

A model for flow in the Chalk unsaturated zone incorporating progressive weathering

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Abstract

Groundwater from unconfined chalk aquifers constitutes a major water resource in the UK. The unsaturated zone in such systems plays a crucial role in the hydrological cycle, determining the timing and magnitude of recharge, and the transport and fate of nutrients. However, despite more than three decades of study, our physical understanding of this system is incomplete. In this research, state of the art instrumentation provided high temporal resolution readings of soil moisture status, rainfall and actual evaporation from two sites in the Pang and Lambourn catchments (Berkshire, UK), for a continuous two year period (2004/5). A parsimonious, physically based model for the flow of water through the Chalk unsaturated zone, including a novel representation of the soil and weathered chalk layers, was developed. The parameters were identified by inverse modelling using field measurements of water content and matric potential. The model was driven by rainfall and evaporation data, and simulated fluxes throughout the profile (including the discrete matrix and fracture components), down to the water table (but not the water table response). Results

showed that the model was able to reproduce closely the observed changes in soil moisture status. Recharge was predominantly through the matrix, and the recharge response was strongly attenuated with depth. However, the activation of fast recharge pathways through fractures in the Chalk unsaturated zone was highly sensitive to rainfall intensity. Relatively modest increases in rainfall led to dramatically different recharge patterns, with potentially important implications for groundwater flooding. The development and migration of zero flux planes with time and depth were simulated. The simulations also provided strong evidence that, for water table depths greater than 5 m, recharge fluxes persist throughout the entire year, even during drought conditions, with important implications for the calculation of specific yield from baseflow estimates and the representation of recharge in groundwater models.

Key words: Chalk; Groundwater recharge; Richards equation; Unsaturated zone; Weathered rock; Zero flux plane

1 Nomenclature

2	C	specific capacity, [m^{-1}]
3	h	hydraulic head, [m]
4	K	hydraulic conductivity, [m/d]
5	K_r	relative hydraulic conductivity, [-]
6	K_s	saturated hydraulic conductivity, [m/d]
7	L	Conductivity exponent parameter, [-]
8	L_{rd}	depth above which 63 % of root density is located, [m]

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9	OF_1	objective function for calibration of $\theta(\psi, z)$ relationship, [-]
10	OF_2	objective function for calibration of $K(\psi)$ parameters, [-]
11	q	Darcian velocity, [m/d]
12	r	pore radius, [m]
13	r_d	root distribution function, [-]
14	r_s	Feddes root stress function, [-]
15	S	sink term in Richards' equation, [d ⁻¹]
16	S_s	specific storage, [m ⁻¹]
17	S_e	effective saturation, [-]
18	$S_{e,1}$	effective saturation related to parameter ψ_1 , [-]
19	$S_{e,2}$	effective saturation related to parameter ψ_2 , [-]
20	t	time, [d]
21	w_f	fracture domain volume fraction, [-]
22	$w_{f,0}$	fracture domain volume fraction at the ground surface, [-]
23	$w_{f,\infty}$	fracture domain volume fraction in the deep Chalk, [-]
24	z	depth below groundlevel, [m]
25	Z_α	CUZ model shape parameter, [m ⁻¹]
26	Z_β	CUZ model shape parameter, [m]
27	α	general relative hydraulic conductivity model parameter, [-]
28	β	general relative hydraulic conductivity model parameter, [-]
29	θ	volumetric moisture content, [m ³ /m ³]
30	θ_s	saturated water content, [m ³ /m ³]
31	θ_r	residual water content, [-]
32	σ	Kosugi parameter, [m]
33	Ψ	pressure head, [m]

34	Ψ_0	Kosugi parameter, [m]
35	Ψ_1	modified Kosugi model parameter, [m]
36	Ψ_2	modified Kosugi model parameter, [m]
37	$\Psi_{1,\infty}$	modified Kosugi model parameter in the deep Chalk, [m]
38	$\Psi_{0,\infty}$	modified Kosugi model parameter at the ground surface, [m]
39	Ψ_{an}	matric potential threshold for anaerobiosis , [m]
40	Ψ_d	matric potential below which plant water stress begins, [m]
41	Ψ_w	wilting point, expressed as a matric potential, [m]

42

43 **1 Introduction**

44 The importance of Chalk aquifers as a water resource in north west Europe (specif-
 45 ically in the UK, northern France, northern Germany and Belgium, Downing et al.,
 46 1993; Kloppmann et al., 1998; Brouyère et al., 2004; Pinault et al., 2005) and as a
 47 potentially significant hydrological pathway for contaminants in Israel (Nativ and
 48 Nissim, 1992; Nativ et al., 1995; Dahan et al., 1998, 1999) is well established. This
 49 paper focuses on a Chalk catchment in south east England where Chalk aquifers
 50 represent approximately 80% of total water supply (Downing, 1998). In this region
 51 much of the Chalk is overlain by a thick unsaturated zone which can be in excess
 52 of 100 m (Jackson et al., 2006).

53 The need for a good hydrological understanding of the Chalk unsaturated zone
 54 (CUZ) is well recognised, as it is the main control for aquifer recharge (Ragab
 55 et al., 1997; Bradford et al., 2002; Ireson et al., 2006; Lee et al., 2006), contami-

56 nant transport (Haria et al., 2003; Jackson et al., 2006, 2007; Mathias et al., 2005,
57 2006, 2007; Gooddy et al., 2007) and groundwater flooding (Jacobs, 2006; Pinault
58 et al., 2005). Chalk comprises a fine-grained porous matrix (high porosity 20-45%,
59 low permeability $< 10^{-2}$ m/day) intersected by a fracture network (low porosity
60 $< 2\%$, higher permeability $> 10^{-2}$ m/day) (Price et al., 1993). Conceptual models
61 of how water and solutes move through the CUZ have evolved over past decades
62 (Smith et al., 1970; Wellings and Bell, 1980; Price et al., 1993, 2000). In unsat-
63 urated conditions pore water pressure is sub-atmospheric. Hence the fractures and
64 matrix may be partially saturated and hydraulic pathways within and between these
65 may be restricted/discontinuous. Wellings and Bell (1980) suggested that the frac-
66 tures become activated when the matric potential exceeds a threshold of -0.5 m.
67 Price et al. (2000) proposed that, in addition to providing a flow pathway, water
68 held by capillary tension on fracture walls could be an important means of stor-
69 ing and supplying water to groundwater. Such conceptual models can be tested
70 using mathematical models to reproduce field observations. The first mathemati-
71 cal models of the CUZ focused on solute transport and assumed steady state flow
72 (Young et al., 1976; Oakes, 1977; Oakes et al., 1981; Barker and Foster, 1981).
73 More recently, there has been an increased interest in transient flow models based
74 on Richards' equation (Mathias et al., 2006; Brouyère, 2006; Van den Daele et al.,
75 2007).

76 Mathias et al. (2006) considered a flow and transport model comprising of a 1D
77 fracture coupled to a 2D matrix block. Flow in the fractures and matrix was de-
78 scribed by Richards' equation. Parameters describing the matrix were obtained
79 from the mercury intrusion data (Price et al., 2000). Fracture parameters were in-

80 ferred from hydraulic conductivity – matric potential relationships observed in the
81 field (Wellings, 1984; Hodnett and Bell, 1990; Cooper et al., 1990; Mahamood-ul-
82 Hassan and Gregory, 2002). In order to reproduce solute profiles which were con-
83 sistent with previous experimental observations (Smith et al., 1970; Oakes et al.,
84 1981; Barraclough et al., 1994) it was necessary to attenuate the input of rainfall,
85 by means of a soil layer. The importance of the soil layer as a control on flow in the
86 CUZ had previously been postulated by Cooper et al. (1990).

87 Mathias (2005, p. 87) demonstrated that if the temporal resolution of the driving
88 rainfall data is daily or coarser, the time for pressure equilibrium between the frac-
89 tures and matrix becomes negligible. Consequently, if solute transport is not con-
90 sidered, flow in the unsaturated Chalk can be represented using the equivalent com-
91 posite medium (ECM) approach of Peters and Klavetter (1988). Richards' equation
92 is solved as for a single porous medium, but the hydraulic properties (i.e. relation-
93 ships of water content, θ , hydraulic conductivity, K , and specific storage, C , with
94 matric potential, ψ) are defined for both domains. The same approach was applied
95 to the CUZ of the Hesbaye region in Belgium by Brouyère (2006). The main dif-
96 ference was that whereas Mathias (following Peters and Klavetter, 1988), obtained
97 composite hydraulic properties by summing the volume averaged characteristics
98 for each domain, Brouyère defined a threshold matric potential, ψ_j , below which
99 the medium is defined by matrix properties, and above which it is defined by frac-
100 ture properties. This difference is merely practical – both models could be used to
101 obtain an identical result, albeit with different parameter values.

102 Van den Daele et al. (2007) used a modelling package called MACRO (Larsbo
103 et al., 2005), under the assumption that chalk fractures are analogous to soil macro-

104 pores and the chalk matrix is analogous to the soil matrix. Flow in the matrix was
105 controlled by Richards' equation, whereas flow in the fractures was represented by
106 the kinematic wave equation; thus assuming that capillary effects in the fracture
107 domain are negligible. Van den Daele et al. (2007) applied this model to the Fleam
108 Dyke lysimeter in Cambridgeshire. The 5 m deep lysimeter contained a 20 cm soil
109 layer overlying weathered chalk which gradually blended into undisturbed chalk at
110 the base. To accommodate this vertical heterogeneity, 5 discrete layers were consid-
111 ered, including a soil layer, weathered chalk layers and an undisturbed chalk layer.
112 The model was calibrated against two years of θ , ψ and drainage flux data. How-
113 ever, the range of $\theta(\psi)$ data for the undisturbed chalk (see their Figure 2) appears to
114 be unsuitable to adequately parameterise a soil moisture characteristic relationship
115 (a problem explicitly acknowledged by Cassiani and Binley, 2005, see their discus-
116 sion of unsaturated flow parameter identification). It was noted that the weathered
117 chalk layers were even harder to characterise. The model was unable to simulate
118 the observed matric potential data (see their Figure 4) although it did exhibit good
119 correspondence with the observed water content and drainage flux data.

120 The use of a kinematic wave in this context essentially represents a simplifying
121 assumption, which is made because of the general problem of a lack of information
122 concerning the hydraulic properties of the fractures/macropores close to saturation
123 (Larsbo et al., 2005). However, this assumption is deemed unnecessary in this study
124 where a high quality field dataset is available to parameterise the fracture properties.
125 Furthermore, there is a well established precedent for using Richards' equation
126 to represent unsaturated fracture flow in fractured rocks as a whole (Wang and
127 Narasimhan, 1985; Peters and Klavetter, 1988; Kwicklis and Healy, 1993; Gerke

128 and van Genuchten, 1993, 1996; Liu et al., 1998, 2003; Doughty, 1999).

129 The three studies described above illustrate an increasingly more sophisticated rep-
130 resentation of near surface properties: a decoupled soil layer (Mathias et al., 2006);
131 a coupled discrete soil layer (Brouyère, 2006); a multi-layered approach which
132 represents both soil and weathered Chalk (Van den Daele et al., 2007). We seek
133 to improve on these by accounting for soil and weathered chalk overlying consol-
134 idated rock using continuous functions to describe the vertical variations in physi-
135 cal and hydraulic properties, using the ECM approach (Mathias, 2005; Peters and
136 Klavetter, 1988). The model will be conditioned and tested against observed field
137 hydrological data, and techniques to assess parameter identifiability and model un-
138 certainty (similar to those used by Cassiani and Binley, 2005) will be applied.

139 **2 The field monitoring scheme**

140 In this study, we exploit data from an extensive set of instrumentation installed
141 in the Pang and Lambourn catchments (Berkshire, UK) during the LOCAR pro-
142 gramme (an overview of which is given by Wheater et al., 2006). In particular, we
143 focus on data from two field sites where there is a deep CUZ above the unconfined
144 Chalk aquifer: West Ilsley (WI) and Warren Farm (WF) (see Ireson et al., 2006, for
145 detailed site descriptions). These sites are located on the Seaford Chalk formation,
146 where marl bands are rare, but flint nodules may be present, and this is overlain by
147 a very shallow soil layer. Frequent measurements of soil moisture status (θ and ψ)
148 were taken over a range of depths down to 3 m, along with measurements of water
149 table response, rainfall and actual evaporation. Discussion of the instrumentation

150 was provided by Ireson et al. (2005) and the data were analysed in detail by Ireson
151 et al. (2006). In this study, we focus on the following aspects of the data:

- 152 • Coincident (in time and depth) readings of θ and ψ are used to construct observed
153 soil moisture characteristic relationships at four depths (0.2, 0.4, 0.6 and 1.0 m),
154 which are then used to optimise a number of model parameters;
- 155 • Rainfall and actual evaporation data are used to drive the model, applied at the
156 upper boundary and over the rooting depth, respectively;
- 157 • Time series of observed θ and ψ data at various depths (≤ 3 m) are used to opti-
158 mise the remaining model parameters (data from 2004), and to gauge the model
159 performance (data from 2005).

160 The water table response is influenced by both drainage from the unsaturated zone
161 and lateral flow processes in the saturated zone. Therefore the observed water table
162 response cannot be reproduced by a one-dimensional model, but can be used to
163 draw inferences about model performance and system behaviour.

164 **3 CUZ model development**

165 *3.1 Conceptualisation of the profile*

166 Figure 1 shows photographs of the upper 2.4 m of the unsaturated zone at West
167 Ilsley. These show major changes in the structure and composition of material in
168 the profile over this depth. In the upper 0.2 m is a soil layer, below which there
169 is chalk which becomes progressively less weathered with increasing depth. The
170 objective of this study is to develop a physically realistic representation of these

171 changing properties in a continuous, progressive and parsimonious manner. The
172 progressive nature of these changes is evident from the statistical properties of the
173 neutron probe θ data in Figure 2. Plant water uptake will strongly affect the water
174 content in the upper 0.2 m. Below this it is reasonable to associate the degree of
175 variation of θ at each depth with the volume of the fracture domain, under the
176 assumption that since the matrix will generally remain saturated by capillary forces
177 (Price et al., 1993), any changes in water content occur in the fractures. Similarly,
178 the trend in the minimum water content can be associated with the volume of the
179 matrix, assuming that under these conditions the matrix potential would be too low
180 for the fractures to hold water, yet higher than the air entry pressure of the matrix
181 (i.e. the matrix is saturated). On this basis it is assumed that as we approach the
182 surface the proportion of rock which is comprised of matrix will reduce, whilst that
183 which is fractures will increase.

184 As well as the changes in the relative proportions of each domain, if we are to rep-
185 resent the soil layer using properties which are scaled as a function of depth, the
186 hydraulic properties of one or both of the domains must also be modified. Weather-
187 ing of the Chalk is conceptualised here as enhanced fracturing of the matrix, with
188 the properties of the intact matrix blocks remaining unchanged. Therefore, it is pro-
189 posed that the fracture pore size distribution should be progressively modified as
190 a function of depth, whilst the matrix pore size distribution be kept constant. In
191 this way, the material in the near surface comprises relatively small, porous matrix
192 blocks, surrounded by fractures and soil macropores which constitute the fracture
193 domain. In the consolidated Chalk the fracture domain is made up of a number
194 of discrete and visible fractures (with fracture apertures of the order of 30 μm ,

195 Wellings, 1984). In summary, the qualitative conceptualisation of the CUZ profile
196 is characterised by the following changes as we approach the surface:

- 197 (1) A reduction in the proportion of the domain taken up by the matrix
- 198 (2) An increase in the proportion of the domain taken up by the fractures
- 199 (3) The fracture pore size distribution is modified such that there are more frac-
200 tures, with a wider range of apertures
- 201 (4) The matrix pore size distribution is unchanged

202 It is intended that this conceptualisation be applied to the entire unsaturated chalk
203 profile, from the water table up to the ground surface, including the soil layer in
204 the top 0.2 m. The physical basis for the extrapolation beyond the weathered chalk
205 and into the soil is questionable. However, the advantages of doing this are consid-
206 erable, including the elimination of a sharp, artificial discontinuity in the hydraulic
207 properties where the soil meets the chalk, and a potentially significant reduction
208 in the number of parameters required to characterise the entire profile. It is further
209 assumed that there are no significant lithological features (such as marl bands) or
210 karst features (such as those described by Allshorn et al., 2007) in the CUZ, which
211 would significantly complicate the recharge processes. This assumption is not un-
212 reasonable at the chosen field sites.

213 3.2 *Quantitative representation of the profile*

214 The CUZ model requires parametric relationships for $\theta(\psi)$, $C(\psi)$ and $K(\psi)$ for
215 both the matrix and the fracture domains, which can be modified as a function of
216 depth in order to achieve the four characteristics identified above. Formulae have

217 been proposed (Kosugi, 1994, 1996) which explicitly relate hydraulic properties to
 218 pore size distributions, based on the assumption that the pore radii are lognormally
 219 distributed and inversely proportional to the negative pore capillary pressure. In
 220 this study the two parameter Kosugi model (Kosugi, 1996) (hereafter referred to
 221 as the KS model) is adopted, as defined by items 5 to 8 in Table 1, where θ_r [-]
 222 and θ_s [-] are the residual and saturated water content, respectively, ψ_0 [L] is the
 223 mode pore capillary pressure, σ [L] is the standard deviation of the $\ln(\psi)$ distribu-
 224 tion, K_s [LT^{-1}] is the saturated hydraulic conductivity and L [-] is a free parameter,
 225 referred to here as the conductivity exponent, whose value may be positive or nega-
 226 tive (Mualem, 1976; Schaap and Leij, 2000). An investigation of different possible
 227 forms of the hydraulic conductivity relationship was performed, which is described
 228 in Appendix A. Based on the results, a non-conventional form of the general re-
 229 lationship for hydraulic conductivity (Hoffmann-Reim et al., 1999; Kosugi, 1999)
 230 was adopted (given by item 8 in Table 1, which can be compared to the Mualem
 231 form presented by Kosugi, 1996).

232 A useful feature of the KS model is that the parameters (ψ_0 and σ) can be obtained
 233 analytically from any two known points on the $S_e(\psi)$ curve. Hence the model can
 234 be defined by arbitrarily selecting two effective saturation values, $S_{e,1}$ and $S_{e,2}$, and
 235 defining the pore capillary pressure at these points, ψ_1 and ψ_2 , which are treated
 236 as the new model parameters. Given that $S_{e,1} = 0.5 + 0.5\text{erf}(-x_1/\sqrt{2})$ (Equation
 237 5 in Table 1), it follows that $x_1 = -\sqrt{2}(\text{erf}^{-1}[2S_{e,1} - 1]) = \ln(\psi_1/\psi_0)/\sigma - \sigma$ (and
 238 likewise for $S_{e,2}$ and x_2). After some further manipulation it can be shown that σ
 239 and ψ_0 can be found from items 3 and 4 of Table 1.

240 Both the matrix and the fracture domains require six parameters. Bulk properties

241 are obtained by summing the volume averaged properties of each domain, as in
 242 items 9 to 11 of Table 1, where w_f is the fracture domain volume fraction (i.e.
 243 the volume of the fracture domain over the total volume, after Gerke and van
 244 Genuchten, 1993) and the superscripts f and m refer the fracture and matrix do-
 245 mains respectively. The fracture domain porosity is taken to be 100 % (i.e. the frac-
 246 ture domain is a void), so $\theta_s^f = 1.0$. The matrix domain volume fraction is given
 247 by $1 - w_f$. Therefore, the basic model structure consists of 13 parameters (6 KS
 248 parameters for each domain, plus w_f). The next step is to introduce vertical hetero-
 249 geneity into the model, by scaling some of the parameters as a function of depth,
 250 consistent with the qualitative conceptualisation of the profile, developed above.

251 In order to scale the size of each domain as a function of depth, z , (the first two scal-
 252 ing objectives above), a relationship for $w_f(z)$ must be established. It is proposed
 253 to use an s-shaped curve to achieve this, as shown in Figure 3a, which is defined by
 254 item 2 of Table 1, and requires four parameters: w_f as $z \rightarrow 0$, w_f as $z \rightarrow \infty$ and two
 255 ‘shape’ parameters (Z_α and Z_β).

256 The modification of the fracture pore size distribution with depth (the third scaling
 257 objective above), can be achieved by scaling either or both of the parameters ψ_1^f
 258 and ψ_2^f . It was decided that ψ_2^f , the matric potential at which the fracture domain is
 259 95% saturated (i.e $S_{e,2} = 0.95$), would be uniform with depth, and ψ_1^f , the matric
 260 potential at which the fracture domain is 5% saturated (i.e $S_{e,2} = 0.05$), would
 261 reduce with increasing depth (see Figure 4). The effect on the fracture aperture
 262 size distribution is that larger fracture apertures remain uniform with depth, whilst
 263 the number of smaller apertures (caused by weathering) increases nearer to the
 264 surface. Thus in the consolidated Chalk the fractures are not active until a certain

265 (relatively high) matric potential threshold is reached (as in the conceptualisation
 266 of Wellings, 1984), whilst in the near surface, the fractures can play a significant
 267 role (representing the macroporous structure of the soil) at lower matric potentials.
 268 ψ_1^f was scaled as a function of depth in the same manner as w_f , that is, using a four
 269 parameter s-shaped curve, given in item 1 of Table 1, and demonstrated in Figure
 270 3b. The parameters include ψ_1^f as $z \rightarrow 0$, ψ_1^f as $z \rightarrow \infty$ and, to minimise the total
 271 number of parameters, the same two ‘shape’ parameters, Z_α and Z_β , as above.

272 The scaling of w_f and ψ_1^f requires an additional 4 parameters, which means that the
 273 model now requires the specification of 17 parameters. Although this appears sub-
 274 stantial, it is worth noting that an equivalent model, which treats the soil/weathered
 275 chalk as a single discrete layer overlying homogeneous consolidated chalk would
 276 require significantly more parameters. For example, the model of Mathias et al.
 277 (2006) employs 20 parameters for the Chalk, 6 parameters to describe the soil and
 278 one additional parameter to specify the soil depth, i.e. 27 parameters in total.

279 3.3 Additional model details

280 The movement of water in unsaturated porous materials, due to gravitational and
 281 capillary forces, can be described using Richards’ equation (Richards, 1931). In
 282 order to achieve good numerical stability for a range of unsaturated and saturated
 283 conditions, the form of Richards’ equation advocated by Tocci et al. (1997) was
 284 selected as the governing equation:

$$285 \left(C(\psi) + S_s S_e(\psi) \right) \frac{\partial \psi}{\partial t} = \frac{\partial}{\partial z} \left(K(\psi) \left(\frac{\partial \psi}{\partial z} - 1 \right) \right) + S \quad (1)$$

286 subject to initial and boundary conditions:

$$\begin{aligned} \psi &= \psi_i(z), & 0 \leq z \leq z_N, & \quad t = 0 \\ K(\psi) \left(\frac{\partial \psi}{\partial z} - 1 \right) &= q_0(t), & z = 0, & \quad t \geq 0 \\ \psi &= 0, & z = z_N, & \quad t \geq 0 \end{aligned} \quad (2)$$

287
288 where the soil hydraulic properties are found using the system of equations in
289 Table 1, S is the sink term which accounts for transpiration losses in the root zone
290 (described below), and S_s [L^{-1}] is the specific storage. Strictly speaking paramete-
291 ter values of S_s should be defined separately for the fracture and matrix domains.
292 However, taking base values of $S_s^f = 10^{-5} \text{ m}^{-1}$ and $S_s^m = 10^{-6} \text{ m}^{-1}$ (Mathias et al.,
293 2006), a sensitivity study varying these parameters over two orders of magnitude
294 found them to be insensitive, and therefore the base values were taken. The gov-
295 erning equation is applied to a one dimensional profile extending from the ground
296 surface ($z = 0$) to just below the water table ($z = z_N$). It is solved numerically using
297 the method of lines, with a block centred finite difference grid in space, integrated
298 in time using the ordinary differential equation solver, ODE15s, available in MAT-
299 LAB. Initial conditions, $\psi_i(z)$, were obtained from observations in the top 4 m, and
300 assuming that below this ψ increases linearly with depth to a value of zero at the
301 water table. As it is not possible to reproduce the water table response with a one-
302 dimensional model of the unsaturated zone, a fixed water table boundary condition
303 was used at the lower boundary, which was located just below the lowest observed
304 water table level. A sensitivity study indicated that over a range of water table
305 depths from 40 to 75 m, the simulated flux above the water table was insensitive
306 to the fixed water table elevation, which gives us some degree of confidence in the

307 use of this assumption when using the CUZ model to simulate recharge patterns.

308 For the Chalk, it is typically assumed that there is no surface runoff, due to the high
309 infiltration capacity of the medium (Smith et al., 1970; Foster, 1975). Therefore,
310 we can disregard the potential for ponding to occur during high intensity rainfall,
311 and hence deal with rainfall as a straightforward specified flux boundary condition
312 at the top of the profile.

313 Actual evaporation, AE, was calculated as the residual of the surface energy balance
314 using measurements of sensible heat flux by eddy correlation, net radiation and soil
315 heat flux over grass, at Warren Farm (a similar procedure was applied by Roberts
316 et al., 2005). There is no way to partition the AE between evaporation (from the
317 soil surface or interception) and transpiration, so it was assumed that all AE comes
318 from transpiration (i.e. all rainfall enters the soil, and all AE is extracted from the
319 soil via the plant roots). The distribution of uptake from the soil was determined
320 using a modified version of the Feddes et al. (1976) root extraction function. The
321 method is based on the root distribution over depth, $r_d(z)$ and a plant water stress
322 function, $r_s(\psi)$, given by (Feddes et al., 1976)

$$r_s(\psi) = \begin{cases} 0, & \psi > \psi_{an} \\ 1, & \psi_{an} \geq \psi > \psi_d \\ 1 - \frac{\psi - \psi_d}{\psi_w - \psi_d}, & \psi_d \geq \psi > \psi_w \\ 0, & \psi_w \geq \psi \end{cases} \quad (3)$$

323

324 where ψ_{an} , ψ_d and ψ_w are water stress thresholds, assumed to have values of -
 325 0.5 m, -4 m and -150 m respectively. It is assumed here that the plant roots are
 326 exponentially distributed with depth, such that

$$327 \quad r_d(z) = \frac{\exp(-z/L_{rd})}{L_{rd}} \quad (4)$$

328 where L_{rd} is the depth above which 63% of plant root density is located, and is
 329 taken to be 0.2 m. To ensure the total volume of AE was removed from the soil, the
 330 water stress function was normalised, and hence the sink term at a given depth and
 331 matric potential, $S(\psi, z)$, is found from:

$$332 \quad S(\psi, z) = \frac{r_s(\psi)r_d(z)}{\int_0^{L_r} r_s(\psi)r_d(z)dz} .AE \quad (5)$$

333 3.4 Parameter identification

334 Of the 17 model parameters, some are amenable to laboratory investigation (e.g.
 335 parameters associated with the chalk matrix, which can be obtained from core sam-
 336 ples) and some are insensitive. The remaining parameters are associated with the
 337 fracture aperture distribution and the nature of the soil and weathering at a partic-
 338 ular site and hence require optimisation. Table 2 summarises how each parameter
 339 value was identified.

340 Parameters requiring optimisation were identified in two stages. The first stage ap-
 341 plies to those parameters associated with the $\theta(\psi, z)$ relationship, for which obser-
 342 vations are available at four discrete depths in the top 1.0 m of the profile. The sec-
 343 ond stage only involves parameters strictly associated with the $K(\psi)$ relationship,
 344 and can only be identified by inverse modelling. The basic calibration strategy was

345 similar to that applied by Cassiani and Binley (2005), that is, a simple Monte Carlo
 346 approach whereby a number of parameter realisations are generated by randomly
 347 selecting parameter values from a uniformly distributed range, specified *a-priori*.
 348 A set of model outputs are generated for each realisation, and an objective function
 349 (the normalised root mean squared error) is used to judge the model performance
 350 against some observed data. A subset of the realisations are classed as having good
 351 performance, and these were ranked in terms of performance to establish 5 and 95%
 352 uncertainty bounds on the model output. Because the first stage simply involves fit-
 353 ting parametric relationship to some data points, a large number of realisations are
 354 possible. The second stage, however, requires solving Richards' equation, and as
 355 such is computationally intensive. Consequently, fewer realisations are possible due
 356 to time constraints.

357 3.4.1 Model calibration stage one

358 For the two sites studied, simultaneous, frequent measurements of water content,
 359 θ^{obs} , and matric potential, ψ^{obs} , were available at 4 depths: 0.2, 0.4, 0.6 and
 360 1.0 m. To quantify the performance of the modelled water content, θ^{mod} for a given
 361 parameter set, the root mean squared error (RMSE) was used as an objective func-
 362 tion, such that

$$363 OF_1 = \sum_{j=1}^4 \sqrt{\left(\frac{1}{N_j} \sum_{i=1}^{N_j} (\theta_{i,j}^{mod} - \theta_{i,j}^{obs})^2 \right)} \quad (6)$$

364 where N_j is the number of $\theta - \psi$ observations for depth j , and OF_1 has units of
 365 water content [-]. Due to there being scatter in the observed data (thought to be
 366 caused by drift in the profile probe calibration) a subset of data was selected, which

367 defines the primary drying curve, by selecting the maximum observed water content
368 in each of a number of discrete ranges of matric potential.

369 There is some variability in the matrix porosity at different depths (indicated by the
370 variation in mean water content with depth in Figure 2), which is largely attributed
371 to the presence of flints. Although this has an effect on the total water content, its
372 effect on flow processes is assumed to be minimal, due to the fact that the matrix
373 generally remains close to or fully saturated, hence the portion of the pore space that
374 may actually dewater is small. For simplicity, a constant effective porosity for the
375 matrix domain, $(\theta_s - \theta_r)$, is assumed to apply throughout the profile (however, note
376 that the matrix domain fraction, $1 - w_f$, is reduced in the near surface), and a value
377 of 35 % was assumed, based on ranges presented by Price et al. (1993). However,
378 in order to correct for the apparent random variation, $\theta_{s,j}^m$, was also treated as a
379 free parameter for each depth during the optimisation. This effectively normalises
380 the observed water content, such that the final matrix residual and saturated water
381 content θ_r and θ_s are 0 and 35 % respectively.

382 *In situ* measurements of hydraulic conductivity have generally indicated the onset
383 of fracture flow to occur at a matric potential of around -0.5 m H₂O (Wellings,
384 1984; Hodnett and Bell, 1990; Cooper et al., 1990; Mahamood-ul-Hassan and Gre-
385 gory, 2002). To be consistent with this it is necessary to set $\psi_2^f \geq -0.5 \geq \psi_{1,\infty}^f$.
386 Following some sensitivity analysis, it was decided that the value of ψ_2^f should be
387 set arbitrarily to -0.1 m (which means in practice that it is assumed that the frac-
388 tures are 95 % saturated at $\psi = -0.1$ m). This reduces the number of parameters
389 in the optimisation from 10 to 9, which significantly reduces the number of model
390 realisations required to accurately describe the objective function surface. To iden-

391 tify the remaining 9 parameters a 100,000 realisation Monte Carlo analysis was
392 performed. The results for both of the sites studied, in terms of model performance,
393 uncertainty and parameter identifiability are shown in Figures 5 and 6 (note, for
394 brevity we do not show the dotted plots for the four $\theta_{s,j}^m$ normalising parameters, as
395 these do not feature in the final model).

396 Plots of each parameter value against objective function, for the five optimised
397 model parameters are shown, for the best 1% of all model realisations. The pa-
398 rameters are somewhat more strongly identifiable at WF than at WI. However, in
399 both cases the optimum parameter set is rather isolated, indicating that potentially
400 a larger number of realisations could be beneficial.

401 The optimised $\theta(\psi, z)$ curves agree well with the observations at both sites, and at
402 all depths, with the exception of the dry range at 0.2 m and 1.0 m depth for Warren
403 Farm, where some observations fall outside the model 5-95% confidence interval.

404 The optimal parameter set identified for each site determines the nature of the frac-
405 ture pore size distribution as a function of depth. These are plotted (as probability
406 distribution functions and cumulative distribution functions at various depths) in
407 Figures 5 and 6. At WF, the parameters imply that the weathered Chalk layers are
408 distributed over 4 m depth, whilst at WI they are distributed over only 1 m. As a
409 result of this, we might expect there to be more attenuation of infiltrating rainfall at
410 WF.

411 3.4.2 Model calibration stage two

412 The remaining model parameters that require optimisation are the matrix and frac-
 413 ture saturated hydraulic conductivities, K_s^m and K_s^f , and the fracture conductivity
 414 exponent, L^f . These parameters were identified by inverse modelling, using the nu-
 415 merical CUZ model driven by observed rainfall and measured actual evaporation
 416 to reproduce observed time series of θ and ψ in the top 3 m of the profile, at the
 417 two sites. The normalised RMSE was used for the objective function, such that

$$418 \quad OF_2 = \sum_{j=1}^6 \sqrt{\left(\frac{1}{N_j} \sum_{i=1}^{N_j} \left[\frac{\psi_{i,j}^{mod} - \psi_{i,j}^{obs}}{\min(\psi_{i=1 \rightarrow N_j, j}^{obs})} \right]^2 \right)} + \sum_{j=1}^4 \sqrt{\left(\frac{1}{N_j} \sum_{i=1}^{N_j} \left[\frac{\theta_{i,j}^{mod} - \theta_{i,j}^{obs}}{\max(\theta_{i=1 \rightarrow N_j, j}^{obs})} \right]^2 \right)} \quad (7)$$

419 where j and i are indices in depth and time respectively and N_j is the number of data
 420 points at each depth, and OF_2 is dimensionless. Frequently logged observations of
 421 ψ are available at $z = 0.2, 0.4, 0.6, 1.0, 2.0$ and 3.0 m and of θ at $z = 0.2, 0.4, 0.6$
 422 and 1.0 m.

423 Using variations of the instantaneous profile method, values of saturated matrix
 424 hydraulic conductivity, ranging from 0.001 to 0.006 m/day and saturated fracture
 425 hydraulic conductivities, ranging from 0.05 to 0.5 m/day have been observed at
 426 various Chalk sites across the UK (e.g. Wellings, 1984; Cooper et al., 1990; Hod-
 427 nett and Bell, 1990; Mahamood-ul-Hassan and Gregory, 2002). Values of saturated
 428 fracture hydraulic conductivities observed from packer testing in the saturated zone
 429 range from 0.01 to 100 m/day (Allen et al., 1997, p. 53 and 59). This significant
 430 increase is believed to be for two reasons: firstly there is uncertainty about whether
 431 the instantaneous profile method ever observed fully saturated fractures; secondly,
 432 fractures in the saturated zone are likely to be enlarged due to continuous exposure

433 to flowing water.

434 The above values apply to the bulk fracture-matrix system, so the fracture hydraulic
435 conductivities need to be divided by w_f and the matrix hydraulic conductivities by
436 $1 - w_f$ to be compatible with item 11 of Table 1. Assuming a 1% fracture porosity, a
437 bulk hydraulic conductivity of 100 m/day requires a fracture hydraulic conductivity
438 of 10,000 m/day. Considering this, the following *a priori* ranges were assumed
439 for the Monte Carlo simulations: $2 \times 10^{-4} \leq K_s^m \leq 2 \times 10^{-3}$ m/day and $1.5 \leq$
440 $K_s^f \leq 1500$ m/day. For both K_s^m and K_s^f parameters were randomly sampled from
441 a log-uniform distribution, because the parameter ranges span one or more order
442 of magnitude. Finally, L^f is an empirical parameter. It was established that in this
443 case, negative values for L^f performed poorly, and therefore the parameter was
444 varied uniformly over an arbitrary range from 0 to 5.

445 We focus initially on WF, where two complete years of field data are available
446 (2004 and 2005). Only data from 2004 were used in the model calibration, so that
447 data from 2005 could be used for model verification. A 2500 realisation Monte
448 Carlo simulation was performed for the above parameter ranges. The results are
449 shown in Figure 7. The model is most sensitive to the value of K_s^m , but there is also
450 some sensitivity to K_s^f , with better performance for values at the low end of the
451 specified range. L^f appeared insensitive for values greater than around 2. There-
452 fore, a further 1000 realisation Monte Carlo simulation was performed for a refined
453 parameter range in terms of K_s^f , which was now varied from 1.5 to 15 m/day. These
454 results simulation are also shown in Figure 7. Similar results were found for K_s^m
455 (optimum value 5.3×10^{-4} m/day) and L^f (optimum value 4.1), albeit with an
456 improved optimum objective function value. This time K_s^f appeared to be largely

457 insensitive, though there was a slight preference for lower values, and an optimum
458 value of 2.8 m/day was found. Both of these analyses demonstrate that the model
459 performance is largely determined by K_s^m , for which the optimum value is around
460 0.5 mm/day.

461 The model performance with the optimum parameter set is plotted in Figure 8. The
462 model generally reproduces well the temporal pattern of ψ at all depths, although
463 it tends to underestimate the magnitude of the peaks at 0.2 m depth. In terms of the
464 cumulative change in water content (*cum.* $\Delta\theta$), again, the model generally performs
465 well in reproducing the general patterns. It is also able to reproduce the wetting
466 events. The main weakness appears to be that during dry periods the modelled $\Delta\theta$
467 tends to be more responsive than the observations, particularly at depths of 0.2 and
468 0.4 m.

469 For WI, observed data were only available for the first part of 2004, so model
470 simulations were only run for this year, and the calibration was performed for the
471 refined parameter range identified from the previous calibration exercises for WF.
472 Figure 9 shows the model performance plotted against parameter values. Again,
473 the dotted plots indicate that K_s^f and L^f are relatively insensitive, whereas K_s^m is
474 strongly identifiable, and hence largely determines the model performance. The
475 optimum value of K_s^m at WI was approximately 1 mm/day, i.e. double that at WF.
476 Figure 10 shows the performance of the optimum solution. For the data that are
477 available, the model performs well in terms of matric potential. In terms of water
478 content, the model performance is somewhat limited at 1.0 m depth, where it tends
479 to underestimate the changes in θ . There is no response in the observed data to
480 the large event on 9 July, which causes a perturbation in the simulated data. This

481 simulated response was caused by a single, high intensity, rainfall event, observed
482 at WF, which saw 48 mm of rain in a 12 hour period. This suggests that the rainfall
483 event recorded at WF did not occur at WI (at least not with the same magnitude).
484 Despite the fact that limited data were available, and concerns about the use of input
485 rainfall data from WF, the overall model performance at WI was considered to be
486 reasonable.

487 **4 CUZ Model application**

488 The CUZ model was conditioned to two field sites, and was seen to perform well
489 in reproducing the soil water dynamics in the top 3 m of the profile. Both config-
490 urations were run with two years of driving data (rainfall and actual evaporation),
491 from 1 January 2004 to 1 January 2006, which were cycled three times to eliminate
492 effects of the initial conditions. The results from the final cycles were analysed to
493 gain insights into the hydrological processes occurring throughout the chalk pro-
494 files.

495 It was noted above that the first stage of the calibration resulted in different patterns
496 of soil/weathered chalk layers at each of the two sites. Figure 11 shows how these
497 differences impact upon the changes in water content throughout the near surface
498 profile. The changes in water content with depth essentially match the pattern of
499 fracture porosity with depth. At WF, the changes are gradual, over a depth of about
500 2 m, whereas at WI the changes occurs rapidly in the top 1 m. In both cases, the
501 profiles of matrix and fracture water content show that, in the consolidated chalk,
502 on average the matrix is close to saturation and the fractures are close to being

503 empty. In the soil and weathered chalk layers, the fractures fill up to absorb the
504 high intensity infiltrating fluxes, whilst the matrix desaturates to satisfy evaporative
505 demand.

506 Since the model includes a fully coupled plant, the development and migration of
507 zero flux planes (ZFPs) throughout the profile over time can be constructed from the
508 high spatial and temporal resolution simulated ψ data. The simulations indicated
509 very similar ZFP patterns for both sites, so the result is only plotted for WF, in Fig-
510 ure 12, against the observed water table elevation. The result is entirely consistent
511 with the schematic pattern originally proposed by Wellings and Bell (1980). During
512 the summer months, ZFPs were developed to significant depths - down to 4.85 m
513 and 4.29 m in 2004, and 6.97 m and 6.25 m in 2005, at WF and WI respectively.
514 The plant roots have access to large amounts of storage in the Chalk matrix, and are
515 hence unlikely to become water stressed. This is supported by recent field measure-
516 ments of actual evaporation, which have shown that Chalk outcrops almost always
517 satisfy potential evaporation (Roberts et al., 2005, in particular see their comparison
518 of potential and actual evaporation at Bridgets Farm, which is the same site as War-
519 ren Farm, in their Figure 3). Below the ZFP water moves downwards to water table,
520 and in this study it was seen that the recharge fluxes were constant over time, despite
521 the fact that there was a significant drought. Traditionally, recharge in groundwater
522 models has been calculated using variations of the Grindley (1969) method (such
523 as MORECS, Thompson et al., 1981), which fails to recognise the presence of a
524 ZFP. This method predicts recharge only when the soil moisture deficit is zero, and
525 in some cases also assuming a bypass of precipitation to the aquifer (e.g. Rushton
526 et al., 1989; Ragab et al., 1997; Andrews et al., 1997; Bradford et al., 2002). The

527 findings of this study indicate that these methods are inappropriate for unconfined
528 Chalk aquifers. Such groundwater models are often calibrated against observed
529 groundwater elevation data, by modifying the aquifer parameters (i.e. transmissiv-
530 ity and specific storage). However, if the recharge is represented erroneously, the
531 aquifer parameters obtained may be wrong. Similarly, estimates of specific yield
532 made on the basis that during a period of sustained groundwater recession there is
533 no recharge are also liable to be wrong (as was found in the study by Lewis et al.,
534 1993a,b, and cited by Price et al., 2000). As such, our understanding of how the
535 aquifer will behave under different conditions (for example floods, droughts, cli-
536 mate change scenarios, new pumping wells, and so on) may be ill founded. The
537 capacity of the Chalk unsaturated zone to perennially supply water to the saturated
538 zone will have particularly important consequences for the catchment during pe-
539 riods of drought, potentially maintaining groundwater levels and river flows at a
540 higher level than would previously have been predicted.

541 A debate which has persisted in the literature, at least since the study by Smith et al.
542 (1970), is the role of the fractures in transmitting flow through the chalk unsaturated
543 zone (perhaps the key contributions to this debate come from Smith et al., 1970;
544 Foster, 1975; Wellings, 1984; Cooper et al., 1990; Hodnett and Bell, 1990; Price
545 et al., 2000; Mahamood-ul-Hassan and Gregory, 2002; Haria et al., 2003). Most
546 workers since Wellings (1984) have concluded that matrix flow is dominant, but
547 instances when conditions likely to initiate fracture flow (based on observations of
548 ψ exceeding the threshold associated with fracture activation) exist near the top of
549 the profile have been reported. Within the CUZ model, it is possible to explicitly
550 differentiate between matrix and fracture flow, since the hydraulic conductivity in

551 both domains is a unique function of ψ , and the head gradient can be extracted from
552 the simulated ψ distribution. Figure 13 shows simulated flow in the matrix and the
553 fractures throughout the profile at WF. Fracture flow occurred down to a depth of
554 1 m, but below this flow was solely transmitted through the matrix. Fracture flow
555 generally occurred in the winter, but also occurred in the top 0.5 m of the profile
556 following the large rainfall event in the summer of 2004. At WI the patterns were
557 very similar, although flow was slightly more attenuated than at WF.

558 At WF, for 2004/5, the water table fluctuated between around 32 and 44 m below
559 ground level. In the model, it was necessary to fix the water table at 40 m. The
560 simulated flux at 35 m, shown in Figure 13, can be assumed to be at least indicative
561 of the actual recharge flux. It can be seen that the recharge flux is attenuated such
562 that it is almost constant over time. At WI, where the water table is about 70 m deep,
563 a similar result is found, with the recharge signal even more strongly attenuated.
564 If the results of these simulations are to be believed, variations in the water table
565 elevation could not have been caused by changes in the recharge flux over time, and
566 hence must have been caused by changes in the lateral head gradient differential in
567 the saturated zone. We speculate that changes in the elevation of the water table
568 at the interfluvium may therefore be largely caused by the propagation of a pressure
569 wave laterally through the saturated zone, initiated by recharge occurring where the
570 unsaturated zone thickness is much less.

571 A serious limitation of the modelling study thus far is that it only considered data
572 from two years, 2004 and 2005. Comparing the rainfall and potential evaporation
573 data with the long term (30 years) record from the area, Figure 14, shows that apart
574 from two wet months (August and October 2004) the period generally received less

575 than average rainfall. In fact, the period of 2004-6 has been described by the De-
576 partment for Environment, Food and Rural affairs, DEFRA, as the worst drought
577 in southeast England in the last 100 years (DEFRA, 2007). Therefore it can be ex-
578 pected that more fracture flow occurs in more typical years. As no data were avail-
579 able to test the model during wetter conditions, a simple sensitivity analysis was
580 carried out. The rainfall intensity was increased progressively from 1% up to 20%
581 (that is, every single rainfall event in the 2 year period was increased by these scale
582 factors). The corresponding change in the occurrence of fracture flow throughout
583 the profile at WI is presented in Figure 15 where it can be seen that the occurrence
584 of fracture flow highly sensitive to the rainfall intensity. Increasing the rainfall by
585 4% resulted in fracture flow at depths of up to 10 m. An increase of 10% leads
586 to fracture flow throughout the entire profile. Figure 16 shows the recharge fluxes
587 (approximated by the simulated flux at 65 m depth) split into matrix flow, fracture
588 flow and total flow, for three rainfall scenarios (normal, 10% and 20% increase in
589 rainfall). For a 10% increase in rainfall 10% of the recharge is transmitted through
590 the fractures. Furthermore, fracture flow is no longer associated with any particu-
591 lar rainfall events, rather it is steady. For a 20% increase in rainfall, flow through
592 the fractures is more variable, and at certain times dominates the recharge signal.
593 In other words, at some point between 10 and 20%, the behaviour of the recharge
594 switched between a completely attenuated, steady response, and a seasonally dy-
595 namic response, comprising temporally discrete rapid recharge events.

596 The same analysis was performed for WF assuming this time that the recharge flux
597 can be approximated by the simulated flux at 30 m depth (see Figure 17). Here,
598 because of the reduced unsaturated zone thickness, there is less attenuation of the

599 recharge flux, and seasonal variations in recharge associated with fracture flow are
600 present for an increase in rainfall of just 10%, although matrix flow still generally
601 dominates the recharge pattern. It is notable that the fracture event for the 10%
602 rainfall scenario is delayed as compared with with 20% rainfall scenario. Hence,
603 the fractures in the CUZ model are not behaving as simple bypass pathways which
604 are either on or off. Rather they allow significant attenuation to occur, depending
605 on the intensity of rainfall. The inference from this work, therefore, is that recharge
606 models that use a simple bypass mechanism to represent fracture flow (Rushton
607 et al., 1989; Ragab et al., 1997; Andrews et al., 1997; Bradford et al., 2002) are of
608 questionable value.

609 **5 Conclusions**

610 In this paper we have presented a methodology for modelling flow in the chalk un-
611 saturated zone, accounting for vertical heterogeneity in the soil and weathered chalk
612 layers, using a parsimonious, Richards' equation based approach, which is able to
613 simulate flow in both the matrix and the fractures. A rational approach was taken
614 to identify appropriate model parameters. Those parameters which are liable to be
615 site specific and are not amenable to laboratory investigation were identified using
616 a Monte Carlo methodology. This was carried out in two stages. Firstly, parame-
617 ters associated with the soil moisture characteristic relationships at various depths
618 in the profile are identified using parametric relationships, which derive from the
619 Kosugi (1996) model. Secondly, parameters associated with hydraulic conductivity
620 are identified by inverse modelling, using a Richards equation based, equivalent
621 composite medium model. For the period studied the model was seen to perform

622 well in reproducing the observed soil moisture status (i.e. water content and ma-
623 tric potential) at the two field sites. However, there were two limitations. Firstly,
624 observed data were only available in the near surface, whereas interpretations were
625 made from the entire profile, down to the water table. Secondly, the period studied
626 was unusually dry, hence the model was not tested for the entire realistic range of
627 field conditions.

628 The model has enabled a number of useful insights to be made. During the period
629 studied zero flux planes developed in the top 7 m of of the profile. Above this, a
630 large amount of storage was potentially available to satisfy evaporative demand.
631 Beneath the ZFP the model suggests that the unsaturated zone drains continuously
632 into the water table. As a result, recharge is continuous throughout the year, albeit
633 with potentially differing rates, which has important implications for the represen-
634 tation of recharge in groundwater models, which have perhaps been largely over-
635 looked in the past. Failure to recognise this is also likely to lead to errors in the
636 estimation of parameters for unconfined chalk aquifers.

637 As the period studied was during a drought, a sensitivity study was carried out to
638 explore the effect of increased rainfall intensity. It was found that the frequency,
639 duration and depth of flow in the fractures, and consequently the recharge rates and
640 patterns were highly sensitive to rainfall intensity. A relatively moderate increase
641 in rainfall intensity (say of the order of 10%) may result in fracture flow being
642 initiated down to greater depths, even right down to the water table, and a seasonal,
643 or event based, recharge pattern may become evident. This non-linear behaviour,
644 is likely to be important when it comes to understanding groundwater flooding in
645 Chalk catchments (see also Pinault et al., 2005). However, the model developed in

646 this study has not yet been verified against field data under wet, or high rainfall,
647 conditions.

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866 **Appendix A: Alternative models for the relative hydraulic conductivity**

867 For variably saturated porous media, the hydraulic conductivity, $K(\psi)$ is given by
868 the product of the saturated hydraulic conductivity, K_s and the relative permeabil-
869 ity, $K_r(\psi)$. Various different forms have been proposed for the $K_r(\psi)$ relationship.
870 Kosugi (1996) considers both the Mualem (1976) and the Burdine (1953) rela-
871 tionships, which are both special cases of the general relationship proposed by
872 Hoffmann-Reim et al. (1999) (cited by Schaap and Leij, 2000)

$$873 \quad K(\psi) = K_s S_e^L \left(\frac{\int_0^{S_e} \psi^{-\alpha} dS_e}{\int_0^1 \psi^{-\alpha} dS_e} \right)^\beta \quad (8)$$

874 where L is a parameter typically related to the tortuosity and pore connectivity
875 (Mualem, 1976), and α and β are parameters which determine the form of the
876 relationship. For the Mualem relationship, $\alpha = 1$ and $\beta = 2$, and for the Burdine
877 relationship $\alpha = 2$ and $\beta = 1$. Furthermore, the Kozeny relationship (Brutsaert,
878 1967), is given by setting $\beta = 0$. The solution to Equation 8 will depend on the
879 form of the $S_e(\psi)$ relationship used. Kosugi (1999) showed that unlike with the van
880 Genuchten (1980) model, using the KS model for $S_e(\psi)$ Equation 8 can be solved
881 for non-integer parameter combinations of α and β , to yield the equation

$$882 \quad K(\psi) = K_s \cdot S_e^L \left(\frac{1}{2} \operatorname{erfc} \left[\operatorname{erfc}^{-1}(2S_e) + \frac{\alpha\sigma}{\sqrt{2}} \right] \right)^\beta \quad (9)$$

883 where σ is a KS model parameter. By treating L and α as fitting parameters, and
884 setting $\beta = 1$, Kosugi was able to obtain an improved description of the observed
885 $K(\psi)$ relationships for a large number of soils, as compared with the results using
886 the Burdine (1953) and Mualem (1976) models. In this study, since we do not have

887 observed $K(\psi)$ data for the Chalk, we examined the performance of the model
888 for a number of integer combinations of these parameters, including each of the
889 established relationships, as well as for the case when $\alpha = 1$ and $\beta = 1$, which
890 we refer to as the ‘modified Mualem’ model. For each model, a 1000 realisation
891 Monte Carlo simulation was performed to identify the parameters K_s^f , K_s^m and L^f
892 (as in the second calibration stage described earlier). Additional model parameters
893 were taken from the first calibration stage for WF, and observed matric potential
894 data from 2004 was used to gauge the model performance. Figure 18 shows the
895 performance for the optimum parameter set for each of the four model configura-
896 tions considered at 0.2, 1.0 and 3.0 m depth. Additionally, the RMSE between each
897 model result and the observed data was calculated. It was found that the modified
898 Mualem model performed significantly better than the alternatives, with an RMSE
899 of 4.4 (the Mualem, Kozeny and Burdine models had RMSE values of 5.15, 5.4
900 and 5.5 respectively), and was consequently adopted for the CUZ model developed
901 in this study. The integrated form of Equation 8 with the Kosugi relationship for
902 $S_e(\psi)$ and $\alpha = \beta = 1$, is given by item 8 in Table 1.

Inputs: ψ , z and 17 parameters:

$$\theta_r^m, \theta_s^m, \theta_r^f, \theta_s^f, w_{f,0}, w_{f,\infty}, \psi_1^m, \psi_2^m, \psi_{1,0}^f, \psi_{1,\infty}^f, \psi_2^f, K_s^m, K_s^f, L^m, L^f, Z_\alpha, Z_\beta$$

Initialise constants:

$$\begin{aligned} x_1 &= -\sqrt{2}(\operatorname{erf}^{-1}[2S_{e,1} - 1]); & S_{e,1} &= 0.05 \\ x_2 &= -\sqrt{2}(\operatorname{erf}^{-1}[2S_{e,2} - 1]); & S_{e,2} &= 0.95 \end{aligned}$$

These constants determine the value of S_e at ψ_1 and ψ_2 to characterise the pore size distribution. See section 3.2

	Description
1. $\psi_1^f = \psi_{1,\infty}^f + \frac{\psi_{1,0}^f - \psi_{1,\infty}^f}{1 + \exp(-Z_\alpha(z - Z_\beta))}$	Depth dependent Fracture pore size dist.(1) and domain size (2)
2. $w_f = w_{f,\infty} + \frac{w_{f,0} - w_{f,\infty}}{1 + \exp(Z_\alpha(z - Z_\beta))}$	

Repeat steps 3–8 for matrix and fracture domains

3. $\sigma = \frac{\ln\left(\frac{\psi_2}{\psi_1}\right)}{x_2 - x_1}$	Transform to get KS model parameters (3&4)
4. $\psi_0 = \frac{\psi_1}{e^{(x_1 + \sigma)\sigma}}$	
5. $S_e = 0.5 + 0.5\operatorname{erf}\left(-\frac{[\ln(\psi/\psi_0)/\sigma - \sigma]}{\sqrt{2}}\right)$	Effective saturation (S_e^m/S_e^f)
6. $\theta = \theta_r + S_e(\theta_s - \theta_r)$	Water content (θ^m/θ^f)
7. $C = \frac{\theta_s - \theta_r}{(2\pi)^{1/2}\sigma(-\psi)} \exp\left(-\frac{[\ln(\psi/\psi_0) - \sigma]^2}{2\sigma^2}\right)$	Specific capacity (C^m/C^f)
8. $K = K_s S_e^L \left[0.5 + 0.5\operatorname{erf}\left(\operatorname{erf}^{-1}[2S_e - 1] - \frac{\sigma}{\sqrt{2}}\right) \right]$	Hydraulic conductivity (K^m/K^f)

9. $\theta(\psi) = w_f \theta^f(\psi) + (1 - w_f) \theta^m(\psi)$	Bulk water content
10. $C(\psi) = w_f C^f(\psi) + (1 - w_f) C^m(\psi)$	Bulk specific capacity
11. $K(\psi) = w_f K^f(\psi) + (1 - w_f) K^m(\psi)$	Bulk hydraulic conductivity

Outputs: θ , C and K

For a description of all the parameters and variables, see the list of notation

Table 1

Summary of the equations used to obtain the hydraulic properties as a function of matrix potential and depth

Parameter ^b	Identification method	Parameter value	
		Warren Farm	West Ilsley
θ_r^m	Price et al. (2000); Mathias et al. (2006)	0	0
θ_s^m	Price et al. (2000); Mathias et al. (2006)	0.35	0.35
θ_r^f	Mathias et al. (2006)	0	0
θ_s^f	Mathias et al. (2006)	1	1
$w_{f,0}$	Optimisation [†]	0.12	0.14
$w_{f,\infty}$	Price et al. (1993); Mathias et al. (2006)	0.01	0.01
Ψ_1^m	Price et al. (2000); Mathias et al. (2006)**	-95.2 m	-95.2 m
Ψ_2^m	Price et al. (2000); Mathias et al. (2006)**	-14.1 m	-14.1 m
$\Psi_{1,0}^f$	Optimisation [†]	-40.1 m	-48.3 m
$\Psi_{1,\infty}^f$	Optimisation [†]	-1.29 m	2.99 m
Ψ_2^f	Fixed [‡]	-0.1 m	-0.1 m
K_s^m	Optimisation [§]	0.53 mm/day	1.01 mm/day
K_s^f	Optimisation [§]	2.83 m/day	1.73 m/day
L^m	Mualem (1976)***	0.5	0.5
L^f	Optimisation [§]	4.08	3.68
Z_α	Optimisation [†]	-1.4 m^{-1}	-9.5 m^{-1}
Z_β	Optimisation [†]	0.89 m	0.49 m

^b See nomenclature for parameter descriptions.

* Note the distinction between the fracture domain saturated water content, θ_s^f which is 100%, and the fracture domain fraction, $w_{f,0}$ and $w_{f,\infty}$, which are equivalent to bulk fracture porosity.

** These parameters were derived from existing (Brooks and Corey, 1966) model parameterisations for the English chalk matrix (Mathias, 2005; Mathias et al., 2006), which are based on the experimental data obtained by Price et al. (2000).

*** The matrix conductivity exponent parameter is insensitive.

[†] These parameters are classified as ‘soil moisture’ parameters, as they can be identified from the $\theta(\psi, z)$ data alone.

[‡] Arbitrarily fixing the value of this parameter did not affect the model performance, and doing so enabled the remaining parameter space to be explored more thoroughly.

[§] These parameters are classified as ‘hydraulic conductivity’ parameters, as they are associated with the hydraulic conductivity, and can be identified by inverse modelling.

Table 2
Summary of all model parameters

Fig. 1. Photographs of the Chalk profile over various depth horizons at West Ilsley

Fig. 2. The Chalk profile at West Ilsley, and statistical properties of the water content data recorded between January and July of 2003

Fig. 3. Parameter distributions with depth, to account for vertical heterogeneity in the near surface

Fig. 4. The effect of changing the fracture pore size distribution scaling parameters on the fracture effective saturation curve

Fig. 5. Stage one calibration results for Warren Farm soil moisture parameters

Fig. 6. Stage one calibration results for West Ilsley soil moisture parameters

Fig. 7. Dotty plots of the model performance (objective function to be minimised) against parameter values for Warren Farm. Note the different units of K_s^m and K_s^f .

Fig. 8. Model performance (optimum simulation) plotted against observed data for Warren Farm

Fig. 9. Dotty plots of the model performance (objective function to be minimised) against parameter values for West Ilsley. Note the different units of K_s^m and K_s^f .

Fig. 10. Model performance (optimum simulation) plotted against observed data for West Ilsley

Fig. 11. Simulated water content in the near surface, demonstrating attenuation by the soil and weathered Chalk layers

Fig. 12. Simulated zero flux planes and observed water table response for Warren Farm

Fig. 13. Simulated matrix and fracture flow throughout the profile at Warren Farm

Fig. 14. Monthly rainfall and potential evaporation recorded at the automatic weather station at Wallingford (SU461189)

Fig. 15. Fracture flow occurrence in the West Ilsley profile as rainfall is increased

Fig. 16. Recharge flux (approximated by the flux at 65 m depth) at West Ilsley as total rainfall is increased

Fig. 17. Recharge flux (approximated by the flux at 30 m depth) at Warren Farm as total rainfall is increased

Fig. 18. Optimum performance of the alternative models for hydraulic conductivity