more than those with molecular weight close to CO_2 .

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CHAPTER 5: HYDRATE STABILITY OF CARBON DIOXIDE

RICH SYSTEM

5.1 Introduction

Hydrate is a compound containing water discovered by Sir Humphrey Davey in 1811. Hydrate is ice-like solid component formed mainly from water and other small molecules. To form hydrate, three things are required: a sufficient amount of water (host), a former (guest) and right combination of pressure and temperature. Light gases such as CH₄, N₂, O₂, Ar, H₂S and CO₂ can form hydrate in the presence of water. There are basically three common types of hydrate, structure I (SI), structure II (SII) and structure H (SH). Primary, the hydrate structure depends on the size of the gas molecules. The mixture of light gasses can result in a change on both the structure and the hydrate stability zone. Suitable thermodynamic conditions of forming hydrate can be found in pipelines; therefore both experimental and modelling works should be carried-out to define the fluid hydrate zone.

With regard to carbon emissions, the most promising solution currently on the table is CO₂ capture and subsurface storage, primarily by injection into ageing oil/gas reservoirs (where it can be used to improve oil/gas recovery) or deep saline aquifers. Conventionally, subsea pipelines are used to transport CO₂ in liquid or dense-phase state from the mainland for disposal. However CO₂ coming from capture processes is generally not pure and can contain impurities such as N₂, H₂, O₂, H₂S, CH₄, CO and water. The presence of these impurities combined with potentially long distance transportation of CO₂ could lead to challenging engineering and flow assurance issues. The presence of water may result in corrosion, ice and/or gas hydrate formation and pipeline blockage. Furthermore, the gaseous CO₂ rich stream is generally compressed to be transported as liquid or dense-phase state in order to avoid two-phase flow and increase the density of the fluid system. Scenario of excepted impurities on CO₂ is reported in [1] and [2], where CO₂ purity varies from 85 to 99.94% .The type and amount of impurities introduced in CO₂ depends on the type of capture process and the fuels beside the type of solvent used. The presence and type of other components may

differ considerably between post-combustion, pre-combustion and oxy-fuel capture processes [3].

Hydrate formation conditions of CCS streams is not as well known as for natural gas. Experience with pipeline transportation of CO₂ shows that CO₂ hydrates form up to temperatures of 283 K at some transportation pressures [4], [5]. Hydrates can form plugs in pipelines, either blocking valves, fouling up instrumentation or in the extreme case block the entire bore of the pipeline at a certain location. During depressurization the acceleration of a hydrate plug can cause structural damage to the pipeline wall in small radius bends [6]. Carbon dioxide is known to form structure I gas hydrates under the appropriate temperature and pressure conditions. As carbon dioxide is sub-critical at hydrate forming conditions and has a relatively low vapour pressure, different phases can be found in the hydrate-carbon dioxide-water system: a hydrate phase, a water rich liquid phase, an ice phase, a carbon dioxide rich vapour phase and a carbon dioxide rich liquid phase as well as two quadruple points. Experimental data for carbon dioxide hydrates have been measured and reported by various authors in the different hydrate regions. Data are widely available for the carbon dioxide/methane system. No experimental data were found in the open literature for CO₂ with oxygen, argon and H₂S. Data are available but limited to 1- 2 sources for CO₂ with nitrogen, hydrogen and carbon monoxide. Table 5.1 gives a list of these data, reporting temperature range and source of the experimental data. The aim of this study is to analyse the risk of hydrate formation in a rich carbon dioxide stream in a pipeline where the stream could be transported at a pressures up to 400 bars.

Table 5.1 Literature data for hydrate equilibria of CO₂ and CO₂ mixtures

Comp.	Temperature Range	Type of equilibrium	CO ₂ /mole fraction	Ref.
Pure CO ₂	273.7 – 282.9	V-L-H		[7]
	283.2 – 292.7	L _{CO₂} -V-H		[8]
	256.8 – 285	I-V-H + V-L-H + L _{CO₂} -V-H		[9]
	151.7– 192.5	I-V-H		[10]
	194.5 – 217.8	I-V-H		[11]
	277.2 – 283.1	V-L-H		[12]
	273.9 – 283.2	V-L-H		[13]
	279.6 – 283.9	V-L-H + L _{CO₂} -V-H		[14]
	271.6 – 283.2	V-L-H		[15]
	274.3 – 282.9	V-L-H		[16]
	274.7 – 279.7	V-L-H		[17]
	276.52 – 283.36	V-L-H + L _{CO₂} -V-H		[18]
	289.73 – 293.97	L _{CO₂} -V-H		[19]
	274.95 – 286.05	V-L-H + L _{CO₂} -V-H		[20]
CO ₂ /CH ₄	264.1 – 275.5	V-L-H	≈0.5	[17]
	275.5 – 285.7	V-L-H	0.274-0.824	[21]
	273.7– 287.6	V-L-H	0.08 to 0.85	[22]
	280.3	V-L-H	0 to 1	[23]
	273.5 – 283.1	V-L-H	0.13 - 0.53	[24]
	273.56 – 285.56	V-L-H	0 to 1	[25]
	272.66 – 285.76	V-L-H	0 to 1	[26]
	273.5 – 283.2	V-L-H	0.95 -0.9658	[27]
	275.9 – 277.7	V-L-H	0.06 – 0.25	[28]
CO ₂ /N ₂	273.1 – 280.2	V-L-H	0.91-0.97	[17]
	273.4 – 284.25	V-L-H	0 to 1	[29]
	272.85 – 284.25	V-L-H	0.07-0.97	[30]
CO ₂ /H ₂	274.15 – 283.23	V-L-H	0.5	[31]
	274.3 – 281.9	V-L-H	0.19-1	[32]
	273.9 – 281.9	V-L-H	0.4-0.83	[33]
CO ₂ /CO	274.15 – 283.2	V-L-H	0.1	[34]

5.2 Experimental Section

5.2.1 Materials

The following compounds were used to make the different synthetic mixtures studied in this work:

- Carbon Dioxide and Methane Supplied by Air Products, Research Grade
- Nitrogen, Argon, Oxygen, Hydrogen and Carbon monoxide Supplied by BOC, Research Grade.
- The water used was ultra pure water produced by an apparatus “PURELAB OPTION” made by ELGA and characterized up to 18.2 M Ω -cm resistance and a maximum organic carbon content of less than 10 ppb.

The binary mixtures of CO₂/impurities system are listed in [Table 5.2](#). The mixtures were prepared gravimetrically (uncertainty ± 0.3 mole %).

Table 5.2 List of impurities composition used in the experiments

Impurity	CO ₂	Mole%
Nitrogen		4.6%
Methane		5.9%
Carbon monoxide	Balance	5.9%
Oxygen		5.3%
Argon		5.0%
Hydrogen		7.1%

5.2.2 Apparatus

An autoclave rig with 300 cm³ volume was used in this work. The basic rig design is comprised of a high pressure cell fitted with a magnetic mixer with a maximum rotation speed of 1,500 revolutions per minute giving good mixing to the cell contents. The rig is surrounded by a jacket through which coolant is circulated using a refrigerated circulator. The rig is well insulated and can be used at temperatures in the range of 233 to 333 K and at pressures up to 60 MPa. The pressure is measured using a strain gauge pressure transducer. The temperature is measured using a Platinum Resistance

Temperature (PRT) Probe. A schematic of the hydrate mixed autoclave rig is shown in Figure 5.1.

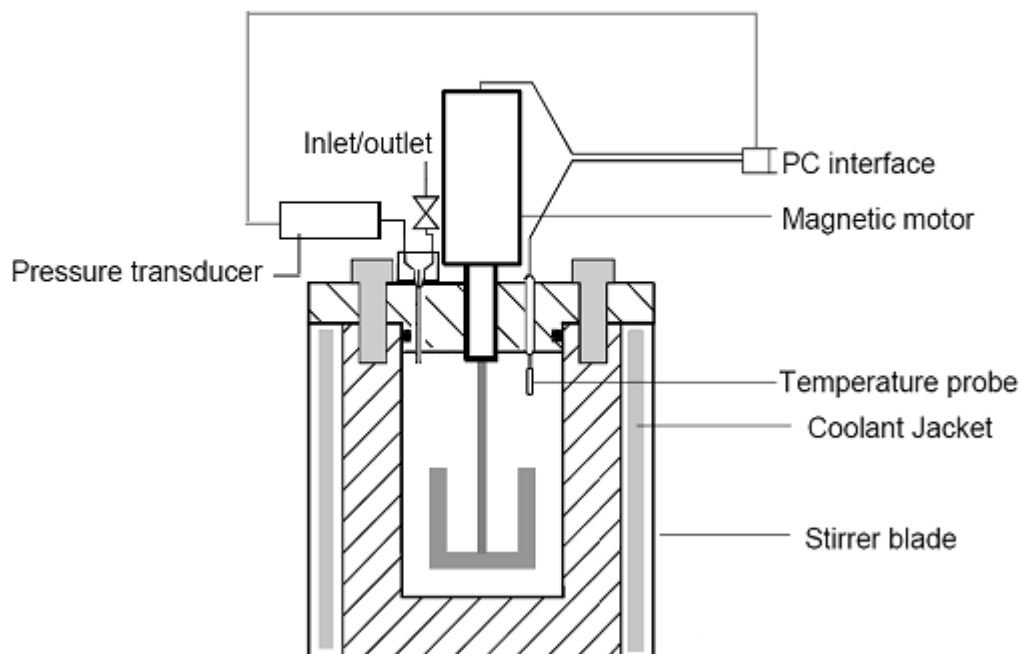


Figure 5.1 Schematic of hydrate mixed autoclave rig

5.2.3 Procedures

The equilibrium cell and its loading lines are vacuumed prior to introduction of about 90 cm³ deionised water. Then, CO₂ or CO₂ rich gas is introduced into the cell directly from a pressurized cylinder (through preliminary evacuated transfer lines) to a pressure corresponding to the pressure of the first measurement. Pressure and temperature are initially set to a point well outside the expected hydrate stability zone for the system under study. After introduction of the gas into the cell, stirring of the cell is started, to accelerate equilibrium in the system. The temperature is then lowered to form hydrates well inside the hydrate stability zone, i.e. ~ -5 °C. After the hydrate formation is observed (sudden pressure drop is observed in the system), the temperature is increased step-wise and the dissociation point is detected by the sudden change (break point) in pressure/temperature profile. The other dissociation points corresponding to the higher pressures were determined by injecting more gas and the above steps are repeated.

5.3 Modelling Section

A general phase equilibrium model based on the equality of component fugacities in all phases has been used to model the phase behaviour of the carbon-dioxide water system. A description of the thermodynamic model can be found elsewhere [35]. In summary, the statistical thermodynamics model uses the Valderrama modification of the Patel and Teja (*VPT-EoS*) equation of state [36] and non-density-dependent (*NDD*) mixing rules [37] for fugacity calculations in all fluid phases. A detail of the VPT EOS is given in Chapter 2. BIPS between CO₂ and N₂, H₂, CO, O₂ Ar and CH₄ were tuned using the data listed in Table 5.3. BIPS between H₂, N₂ and CO were set to zero. Binary interaction parameters (BIPS) between water and some of these components can be found elsewhere [38].

Table 5.3 VLE and BP data for binary CO₂/impurities system

Mixtures	Source
CO ₂ /N ₂	[39]-[54]
CO ₂ /CH ₄	[42], [43], [46], [49], [52], [55]-[63]
CO ₂ /O ₂	[51], [64]
CO ₂ /H ₂	[41],[51], [65]- [68]
CO ₂ /Ar	[69], [70]
CO ₂ /CO	[71], [72]

The hydrate phase is modelled using the solid solution theory of van der Waals and Platteuw [73], as developed by Parrish and Prausnitz [74]. The equation recommended by Holder *et al.* [75] is used to calculate the heat capacity difference between the empty hydrate lattice and pure liquid water. The Kihara model for spherical molecules is applied to calculate the potential function for compounds forming hydrate phases [76]. The fugacity of water in the hydrate phase is given by:

$$f_w^H = f_w^\beta \exp\left(-\frac{\Delta\mu_w^{\beta-H}}{RT}\right) \quad (5.1)$$

Where H , β and μ refer to hydrate, empty hydrate lattice and chemical potential respectively. $\Delta\mu_w^{\beta-H}$ is the chemical potential difference of water between the empty hydrate lattice (μ_w^β) and the hydrate phase (μ_w^H). It is defined as following:

$$\Delta\mu_w^{\beta-H} = \mu_w^\beta - \mu_w^H = RT \sum_m \bar{v}_m \ln \left(1 + \sum_j C_{mj} f_j \right) \quad (5.2)$$

Where \bar{v}_m is the number of cavities of type m per water molecule. f_j is the fugacity of the gas component j . C_{mj} is the Langmuir constant, which accounts for the gas-water interaction in the cavity. The relation for the Langmuir constant can be developed from the potential energy and numerical values for the Langmuir constant can be calculated by choosing a model for the guest-host interaction:

$$C_{mj}(T) = \frac{4\pi}{kT} \int_0^\infty \exp \left[-\frac{w(r)}{kT} \right] r^2 dr \quad (5.3)$$

Where k and $w(r)$ are Boltzmann's constant and the spherically symmetric cell potential in the cavity with r measured from the centre. Holder et al. [75] showed that the chemical potential difference of water between the empty hydrate lattice and the liquid phase can be defined as following:

$$\frac{\Delta\mu_w^{\beta-L}}{RT} = \frac{\Delta\mu_w^0}{RT_0} - \int_{T_0}^T \frac{\Delta h_w^{\beta-L}}{RT^2} dT + \int_{P_0}^P \frac{\Delta v_w^{\beta-L}}{RT} dP \quad (5.4)$$

Where 0 superscript stands for reference property. v and h refer to molar volume and enthalpy respectively. $\Delta\mu_w^0$ is the reference chemical potential difference between water in the empty hydrate lattice and pure water at 273.15 K.

5.4 Results & Discussion

The hydrate formation pressures of pure CO₂ in water were measured and compared with the data reported in the above literature (Table 5.1) in order to validate the

experimental apparatus and procedures. Figure 5.2 shows the experimental and predicted phase boundaries for carbon dioxide hydrates. As can be seen, the model predictions are in good agreement with our experimental data and the data collected from the open literature, demonstrating the validation to our experimental work and the reliability of the thermodynamic model. Our experimental data were found to be consistent with the data reported by Deaton and Frost [7] below the system saturation pressure and are also consistent with Ng and Robinson [14] data above the saturation pressure.

The three phases HLV equilibria of the ternary $\text{CO}_2 + (\text{N}_2 \text{ or } \text{CH}_4 \text{ or } \text{O}_2 \text{ or } \text{Ar} \text{ or } \text{H}_2 \text{ or } \text{CO}) + \text{water}$ systems were determined at a constant composition (4.56% N_2 , 7.09% H_2 , 5.85% CH_4 , 5.87% CO , 5.03% Ar , and 5.34% O_2 in dry basis). The mixture hydrates formed over the wide temperature and pressure ranges of 276-288 K and 2-15 MPa.

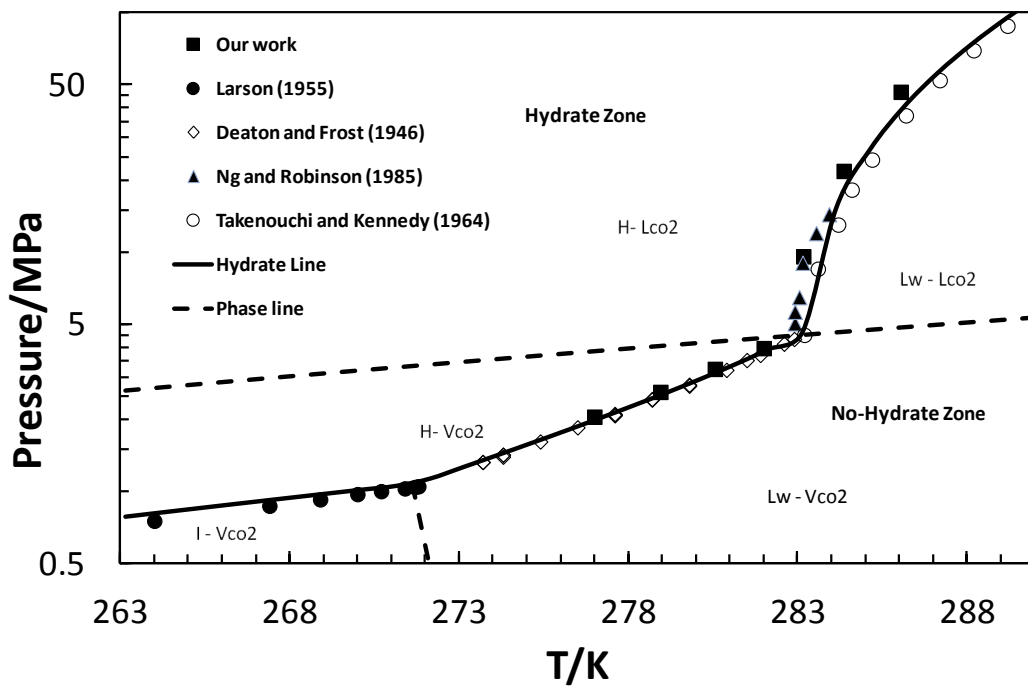


Figure 5.2 Hydrate phase equilibria of the binary $\text{CO}_2 + \text{water}$ mixtures

Table 5.4 Equilibrium hydrate dissociation pressure for the CO₂/impurities mixtures

Pure CO ₂			CO ₂ /N ₂			CO ₂ /H ₂		
T/K	P/MPa	Aqueous Fraction	T/K	P/MPa	Aqueous fraction	T/K	P/MPa	Aqueous fraction
278.95	2.60	0.94	276.91	2.05	0.96	274.45	1.64	0.97
277.00	2.05	0.96	279.65	2.82	0.94	275.45	1.85	0.96
280.56	3.24	0.93	281.23	3.66	0.92	280.55	3.33	0.93
282.00	3.96	0.91	283.64	5.72	0.86	284.05	6.83	0.85
283.17	9.58	0.53	287.40	40.82	0.51	285.65	12.69	0.56
284.36	21.79	0.52	288.55	55.11	0.50	286.75	16.51	0.55
286.05	46.65	0.50						
CO ₂ /CH ₄			CO ₂ /CO			CO ₂ /O ₂		
T/K	P/MPa	Aqueous Fraction	T/K	P/MPa	Aqueous Fraction	T/K	P/MPa	Aqueous Fraction
276.00	1.82	0.96	273.15	1.38	0.97	276.75	2.05	0.96
279.20	2.68	0.94	278.35	2.63	0.94	278.85	2.68	0.94
281.35	3.61	0.92	280.75	3.64	0.92	281.75	4.05	0.91
284.15	5.81	0.57	283.85	6.69	0.83	284.01	7.04	0.56
285.75	12.25	0.54	284.75	11.63	0.55	285.25	13.19	0.54
286.95	19.97	0.53	285.85	21.30	0.53	285.75	18.33	0.53
285.75	12.25	0.54						
CO ₂ /Ar								
T/K	P/MPa	Aqueous Fraction						
275.45	1.72	0.97						
280.55	3.22	0.93						
284.05	5.81	0.85						
285.25	11.26	0.55						
285.65	16.07	0.53						

The measured equilibrium data together with their predictions are plotted in [Figure 5.8](#). In the same plots, the vapour liquid equilibrium lines for all binary mixtures are plotted together with the HL_V and HL_L lines produced by VPT EOS model. The quadruple points where ice, liquid vapour, and hydrate (I + V + L + H) coexist for all

systems has been defined by the model. The quadruple point is noted by the phases that are in equilibrium and it has been evaluated by the model. The experimental results of the equilibrium hydrate dissociation pressure for the CO₂/impurities mixtures are also listed in [Table 5.4](#)

Table 5.4. The aqueous fraction is the water ratio of water to the fluid mixture in the cell. As simple hydrates, CO₂ and N₂ form SI and SII, respectively [78]. The mixed hydrate structure is considered to be either SI or SII depending on the relative ratio of these two different gas molecules occupied in the small and large cavities. According to the literature [77], 15 mol% of CO₂ gaseous mixture was known as the boundary of coexisting SI and SII hydrate. For rich N₂, Seo and Lee [78] had the same conclusion where SI is crystallized at the gas mixture composition range of 3-20 mol % CO₂ and transformed to SII at 1 mol%. As generally expected, the hydrate formation line is located between two pure CO₂ and N₂ HLV lines, and the overall experimental data were well predicted by the thermodynamic model. The binary mixture containing 4.56 mol% N₂ is found slightly higher than that of pure CO₂ HLV phase. Phase equilibrium relations for the ternary system of H₂+CO₂+H₂O in the presence of hydrate phase have been investigated. It is well known that H₂ is too small to stabilize cavities any hydrates by itself except for a high-pressure region of the GPa order where it can form SII clathrate hydrates [79-80], that is, H₂ almost never contribute to the stability of hydrate cage. The gas hydrate obtained from the H₂+CO₂ mixture has been analyzed previously [81] by use of Raman spectroscopy. The results reveal that H₂ is not enclathrated in the hydrate-cages at low pressure. That is, H₂ behaves only like a diluent gas toward the formation of CO₂ hydrate. The experimental results obtained in the current study ([Table 5.4](#) and [Figure 5.4](#)) indicate that the hydrate dissociation pressures (HLV) from the CO₂/H₂ binary mixture are greater than those for pure CO₂. Kumar et al. [33] found that the increase of hydrogen in the gas mixture tends to shift the equilibrium to higher pressures at a given temperature. They reported that the mixture CO₂/H₂ forms SI like that formed by pure CO₂.

Both methane and carbon dioxide form SI hydrate. Phase diagram for CH₄/CO₂ hydrate is shown next to the previous mentioned mixtures in [Figure 5.5](#). The HLV curve of the

mixture was found to locate between the carbon dioxide and methane hydrates curves. At a fixed temperature, the phase boundary (HLV) of CO₂ hydrate is shifted to higher pressures by adding 5.85% CH₄. Seo and Lee [25] indicated that in rich CO₂ mixture (above 60 mol% CO₂) the HLV curves of the mixed hydrates closely moved towards that of the simple carbon dioxide hydrate. Three HLV phase equilibrium data have been measured by Mohammadi et.al [34] for 10%CO₂+CO mixture. The mixture formed is SI and they plotted the data together with both pure CO₂ and pure CO. No literature was found for the other ternary CO₂/impurities (Ar, O₂)-water mixtures. According to [82], formation of SH is established for the argon-water system by using differential thermal analysis while O₂ forms SII [83, 84]. In our calculations, the rich CO₂-water mixture is assumed to be structure I under the temperature conditions investigated in this work.

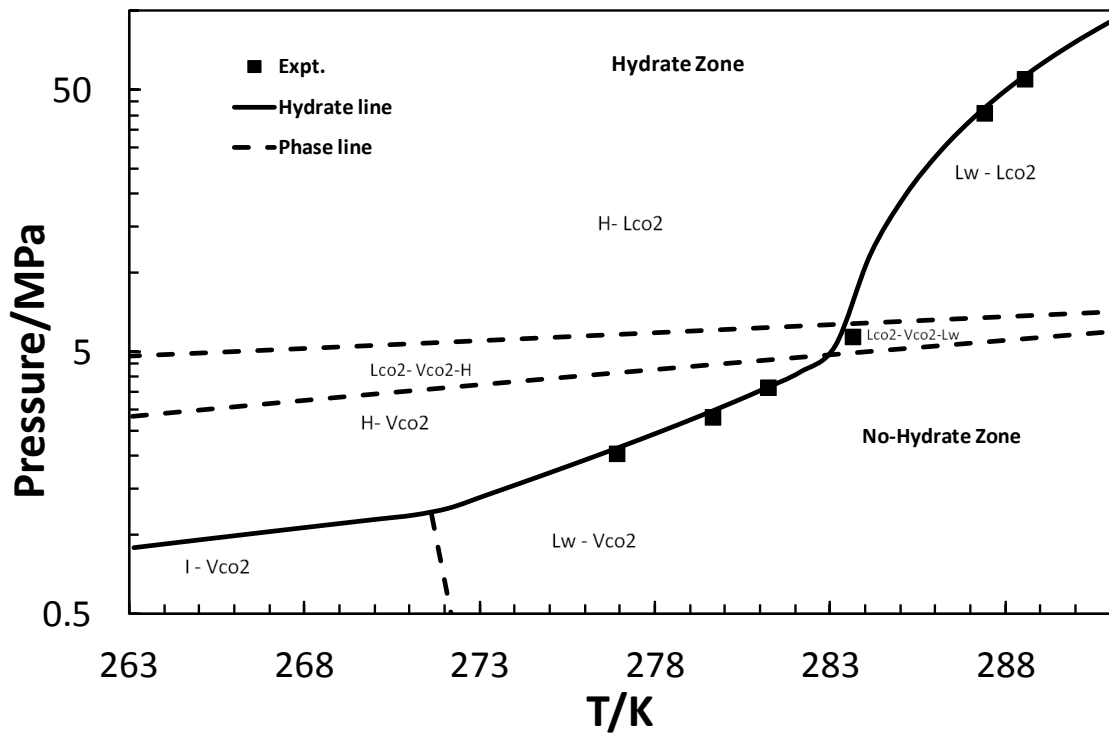


Figure 5.3 Hydrate phase equilibria of the ternary CO₂ + N₂ + water mixtures

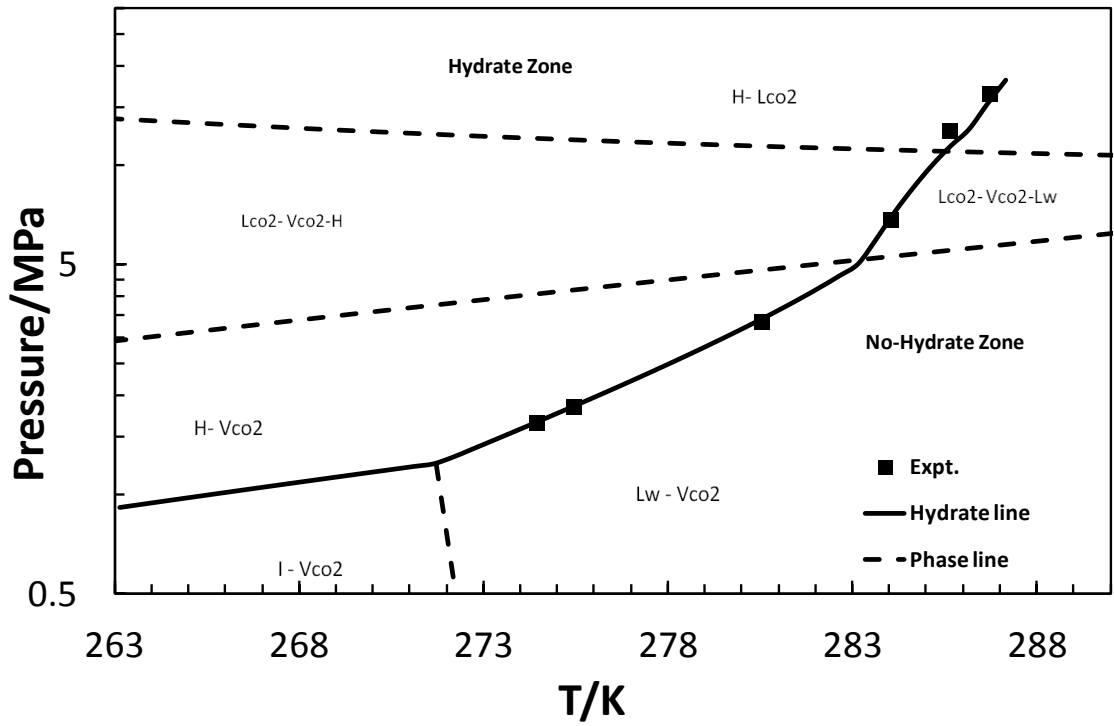


Figure 5.4 Hydrate phase equilibria of the ternary $\text{CO}_2 + \text{H}_2 + \text{water}$ mixtures

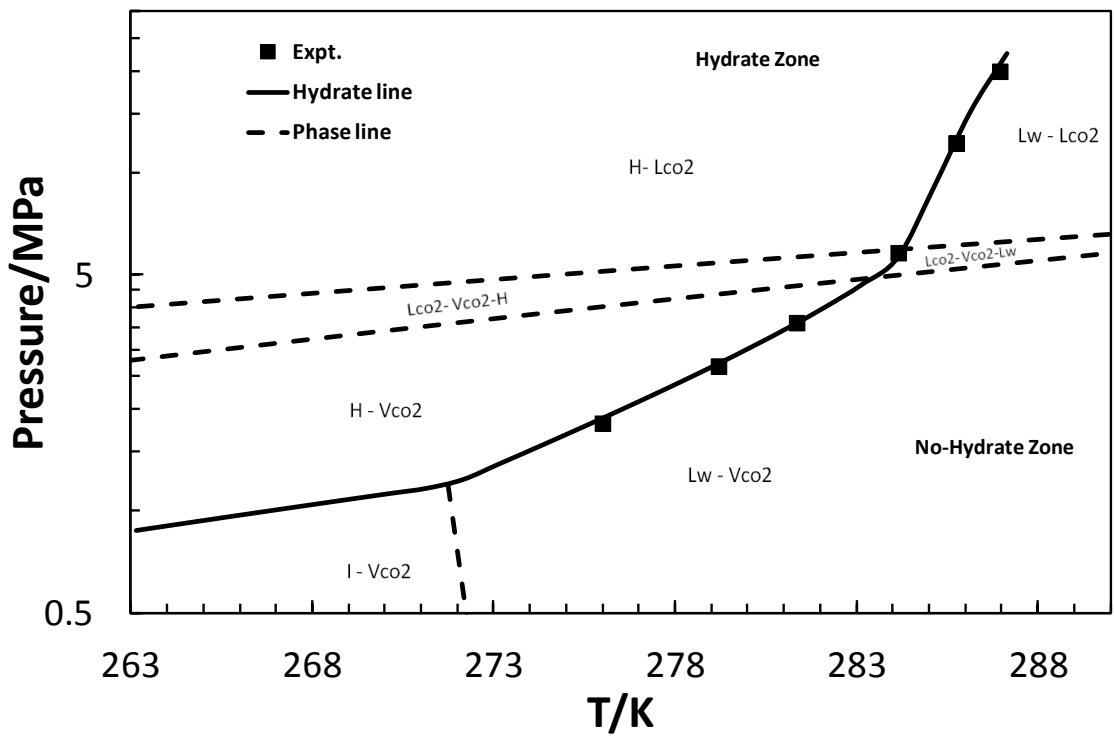


Figure 5.5 Hydrate phase equilibria of the ternary $\text{CO}_2 + \text{CH}_4 + \text{water}$ mixtures

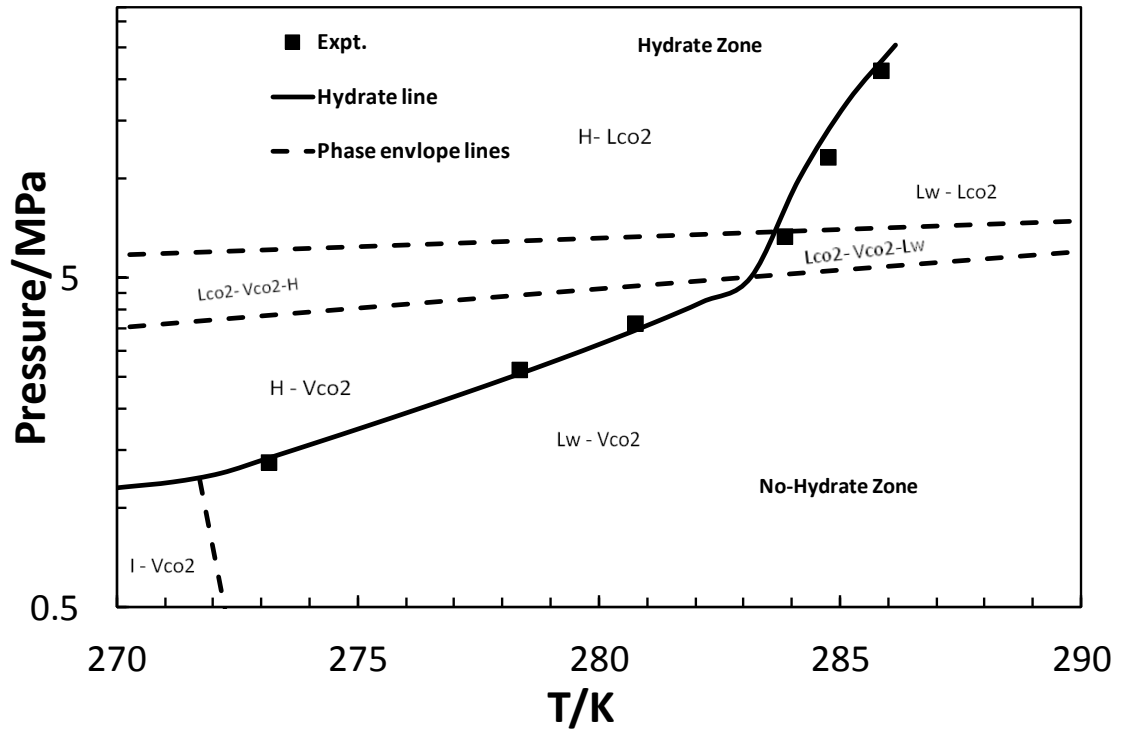


Figure 5.6 Hydrate phase equilibria of the ternary CO₂ + CO + water mixtures

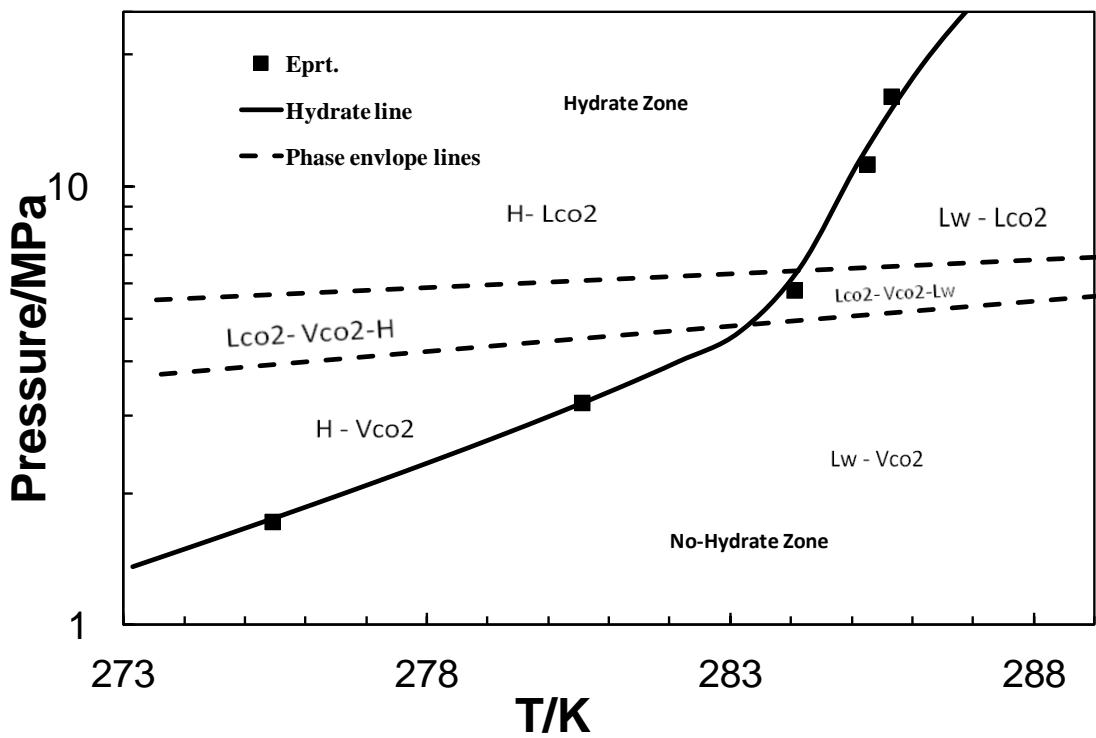


Figure 5.7 Hydrate phase equilibria of the ternary CO₂ + Ar + water mixtures

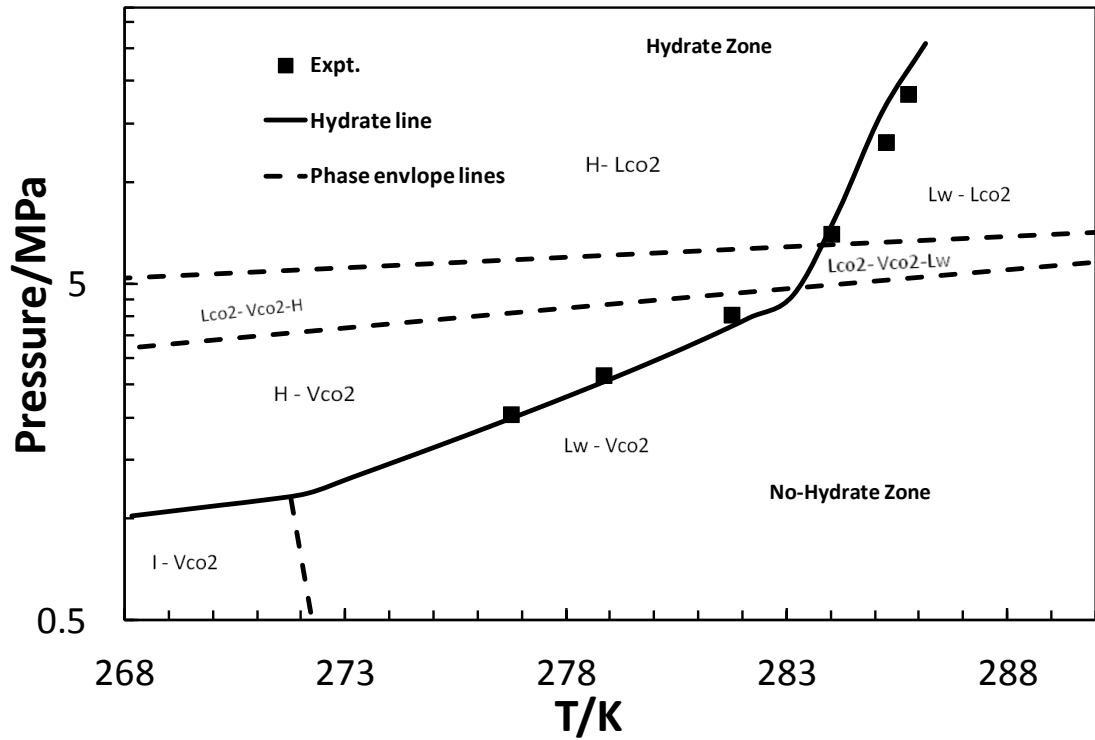


Figure 5.8 Hydrate phase equilibria of the ternary $CO_2 + O_2 + \text{water}$ mixtures

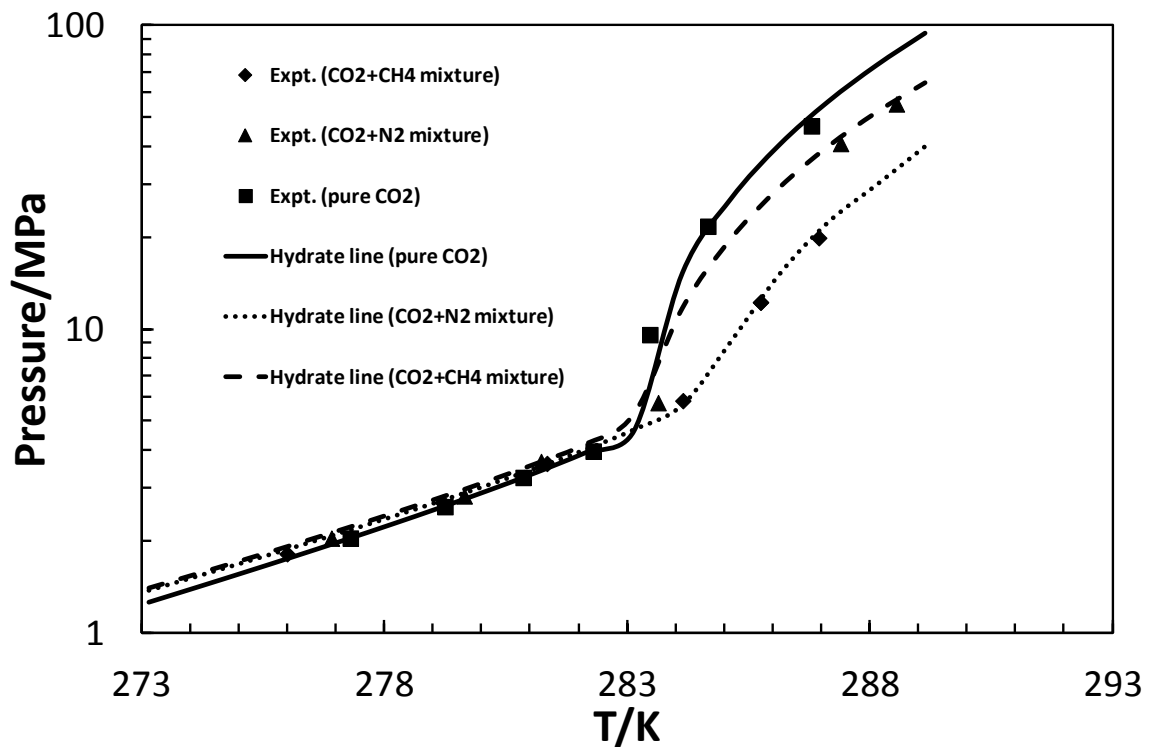


Figure 5.9 Effects of impurities on the hydrate phase equilibria of CO_2

At the composition of the gas mixtures systems, HLV equilibrium lines compared to the pure CO₂ were shifted slightly to higher pressure and lower temperature domain. In the other hand at higher pressure, above the bubble point of the systems, HLL equilibrium lines were shifted to lower pressure and higher temperature domain. The increase stability of the hydrate zone of CO₂ is higher in the presence of methane. The effect is less significant in the presence of nitrogen as shown in [Figure 5.9](#). The other impurities (Ar, O₂, CO and H₂) seem to follow closely the hydrate phase path of the CO₂/N₂ system. In the vapour region the increase in the hydrate stability pressure of pure CO₂ is within 5%. In liquid region, at 288 K the reduction of the hydrate stability pressure can reach 60% in the CO₂/CH₄ mixture. At the same condition, the reduction in pressure of the hydrate stability zone for the other impurities excluding methane is between 28 to 35%.

5.5 Conclusions

In the work, the effects of impurities (CH₄, N₂, O₂, H₂, CO and Ar) on the hydrate stability conditions of CO₂ were determined. The hydrate equilibrium pressures were increased apparently with addition of the tested impurities in the vapour region. The effect of the tested impurities can be classified into two groups. Methane, as one group, tends to increase the hydrate pressure at constant temperature by double the increase achieved by the other impurities (N₂, O₂, H₂, CO and Ar) as second group. In the work, the VPT equation of state in conjunction with the van der Waals–Platteeuw model was used to predict successfully the hydrate dissociation conditions of CO₂ and six CO₂ binary mixtures. A good agreement between the predicted and experimental data was found.

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CHAPTER 6: VISCOSITY OF CARBON DIOXIDE RICH SYSTEMS

6.1 Introduction

In many engineering processes, including transport phenomena, understanding viscosity of liquid and liquid mixtures is necessary. Transporting fluids including liquid CO₂ are subjected to various resistances due to friction. Fluid viscosity causes friction which causes pressure drop according to Darcy-Weisbach equation. Understanding the effect of impurities on CO₂ viscosity leads us defining the capacity of pumps and compressors and hence the cost.

Viscosity data of pure CO₂ gases and impure gases are readily available in literature. Li et al. [1], have published recently an extensive review of available experimental data and models for the transport properties of CO₂-mixtures relevant to CO₂ capture, transport and storage. They concluded that no data were available for CO₂ viscosity in the presence of impurities in the liquid and supercritical regions. A summary of the experimental data available in literature for CO₂/impurities systems is listed in Table 6.1. Several models [2-5] ranging from those based on kinetic theories to completely empirical have been proposed. Vesovic et al. (1990) [5] proposed a viscosity correlation that covered a wide range of pressure and temperature. They found that their model can predict liquid and supercritical CO₂ in the range from 5 to 7%. Fenghour et al. [6] extended Vesovic et al. [5] work to improve viscosity of liquid CO₂. The uncertainties associated with the proposed representation vary from 0.3% for dilute gas near room temperature to 5.0% at the highest pressures. It has a very good accuracy for pure CO₂. However, it cannot be applied for mixtures without modification. Quinones-Cisneros et al. [3] proposed what is called a “f-theory” model of viscosity for both pure components and their mixtures. The model is modified to be used in conjunction with the SRK, PR and PRSV cubic equations of state in order to develop three one parameter general models for viscosity prediction [7]. In a different approach, the Pedersen correlation which is a modified form of Ely and Hanley model [8] employed the corresponding states principle to predict the viscosity of pure components or their mixtures, knowing

the viscosity of a reference substance (methane). The method of Lohrenz–Bray–Clark (LBC), which relates the residual viscosity to the reduced density, is the most widely used engineering tool to predict the viscosity of reservoir fluids. It is a predictive model for gas or liquid viscosity. The LBC correlation is a modification of Jossi et al [9] correlation for pure components, and is suitable for gases and light oils. Pedersen, LBC and f-theory (general form of f-theory) models were tested by Al-Marri [10]. The author concluded that for supercritical CO₂ all models are very good to represent the viscosities when the pressure is less than 10.34 MPa. Above that pressure, CSP over predicts while the LBC underpredicts the viscosities but the f-theory and GF-f-theory are much better to predict the viscosities. The ADD% is 1.25 for f-theory, 1.39 for GF-f-theory, 15.15 for CSP (Pedersen model) and 6.38 for LBC. The f-theory model is rarely used in oil industries compare to the LBC and Pedersen models. However, the later models cannot be used to predict CO₂ and CO₂ mixtures without modification. Therefore, the viscosity models considered in this paper work are the Pedersen model (with CO₂ as a reference) and LBC (tuned) based on PR-EOS.

Table 6.1 Source of viscosity data for CO₂/impurities system

System	T Range/K	P/MPa Range	Ref.
CO ₂ + CH ₄	298	>0.1	[11]
CO ₂ + CH ₄	323–473	3.4–68	[12]
CO ₂ + CH ₄	293–303	0.1–2.6	[13]
CO ₂ + H ₂	300–551	0.1	[14]
CO ₂ + H ₂	291	0.1	[15]
CO ₂ + H ₂	298	0.1	[16]
CO ₂ + H ₂	295–303	0.098	[17]
CO ₂ + H ₂	295–303	0.1	[18]
CO ₂ + H ₂	500–1100	0.3	[19]
CO ₂ + N ₂	293	0.1–2.14	[20]
CO ₂ + N ₂	293–303	0.1–2.6	[21]
CO ₂ + N ₂	295–303	0.098	[17]
CO ₂ + N ₂	289	2–12	[22]

CO ₂ + N ₂	298–873	0.1	[23]
CO ₂ + Ar	293–303	0.1–2.6	[21]
CO ₂ + Ar	310–521	0.1	[24]
CO ₂ + Ar	298–873	>0.1	[23]
CO ₂ + CO	298–473	0.1	[25]
CO ₂ + O ₂	298–674	0.1	[26]
CO ₂ + O ₂	295–303	0.098	[17]
CO ₂ + O ₂ + CO + H ₂ + CH ₄ + N ₂	293–1287	0.1	[27]
O ₂ + N ₂ + CO ₂	317–1161	0.1	[28]
CO ₂ + O ₂ + N ₂ , CO ₂ + O ₂ + H ₂	295–303	0.098	[17]
Ar + CO ₂ + N ₂ , CH ₄ + CO ₂ + N ₂	298–873	0.1	[23]
Ar + N ₂ + CO ₂ , CH ₄ + N ₂ + CO ₂	298–873	0.1	[29]
N ₂ + O ₂ + CO ₂ + N ₂ O	313–413	>0.1	[30]

In this work, the viscosity of pure CO₂ and five binary CO₂/impurities mixtures were measured from 280 to 300.15 K in the liquid phase and from 308.15 to 343.15 K in the supercritical phase. Both measurements were run for pressure up to 41 MPa in liquid and supercritical regions. This work was essentially presented on a steering meeting held on 3rd May 2012 in Heriot-Watt University (Alsiyabi I. and Nazeri M.).

6.2 Experimental Methods and Equipment

6.2.1 Materials

The following compounds were used to make the different synthetic mixtures studied in this work:

- Carbon dioxide, Supplied by Air Products, Research Grade
- Methane, Supplied by Air Products, Grade N4.5
- Nitrogen, Supplied by BOC, Research Grade
- Argon, Supplied by BOC, Research Grade
- Carbon monoxide, Supplied by BOC, Research Grade

- Oxygen, Supplied by BOC, Research Grade
- Hydrogen, Supplied by BOC Gases, $\geq 99.995\%$

The systems listed in [Table 6.2](#) were prepared gravimetrically and were used to conduct the viscosity tests. A multi component mixture, MIX1, (composition given in [Table 6.3](#)) supplied by BOC was also used to carry out the tests.

Table 6.2 The Composition of mixtures used for the experiments

Component	CO ₂ %mole	Impurity %mole (± 0.3)
Pure CO ₂	100	0
CO ₂ + methane	95	5
CO ₂ + nitrogen	95	5
CO ₂ + argon	95	5
CO ₂ + carbon	95	5
CO ₂ + oxygen	95	5
CO ₂ + hydrogen	95	5

Table 6.3 The mixture composition (MIX1)

Mixture Components	% mole(± 0.05)
Carbon Dioxide	Balance
Methane	0.67
Hydrogen	0.82
Nitrogen	1.41
Carbon m	0.21
Argon	1.21
Oxygen	0.08
Total	100

6.2.2 Experimental Equipment

All viscosity measurements were conducted in an in-house designed and constructed set-up. A schematic view of the setup is shown in [Figure 6.1](#). This setup has been designed to have 200 MPa maximum working pressure and 520 K maximum working

temperature. The set-up is comprised of two small cylinders, with volumes of 100 and 25 cm³, connected to each other through a capillary tube with a measured length of 14.781 metres and a calculated internal diameter of 0.296 mm. A three-way valve is installed on top of one of the cylinders to inject the sample inside the cylinders and tube system. All of the cylinders and tube are in a temperature controlled bath filled with water. The base side of the two cylinders are connected to the opposite sides of a push-pull, motor driven mercury pump. This pump can move the sample fluid forwards and backwards between the two cylinders. There is also a hand pump connected to the system to control the pressure of the entire fluid system by injection and withdrawal of mercury. Both the opposed piston pump and the hand pump are fitted with Mitutoyo linear transducers readable to 0.005 mm on Mitutoyo SD-D1E readouts. As a 1 mm movement represents 0.151 cm³ displacement in both pumps, the readability is 0.000755 cm³. The opposed piston pump has a variable control with which the speed can be adjusted to a maximum of 5cm³/sec. The rate can be set with a margin of error of ± 0.00003 cm³/sec.

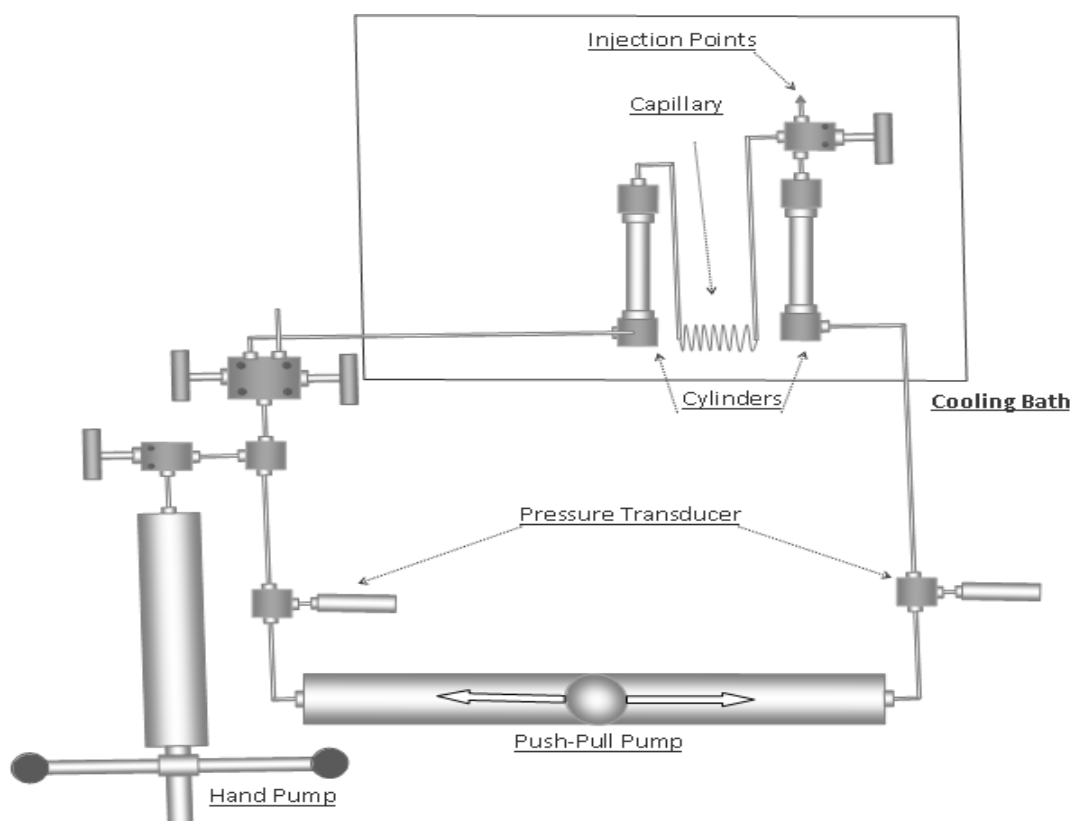


Figure 6.1 schematic view of the viscosity experiments setup

6.2.3 Experimental Procedure

The capillary tube viscosity measurement method has been employed to measure the viscosity of CO₂ systems with impurities above the critical point. In each test, one of the cylinders was loaded with the prepared synthetic mixtures after vacuuming the entire system. Then after disconnecting the sample cylinder from the system, the sample fluid was pushed through the capillary tube into the other cylinder using the push-pull mercury pump. The temperature of the system was set to the desired condition and the desired pressure was set using the hand pump. Once conditions had stabilized the sample was pumped through the capillary tube at a set flow rate. Pumping the sample fluid through capillary tube by the piston pump results in a dynamic differential pressure that was monitored and recorded until stable. Then the pump was stopped to record the static differential pressure. The difference between the dynamic and static differential pressure was used as the pressure drop across the tube. To ensure laminar flow conditions, the Reynolds number was checked for the flow rates in which the measurements were performed. The Poiseuille equation, below, can relate the pressure drop across the capillary tube to the viscosity, tube characteristics and also flow rate for laminar flow:

$$\Delta P = \frac{128 L Q \eta}{C \pi D^4} \quad (6.1)$$

Where, ΔP is differential pressure across the capillary tube viscometer in psi, Q represents flow rate in cm³/sec, L is length of the capillary tube in cm, D refers to internal diameter of the capillary tube in cm equals 0.0296 cm, η is the viscosity of the flown fluid in cP and C is a unit conversion factor equal to 6894757 if the above units are used. The internal diameter of the tube was calculated by knowing the length and volume of the tube. The tube length changes with temperature but this had no noticeable influence on the obtained viscosity. The set flow rate has no effect on the accuracy of the viscosity measurement. Only differential pressure is a variable in the above formulation and can cause error in the viscosity measurements. The usual uncertainty in differential pressure measurement is 0.01 psi and this leads to $\pm 1\%$ of error in the calculated viscosity for those measured in this study.

6.3 Modelling

Two models are suggested in this work to predict the viscosity of CO₂ dominated systems with impurities: CO₂-Pedersen and Lohrenz–Bray–Clark (LBC). The CO₂-Pedersen model predicts the viscosity using the corresponding states theory. The LBC correlation is a fourth-degree polynomial equation in the reduced density which can be tuned to match to experimental data.

6.3.1 CO₂-Pedersen Model

The model has been described in Pedersen and Christensen [31]. According to the corresponding states principles applied to viscosity, the reduced viscosity, $\eta_r = \frac{\eta(P,T)}{\eta_c}$, for two components at the same reduced pressure, $P_r = \frac{P}{P_c}$, and reduced temperature, $T_r = \frac{T}{T_c}$, will be same.

$$\eta_r = f(P_r, T_r) \quad (6.2)$$

Based on the dilute gases considerations, viscosity at critical point can be approximated as:

$$\eta_c \approx \frac{P_c^{2/3} M^{1/2}}{T_c^{1/6}} \quad (6.3)$$

Where, M denotes the Molecular weight. Thus, the reduced viscosity can be expressed as:

$$\eta_r = \frac{\eta(P, T)}{\eta_c} = \frac{\eta(P, T) T_c^{1/6}}{P_c^{2/3} M^{1/2}} \quad (6.4)$$

For one component as a reference component if the function, f , in Equation 6.2 is known, it is possible to calculate the viscosity of any other components, such as component x, at any pressure and temperature. Thus,

$$\eta_x(P, T) = \frac{\left(\frac{P_{cx}}{P_{c0}}\right)^{\frac{2}{3}} \left(\frac{M_x}{M_0}\right)^{\frac{1}{2}}}{\left(\frac{T_{cx}}{T_{c0}}\right)^{\frac{1}{6}}} \eta_0 \left(\frac{P_{c0}}{P_{cx}}, \frac{T_{c0}}{T_{cx}} \right) \quad (6.5)$$

Where, 0 refers to the reference component methane with the viscosity data published by Hanley et al. [32] has been selected as the reference fluid in the original Pedersen Model. In this work, carbon dioxide with the viscosity data published by Fenghour et al. [6] has been selected as the reference fluid for the CO₂ systems including impurities. The viscosity of CO₂ as a function of density and temperature can be calculated from the following equation:

$$\eta(\rho, T) = \eta_0(T) + \Delta\eta(\rho, T) \quad (6.6)$$

Where, $\eta_0(T)$ is the zero-density viscosity and can be obtained from the following equation:

$$\eta_0(T) = \frac{1.00697 T^{1/2}}{\Psi_{\eta}^*(T^*)} \quad (6.7)$$

In this equation, the zero-density viscosity is in units of $\mu\text{Pa}\cdot\text{s}$ and temperature, T , in K. The reduced effective cross section, $\Psi_{\eta}^*(T^*)$, is represented by the empirical equation.

$$\ln\Psi_{\eta}^*(T^*) = \sum_{i=0}^4 a_i (\ln T^*)^i \quad (6.8)$$

Where the reduced temperature, T^* , is given by

$$T^* = kT/\varepsilon \quad (6.9)$$

And the energy scaling parameter, $\frac{\varepsilon}{k} = 251.196$ K. Where k is the Boltzmann's constant and ε is the minimum of the pair-potential energy. The coefficients, a_i , are listed in the

Table 6.4. The second contribution in Equation 6.6 is the excess viscosity, $\Delta\eta(\rho, T)$, which describes how the viscosity can change as a function of density outside of the critical region. The excess viscosity correlation can be correlated as follows:

$$\Delta\eta(\rho, T) = d_{11}\rho + d_{21}\rho^2 + \frac{d_{64}\rho^6}{T^{*3}} + d_{81}\rho^8 + \frac{d_{82}\rho^8}{T^*} \quad (6.10)$$

Where, the temperature is in Kelvin, the density in kg/m^3 and the excess viscosity in $\mu\text{Pa}\cdot\text{s}$. The coefficients are shown in Table 6.5.

Table 6.4 Values of coefficients, a_i , in Equation 6.8 for CO_2

i	a_i
0	0.235156
1	-0.491266
2	5.211155×10^{-2}
3	5.347906×10^{-2}
4	-1.537102×10^{-2}

Table 6.5 Values of coefficients, d_{ij} , in Equation 6.10

d_{ij}	Value
d_{11}	0.4071119×10^{-2}
d_{21}	0.7198037×10^{-4}
d_{64}	$0.2411697 \times 10^{-16}$
d_{81}	$0.2971072 \times 10^{-22}$
d_{82}	$-0.162788 \times 10^{-22}$

The corresponding states principle expressed in Equation 6.5 for the viscosity of pure components works well for mixtures. Pedersen et al. [4] have expressed the following expression to calculate the viscosity of mixtures at any pressure and temperature.

$$\eta_{\text{mix}}(P, T) = \frac{\left(\frac{P_{\text{cmix}}}{P_{c0}}\right)^{\frac{2}{3}} \left(\frac{M_{\text{mix}}}{M_0}\right)^{\frac{1}{2}} \left(\frac{\alpha_{\text{mix}}}{\alpha_0}\right)}{\left(\frac{T_{\text{cmix}}}{T_{c0}}\right)^{\frac{1}{6}}} \eta_0(P_0, T_0) \quad (6.11)$$

Where

$$P_0 = \frac{P P_{c0}}{P_{cmix}} \frac{\alpha_0}{\alpha_{mix}}; \quad T_0 = \frac{T T_{c0}}{T_{cmix}} \frac{\alpha_0}{\alpha_{mix}} \quad (6.12)$$

The critical temperature and pressure for mixtures, according to recommended mixing rules by Murad and Gubbins [33], can be found from

$$T_{cmix} = \frac{\sum_{i=1}^N \sum_{j=1}^N Z_i Z_j \left[\left(\frac{T_{ci}}{P_{ci}} \right)^{\frac{1}{3}} + \left(\frac{T_{cj}}{P_{cj}} \right)^{\frac{1}{3}} \right]^3 \sqrt{T_{ci} T_{cj}}}{\sum_{i=1}^N \sum_{j=1}^N Z_i Z_j \left[\left(\frac{T_{ci}}{P_{ci}} \right)^{\frac{1}{3}} + \left(\frac{T_{cj}}{P_{cj}} \right)^{\frac{1}{3}} \right]^3} \quad (6.13)$$

$$P_{cmix} = \frac{8 \sum_{i=1}^N \sum_{j=1}^N Z_i Z_j \left[\left(\frac{T_{ci}}{P_{ci}} \right)^{\frac{1}{3}} + \left(\frac{T_{cj}}{P_{cj}} \right)^{\frac{1}{3}} \right]^3 \sqrt{T_{ci} T_{cj}}}{\left(\sum_{i=1}^N \sum_{j=1}^N Z_i Z_j \left[\left(\frac{T_{ci}}{P_{ci}} \right)^{\frac{1}{3}} + \left(\frac{T_{cj}}{P_{cj}} \right)^{\frac{1}{3}} \right]^3 \right)^2} \quad (6.14)$$

The mixture molecular weight is found from

$$M_{mix} = 1.304 \times 10^{-4} \left(\bar{M}_w^{2.303} - \bar{M}_n^{2.303} \right) + \bar{M}_n \quad (6.15)$$

where \bar{M}_w and \bar{M}_n are the weight average and number average molecular weights, respectively.

$$\bar{M}_w = \frac{\sum_{i=1}^N Z_i M_i^2}{\sum_{i=1}^N Z_i M_i} \quad (6.16)$$

$$\bar{M}_n = \sum_{i=1}^N Z_i M_i \quad (6.17)$$

The parameter α for mixtures in Equation 6.11 can be found from

$$\alpha_{\text{mix}} = 1 + A\rho_r^B M_{\text{mix}}^C \quad (6.18)$$

Table 6.6 The original Pedersen mixture tuning parameters

Parameters	Value
A	7.378×10^{-3}
B	1.847
C	0.5173

The tuning parameters of A, B and C can be found from Table 6.6. α for the reference fluid can be found from the Equation 6.18 by replacing the molecular weight of the mixture with that of the reference fluid, carbon dioxide. The reduced density, ρ_r , defines as

$$\rho_r = \frac{\rho_0 \left(\frac{PP_{\text{CO}_2}}{P_{\text{cmix}}}, \frac{T T_{\text{CO}_2}}{T_{\text{cmix}}} \right)}{\rho_{\text{CO}_2}} \quad (6.19)$$

The critical density of carbon dioxide, ρ_{CO_2} , is equal to 467.69 kg/m³. The Modified Benedict–Webb–Rubin (MBWR) equation of state is used for computing the reference fluid density, ρ_0 , at the desired pressure and temperature of $\frac{PP_{\text{CO}_2}}{P_{\text{cmix}}}, \frac{T T_{\text{CO}_2}}{T_{\text{cmix}}}$. The mathematical equation of the MBWR has been presented by Younglove et al. [4]. The procedure below should be followed to calculate the viscosity of CO₂ systems with impurities by the corresponding state principle of Pedersen Model:

1. Calculate the T_{cmix} , P_{cmix} and M_{mix} from Equations 6.13, 6.14 and 6.15 respectively.
2. Obtain the CO₂ density at $\frac{PP_{\text{CO}_2}}{P_{\text{cmix}}}, \frac{T T_{\text{CO}_2}}{T_{\text{cmix}}}$ from the MBWR EOS and calculate the reduced density from Equation 6.19
3. The mixture parameter, α_{mix} , and α_0 should be calculated from Equation 6.18

4. The reference pressure and temperature, P_0 and T_0 , should be calculated from [Equation 6.12](#)
5. Calculate the CO₂ reference fluid, $\eta_0(P_0, T_0)$, in [Equation 6.11](#) from [Equation 6.6](#)
6. Calculate the mixture viscosity from [Equation 6.11](#)

6.3.2 Lohrenz–Bray–Clark (LBC) model

The LBC correlation is a fourth-degree polynomial equation in the reduced density, $\rho_r = \frac{\rho}{\rho_c}$, as presented in 1964. Lohrenz et al. [2] extended the JST method [9] for calculating the viscosity of mixtures of naturally occurring hydrocarbons. The equation is:

$$[(\eta - \eta^*)\xi + 10^{-4}]^{1/4} = a_1 + a_2 \rho_r + a_3 \rho_r^2 + a_4 \rho_r^3 + a_5 \rho_r^4 \quad (6.20)$$

where the constants a_1 to a_5 are listed in [Table 6.7](#). These parameters can be tuned to match to the experimental data of CO₂ systems with impurities. In [Equation 6.20](#), ξ is the viscosity reducing parameter, which for a mixture is given by

$$\xi = \frac{[\sum_{i=1}^N Z_i T_{ci}]^{1/6}}{[\sum_{i=1}^N Z_i M_i]^{1/2} [\sum_{i=1}^N Z_i P_{ci}]^{2/3}} \quad (6.21)$$

Table 6.7 Parameters in the original LBC Viscosity Correlation

LBC Parameters	Original Constants	New Constants
a_1	0.10230	0.07960
a_2	0.023364	0.142735
a_3	0.058533	-0.113846
a_4	-0.040758	0.056084
a_5	0.0093324	-0.0090435

The critical density of mixtures can be determined from the critical molar volume.

$$\rho_c = \frac{1}{V_c} = \frac{1}{\sum_{i=1}^N Z_i V_{ci}} \quad (6.22)$$

The density of mixture for CO₂ systems with impurities can be calculated by Peng-Robinson equation of state. In this work, the density of the modified PR EOS demonstrated in [Chapter 4](#) was used to calculate the mixture density. Therefore, a new set of parameters were obtained by minimizing the error with the experimental data of viscosity. The new parameters are listed together with the original ones in [Table 6.7](#). Low pressure mixture viscosity, η^* , can be determined from the following equation presented by Herning et al. [\[27\]](#).

$$\eta^* = \frac{\sum_{i=1}^N Z_i \eta_i^* \sqrt{M_i}}{\sum_{i=1}^N Z_i \sqrt{M_i}} \quad (6.23)$$

The dilute component viscosity, η_i^* , for each component has been expressed as a function of reduced temperature, T_r , by Stiel and Thodos [\[35\]](#) as follows:

$$\eta_i^* = 34 \times 10^{-5} \frac{1}{\xi_i} T_{ri}^{0.94} ; \quad \text{for } T_{ri} < 1.5 \quad (6.24)$$

$$\eta_i^* = 17.78 \times 10^{-5} \frac{1}{\xi_i} (4.58 T_{ri} - 1.67)^{5/8} ; \quad \text{for } T_{ri} > 1.5 \quad (6.25)$$

where,

$$\xi_i = \frac{T_{ci}^{1/6}}{M_i^{1/2} P_{ci}^{2/3}} \quad (6.26)$$

6.4 Results and Discussion

6.4.1 Experimental Results

The viscosity of six binary CO₂ mixtures with 5% impurity (H₂, CH₄, N₂, O₂, CO and Ar) plus a multi-components mixture (MIX1) with 4.36% impurities are presented in this work. The measurements were performed at temperatures of 280, 288.15, 300.15, 308.15, 323.15 and 343.15 K. The experimental and modelling results for the viscosity of CO₂ and CO₂/impurities systems are given in [Table 6.8](#) to [Table 6.13](#) and [Figure 6.2](#) to [Figure 6.8](#). The experiments for each binary mixture were conducted in the single liquid phase region (above saturation) or in the supercritical region. In order to validate our experimental procedure and setup, the viscosity of CO₂ at 288 K obtained from Fenghour et al. [6] and at 323.15 K and 343.15 K from Pensado et al. [36] were compared to our experimental results. [Figure 6.2](#) shows that our results are in a good agreement with these literature sources. The viscosity of carbon dioxide is a function of pressure and temperature; it increases as pressure increases and decreases as temperature increases. The viscosity of impure CO₂ was tested under the same condition of pressure and temperature of pure CO₂. The viscosity of liquid and supercritical CO₂ systems show an increase as pressure increases and the change becomes less significant as the system moves away from the saturation line. At constant pressure, the increase of temperature tends to reduce the viscosity of CO₂/impurities systems. As the system approaches the critical pressure and temperature, the viscosity change becomes significant. For example at the same pressure condition, 5% N₂ can reduce the viscosity of CO₂ up to 35% and 13% near the critical temperature (300.15 K) and at 288.15 K respectively.

The viscosity of CO₂ is highly reduced in the presence of H₂ compare to the other impurities tested. This is mainly due to the lower molecular weight of hydrogen, i.e. lower CO₂ mixture density. The other CO₂ systems are found to have a similar effect on the viscosity. [Figure 6.2](#) to [Figure 6.8](#) demonstrate the effect of Ar and H₂ on CO₂ viscosity at 280, 288.15, 300.15, 308.15, 323.15 and 343.15 K. The presence of H₂ and Ar strongly reduce the liquid CO₂ viscosity. The viscosity of CO₂/H₂ mixture is around 20% lower than those of pure CO₂ near the saturation at 288.15 K while the viscosity of CO₂/Ar mixture is lower by 10% at the same condition. In the supercritical region, the viscosity is much lower than in the liquid region of both pure and impure CO₂ systems.

For example, at 323.15 K, the viscosity of pure CO₂ drops as low as 30% and 22% in the presence of H₂ and Ar respectively. This reduction occurs at a pressure around 11 MPa for CO₂/Ar and CO₂/H₂ systems where their critical pressures are 7.96 and 8.93 MPa respectively (VPT EOS model). Therefore, the critical change in CO₂ viscosity seems to take place by 2 to 3 MPa above the supercritical boundary of the impurities system.

Table 6.8 Experimental and modelling results of the 95% CO₂ + 5% CH₄ binary system

Temp. K (±0.01)	Pressure MPa (±0.05)	Exp. Visc. μPa.s (±1%)	Modified- Pedersen Dev%	Modified- LBC Dev%
280	6.50	80.8	1.3	3.1
280	10.19	91.1	0.1	0.7
280	19.07	107.5	0.1	7.2
280	23.21	113.9	0.1	8.7
280	28.21	121.1	0.1	9.9
280	32.59	127.0	0.1	10.5
280	42.03	139.0	0.1	10.6
280	45.49	143.2	0.1	10.3
280	48.72	147.0	0.1	10.0
288.15	6.81	66.5	2.3	0.2
288.15	9.01	74.5	1.4	1.9
288.15	12.88	84.5	0.5	1.6
288.15	22.67	102.6	0.5	0.2
288.15	26.92	108.9	0.5	0.4
288.15	31.21	114.7	0.4	0.5
288.15	36.31	121.3	0.5	0.6
288.15	41.83	128.1	0.5	0.8
288.15	48.19	135.5	0.5	1.2
300.15	8.44	49.5	3.0	1.1
300.15	10.79	61.9	1.7	4.1
300.15	13.89	71.9	0.2	3.2
300.15	18.93	82.1	0.2	2.5
300.15	25.33	92.6	0.4	1.4

300.15	29.01	98.3	0.9	0.7
300.15	34.43	105.3	0.8	0.8
300.15	39.13	110.9	0.7	1.0
300.15	45.75	118.4	0.5	1.2
308.15	9.16	35.8	2.5	5.9
308.15	12.50	57.3	1.0	4.6
308.15	16.58	68.8	0.0	3.5
308.15	19.69	75.6	0.8	2.4
308.15	24.66	84.3	1.4	2.1
308.15	36.70	100.8	1.6	2.2
308.15	36.93	101.0	1.5	2.2
308.15	42.37	108.0	2.1	1.6
308.15	46.50	112.8	2.2	1.2
323.15	10.71	30.0	6.6	1.0
323.15	14.95	48.9	0.4	0.5
323.15	20.13	62.2	0.1	1.0
323.15	27.95	76.1	1.3	1.2
323.15	32.46	81.8	0.9	2.4
323.15	37.44	88.1	1.1	2.7
323.15	43.62	95.1	1.1	3.1
323.15	48.66	100.1	0.8	3.5
343.15	50.10	88.8	2.3	1.7
343.15	45.03	83.4	2.0	1.5
343.15	34.46	71.5	1.8	0.1
343.15	31.60	67.7	1.6	0.5
343.15	28.10	62.4	0.9	0.7
343.15	23.90	55.7	0.8	1.6
343.15	19.62	47.0	0.7	1.6
343.15	13.95	31.1	1.8	2.9
343.15	9.24	21.0	0.9	2.0

Table 6.9 Experimental and modelling results of the 95% CO₂ + 5% Ar binary system

Temp. K (±0.01)	Pressure MPa (±0.05)	Exp.Visc. μPa.s (±1%)	Modified- Pedersen Dev%	Modified- LBC Dev%
280	7.75	87.1	2.9	3.6
280	10.34	93.3	2.4	2.9
280	13.86	100.7	2.4	3.0
280	24.14	118.0	2.1	3.5
280	28.68	124.9	2.3	3.9
280	33.41	131.3	2.1	3.9
280	41.75	142.3	2.2	4.5
280	47.90	150.4	2.4	5.4
288.15	8.80	74.5	0.7	0.9
288.15	11.55	82.2	0.8	0.1
288.15	15.01	90.2	1.3	0.7
288.15	19.83	99.9	2.2	2.0
288.15	24.33	107.1	2.1	2.2
288.15	29.97	115.2	2.0	2.2
288.15	34.17	120.9	2.0	2.2
288.15	41.35	130.0	1.9	2.3
288.15	46.46	136.2	1.9	2.6
300.15	9.15	54.9	2.8	3.0
300.15	13.63	73.0	3.5	0.7
300.15	17.03	81.6	4.2	1.7
300.15	32.65	105.2	2.5	1.2
300.15	38.15	112.0	2.3	0.9
300.15	45.08	120.9	2.7	1.2
300.15	49.34	126.2	3.0	1.6
308.15	8.93	29.0	4.0	1.2
308.15	10.63	48.2	2.5	1.4
308.15	11.82	54.0	0.4	2.4

308.15	21.74	81.5	3.8	1.0
308.15	26.59	89.5	3.8	0.7
308.15	32.10	96.9	3.3	0.0
308.15	37.61	104.2	3.4	0.0
308.15	43.15	110.3	2.9	0.4
308.15	48.03	115.8	2.8	0.3
323.15	12.16	36.7	5.0	0.2
323.15	16.86	55.7	2.9	2.7
323.15	24.06	71.5	3.4	2.0
323.15	26.59	76.0	3.8	2.0
323.15	32.10	83.8	3.6	0.8
323.15	37.61	90.7	3.4	0.0
323.15	43.21	97.4	3.5	0.2
343.15	10.36	23.6	4.7	2.3
343.15	16.00	39.1	7.2	4.6
343.15	22.48	54.0	3.0	4.0
343.15	26.89	62.6	4.1	4.5
343.15	32.03	69.4	3.1	2.3
343.15	40.06	80.1	4.1	1.8
343.15	43.15	82.9	3.3	0.6
343.15	50.39	90.2	3.1	0.5

Table 6.10 Experimental and modelling results of the 95% CO₂ + 5% CO binary system

Temp. K (±0.01)	Pressure MPa (±0.05)	Exp.Visc. μPa.s (±1%)	Modified- Pedersen Dev%	Modified- LBC Dev%
280	8.90	85.4	2.4	0.9
280	10.07	88.3	2.4	0.8
280	14.39	97.7	2.5	1.3
280	18.45	104.9	2.3	1.5
280	25.36	115.4	1.9	1.4
280	29.39	121.8	2.4	2.1

280	33.72	127.5	2.2	2.1
280	38.23	133.8	2.5	2.7
280	45.57	142.6	2.2	3.1
288.15	9.24	72.4	1.5	0.4
288.15	10.79	76.6	1.0	1.3
288.15	15.21	88.4	3.0	0.7
288.15	18.87	95.0	2.7	0.7
288.15	24.06	103.2	2.4	0.7
288.15	28.34	109.3	2.2	0.6
288.15	34.39	117.2	1.9	0.4
288.15	39.21	123.1	1.7	0.3
288.15	47.83	133.8	2.0	1.0
300.15	9.96	54.6	0.7	2.2
300.15	13.49	67.7	1.1	3.1
300.15	18.30	79.2	2.1	1.6
300.15	20.78	83.8	2.2	1.2
300.15	26.66	93.7	2.8	0.1
300.15	29.93	98.2	2.6	0.2
300.15	36.52	106.5	2.2	0.7
300.15	41.60	112.5	2.1	0.9
300.15	49.97	122.5	2.4	0.6
300.15	53.79	126.4	2.1	0.7
308.15	9.17	28.8	1.7	6.7
308.15	14.32	61.7	3.1	1.1
308.15	18.89	72.7	3.1	0.8
308.15	21.59	77.8	3.1	0.9
308.15	26.97	86.5	3.1	1.2
308.15	32.91	94.8	3.1	1.5
308.15	36.69	99.9	3.4	1.4
308.15	45.75	110.8	3.5	1.0
308.15	50.19	116.4	4.1	0.2
323.15	8.94	22.2	3.9	2.2
323.15	13.71	42.9	4.6	1.5
323.15	18.59	58.5	3.7	2.5

323.15	22.95	67.5	3.7	1.8
323.15	26.17	73.0	3.9	1.3
323.15	31.43	80.7	3.9	0.4
323.15	35.85	86.7	4.2	0.0
323.15	40.21	91.7	3.9	0.7
323.15	46.35	98.2	3.5	1.4
343.15	9.19	21.8	4.4	6.1
343.15	13.26	28.8	3.6	2.4
343.15	19.21	44.9	2.4	2.2
343.15	23.07	54.1	4.0	4.5
343.15	26.28	59.6	3.7	3.8
343.15	31.04	66.5	3.5	2.5
343.15	36.01	72.7	3.3	1.1
343.15	40.70	77.9	3.0	0.0
343.15	45.52	82.8	2.8	1.1
343.15	49.35	86.5	2.6	1.7

Table 6.11 Experimental and modelling results for the 95% CO₂ + 5% O₂ binary system

Temp. K (±0.01)	Pressure MPa (±0.05)	Exp. Visc. μPa.s (±1%)	Modified- Pedersen Dev%	Modified- LBC Dev%
280	9.10	86.6	1.7	1.0
280	12.72	94.7	1.3	0.6
280	17.94	104.4	1.0	0.2
280	22.75	113.1	0.0	1.5
280	26.20	118.2	0.0	1.6
280	30.53	124.3	0.0	1.8
280	35.82	131.3	0.0	2.1
280	40.10	136.8	0.0	2.5
280	44.84	142.6	0.0	2.9
280	48.29	146.7	0.1	3.2
288.15	9.39	76.0	0.3	0.2
288.15	11.92	82.7	0.4	0.2

288.15	17.86	94.6	0.4	0.2
288.15	21.30	100.4	0.5	0.5
288.15	25.17	106.3	0.4	0.7
288.15	29.56	112.6	0.5	0.8
288.15	34.56	119.2	0.4	0.8
288.15	39.84	125.8	0.4	0.9
288.15	44.13	131.0	0.4	1.1
288.15	47.77	135.2	0.3	1.3
300.15	8.75	49.7	1.5	0.4
300.15	11.53	64.9	0.8	1.3
300.15	16.92	78.6	1.0	1.4
300.15	21.66	87.2	0.9	0.7
300.15	24.21	91.3	1.0	0.5
300.15	29.65	99.1	0.8	0.3
300.15	35.19	106.4	0.8	0.4
300.15	39.05	111.1	0.8	0.5
300.15	42.92	115.7	0.8	0.5
300.15	47.10	120.4	0.7	0.6
308.15	10.60	47.3	0.7	0.0
308.15	13.35	60.4	1.3	1.7
308.15	16.52	69.1	1.4	1.3
308.15	20.70	77.5	1.4	1.3
308.15	25.72	85.7	1.3	1.6
308.15	28.35	89.5	1.2	1.8
308.15	32.79	95.5	1.2	2.1
308.15	38.12	102.1	1.1	2.2
308.15	41.62	106.2	1.1	2.2
308.15	47.47	112.7	0.9	2.0
323.15	9.15	22.3	1.3	5.8
323.15	14.63	46.3	1.2	1.5
323.15	18.17	58.6	2.0	1.9
323.15	19.92	62.6	2.0	1.6
323.15	26.72	74.6	2.0	0.2
323.15	28.46	77.1	1.9	0.2

323.15	32.56	82.7	1.8	0.8
323.15	37.17	88.3	1.6	1.6
323.15	41.61	93.4	1.5	2.0
323.15	44.97	97.1	1.5	2.2
343.15	8.96	21.0	1.3	3.8
343.15	12.66	27.0	1.0	5.0
343.15	16.43	37.3	0.6	2.4
343.15	20.62	48.1	0.7	0.5
343.15	24.98	57.2	0.7	1.7
343.15	29.59	65.3	2.3	2.3
343.15	33.68	70.7	2.1	1.3
343.15	36.21	74.0	2.4	1.0
343.15	41.37	79.8	2.2	0.1
343.15	47.28	85.9	2.1	1.0

Table 6.12 Experimental and modelling results of the 95% CO₂ + 5% H₂ binary system

Temp. K (±0.01)	Pressure MPa (±0.05)	Exp. Visc. μPa.s (±1%)	Modified- Pedersen Dev%	Modified- LBC Dev%
280	8.73	86.0	16.3	10.7
280	13.10	93.5	11.3	6.4
280	18.47	102.3	9.1	5.0
280	20.79	106.0	8.7	4.8
280	24.75	112.0	8.3	4.7
280	28.35	117.2	8.0	4.6
280	33.00	123.5	7.7	4.6
280	37.00	128.7	7.6	4.7
280	40.53	133.0	7.3	4.7
280	48.26	141.7	6.6	4.9
288.15	12.57	77.6	8.8	2.4
288.15	17.20	87.4	7.7	2.3
288.15	21.75	96.1	7.9	3.1
288.15	26.85	104.8	8.3	3.9

288.15	29.37	108.7	8.4	4.1
288.15	33.20	114.1	8.4	4.1
288.15	36.78	118.6	8.1	3.9
288.15	41.04	123.2	7.4	3.4
300.15	14.17	65.9	8.6	0.4
300.15	16.86	72.2	7.5	0.3
300.15	19.80	78.6	7.6	1.2
300.15	22.38	83.8	8.1	2.2
300.15	26.23	90.8	8.7	3.3
300.15	31.57	99.2	9.1	3.9
300.15	41.31	110.5	7.8	2.5
300.15	47.24	118.0	8.2	2.9
308.15	11.13	40.8	12.2	5.6
308.15	14.86	58.5	8.7	0.1
308.15	19.86	70.4	7.3	0.5
308.15	23.66	77.8	7.6	1.0
308.15	30.12	88.4	8.2	1.4
308.15	34.39	94.2	8.1	1.2
308.15	38.42	99.6	8.3	1.4
308.15	45.41	107.7	8.0	1.3
323.15	9.31	23.0	7.8	1.4
323.15	11.96	31.0	10.0	3.5
323.15	15.14	44.2	9.2	1.3
323.15	19.82	57.5	7.7	3.0
323.15	23.70	65.6	7.9	3.2
323.15	28.88	74.1	7.9	2.7
323.15	34.46	82.0	8.0	2.2
323.15	38.00	86.5	8.1	1.9
323.15	43.78	93.2	8.0	1.4
343.15	11.32	23.9	4.2	0.6
343.15	14.15	29.7	6.4	1.3
343.15	18.33	40.9	8.1	3.5
343.15	23.03	51.1	7.3	4.9
343.15	31.98	65.3	7.0	4.0

343.15	36.19	71.1	7.6	3.8
343.15	40.77	77.2	8.5	4.0
343.15	44.73	82.1	9.1	4.1

Table 6.13 Experimental and modelling results for MIX1

Temp. K (± 0.01)	Pressure MPa (± 0.05)	Exp. Visc. $\mu\text{Pa}\cdot\text{s}$ ($\pm 1\%$)	Modified- Pedersen Dev%	Modified- LBC Dev%
280	13.14	94.0	0.6	1.8
280	15.70	99.2	0.2	1.0
280	20.68	106.8	0.9	1.1
280	25.95	115.3	0.3	0.2
280	32.19	123.5	0.6	0.1
280	38.44	131.2	0.9	0.1
280	45.41	139.8	0.8	0.9
288.15	10.68	80.4	4.0	1.2
288.15	12.89	85.1	3.1	0.4
288.15	17.37	93.6	2.4	0.4
288.15	22.23	101.5	2.0	0.6
288.15	27.47	108.3	1.1	0.2
288.15	32.12	115.1	1.4	0.4
288.15	34.62	117.8	0.9	0.1
288.15	40.47	125.4	1.1	0.4
288.15	47.23	131.7	0.2	0.5
300.15	11.07	59.9	1.6	6.6
300.15	13.81	68.6	0.8	6.0
300.15	18.47	79.7	0.7	3.3
300.15	21.01	84.1	0.6	2.9
300.15	25.92	92.4	1.2	1.6
300.15	28.37	95.8	1.1	1.6
300.15	33.81	104.2	2.1	0.4
300.15	39.09	110.0	1.4	1.2
300.15	45.50	117.2	1.1	1.5

308.15	12.26	55.3	2.3	3.7
308.15	14.95	64.8	3.1	2.0
308.15	18.52	72.1	1.7	2.8
308.15	20.12	74.8	1.2	3.3
308.15	23.75	82.3	2.7	1.7
308.15	27.23	88.2	3.3	1.1
308.15	31.63	93.3	2.1	2.5
308.15	33.88	96.6	2.5	2.1
308.15	40.67	104.1	1.7	2.8
308.15	45.16	110.1	2.4	1.8
308.15	50.66	115.2	1.5	2.3
323.15	12.24	35.9	4.1	3.2
323.15	14.79	45.8	1.0	3.8
323.15	17.60	56.2	2.4	0.4
323.15	21.49	63.6	0.6	1.8
323.15	29.00	76.3	1.5	2.1
323.15	32.55	81.6	2.1	1.9
323.15	36.06	86.4	2.5	1.8
323.15	44.19	95.7	2.4	2.4
323.15	51.21	103.5	2.7	2.1
343.15	50.32	88.1	2.4	2.0
343.15	42.52	80.2	2.5	1.1
343.15	34.15	70.6	2.6	0.4
343.15	30.22	65.6	2.8	1.5
343.15	26.03	58.8	2.0	1.5
343.15	22.49	52.4	1.8	1.7
343.15	18.47	43.6	2.5	1.3
343.15	14.26	31.0	1.2	4.8
343.15	12.33	26.6	3.2	3.2

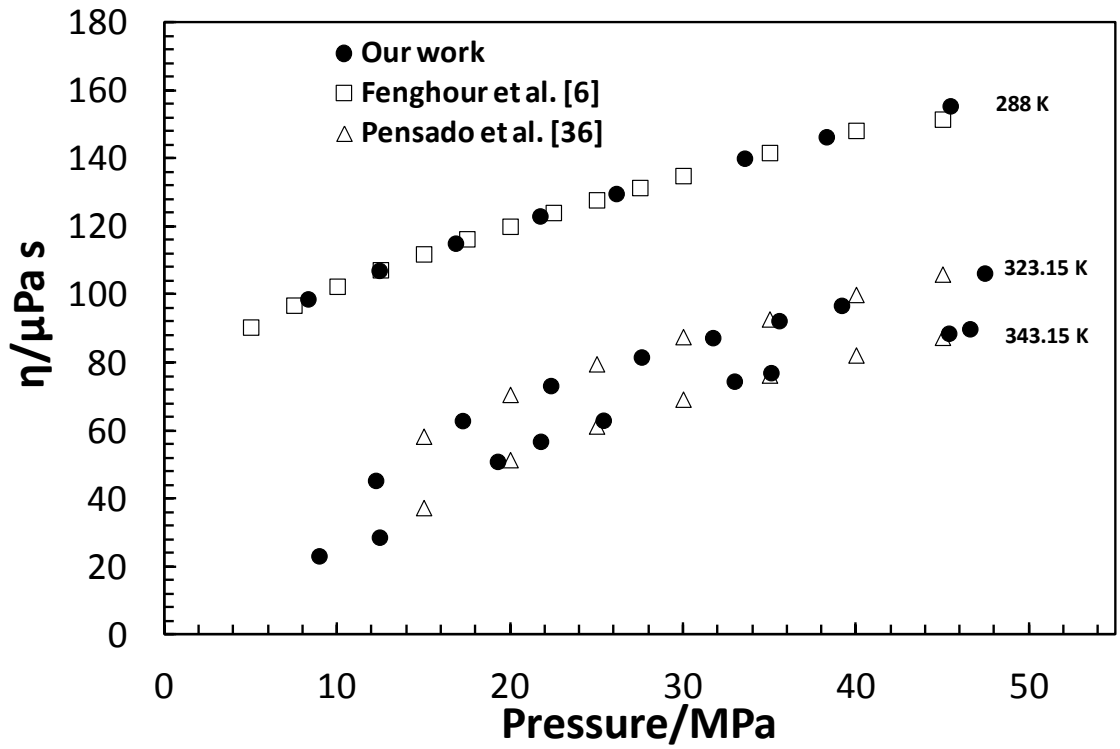


Figure 6.2 Our viscosity data of pure CO_2 together with some literature data at three isothermal conditions

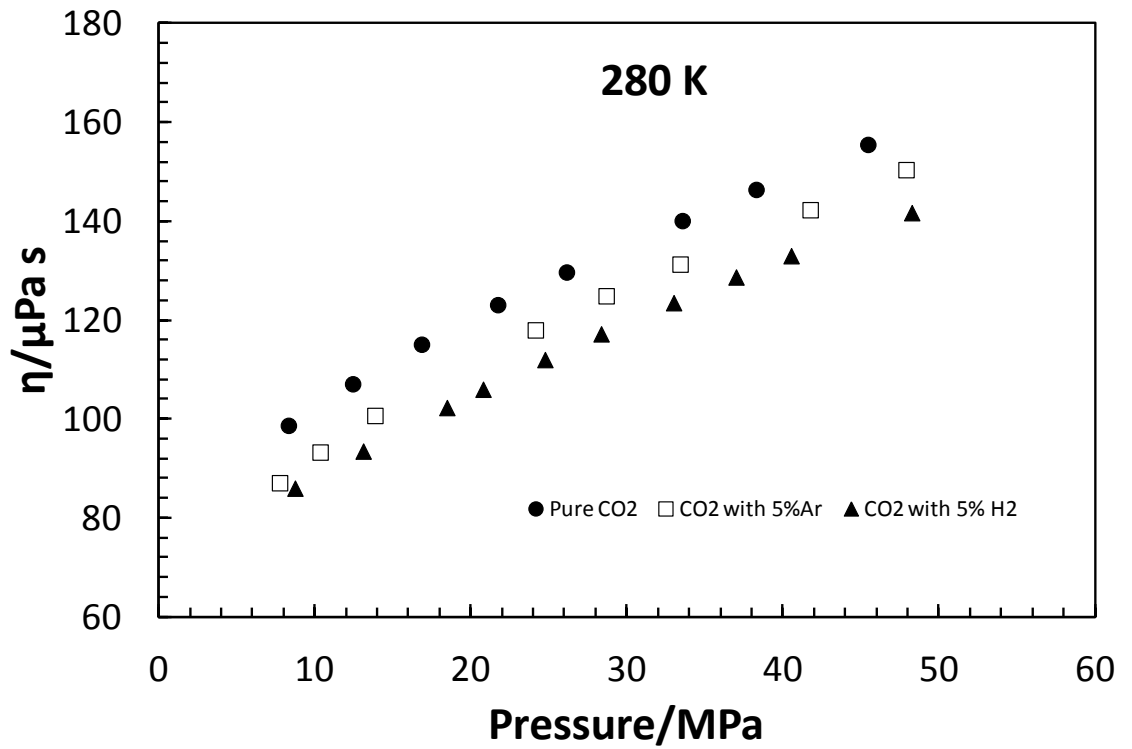


Figure 6.3 The effect of Ar and H_2 on CO_2 viscosity at 280 K

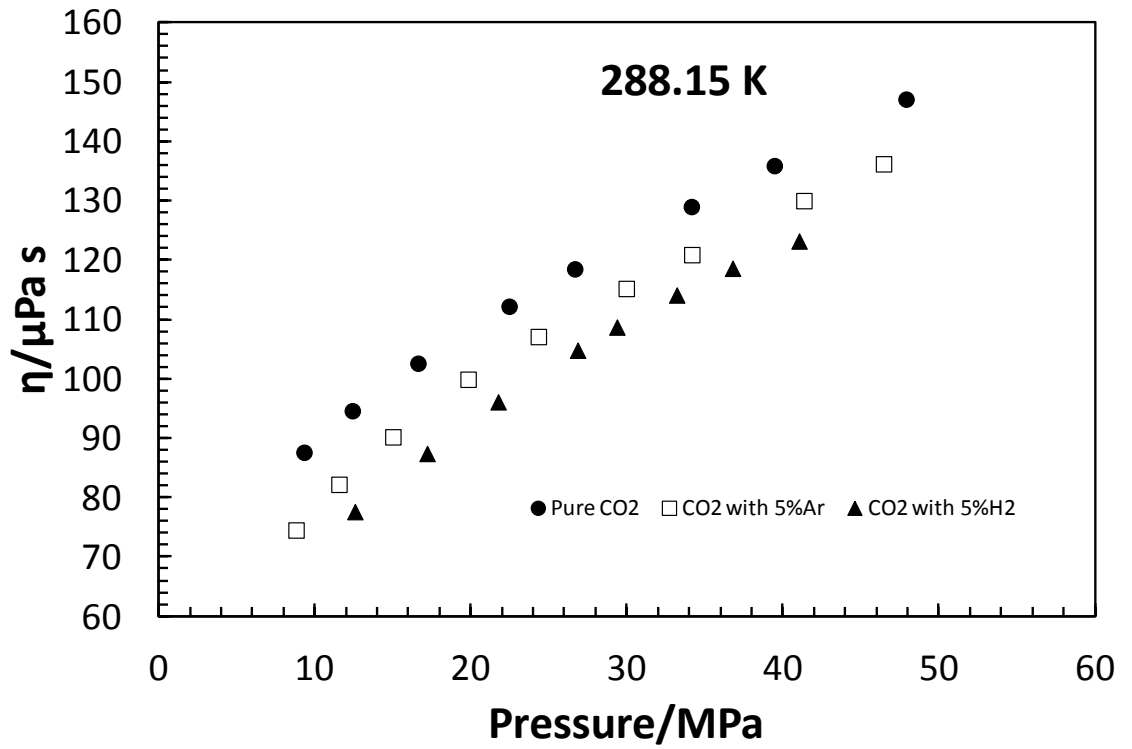


Figure 6.4 The effect of Ar and H₂ on CO₂ viscosity at 288.15 K

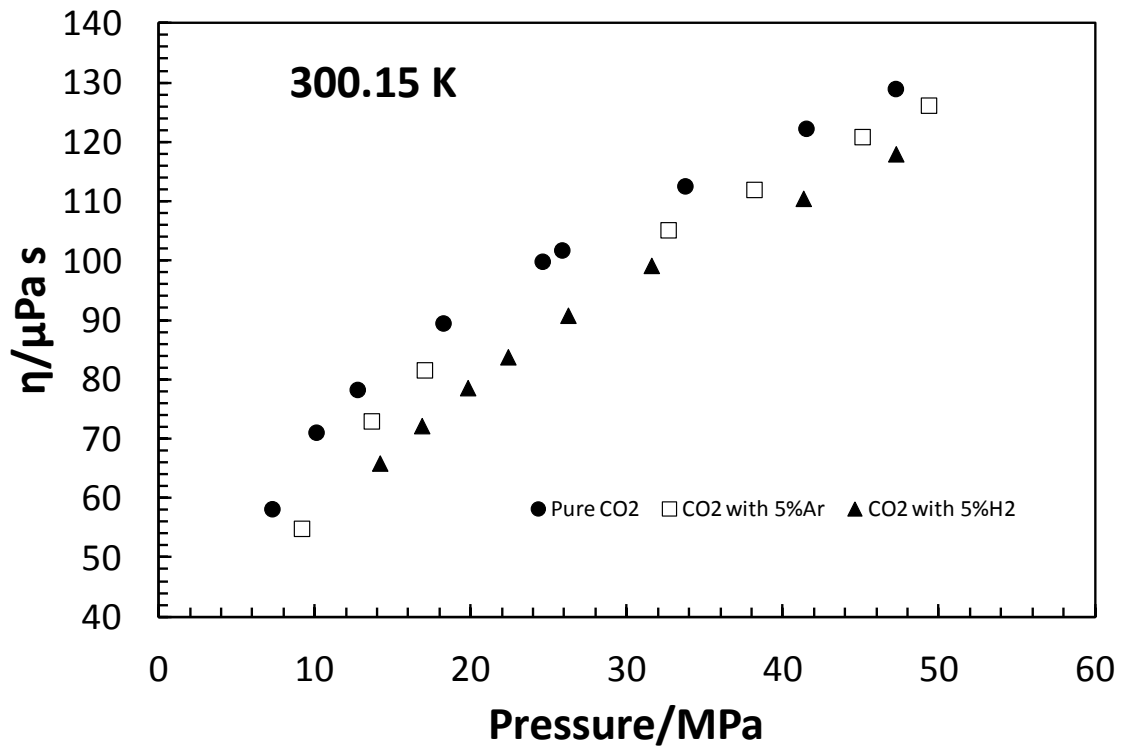


Figure 6.5 The effect of Ar and H₂ on CO₂ viscosity at 300.15 K

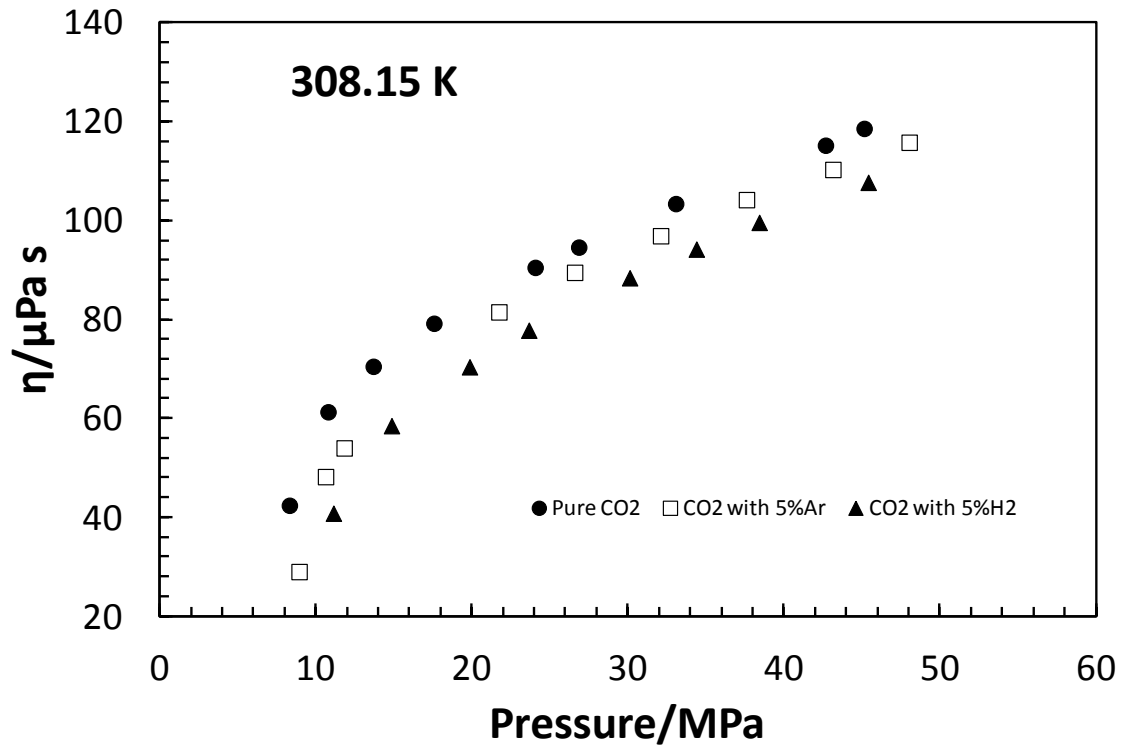


Figure 6.6 The effect of Ar and H_2 on CO_2 viscosity at 308.15 K

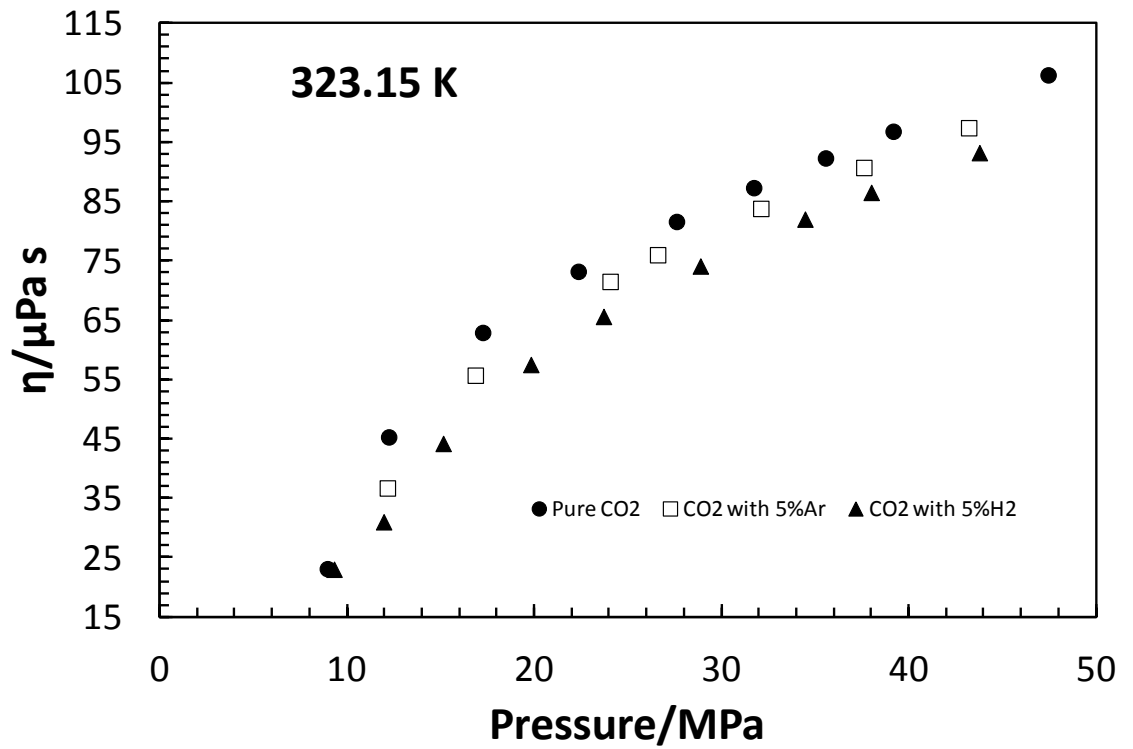


Figure 6.7 The effect of Ar and H_2 on CO_2 viscosity at 323.15 K

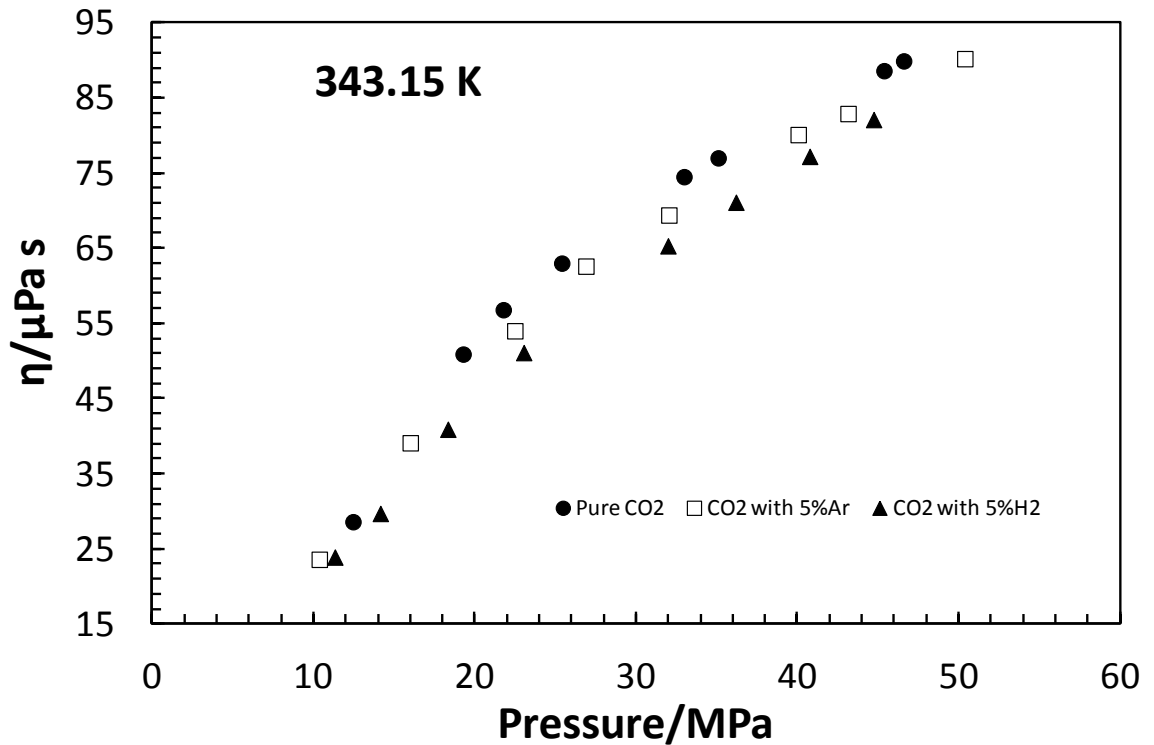


Figure 6.8 The effect of Ar and H₂ on CO₂ viscosity at 343.15 K

6.4.2 Modelling Results

The procedures described in Section 6.3 for the LBC and CO₂-Pedersen models were used to predict the viscosity of a number of binary and multi-component mixtures (CO₂ and impurities) over wide ranges of temperature and pressure. The capability of the models was investigated in predicting the viscosity of pure CO₂ as shown in Figure 6.9 and Figure 6.10. The CO₂-Pedersen model gives accurate results with less than 1.3% maximum absolute deviation. In the other hand, the modified LBC is over predicting at low pressure near the saturated and supercritical boundaries. At low temperature and high pressure measurements, the model is under predicting and the deviation can reach 7%. The viscosity of each conducted test was calculated using these two models. The results are also shown in Table 6.8 to Table 6.13. The average absolute deviation associated with the models for each system can be found in Table 6.14. From the table, the CO₂-Pedersen model predicts the viscosity of pure CO₂, CO₂/CH₄ and CO₂/O₂ systems with greater accuracy than the LBC model. However, the overall predictions of the mixture viscosity shows that the modified LBC has lower deviation compare to the modified Pedersen model, i.e. AADs are 2.4% and 2.8% with the LBC and CO₂-Pedersen models respectively. The LBC correlation was tuned to match the

experimental data of CO₂ systems. The average absolute deviation for the LBC correlation before tuning was 9.3%, which reduced to 2.3% after tuning to the experimental data. To compare the original Pedersen model in which methane is the reference fluid, with the CO₂-Pedersen model, the viscosity data for a 0.95%CO₂ + 5% Ar has been compared. The results show that the CO₂-Pedersen model can predict the mixture viscosity with less than 1.9% deviation compare to 9.15% in the original Pedersen model.

Table 6.14 Average Absolute Deviation (AAD) for the different studied systems

Component	CO2-Pedersen AAD%	CO2-LBC AAD%
Pure CO ₂	0.3	3.2
CO ₂ + methane	1	2.8
CO ₂ + nitrogen	4.3	3.1
CO ₂ + argon	2.9	1.9
CO ₂ + carbon monoxide	2.8	1.5
CO ₂ + oxygen	1	1.4
CO ₂ + hydrogen	8.3	3.1
MIX 1	1.8	1.8
AAD%	2.8	2.3

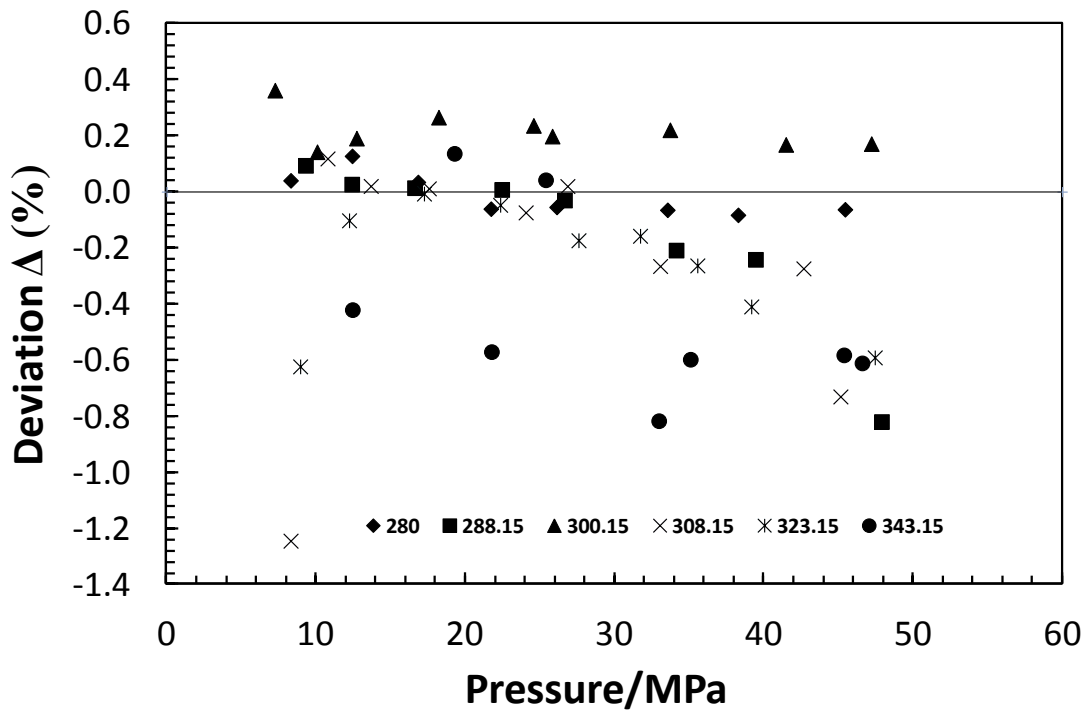


Figure 6.9 Deviation of the CO₂-Pedersen model from the experimental results for pure CO₂

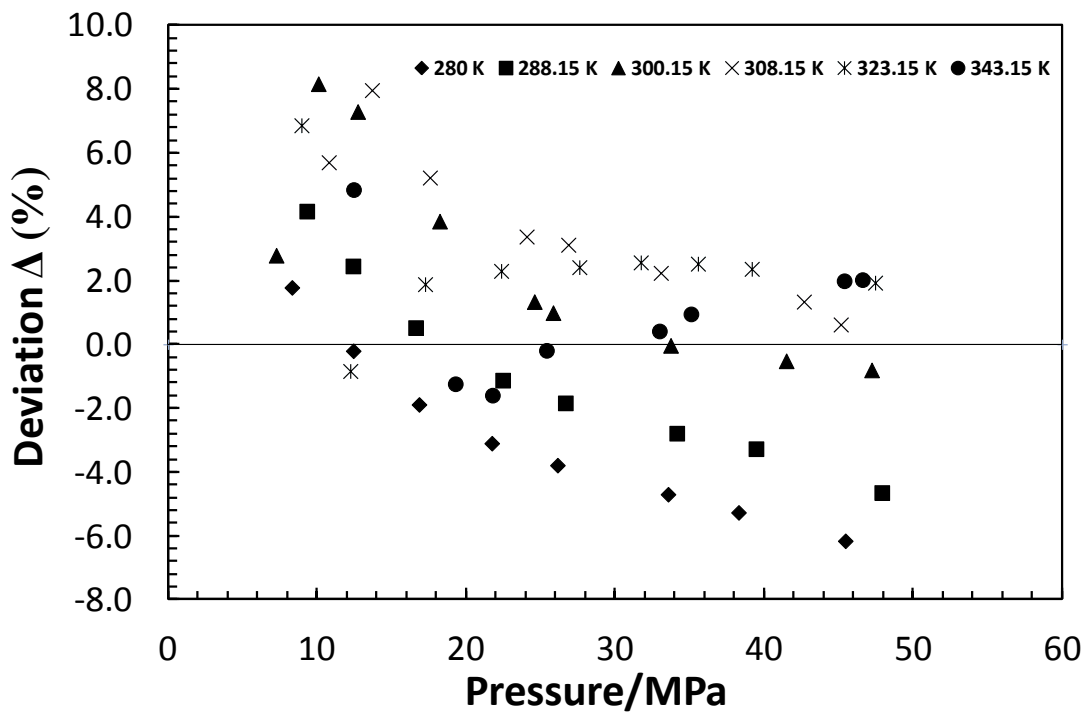


Figure 6.10 Deviation of the LBC model from the experimental results for pure CO₂

6.5 Conclusions

The capillary tube viscosity measurement method has been employed to measure the viscosity of CO₂ systems with impurities such as methane, nitrogen, argon, carbon monoxide, oxygen and hydrogen. The experiments were conducted in an in-house designed and constructed set-up at pressures ranging from 7.4 to 50 MPa and at 280, 288.15, 300.15, 308.15, 323.15 and 343.15 K in the liquid and supercritical region. The results show that all of the tested impurities caused a reduction of viscosity compared to pure CO₂. In addition, two models were modified in order to predict the viscosity of pure and impure CO₂ systems. First, the predictive model of Pedersen was modified by substituting CO₂ as a reference fluid to the original reference fluid of methane. Then, the LBC correlation was tuned to match the obtained experimental data. The modelling results show that the models of CO₂-Pedersen and CO₂-LBC have an average absolute deviation (AAD) of 2.8% and 2.3%, respectively. However, more measurements on high pressure high temperature viscosity should be carried-out in order to compare the capability of the models in wide range of pressure and temperature.

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CHAPTER 7: IFT, SWELLING FACTOR AND MMP OF CO₂/N-DECANE SYSTEM

7.1 Introduction

In order to minimize the emission of CO₂ to the atmosphere, CO₂ can potentially be either sequestered into ground or utilized for enhanced oil recovery (EOR). The aim of CO₂-EOR process is to increase oil recovery by which CO₂ is injected into mature oil fields. Oil is recovered by CO₂ by either miscible or immiscible displacement processes, i.e. above or below the Minimum Miscibility Pressure (MMP) respectively. Above the MMP, the injected CO₂ completely mixes with the reservoir oil and the interfacial tension (IFT) between the CO₂ and the reservoir oil becomes zero. Both the density and viscosity are reduced and therefore the oil can be easily displaced. Below the MMP, where pressure is not enough to push IFT to zero, combination of viscosity reduction, oil swelling, and lowering of interfacial tension is enough to play a role in the CO₂/oil recovery process [1]. By dissolving CO₂ into oil, the IFT reduces [2]. This reduction significantly influences the relative permeability curves i.e. fluids act as single phase and trapping of fluids in pores is impossible [3]. Swelling (extraction) occurs when CO₂ dissolves into reservoir oil. It is a function of three parameters: pressure, temperature and oil composition when pure CO₂ is a solvent [2].

However CO₂ coming from capture processes is generally not pure and can contain impurities such as N₂, H₂, O₂, H₂S, CH₄, CO and water. Major properties in CO₂/oil fluid are influenced by the presence of impurities in the CO₂ stream. For instance, the MMP of CO₂ can increase in the presence of methane and nitrogen and decrease substantially in the presence of hydrogen sulphide and intermediate n-alkanes [4]. Metcalfe et al. [5] conducted slim tube experiments for CO₂ containing H₂S, methane and other hydrocarbons (C₂, C₃). The CO₂ streams containing H₂S and/or LPG components have lower MMPs than those of pure CO₂ systems. Conversely, C₁ in the CO₂ stream raises the MMP. Their findings also stated that, C₂ is as effective in reducing MMP as H₂S and (C₃, C₄) are even more effective. The effects of CO₂/impurities components on CO₂/oil MMP were also found by Shokir [6] in the

following order in terms of their impact: N_2 , C_1 , hydrocarbon components (C_2 - C_4) and H_2S . N_2 and C_1 have a negative impact while the others have a positive impact. Nguyen et al. [7] covered the effect of nitrogen in different impurity range on the properties of CO_2 /oil. They concluded that the presence of nitrogen in the CO_2 stream can potentially reduce diffusivity, swelling factor, solubility and oil recovery. Oil permeability and oil swelling are higher while IFT property of oil is lower in CO_2 injection compared to N_2 injection [4]. Effects on the solubility of CO_2 /crude oil mixtures of adding N_2 contaminant was reported by Monger [8], Nguyen et al. [7] and Spivak and Chima [9]. The authors concluded that N_2 has an adverse effect on solubility and therefore on the immiscible carbon dioxide process mechanisms.

Several correlations were derived for predicting the Minimum Miscibility Pressure (MMP) pure CO_2 and CO_2 containing impurities with reservoir oil and reported in the literature [5, 10-17]. Some of those correlations require few input parameters (molecular weight of light components, MC_5^+ , volatile oil fraction, intermediate oil fraction, °API gravity and oil weight) which are useful when detailed oil characterizations are not available, however they might lead to less accurate in predicting MMP. Computational methods to estimate the MMP for gas-oil systems comprise: Method of the Characteristics (MOC) [18, 19], 1D slim tube simulation [20, 21] and (one or multiple) mixing cell method (MCM) [22-25]. In this work, a modified version of the model proposed by Jaubert et al. [26] (which is based on Metcalfe et al. and Zick [24, 27] approaches) is used.

A wide range of IFT measurements methods are discussed by Rusanov and Prokhorov [28] and Drelich et al. [29]. These methods are capillary rise, pendant drop, drop weight, spinning drop, sessile drop and Wilhelmy plate methods. Two common IFT techniques are available for elevated pressure measurements, pendant drop and capillary rise methods. The later is rarely used in petroleum studies and it is not commercially available. However, it is one of the most and best accurate IFT methods [29, 30]. Ayirala [31] conducted IFT experiments on oil- CO_2 systems using the capillary rise method. He demonstrated that the technique can be easily adapted to high pressures and temperatures and is well suited to measure low interfacial tensions (0.044 dynes/cm).

In this work, an in-house experimental apparatus is adopted to measure the IFTs and swelling factors of pure and impure CO_2 gas with n-decane oil. Four CO_2 binary

mixtures (CO_2/N_2 , CO_2/O_2 , CO_2/H_2 , and CO_2/CH_4) and two multi-components mixtures are investigated experimentally at 310.95 K. The MMPs of the systems are estimated experimentally using the Vanishing Interfacial Tension Approach (VIT) and predicted using a developed in-house modified algorithm based on the multiple-mixing cell approach.

7.2 Experimental Section

7.2.1 Materials

The following compounds were used to make the different synthetic mixtures studied in this work:

- Carbon Dioxide, Supplied by Air Products, Research Grade
- Methane, Supplied by Air Products, Grade N4.5
- Nitrogen, Supplied by BOC, Research Grade
- Oxygen, Supplied by BOC, Research Grade
- Hydrogen, Supplied by BOC Gases, $\geq 99.995\%$
- N-decane, Supplied by ACROS-ORGANIC $> 99\%$

The systems listed in [Table 7.1](#) were prepared gravimetrically and used to conduct density tests. Multi component mixtures, MIX1 and MIX2 (compositions given in [Table 7.2](#) supplied by BOC were also used to carry out the tests.

Table 7.1 Injection gas compositions

Component	$\text{CO}_2\%$ mole	Impurity %mole (± 0.3)
Pure CO_2	100	0
CO_2+CH_4	95	5
CO_2+N_2	95	5
CO_2+O_2	95	5
CO_2+H_2	95	5
MIX1	95.64	See Table.7.2
MIX2	89.83	See Table.7.2

Table 7.2 Mixture compositions MIX1 and MIX2

Mixture Components	% mole (MIX1)	% mole (MIX2)
Carbon Dioxide	Balance	Balance
Methane	0.6261	-
Hydrogen	0.8175	-
Nitrogen	1.41	5.05
Carbon monoxide	0.2127	-
Argon	1.21	2.05
Oxygen	0.08	3.07
Total	100	100

7.2.2 Equipment and procedures

The basic theory for the capillary rise technique is to measure the meniscus height in a round glass tube with known inner radius. The equation governing the capillary rise in the tube is given by:

$$IFT = \frac{rh(\rho_{liquid} - \rho_{gas})g}{2\cos(\theta)} \quad (7.1)$$

Where:

IFT = interfacial tension in mN/m; r = pore throat radius in cm; h = capillary rise in cm

ρ_{liquid} = density of liquid phase in g/cm³; ρ_{gas} = density of vapour phase in g/cm³

θ = equilibrium contact angle in degrees (equals to zero, glass completely wet)

g = acceleration due to gravity (980 cm/s²)

To achieve high accuracy, the tube must be absolutely vertical, its cross-section uniform and the liquid should be completely wetting the wall. In the presence of a vapour phase, the contact angle (θ) in the above equation can be set to zero for CO₂-oil systems, since liquids wet the glass surfaces completely [31]. The schematic diagram of the capillary rise technique used is shown in Figure 7.1. The relation between the tube radius, r , and the capillary rise, h , given in Equation 7.1 is inversely proportional to the interfacial surface tension i.e. as the radius becomes smaller the rise can be measured precisely,

especially at low IFT. Therefore, a radius of 0.2 mm tube was installed in this experiment, four and half times smaller than the tube radius used by Ayirala [31] work.

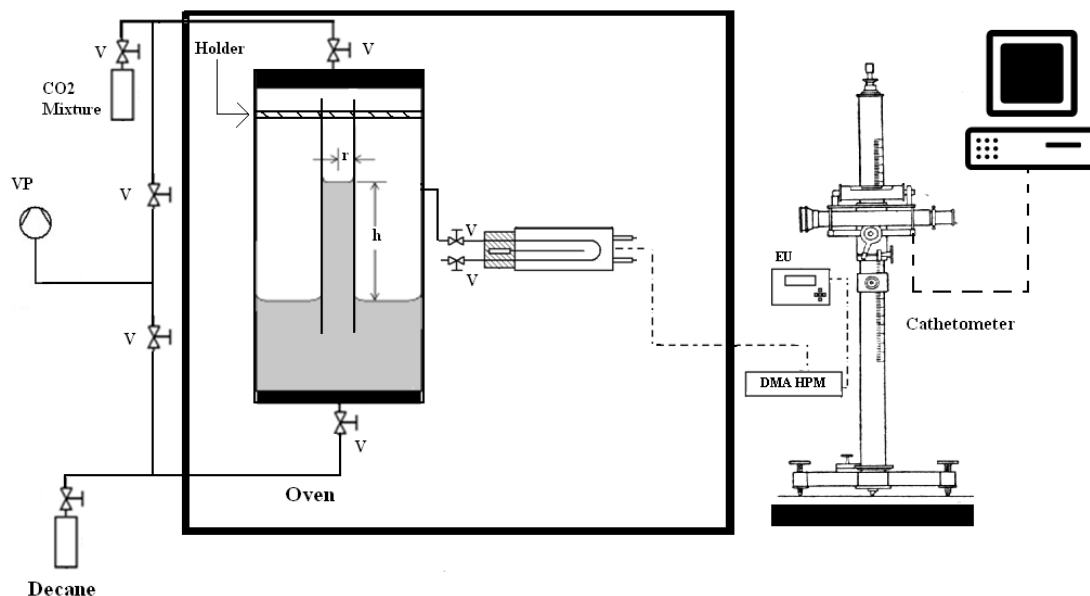


Figure 7.1 Schematic of capillary rise technique used for $\text{CO}_2/\text{n-decane}$ system

At first, the capillary tube was carefully fitted into a special holder and is placed inside the equilibrium cell. The cell has a volume of 100 cm^3 fitted with a glass window and has a pressure limit up to 70 MPa. Since this test is very sensitive to trace contaminants, the cell should be cleaned. Heptane is introduced to the cell and followed by nitrogen purging. The cell and its loading lines are vacuumed prior to introduction of any fluid. The cell was then filled with pure CO_2 gas and was heated to 310.95 K using the oven. “n-decane” in a pressurized transfer vessel was injected into the cell so that the cell was filled with fluids at a fixed initial volumetric liquid oil ratio (80% CO_2 20% n-decane). The cell pressure is regulated by injecting more CO_2 or $\text{CO}_2/\text{impurities}$. The compressed gas density data of CO_2 and $\text{CO}_2/\text{impurities}$ systems were measured with a high temperature and pressure vibrating tube densitometer, Anton Paar DMA-HPM, which consists of a measuring cell and an interface module as shown in Figure 7.1. The liquid density was reported to increase slightly by increasing the pressure in n-decane/ CO_2 system at different temperatures [32, 33]. For example at 344.3 K and pressure up to 12.4 MPa, the increase in the liquid phase density of n-decane/ CO_2 system is about 1.4% compare to the density of pure n-decane. Therefore, the liquid phase density of $\text{CO}_2/\text{n-decane}$ is measured at different pressures; however the liquid density of the

mixtures containing impurities is interpolated with data of CO₂/n-decane system. So the densitometer is only used to measure the density of vapour phase in the tested systems including impurities systems. The densitometer is calibrated with two fluids: pure CO₂ and water at 310.95 K. Two hours are given for the fluid phases to equilibrate in the cell. The capillary rise observed in the glass tube was then measured using the cathetometer. The height was measured precisely using the magnification system of the camera and a computer. The precision of the reading is better than 0.01 mm. Two heights, capillary rise and the oil level in the equilibrium cell, including the density are measured at each pressure. The above steps are repeated for the other systems containing impurities. Examples of images captured from the apparatus for pure CO₂ as injection gas are shown in [Figure 7.2](#).

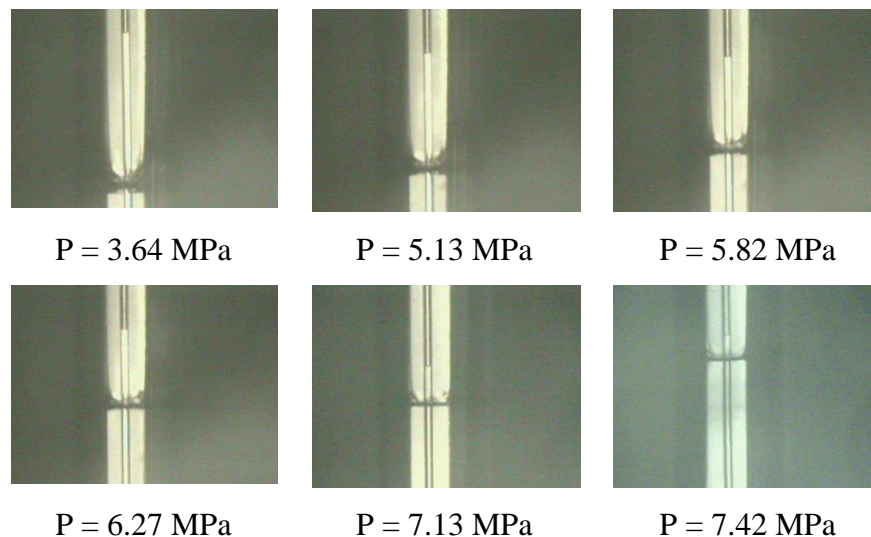


Figure 7.2 Images of CO₂/n-decane at 310.95 K, shown that the capillary rise reduces as the system approaches the MMP by increasing the injection gas pressure

7.3 Results and Discussion

7.3.1 IFT

In order to validate the experimental setup, the IFTs results obtained by the capillary rise method are plotted together with the data available in literature. [Figure 7.3](#) shows that the IFT measurements of CO₂/n-decane generated using the apparatus setup were in excellent agreement with literature data [31]. The Interfacial tension (IFT) between n-decane and CO₂ phase decreases as the pressure increases. The pendant drop results

show that the apparatus cannot measure IFT lower than 3.75 mN/m , making the capillary rise technique better suited for low IFT measurements.

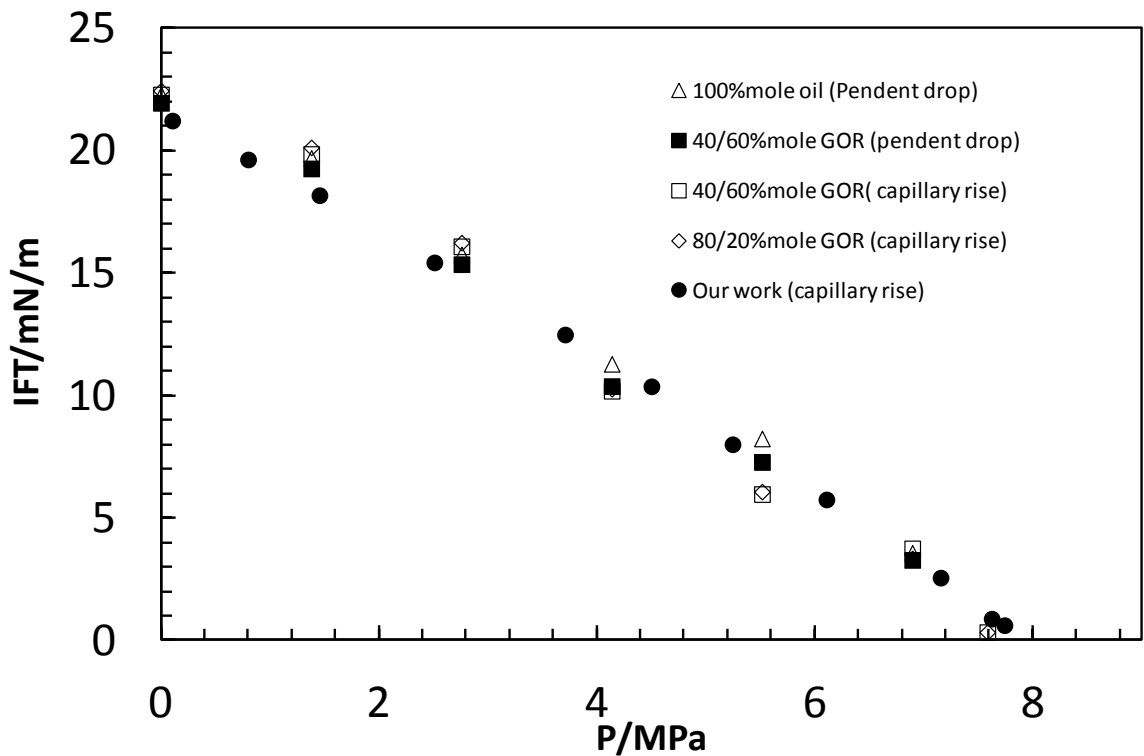


Figure 7.3 Experimental data for $\text{CO}_2/\text{n-decane}$ system at 310.95 K , literature data [31]

The effect of impurities on this property is investigated and listed in Table 7.3. Four binary mixtures are used with 5% mole fraction of impurity, CO_2/N_2 , CO_2/H_2 , CO_2/CH_4 , CO_2/O_2 and two multi-components mixtures MIX1 and MIX2 listed in Table 7.1 and Table 7.2. The IFT values in $\text{CO}_2/\text{impurities}$ systems are shifted to higher pressures making the displacement of n-decane less efficient. As the system pressure increased, the IFTs difference between pure and impure CO_2 in n-decane becomes higher. For the binary systems, hydrogen seems to have the highest impact of the four impurities and methane has the lowest impact. The vapour phase and organic phase (rich decane) IFTs with the three mixtures, CO_2/H_2 , CO_2/O_2 and MIX1, almost behave in a similar manner as pressure increases while MIX2 with 10.17mole% impurity fraction has the highest impact compare to the tested systems. For example, the IFTs at around 2.7 mN/m occurs at pressure 7.2, 8.1, 8 and 9.2 MPa for pure CO_2 , CO_2/H_2 , MIX1 and MIX2 respectively. It means that, both mole fraction and the type of impurities affect the IFT for certain oil such as n-decane used in this work.

Table 7.3 Equilibrated fluid densities, capillary rise heights and IFTs Measured in *n*-decane/ CO_2 System and *n*-decane/impure CO_2 Systems at 310.95 K

Pure CO_2				
P/MPa ($\pm 0.1\%$)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (± 0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (± 0.05)	Height/cm (± 0.002)	IFT/ $\text{mN}\cdot\text{m}^{-1}$ (± 0.05)
0.10	0.7171	0.0018	3.028	21.23
0.80	0.7177	0.0141	2.848	19.64
1.45	0.7186	0.0265	2.68	18.18
2.51	0.7203	0.0491	2.346	15.43
3.71	0.7223	0.0778	1.977	12.49
4.50	0.7236	0.1	1.697	10.37
5.25	0.7247	0.1246	1.361	8
6.11	0.7258	0.1606	1.038	5.75
7.16	0.7262	0.2266	0.522	2.56
7.63	0.7258	0.2751	0.199	0.88
7.74	0.7257	0.2964	0.146	0.61
CO_2/N_2				
P/MPa ($\pm 0.1\%$)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (± 0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (± 0.05)	Height/cm (± 0.002)	IFT/ $\text{mN}\cdot\text{m}^{-1}$ (± 0.05)
0.10	0.7171	0.0018	2.997	21.01
0.99	0.718	0.0173	2.789	19.15
2.82	0.7208	0.0537	2.329	15.23
4.47	0.7235	0.0954	1.836	11.3
5.27	0.7248	0.1184	1.543	9.17
6.71	0.7262	0.1715	1.098	5.97
7.37	0.7261	0.2041	0.696	3.56
7.82	0.7255	0.2322	0.486	2.35
8.22	0.7246	0.264	0.291	1.31
CO_2/CH_4				
P/MPa ($\pm 0.1\%$)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (± 0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (± 0.05)	Height/cm (± 0.002)	IFT/ $\text{mN}\cdot\text{m}^{-1}$ (± 0.05)
0.10	0.7171	0.0023	3.043	21.32
0.81	0.7177	0.0137	2.847	19.64

1.90	0.7193	0.0342	2.579	17.32
3.16	0.7214	0.0611	2.179	14.1
4.85	0.7241	0.1046	1.647	10
6.06	0.7257	0.1458	1.215	6.91
6.78	0.7262	0.1769	0.852	4.59
7.38	0.7261	0.21	0.614	3.11
7.80	0.7256	0.2408	0.478	2.27
CO ₂ /H ₂				
P/MPa (±0.1%)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (±0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (±0.05)	Height/cm (±0.002)	IFT/mN.m ⁻¹ (±0.05)
0.10	0.7171	0.0017	3.02	21.17
1.28	0.7184	0.0222	2.731	18.63
3.08	0.7213	0.0583	2.244	14.58
4.39	0.7234	0.0892	1.834	11.4
5.93	0.7256	0.1324	1.319	7.67
6.59	0.7261	0.1561	1.12	6.26
7.03	0.7262	0.1735	0.968	5.24
7.73	0.7257	0.202	0.671	3.44
8.13	0.7248	0.2247	0.549	2.69
8.54	0.7234	0.2515	0.396	1.83
CO ₂ /O ₂				
P/MPa (±0.1%)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (±0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (±0.05)	Height/cm (±0.002)	IFT/mN.m ⁻¹ (±0.05)
0.10	0.7171	0.0018	3.018	21.16
1.22	0.7183	0.0217	2.713	18.52
2.99	0.7211	0.0579	2.257	14.67
4.48	0.7236	0.0954	1.85	11.39
5.94	0.7256	0.1419	1.368	7.83
7.12	0.7262	0.1933	0.921	4.81
8.32	0.7242	0.282	0.492	2.13
MIX1				
P/MPa (±0.1%)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (±0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (±0.05)	Height/cm (±0.002)	IFT/mN.m ⁻¹ (±0.05)

0.10	0.7171	0.0017	3.022	21.19
1.27	0.7184	0.0221	2.733	18.65
3.17	0.7214	0.0614	2.232	14.44
5.38	0.7249	0.1198	1.544	9.16
6.98	0.7262	0.1811	0.981	5.24
8.01	0.7251	0.2417	0.578	2.74
8.54	0.7234	0.2918	0.369	1.56
MIX2				
P/MPa (±0.1%)	$\rho_{\text{liquid}}/\text{g}\cdot\text{cm}^{-3}$ (±0.0001)	$\rho_{\text{gas}}/\text{g}\cdot\text{cm}^{-3}$ (±0.05)	Height/cm (±0.002)	IFT/mN.m ⁻¹ (±0.05)
0.10	0.7171	0.0017	3.027	21.22
1.35	0.7185	0.0235	2.739	18.65
3.64	0.7222	0.0706	2.175	13.89
5.36	0.7249	0.114	1.703	10.2
7.06	0.7262	0.1708	1.231	6.7
8.39	0.724	0.2382	0.831	3.96
9.17	0.7198	0.289	0.645	2.72

7.3.2 Swelling factor

When CO₂ comes into contact with oil, a process of dissolution occurs thereby causing swelling. The degree of swelling depends on pressure, temperature and oil composition [34]. Swelling is important for two reasons: firstly, the residual oil saturation is inversely proportional to the swelling factor. The residual oil saturation is an important point in relative permeability curves and determines the ultimate recovery. In addition, swollen oil droplets can push fluids out of the pores, creating a drainage process [35]. The oil swelling increases oil saturation therefore increases oil relative permeability too [3]. The governing equation used in oil and gas industries is defined as:

$$\text{swelling factor} = \frac{\text{oil volume at P and T}}{\text{oil volume at 1 atm and T}} \quad (7.2)$$

The relation between the volume is linearly proportional with distance, see Figure 7.4, (measured by cathetometer), i.e. the cross section of the cell can be assumed constant. In

this work, the volume of the cell was calibrated with distance. Therefore the swelling factor can be re-defined as:

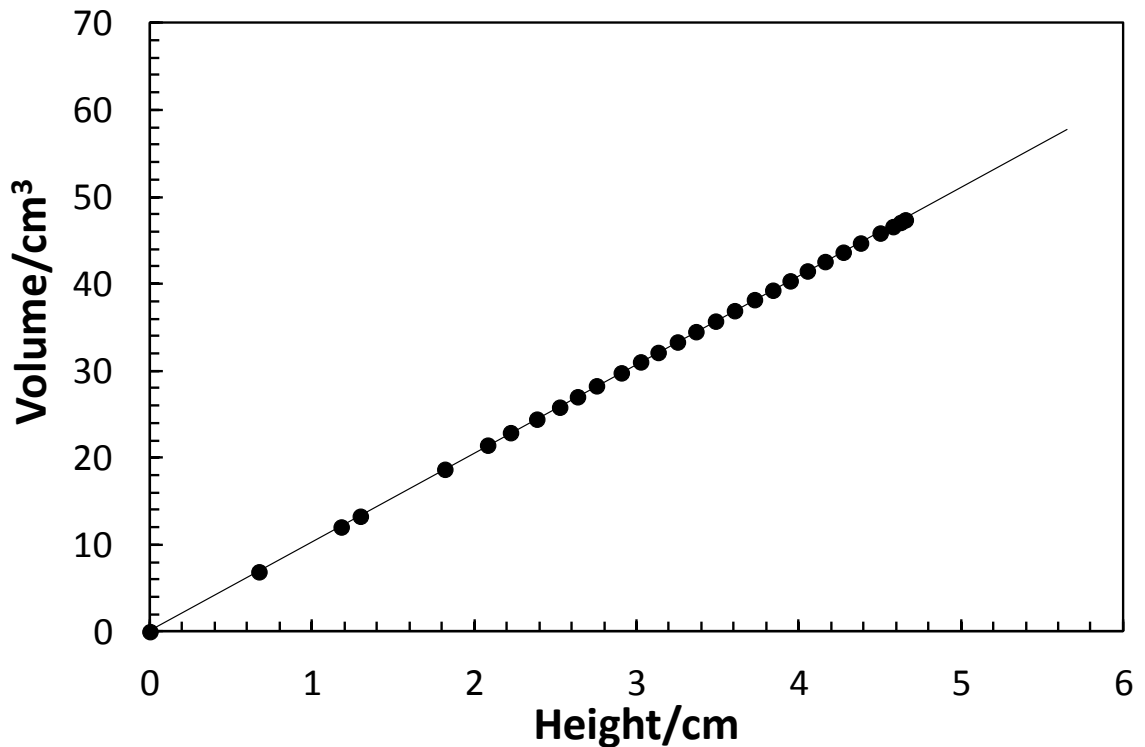


Figure 7.4 The cell volume Calibration in terms of the vertical distance (height)

$$\text{swelling factor} = \frac{Z_{\text{oil}} \text{ at P and T}}{Z_{\text{oil}} \text{ at 1 atm and T}} ; \text{ Where } Z_{\text{oil}} \text{ is oil height.} \quad (7.3)$$

The volume of oil in the liquid phase increases with pressure as CO_2 dissolves in and swells the oil. The isothermal measurements of the swelling factor for the pure and impure CO_2 /n-decane systems at 310.95 K are shown in Figure 7.5. Swelling is temperature independent as was suggested previously [8] when a certain sample is used. Therefore, running swelling factor experiments at 310.95 K can give a good estimation for the effect of impure CO_2 in n-decane at different temperatures. Pure CO_2 swells n-decane to higher value compared to CO_2 containing impurities. For example, at pressure 7.3 MPa the swelling factors are 2.53, 1.73, 1.72 and 1.44 for pure CO_2 , CO_2/N_2 , MIX1 and MIX2 respectively. This means 10% mole fraction of impurities (as the case of MIX2) can potentially reduce swelling by 43% at the pressure given in the above

example. The effect of impurities on binary mixtures is insignificant below 6 MPa, and then the difference becomes wider as the pressure increases.

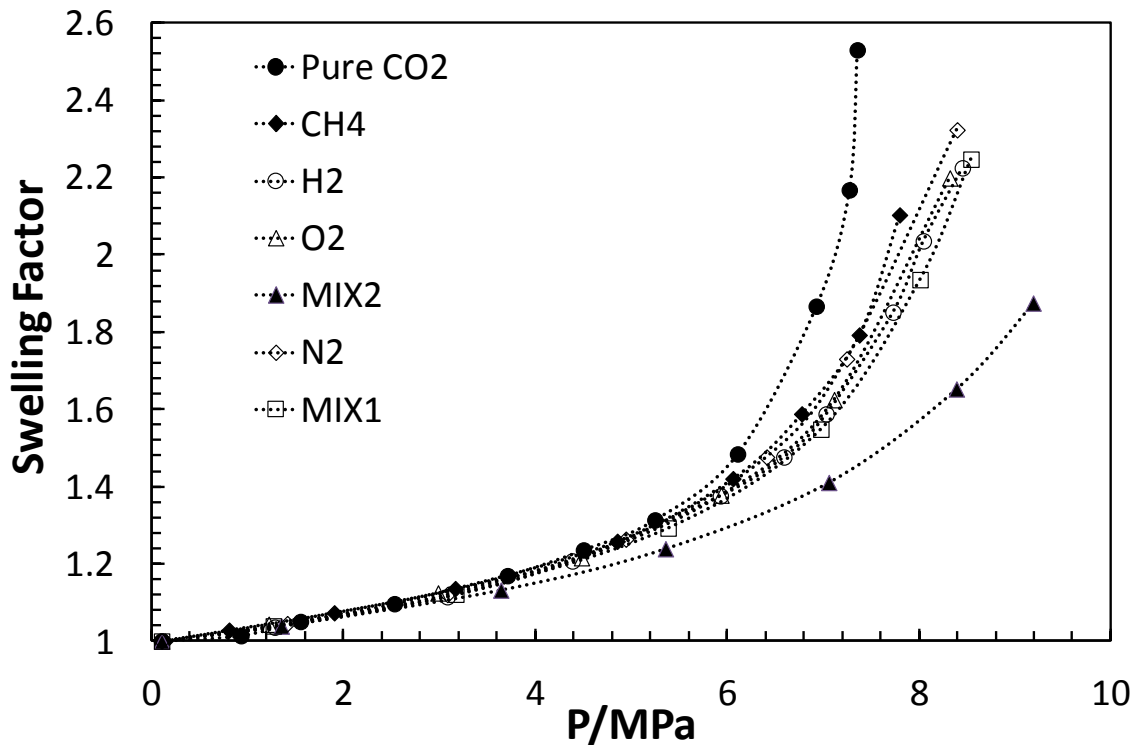


Figure 7.5 Swelling factor measurements in CO_2 and CO_2 /impurities systems with *n*-decane at 310.95 K

7.3.3 MMP

7.3.3.1 Vanishing Interfacial Tension Approach (VIT)

Recently [36] an experimental method, called the vanishing interfacial tension (VIT) technique, has been used to determine the miscibility conditions of different oil- CO_2 systems. The method is based on measuring the IFTs against one of the independent variables, pressure or enrichment levels of the gas phase. The MMP is determined by linearly extrapolating the measured equilibrium IFT versus equilibrium pressure data (in our case) to zero equilibrium IFT, i.e. the corresponding pressure at $\text{IFT} = 0$. The technique is based on the concept that the interfacial tension (IFT) between a crude oil and CO_2 becomes zero when they are miscible. Figure 7.6 demonstrates the VIT approach with pure and impure CO_2 at 310.95 K. The mixtures with their corresponding MMPs calculated by the method are given in Figure 7.7. This CO_2 /*n*-decane system has

a reported slim-tube MMP of 8.2 MPa to 8.6 MPa a rising-bubble MMP of 8.8 MPa at 310.95 K [37]. A new method based on magnetic resonance imaging (MRI) [38] and VIT method [31] also provided MMP values of 7.8 and 7.79 MPa respectively for the same system. The MMPs values obtained with slim tube and rising-bubble apparatus are in good agreement with the MMP estimated in our work using the VIT approach (8.12 MPa). From Figure 7.7, the presence of impurities in CO₂ stream increase the MMP of CO₂/n-decane system to 8.93, 8.71, 9.28, 9.27, 9.23 and 10.36 MPa when the injection gas mixtures CO₂/N₂, CO₂/CH₄, CO₂/H₂, CO₂/O₂, MIX1 and MIX2 are used respectively.

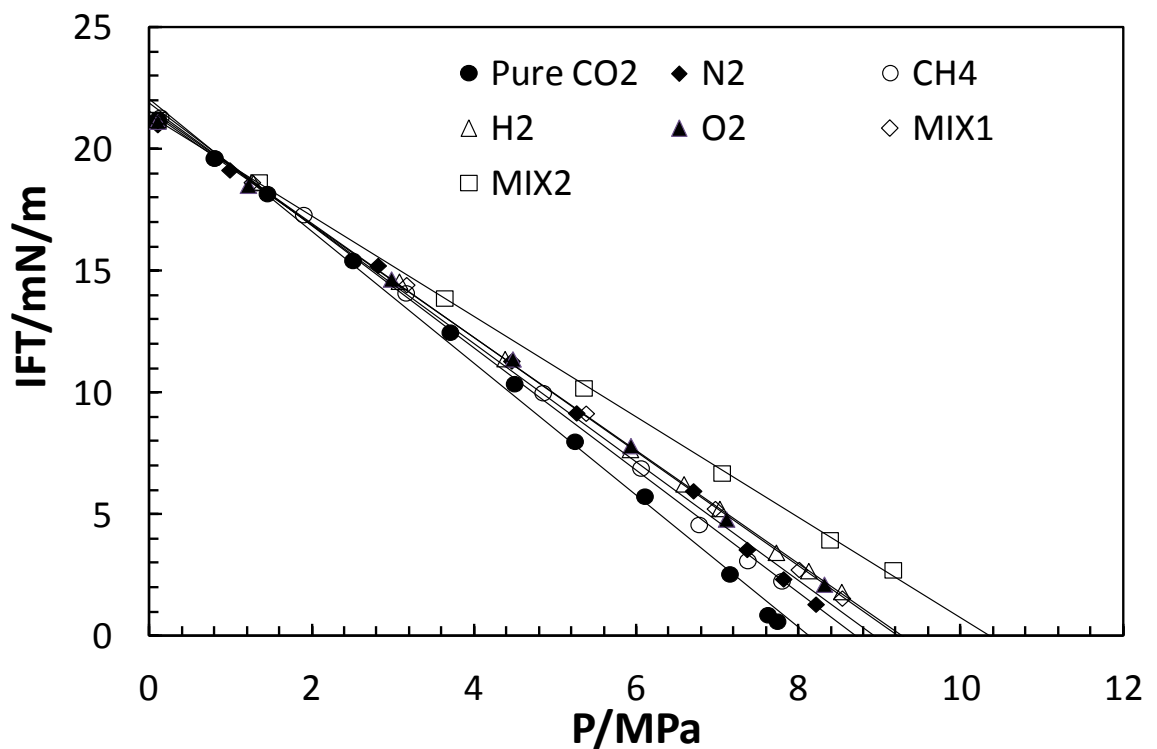


Figure 7.6 Determination of VIT Miscibility in n-Decane-CO₂ System and n-Decane impure CO₂ System at 310.95 K

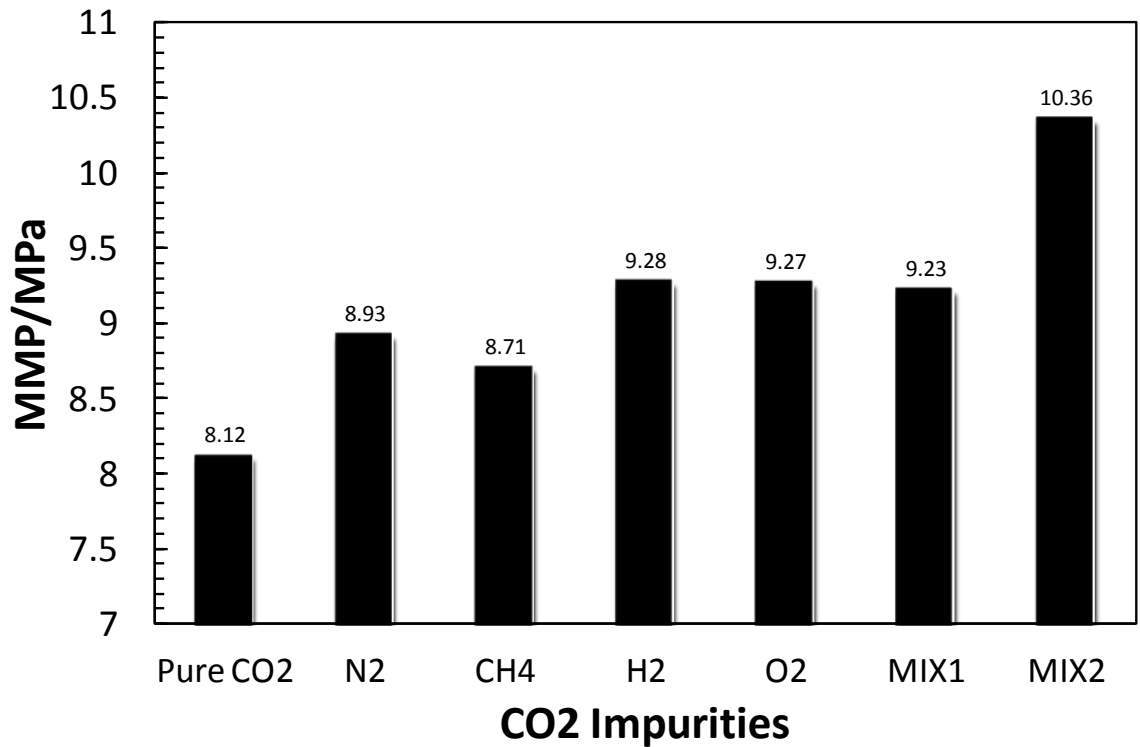


Figure 7.7 The Effect of impurities on MMP of CO₂/n-decane system at 310.95 K

7.3.3.2 Multiple Mixing Cell Algorithm for predicting the MMP

The multiple mixing cell algorithm used in the developed in-house model, is a modification of the algorithm proposed by Jaubert [26] (which is based in the method used in [24] and [27]). The multicell method used is the same proposed by Metcalfe et al. [24] and described in Figure 7.8, where a Batch correspond to the multiple-contact achieved, by maintaining a constant volume after equilibrium (after every flash calculation), transferring the excess in volume from cell to cell, throughout all the cells from the first one to the last one, using 20% of the total amount of gas required for displacing totally the volume of one cell containing 1PV of initial oil (1.2PV of gas). Therefore, knowing that 6 injections (1.2PV/20%) of gas (batches) are required for displacing one cell of initial oil, 300 batches are required for displacing 50 cells, as in our case, ($50 \times 1.2PV / 0.2 = 300$). The equilibrium of each cell is obtained by flash calculation, using the parameters for Peng Robinson EOS given in Table 7.4 and 7.5.

Table 7.4 Pure components properties

comp.	T _c / K	P _c /MPa	ω	MW/g.mol ⁻¹
CO ₂	304.20	7.38	0.23	44.01
CO	132.92	3.50	0.07	28.01
N ₂	126.05	3.39	0.04	28.01
O ₂	154.58	5.04	0.02	31.99
Ar	150.86	4.90	0.00	39.95
H ₂	33.18	1.33	-0.22	2.02
CH ₄	190.58	4.60	0.01	16.04
nC ₁₀	617.70	2.11	0.49	142.29

Table 7.5 Binary Interaction Parameters (BIPs) used in PR EOS

	Binary Interaction Parameters (BIPs)							
	CO ₂	CO	N ₂	O ₂	Ar	H ₂	CH ₄	nC ₁₀
CO ₂	-	0.005	-0.023	0.111	0.129	0.02	0.099	0.104
CO	-0.042	-	0.005	0.000	0.007	-0.016	0.022	0.000
N ₂	-0.048	0.005	-	-0.013	-0.007	-0.020	0.032	0.112
O ₂	0.091	0.000	-0.013	-	0.000	0.000	0.000	0.000
Ar	0.113	0.007	-0.007	0.000	-	0.000	0.026	0.000
H ₂	0.0202	-0.016	-0.02	0.000	0.000	-	0.000	0.000
CH ₄	0.092	0.022	0.032	0.000	0.026	0.000	-	0.041
nC ₁₀	0.104	0.000	0.112	0.000	0.000	0.000	0.041	-

In order to displace the oil of every cell (1 PV= 1 Pore Volume of 1 cell, in our case, each cell contains 1 bbl of initial oil), it is injected 1.2 PV of gas per cell of oil used in the simulation. This injection is made by batches of 20% of the gas (1.2PV/0.2 = 6 injections of gas (batch) per cell of oil to be displaced) as proposed by Metcalfe et al. [24], instead of using batches of 33% of gas (1.2PV/0.33≈ 4 injections) as used by Jaubert et al. [26]. This modification increases the amount of batch used, and therefore the total amount of flash calculations, which refined at every pressure, the recovery factor determined. As mentioned above, in the case of 50 cells, the total amount of batches is: 300 batches, and therefore the total amount of flash calculations: is 300 batch* 50 cells = 15000 flashes calculations, at each pressure supposed. The second modification used, corresponds to a more cautious procedure (more pressure steps at

high recovery factors values) required for generating the curve (Pressure vs Recovery Factor (RF)), in order to obtain the MMP at the value of 0.95% of recovery factor (RF). As it will be explained later, the experimental work is carried-out at low reservoir temperature; therefore drastic changes of the recovery factors with pressure are expected. Due to the higher amount of pressure steps required and the higher amount of flash calculation at every pressure step supposed, a compromise is required in order to calculate the MMP within a considerable computing time. Thus, in this work 50 cells have been used, showing a good accuracy for the prediction of MMP of CO₂/impurities systems.

The mobility method used in this algorithm, is the moving-excess-oil option proposed by Metcalfe et al.[24] in which the excess volume (either gas or liquid or both) is transferred from cell to cell, while the cell volume in each cell is kept constant. [Figure 7.8](#) depicts clearly the multi-contact displacement method used. Before the first gas injection, all the cells contain the initial oil. During the every batch (gas injection), every transfer of excess volume from cell to cell, increases the concentration of volatile components (from the injection gas) in the residual oil. Nevertheless, the excess volume injected becomes heavier as the number of mobility steps (transfer from cell to cell) increases. Thus, on one hand, after reaching the equilibrium during the first batch, the residual oil in all cells becomes lighter than the initial oil and during the second batch, the residual oil in all cells becomes lighter than the residual oil in the first batch and so on, as the number of batches increases.

When the supposed pressure is near to the MMP, the residual fluid becomes more volatile with the increasing number of batches, until the oil in the initial cells is completely displaced (i.e.: the first cell is completely occupied by the injection gas) therefore, the excess volume of these cells is the same as the injection gas volume. On the other hand, during each batch, the initial gas injection is added uniquely to the first cell, and after reaching equilibrium, the excess volume, which is heavier than the initial injection gas, is injected to the second cell, which in turn, after reaching equilibrium, its excess volume is transferred to the third cell, making it heavier than the excess volume injected in the second cell, and so on, as the number of cells increase. In summary, the composition of residual oil becomes heavier with the number of cells, and becomes lighter with the number of batches, thus, the maximum volume of oil recovered (excess volume in the selected cell) is found in the last cell during the first batch.

The recovery factor is calculated in 5 selected cells, which amount to the total number of simulation cells (in our case, the cells: 10, 20, 30, 40 and 50), and is defined as the ratio between the sum of the excess volumes (at standard conditions) during the total number of batches in the given selected cell (i.e.: cell number 10), and the volume of initial oil (at standard conditions) in the number of cells selected to be displaced (i.e.: 10 cells). As it is depicted in Figure 7.8, only the first gas injection (first batch) will be in contact with the initial oil, and the following batches will be in contact with the residual oil obtained from the previous batch after removing the excess volume at equilibrium.

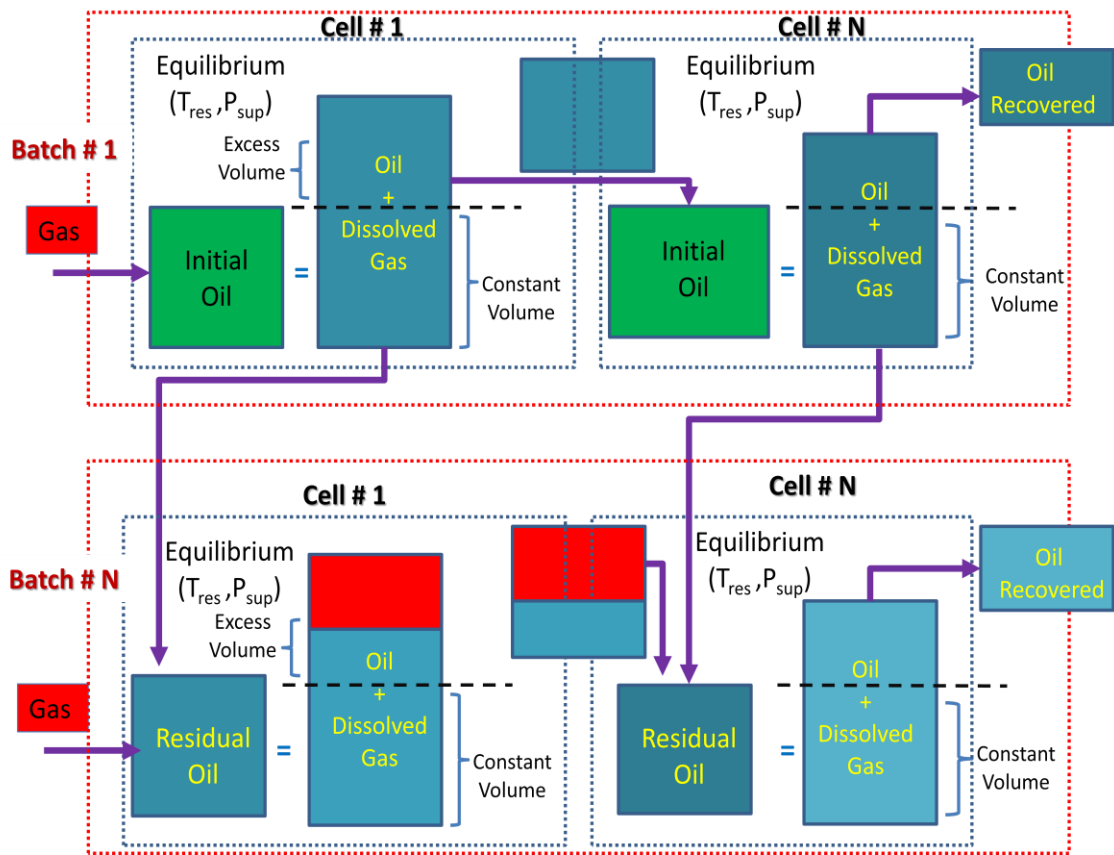


Figure 7.8 Multiple-mixing cell mechanism (constant volume cells, moving-excess-oil option) proposed in [25]

In order to eliminate the high dependence on the number of cells used in the simulation, generating either physical or numerical dispersion, Stalkup [39] proposed to extrapolate linearly the plot of ultimate recovery at 1.2 PV ($RF_{1.2}$) against $(1/\sqrt{N})$ where N is the number of cells. Thus, at a given supposed pressure the recovery factor is calculated for 5 cells (i.e.: 10, 20, 30, 40 and 50, in our case), and the plot of ultimate recovery against $(1/\sqrt{N})$ is elaborated, and finally, extrapolate to infinite number of cells, where $(1/\sqrt{N})$

= 0, in order to obtain a recovery factor independent of the number of cells (RF infinite) at the supposed pressure.

Finally, using the previous data, a curve of recovery factor (RF) vs. Pressure can be generated and the MMP can be found by extrapolating at 95% of recovery factor. Jaubert et al. [26] proposed to use (3 to 5) supposed pressures in the range from 10% to 97% of oil recovery, reporting good results for the prediction of the MMP in several systems. This methods have been used by Hernandez et al. [41] validating the MMP predicted with literature data ([26, 42, 43 and 44]), obtaining a maximum relative error of 1.6%. It is important to note that all the previous cases used temperatures higher than 339.15 K, which is considered by Holm [40] as the range of high temperatures, where the recovery factor increase gradually with pressure. On the contrary, where the temperature is lower than 339.15 K, Holm [40] described a sharp reduction on oil recovery from pressures lower than the MMP.

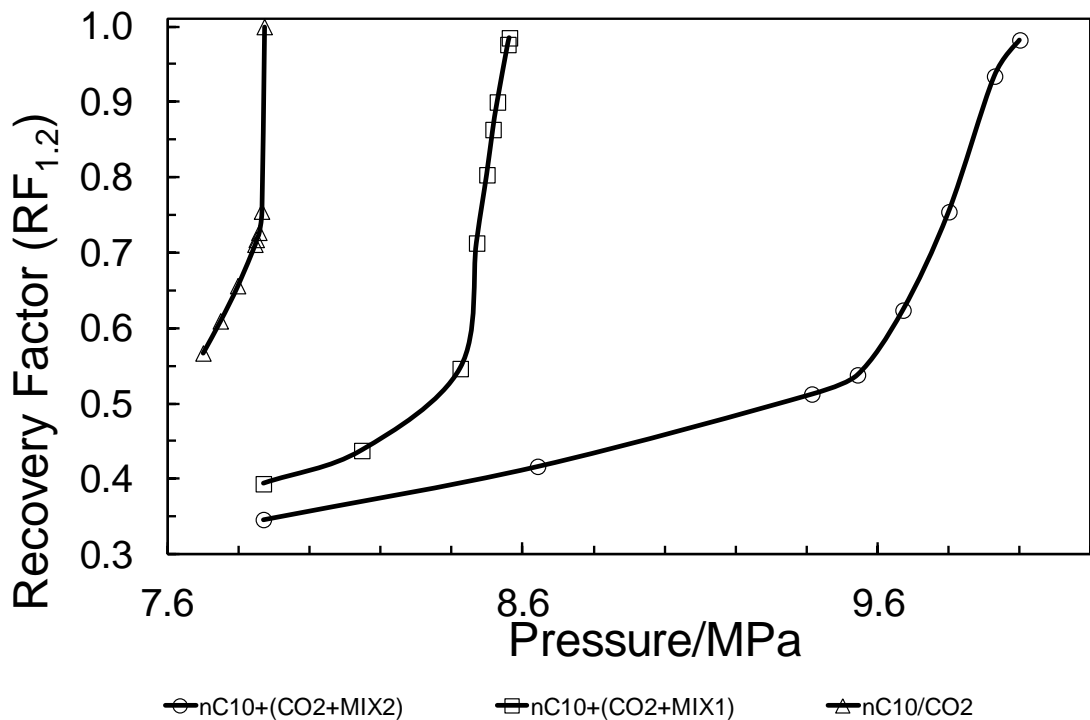


Figure 7.9 Recovery factor vs. Pressure for CO_2 and CO_2 + impurities with *n*-decane at 310.95 K

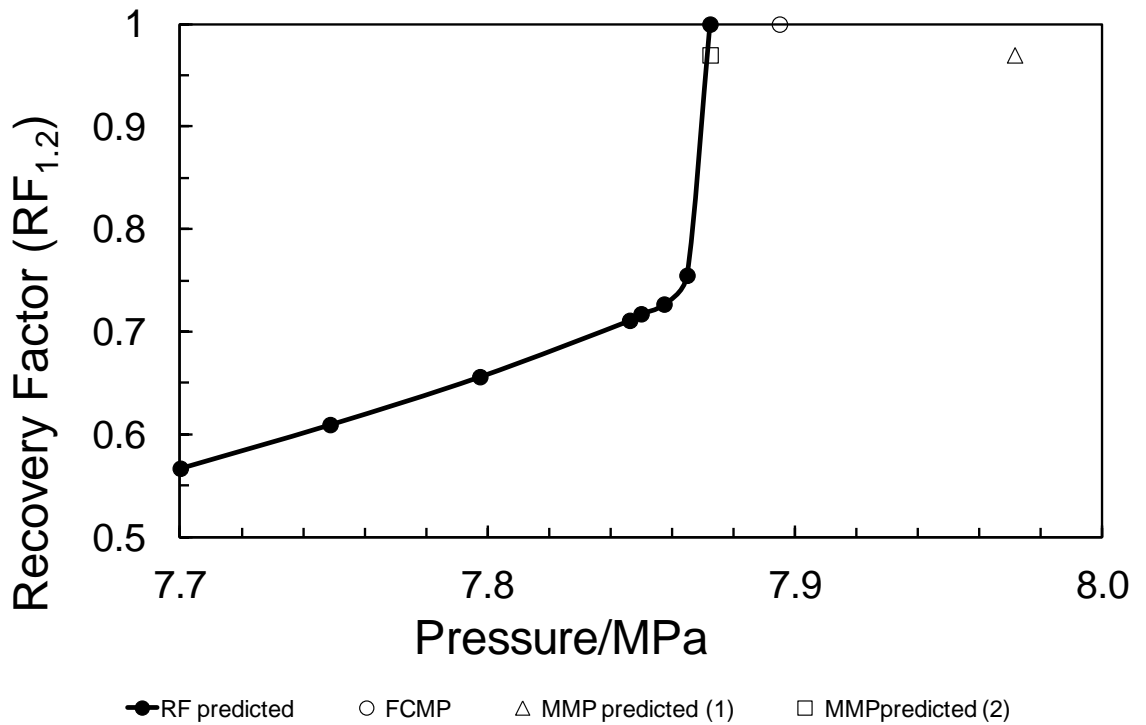


Figure 7.10 Recovery Factor (RF) vs. pressure supposed, using 50 cells. MMP predicted (1): using the range of RF (0.1-0.95) and MMP predicted (2): using the range of RF (0.7:0.99), for $\text{CO}_2+n\text{C}10$ at 310.95 K

This work uses a reservoir temperature of 310.95 K, which is considered as lower temperature, and which behavior described by Holm [40], can be confirmed on Figure 7.10, showing a drastically increase of Recovery Factor from (RF_{1,2}) greater than 75%. Consequently, using the method proposed by Jaubert et al. [26], for extrapolating RF_{1,2} curves against pressure gives a far overestimated MMP, even greater than the FCMP as depicted in Figure 7.10. Therefore a more careful method for selecting the supposed pressures required for generating the curve of ultimate recovery vs. pressure is proposed in this work. A plot of RF_{1,2} vs. Pressure is generated, and only the RF_{1,2} in the range from 75% to 99% are used for interpolating (instead of extrapolating) the MMP to the value of 97% of RF_{1,2}. This modification requires a higher pressure supposed (more than 12) in order to avoid an inaccurate description of the curve of RF_{1,2} vs. Pressure, and therefore a bad prediction of MMP. The sharp increase of recovery with pressure is shown in Figure 7.10 for pure CO₂ in n-decane at 310.95 K. It can be seen that extrapolating the curve in the range of (10 to 95)% will produce an MMP equal to 7.97 MPa which is not acceptable since it is greater than the FCMP, 7.90 MPa. However, using the method proposed in this work; interpolating in the range of (70-99) % of oil

recovery at the value of 97%, gives a good and logical result, MMP equal 7.87 MPa. [Figure 7.9](#) shows the Recovery Factor for three gases with n-decane, pure CO₂, CO₂ + MIX1 and CO₂ + MIX2. The presence of MIX1 and MIX2 in the CO₂ stream can reduce the recovery by 45 and 51% respectively at 7.85 MPa. The effect of the CO₂/impurities near the miscibility displacement is obvious from the trend of the injection of the three gasses. By adding more impurities on CO₂, higher pressures are required in order to obtain similar recovery factor, and therefore higher MMPs are required with the presence of impurities on CO₂. [Table 7.6](#), shows the comparison between the predicted and experimental MMP, obtained with the in-house multi-cell algorithm and the VIT experimental approach for the seven systems studied in this work. The average absolute deviation obtained is lower than 4.3%.

Table 7.6 Comparison of experimental and predicted MMP for the studies systems in this work

Gas injected to nC ₁₀	MMP			
	FCMP	Predicted	Experimental	Dev%
CO ₂	7.89	7.88	8.12	3.0
CO ₂ +N ₂	10.11	8.72	8.93	2.3
CO ₂ +CH ₄	8.85	8.41	8.71	3.4
CO ₂ +H ₂	10.90	8.86	9.28	4.5
CO ₂ +O ₂	9.31	8.65	9.27	6.7
CO ₂ +MIX1	9.53	8.63	9.23	6.5
CO ₂ +MIX2	12.50	9.99	10.36	3.5

The steps for evaluating the MMPs at low temperature (lower than 339.15 K) are described as follows:

- 1) Define the initial composition of oil in all cells.
- 2) Determine the volume of oil and gas at Standard conditions.
- 3) Determine the moles of gas injected; knowing that at each batch 20% of the 1.2 Pore Volume of Oil is injected.
- 4) Calculate the total moles in the cell.
- 5) Calculate the new composition of the cell.

- 6) Flash calculation at reservoir temperature, and pressure from P_1 to P_6 when 20% of PVI is proposed.
- 7) Calculate the new equilibrium volume of the cell.
- 8) Calculate the equilibrium excess volume of oil & gas, while the cell volume is kept constant.
- 9) Evaluate the recovery factor ($RF_{1,2}$) at the end of each set of cells, by dividing the excess oil volume exiting cell (N_{10} to N_{50}) to the initial oil volume. The oil recovered is calculated at surface condition (15 °C, 1 atm).
- 10) To evaluate the $RF_{1,2}$ at infinite number of cells, a linear extrapolation is plotted for $RF_{1,2}$ against $1/\sqrt{}$ (Chosen cells).
- 11) The values of $\ln(RF_{1,2}$ at infinity) are plotted with their corresponding pressures, in our calculation a perfect linear fitting is found for $RF_{1,2}$ at infinite from 0.7 to 0.99.
- 12) The MMP is interpolated for $RF_{1,2}$ at infinity equal to 0.97.

7.4 Conclusions

IFTs of CO₂ (pure and impure) in n-decane were measured with the capillary rise method. The method can be used to measure low IFTs (0.044 dynes/cm). The apparatus is also used to evaluate the swelling factors of the systems. The VIT approach can give a good estimation of the MMP of the tested systems when the IFTs measurements are fitted linearly. Overall, the effects of the tested impurities (N₂, H₂, O₂, H₂S, CH₄, CO) on CO₂ stream of injection of gas, for miscible displacement of n-decane are:

- 1) Shift the IFTs to higher values
- 2) Reduce the swelling factors
- 3) Increase the MMPs

Finally, a good agreement was obtained between the predicted MMPs using the proposed in-house model for low reservoir temperature, and the MMP obtained with VIT approach.

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CHAPTER 8: CONCLUSIONS AND RECOMMENDATIONS FOR FUTURE WORK

8.1 Introduction

- 1) CO₂ could potentially be captured from flue gasses power plants and other large industrial stations and then transported for either sequestration or EOR purpose. The purified CO₂ would not be pure and may contain some impurities which potentially could affect the behaviour of the CO₂ stream. In this work, several transport and physical properties of CO₂ in the presence of impurities were investigated. The properties that were covered are VLE, critical data, density, speed of sound, isothermal compressibility and hydrate formation. In addition the effect of impurities on other properties such as interfacial tension, swelling factor, MMP, saturation pressure and speed of sound in CO₂/alkane mixture were also investigated. The present research is aimed to evaluate both experimental and modelling aspects of the above mentioned properties on CO₂ and CO₂/impurities mixtures. A summary of the main outcomes of the work are highlighted below.
- 2) A large experimental database from literature of VLE of CO₂ and binary CO₂ rich mixtures was gathered. Three EOSs, PR, SRK and VPT, were tuned in a wide range of pressure and temperature ([Chapter 2](#)). The binary interaction parameters of the equations were evaluated for the binary CO₂/impurities systems. Eight impurities were investigated in this work, H₂, Ar, CO, SO₂, H₂S, O₂, N₂ and CH₄. Both maximum and average deviations in the saturated pressure and vapour composition were plotted for the CO₂/impurities mixture. The work in [Chapter 2](#) also includes the evaluation of the critical data of the CO₂/impurities mixtures.
- 3) Saturated pressures were measured experimentally for pure and impure CO₂ in a synthetic alkane mixture at 344.3 K ([Chapter 2](#)). The alkane mixture contains C10, C13 and C16 while the CO₂ impurities used were Ar, H₂, CO and CH₄. The VPT EOS was used to predict the measured data.
- 4) A series of new isothermal compressibility and speed of sound data of CO₂ and CO₂/impurities systems were generated above the saturation pressure from 268.15 to

301.15 K (Chapter 3). A new modified volume correction is presented in this work based on the PR EOS for CO₂ and CO₂/impurities in order to get better prediction of liquid isothermal compressibility in CO₂ and CO₂ rich systems. The speed of sound in liquid alkane/CO₂ mixture (Chapter 2) was measured. The sound property in pure CO₂ mixture was fitted with a new correlation to match the experimental data in order to compare with the mixture contains impurities.

- 5) Seven new isothermal density measurements were generated for CO₂ and CO₂ rich systems from 393.15 to 423.15 K. The measurements were carried-out with an Anton Paar densitometer in both liquid and supercritical CO₂ regions. A new modelling approach is adopted in the study based on mixing volume obtained from simple EOSs such as PR and SRK with the CO₂-MBWR in order to predict the density of both pure and impure CO₂ mixtures (Chapter 4).
- 6) Experimental and calculated phase equilibrium conditions of pure and impure CO₂ hydrate with water were measured in a wide range of pressure and temperature (Chapter 5). An in-house code based on the VPT EOS and the solid theory of van der Waals and Platteeuw was used to predict the hydrate stability of CO₂ and CO₂ rich mixtures in both gas and liquid regions. The impurities covered in the work were N₂, CH₄, CO, O₂, Ar and H₂.
- 7) Viscosity tests to study the impact of impurities on CO₂ liquid viscosity were conducted from 280 to 343.15 K (Chapter 6). The measurements were generated using an in-house capillary tube viscometer. Two viscosity models LBC and Pedersen were modified in order to minimize their errors with the generated experimental data.
- 8) A new IFT apparatus was developed in this work based on the capillary rise method (Chapter 7). The apparatus was used to quantify IFT of CO₂ and n-decane system at 310.95 K. The measurements were extended to cover systems that contain impurities in CO₂ stream. In this work, the systems investigated were binary CO₂ rich mixtures containing CH₄, N₂, O₂ and H₂ and two multi-components mixtures. The apparatus was also used to determine the swelling factor property simultaneously with IFT measurements. The Vanishing Interfacial Tension (VIT) approach was used to estimate the Minimum Miscibility Pressure (MMP) of the above mentioned systems. In addition, a Multiple Mixing Cell (MMC) was employed for computing the MMP of the investigated systems.

A number of important conclusions and recommendations are outlined in this chapter from the preceding chapters of this thesis. This includes the experimental and modelling works plus a summary of the impact of impurities on various properties of CO₂ stream mentioned in details previously (Chapter 1 to Chapter 7). In the light of this work, these conclusions and recommendations are stated below:

8.2 Conclusions from Literature Survey

Literature review is included previously in each chapter. In reviewing the literature, no or very little data were found on some CO₂/impurities properties. Whereas the effect of some impurities were covered widely, for others no data were found. A literature survey on CO₂ rich mixtures was generated by Li et al. [1], associated with the transport properties such as viscosity and thermal conductivity. One main conclusion was that no data were available for transport properties in the liquid phase. In this work, eight components were considered, H₂, Ar, CH₄, N₂, O₂, CO, SO₂ and H₂S. The main points regarding the literature review are summarized below

- ***VLE***

Experimental data are available in the literature for the binary systems of CO₂/H₂S, CO₂/H₂, CO₂/CH₄ and CO₂/N₂. However, there were still some gaps in the literature for experimental data in the other impurities. CO₂/Ar and CO₂/SO₂ have very little information in VLE, therefore the interaction parameters generated for the EOSs may not sufficient accurate in wide range of pressure and temperature. Limited data regarding ternary and multi-components VLE data are available except for the N₂/CH₄/CO₂ system. Critical data of rich CO₂ mixtures are limited in literature for CO₂/N₂, CO₂/H₂S and CO₂/CH₄; and none for the other impurities systems.

- ***Density***

Limited density measurements are available for CO₂ + (O₂, N₂, CH₄, SO₂) and other impurities. Recently, one experimental work on supercritical density [2] was aimed to analysis the effect of three binary CO₂/impurities includes O₂, Ar and N₂.

- ***Viscosity***

No viscosity data are available for CO₂/impurities in both the liquid and supercritical regions. Limited data are available in the gas phase for systems containing SO₂ or CO compare to the other impurities. However no research work has deeply looked at the effect of impurities in rich CO₂ system on viscosity.

- ***Speed of Sound & Isothermal Compressibility***

No experimental data for speed of sound and isothermal compressibility properties for CO₂/impurities was found in the literature in the liquid and supercritical phases. Few data points were published on gas phase.

- ***Hydrate***

Data are widely available for pure CO₂ and the CO₂/CH₄ systems. Few data are available for CO₂ with N₂, H₂ and CO. No experimental data were found in the open literature for CO₂ with O₂, Ar and H₂S.

As mentioned in the previous chapters, CO₂ is used for Enhanced Oil Recovery (EOR), where CO₂/oil properties such as Minimum Miscibility Pressure (MMP), VLE, IFT and swelling factors are important. The literature collected for those properties is stated below.

- ***MMP ,IFT and Swelling Factor***

Few experimental data were available on the impact of impurities CH₄, N₂ and H₂S on MMP and swelling factor measurements. However no data was found for the other impurities effect.

- ***VLE of Impure CO₂/Oil Mixture***

No data are available regarding the effect of impurities on VLE of impure CO₂/oil.

8.3 Experimental Apparatus

In the event that no data are available in the open literature, few measurements were carried to define the effect of impurities on CO₂ properties such as speed of sound, density, viscosity and IFT. Several apparatus were used in this study for measuring CO₂

and CO₂/impurities properties includes an in-house acoustic apparatus (Chapter 2 and Chapter 3 Anton Paar densitometer (Chapter 4) hydrate mixed autoclave rig (Chapter 5), in-house capillary tube viscometer (Chapter 6) and in-house IFT capillary rise apparatus (Chapter 7). The acoustic apparatus can be used effectively to measure the liquid speed of sound of CO₂ systems within less than 2 m/s⁻¹ in accuracy. The travelling time versus pressure generated from the setup can be used successfully to quantify the saturation pressure of the CO₂ systems. The calibration method based on one substance (CO₂ in this work) at two different pressure conditions (see Chapter 4) proposed in the work for Anton Paar densitometer is suitable and more reliable when impurities are considered. Slim tube viscometer operated by mercury pump is based on the fundamental Poiseuille's law mercury, therefore the viscosity measurements are found in a good agreement with other published data. However, a caution should be taken at high temperature where the system pressure may not be stable. The capillary rise setup is a new in-house setup added to this work to measure the IFT between oil and gas. This work is a complement of previous work by Ayirala [2]. Moreover, in order to achieve more accurate measurements, a tube with 4½ times smaller than his work was used. The apparatus can be suited to measure low interfacial tensions (0.044 mNm⁻¹). The apparatus has been successfully used to measure swelling factor.

8.4 Modelling Results

A number of models were implemented to the research to improve the prediction of CO₂/impurities properties. This includes a volume correction based on PR EOS for isothermal compressibility, a CO₂-MBWR and general EOS mixing volume approach for density, a thermodynamic approach was employed to model the phase equilibria, VIT and MOC approaches for MMP and modified LBC and Pedersen models for viscosity. The volume correction used in Chapter 3 is very accurate (less than 0.1% in AAD) prediction for CO₂ liquid density from 228 K to the critical temperature and pressure up to 50 MPa. The modified model results are also in a good agreement with the generated isothermal compressibility results of both pure and impure CO₂. The hydrate model was used to predict the hydrate stability zones of the CO₂/impurities systems in the presence of free water. Good agreement between predicted values and experimental data was found. For viscosity calculations, two models, LBC and Pedersen, were modified and their results show better prediction compare to the original

form of the equations. The results show that the CO₂-Pedersen model can predict the mixture viscosity with less than 2% deviation compare to 9 % in the original Pedersen model and from 9% to 2% in the LBC model. For MMP, the Multiply Mixing Cell model is within 4% of the experimental data at the temperature studied in this work. The tuned results of the EOSs on VLE of CO₂ impurities show that the three EOSs predict the VLE of binary mixtures with the same range of accuracy (3.85% to 4.14%) for all the tested impurities systems. However, the presence of H₂ in CO₂ stream makes the EOSs reduces the accuracy compare to the other impurities. The phase behaviours of the CO₂/ impurities are well described by the EOSs in a wide range of pressure and temperature; however the models failed in accuracy as the system approaches the critical point.

8.5 The Effect of Impurities on CO₂ Properties & Future Work Recommendations

The overall conclusions on the effect of the studied impurities on CO₂ properties can be summarized as following:

- ***Critical Data***

The presence of impurities in CO₂ rich system tends to raise the critical pressure of CO₂ except for H₂S. The impurities H₂, CH₄, CO, Ar, O₂ and N₂ have a negative impact as the CO₂ critical temperature is increased while SO₂ and H₂S have a positive impact. Few experimental data on critical points of CO₂/impurities were found in the literature; therefore more VLE and critical data of the systems are required. Moreover, the prediction of the VLE from the cubic EOSs with the classic mixing rules is still weak especially near the critical points thus testing other mixing rules is recommended for any future work. Another option is to modify the mixing rule in the EOSs in order to improve the prediction of CO₂ mixtures critical data.

- ***Speed of Sound & Isothermal Compressibility***

The presence of impurities tested in this work tends to reduce the speed of sound. The reduction in the property strongly depends on pressure compared to temperature variation. As the pressure increases, the isothermal lines of the change will intersect at a pressure \approx 15 MPa. This finding was observed in all the investigated mixtures. In the

other hand, the presence of impurities increases the isothermal compressibility of CO₂ stream. The effect is much higher as the pressure approaches the critical temperature of the mixture.

- ***Density***

The density of pure CO₂ and the CO₂/impurities mixtures have a stronger dependency on temperature compared to pressure at lower temperatures below the critical temperature. Above the critical temperature, a maximum reduction on the density capacity at a certain pressure under a given temperature was found for each CO₂/impurity system. Overall, the light MW impurities tend to reduce CO₂ density much more than those with MW close to CO₂. Other impurities such as NO_x, SO_x and H₂S must be investigated for future work. The pressure and temperature range might be extended to cover higher pressure and lower temperature conditions.

- ***Hydrate***

The hydrate equilibrium pressures were decreased apparently with addition of the tested impurities in the HLL phase. Methane tends to increase the hydrate zone by twice as much as the other impurities (N₂, O₂, H₂, CO and Ar). Investigation other impurities such as H₂S and SO₂ plus increasing the fraction of impurities are recommended for future work in this part.

- ***Viscosity***

The viscosity of CO₂ and CO₂/impurities is a function of pressure and temperature; it increases as pressure increases and decreases as temperature increases. The experimental results show that all of the tested impurities caused a reduction of viscosity compared to pure CO₂. The reduction was found high for light impurities such as H₂ and low as with Ar. As the system approaches the critical pressure and temperature, the viscosity change becomes significant.

- ***Conclusions on CO₂/oil mixture***

The investigated impurities shift the CO₂ IFTs to higher values. Impurities such as H₂ have high impact on the IFT compare to the others. Higher concentrations of impurities

on the mixture increase the IFT. The investigated non-condensable impurities reduce the swelling factor and increase the Minimum Miscibility Pressure of CO₂.

References

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