

## Particle flux in the oceans: Challenging the steady state assumption

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### Key Points

- Increased biogenic silica fluxes with depth imply non-steady state conditions in the NE Atlantic in August 2009
- Taxonomic flux analysis and simple modelling support the non-steady state hypothesis
- By assuming steady state, net organic carbon supply to the dark ocean may be wrong by  $\leq 25\%$

**Abstract** Atmospheric carbon dioxide levels are strongly controlled by the depth at which the organic matter that sinks out of the surface ocean is remineralized. This depth is generally estimated from particle flux profiles measured using sediment traps. Inherent in this analysis is a steady state assumption; that export from the surface does not significantly change in the time it takes material to reach the deepest trap. However, recent observations suggest that a significant fraction of material in the mesopelagic zone sinks slowly enough to bring this into doubt. We use data from a study in the North Atlantic during July/August 2009 to challenge the steady state assumption. An increase in biogenic silica flux with depth was observed which we interpret, based on vertical profiles of diatom taxonomy, as representing the remnants of the spring diatom bloom sinking slowly ( $<40 \text{ m d}^{-1}$ ). We were able to reproduce this behaviour using a simple model using satellite-derived export rates and literature-derived remineralization rates. We further provide a simple equation to estimate 'additional' (or 'excess') POC supply to the dark ocean during non-steady state conditions, which is not captured by traditional sediment trap deployments. In seasonal systems, mesopelagic net organic carbon supply could be wrong by as much as 25% when assuming steady state. We conclude that the steady state assumption leads to misinterpretation of particle flux profiles

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when input fluxes from the upper ocean vary on the order of weeks, such as in temperate and polar regions with strong seasonal cycles in export.

### **Keywords**

Non-steady state, slow sinking particles, carbon export, biogenic silica, particle flux, sediment trap

### **Running Title**

Challenging the steady state assumption

### **Index terms**

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4806 Carbon cycling

0439 Ecosystems, structure and dynamics (4815)

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## **1. Introduction**

The depth at which carbon exported from the surface ocean is recycled is a key control over ocean atmosphere carbon dioxide partitioning [Kwon et al., 2009]. Some of the earliest systematic observations of particle flux were made by Martin et al. [1987] who deployed sediment traps off the west coast of the US in the mesopelagic zone (100-1000 m) and found that particle flux decreased with depth and was best described by a power-law function, now commonly known as the Martin curve. Even though this function was chosen purely as the best fit, with no underlying theoretical justification, it is now widely used to model flux attenuation in ocean models [e.g. in all of the models used in the Ocean-Carbon Cycle Model Intercomparison Project; Sarmiento and LeQuere, 1996; Doney et al., 2004] and to extrapolate measured fluxes to different depths.

Flux profiles acquired using sediment traps or estimated using the Martin curve are commonly assumed to be at steady state, partly because of reported rapid sinking speeds of aggregates: up to 400 m d<sup>-1</sup> for marine snow and 1800 m d<sup>-1</sup> for faecal pellets [see review by Turner, 2002]. Imagine a scenario of two traps collecting material at 50 m and 850 m depth. At an average sinking speed of 400 m d<sup>-1</sup>, material collected at 850 m would have been at 50 m around 2 days before. Unless there was a significant change in export at 50 m over 2 days then the steady state assumption is valid. Accurate simultaneous flux measurements are necessary to understand the

ocean carbon cycle, for example, when comparing the observed loss of sinking particulate carbon in the mesopelagic zone with estimates of heterotrophic respiration [Giering et al., 2014].

However, recent observations [e.g. Alonso-Gonzalez et al., 2010; Villa-Alfageme et al., 2014; Giering et al., 2016], including some from the cruise reported here [Riley et al., 2012], suggest that a significant fraction of the particle flux in the oceans may sink at much slower speeds ( $<40 \text{ m d}^{-1}$ ). If this is the case, particles may take several weeks to sink from the surface ocean to the deep ocean. For the steady state assumption to be valid, the timescale over which the flux leaving the surface ocean needs to be constant must therefore be considerably longer, extending from a few days to several months. Yet, a large proportion of the ocean exhibits seasonal variability in primary production and phytoplankton standing stocks, and may therefore also display variability in export flux. We hypothesize that in such systems the steady state assumption may not be valid.

We first undertake some theoretical modelling of how non-steady state situations coupled with slowly sinking particles may affect vertical particle flux profiles. We then discuss data from a field site which shows the predicted pattern. Finally, we derive a simple equation to estimate the 'additional' (or 'excess') POC supply to the dark ocean during non-steady state conditions, which is not captured by traditional analyses of sediment trap deployments.

## **2. Methods**

### **2.1 Non-Steady State model**

Some simple calculations help to examine how the presence of remnant particles at depth, owing to slow sinking speeds, affects particle flux profiles. The calculations are based on the assumption that the sinking speed and particle remineralization rates are constant with depth and during the time period considered. Commonly, particle flux is described by assuming a power-law function [Martin et al., 1987]. This function relates the flux at a target depth  $z$  ( $F_z$ ) to the flux at a shallower reference depth  $z_0$  using the equation  $F_z = F_{z_0} (z/z_0)^{-b}$ , where  $b$  describes the rate of flux attenuation. In this case, there is an implicit assumption that either sinking speed increases with depth or remineralization rate decreases as the particle sinks [e.g. Gehlen et al., 2006; Villa-Alfageme et al., 2014]. Hence, to use the Martin curve would be inconsistent with our assumptions. An assumption of constant remineralisation rate and sinking speed gives an exponential relationship rather than a power law one, and therefore this was used when

describing flux with depth. For the exponential model the rate of flux attenuation with depth is set by the choice of the two constants representing sinking speed and remineralisation rate.

It is however important to note that particles in the oceans do not have a universal sinking speed or remineralization rate [e.g. Alonso-Gonzalez et al., 2010; McDonnell and Buesseler 2010]. Rather, particles experience diversity in both rates, which could in itself cause the shape of idealized flux curves or departures therefrom [e.g. Boyd and Trull, 2007]. Along these lines, a previous study measured the distribution of sinking rates and demonstrated that the variability in these could have produced an apparent Martin curve without the any changes in sinking rates or remineralization rates because slow-sinking particles are lost at a shallower depths and only fast-sinking particles reach deeper waters [Trull et al., 2008]. Such selective losses would also explain the overall faster bulk sinking rates observed with increasing depth [Berelson 2002; Villa-Alfageme et al., 2014, 2016]. It is, however, widely accepted that the processes that shape particle flux profiles are more complex and involve, amongst others, physical or zooplankton-mediated fragmentation and aggregation [e.g. Conte et al., 2001; Trull et al., 2008; Burd and Jackson, 2009; Giering et al., 2014]. In order to explore the effect of non-steady state conditions, we refrain from adding these complexities and focus on a simple particle-flux scenario by assuming a fixed sinking speed and remineralization rate.

We simulated a series of downward flux profiles at 8-day intervals over a 60-day period. The input of matter from the mixed layer ('export flux',  $F_{in}$ ) was chosen to decrease over time at a rate ( $\alpha$ ) of  $1.2 \text{ mg m}^{-2} \text{ d}^{-1}$ , simulating declining phytoplankton standing stocks during the post-bloom phase [ $-1$  to  $-1.6\% \text{ d}^{-1}$ ; from Fig. 1b in Henson et al., 2013]. Export flux on our simulated observation date (at the end of the 60-day period) was  $28 \text{ mg m}^{-2} \text{ d}^{-1}$  ( $F_{in,0} = 28 \text{ mg m}^{-2} \text{ d}^{-1}$ ).  $F_{in}$  at a time  $\Delta t$  earlier than the simulated profile can thus be calculated as

$$F_{in,\Delta t} = F_{in,0} + \alpha \Delta t \quad (1),$$

Flux attenuation with depth was calculated by applying an exponential decay model with both a weak remineralization rate ( $r = 0.01 \text{ d}^{-1}$ ) and a strong remineralization rate ( $r = 0.20 \text{ d}^{-1}$ ). The flux observed at depth  $z_{\Delta t}$  will have left the surface at a time  $\Delta t = z_{\Delta t}/v$  days earlier, where  $v$  is the particle sinking speed ( $\text{m d}^{-1}$ ). The relation between the flux leaving the surface and that observed at depth  $z$  is therefore

$$F(z_{\Delta t}) = F_{in,\Delta t} \exp(-rz_{\Delta t}/v) = F_{in,\Delta t} \exp(-r\Delta t) \quad (2),$$

where  $F_{in,\Delta t}$  is the export flux  $\Delta t$  days before the observation (see Equation 1) and  $r$  is the flux attenuation coefficient ( $d^{-1}$ ). For our simulation, we assume that particles sink at a speed of 10  $m d^{-1}$  [Alonso-Gonzalez et al., 2010, Riley et al., 2012]. For a particle profile collected on one day, a trap deployed at 200 m below the mixed layer would capture particles that were exported from the surface 20 days ago, whilst a trap at 600 m below the mixed layer would collect 60-day-old particles. This ‘observed’ profile is hereafter referred to as ‘composite profile’.

## 2.2 Potential error when estimating net POC supply

We can further use Equations 1 and 2 to calculate both the apparent loss of POC from a composite profile - as traditionally observed during deployments of a fleet of 5 sediment traps - and the actual loss of POC. ‘Loss’ is here synonymous with the net POC supply between the top and bottom trap depths. We assume again that remineralization rate and sinking speed do not change over time.

To calculate the apparent loss between two depths a vertical distance  $\Delta z = v \times 1 \text{ day}$  apart, we simply subtract the observed flux at the deeper depth from the observed flux at the shallower depth (e.g.  $F_{50m} - F_{100m}$ ). Let us consider the two depths corresponding to those which material would sink to in  $\Delta t$  and  $\Delta t - 1$  days. The apparent loss ( $L_{app}$ ) is hence calculated using Equation 1 as

$$L_{app} = F_{in,\Delta t-1} \exp(-z_{\Delta t-1} r/v) - F_{in,\Delta t} \exp(-z_{\Delta t} r/v) \quad (3).$$

The actual loss ( $L_{act}$ ), on the other hand, is the difference between the observed flux at one reference depth and the flux at a shallower depth that originated from the same input flux,  $F_{\Delta t}$ :

$$L_{act} = F_{in,\Delta t} [\exp(-z_{\Delta t-1} r/v) - \exp(-z_{\Delta t} r/v)] \quad (4).$$

We can calculate the difference between the actual loss and the apparent loss at depth  $z$  ( $D_z = L_{act} - L_{app}$ ) by combining Equations 3 and 4, substituting with Equation 1, and simplifying to:

$$D_z = \alpha \exp(-z_{\Delta t-1} r/v) \quad (5).$$

Using the sum for finite geometric series, the sum throughout the water column is

$$S = \alpha [ 1 - \exp( -r * t_{max} ) ] / [ 1 - \exp( -r ) ] \quad (6),$$

where  $t_{max}$  is the length of time corresponding to the depth (relative to the mixed layer depth or photic zone depth) of the deepest sample divided by the assumed sinking speed ( $t_{max} = z_{max}/v$ ). This sum,  $S$ , allows quantification of how much the net organic carbon supply to the mesopelagic zone may be under- or overestimated when assuming steady state. This calculation can easily be applied wherever there is information about  $\alpha$ ,  $r$  and  $v$ .

### **2.2.1 Case study: Porcupine Abyssal Plain (PAP) site**

The theoretical model described above was applied to the PAP data (see section 2.3) and a sensitivity analysis was performed. The parameters were chosen to reflect likely conditions at the PAP site. During the 60 days prior to our cruise, POC export decreased according to satellite-derived data at a rate of  $1.1 \text{ mg C m}^{-2} \text{ d}^{-1}$  (linear regression;  $p < 0.01$ ,  $R^2 = 0.73$ ,  $n = 8$ ). POC has been shown to be remineralized at  $0.08\text{-}0.20 \text{ d}^{-1}$  [Iversen and Ploug, 2010], and average sinking speeds at the PAP site were likely between  $20$  and  $100 \text{ m d}^{-1}$  [Riley et al., 2012; Villa-Alfageme et al., 2014]. To represent uncertainties, we thus varied the parameters for temporal decrease in organic matter export from the mixed layer ( $\alpha$ ;  $0.5 - 1.3 \text{ mg C m}^{-2} \text{ d}^{-1}$ ), remineralization rate ( $r$ ;  $0.01 - 0.50 \text{ d}^{-1}$ ) and sinking speed ( $v$ ;  $10 - 200 \text{ m d}^{-1}$ ).

### **2.2.2 Seasonal examples using Atlantic and Pacific transects**

To illustrate how the potential range of error when assuming steady state may vary geographically and seasonally, we applied Equation 6 to satellite-derived export estimates (see section 2.4). We use climatological mean (2003-2014) estimates with temporal resolution of 8 days and spatial resolution of  $1^\circ$ . The change in export flux from the upper 100 m ( $\alpha$ ) was calculated at every grid cell for each time step as the gradient over an 8-day period.  $S$  was then calculated using a sinking speed  $v$  of  $40 \text{ m d}^{-1}$ ,  $z_{max}$  of  $400 \text{ m}$  (i.e.  $500 \text{ m}$  depth), and a remineralization rate scaled with a temperature coefficient ( $Q_{10}$ ) of  $3.5$  [Iversen & Ploug, 2013]:

$$r = 0.03 \text{ d}^{-1} \times Q_{10}^{[(T-4^\circ\text{C})/10^\circ\text{C}]} \quad (7),$$

where  $0.03 \text{ d}^{-1}$  and  $4^\circ\text{C}$  are the reference rates [Iversen & Ploug, 2013], and  $T$  is the satellite-derived surface temperature in  $^\circ\text{C}$ . We chose a  $Q_{10}$  of  $3.5$  as this value is based on measurements of carbon-specific respiration on diatom aggregates [Iversen & Ploug, 2013] and is similar to the  $Q_{10}$  measured for growth rates of pelagic bacteria [ $Q_{10} = 3.3$ ; White et al. 1991]. To test the effect of this choice, we also recalculated our analysis using a  $Q_{10}$  of  $2$ . The patterns are the same

for the two  $Q_{10}$  values, but a  $Q_{10}$  of 2 suggests a more extreme over/underestimate of global export. This is because temperature decreases with depth, which likely slows down remineralization rates. A higher  $Q_{10}$  thus induces a more pronounced decrease of the remineralization rate, leading to a stronger non-steady state effect (see Figure 7 and Table 1). To test the sensitivity of  $S$ , we calculated  $S$  at four 'example' sites in both the Atlantic and Pacific (50°N 35°W; 40°S 35°W; 50°N 175°W; and 40°S 175°W) using a range of sinking speeds (10, 40 and 200 m d<sup>-1</sup>) and remineralization rates (0.01, 0.03 and 0.09 d<sup>-1</sup>). As we are interested in illustrating the relative uncertainty in sinking flux due to non-steady state effects,  $S$  was normalized to the daily export flux.

To test the sensitivity of our results to the choice of satellite-derived PP and e-ratio algorithm, we also calculated  $S$  normalised to daily export flux at the four 'example' sites using 12 permutations of commonly used algorithms. We used PP algorithms of Carr [2001], Marra et al. [2003] and the Vertically Generalised Production Model (VGPM [Behrenfeld and Falkowski, 1997]). Export production was estimated using the empirical algorithms by Dunne et al. [2005], Henson et al. [2011] and Laws et al. [2000], in addition to the food web model of Laws et al. [2000]. The procedure followed is the same as in the supplementary information of Henson et al. [2012].

### 2.3 Case study sampling site and particle collection

Samples were collected at the PAP site (49.0°N, 16.5°W) from 8<sup>th</sup> July to 13<sup>th</sup> August 2009 aboard the RRS *Discovery* (D341 Cruise report [Sanders, 2009]). The mixed-layer depth was ~50 m throughout the study period [Giering et al., 2014], and this depth is used as the reference depth for export throughout this study. Sinking flux was measured using free-drifting, neutrally-buoyant sediment traps (PELAGRA [Lampitt et al., 2008]). PELAGRA traps were deployed at depths between 50 and 670 m over a period of 4 weeks (Fig. 1). A total of 16 successful PELAGRA deployments were conducted to acquire four depth profiles of particle flux.

Sample cups for each trap contained filtered seawater (5 ppt excess salinity) and chloroform (at saturating concentrations). Sample cups were closed during the first 18-24 hours of each deployment until traps reached, and stabilised at, the programmed depth. Then, cups opened to collect sinking material for 48-132 hours, before closing again for the ascent and trap recovery. For each trap, the contents of two or three of the four sample cups were screened (350- $\mu$ m mesh) to remove swimmers and split into eight aliquots for analyses of biogenic silica (bSiO<sub>2</sub>), community composition and other constituents of the vertical flux.

For bSiO<sub>2</sub>, aliquots were filtered through 0.8-µm pore polycarbonate membranes (Nuclepore, 25-mm diameter) while at sea and stored at -20 °C. Filters were dried (60°C, 24 hours), then processed and analysed using a method adapted from Ragueneau and Treguer [1994]. Filters were heated with 0.2 M sodium hydroxide for ninety minutes to dissolve biogenic silica, then cooled and neutralised with 0.1 M hydrochloric acid. The resulting solution was separated from remaining particulate material by centrifugation, and dissolved silicate was measured on a SEAL QuAATro autoanalyser. Results were converted to mass of bSiO<sub>2</sub> by assuming a molecular mass of 60. No replicate analyses for bSiO<sub>2</sub> were carried out during this work, but previous replicate measurements by PELAGRA traps [Salter et al., 2007; Martin et al., 2011] suggest a relative standard deviation of 10% between replicate aliquots.

For community composition, aliquots were preserved in 2% borax-buffered formaldehyde and stored in a cool, dark place until further analysis. 1 mL of the preserved flux material was diluted with alkaline de-ionized water (pH 9) to make up 10 mL. The solution was split in half and filtered onto cellulose nitrate filters (0.8-µm nominal pore size, 25-mm diameter), rinsed with alkaline de-ionized water (pH 9), and dried (30°C, 8 hours). Half of each filter was mounted on a glass slide with cover slit using Norland Optical Adhesive (Technoptics) and cured in UV light. Diatoms were counted under plain light at x400 magnifications using a polarising light microscope (Brunel SP200) and identified to species level.

In addition, 30 marine snow aggregates were collected from the base of the mixed layer (~50 m) using the Marine Snow Catcher (MSC) [Riley et al., 2012] and analysed for diatom composition following the method above. The taxonomic composition of diatoms in the mixed layer was determined from 100 mL water samples collected at 5, 25, and 50 m depth using Niskin bottles. Samples were preserved with 2 mL Lugol's iodine solution and stored in a dark, cool place. On shore, cells were allowed to settle for 12 hours, and diatoms were counted and identified using an inverted microscope (Brunel SP951).

#### **2.4 Satellite-derived fluxes at the PAP site**

Fluxes of bSiO<sub>2</sub> at 50 m ('export flux') were estimated from satellite-derived POC flux data. We estimated downward POC flux from the surface ocean using satellite data following Henson et al. [2011]. An empirical relationship between the export ratio (export/primary production) and sea surface temperature (SST) was derived from a database of 306 observations of particle export:  $e\text{-ratio} = 0.23 * \exp(-0.08 * \text{SST})$ . We applied this equation to satellite-derived SST



(MODIS Aqua) at 8-day resolution for a 1-degree region around the PAP site. The e-ratio estimates were then multiplied by satellite-derived primary production (PP) calculated using the algorithm by Carr [2001], again at 8-day temporal resolution, giving an estimate of export in  $\text{mg C m}^{-2} \text{d}^{-1}$ .

POC fluxes were converted to  $\text{bSiO}_2$  fluxes using a  $\text{bSiO}_2$ :POC ratio that varies with the time of year. We extracted this information for the region of the PAP site (47-49 °N, 16-21 °W) from data sets of sediment trap observations during the North Atlantic Bloom experiment [Honjo et al., 1992; Lampitt et al., 2001; Torres-Valdes et al., 2014]. The traps ( $n=26$ ) were deployed between 1000-1200 m depth between the years 1989 – 1990. The data reveal a seasonal pattern in  $\text{bSiO}_2$ :POC (Fig. 2a), which we fitted with a cyclic spline smoother ( $R^2=0.76$ ,  $n=26$ ). The  $\text{bSiO}_2$ :POC ratios at ~1000 m depth are likely to be higher than the corresponding surface ratios due to preferential recycling of POC [Martin et al., 2011]. During our cruise, for example, average  $\text{bSiO}_2$ :POC ratios of slow-sinking particles and total flux at 50 m were, respectively,  $0.065 \pm 0.031 \text{ mg mg}^{-1}$  [Riley et al., 2012] and  $0.021 \pm 0.013 \text{ mg mg}^{-1}$  [this study, using POC measurements from the same trap samples], whereas the deep sediment trap record shows a  $\text{bSiO}_2$ :POC ratio of  $\sim 1.4 \text{ mg mg}^{-1}$  (Fig. 2a). We therefore scaled the time series accordingly ( $\text{bSiO}_2$ :POC at surface = scaling\_factor  $\times$   $\text{bSiO}_2$ :POC at 1000 m) using an average scaling factor ( $\text{bSiO}_2$ :POC<sub>surface</sub>/  $\text{bSiO}_2$ :POC<sub>1000m</sub>) of 0.03 and bracketed this value using lower and higher scaling factors of, respectively, 0.015 and 0.046. We acknowledge that the uncertainties of this ratio are large and that the ratio of  $\text{bSiO}_2$ :POC<sub>surface</sub>/ $\text{bSiO}_2$ :POC<sub>1000m</sub> likely varies throughout the year; however, we consider it sufficient for the purpose of this study, which is to discuss general assumptions rather than to calculate absolute rates.

### **3. Results and Discussion**

#### **3.1 Increase of fluxes with depth**

The results of the two non-steady state simulations (using  $r=0.01 \text{ d}^{-1}$  or  $r=0.20 \text{ d}^{-1}$ ) show that a combination of slow sinking speeds and weak remineralization during a period of declining surface export can lead to an apparent increase of particle flux with depth (Fig. 3a-c). On the other hand, a combination of slow sinking speeds, strong remineralization and declining surface export results in profiles following the expected decrease with depth (Fig. 3a, d-e) because the greater remineralisation rate compensates for the effects of weak remineralisation and decreasing export.

We now present data from the PAP site in the North Atlantic, where we observed an increase of bSiO<sub>2</sub> fluxes with depth. A thorough analysis of the site suggests that this apparent increase could have indeed been caused by non-steady state conditions.

### 3.2 Case study: PAP site

The sediment trap deployments showed that bSiO<sub>2</sub> fluxes increased systematically with depth (Fig 2c). Fluxes ranged between 0.13 – 1.3 mg bSiO<sub>2</sub> m<sup>-2</sup> d<sup>-1</sup> at 50 m and increase to 2.5 – 9.2 mg bSiO<sub>2</sub> m<sup>-2</sup> d<sup>-1</sup> below 500 m. The observed increase of bSiO<sub>2</sub> fluxes with depth is surprising as bSiO<sub>2</sub> is undersaturated in the oceans, with an estimated 55-60% of sinking bSiO<sub>2</sub> dissolving in the upper 100 m on a global scale [Nelson et al. 1995]. The increase with depth has four possible explanations: either (a) there is a significant internal source of bSiO<sub>2</sub>, (b) there is a lateral advection of material, (c) sediment traps near the surface fail to collect the entirety of the flux, or (d) the steady state assumption is incorrect.

bSiO<sub>2</sub> (commonly referred to 'opal') is produced from silicate by specific organisms, including diatoms, radiolarian, silicoflagellates, and siliceous sponges as part of their skeletal structures. Though the production of bSiO<sub>2</sub> is not dependent on light availability, most occurs in the surface ocean with diatoms being the dominant producer [Nelson et al., 1995; Ragueneau et al., 2000]. We therefore assume that the observed increase of bSiO<sub>2</sub> flux at depth is not driven by internal sources of bSiO<sub>2</sub>.

Another potential source of bSiO<sub>2</sub> at depth is laterally advected material. Alonso-González et al. [2009] found that lateral transport of POC from the shelf in the subtropical Northeast Atlantic can reach further than 1000 km off-shore and was up to 3 orders of magnitude larger than vertical fluxes. For the PAP site, however, flow-field analysis (based on satellite-derived near-surface velocities) showed that no water mass reaching the site during our study period passed over the continental shelf during the 3 months preceding the study [Giering et al., 2014].

Even when using Lagrangian sediment traps, such as PELAGRA, source and collection sites can vary. At the study site ALOHA in the subtropical Pacific Ocean, for example, the source of the particles captured in the traps at 500 m depth was likely up to 25 km from the trap location [Siegel et al., 2008]. Although we do not have *in situ* current data, an analysis of modelled surface currents (from ¼°, 5-day NEMO) suggests that the source region of particles captured by the sediment traps was at most 100 km distant. Satellite Chl data from this region and during the deployment period suggest relatively homogeneous surface conditions, with Chl ranging

from 0.45-0.75 mg Chl m<sup>-3</sup> along the tracks. Moreover, during our study the PELAGRAS collected particles from an area covering ~9100 km<sup>2</sup> over a period of 25 days; yet the increase in bSiO<sub>2</sub> flux with depth was apparent across all profiles (Fig. 2c). We therefore conclude that lateral advection of bSiO<sub>2</sub> was likely negligible during our study.

We have also considered the possibility that the collection efficiency of PELAGRA increases with depth [Buesseler et al., 2007]. Low collection efficiency at shallow depths coupled to a slow decline in flux with depth could lead to an apparent increase of flux with depth. During our study cruise, we conducted an extensive comparison of PELAGRA-derived fluxes out of the mixed layer (~50 m) with those obtained from the <sup>234</sup>Th technique, MSC [Riley et al., 2012] and satellites (Fig. 2b). Estimated POC fluxes were 84±2, 99±41, 146±26 and 88±5 mg C m<sup>-2</sup> d<sup>-1</sup> according to PELAGRA, <sup>234</sup>Th, MSC and satellite-derived estimates, respectively. On the basis of this comparison, we do not believe that there was a systematic under-tapping by PELAGRA in the surface.

It follows that a likely explanation for the observed increased bSiO<sub>2</sub> flux at depth is the invalidity of the steady state assumption. We hypothesise that the enhanced bSiO<sub>2</sub> flux at depth is the remnant of an earlier large flux of slowly sinking particles from the surface.

### 3.3 Challenging the steady state assumption

The simulation suggested that three conditions are needed in order to see an increase of flux with depth: (1) slow sinking speeds, (2) a period of declining surface export, and (3) low remineralization rates.

*(1) Slow sinking speeds.* During our cruise, Villa-Alfageme et al. [2014] calculated bulk sinking speeds using <sup>210</sup>Po as a tracer and inverse modelling, and estimated that particles sank, on average, at 88 (±47 S.D.) m d<sup>-1</sup> through the mesopelagic zone. This estimate agrees well with contemporaneous, indirect observations by Riley et al. [2012] via the collection of fast- and slow-sinking particles at 50 m, which suggested that ~63% of the POC flux was sinking at <20 m d<sup>-1</sup> and 37% was sinking at 181 m d<sup>-1</sup> (resulting in a bulk sinking speed of ~73±35 m d<sup>-1</sup>). To support these indirect observations that slowly sinking particles can reach the lower mesopelagic zone, we ideally require a series of tracers that exit the surface ocean successively and thus indicate the age of particles at depth. The diatom population conveniently serves as such a tracer at the PAP site as it changes its taxonomic composition seasonally. Near-surface observations of the community structure are collected throughout the year as part of the

Continuous Plankton Recorder (CPR) survey [www.sahfos.ac.uk]. We compared the taxonomic composition of the CPR samples, determined via microscopy [Richardson et al., 2006], with the taxonomic composition of diatoms in our sediment traps, in the mixed layer and in marine snow aggregates collected at the bottom of the mixed layer.

The change in the species composition with depth in the trap samples was similar to the seasonal change in surface diatom composition. The deepest trap samples analysed for species composition (630 m) contained diatoms of the genus *Bacteriastrium* (cell diameter: 5-56  $\mu\text{m}$  [Tomas 1997]), which had only occurred in the upper ocean  $\sim 3$  months prior to the cruise (Fig. 4). Considering that *Bacteriastrium* spp. must have sunk out of the mixed layer  $\sim 40$ -70 days prior to the deployment of the deepest trap, this suggests that *Bacteriastrium* spp. sank at a speed of 8-15  $\text{m d}^{-1}$ . The diatom *Cylindrotheca closterium* (cell length: 30-400  $\mu\text{m}$ , cell width: 2-8  $\mu\text{m}$  [Tomas 1997]), on the other hand, only occurred in surface waters in July, yet were present in all traps during our cruise, including the deepest trap which had been deployed in mid-July (Fig. 4). To reach this depth within 14 days, *Cylindrotheca* spp. would have to sink at  $\sim 40$   $\text{m d}^{-1}$ . It follows that the effective sinking speed of these diatoms was on the order of 8-40  $\text{m d}^{-1}$ , which agrees with the observation that a significant fraction of particle flux is slow-sinking [Alonso-Gonzalez et al., 2010; Riley et al., 2012; Villa-Alfagame et al., 2014; Giering et al., 2016].

(2) *A period of declining surface export.* Satellite-derived export estimates show that export flux was indeed declining in the weeks before our study. Throughout 2009, estimated export flux of POC and bSiO<sub>2</sub> at the PAP site ranged over one order of magnitude from  $\sim 20$   $\text{mg POC m}^{-2} \text{d}^{-1}$  and  $\sim 1$   $\text{mg bSiO}_2 \text{m}^{-2} \text{d}^{-1}$  in winter to 140  $\text{mg POC m}^{-2} \text{d}^{-1}$  and  $7 \pm 3.5$   $\text{mg bSiO}_2 \text{m}^{-2} \text{d}^{-1}$  at the peak of the bloom (Fig. 2b). For the two weeks immediately prior to the first sediment trap deployments (1<sup>st</sup> Jul – 15<sup>th</sup> Jul) export was fairly constant, whereas during the preceding month (1<sup>st</sup> Jun – 1<sup>st</sup> Jul) export had declined rapidly (Fig. 2b).

(3) *Low remineralization rates.* Unfortunately, we have no direct measurements for the remineralization/dissolution rate of bSiO<sub>2</sub> at our study site. Yet, we can estimate these rates indirectly. Rickert et al. [2002] measured dissolution kinetics of bSiO<sub>2</sub> collected from the surface, sediment traps and sediments in different oceanic regimes, including material collected by sediment traps at 1000 m and 2500 m depth from the Norwegian Basin (70°N, 4°E). The data suggested that dissolution rates of bSiO<sub>2</sub> vary between the different regions, with natural samples dissolving fastest in the Norwegian Basin. Assuming that the bSiO<sub>2</sub> dissolution kinetics in the Norwegian Basin are similar to those at our study site, potential dissolution rates can be

estimated using the reported equations for temperature-corrected solubility and dissolution rate coefficient ( $k$ ) [Rickert et al., 2002]. Temperatures at the PAP site were  $\sim 15.5^{\circ}\text{C}$  at the surface and dropped to as low as  $10^{\circ}\text{C}$  at 600 m depth, with an average temperature of  $11.6^{\circ}\text{C}$  between 50 – 600 m. Solubility of  $\text{bSiO}_2$  at the PAP site was thus likely  $\sim 1260 \mu\text{M}$  [Equation 11 by Rickert et al., 2002] and  $k$  was likely between  $9\text{-}15 \mu\text{mol g}^{-1} \text{h}^{-1}$  [Equation 10 by Rickert et al., 2002]. At silicate concentrations of  $\sim 4 \mu\text{M}$  during our study [Stinchcombe, pers. comm.], the daily dissolution rate of  $\text{bSiO}_2$  [Equation 8 by Rickert et al., 2002] was thus probably in the order of  $0.01 - 0.02 \text{ d}^{-1}$ .

It appears that all of the three conditions that can lead to an apparent increase of flux with depth – slow sinking speeds, a period of declining surface export, and low remineralization rates – were present during our study. This strongly supports our hypothesis that parts of the deep flux are the remnants of export events that occurred several weeks previously.

Taking this a step further, we combine the derived rates of surface  $\text{bSiO}_2$  export (Fig. 2b), sinking speed ( $10 \text{ m d}^{-1}$ ) and a remineralization rate of  $0.01 \text{ d}^{-1}$  to produce depth profiles of particle flux as described for the non-steady state model (Equation 2). The resulting flux profile matches the observed  $\text{bSiO}_2$  flux reasonably well (Fig. 5), supporting the hypothesis that the flux observed at depth could have been the remnants of the spring bloom. This calculation is, however, based on estimates with relatively large uncertainties and thus provides an illustrative example only.

We conclude that steady state conditions did not occur at the PAP site during our study, likely leading to the observed increase of  $\text{bSiO}_2$  flux with depth. Non-steady state conditions may lead to misinterpretation of particle attenuation rates both at PAP and at other sites where input fluxes from the upper ocean vary on the order of weeks, such as in polar and temperate regions with a strong seasonal cycle in export [Henson et al., 2015].

### **3.3 Potential error when estimating net carbon supply to the mesopelagic zone**

Increases in  $\text{bSiO}_2$  flux, and also iron or manganese fluxes, with depth have been observed previously in higher latitude regions with pronounced seasonal blooms, such as at K2 ( $47^{\circ}\text{N}$ ,  $160^{\circ}\text{E}$ ) during the VERTIGO project [Lamborg et al., 2008]. These increases were attributed to lateral advection of suspended particulate matter. However, we here suggest that the increased flux with depth at these sites could have been, as was the case here, remnant material of the bloom, which had occurred 50-60 days before sampling took place [Buesseler et al., 2008]. We

believe that our observation of non-steady state conditions is an example of a more widespread phenomenon. This has significant implications for the interpretation of POC flux profiles and attempts to balance the net supply and demands of organic matter to the dark ocean.

We derived an equation that allows estimation of the magnitude of 'misestimated' organic matter supply when assuming steady state (Equation 6). It is important to note that this 'misestimate' can be either more or less organic matter supply than inferred when assuming steady state. For example, underestimates of net carbon supply increase when remineralization rates decrease (Table 1), sinking speeds slow down (Table 1.1), and/or the decrease in export flux accelerates (Table 1.2). Our equation can be easily applied to any situation where the time history of export flux prior to sampling is known.

For the PAP site in August 2009, the loss of POC was likely underestimated by 5 – 20 mg C m<sup>-2</sup> d<sup>-1</sup> across the mesopelagic zone (mixed layer depth to 1000 m; Table 1). This is on the order of 5 – 22% of the net organic carbon supply at this site when assuming steady state [92 mg C m<sup>-2</sup> d<sup>-1</sup>; Giering et al., 2014]. This analysis demonstrates quantitatively that ignoring temporal variability in flux could lead to some of the uncertainties surrounding imbalanced carbon budgets [Burd et al., 2010].

We next illustrate on a global scale the potential range of error in estimated net POC supply to the upper mesopelagic zone (100 – 500 m depth) when assuming steady state (Fig. 6). For this illustration, we assumed average sinking speeds of 40 m d<sup>-1</sup> and a temperature-corrected remineralization rate (see Section 2.2), though variations in these parameters will affect the predictions in a similar fashion as explored in Table 1 (Fig. 7). We further used several combinations of algorithms for PP and e-ratio to see whether our estimates for *S* are robust. The median deviation away from the combination of algorithms used in the main text [Carr, 2001; Henson et al., 2003] is ~5% and, although some combinations of algorithms show larger deviations, all show the same temporal evolution (Fig. 8).

Largest errors in estimates of net POC supply to the mesopelagic zone are predicted for the high latitudes (50-65°N and S) in both the Atlantic and Pacific (Fig. 6). In our scenario (sinking speed of 40 m d<sup>-1</sup>; remineralization rate dependent on temperature), daily net POC supply is likely overestimated by up to 20 mg C m<sup>-2</sup> d<sup>-1</sup> before the bloom, whereas it is likely underestimated by the same amount after the bloom (Fig. 6a,c). As expected, estimates from regions with near

constant export rates throughout the year (30°N to 30°S) are little affected when assuming steady state ( $<5 \text{ mg C m}^{-2} \text{ d}^{-1}$ ; Fig. 6a, c).

A similar global pattern emerges when normalizing  $S$  to daily export (Fig. 6b,d). Errors in the tropics (30°N to 30°S) are low ( $<10\%$ ), and field campaigns in these regions (e.g. at ALOHA or BATS) are likely unaffected when assuming steady state. Between 30-65°N and S, two phases are prevalent: spring/summer before the peak of the bloom when net POC supply could be overestimated by  $\leq 25\%$ , and autumn after the peak of the bloom when net POC supply could be underestimated by  $\leq 25\%$ . Field campaigns in these regions (e.g. at PAP or K2) are thus likely to sample during non-steady state conditions and are more susceptible to misinterpretation of particle flux profiles. For high-latitude regions ( $>65^\circ\text{N}$  and S) the estimates of POC loss in the model are poorly constrained. Nevertheless, potential errors are likely comparable to those of the temperate regions (30-65°) – if not higher – owing to the strong seasonality in these regions. Field campaigns investigating particle flux in high-latitude regions will have to pay careful attention when assuming steady state.

#### **4. Conclusion**

Sediment traps have been used extensively for determining particle flux and remineralization rates. The latter is calculated by comparing fluxes collected at different depths, which implies that the source (both composition and quantity) of the material caught in the traps is identical for these traps. However, our data suggests that this is likely not the case in temperate regions, where seasonal cycles cause large changes in export rates and composition throughout the year. As a large fraction of the sinking material may be sinking at slow rates ( $<40 \text{ m d}^{-1}$ ), these ecosystems do not operate at steady state, and the resulting interpretation of the flux profile is misleading. Net POC supply to the twilight zone may be wrong by as much as 25% when assuming steady state. Nevertheless, many flux campaigns take place in the post bloom period assuming that the system is at steady state. We here provide a simple equation to estimate the potential error when calculating carbon supply to the dark ocean during non-steady state conditions. Future studies and methods (e.g. radio-tracer methods using  $^{234}\text{Th}$  and  $^{210}\text{Po}$ ) should consider this when interpreting flux profiles.

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satellite products used in the paper can be obtained by emailing [s.henson@noc.ac.uk](mailto:s.henson@noc.ac.uk). We thank four anonymous reviewers for their constructive criticism, and the Natural Environmental Research Council (NERC) for support through National Capability funding.

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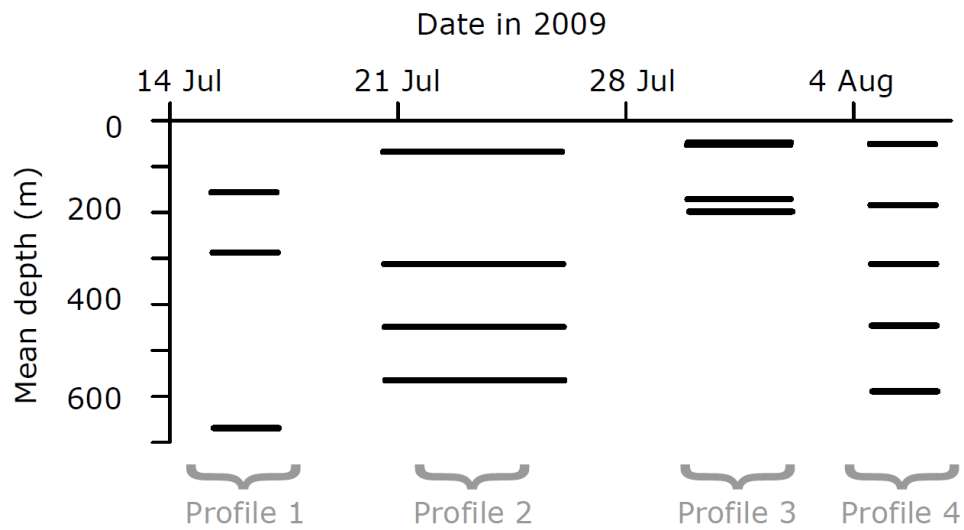
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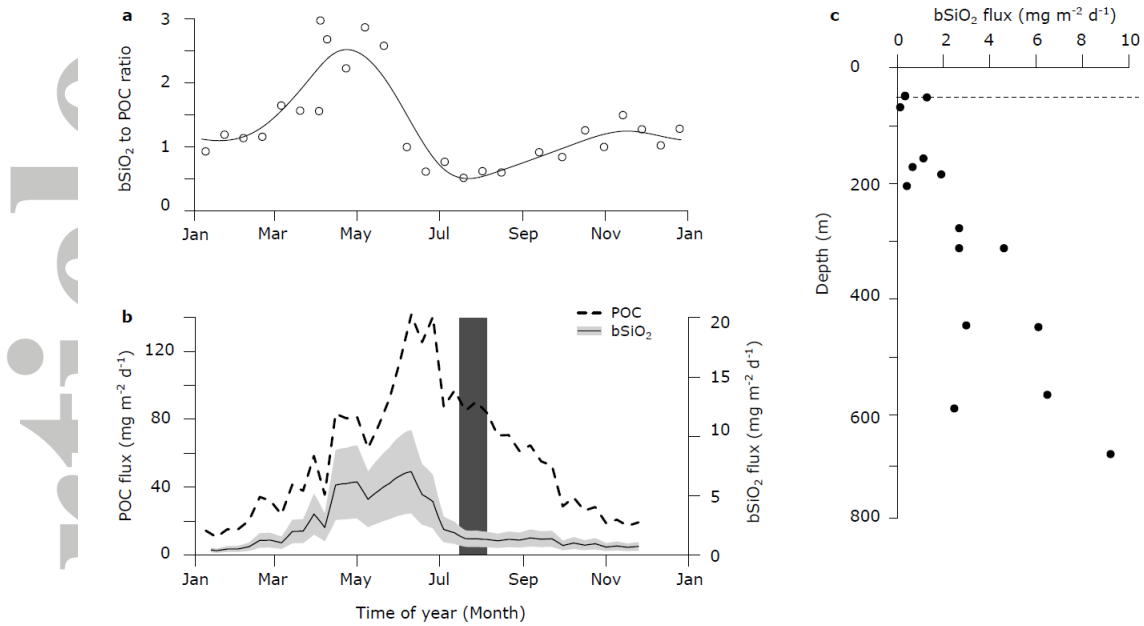
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## Figures

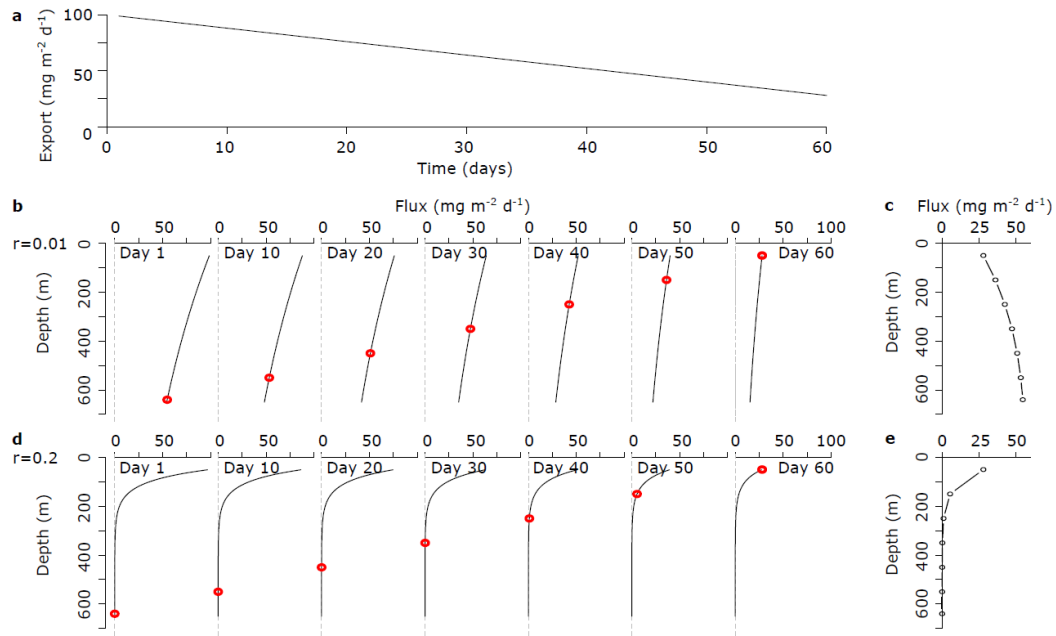


**Figure 1.** Details for the four PELAGRA profiles. Deployment period and mean depth for each PELAGRA deployment is indicated by a horizontal line, which marks the opening and closing of the sample cups.

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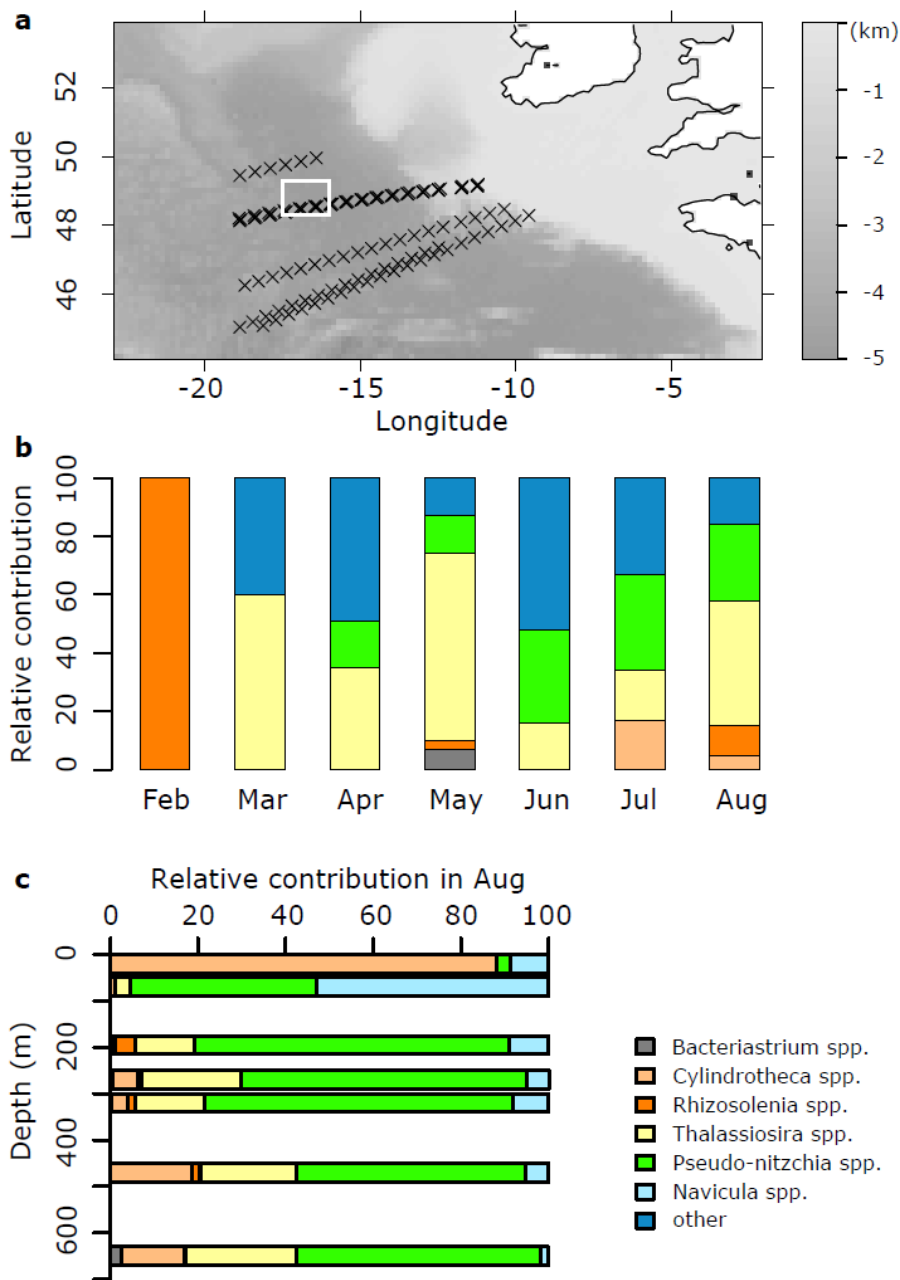


**Figure 2.** (a) Ratio of bSiO<sub>2</sub> to POC at the PAP site. Flux was collected using sediment traps at 1000 – 1200 m depth during 1989 – 1990 at the PAP site [Honjo et al., 1992, Lampitt et al., 2001; Torres-Valdes et al., 2014] and bSiO<sub>2</sub>:POC ratio calculated (mg SiO<sub>2</sub> mg<sup>-1</sup> POC). Line shows fit with cyclic spline smoother ( $p < 0.001$ ,  $R^2 = 0.76$ ,  $n = 26$ ). (b) Satellite time series of particulate organic carbon (POC) export (dashed line) and bSiO<sub>2</sub> export (solid line). Shaded area (light grey) shows uncertainty of bSiO<sub>2</sub> export. The time of the cruise is highlighted in dark grey. (c) Depth profile of bSiO<sub>2</sub> fluxes as measured using PELAGRA during our study. Dashed line: Average mixed layer depth during the study.

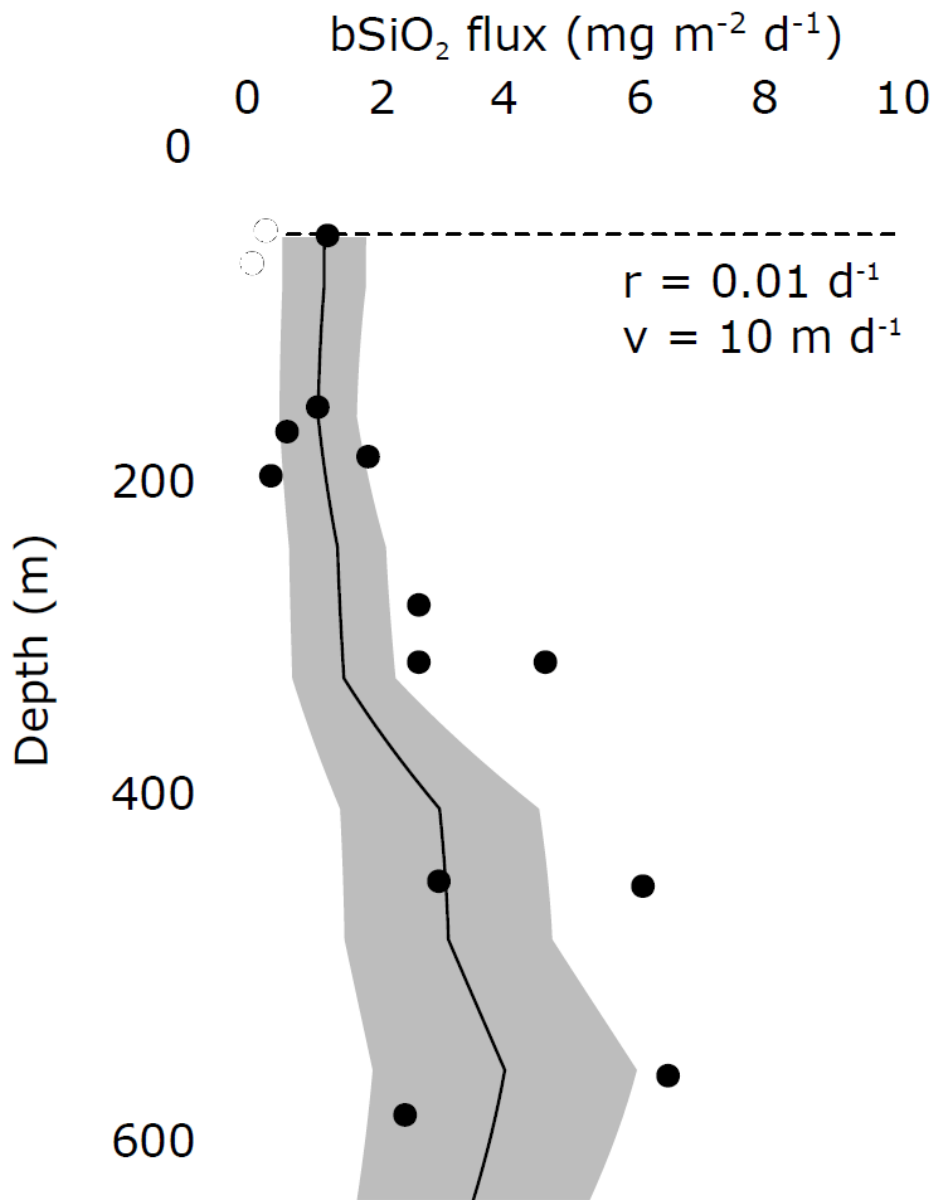


**Figure 3.** Simulation of how the presence of remnant particles at depth affect particle flux profiles. **(a)** Export flux at 50 m for the simulation period. **(b,d)** Depth profiles of particle flux between days 1 to 60 when assuming remineralization rate  $r$  of 0.01 and 0.20 (in panel b and d, respectively). Red circles show the flux sampled by a hypothetical sediment trap deployed at day 60 if particles sank at 10 m d<sup>-1</sup>. **(c,e)** Depth profiles of particle flux as seen by a hypothetical sediment trap deployed at day 60 if particles sank at 10 m d<sup>-1</sup> and  $r=0.01$  d<sup>-1</sup> or  $r=0.20$  d<sup>-1</sup> (in panel c and e, respectively).

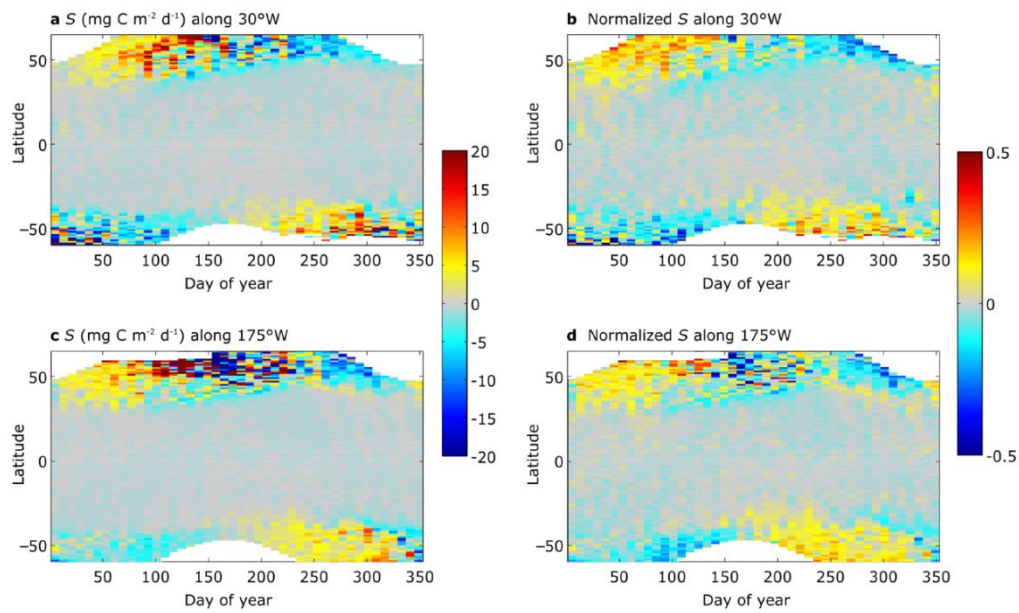




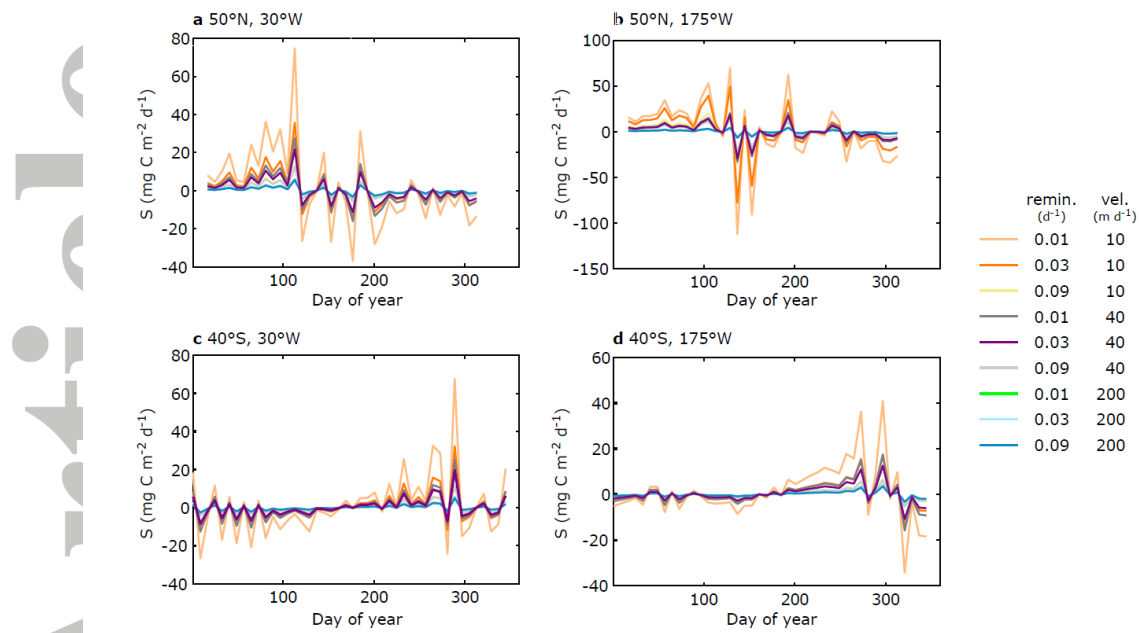
**Figure 4.** Changes in diatom population at the PAP site. **(a)** Sampling locations of the Continuous Plankton Recorder (CPR) survey in Feb-Aug 2009. White box shows PAP site area. **(b)** Relative contribution of different diatom species to total diatom abundance (based on cell counts in the CPR samples) in the upper ocean during Feb-Aug 2009. Colours represent species as in the key. **(c)** Changes of relative contribution with depth within the PAP site area in Aug 2009. Samples were collected using CPR (surface), Marine Snow Catcher (50 m) and PELAGRA sediment traps (175-630 m). Colours represent species as in the key.



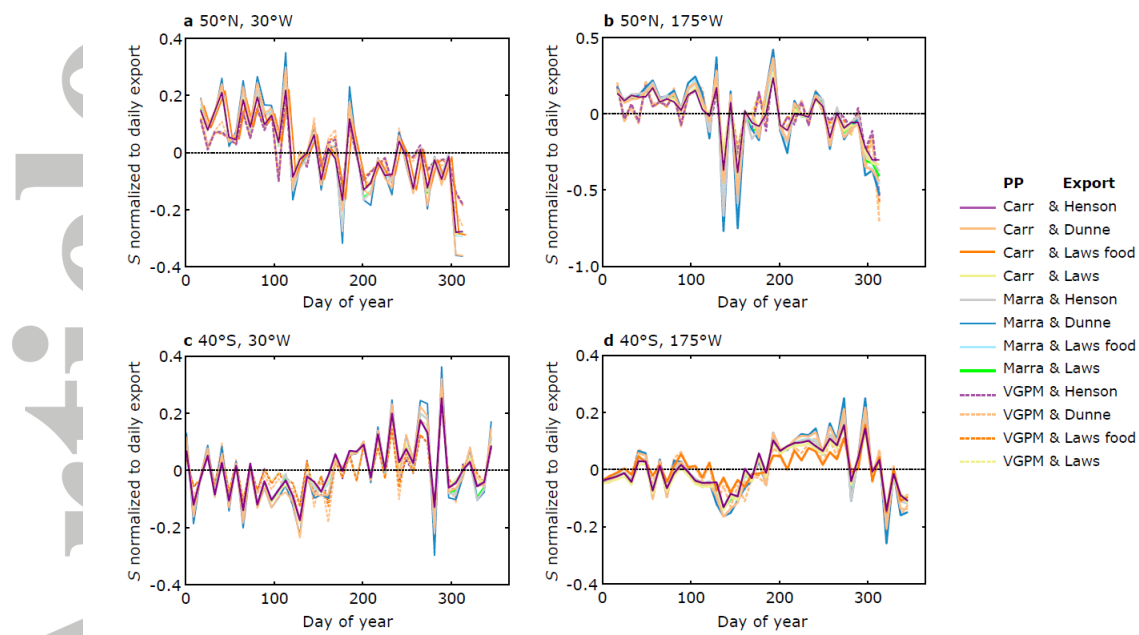
**Figure 5.** Comparison of observed bSiO<sub>2</sub> flux (mg bSiO<sub>2</sub> m<sup>-2</sup> d<sup>-1</sup>) at the PAP site during 6<sup>th</sup> Jun – 5<sup>th</sup> Aug 2009 (circles) and bSiO<sub>2</sub> flux profile calculated using satellite-derived estimates of bSiO<sub>2</sub> export and assuming sinking speeds of 10 m d<sup>-1</sup> and remineralization rates of 0.01 d<sup>-1</sup> (solid line). Grey area: Uncertainty envelope expressing uncertainties in the scaling factor used for calculating bSiO<sub>2</sub>:POC ratios. Dashed line: Average mixed layer depth during the study.



**Figure 6.** Seasonal variation in the potential error of estimated net POC supply to the upper mesopelagic zone (100 – 500 m depth) when assuming steady state ( $S$ ). **(a,c)**  $S$  ( $\text{mg C m}^{-2} \text{d}^{-1}$ ) was calculated for two transects along the Atlantic ( $30^\circ\text{W}$ ; panel a) and the Pacific ( $175^\circ\text{W}$ ; panel c) using a satellite-derived climatology of export estimates, average sinking speeds of  $40 \text{ m d}^{-1}$ , and a temperature-corrected remineralization rate. **(b,d)**  $S$  was then normalized to daily export (panel b and d for the Atlantic and Pacific transect, respectively).



**Figure 7.** Sensitivity analysis of  $S$  (error in estimated net POC supply when assuming steady state) at four sites in the **(a)** North Atlantic (50°N, 30°W), **(b)** North Pacific (50°N, 175°W), **(c)** South Atlantic (40°S, 30°W), and **(d)** South Pacific (40°S, 175°W).  $S$  increases with decreasing remineralization rate ( $r$  in d<sup>-1</sup>) and/or decreasing sinking speed ( $v$  in m d<sup>-1</sup>). Combinations of  $v$  and  $r$  are presented as shown in legend.



**Figure 8.** Time series of  $S$  normalised to daily export calculated using combinations of different algorithms for PP [Carr, 2001; Marra et al., 2003; Behrenfeld and Falkowski, 1997 ('VGPM')] and e-ratio [Dunne et al., 2005; Henson et al., 2011; Laws et al., 2000] for four sites in the **(a)** North Atlantic ( $50^{\circ}\text{N}$ ,  $30^{\circ}\text{W}$ ), **(b)** North Pacific ( $50^{\circ}\text{N}$ ,  $175^{\circ}\text{W}$ ), **(c)** South Atlantic ( $40^{\circ}\text{S}$ ,  $30^{\circ}\text{W}$ ), and **(d)** South Pacific ( $40^{\circ}\text{S}$ ,  $175^{\circ}\text{W}$ ). Lines show permutations of the algorithms as shown in legend.

**Table 1.1:**  $\alpha = 1.1 \text{ mg C m}^{-2} \text{ d}^{-1}$ 

$v$	$\Delta t$	$r$					
		-0.01	-0.02	-0.05	-0.1	-0.2	-0.5
10	95	68	47	22	12	6	3
20	48	42	34	20	11	6	3
50	19	19	18	14	10	6	3
100	10	10	10	9	7	5	3
200	5	5	5	5	4	4	3

**Table 1.2:**  $v = 50 \text{ m d}^{-1}$ .

$\alpha$	$r$					
	-0.01	-0.02	-0.05	-0.1	-0.2	-0.5
0.5	9	8	6	4	3	1
0.7	12	11	9	6	4	2
0.9	16	14	11	8	5	2
1.1	19	18	14	10	6	3
1.3	23	21	16	12	7	3

**Table 1.** Potential underestimation of organic matter supply to the dark ocean, integrated between 0-950 m below the mixed layer (50 m deep), as a function of remineralization rate ( $r$ ;  $\text{d}^{-1}$ ), sinking speed ( $v$  in  $\text{m d}^{-1}$ ; and associated  $\Delta t$  in days), and decrease in export flux over time ( $\alpha$ ; here in  $\text{mg C m}^{-2} \text{ d}^{-1}$ ). Light and dark grey shaded areas show, respectively, likely and most likely combinations at the PAP site during August 2009.