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# Comment on “Predicting Narrow States in the Spectrum of a Nucleus beyond the Proton Drip Line”

## **Disciplines**

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## **Comments**

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**Comment on “Predicting Narrow States in the Spectrum of a Nucleus beyond the Proton Drip Line”**

A recent Letter [1] presented calculations of several resonances in  $^{14}\text{C} + n$  and  $^{14}\text{O} + p$ , including three negative-parity states, for which they used the structure of a  $0p_{1/2}$  nucleon coupled to  $0_2^+$  and  $2_2^+$  core states in  $^{14}\text{C}$  and  $^{14}\text{O}$ . These negative-parity states in  $^{15}\text{C}$  are nearly pure  $(sd)^2$  neutrons coupled to the ground state (g.s.) of  $^{13}\text{C}$ , with the pair of neutrons having  $J = 0$  or  $2$ , as has been suggested [2] by the reaction  $^{13}\text{C}(t, p)$ . This configuration has a large overlap with that of Ref. [1], but there appears to be a problem with the widths in both  $^{15}\text{C}$  and  $^{15}\text{F}$  and with the energy shifts between  $^{15}\text{C}$  and  $^{15}\text{F}$ .

In Ref. [1], the  $1/2^-$  state at 3.10 MeV in  $^{15}\text{C}$  has a calculated width of 2 keV, but its experimental width is 42 keV [3]. This width can come only from neutron decay to  $^{14}\text{C}(\text{g.s.})$ . A width of 42 keV, combined with an  $\ell = 1$  single-particle (SP) width of 1.3 MeV results in a spectroscopic factor  $S$  of 0.033, where we have used the relationship  $\Gamma_{\text{expt}} = C^2 S \Gamma_{\text{SP}}$ , with  $C^2 = 1$  here. The SP width, and hence  $S$ , can be sensitive to the details of the SP calculation. However, here we are primarily interested in the ratio  $\Gamma(^{15}\text{F})/\Gamma(^{15}\text{C})$  for mirror states, and that ratio is very insensitive to those details.

For the mirror state in  $^{15}\text{F}$ , we can compute its energy using the configuration  $^{13}\text{N}(\text{g.s.}) \times (sd)_{0+}^2$ , with the mixture of  $s^2$  and  $d^2$  from [4]. The result is  $E_p = 4.63$  MeV, not very close to 5.49 MeV in Ref. [1]. With good isospin, the spectroscopic factor in  $^{15}\text{F}$  is the same as in  $^{15}\text{C}$ , so we can compute the expected width of this  $1/2^-$  state in  $^{15}\text{F}$  from the expression  $\Gamma = S \Gamma_{\text{SP}}$ , where now  $\Gamma_{\text{SP}}$  is the  $\ell = 1$  SP width for proton decay. Our calculated value for this SP width for a state at our calculated energy is about 1.6 MeV, so that we expect  $\Gamma(^{15}\text{F}, 1/2^-) \approx 55$  keV, significantly larger than the width of 5 keV in Ref. [1]. If the state is at the energy computed in Ref. [1], its width would be  $\geq 65$  keV. These values are summarized in Table I.

The  $5/2^-$  and  $3/2^-$  energies in  $^{15}\text{C}$  are 4.22 and 4.66 MeV, respectively. The  $3/2^-$  state has considerable width—perhaps (by inspection of the spectrum in [2]) as much as 100–150 keV, similar to the calculated width of 90 keV in Ref. [1]. With the configuration of  $(sd)_{2+}^2$  coupled to the  $^{13}\text{C}$  (or  $^{13}\text{N}$ ) g.s., we get energies in  $^{15}\text{F}$  of  $E_p = 5.92$  and 6.30 MeV, for  $5/2^-$  and  $3/2^-$ , respectively. Reference [1] has these two states at 6.88 and 7.25 MeV. Their width for the  $3/2^-$  state in  $^{15}\text{F}$  is 40 keV. It is very difficult to understand how the width in  $^{15}\text{F}$  could be less than in  $^{15}\text{C}$ . From their  $n$  width in  $^{15}\text{C}$ , we estimate a  $3/2^-$  width in  $^{15}\text{F}$  of about 180 keV for a state at our energy and about 200 keV if at their energy. If the width in  $^{15}\text{C}$  is 150 keV, these change to 300 and 325 keV. These are also listed in Table I.

We have not found an estimate of the experimental neutron width of the  $5/2^-$  state in  $^{15}\text{C}$ , for which the

TABLE I. Energies (MeV) and widths (keV) of three lowest negative-parity states in  $^{15}\text{C}$  and  $^{15}\text{F}$ .

$J^\pi$	$E_x$	$^{15}\text{C}$		$^{15}\text{F}$			
		Ref. [1]	Other	Ref. [1] $E_p$	$\Gamma$	Present $E_p$	$\Gamma$
$1/2^-$	3.10 <sup>a</sup>	2	42 <sup>b</sup>	5.49	5	4.63	55 <sup>c</sup>
$5/2^-$	4.22 <sup>d</sup>	2	$\leq 14^a$	6.88	10	5.92	6 <sup>e</sup>
$3/2^-$	4.66 <sup>d</sup>	90	100–150 <sup>f</sup>	7.25	40	6.30	180 <sup>g</sup>

<sup>a</sup>Ref. [5].<sup>b</sup>Ref. [3].<sup>c</sup>If  $E_p$  is 5.49 MeV,  $\Gamma$  is  $\geq 65$  keV.<sup>d</sup>Ref. [2].<sup>e</sup>This value is for  $\Gamma(^{15}\text{C}) = 2$  keV and  $E_p = 5.92$  MeV. If  $E_p$  is 6.88 MeV, we get  $\Gamma = 10$  keV.<sup>f</sup>By inspection of the spectrum of Ref. [2].<sup>g</sup>Using  $\Gamma(^{15}\text{C}) = 90$  keV. If  $E_p$  is 7.25 MeV,  $\Gamma$  is  $\geq 200$  keV. If  $\Gamma(^{15}\text{C})$  is 150 keV,  $\Gamma$  is 300–325 keV.

compilation [5] gives  $\leq 14$  keV. Reference [1] lists 2 keV for the calculated value of this quantity. If this value is correct, the width of the mirror state in  $^{15}\text{F}$  would be 6 keV if it is at our energy, 10 keV if at the energy of Ref. [1].

In addition, a second  $1/2^-$  state in  $^{15}\text{C}$  at 5.87 MeV, with a width of about 100 keV, is within the range of energies considered by Ref. [1].

We agree with Ref. [1] that narrow resonances are to be expected in  $^{14}\text{C} + n$  and  $^{14}\text{O} + p$  in the energy range discussed, but it would appear that the energies and widths of the negative-parity resonances will be considerably different from the ones calculated in Ref. [1].

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