

The effect of grain size on the twin initiation stress in a TWIP steel

K.M. Rahman^a, V.A. Vorontsov^a, D. Dye^a

^a*Department of Materials, Royal School of Mines, Imperial College London, Prince Consort Road, London SW7 2BP, UK*

Abstract

The influence of grain size on the twinning stress of an Fe-15Mn-2Al-2Si-0.7C twinning induced plasticity (TWIP) steel has been investigated. Five grain sizes were obtained using a combination of cold rolling and annealing. Electron backscatter diffraction (EBSD) analysis revealed that the material exhibited a typical cold rolled and annealed texture. Tensile testing showed the yield stress to increase with decreasing grain size, however, the ductility of the material was not substantially affected by a reduction in grain size. Cyclic tensile testing at sub-yield stresses indicated the accumulation of plastic strain with each cycle, consequently the nucleation stress for twinning was determined. The twin stress was found to increase with decreasing grain size. Furthermore, the amount of strain accumulated was greater in the coarser grain material. It is believed that this is due to a difference in the twin thickness, which is influenced by the initial grain size of the material. SEM and TEM analysis of the material deformed to 5% strain revealed thinner primary twins in the fine grain material compared to the coarse grain. TEM examination also showed the dislocation arrangement is affected by the grain size. Furthermore, a larger fraction of stacking faults was observed in the coarse-grained material. It is concluded that the twin nucleation stress and also the thickness of the deformation twins in a TWIP steel, is influenced by the initial grain size of the material. In addition grain refinement results in a boost in strength and energy absorption capabilities in the material.

Key words: Twinning, Grain Size, Austenitic Steel, Yield Phenomena, Annealing

1. Introduction

High manganese Twinning Induced Plasticity (TWIP) steels have been attracting significant research interest in recent years owing to their high strength (up to 800 MPa) and superior formability (up to 95% ductility) [1, 2, 3, 4]. These excellent characteristics arise from a high work hardening capacity in the steel, which is due to the continuous formation of mechanical twins during deformation. Hence, these properties make the alloys ideal candidate materials

8 for energy absorption applications, including military vehicle armour and auto-
9 motive crash safety. However, the widespread use of TWIP steels, particularly
10 for automotive applications, has been limited. This is partly due to their rela-
11 tively low yield strength, when compared to other advanced high strength steels
12 (AHSS).

13 The TWIP mechanism is observed in alloys which have a medium stack-
14 ing fault energy (SFE), typically in the range of 18-45 mJ m⁻² [5, 4] and is
15 characterised by the formation of discrete sheared grain subregions containing
16 a mirror plane at the interface, *i.e.* nanometer thick deformation twins. The
17 impressive strain hardening exhibited in TWIP steels is largely attributed to a
18 dynamic Hall-Petch effect [6, 7, 8]. As deformation progresses and twins nu-
19 cleate, they act as obstacles for gliding dislocations, effectively resulting in a
20 continuous grain refinement process. Consequently, this leads to a reduction in
21 the dislocation mean free path, thus producing the characteristic high harden-
22 ing rate observed. Although several mechanisms have been proposed to explain
23 the formation of deformation twins [7], it is generally considered to be a process
24 which proceeds *via* a dislocation mechanism, whether by a pole mechanism [9],
25 a deviation process [10], or by twin nucleation through the formation of stacking
26 faults [11].

27 The stress required to generate twinning, known as the ‘twinning stress’,
28 can be considered to be a combination of two separate terms. Firstly, a stress
29 is required for twin nucleation followed by a further stress for twin growth, to-
30 gether defining the twinning stress. However, determining the stress required
31 to nucleate a twin experimentally is extremely difficult [12]. Consequently, it
32 is generally considered that the nuclei for twins already exist within the materi-
33 al *e.g.* stacking faults, and that the twinning stress which is experimentally
34 determined is actually the stress required for twin growth.

35 The morphology and thickness of deformation twins is controlled by the SFE
36 as proposed by Friedel [13], which has been extended by Allain *et al.* [14] who
37 defined a linear relationship between SFE and twin thickness. Similarly, twin
38 thickness is also affected by the initial grain size of the material [14]. Once the
39 first twin system is activated, the twins must develop through the whole grain.
40 However, once secondary twin systems become active, the twins only need to
41 develop between one twin boundary to another since twin boundaries are strong
42 obstacles and comparable to grain boundaries. Thus, secondary twins are much
43 thinner than the primary twins and larger grain sizes promote the growth of
44 thicker primary twins.

45 The relatively low yield strength of TWIP steels has been an obstacle in-
46 hibiting the widespread use of the alloys in industrial applications. However
47 this problem can be resolved using a range of methods. Precipitation strength-
48 ening is one such method. However, the high concentration of carbon, which is
49 generally alloyed into TWIP steels, can lead to the formation of carbides. Fur-
50 thermore, for longer annealing periods the formation of pearlite can occur [6].
51 Another method available for improving the strength of an alloy is grain re-
52 finement. This is attractive since it does not involve changing the chemical
53 composition of the material. Bouaziz *et al.* [6] have predicted that a yield stress

54 of 700 MPa can be obtained in an Fe-22Mn-0.6C TWIP steel with a grain size
55 of $\sim 1 \mu\text{m}$. Similarly, Santos *et al.* [15] and Kang *et al.* [16] have investigated the
56 effect of annealing temperature on recrystallisation in TWIP steels, concluding
57 that specimens exhibiting a finer grain size also display higher yield strengths.

58 Deformation twinning is strongly dependent on crystallographic grain orien-
59 tation and the average grain size of the material [17, 18]. However, only a few
60 studies have been conducted which investigate the influence of grain size on the
61 strain hardening and twinning behaviour in TWIP steels [19, 18, 20, 21, 22].
62 Gutierrez-Urrutia *et al.* [19] have attempted to elucidate the role of grain size
63 on the strain hardening behaviour of a TWIP steel by investigating the disloca-
64 tion and twin substructures in the material using electron channeling contrast
65 imaging (ECCI). The authors concluded that the fine-grained material investi-
66 gated exhibited a different hardening behaviour compared to the coarse-grained
67 material. This behaviour is explained by the existence of a loose dislocation
68 arrangement in the fine-grained material compared to a cell block structure in
69 the coarse-grained. This leads to the formation of a single twin type, lamellar
70 twin structure in the fine-grained material and a multiple twin substructure in
71 the coarse-grained consisting of two active twin types.

72 In a separate study, Gutierrez-Urrutia *et al.* [18] found that grain refinement
73 within the micrometer range does not suppress deformation twinning, although
74 it does become more difficult and a reduction in twin area fraction occurs in
75 the finer grain material. The authors also found that a Hall-Petch relation
76 provided a good estimate for the effect of grain size on the twinning stress and
77 the experimental evidence suggested that the effect of the grain size on twinning
78 stress was similar to the effect on the yield stress of the material. However, Ueji
79 *et al.* [21] suggest that deformation twinning is strongly suppressed by grain
80 refinement. The contrasting conclusions from the two studies may be due to the
81 influence of the stacking fault energy of the alloys investigated in each study.
82 Gutierrez-Urrutia *et al.* [18] used an alloy with a low stacking fault energy
83 ($\sim 24 \text{ mJ m}^{-2}$ [23, 24, 25]), whereas Ueji *et al.* [21] utilised a material alloyed
84 with aluminium and silicon, which had a higher SFE ($\sim 42 \text{ mJ m}^{-2}$).

85 Bouaziz *et al.* [22] have studied the effect of grain and twin boundaries on
86 the hardening mechanisms of TWIP steels, with particular emphasis on the
87 Bauschinger effect during reverse strain testing. Here the authors concluded
88 using a physically based model that the twin nucleation stress was independent
89 of the grain size and was approximately 550 MPa for the grain sizes investigated
90 (between 1.3 - 25 μm). They consequently inferred that the twin initiation strain
91 increased with grain size and was 12 % true strain for a 25 μm grain size.

92 A complimentary study by the authors on a TWIP steel using X-ray syn-
93 chrotron diffraction measurements has shown evidence of the nucleation of de-
94 formation twins before the macroscopic yield point of the material. This ob-
95 servation was rationalised using *ex-situ* cyclic tensile testing at stresses below
96 the macroscopic yield stress. The accumulation of strain with each cycle was
97 observed and post deformation microscopy revealed the presence of fine defor-
98 mation twins in the sample.

99 In the present work, the effect of austenite grain size on the hardening be-

100 haviour and twin initiation stress of a TWIP steel has been investigated by
 101 performing a variety of tensile tests. A range of grain sizes has been obtained
 102 by varying annealing time. The effect of grain size on the mechanical properties
 103 and hardening behaviour has been determined and the twin initiation stress
 104 has been investigated using cyclic tensile testing. Finally, the mechanical test-
 105 ing results have been rationalised and augmented using a range of microscopy
 106 techniques.

107 2. Experimental Procedures

108 2.1. Material

109 The TWIP steel used in this study had a nominal composition of 15Mn-2Al-
 110 2Si-0.7C wt. % and was provided in 3 mm rolled sheet form by Tata Steel Strip
 111 Mainland Europe. The average grain size of the material was $10 \pm 6 \mu\text{m}$ and
 112 the stacking fault energy was calculated to be $30 \pm 10 \text{ mJ m}^{-2}$ using thermody-
 113 namics based models and data [23, 24, 25].

114 2.2. Cold rolling and annealing procedure

115 Strips measuring $25 \times 80 \text{ mm}$ were cold rolled parallel to the rolling direction
 116 of the as-received plate at $\sim 10\%$ reduction per pass to a final thickness reduction
 117 of 50% ; thereby achieving a 1.5 mm final strip thickness.

118 In order to obtain a range of grain sizes the strips were annealed at 850°C ,
 119 employing different soaking times in the furnace to achieve the final grain sizes.
 120 Samples were subsequently either quenched in cooled brine, water or were air
 121 cooled. The experimental annealing schedule is summarised in Table 1.

Table 1: Annealing schedule used to obtain different grain sizes from the cold rolled TWIP steel.

Sample	Annealing temp ($^\circ\text{C}$)	Annealing time (min/h)	Cooling conditions
1	850	1 min	-20°C , brine quench
2	850	2 min	Water quench
3	850	24 h	Air cool
4	850	96 h	Air cool

122 2.3. Microscopy

123 Samples for light microscopy (LM) and electron backscatter diffraction (EBSD)
 124 were prepared following a standard metallographic schedule. Specimens for LM
 125 were etched using a solution of 4% Nital to reveal the grain boundaries.

126 EBSD was performed for grain size analysis on a JEOL JSM6400 SEM fitted
 127 with an Oxford Instruments HKL Nordlys EBSD detector. Step sizes ranging
 128 between $0.15\text{-}1 \mu\text{m}$ were selected for indexing. Backscattered imaging of the
 129 fine twins was conducted on a Zeiss Auriga FEGSEM. Samples for transmission

130 electron microscopy were prepared using Focussed Ion Beam (FIB) milling on
131 a FEI Helios NanoLab 50 series DualBeam microscope and TEM examinations
132 were conducted on a JEOL 2000FX microscope.

133 2.4. Texture

134 After cold rolling and annealing, the texture was characterised from EBSD
135 measurements where a minimum of 1000 grains had been indexed. The data
136 was then used to reconstruct a complete orientation distribution function (ODF)
137 using spherical harmonics. A Williams-Imhof-Matthies-Vinel (WIMV) [27] cal-
138 culation was then performed to remove any ‘ghost’ points. This involves fitting
139 a minimum-curvature orientation distribution (based on the weight of each Eu-
140 ler angle triplet) to the spherical harmonics pole figure. The WIMV calculated
141 ODF is then used to reconstruct the final set of experimental pole figures. These
142 are then visualised using the software Pod2k. The WIMV calculated ODF is
143 also used to determine the texture index (TI). The texture index is useful pa-
144 rameter to compare the texture strength of a sample without regard for the
145 individual components of that texture, where the TI is the mean square value
146 of the orientation distribution. Therefore, a random material has a TI equal to
147 unity, while textured samples have higher values.

148 2.5. Tensile testing

149 Tensile testing was conducted on a Zwick Roell 100 kN load frame using a
150 10 mm gauge length extensometer. Testing was conducted at a nominal strain
151 rate of 10^{-3} s^{-1} on samples with gauge dimensions of $19 \times 1.5 \times 1.5 \text{ mm}$. The
152 tensile axis was aligned to the rolling direction of the plate.

153 2.6. Cyclic testing and twin stress determination

154 In a complimentary study by the authors evidence was provided using mul-
155 tiple methods that twinning was occurring at sub yield stresses, this included
156 X-ray synchrotron lattice strain, peak width and intensity evolution in addition
157 to *ex-situ* cyclic tensile testing.

158 Since it has been shown that the current experimental TWIP steel twins
159 at stresses below the macroscopic yield point, it is possible to experimentally
160 determine the twin initiation stress for various grain sizes using a series of cyclic
161 tensile tests at different target stresses which are below the yield stress. The
162 total accumulated strain (ε_t) can be determined after a set number of loading
163 cycles (N). This can then be used to calculate the amount of microstrain in-
164 duced per cycle (ε_{cycle}) for a given target stress, *i.e.* $\varepsilon_{cycle} = \varepsilon_t / N$. Once the
165 accumulated strain at different stresses is determined for each sample, a linear
166 relationship can be used to determine the twin initiation stress.

167 Cyclic testing was conducted between a threshold stress of 10 MPa and a
168 selected target stress for 50 cycles. The initial target stress was 100 MPa, this
169 value was increased by an additional 100 MPa upon completion of every 50 cycles
170 up to the yield stress for each experimental sample. Each test was conducted on
171 separate tensile specimens with gauge dimensions of $19 \times 1.5 \times 1.5 \text{ mm}$, the tensile

172 axis was aligned to the rolling direction of the plate. Testing was conducted
 173 under position control at a nominal strain rate of 10^{-3} s^{-1} . The accumulation
 174 of permanent strain per cycle was also confirmed using samples with strain
 175 gauges bonded on for extension measurements instead of an extensometer.

176 3. Results and Discussion

177 3.1. Microstructure characterisation

178 The microstructure of the test material after cold rolling and annealing was
 179 fully austenitic for all the experimental annealing times. Electron backscatter
 180 diffraction (EBSD) revealed the existence of numerous $\Sigma 3$ annealing twins.
 181 However, no evidence of ϵ -martensite was found, Figure 1(a-e). EBSD was
 182 also used to determine the average grain size, Figure 1(f). Here twin bound-
 183 aries were excluded from the grain size analysis. The cumulative distribution
 184 function (CDF) from the raw EBSD data was smoothed using a Weibull fit-
 185 ting function. Subsequently, the final number average grain size distribution
 186 function was obtained from the derivative of the Weibull function.

187 A unimodal grain size distribution was observed in all the samples (Ap-
 188 pendix), Figure 1(f), and average grain sizes between 0.7 - 84 μm were obtained

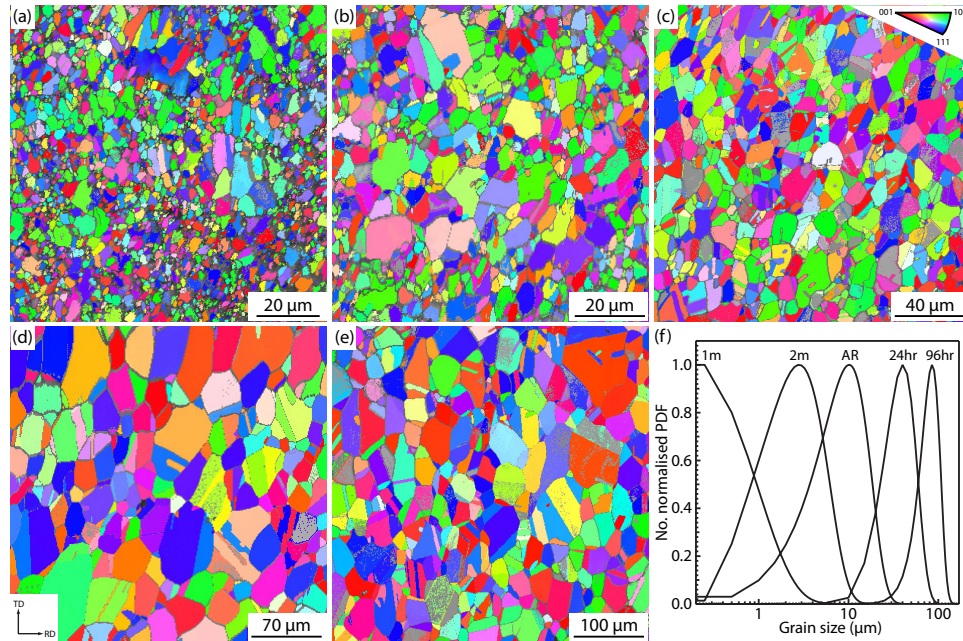


Figure 1: EBSD maps with IPF colouring relative to the rolling direction of specimens cold rolled to 50 % reduction and annealed for; (a) 1 min followed by brine quench [84 % indexing], (b) 2 min then water quench [88 % indexing], (c) as-received material [85 % indexing], (d) 24 h then air cooled [92 % indexing] and (e) 96 h air cooled [88 % indexing]. Number average grain size distribution smoothed using a Weibull fitting function (f).

189 using the different annealing schedules, Table 2. The experimental steel had
 190 an average grain size of 10 μm in the as-received material condition. Therefore,
 191 this sample was not subjected to additional cold rolling and annealing.

Table 2: Average grain size determined from EBSD after cold rolling to 50 % reduction and annealing at 850 $^{\circ}\text{C}$.

Sample	Annealing time (min/h)	Cooling conditions	Scan Area (μm)	No. of Sampled Grains	Average grain size (μm)
1	1 min	-20 $^{\circ}\text{C}$, brine quench	50 \times 50	1571	0.7 \pm 0.5
2	2 min	Water quench	200 \times 150	1077	4.3 \pm 2.4
AR	—	—	500 \times 500	1683	10 \pm 6.0
3	24 h	Air cool	2500 \times 1500	1320	45 \pm 2.0
4	96 h	Air cool	2500 \times 2500	1142	84 \pm 1.0

192 The texture of the investigated alloy after cold rolling and annealing was de-
 193 termined using EBSD after a minimum of 1000 grains were indexed, Figure 2.
 194 A typical cold rolled and annealed texture is exhibited, consisting of three main
 195 components *i.e.* Brass, Goss and Copper. The annealed and recrystallised tex-
 196 ture is similar to that which would be expected in a cold rolled sample [28] with
 197 the exception of weakening of the texture intensities. Since the recrystallised
 198 texture shares the main components to a typical cold rolled f.c.c. texture, it
 199 may indicate that recrystallisation occurs *via* a site saturated nucleation mech-
 200 anism as suggested by Bracke *et al.* [28]. The texture seen in Figure 2(AR) is
 201 weak and is essentially random texture. This is the 10 μm grain size specimen
 202 which is tested in the as received material condition, and has not been further
 203 cold rolled and annealed. Retention of cold rolling texture does not always oc-
 204 cur and randomisation is possible. This may be due to the fragmentation of
 205 coarse grains during cold rolling and the profuse formation of annealing twins,
 206 particularly at higher annealing temperatures.

207 3.2. Mechanical characterisation

208 The tensile behaviour of the five grain sizes investigated in this study can be
 209 seen in Figure 3(a). A significant influence of grain size on the yield strength of
 210 the steel is clearly seen and, as expected, the yield strength increases with de-
 211 creasing grain size. Similarly, an increase in the ultimate tensile strength (UTS)
 212 is exhibited with decreasing grain size. However, an interesting observation is
 213 that the elongation to failure for all the experimental samples is relatively sim-
 214 ilar and a decreasing grain size appears to have little effect on the strain to
 215 failure. It has been reported elsewhere [29] that although high stacking fault
 216 energy (SFE) f.c.c. and b.c.c. alloys tend to display high yield strengths with a
 217 reduced grain size such as 1 μm , they also tend to exhibit a substantial loss in
 218 ductility. The observations in the experimental material may be ascribed to the
 219 relatively low SFE of the alloy and also due to a minor influence on deformation
 220 twinning caused by a reduced grain size.

221 A remarkably high strain hardening rate is observed for all the grain sizes
 222 investigated, which is characteristic of f.c.c. steels which deform *via* twinning.
 223 However, the hardening behaviour of the material is also affected by the grain

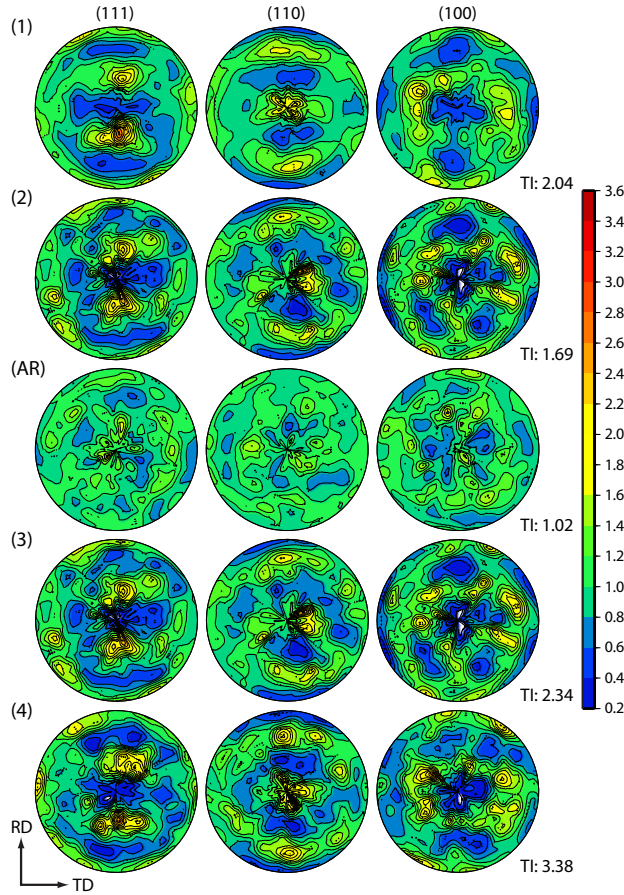


Figure 2: Characteristic cold rolled and annealed texture determined using EBSD for the different annealing conditions along with the texture index (TI); (1) $0.7\ \mu\text{m}$ material, (2) $4.3\ \mu\text{m}$, (AR) $10\ \mu\text{m}$ material in the as-received condition, (3) $45\ \mu\text{m}$ and (4) $84\ \mu\text{m}$ grain size TWIP steel.

224 size where an extra work hardening region is exhibited in the hardening curve of
 225 the fine grain (FG) specimens *i.e.* 0.7 and $4.3\ \mu\text{m}$ when compared to the three
 226 hardening regimes exhibited by the coarse grain (CG) material, Figure 3(b).

227 The hardening behaviour exhibited by the fine grain material initiates with
 228 a decrease in the hardening rate with the onset of straining (region 1). The
 229 decrease in hardening observed in region 1 is non-linear. This suggests a non-
 230 equilibrium between the accumulation and recovery of dislocations. The conse-
 231 quent implication is that a certain level of strain is required before the deforma-
 232 tion twins being formed are thick and frequent enough to affect the hardening
 233 rate. The first stage of the hardening behaviour extends to a minimum with
 234 further strain, finally transiting to region 2. Region 2 initiates with an increase
 235 in the hardening rate. This is caused by an increase in the deformation twinning

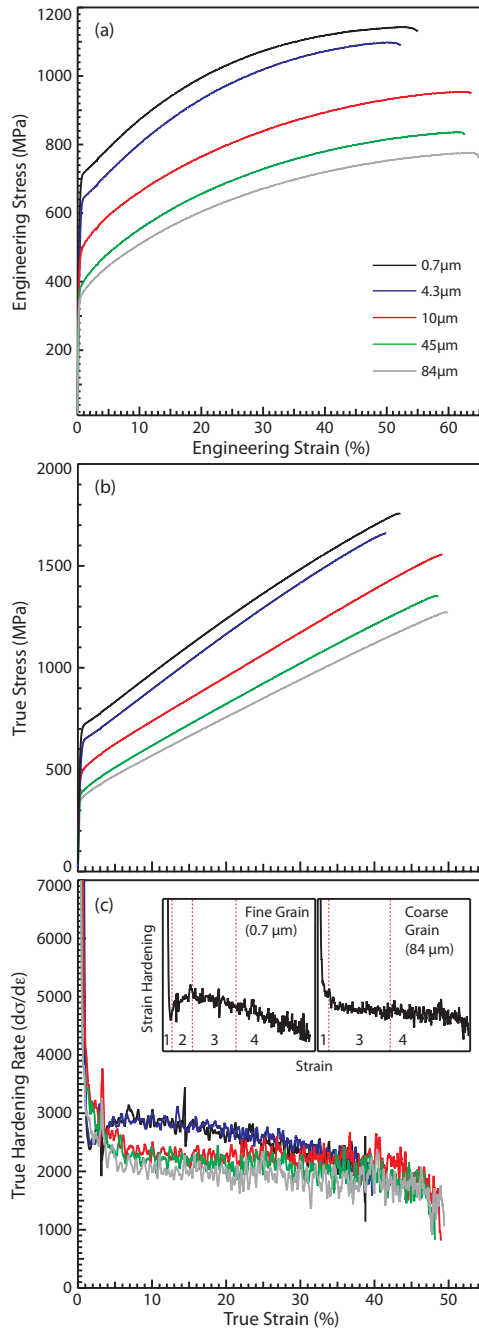


Figure 3: Influence of grain size on the mechanical behaviour of the experimental TWIP steel. (a) engineering stress-strain response, (b) true stress-strain behaviour and strain hardening behaviour.

236 rate. Furthermore, it has been suggested that the increase in the hardening rate
237 may be due to the activation of secondary twin systems [30, 31]. The second
238 hardening stage extends to a maximum after which the third region in the hard-
239 ening behaviour begins. Region 3 extends to higher strains ($\sim 30\%$ strain) and
240 proceeds over a larger strain window compared to regions 1 and 2. As region
241 3 extends with further straining, a subtle and gradual decrease in the gradient
242 of the hardening rate is observed ($\sim 30\%$), which leads to the onset of region 4.
243 Region 4 continues with a decrease in the work hardening rate until the UTS
244 is reached. The initial decrease in the hardening rate is likely to be caused by
245 a reduction in the twinning activity. Since previously formed twins will now be
246 present in the microstructure, they will have effectively reduced the grain size
247 of the material. Therefore, higher stresses will be required to generate further
248 twins. As deformation progresses, the twin volume fraction will inevitably in-
249 crease and consequently the twin bundles will become denser and thicker. The
250 fraction of newly twinned grains will saturate and the nucleation of new twins
251 will become more difficult in the already strain hardened parent grains. Any
252 grains which remain free of twins will almost certainly be in an orientation
253 which is unfavourable for deformation twinning. Hence, a high work hardening
254 rate cannot be maintained with further strain and, as a consequence, a gradual
255 decrease in the hardening rate is observed until fracture.

256 The hardening behaviour exhibited by the coarse grain specimens demon-
257 strates only three distinct work hardening regions compared to the four seen in
258 the fine grain material *i.e.* region 2 is not observed in the coarse grain hardening
259 behaviour. Instead, a gradual decrease in the hardening rate is observed directly
260 between region 1 and region 3. Similar to the fine grain behaviour, a delicate
261 change in the hardening rate at higher strains ($\sim 25\%$ strain) is observed *i.e.*
262 a transition between region 3 and 4. However, the reduction in the hardening
263 rate between regions 3 and 4 ($\sim 50\%$) is more defined when compared to the
264 transition seen in the fine grain behaviour. The hardening behaviour demon-
265 strated by the coarse grain material is often observed in partially recrystallised
266 microstructures. In such cases, an increase in the hardening rate during early
267 deformation is not observed due to the difficulty of twin formation in the recover-
268 ed microstructure [15]. However, since the samples investigated in this study
269 are fully recrystallised, the lack of increasing hardening rate during the early
270 deformation may be due to the need for a longer strain window in order to
271 saturate the larger grains with a sufficient volume fraction of deformation twins
272 compared to the fine grain material. Thus an increase in hardening rate is not
273 observed. A further cause for the hardening behaviour observed may be due to
274 a lower primary deformation twinning rate in the coarser grain material. Idrissi
275 *et al.* [32] have also reported different hardening behaviours observed for TWIP
276 steels with different chemical compositions. The authors suggested that the dif-
277 ferent hardening rates were resultant of different thickness sessile dislocations
278 which are stored at the twin-matrix interface. In addition, the twins formed in
279 the grain sizes showing an extra hardening region were thinner and contained a
280 larger density of sessile defects, thus making the twins stronger. Therefore, the
281 different hardening behaviour observed in the current study may arise due to a

282 difference in twin thickness, which is influenced by the grain size.

283 The increase in the strength of the material with decreasing grain size is well
 284 represented by a Hall-Petch type relation:

$$\sigma_y = \sigma_0 + \frac{K^{HP}}{\sqrt{D}} \quad (1)$$

285 where σ_y is the yield stress, σ_0 is the lattice friction stress, K^{HP} is the strength-
 286 ening coefficient and D is the grain size, Figure 4. It can be seen from Figure
 287 4 that the experimental values are consistently higher than the predicted
 288 behaviour using values for the Hall-Petch constants from the literature for a
 289 Fe-22Mn-0.6C TWIP steel ($\sigma_0 = 132$ MPa and $K^{HP} = 449$ MPa $\mu\text{m}^{1/2}$) [6, 33].
 290 This would suggest that the addition of aluminium and silicon in the present
 291 experimental alloy has a significant strengthening influence on the alloy. This
 292 will also alter the SFE and consequently the morphology of the deformation
 293 twins. Thus, using Hall-Petch constants fitted for the experimental data (σ_0
 294 = 305 MPa and $K^{HP} = 330$ MPa $\mu\text{m}^{1/2}$) shows a better agreement with the
 295 experimental observations.

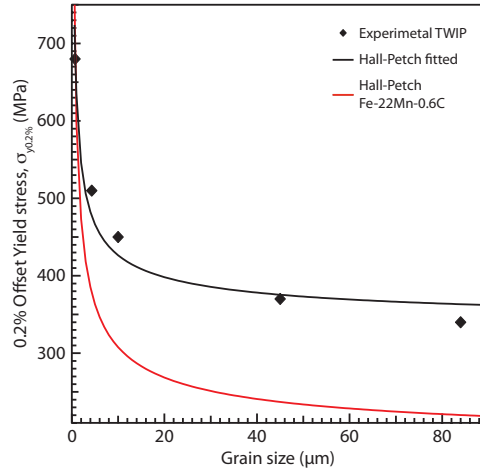


Figure 4: Influence of grain size on the yield strength of the material can be represented using a Hall-Petch type relationship. The addition of silicon and aluminium in the experimental steel suggests additional strengthening (black line) compared to an FeMnC TWIP steel (red line).

296 3.3. Influence of grain size on twin stress

297 Cyclic tensile testing revealed that an accumulation of plastic strain occurred
 298 with each tensile cycle at selected stresses, which were below the yield stress of
 299 the test specimen. This is illustrated in Figure 5, which shows the accumulation
 300 of strain in the 10 μm grain size material, this behaviour was exhibited by all the
 301 different grain size specimens investigated. The stress at which the accumulation

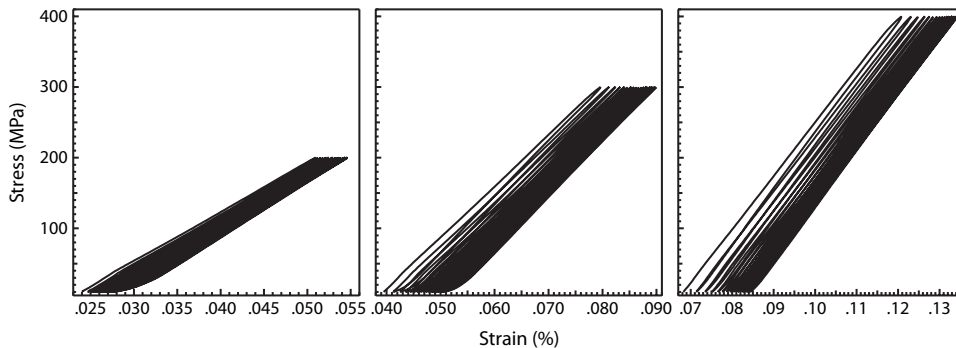


Figure 5: Accumulation of permanent strain during cyclic tensile testing at a range of stresses below the yield point is exhibited as shown in the 10 μm grain size material. This behaviour is characteristic of all the experimental grain sizes.

of strain initiated was lower in the coarse grain material, which suggests that a
 302 coarse grain size promotes early deformation twinning.
 303

304 The twin initiation stress was calculated by plotting the strain accumulated
 305 per cycle at different stresses, fitting a linear relationship to the experimental
 306 data and finding the intercept, Figure 6. The results indicate that grain size
 307 refinement increases the twin initiation stress in the material, Table 3. Further-
 308 more, less strain is accumulated in the fine grain material compared to the coarse
 309 grain which suggests that grain refinement suppresses either the formation of
 310 twins or the subsequent thickening of the nucleated twins. Using the experi-
 311 mentally determined twin stress for each grain size, it is possible to estimate
 312 the critical twin stress for twin nucleation at the single crystal limit (*i.e.* $1/d =$
 313 0). This is achieved by plotting the twin initiation stress against the reciprocal
 314 square root of the grain size. Fitting a linear relationship to the data enables
 315 the limit of large grain size to the twin stress to be calculated, Figure 7(a).
 316 The experimental data suggests the critical twin nucleation stress for an infinite
 317 grain size to be ~ 50 MPa.

318 It is generally accepted that twinning in pure metals and alloys is initiated by
 319 pre-existing dislocations that dissociate into multi-layered stacking faults which
 320 creates a twin nucleus. Several dislocation based models have been proposed
 321 for twin nucleation in f.c.c. materials [34, 35, 11]. All involve the glide of

Table 3: Calculated twin initiation stress for different grain sizes.

Sample	Average grain size (μm)	Twin Stress (MPa)	0.5% Yield Stress (MPa)
1	0.7 ± 0.5	316	720
2	4.3 ± 2.4	184	640
AR	10 ± 6.0	115	490
3	45 ± 2.0	77	390
4	84 ± 1.0	62	350

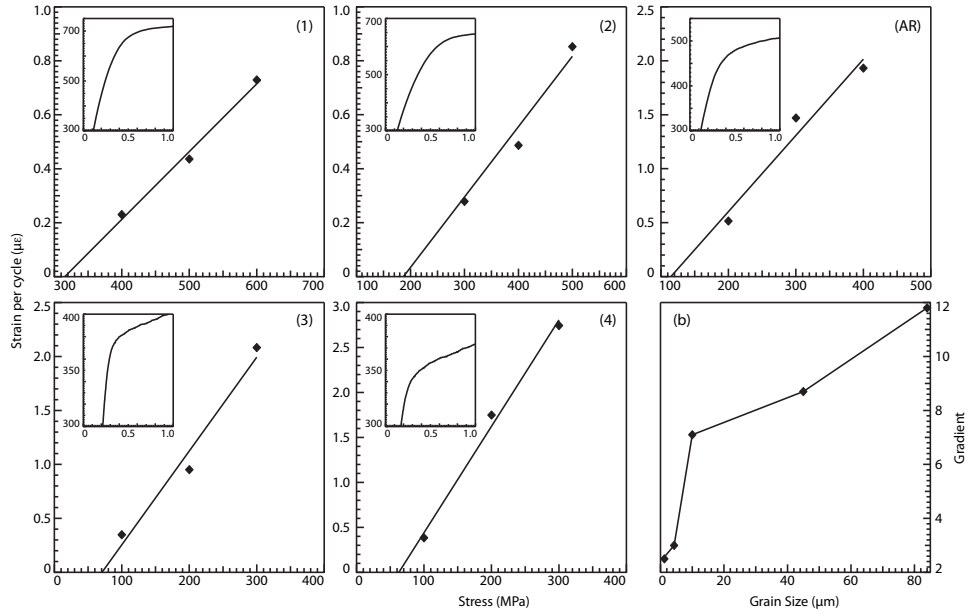


Figure 6: Twin initiation stress determined by plotting accumulated strain against stress and fitting a linear relationship to the experimental data along with a close view of the macroscopic yield transition (inset). (1) 0.7 μm grain size material, (2) 4.3 μm , (AR) 10 μm as-received material, (3) 45 μm and (4) 84 μm and (b) influence of grain size on the gradient of the fitted lines seen in (1-4).

322 Shockley partial dislocations with Burgers vector $a/6\langle 112 \rangle$ on successive $\{111\}$
 323 planes. Since twinning is influenced by the SFE, Venables [35] proposed a
 324 phenomenological relationship between the SFE and the twinning stress where
 325 the influence of the SFE on the twinning stress is proportional. Based on the
 326 analysis of several f.c.c. metals and alloys Narita and Takamura [36] determined
 327 that the SFE and twin stress were proportional such that

$$\tau_{twin} = \frac{\gamma_{SF}}{Kb_s} \quad (2)$$

328 where τ_{twin} is the critical resolved shear stress to separate a leading Shockley
 329 partial from the trailing partial and thus create a twin, γ_{SF} is the stacking fault
 330 energy, K is a fitting parameter which was determined to be 2 by Narita and
 331 Takamura [36] and b_s is the Burgers vector for a Shockley partial dislocation.

332 Since τ_{twin} can be considered to be the twinning stress for a single crystal,
 333 Equation 2 can be used to calculate a critical twinning stress for the exper-
 334 imental alloy using the SFE of the material. Considering the standard deviation
 335 which arises from the thermodynamic derivation of the SFE in the experimen-
 336 tal alloy, a critical twin stress as low as ~ 67 MPa is predicted. The calculated stress
 337 is remarkably close to the experimental prediction for the critical twin stress for
 338 the single crystal limit.

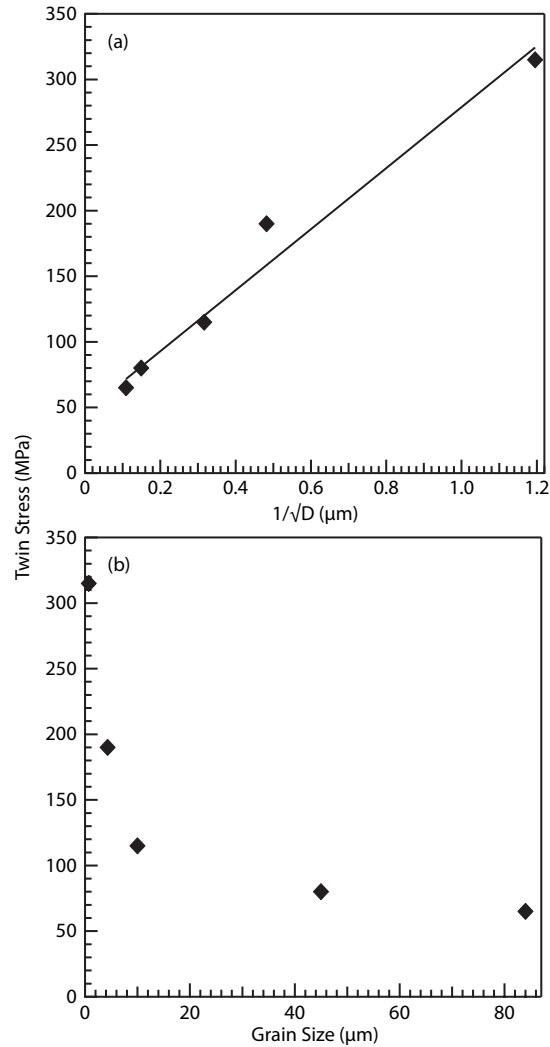


Figure 7: (a) Critical twin initiation stress for an infinite grain size TWIP steel determined from experimental data and (b) effect of grain size on the twin initiation stress shows a Hall-Petch type behaviour.

339 It should be noted that such phenomenological relations are limited. There-
 340 fore, a degree of uncertainty is expected. Kibey *et al.* [37] have shown that
 341 the true twinning stress depends on the entire generalised planar fault energy,
 342 including the unstable twin stacking fault energy and not just the intrinsic stack-
 343 ing fault energy. Similarly, Meyers *et al.* [38] have discussed how the equation
 344 proposed by Venables [35] does not always predict the twin stress correctly, even
 345 though the relationship is accurate for most f.c.c. metals. An example of this is
 346 the case of some copper alloys for which the twin stress varies with the square

347 root of the SFE.

348 The deformation and strain hardening behaviour of low SFE materials is
349 strongly influenced by grain size and consequently the twin stress is dependent
350 on the initial grain size of the material. The length scale of twinning is also
351 expected to have a significant effect on the twinning stress which, in turn, is
352 also affected by the length scale for homogeneous slip. It has been shown by El-
353 Danaf *et al.* [39] that the average slip length during straining in low SFE large
354 grain f.c.c. metals does not change significantly. However, in fine grain materials
355 the average slip length decreases with strain. Although these reductions are
356 relatively small, it is nevertheless enough to inhibit the build up of sufficient
357 dislocations that are necessary to nucleate a nano-sized twin. Furthermore, a low
358 SFE in f.c.c. materials hinders the development of in-grain misorientations. This
359 allows the slip length to remain close to initial grain size, *i.e.* before deformation
360 twinning occurs. Hence, a higher dislocation density and larger average slip
361 length are promoted in a large initial grain size. Therefore, the twinning stress
362 is expected to increase with grain refinement since the slip length and dislocation
363 density are reduced, thus making twin nucleation more difficult.

364 The predicted twin stress for each grain size investigated was determined to
365 be below the bulk yield stress of the material, which is contrary to observations
366 made by other authors [22, 18]. Bouaziz *et al.* [22] have suggested that the twin
367 stress is not affected by the grain size and remains constant at ~ 550 MPa. The
368 authors further suggest that the initiation strain for twinning evolves linearly
369 with grain size, implying that a coarser grain size requires a higher initiation
370 strain. However, one would expect a larger grain size to promote twinning more
371 readily since the slip length is greater. Furthermore, the model employed by
372 Bouaziz *et al.* [22] assumes an average twin thickness which is not influenced
373 by the initial grain size. Conversely, Gutierrez-Urrutia *et al.* [18] have reported
374 the twin stress is strongly influenced by the yield stress which is controlled by
375 the grain size. Consequently, the authors found that a Hall-Petch type relation
376 provided a good estimate of the grain size on the twin stress. However, in both
377 these studies the twin stress is conflated with the yield stress; this confusion
378 most likely underlies the disagreement between these authors' interpretations of
379 their data, and further, with Venables' theory [35].

380 The influence of grain size on the twin stress in the current experimental
381 material indicates that the twin stress decreases with increasing grain size fol-
382 lowing a Hall-Petch type relationship, Figure 7(b). However, the calculated
383 twin stresses are all below the yield stress of the material. Consequently, the
384 twinning constant in a Hall-Petch type model would need to be a lower value
385 than that required for slip. Gutierrez-Urrutia *et al.* in ref [18] have used a Hall-
386 Petch type relationship for determining the grain size dependence on twin stress.
387 Using the Hall-Petch constant value for twinning identical to that for slip, the
388 authors found that the relationship overestimates the twin stress compared to
389 experimental observations. This suggests a lower twin constant compared to
390 slip which the experimental observations for the current TWIP steel support.
391 This also indicates that although the twin stress increases with grain refinement,
392 twinning is not suppressed.

393 A noteworthy observation is the calculated twin stress in the as-received con-
394 dition 10 μm grain size material which is predicted to be ~ 115 MPa. This is very
395 similar to the stress at which deformation twinning was initiated in the same
396 material based on X-ray synchrotron diffraction data in our complementary in-
397 vestigation. This is a reassuring observation, since it has been shown by the
398 authors that deformation twins are observed in TEM foils prepared from spec-
399 imens with the (111) orientation aligned to the tensile axis that have been sub-
400 jected to cyclic tensile testing at 200 MPa. It has also been reported by several
401 authors [18, 30, 40] that during the early stages of deformation in TWIP steels,
402 twinning occurs predominantly in grains orientated close to the $\langle 111 \rangle //$ tensile
403 axis.

404 Once the critical stress twin nucleation is attained, any further stress only
405 serves to thicken the already nucleated perfect f.c.c. twins. Therefore, the
406 initial grain size is expected to influence twin morphology, whereby a larger
407 grain size promotes the formation of thicker twins, since the twin needs to grow
408 over a larger distance. Renard *et al.* [41] have recently shown that a greater
409 twin thickness produces easier internal plasticity of the twins. Therefore, it
410 would be expected that upon cyclic loading, the coarse grain material would be
411 able to accommodate greater plastic strain per cycle since the twin thickness is
412 assumed to be greater. As a consequence, the gradient of the line of best fit in
413 Figures 6(a-e) would be expected to increase with grain size, which is seen in
414 Figure 6(f).

415 Examination of the microstructure after each grain size material is strained
416 to 5% engineering strain reveals that the coarser grain material does indeed
417 contain thicker deformation twins compared to the finer grain sizes, Figure 8.
418 Further investigation of the finest and coarsest grain size specimens using trans-
419 mission electron microscopy reinforces this observation, Figure 9. Here we can
420 see that the relative twin thickness in the 84 μm is distinctly thicker than those
421 observed in the 0.7 μm material even though both specimens have been de-
422 formed to the same strain. Observations from the fine grain material suggest
423 the presence of primary twin less than 10 nm thick within the material. TEM
424 observations also reveal a contrast in the dislocation arrangement between the
425 fine and coarse grain samples. The the dislocation arrangement within the fine
426 grain exhibits a relatively loose configuration and lower volume fraction com-
427 pared to the coarse grain. Many of the dislocations observed in the large grain
428 material appear to exist in dissociated pairs, Figure 9(e). The volume fraction
429 of stacking faults observed in the two samples also appear to be affected by
430 the grain size. A greater fraction of stacking faults were observed in the coarse
431 grain material, Figure 9(f). Since stacking faults essentially operate as nucle-
432 ation sites for twin growth, the many more faults observed within the coarse
433 grain material may explain the lower nucleation stress required to initiate twin-
434 ning in TWIP steels exhibiting larger grain sizes. The varying microstructure
435 observations made from the fine and coarse grain samples suggest the harden-
436 ing behaviour variation seen during mechanical testing is due to different twin
437 thicknesses which is affected by the grain size.

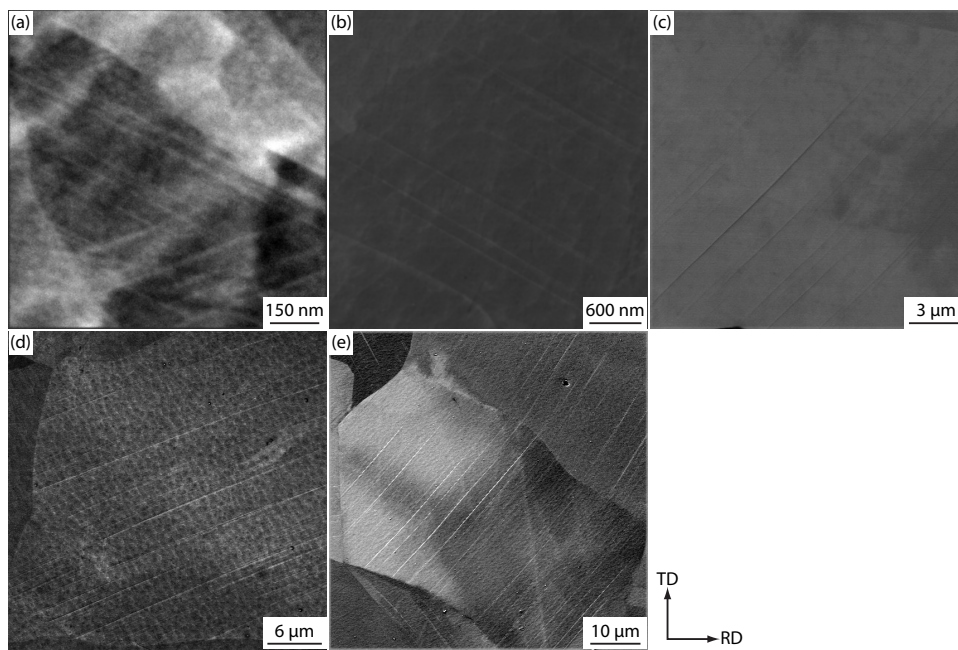


Figure 8: Backscattered electron images of the TWIP steel exhibiting different grain sizes after deformation to 5% engineering strain showing the relative twin thickness is influenced by the initial grain size. (a) 0.7 μm , (b) 4.3 μm , (c) 10 μm , (d) 45 μm and (e) 84 μm grain size material.

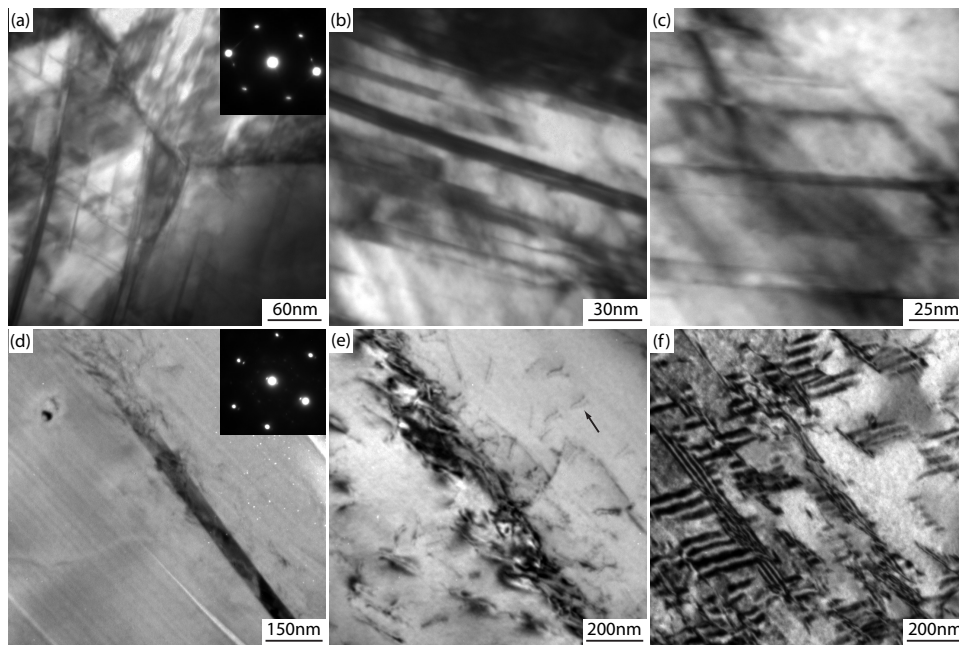


Figure 9: Bright field TEM micrographs and selected area diffraction patterns of the 0.7 μm (a-c) and 84 μm (d-f) TWIP steel after deformation to 5% strain. The deformation twins in the fine grain material are thinner than the twins present in the coarse grain. Numerous dissociated dislocation pairs are observed in the coarse grain material as indicated by the arrow in (e). A large number of stacking faults are also present in the 84 μm grain size material (f).

438 4. Conclusions

439 The effect of initial grain size on the mechanical behaviour and microstructure of a TWIP steel have been investigated using a range of tensile testing
440 and microscopy techniques, including transmission electron microscopy. The
441 following conclusions can be drawn from the investigation:
442

- 443 1. A 15Mn-2Al-2Si-0.7C wt % TWIP steel was produced with 5 different grain
444 sizes ranging between 0.7 - 84 μm using a combination of cold rolling and
445 annealing.
- 446 2. The texture of the material represented a characteristic cold rolled and
447 annealed texture comprising of Brass, Goss and Copper components. The
448 10 μm TWIP was obtained from the as-received condition material and has
449 a random texture.
- 450 3. The tensile behaviour of the material showed an increase in the the yield
451 stress with decreasing grain size. This behaviour was represented well using
452 a Hall-Petch type relationship. Therefore grain refinement is found to result
453 in an overall strength and energy absorption boost.
- 454 4. The strain hardening behaviour is affected by the grain size, whereby an
455 additional hardening region is observed in the fine grained materials.
- 456 5. Cyclic tensile testing of the different grain size specimens at stresses below
457 the yield stress revealed the accumulation of strain with each cycle. This
458 was also used to determine the effect of grain size on the twin nucleation
459 stress.
- 460 6. The progression of sub yield strain accumulation proceeds in a manner
461 consistent with sub yield twinning being hardened in a conventional Hall-
462 Petch manner.
- 463 7. The twin nucleation stress was found to increase with decreasing grain size.
464 The critical twin stress at the single crystal limit was determined to be
465 ~ 50 MPa.
- 466 8. A larger amount of strain is accumulated per cycle in the coarse grain
467 material compared to the fine grain material. It is suggested that this is
468 due to the formation of thicker twins in the coarse grain material.
- 469 9. SEM analysis of each grain size material deformed to 5% engineering strain
470 revealed thicker deformation twins present in the coarser grain material.
- 471 10. TEM examination of the finest and coarsest grain size specimens reinforced
472 the SEM observations. The dislocation arrangement was also determined
473 to be affected by the grain size.
- 474 11. A larger fraction of stacking faults was observed in the coarse grain material
475 indicating the relative ease for twin formation in coarse grained TWIP steels
476 compared to a fine grain material.

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483 **A. Weibull Cumulative Distribution Functions**

484 It is common practice to plot grain size distributions from EBSD data in the
485 form of binned histograms. However, this does not allow the easy interpretation
486 of the grain size distribution in a statistically meaningful manner *i.e.* the mean,
487 standard deviation or kurtosis. Furthermore, the number of grains sampled
488 and hence the significance of any anomalies which for example may reveal a
489 multimodal grain distribution are also unclear. Finally, the unit of the frequency
490 axis is also often unclear, *i.e.* whether it is number or area normalised, making
491 it difficult to compare distributions.

492 In the present study, we fit a distribution function to the cumulative distri-
493 bution function (CDF) using a Weibull smoothing method, Figure 10(a). This
494 allows the probability distribution to be plotted in a manner that permits the
495 comparison between microstructures in both a visual and statistical fashion,
496 Figure 10(b). We have chosen to use a Weibull function in the current analysis,
497 but we acknowledge that the choice of function should ultimately be placed on
498 a theoretically sound foundation, which would be a useful topic for further work
499 based, *e.g.* on recrystallisation modelling [42].

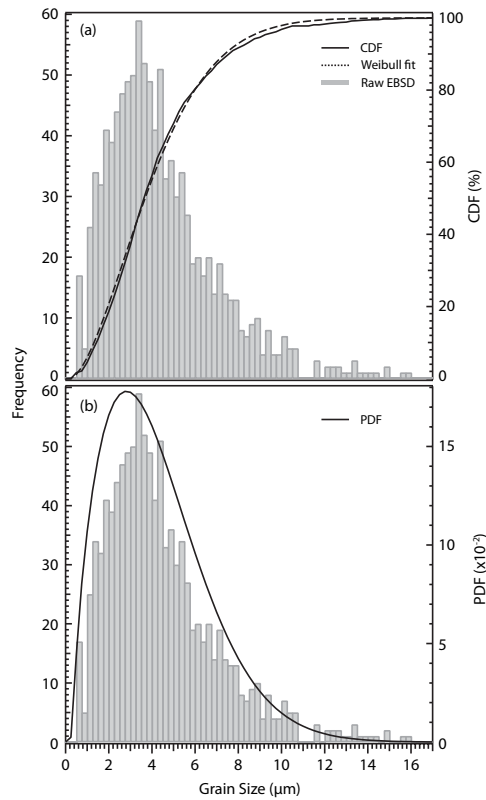


Figure 10: Weibull smoothing procedure where (a) the cumulative distribution function (CDF) of the raw EBSD data is fitted using a Weibull function and (b) smoothed probability density function (PDF) is plotted using a derivative of the fitted Weibull.

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