## Original citation:

Korepanov, Alexey, Kosloff, Zemer and Melbourne, Ian. (2017) Averaging and rates of averaging for uniform families of deterministic fast-slow skew product systems. Studia Mathematica, 238. pp. 59-89.

## Permanent WRAP URL:

http://wrap.warwick.ac.uk/85562

## Copyright and reuse:

The Warwick Research Archive Portal (WRAP) makes this work by researchers of the University of Warwick available open access under the following conditions. Copyright © and all moral rights to the version of the paper presented here belong to the individual author(s) and/or other copyright owners. To the extent reasonable and practicable the material made available in WRAP has been checked for eligibility before being made available.

Copies of full items can be used for personal research or study, educational, or not-for-profit purposes without prior permission or charge. Provided that the authors, title and full bibliographic details are credited, a hyperlink and/or URL is given for the original metadata page and the content is not changed in any way.

## A note on versions:

The version presented here may differ from the published version or, version of record, if you wish to cite this item you are advised to consult the publisher's version. Please see the 'permanent WRAP URL' above for details on accessing the published version and note that access may require a subscription.

For more information, please contact the WRAP Team at: wrap@warwick.ac.uk

# Averaging and Rates of Averaging for Uniform Families of Deterministic Fast-Slow Skew Product Systems 

A. Korepanov<br>Z. Kosloff<br>I. Melbourne<br>Mathematics Institute, University of Warwick, Coventry, CV4 7AL, UK

13 January 2017


#### Abstract

We consider families of fast-slow skew product maps of the form $$
x_{n+1}=x_{n}+\epsilon a\left(x_{n}, y_{n}, \epsilon\right), \quad y_{n+1}=T_{\epsilon} y_{n},
$$ where $T_{\epsilon}$ is a family of nonuniformly expanding maps, and prove averaging and rates of averaging for the slow variables $x$ as $\epsilon \rightarrow 0$. Similar results are obtained also for continuous time systems $$
\dot{x}=\epsilon a(x, y, \epsilon), \quad \dot{y}=g_{\epsilon}(y) .
$$

Our results include cases where the family of fast dynamical systems consists of intermittent maps, unimodal maps (along the Collet-Eckmann parameters) and Viana maps.


## 1 Introduction

The classical Krylov-Bogolyubov averaging method [32] deals with skew product flows of the form

$$
\dot{x}=\epsilon a(x, y, \epsilon), \quad \dot{y}=g(y) .
$$

Let $\nu$ be an ergodic invariant probability measure for the fast flow generated by $g$. Under a uniform Lipschitz condition on $a$, it can be shown that solutions to the slow $x$ dynamics, suitably rescaled, converge almost surely to solutions of an averaged ODE $\dot{X}=\bar{a}(X)$ where $\bar{a}(x)=\int a(x, y, 0) d \nu(y)$.

A considerably harder problem is to handle the fully-coupled situation

$$
\dot{x}=\epsilon a(x, y, \epsilon), \quad \dot{y}=g(x, y, \epsilon)
$$

Here it is supposed that there is a distinguished family of ergodic invariant probability measures $\nu_{x, \epsilon}$ for the fast vector fields $g(x, \cdot, \epsilon)$ and the averaged vector field is given by $\bar{a}(x)=\int a(x, y, 0) d \nu_{x, 0}(y)$. The first results on averaging for fully-coupled systems were due to Anosov [6] who considered the case where the fast vector fields are Anosov with $\nu_{x, \epsilon}$ absolutely continuous. Convergence here is in the sense of convergence in probability with respect to Lebesgue measure.

Kifer [23, 24] extended the results of [6] to the case where the fast vector fields are Axiom A (uniformly hyperbolic) with SRB measures $\nu_{\epsilon}$. More generally, Kifer considers the case where $x \mapsto \nu_{x, 0}$ is sufficiently regular so that $\bar{a}$ is Lipschitz, and gives necessary and sufficient conditions for averaging to hold. However, the only situations where the conditions in [23, 24] are verified are in the Axiom A case, even though it is hoped [24] that the conditions are verifiable for nonuniformly hyperbolic examples. Analogous results for the discrete time case are obtained in [22]. See also [14, Theorem 5] for certain partially hyperbolic fast vector fields.

Here, we consider an intermediate class of examples that lies between the classical uncoupled situation and the fully coupled systems of [6, 23], namely families of skew products of the form

$$
\begin{equation*}
\dot{x}=\epsilon a(x, y, \epsilon), \quad \dot{y}=g(y, \epsilon), \tag{1.1}
\end{equation*}
$$

with distinguished family of ergodic invariant measures $\nu_{\epsilon}$ and averaged vector field $\bar{a}(x)=\int a(x, y, 0) d \nu_{0}(y)$. Notice that in this way we avoid issues concerned with the regularity of the averaged vector field $\bar{a}$, but we still have to deal with the $\epsilon$-dependence of the measures $\nu_{\epsilon}$ as well as the fast vector fields. In other words, linear response (differentiability) of the invariant measures is replaced by statistical stability (weak convergence) which is more tractable. Indeed one aspect of the general framework in this paper is that our averaging theorems hold in a similar generality to the methods of Alves \& Viana [2, 5] for proving statistical stability.

Hence, we obtain results on averaging and rates of averaging for a large class of families of skew products (1.1), going far beyond the uniformly hyperbolic setting, both in discrete and continuous time. Our examples include situations where the fast dynamics is given by intermittent maps with arbitrarily poor mixing properties, unimodal maps where linear response fails, and flows built as suspensions over such maps.

We obtain results also on rates of averaging. In the very simple situation $\dot{x}=\epsilon a(y)$, $\dot{y}=g(y)$, where $g$ is a uniformly expanding semiflow or uniformly hyperbolic flow, it is easily seen that the optimal rate of averaging in $L^{1}$ is $O\left(\epsilon^{1 / 2}\right)$. For systems of the form (1.1), we often obtain the essentially optimal rate $\left.O\left(\epsilon^{\frac{1}{2}-}\right)\right|_{-} ^{1}$

We have chosen to focus in this paper on the case of noninvertible dynamical systems. In this situation, the measures of interest are absolutely continuous and we are able to present the main ideas without going into the technical issues presented by

[^0]dealing with nonabsolutely continuous measures as required in the invertible setting. The invertible case will be covered in a separate paper.

Even in the noninvertible setting, our results depend strongly on extensions and clarifications of the classical second order averaging theorem. These prerequisites are presented in Appendix A and may be of independent interest.

The remainder of the paper is organised as followed. In Section 2, we set up the averaging problem for families of fast-slow skew product systems in the discrete time case, leading to a general result Theorem 2.2 for such systems. In Section 3, we show that Theorem 2.2 leads easily to averaging when the fast dynamics is a family of uniformly expanding maps. Section 4 is the heart of the paper and deals with the case when the fast dynamics is a family of nonuniformly expanding maps. Our main examples are presented in Section 5. In Section 6, we show how the continuous time case reduces to the discrete time case. In Section 7, we present a simple example to show that almost sure convergence fails for families of skew products.

## 2 General averaging theorem for families of skew products

Let $T_{\epsilon}: M \rightarrow M, 0 \leq \epsilon<\epsilon_{0}$, be a family of transformations defined on a measurable space $M$. For each $\epsilon \in\left[0, \epsilon_{0}\right)$, let $\nu_{\epsilon}$ denote a $T_{\epsilon}$-invariant ergodic probability measure on $M$.

We consider the family of fast-slow systems

$$
\begin{align*}
x_{n+1}^{(\epsilon)} & =x_{n}^{(\epsilon)}+\epsilon a\left(x_{n}^{(\epsilon)}, y_{n}^{(\epsilon)}, \epsilon\right), \quad x_{0}^{(\epsilon)}=x_{0} \\
y_{n+1}^{(\epsilon)} & =T_{\epsilon} y_{n}^{(\epsilon)}, \quad y_{0}^{(\epsilon)}=y_{0} \tag{2.1}
\end{align*}
$$

where the initial condition $x_{0}^{(\epsilon)}=x_{0}$ is fixed throughout. The initial condition $y_{0} \in$ $M$ is chosen randomly with respect to various measures that are specified in the statements of the results. Here $a: \mathbb{R}^{d} \times M \times\left[0, \epsilon_{0}\right) \rightarrow \mathbb{R}^{d}$ is a family of functions satisfying certain regularity hypotheses.

Define $\bar{a}(x)=\int_{M} a(x, y, 0) d \nu_{0}(y)$ and consider the ODE

$$
\begin{equation*}
\dot{X}=\bar{a}(X), \quad X(0)=x_{0} . \tag{2.2}
\end{equation*}
$$

We are interested in the convergence, and rate of convergence, of the slow variables $x_{n}^{(\epsilon)}$, suitably rescaled, to solutions $X(t)$ of this ODE. More precisely, define $\hat{x}^{(\epsilon)}$ : $[0,1] \rightarrow \mathbb{R}^{d}$ by setting $\hat{x}^{(\epsilon)}(t)=x_{[t / \epsilon]}^{(\epsilon)}$. We study convergence of the difference

$$
z_{\epsilon}=\sup _{t \in[0,1]}\left|\hat{x}^{(\epsilon)}(t)-X(t)\right| .
$$

Remark 2.1 The restriction to the time interval $[0,1]$ entails no loss of generality: the results apply to arbitrary bounded intervals by rescaling $\epsilon$.

Regularity assumptions Given a function $g: \mathbb{R}^{d} \rightarrow \mathbb{R}^{n}$, we define $\|g\|_{\text {Lip }}=$ $\max \left\{|g|_{\infty}, \operatorname{Lip} g\right\}$ where $\operatorname{Lip} g=\sup _{x \neq x^{\prime}}\left|g(x)-g\left(x^{\prime}\right)\right| /\left|x-x^{\prime}\right|$ and $\left|x-x^{\prime}\right|=$ $\max _{i=1, \ldots, n}\left|x_{i}-x_{i}^{\prime}\right|$.

In this section, and also in Appendix A, we consider functions $g: \mathbb{R}^{d} \times M \times$ $\left[0, \epsilon_{0}\right) \rightarrow \mathbb{R}^{n}$ where there is no metric structure assumed on $M$. In that case, $\|g\|_{\text {Lip }}=$ $\sup _{y \in M} \sup _{\epsilon \in\left[0, \epsilon_{0}\right)}\|g(\cdot, y, \epsilon)\|_{\text {Lip }}$. If $E \subset \mathbb{R}^{d}$, then $\left\|\left.g\right|_{E}\right\|_{\text {Lip }}$ is computed by restricting to $x, x^{\prime} \in E\left(\right.$ and $\left.y \in M, \epsilon \in\left[0, \epsilon_{0}\right)\right)$.

Throughout, we write $D=\frac{\partial}{\partial x}$. If $g: \mathbb{R}^{d} \times M \times\left[0, \epsilon_{0}\right) \rightarrow \mathbb{R}^{n}$, then $D g: \mathbb{R}^{d} \times$ $M \times\left[0, \epsilon_{0}\right) \rightarrow \mathbb{R}^{n \times d}$ and $\|D g\|_{\text {Lip }}$ is defined accordingly. Similarly for $\left\|\left.D g\right|_{E}\right\|_{\text {Lip }}$ when $E \subset \mathbb{R}^{d}$.

Below, $L_{1}, L_{2}, L_{3} \geq 1$ are constants. We require that $a$ is globally Lipschitz in $x$ :

$$
\begin{equation*}
\|a\|_{\text {Lip }} \leq L_{1} \tag{2.3}
\end{equation*}
$$

Set $E=\left\{x \in \mathbb{R}^{d}:\left|x-x_{0}\right| \leq L_{1}\right\}$. We assume that $\left.a\right|_{E}$ is differentiable as a function of $x$ with Lipschitz derivative:

$$
\begin{equation*}
\left\|\left.D a\right|_{E}\right\|_{\text {Lip }} \leq L_{2} \tag{2.4}
\end{equation*}
$$

and that

$$
\begin{equation*}
\sup _{x \in E} \sup _{y \in M}|a(x, y, \epsilon)-a(x, y, 0)| \leq L_{3} \epsilon \tag{2.5}
\end{equation*}
$$

In the sequel we let $L=\max \left\{L_{1}, L_{2}, L_{3}\right\}$.

### 2.1 Order functions and a general averaging theorem

Define $\bar{a}(x, \epsilon)=\int_{M} a(x, y, \epsilon) d \nu_{\epsilon}(y)$ and let $v_{\epsilon, x}(y)=a(x, y, \epsilon)-\bar{a}(x, \epsilon)$. Set

$$
\begin{aligned}
& \delta_{1, \epsilon}=\sup _{x \in E} \sup _{1 \leq n \leq 1 / \epsilon} \epsilon\left|v_{\epsilon, x, n}\right|, \quad \text { where } \quad v_{\epsilon, x, n}=\sum_{j=0}^{n-1} v_{\epsilon, x} \circ T_{\epsilon}^{j} \\
& \delta_{2, \epsilon}=\sup _{x \in E} \sup _{1 \leq n \leq 1 / \epsilon} \epsilon\left|\mathbf{V}_{\epsilon, x, n}\right|, \quad \text { where } \quad \mathbf{V}_{\epsilon, x, n}=\sum_{j=0}^{n-1}\left(D v_{\epsilon, x}\right) \circ T_{\epsilon}^{j} .
\end{aligned}
$$

Then we define the order function $\delta_{\epsilon}=\delta_{1, \epsilon}+\delta_{2, \epsilon}: M \rightarrow \mathbb{R}$.
Finally, define $S_{\epsilon}=\sup _{x \in E}\left|\int_{M} a(x, y, 0)\left(d \nu_{\epsilon}-d \nu_{0}\right)(y)\right|+\epsilon$.
Theorem 2.2 If $\delta_{\epsilon} \leq \frac{1}{2}$, then $z_{\epsilon} \leq 6 e^{2 L}\left(\delta_{\epsilon}+S_{\epsilon}\right)$.
The proof of Theorem 2.2 is postponed to Appendix A.
Remark 2.3 For averaging without rates, it suffices instead of condition (2.5) that $\lim _{\epsilon \rightarrow 0}|a(x, y, \epsilon)-a(x, y, 0)|=0$ for all $x \in E, y \in M$.

Remark 2.4 As shown in Section 7, almost sure convergence in the averaging theorem is not likely to hold for fast-slow systems of type (2.1). Hence we consider convergence in $L^{q}$ with respect to certain absolutely continuous probability measures on $M$. Since $z_{\epsilon} \leq 2 L$ and $\delta_{\epsilon} \leq 4 L$, convergence in $L^{p}$ is equivalent to convergence in $L^{q}$ for all $p, q \in(0, \infty)$. For brevity, we restrict statements to convergence in $L^{1}$ except when speaking of rates.

Remark 2.5 If condition (2.4) fails, then all of our results without rates go through unchanged. Moreover, it is still possible to obtain results with rates but usually with weaker rates of convergence (the best rates are $O\left(\epsilon^{\frac{1}{4}-}\right)$ instead of $O\left(\epsilon^{\frac{1}{2}-}\right)$ ). These results are obtained by using $\delta_{1, \epsilon}$ (first order averaging) instead of $\delta_{\epsilon}=\delta_{1, \epsilon}+\delta_{2, \epsilon}$ (second order averaging) and can be found in the first (much longer) version of this paper [26].

According to Theorem 2.2, results on averaging reduce to estimating the scalar quantity $S_{\epsilon}$ and the random variable $\delta_{\epsilon}=\delta_{\epsilon}\left(y_{0}\right)$. These quantities are discussed below in Subsections 2.2 and 2.3 respectively.

### 2.2 Statistical stability

In this subsection, we suppose that $M$ is a topological space and that the $\sigma$-algebra of measurable sets is the $\sigma$-algebra of Borel sets. Recall that the family of measures $\nu_{\epsilon}$ is statistically stable at $\epsilon=0$ if $\nu_{0}$ is the weak limit of $\nu_{\epsilon}$ as $\epsilon \rightarrow 0\left(\nu_{\epsilon} \rightarrow_{w} \nu_{0}\right)$. This means that $\int_{M} \phi d \nu_{\epsilon} \rightarrow \int_{M} \phi d \nu_{0}$ for all continuous bounded functions $\phi: M \rightarrow \mathbb{R}$.

In the noninvertible setting, often a stronger property known as strong statistical stability holds. Let $m$ be a reference measure on $M$ and suppose that $\nu_{\epsilon}$ is absolutely continuous with respect to $m$ for all $\epsilon \geq 0$. Then $\nu_{0}$ is strongly statistically stable if the densities $\rho_{\epsilon}=d \nu_{\epsilon} / d m$ satisfy $\lim _{\epsilon \rightarrow 0} R_{\epsilon}=0$ where $R_{\epsilon}=\int_{M}\left|\rho_{\epsilon}-\rho_{0}\right| d m$. We note that $S_{\epsilon} \leq L R_{\epsilon}+\epsilon$.

Proposition 2.6 If $\nu_{\epsilon} \rightarrow_{w} \nu_{0}$, then $\lim _{\epsilon \rightarrow 0} S_{\epsilon}=0$.
Proof Let $A_{\epsilon}(x)=\int_{M} a(x, y, 0) d \nu_{\epsilon}(y)-\int_{M} a(x, y, 0) d \nu_{0}(y)$. Let $\delta>0$. Since $\nu_{\epsilon} \rightarrow_{w}$ $\nu_{0}$, we have that $A_{\epsilon}(x) \rightarrow 0$ for each $x$, so there exists $\epsilon_{x}>0$ such that $\left|A_{\epsilon}(x)\right|<\delta$ for all $\epsilon \in\left(0, \epsilon_{x}\right)$. Moreover, $\left|A_{\epsilon}(z)\right|<2 \delta$ for all $\epsilon \in\left(0, \epsilon_{x}\right)$ and $z \in B_{\delta /(2 L)}(x)$. Since $E$ is covered by finitely many such balls $B_{\delta /(2 L)}(x)$, there exists $\bar{\epsilon}>0$ such that $\sup _{x \in E}\left|A_{\epsilon}(x)\right|<2 \delta$ for all $\epsilon \in(0, \bar{\epsilon})$. Hence $\int_{M} a(x, y, 0)\left(d \nu_{\epsilon}-d \nu_{0}\right)(y)$ converges to zero uniformly in $x$.

Hence for proving averaging theorems, statistical stability takes care of the term $S_{\epsilon}$ in Theorem 2.2. In specific examples, we are able to appeal to results on statistical stability with rates, yielding effective estimates.

Proposition 2.7 Let $q \geq 1$. There is a constant $C>0$ such that

$$
\left|z_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)} \leq C\left(\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}+S_{\epsilon}\right),
$$

for all $\epsilon \in\left[0, \epsilon_{0}\right)$.
If the measures $\nu_{\epsilon}$ are absolutely continuous with respect to $m$, then there is a constant $C>0$ such that

$$
\left|z_{\epsilon}\right|_{L^{q}\left(\nu_{0}\right)} \leq C\left(\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}+R_{\epsilon}^{1 / q}+\epsilon\right)
$$

for all $\epsilon \in\left[0, \epsilon_{0}\right)$.
Proof Let $A_{\epsilon}=\left\{y \in M: \delta_{\epsilon}(y) \leq \frac{1}{2}\right\}$. Then Theorem 2.2 applies on $A_{\epsilon}$ and

$$
\begin{aligned}
\int_{M} z_{\epsilon}^{q} d \nu_{\epsilon} & =\int_{M \backslash A_{\epsilon}} z_{\epsilon}^{q} d \nu_{\epsilon}+\int_{A_{\epsilon}} z_{\epsilon}^{q} d \nu_{\epsilon} \leq(2 L)^{q} \nu_{\epsilon}\left(\delta_{\epsilon}>\frac{1}{2}\right)+\left(6 e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+S_{\epsilon}\right)^{q} d \nu_{\epsilon} \\
& \leq(4 L)^{q} \int_{M} \delta_{\epsilon}^{q} d \nu_{\epsilon}+\left(6 e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+S_{\epsilon}\right)^{q} d \nu_{\epsilon} \leq\left(12 e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+S_{\epsilon}\right)^{q} d \nu_{\epsilon}
\end{aligned}
$$

Hence

$$
\left|z_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)} \leq 12 e^{2 L}\left|\delta_{\epsilon}+S_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)} \leq 12 e^{2 L}\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}+12 e^{2 L} S_{\epsilon},
$$

yielding the first estimate.
Next,

$$
\begin{aligned}
\int_{M} z_{\epsilon}^{q} d \nu_{0} & =\int_{M} z_{\epsilon}^{q} d \nu_{\epsilon}+\int_{M} z_{\epsilon}^{q}\left(d \nu_{0}-d \nu_{\epsilon}\right) \leq\left(12 e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+S_{\epsilon}\right)^{q} d \nu_{\epsilon}+(2 L)^{q} R_{\epsilon} \\
& \leq\left(12 e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+L R_{\epsilon}+\epsilon\right)^{q} d \nu_{\epsilon}+(2 L)^{q} R_{\epsilon} \\
& \leq\left(12 L e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+R_{\epsilon}^{1 / q}+\epsilon\right)^{q} d \nu_{\epsilon}+(2 L)^{q} R_{\epsilon} \\
& \leq\left(24 L e^{2 L}\right)^{q} \int_{M}\left(\delta_{\epsilon}+R_{\epsilon}^{1 / q}+\epsilon\right)^{q} d \nu_{\epsilon}
\end{aligned}
$$

Hence

$$
\left|z_{\epsilon}\right|_{L^{q}\left(\nu_{0}\right)} \leq 24 L e^{2 L}\left|\delta_{\epsilon}+R_{\epsilon}^{1 / q}+\epsilon\right|_{L^{q}\left(\nu_{\epsilon}\right)} \leq 24 L e^{2 L}\left(\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}+R_{\epsilon}^{1 / q}+\epsilon\right),
$$

yielding the second estimate.

Corollary 2.8 (a) Assume statistical stability and that $\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{\epsilon} d \nu_{\epsilon}=0$. Then $\lim _{\epsilon \rightarrow 0} \int_{M} z_{\epsilon} d \nu_{\epsilon}=0$.
(b) Assume in addition strong statistical stability and that $\mu$ is a probability measure on $M$ with $\mu \ll \nu_{0}$. Then $\lim _{\epsilon \rightarrow 0} \int_{M} z_{\epsilon} d \mu=0$.

Proof Part (a), and part (b) in the special case $\mu=\nu_{0}$, are immediate from Proposition 2.7. To prove the general case of part (b), suppose for contradiction that $\int_{M} z_{\epsilon_{k}} d \mu \rightarrow b>0$ along some subsequence $\epsilon_{k} \rightarrow 0$. Since $\int_{M} z_{\epsilon_{k}} d \nu_{0} \rightarrow 0$, by passing without loss to a further subsequence, we can suppose also that $z_{\epsilon_{k}} \rightarrow 0$ on a set of full measure with respect to $\nu_{0}$ and hence with respect to $\mu$. By the bounded convergence theorem, $\int_{M} z_{\epsilon_{k}} d \mu \rightarrow 0$ which is the desired contradiction.

The next result is useful in situations where $\nu_{0}$ is absolutely continuous but whose support is not the whole of $M$.

Corollary 2.9 Assume strong statistical stability and that $\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{\epsilon} d \nu_{0}=0$. Suppose further that each $T_{\epsilon}$ is nonsingular with respect to $m$ and that for almost every $y \in M$, there exists $N \geq 1, \epsilon_{1} \in\left(0, \epsilon_{0}\right)$ such that $T_{\epsilon}^{N} y \in \operatorname{supp} \nu_{0}$ for all $\epsilon \in\left[0, \epsilon_{1}\right]$. Then $\lim _{\epsilon \rightarrow 0} \int_{M} z_{\epsilon} d \mu=0$ for every probability measure $\mu$ on $M$ with $\mu \ll m$.

Proof First, we note that for all $N \geq 1, \epsilon \geq 0$,

$$
\begin{equation*}
\left|\delta_{\epsilon} \circ T_{\epsilon}^{N}-\delta_{\epsilon}\right|_{\infty} \leq 8 L N \epsilon \tag{2.6}
\end{equation*}
$$

By the arguments in the proof of Corollary 2.8, it suffices to prove that $\int_{M} \delta_{\epsilon} d m \rightarrow 0$. Suppose this is not the case. By Corollary 2.8(b), $\int_{\text {supp } \nu_{0}} \delta_{\epsilon} d m \rightarrow 0$. Hence there exists a subsequence $\epsilon_{k} \rightarrow 0$ and a subset $A \subset \operatorname{supp} \nu_{0}$ with $m\left(\operatorname{supp} \nu_{0} \backslash\right.$ $A)=0$ such that (i) $\delta_{\epsilon_{k}} \rightarrow 0$ pointwise on $A$ and (ii) $\int_{M} \delta_{\epsilon_{k}} d m \rightarrow b>0$.

Since each $T_{\epsilon}$ is nonsingular, there exists $M^{\prime} \subset M$ with $m\left(M^{\prime}\right)=1$ such that $M^{\prime} \cap T_{\epsilon_{k}}^{-n}\left(\operatorname{supp} \nu_{0} \backslash A\right)=\emptyset$ for all $k, n \geq 1$. By the hypothesis of the result, there is a subset $M^{\prime \prime} \subset M^{\prime}$ with $m\left(M^{\prime \prime}\right)=1$ such that for any $y \in M^{\prime \prime}$ there exists $N \geq 1$ such that $T_{\epsilon_{k}}^{N} y \in A$ for all $k$ sufficiently large. Hence it follows from (i) and (2.6) that $\delta_{\epsilon_{k}} \rightarrow 0$ pointwise on $M^{\prime \prime}$. By the bounded convergence theorem $\int_{M} \delta_{\epsilon_{k}} d m \rightarrow 0$. Together with (ii), this yields the desired contradiction.

Remark 2.10 The hypotheses of Corollary 2.9 are particularly straightforward to apply when $\operatorname{supp} \nu_{0}$ has nonempty interior. This property is automatic for large classes of nonuniformly expanding maps, see [3, Lemma 5.6] (taking $G=\operatorname{supp} \nu_{0} \cap H(\sigma)$, it follows that $\bar{G}$ contains a disk).

To be more specific, we consider the examples of the logistic family restricted to Collet-Eckmann parameters and Viana maps, referring forwards to Examples 5.2 and 5.4 respectively. In both cases the attractors $\Lambda_{\epsilon}=\operatorname{supp} \nu_{\epsilon}$ have nonempty interior (for the logistic family, it is standard that $\Lambda_{\epsilon}$ is a finite union of intervals; for the Viana maps, we apply [3, Lemma 5.6]).

In the case of the logistic family, defined on $M=[-1,1]$, almost every initial condition $y \in M$ is attracted under $T_{\epsilon}$ to $\Lambda_{\epsilon}$. The same is true for the Viana maps on restricting to $M=S^{1} \times I$.

Choose a nonempty open set $U \subset \Lambda_{0}$. Almost every point in $M$ is attracted to $\Lambda_{0}$ and almost every point in $\Lambda_{0}$ has an iterate inside $U$. Hence for almost every fixed
$y \in M$, we can choose $N$ so that $T_{0}^{N} y \in U$. By continuity, we can choose $\epsilon_{1}$ so that $T_{\epsilon}^{N} y \in U$ for all $\epsilon \in\left[0, \epsilon_{1}\right]$. Hence the hypotheses of Corollary 2.9 are satisfied, and we conclude that $\lim _{\epsilon \rightarrow 0} \int_{M} z_{\epsilon} d \mu=0$ for every absolutely continuous measure $\mu$ on $M$.

### 2.3 Estimating the order function

By Proposition 2.7 and Corollary 2.8, it remains to deal with the order function $\delta_{\epsilon}$. This is a random variable depending on the initial condition $y_{0}$. Here we give a useful estimate.

Since $\delta_{\epsilon}=\delta_{1, \epsilon}+\delta_{2, \epsilon}$ and the definition of $\delta_{2, \epsilon}$ is identical to that of $\delta_{1, \epsilon}$ with $a$ replaced by $D a$, it suffices to consider $\delta_{1, \epsilon}$. From now on, $\Gamma$ denotes a constant that only depends on $d, p, L$ and whose value may change from line to line.

Lemma 2.11 Let $\mu$ be a probability measure on $M$. Then for all $p \geq 0, \epsilon \in\left[0, \epsilon_{0}\right)$,

$$
\int_{M} \delta_{1, \epsilon}^{p+d+1} d \mu \leq \Gamma \epsilon^{p} \sup _{x \in E} \sup _{1 \leq n \leq 1 / \epsilon} \int_{M}\left|v_{\epsilon, x, n}\right|^{p} d \mu .
$$

Proof For most of the proof, we work pointwise on $M$ suppressing the initial condition $y_{0} \in M$. There exist $\tilde{x} \in E, \tilde{n} \in[0,1 / \epsilon]$ such that $\delta_{1, \epsilon}=\epsilon \mid v_{\epsilon, \tilde{x}, \tilde{n} \mid}$. Observe that

$$
\left|v_{\epsilon, x, n}-v_{\epsilon, \tilde{x}, \tilde{n}}\right| \leq\left|v_{\epsilon, x, n}-v_{\epsilon, \tilde{x}, n}\right|+\left|v_{\epsilon, \tilde{x}, n}-v_{\epsilon, \tilde{x}, \tilde{n}}\right| \leq 2 L \epsilon^{-1}|x-\tilde{x}|+2 L|n-\tilde{n}|
$$

for every $x \in E$ and $n \leq 1 / \epsilon$. Define

$$
A=\left\{x \in E:|x-\tilde{x}| \leq \frac{1}{8 L} \delta_{1, \epsilon}\right\}, \quad B=\left\{n \in\left[0, \epsilon^{-1}\right]: \epsilon|n-\tilde{n}| \leq \frac{1}{8 L} \delta_{1, \epsilon}\right\}
$$

Then for every $x \in A$ and $n \in B$ we have $\epsilon\left|v_{\epsilon, x, n}\right| \geq \delta_{1, \epsilon} / 2$. Moreover, since $\delta_{1, \epsilon} \leq 2 L$, we have $\operatorname{Leb}(A) \geq \Gamma \delta_{1, \epsilon}^{d}$ and $\# B \geq \epsilon^{-1} \delta_{1, \epsilon} / 8 L$. Hence

$$
\epsilon^{p} \sum_{n=0}^{[1 / \epsilon]-1} \int_{E}\left|v_{\epsilon, x, n}\right|^{p} d x \geq(\# B) \operatorname{Leb}(A)\left(\frac{\delta_{1, \epsilon}}{2}\right)^{p} \geq \Gamma \epsilon^{-1} \delta_{1, \epsilon}^{p+d+1}
$$

Finally,

$$
\int_{M} \delta_{1, \epsilon}^{p+d+1} d \mu \leq \Gamma \epsilon^{p+1} \sum_{n=0}^{[1 / \epsilon]-1} \int_{E} \int_{M}\left|v_{\epsilon, x, n}\right|^{p} d \mu d x \leq \Gamma \epsilon^{p} \sup _{x \in E} \sup _{1 \leq n \leq 1 / \epsilon} \int_{M}\left|v_{\epsilon, x, n}\right|^{p} d \mu
$$

as required.
Remark 2.12 Often, estimating $\int_{M}\left|v_{\epsilon, x, n}\right| d \mu$ leads to an essentially identical estimate for $\int_{M} \sup _{1 \leq n \leq 1 / \epsilon}\left|v_{\epsilon, x, n}\right| d \mu$. In this case, slightly better convergence rates for $\delta_{1, \epsilon}$ can be obtained using the estimate

$$
\begin{equation*}
\int_{M} \delta_{1, \epsilon}^{p+d} d \mu \leq \Gamma \epsilon^{p} \sup _{x \in E} \int_{M} \sup _{1 \leq n \leq 1 / \epsilon}\left|v_{\epsilon, x, n}\right|^{p} d \mu \tag{2.7}
\end{equation*}
$$

for all $p \geq 0, \epsilon \in\left[0, \epsilon_{0}\right)$.

## 3 Examples: Uniformly expanding maps

Let $T_{\epsilon}: M \rightarrow M$ be a family of maps defined on a metric space $\left(M, d_{M}\right)$ with invariant ergodic Borel probability measures $\nu_{\epsilon}$. Let $P_{\epsilon}$ denote the corresponding transfer operators, so $\int_{M} P_{\epsilon} v w d \nu_{\epsilon}=\int_{M} v w \circ T_{\epsilon} d \nu_{\epsilon}$ for all $v \in L^{1}\left(\nu_{\epsilon}\right), w \in L^{\infty}\left(\nu_{\epsilon}\right)$.

From now on, we require Lipschitz regularity in the $M$ variables in addition to the $\mathbb{R}^{d}$ variables as was required in assumptions (2.3) and (2.4). So for $g: \mathbb{R}^{d} \times$ $M \times\left[0, \epsilon_{0}\right) \rightarrow \mathbb{R}^{n}$ we define $\|g\|_{\text {Lip }}=|g|_{\infty}+\sup _{\epsilon \in\left[0, \epsilon_{0}\right)} \operatorname{Lip} g(\cdot, \cdot, \epsilon)$ where $\operatorname{Lip} g(\cdot, \cdot, \epsilon)=$ $\sup _{x \neq x^{\prime}} \sup _{y \neq y^{\prime}}\left|g(x, y, \epsilon)-g\left(x^{\prime}, y^{\prime}, \epsilon\right)\right| /\left(\left|x-x^{\prime}\right|+d_{M}\left(y, y^{\prime}\right)\right)$. We continue to assume conditions (2.3) - (2.5) with this modified definition of $\left\|\|_{\text {Lip }}\right.$.

Proposition 3.1 Suppose that there is a sequence of constants $a_{n} \rightarrow 0$ such that $\int_{M}\left|P_{\epsilon}^{n} v-\int_{M} v d \mu_{\epsilon}\right| \leq a_{n}\|v\|_{\text {Lip }}$ for all Lipschitz $v: M \rightarrow \mathbb{R}$ and all $n \geq 1, \epsilon \geq 0$. Then $\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{\epsilon} d \nu_{\epsilon}=0$.

Proof We prove the result for $\delta_{1, \epsilon}$ and $\delta_{2, \epsilon}$ separately. By Remark 2.4, we can work in $L^{q}$ for any choice of $q$ and we take $q=d+3$.

For every Lipschitz $v$, we have

$$
\begin{aligned}
\int_{M}\left(\sum_{j=0}^{n-1} v \circ T_{\epsilon}^{j}\right)^{2} d \nu_{\epsilon} & =\sum_{j=0}^{n} \int_{M} v^{2} \circ T_{\epsilon}^{j} d \nu_{\epsilon}+2 \sum_{0 \leq i<j \leq n-1} \int_{M} v \circ T_{\epsilon}^{i} v \circ T_{\epsilon}^{j} d \nu_{\epsilon} \\
& =n \int_{M} v^{2} d \nu_{\epsilon}+2 \sum_{1 \leq k \leq n-1}(n-k) \int_{M} v v \circ T_{\epsilon}^{k} d \nu_{\epsilon} \\
& =n \int_{M} v^{2} d \nu_{\epsilon}+2 \sum_{1 \leq k \leq n-1}(n-k) \int_{M} P_{\epsilon}^{k} v v d \nu_{\epsilon}
\end{aligned}
$$

Hence for $v$ Lipschitz and mean zero,

$$
\int_{M}\left(\sum_{j=0}^{n-1} v \circ T_{\epsilon}^{j}\right)^{2} d \nu_{\epsilon} \leq b_{n}\|v\|_{\text {Lip }}^{2}
$$

where $b_{n}=n+2 n \sum_{1 \leq k \leq n-1} a_{k}=o\left(n^{2}\right)$ by the assumption on $a_{n}$.
By condition (2.3), $\left\|v_{\epsilon, x}\right\|_{\text {Lip }} \leq 2 L$ for all $\epsilon, x$, $\operatorname{so~}_{\sup }^{x \in E}$ $\sup _{1 \leq n \leq 1 / \epsilon} \int_{M}\left|v_{\epsilon, x, n}\right|^{2} d \nu_{\epsilon} \leq$ $4 L^{2} b_{n}$. By Lemma 2.11, it follows that $\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{1, \epsilon}^{d+3} d \nu_{\epsilon}=0$. Similarly, by condition (2.4), $\left\|\mathbf{V}_{\epsilon, x}\right\|_{\text {Lip }} \leq 2 L$, so $\sup _{x \in E} \sup _{1 \leq n \leq 1 / \epsilon} \int_{M}\left|\mathbf{V}_{\epsilon, x, n}\right|^{2} d \nu_{\epsilon} \leq 4 L^{2} b_{n}$ and hence $\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{2, \epsilon}^{d+3} d \nu_{\epsilon}=0$.

Remark 3.2 The proof uses only that $n^{-1} \sum_{k=1}^{n} a_{n} \rightarrow 0$.
Proposition 3.1 is useful in situations where $T_{\epsilon}$ is a family of (piecewise) uniformly expanding maps. A general result of Keller \& Liverani [21] guarantees uniform spectral properties of the transfer operators $P_{\epsilon}$ under mild conditions, and consequently
$\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{\epsilon}^{q} d \nu_{\epsilon}=0$ for all $q$. We mention two situations where this idea can be applied. Again, for brevity we work in $L^{1}$ except when discussing convergence rates (see Remark 2.4).

Example 3.3 (Uniformly expanding maps) Suppose that $M=\mathbb{T}^{k} \cong \mathbb{R}^{/} \mathbb{Z}^{k}$ is a torus with Haar measure $m$ and distance $d_{M}$ inherited from Euclidean distance on $\mathbb{R}^{k}$ and normalised so that $\operatorname{diam} M=1$. We say that a $C^{2} \operatorname{map} T: M \rightarrow M$ is uniformly expanding if there exists $\lambda>1$ such that $\left|(D T)_{y} v\right| \geq \lambda|v|$ for all $y \in M$, $v \in \mathbb{R}^{k}$. There is a unique absolutely continuous invariant probability measure, and the density is $C^{1}$ and nonvanishing.

If $T_{\epsilon}: M \rightarrow M, \epsilon \in[0,1]$, is a continuous family of $C^{2}$ maps, each of which is uniformly expanding, with corresponding probability measures $\nu_{\epsilon}$, then it follows from [21] that we are in the situation of Proposition 3.1, and so $\lim _{\epsilon \rightarrow 0} \int_{M} \delta_{\epsilon} d \nu_{\epsilon}=0$.

Moreover, it is well-known that $\nu_{0}$ is uniformly equivalent to $m$ and is strongly statistically stable. Hence by Corollary 2.8 we obtain the averaging result $\lim _{\epsilon \rightarrow 0} \int_{M} z_{\epsilon} d \nu_{\epsilon}=0$ and $\lim _{\epsilon \rightarrow 0} \int_{M} z_{\epsilon} d \mu=0$ for every absolutely continuous probability measure $\mu$.

Suppose further that $T_{\epsilon}: M \rightarrow M, \epsilon \in[0,1]$, is a $C^{k}$ family of $C^{2}$ maps, for some $k \in(0,1]$. By standard results (for instance [21] with Banach spaces $C^{0}$ and $C^{1}$ ), $R_{\epsilon}=\int_{M}\left|\rho_{\epsilon}-\rho_{0}\right| d m=O\left(\epsilon^{k}\right)$. If $k \in\left(0, \frac{1}{2}\right)$, then by Remark 4.4 below we obtain the convergence rate $O\left(\epsilon^{k}\right)$ for $z_{\epsilon}$ in $L^{q}\left(\nu_{\epsilon}\right)$ and $L^{q}(m)$ for all $q>0$. If $k \geq \frac{1}{2}$, then the convergence rate for $z_{\epsilon}$ is $O\left(\epsilon^{\frac{1}{2}-}\right)$.

Example 3.4 (Piecewise uniformly expanding maps) Let $M=[-1,1]$ with Lebesgue measure $m$. We consider continuous maps $T: M \rightarrow M$ with $T(-1)=$ $T(1)=-1$ such that $T$ is $C^{2}$ on $[-1,0]$ and $[0,1]$. We require that there exists $\lambda>1$ such that $T^{\prime} \geq \lambda$ on $[-1,0)$ and $T^{\prime} \leq-\lambda$ on $(0,1]$. There exists a unique absolutely continuous invariant probability measure with density of bounded variation.

The results are analogous to those in Example 3.3. Let $\epsilon \rightarrow T_{\epsilon}$ be a continuous family of such maps on $[-1,1]$ with associated measures $\nu_{\epsilon}$. We assume that $T_{0}$ is topologically mixing on the interval $\left[T_{0}^{2}(0), T_{0}(0)\right]$ and that 0 is not periodic (which guarantees that $T_{\epsilon}$ is mixing for $\epsilon$ small).

Then $\nu_{0}$ is strongly statistically stable, so by Corollaries 2.8 and 2.9 we obtain averaging in $L^{1}\left(\nu_{\epsilon}\right)$ and also in $L^{1}(\mu)$ for every absolutely continuous probability measure $\mu$.

Now suppose that $\epsilon \rightarrow T_{\epsilon}$ is a $C^{1}$ family of such maps on $[-1,1]$ with densities $\rho_{\epsilon}=d \nu_{\epsilon} / d m$. Keller [20] showed that $\epsilon \rightarrow \rho_{\epsilon}$ is $C^{1-}$ as a map into $L^{1}$ densities. Hence we obtain the convergence rate $O\left(\epsilon^{\frac{1}{2}-}\right)$ for $z_{\epsilon}$ in $L^{q}\left(\nu_{\epsilon}\right)$, for all $q>0$ and in $L^{1}\left(\nu_{0}\right)$.

More precisely, [20] shows that $\int_{M}\left|\rho_{\epsilon}-\rho_{0}\right| d m=O\left(\epsilon \log \epsilon^{-1}\right)$. By [8], this estimate is optimal, so this is a situation where linear response fails in contrast to Example 3.3.

## 4 Families of nonuniformly expanding maps

In this section, we consider the situation where the fast dynamics is generated by nonuniformly expanding maps $T_{\epsilon}$, such that the nonuniform expansion is uniform in the parameter $\epsilon$.

In Subsection 4.1, we recall the notion of nonuniformly expanding map. In Subsection 4.2, we describe the uniformity criteria on the family $T_{\epsilon}$ and state our main result on averaging, Theorem 4.3, for such families. In Subsections 4.3 and 4.4 we establish some basic estimates for uniformly and nonuniformly expanding maps. In Subsection 4.5 we prove Theorem 4.3.

### 4.1 Nonuniformly expanding maps

Let $\left(M, d_{M}\right)$ be a locally compact separable bounded metric space with finite Borel measure $m$ and let $T: M \rightarrow M$ be a nonsingular transformation for which $m$ is ergodic. Let $Y \subset M$ be a subset of positive measure, and let $\alpha$ be an at most countable measurable partition of $Y$ with $m(a)>0$ for all $a \in \alpha$. We suppose that there is an $L^{1}$ return time function $\tau: Y \rightarrow \mathbb{Z}^{+}$, constant on each $a$ with value $\tau(a) \geq 1$, and constants $\lambda>1, \eta \in(0,1], C_{0}, C_{1} \geq 1$ such that for each $a \in \alpha$,
(1) $F=T^{\tau}$ restricts to a (measure-theoretic) bijection from $a$ onto $Y$.
(2) $d_{M}(F x, F y) \geq \lambda d_{M}(x, y)$ for all $x, y \in a$.
(3) $d_{M}\left(T^{\ell} x, T^{\ell} y\right) \leq C_{0} d_{M}(F x, F y)$ for all $x, y \in a, 0 \leq \ell<\tau(a)$.
(4) $\zeta=\frac{\left.d m\right|_{Y}}{\left.d m\right|_{Y} \circ F}$ satisfies $|\log \zeta(x)-\log \zeta(y)| \leq C_{1} d_{M}(F x, F y)^{\eta}$ for all $x, y \in a$.

Such a dynamical system $T: M \rightarrow M$ is called nonuniformly expanding. We refer to $F=T^{\tau}: Y \rightarrow Y$ as the induced map. (It is not required that $\tau$ is the first return time to $Y$.) It follows from standard results (recalled later) that there is a unique absolutely continuous ergodic $T$-invariant probability measure $\nu$ on $M$.

Remark 4.1 The uniformly expanding maps in Example 3.3 are clearly nonuniformly expanding: take $Y=M, \eta=1, \tau=1$. Then conditions (1) and (2) are immediate, condition (3) is vacuously satisfied, and condition (4) holds with $C_{1}=\sup _{x, y \in M, x \neq y}\left|\left(D T_{\epsilon}\right)_{x}-\left(D T_{\epsilon}\right)_{y}\right| / d_{M}(x, y)$.

### 4.2 Uniformity assumptions

Now suppose that $T_{\epsilon}: M \rightarrow M, \epsilon \in\left[0, \epsilon_{0}\right)$, is a family of nonuniformly expanding maps as defined in Subsection 4.1, with corresponding absolutely continuous ergodic invariant probability measures $\nu_{\epsilon}$.

Definition 4.2 Let $p>1$. We say that $T_{\epsilon}: M \rightarrow M$ is a uniform family of nonuniformly expanding maps (of order $p$ ) if
(i) The constants $C_{0}, C_{1} \geq 1, \lambda>1, \eta \in(0,1]$ can be chosen independent of $\epsilon \in\left[0, \epsilon_{0}\right)$.
(ii) The return time functions $\tau_{\epsilon}: Y_{\epsilon} \rightarrow \mathbb{Z}^{+}$lie in $L^{p}$ for all $\epsilon \in\left[0, \epsilon_{0}\right)$, and moreover $\sup _{\epsilon \in\left[0, \epsilon_{0}\right)} \int_{Y_{\epsilon}}\left|\tau_{\epsilon}\right|^{p} d m<\infty$.

We can now state our main result for this section. Recall the set up in Section 2.
Theorem 4.3 If $T_{\epsilon}: M \rightarrow M$ is a uniform family of nonuniformly expanding maps of order $p$, then there is a constant $C>0$ such that for all $\epsilon \in\left[0, \epsilon_{0}\right)$

$$
\int_{M} \delta_{\epsilon}^{p+d-1} d \nu_{\epsilon} \leq \begin{cases}C \epsilon^{(p-1) / 2}, & p>2 \\ C \epsilon^{(p-1)^{2} / p}, & p \in(1,2]\end{cases}
$$

Remark 4.4 In the case $p>2$, it follows from Theorem 4.3 that

$$
\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}=O\left(\epsilon^{(p-1) /(2(p+d-1))}\right),
$$

for all $q \leq p+d-1$. Since $\delta_{\epsilon}$ is uniformly bounded, $\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}=O\left(\epsilon^{(p-1) /(2 q)}\right)$ for all $q>p+d-1$. Similar comments apply for $p \in(1,2]$.

In particular, if $p$ can be taken arbitrarily large in Definition 4.2, then we obtain that $\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}=O\left(\epsilon^{\frac{1}{2}-}\right)$ for all $q>0$.

If in addition $\nu_{0}$ is strongly statistically stable with $R_{\epsilon}=\int_{M}\left|\rho_{\epsilon}-\rho_{0}\right| d m=O\left(\epsilon^{\frac{1}{2}-}\right)$, then by Proposition 2.7 we obtain $\left|z_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}=O\left(\epsilon^{\frac{1}{2}-}\right)$ for all $q>0$ and $\left|z_{\epsilon}\right|_{L^{1}\left(\nu_{0}\right)}=$ $O\left(\epsilon^{\frac{1}{2}-}\right)$.

Remark 4.5 Alves \& Viana [5] prove strong statistical stability for a large class of noninvertible dynamical systems. These maps are uniform families of nonuniformly expanding maps in a sense that is very similar to our definition. In fact it is almost the case that their definition includes our definition, so it is almost the case that verifying the assumptions of [5] is sufficient to obtain averaging and rates of averaging via Theorem 4.3.

To be more precise, let us momentarily ignore assumption (3) in Subsection 4.1. Then Definition 4.2(i) with $\eta=1$ is immediate from [5, (U1)], and Definition 4.2(ii) is immediate from [5, (U2')] which follows from their conditions (U1) and (U2).

Hence it remains to discuss assumption (3). This assumption is not explicitly mentioned in [5] since it is not required for the statement of their main results. However, in specific applications, the hypotheses in [5] are often verified via the method of hyperbolic times [1]. When the return time function $\tau_{\epsilon}$ is a hyperbolic time, then it is automatic that $C_{0}=1$ (see for example [2, Proposition 3.3(3)]).

Alves et al. 44 introduced a general method for constructing inducing schemes, where $\tau_{\epsilon}$ is not necessarily a hyperbolic time but is close enough that $C_{0}$ can still be chosen uniformly. Alves [2] combined the methods of [4] and [5] to prove statistical stability for large classes of examples. We show now that in the situation discussed by [2], assumption (3) holds with uniform $C_{0}$ and hence our main results hold. Certain quantities $\delta_{1}>0$ and $N_{0} \geq 1$ are introduced in [2, Lemma 3.2] and [2, eq. (16)] respectively, and are explicitly uniform in $\epsilon$. Moreover $\tau_{\epsilon}=n+m$ where $n$ is a hyperbolic time and $m \leq N_{0}$ (see [2, Section 4.3]), so $C_{0}$ depends only on at most $N_{0}$ iterates of $T_{\epsilon}$. The construction in [2] (see in particular the proof of [2, Lemma 4.2]) ensures that the derivative of $T_{\epsilon}$ is bounded along these iterates, so assumption (3) holds and $C_{0}$ is uniform in $\epsilon$.

We mention also the extension of [4] due to Gouëzel [18] where $C_{0}=1$ (see [18, Theorem 3.14 )].)

Finally, we note that when [4] is used to obtain polynomial decay of correlations with rate $O\left(1 / n^{\beta}\right), \beta>0$, the resulting uniform family is of order $p=\beta+1-$. (Uniformity in $\epsilon$ in Definition 4.2(ii) follows from [2, Lemma 5.1].)

### 4.3 Explicit estimates for uniformly expanding maps

Throughout this subsection, we work with a fixed uniformly expanding map $F: Y \rightarrow$ $Y$ satisfying conditions (1), (2) and (4). Some standard constructions and estimates are described. The main novelty is that we stress the dependence of various constants on the underlying constants $C_{1}, \lambda$ and $\eta$. For convenience, we normalise the metric $d_{M}$ so that $\operatorname{diam} M=1$.

For $\theta \in(0,1)$, we define the symbolic metric $d_{\theta}(x, y)=\theta^{s(x, y)}$ where the separation time $s(x, y)$ is the least integer $n \geq 0$ such that $F^{n} x$ and $F^{n} y$ lie in distinct partition element. It is assumed that the partition $\alpha$ separates orbits of $F$, so $s(x, y)$ is finite for all $x \neq y$ guaranteeing that $d_{\theta}$ is a metric.

Given $\phi: Y \rightarrow \mathbb{R}^{d}$, we define $\|\phi\|_{\theta}=|\phi|_{\infty}+|\phi|_{\theta}$ where $|\phi|_{\theta}=\sup _{x \neq y} \mid \phi(x)-$ $\phi(y) \mid / d_{\theta}(x, y)$. Then $\phi$ is $d_{\theta}$-Lipschitz if $\|\phi\|_{\theta}<\infty$.

The assumptions on $F$ guarantee that there exists a unique absolutely continuous $F$-invariant probability measure $\mu$ on $Y$. Let $P: L^{1}(Y) \rightarrow L^{1}(Y)$ denote the (normalised) transfer operator corresponding to $F$ and $\mu$, so $\int_{Y} \phi \circ F \psi d \mu=\int_{Y} \phi P \psi d \mu$ for all $\phi \in L^{\infty}$ and $\psi \in L^{1}$. Define $g: Y \rightarrow \mathbb{R}$ by setting $\left.g\right|_{a}=\left.d \mu\right|_{a} / d\left(\left.\mu \circ F\right|_{a}\right)$ for $a \in \alpha$. Then $(P \phi)(y)=\sum_{a \in \alpha} g\left(y_{a}\right) \phi\left(y_{a}\right)$ where $y_{a}$ is the unique preimage of $y$ under $F$ lying in $a$.

Lemma 4.6 There exist constants $\theta, \gamma \in(0,1), C_{2} \geq 1$ depending continuously on $\lambda, \eta$ and $C_{1}$ such that

$$
\text { (a) } d_{M}(x, y)^{\eta} \leq d_{\theta}(x, y) \text { for all } x, y \in Y \text {. }
$$

(b) For all $x, y \in a, a \in \alpha$,

$$
\begin{equation*}
g(y) \leq C_{2} \mu(a) \quad \text { and } \quad|g(x)-g(y)| \leq C_{2} \mu(a) d_{\theta}(x, y) \tag{4.1}
\end{equation*}
$$

(c) Let $\phi: Y \rightarrow \mathbb{R}$ be $d_{\theta}$-Lipschitz with $\int \phi d \mu=0$. Then

$$
\left|P^{n} \phi\right|_{\infty} \leq C_{2} \gamma^{n}\|\phi\|_{\theta} \text { for all } n \geq 1
$$

Proof Choose $\theta=\lambda^{-\eta}$. Let $n=s(x, y)$. By condition (2),

$$
1 \geq \operatorname{diam} Y \geq d_{M}\left(F^{n} x, F^{n} y\right) \geq \lambda^{n} d_{M}(x, y)=\left(\theta^{1 / \eta}\right)^{-n} d_{M}(x, y)
$$

Hence $d_{M}(x, y)^{\eta} \leq \theta^{n}=d_{\theta}(x, y)$ proving (a).
By condition (4) and [27, Proposition 2.3], there is a constant $K$ depending continuously on $\lambda, \eta$ and $C_{1}$ such that $|\log g|_{\theta} \leq K$. Hence for $y \in a, a \in \alpha$,

$$
\mu(a)=\int_{a} g d \mu \circ F \geq\left.\inf _{a} g\right|_{a} \mu(F a)=\left.\inf g\right|_{a} \geq e^{-K} g(y)
$$

so $\left.g\right|_{a} \leq e^{K} \mu(a)$. Next, we note the inequality $t-1 \leq t \log t$ which is valid for all $t \geq 1$. Let $x, y \in a$ and suppose without loss that $g(y) \leq g(x)$. Setting $t=g(x) / g(y) \geq 1$,

$$
\frac{g(x)}{g(y)}-1 \leq \frac{g(x)}{g(y)} \log \frac{g(x)}{g(y)} \leq e^{K} K d_{\theta}(x, y)
$$

Hence $g(x)-g(y) \leq g(y) e^{K} K d_{\theta}(x, y) \leq e^{2 K} K \mu(a) d_{\theta}(x, y)$. Hence part (b) holds with $C_{2}=e^{2 K} K$.

Finally, part (c) follows for example from [27, Proposition 2.5].

### 4.4 Explicit estimates for nonuniformly expanding maps

Throughout this subsection, we work with a fixed nonuniformly expanding map $T$ satisfying assumptions (1)-(4) and such that $\tau \in L^{p}$ for some $p>1$.

There is a standard procedure to pass from the $F$-invariant ergodic absolutely continuous probability measure $\mu$ on $Y$ to a $T$-invariant ergodic absolutely continuous probability measure $\nu$ on $M$. We briefly describe this procedure, since the construction is required in the proof of Lemma 4.10. Define the Young tower [35]

$$
\begin{equation*}
\Delta=\{(y, \ell) \in Y \times \mathbb{Z}: 0 \leq \ell<\tau(y)\} \tag{4.2}
\end{equation*}
$$

with probability measure $\mu_{\Delta}=\mu \times\{$ counting $\} / \int_{Y} \tau d \mu$. Define $\pi_{\Delta}: \Delta \rightarrow M$, $\pi_{\Delta}(y, \ell)=T^{\ell} y$. Then $\nu=\left(\pi_{\Delta}\right)_{*} \mu_{\Delta}$ is the desired probability measure on $M$.

In the remainder of this subsection, $L^{q}$ norms of functions defined on $Y$ are computed using $\mu$. For functions on other spaces, the measures are indicated explicitly in the notation.

Given an observable $v: M \rightarrow \mathbb{R}^{d}$, we define the induced observable $V: Y \rightarrow \mathbb{R}$,

$$
V(y)=\sum_{\ell=0}^{\tau(y)-1} v\left(T^{\ell} y\right)
$$

If $v: M \rightarrow \mathbb{R}^{d}$ satisfies $\int_{M} v d \nu=0$, then $\int_{Y} V d \mu=0$.
Proposition 4.7 If $v: M \rightarrow \mathbb{R}^{d}$ is $d_{M}$-Lipschitz, then $P V: Y \rightarrow \mathbb{R}^{d}$ is $d_{\theta}$-Lipschitz.
Moreover,
$|V(y)| \leq \tau(a)|v|_{\infty}, \quad|V(x)-V(y)| \leq C_{0} \theta^{-1} \tau(a) \operatorname{Lip} v d_{\theta}(x, y)$, for all $x, y \in a, a \in \alpha$, and

$$
|P V|_{\infty} \leq C_{2}|\tau|_{1}|v|_{\infty}, \quad|P V|_{\theta} \leq C_{0} C_{2} \theta^{-1}|\tau|_{1}\|v\|_{\text {Lip }}
$$

Proof The estimate for $V(y)$ is immediate. By condition (3) and Lemma 4.6(a),

$$
\begin{aligned}
\mid V(x) & -V(y) \mid \leq \operatorname{Lip} v \sum_{\ell=0}^{\tau(a)-1} d_{M}\left(T^{\ell} x, T^{\ell} y\right) \leq C_{0} \operatorname{Lip} v \sum_{\ell=0}^{\tau(a)-1} d_{M}(F x, F y) \\
& =C_{0} \tau(a) \operatorname{Lip} v d_{M}(F x, F y) \leq C_{0} \tau(a) \operatorname{Lip} v d_{\theta}(F x, F y)^{1 / \eta} \\
& \leq C_{0} \tau(a) \operatorname{Lip} v d_{\theta}(F x, F y)=C_{0} \theta^{-1} \tau(a) \operatorname{Lip} v d_{\theta}(x, y)
\end{aligned}
$$

completing the estimates for $V$.
Next, $(P V)(y)=\sum_{a \in \alpha} g\left(y_{a}\right) V\left(y_{a}\right)$, so by (4.1),

$$
|P V|_{\infty} \leq C_{2} \sum_{a \in \alpha} \mu(a) \tau(a)|v|_{\infty}=C_{2}|\tau|_{1}|v|_{\infty}
$$

Also,

$$
\begin{aligned}
& |(P V)(x)-(P V)(y)| \leq \sum_{a \in \alpha}\left|g\left(x_{a}\right)-g\left(y_{a}\right)\right|\left|V\left(x_{a}\right)\right|+\sum_{a \in \alpha} g\left(y_{a}\right)\left|V\left(x_{a}\right)-V\left(y_{a}\right)\right| \\
& \quad \leq C_{2} \sum_{a \in \alpha} \mu(a) d_{\theta}(x, y) \tau(a)|v|_{\infty}+C_{2} \sum_{a \in \alpha} \mu(a) C_{0} \theta^{-1} \tau(a) \operatorname{Lip} v d_{\theta}(x, y),
\end{aligned}
$$

yielding the estimate for $|P V|_{\theta}$.
Proposition 4.8 Let $p \geq 1$. There exists $m \in L^{p}\left(Y, \mathbb{R}^{d}\right)$ and $\chi \in L^{\infty}\left(Y, \mathbb{R}^{d}\right)$ such that $V=m+\chi \circ F-\chi$, and $m \in \operatorname{ker} P$. Moreover,

$$
|m|_{p} \leq 3 C_{3}|\tau|_{p}\|v\|_{\text {Lip }}, \quad \text { and } \quad|\chi|_{\infty} \leq C_{3}|\tau|_{1}\|v\|_{\text {Lip }}
$$

where $C_{3}=2 C_{0} C_{2}^{2} \theta^{-1}(1-\gamma)^{-1}$.

Proof By Proposition 4.7 and Lemma 4.6(c) with $\phi=P V$, for $n \geq 1$,

$$
\left|P^{n} V\right|_{\infty} \leq C_{2} \gamma^{n-1}\|P V\|_{\theta} \leq 2 C_{0} C_{2}^{2} \theta^{-1} \gamma^{n-1}|\tau|_{1}\|v\|_{\text {Lip }} .
$$

It follows that $\chi=\sum_{k=1}^{\infty} P^{k} V$ lies in $L^{\infty}$ and $|\chi|_{\infty} \leq C_{3}|\tau|_{1}\|v\|_{\text {Lip }}$.
Write $V=m+\chi \circ F-\chi$; then $m \in L^{p}$ and $P m=0$. Finally, $|m|_{p} \leq|V|_{p}+\left.2\right|_{\chi} \leq$ $|\tau|_{p}|v|_{\infty}+2|\chi|_{\infty} \leq 3 C_{3}|\tau|_{p}\|v\|_{\text {Lip }}$.

Corollary 4.9 Define $V_{n}=\sum_{j=0}^{n-1} V \circ F^{j}$. Let $p>1$. There exists a constant $C_{4} \geq 1$ depending only on $p$ and $C_{3}$ such that

$$
\left|\max _{1 \leq j \leq n}\right| V_{j}| |_{p} \leq C_{4}|\tau|_{p}\|v\|_{\text {Lip }} n^{\max \{1 / 2,1 / p\}}
$$

Proof First note that $V_{n}=m_{n}+\chi \circ F^{n}-\chi$ where $m_{n}=\sum_{j=0}^{n-1} m \circ F$. Since $m \in \operatorname{ker} P$, an application of Burkholder's inequality [12] shows that

$$
\left|\max _{1 \leq j \leq n}\right| m_{j}| |_{p} \leq C(p)|m|_{p} n^{\max \{1 / 2,1 / p\}}
$$

(see for example the proof of [30, Proposition 4.3]). Hence

$$
\left|\max _{1 \leq j \leq n}\right| V_{j}| |_{p} \leq C(p)|m|_{p} n^{\max \{1 / 2,1 / p\}}+2|\chi|_{\infty}
$$

The result follows from Proposition 4.8 with $C_{4}=5 C_{3} C(p)$.
Lemma 4.10 Let $p>1$. Let $v_{n}=\sum_{j=0}^{n-1} v \circ T^{j}$. Then

$$
\left.\left|\max _{j \leq n}\right| v_{j}\right|_{L^{p-1}(\nu)} \leq 5 C_{4}|\tau|_{p}^{p /(p-1)}\|v\|_{\text {Lip }} n^{\max \{1 / 2,1 / p\}}
$$

Proof Let $q=p-1$. Define the tower $\Delta$ as in (4.2) with tower map $f: \Delta \rightarrow \Delta$ where

$$
f(y, \ell)=\left\{\begin{array}{ll}
(y, \ell+1), & \ell \leq \tau(y)-2 \\
(F y, 0), & \ell=\tau(y)-1
\end{array} .\right.
$$

Recall that $\mu_{\Delta}=\mu \times$ counting $/ \bar{\tau}$ on $\Delta$ where $\bar{\tau}=\int_{Y} \tau d \mu$. Also, $\nu=\left(\pi_{\Delta}\right)_{*} \mu_{\Delta}$ where $\pi_{\Delta}: \Delta \rightarrow M$ is the projection $\pi_{\Delta}(y, \ell)=T^{\ell} y$.

Let $\hat{v}=v \circ \pi_{\Delta}$ and define $\hat{v}_{n}=\sum_{j=0}^{n-1} \hat{v} \circ f^{j}$. Then $\int_{M}\left|v_{n}\right|^{q} d \nu=\int_{\Delta}\left|\hat{v}_{n}\right|^{q} d \mu_{\Delta}$.
Next, let $N_{n}: \Delta \rightarrow\{0,1, \ldots, n\}$ be the number of laps by time $n$,

$$
N_{n}(y, \ell)=\#\left\{j \in\{1, \ldots, n\}: f^{j}(y, \ell) \in Y \times\{0\}\right\}
$$

Then

$$
\hat{v}_{n}(y, \ell)=V_{N_{n}(y, \ell)}(y)+H \circ f^{n}(y, \ell)-H(y, \ell)
$$

where $H(y, \ell)=\hat{v}_{\ell}(y, 0)$. Note that $|H(y, \ell)| \leq|v|_{\infty} \tau(y)$ for all $(y, \ell) \in \Delta$.
Now $f^{n}(y, \ell)=\left(F^{N_{n}(y, \ell)} y, \ell+n-\tau_{N_{n}(y, \ell)}(y)\right)$, so

$$
\begin{aligned}
& \max _{j \leq n}\left|H \circ f^{j}(y, \ell)\right| \leq|v|_{\infty} \max _{j \leq n} \tau\left(F^{N_{j}(y, \ell)} y\right) \leq|v|_{\infty} \max _{j \leq n} \tau\left(F^{j} y\right) \\
& \quad \leq|v|_{\infty} \tau(y)+|v|_{\infty} \max _{1 \leq j \leq n} \tau\left(F^{j} y\right)=|v|_{\infty} \hat{\tau}(y, \ell)+|v|_{\infty} \max _{1 \leq j \leq n} \hat{\tau}\left(F^{j} y, \ell\right),
\end{aligned}
$$

where $\hat{\tau}: \Delta \rightarrow \mathbb{Z}^{+}$is given by $\hat{\tau}(y, \ell)=\tau(y)$.
We estimate the first term in $L^{q}\left(\mu_{\Delta}\right)$ and the second term in $L^{p}\left(\mu_{\Delta}\right)$. Using the definition of $\mu_{\Delta}$ and the fact that $\bar{\tau} \geq 1$,

$$
\int_{\Delta} \hat{\tau}^{q} d \mu_{\Delta}=(1 / \bar{\tau}) \int_{Y} \tau^{q+1} d \mu \leq \int_{Y} \tau^{p} d \mu
$$

so $|\hat{\tau}|_{L^{q}\left(\mu_{\Delta}\right)} \leq|\tau|_{p}^{p /(p-1)}$. Also,

$$
\begin{gathered}
\int_{\Delta} \max _{1 \leq j \leq n} \hat{\tau}\left(F^{j} y\right)^{p} d \mu_{\Delta}=(1 / \bar{\tau}) \int_{Y} \tau \max _{1 \leq j \leq n} \tau^{p} \circ F^{j} d \mu \leq(1 / \bar{\tau}) \sum_{j=1}^{n} \int_{Y} \tau \tau^{p} \circ F^{j} d \mu \\
=(1 / \bar{\tau}) \sum_{j=1}^{n} \int_{Y} P \tau \tau^{p} \circ F^{j-1} d \mu \leq(1 / \bar{\tau}) \sum_{j=1}^{n}|P \tau|_{\infty} \int_{Y} \tau^{p} \circ F^{j-1} d \mu \\
=(1 / \bar{\tau}) n|P \tau|_{\infty} \int_{Y} \tau^{p} d \mu \leq C_{2} n \int_{Y} \tau^{p} d \mu
\end{gathered}
$$

where we used Proposition 4.7 with $v=1$ (and hence $V=\tau$ ) for the final inequality. Hence

$$
\left|\max _{1 \leq j \leq n} \hat{\tau}\left(F^{j} y\right)\right|_{L^{q}\left(\mu_{\Delta}\right)} \leq\left|\max _{1 \leq j \leq n} \hat{\tau}\left(F^{j} y\right)\right|_{L^{p}\left(\mu_{\Delta}\right)} \leq C_{2}^{1 / p} n^{1 / p}|\tau|_{p}
$$

Combining these estimates, we obtain that

$$
|H|_{L^{q}\left(\mu_{\Delta}\right)} \leq\left.\left|\max _{j \leq n}\right| H \circ f^{j}\right|_{L^{q}\left(\mu_{\Delta}\right)} \leq 2 C_{2}^{1 / p}|v|_{\infty}|\tau|_{p}^{p /(p-1)} n^{1 / p}
$$

Next, using Hölder's inequality,

$$
\begin{aligned}
\int_{\Delta} \max _{j \leq n}\left|V_{N_{j}(y, \ell)}(y)\right|^{q} d \mu_{\Delta} & \leq \int_{\Delta} \max _{j \leq n}\left|V_{j}(y)\right|^{q} d \mu_{\Delta}=(1 / \bar{\tau}) \int_{Y} \tau \max _{j \leq n}\left|V_{j}\right|^{q} d \mu \\
& \leq\left.\left.|\tau|_{p}\left|\max _{j \leq n}\right| V_{j}\right|^{q}\right|_{p / q}=\left.|\tau|_{p}\left|\max _{j \leq n}\right| V_{j}\right|_{p} ^{q}
\end{aligned}
$$

By Corollary 4.9,

$$
\left|\max _{j \leq n}\right| V_{N_{j}}| |_{L^{q}\left(\mu_{\Delta}\right)} \leq C_{4}|\tau|_{p}^{p /(p-1)}\|v\|_{\text {Lip }} n^{\max \{1 / 2,1 / p\}}
$$

By the triangle inequality, using that $C_{2}^{1 / p} \leq C_{4}$,

$$
\left|\max _{j \leq n}\right| v_{n}| |_{L^{p-1}(\nu)}=\left.\left|\max _{j \leq n}\right| \hat{v}_{n}\right|_{L^{q}\left(\mu_{\Delta}\right)} \leq 5 C_{4}|\tau|_{p}^{p /(p-1)}\|v\|_{\text {Lip }} n^{\max \{1 / 2,1 / p\}}
$$

as required.

### 4.5 Proof of Theorem 4.3

Define $v_{\epsilon, x}$ and $v_{\epsilon, x, n}$ as in Section 22. Note that $\operatorname{Lip} v_{\epsilon, x} \leq 2 L$ for all $\epsilon, x$ and $\int_{M} v_{\epsilon, x} d \nu_{\epsilon}=0$.

It follows from Lemma 4.10 that

$$
\left.\left|\max _{j \leq n}\right| v_{\epsilon, x, j}\right|_{L^{p-1}\left(\nu_{\epsilon}\right)} \leq 10 C_{4} L\left|\tau_{\epsilon}\right|_{p}^{p /(p-1)} n^{\max \{1 / 2,1 / p\}}
$$

for all $\epsilon \geq 0, x \in \mathbb{R}, n \geq 1$. By (2.7),

$$
\begin{aligned}
\int_{M} \delta_{1, \epsilon}^{p+d-1} d \nu_{\epsilon} & \leq \Gamma \epsilon^{p-1} \sup _{x \in E} \int_{M} \max _{n \leq 1 / \epsilon}\left|v_{\epsilon, x, n}\right|^{p-1} d \nu_{\epsilon} \leq \Gamma C_{4}^{p-1}\left|\tau_{\epsilon}\right|_{p}^{p} \epsilon^{p-1} \epsilon^{-(p-1) \max \{1 / 2,1 / p\}} \\
& =\Gamma C_{4}^{p-1}\left|\tau_{\epsilon}\right|_{p}^{p} \epsilon^{p-1} \epsilon^{\min \{-(p-1) / 2,-(p-1) / p\}}=\Gamma C_{4}^{p-1}\left|\tau_{\epsilon}\right|_{p}^{p} \epsilon^{\min \left\{(p-1) / 2,(p-1)^{2} / p\right\}}
\end{aligned}
$$

We obtain the same estimate for $\delta_{2, \epsilon}$ replacing $v_{\epsilon, x}, v_{\epsilon, x, n}$ by $\mathbf{V}_{\epsilon, x}, \mathbf{V}_{\epsilon, x, n}$.

## 5 Examples: Nonuniformly expanding maps

Example 5.1 (Intermittent maps) Let $M=[0,1]$ with Lebesgue measure $m$ and consider the intermittent maps $T: M \rightarrow M$ given by

$$
T x= \begin{cases}x\left(1+2^{a} x^{a}\right), & x \in\left[0, \frac{1}{2}\right]  \tag{5.1}\\ 2 x-1, & x \in\left(\frac{1}{2}, 1\right]\end{cases}
$$

These were studied in [28]. Here $a>0$ is a parameter. For $a \in(0,1)$ there is a unique absolutely continuous invariant probability measure with $C^{\infty}$ nonvanishing density.

We consider a family $T_{\epsilon}: M \rightarrow M, \epsilon \in\left[0, \epsilon_{0}\right)$, of such intermittent maps with parameter $a_{\epsilon} \in(0,1)$ depending continuously on $\epsilon$. Let $\nu_{\epsilon}$ denote the corresponding family of absolutely continuous invariant probability measures.

For each $\epsilon$, we take $Y=\left[\frac{1}{2}, 1\right]$ and let $\tau_{\epsilon}: Y \rightarrow \mathbb{Z}^{+}$be the first return time $\tau_{\epsilon}(y)=\inf \left\{n \geq 1: T_{\epsilon}^{n} y \in Y\right\}$. Define the first return map $F_{\epsilon}=T_{\epsilon}^{\tau_{\epsilon}}: Y \rightarrow Y$. Let $\alpha_{\epsilon}=\left\{Y_{\epsilon}(n), n \geq 1\right\}$ where $Y_{\epsilon}(n)=\left\{y \in Y: \tau_{\epsilon}(y)=n\right\}$.

It is standard that each $T_{\epsilon}$ is a nonuniformly expanding map in the sense of Section 4.1 with $\tau_{\epsilon} \in L^{p}$ for all $p<1 / a_{\epsilon}$. Fix $p \in\left(1,1 / a_{0}\right)$ and choose $0<a_{-}<a_{0}<$ $a_{+}<1$ such that $p<1 / a_{+}$. Without loss we can shrink $\epsilon_{0}$ so that $a_{\epsilon} \in\left[a_{-}, a_{+}\right]$for all $\epsilon \in\left[0, \epsilon_{0}\right)$. We show that $T_{\epsilon}, \epsilon \in\left[0, \epsilon_{0}\right)$, satisfies Definition 4.2 for this choice of $p$.

Since $T_{\epsilon}^{\prime} \geq 1$ on $M$ and $T_{\epsilon}^{\prime}=2$ on $Y$, it is immediate that conditions (1)-(3) in Section 4.1 are satisfied with $\lambda=2$ and $C_{0}=1$.

Recall that $\zeta_{\epsilon}=\frac{\left.d m\right|_{Y}}{\left.d m\right|_{Y} \circ F_{\epsilon}}$. Note that $\zeta_{\epsilon}(y)=1 / F_{\epsilon}^{\prime}(y)$. For each $Y_{\epsilon}(n) \in \alpha_{\epsilon}$, define the bijection $F_{\epsilon, n}=\left.F_{\epsilon}\right|_{Y_{\epsilon}(n)}: Y_{\epsilon}(n) \rightarrow M$. Let $G_{\epsilon, n}=\left(F_{\epsilon, n}^{-1}\right)^{\prime}=\zeta_{\epsilon} \circ F_{\epsilon, n}^{-1}$. By [25, Assumption A2 and Theorem 3.1], there is a constant $K$, depending only on $a_{-}$and
$a_{+}$, such that $\left|\left(\log G_{\epsilon, n}\right)^{\prime}\right| \leq K$ for all $\epsilon \in\left[0, \epsilon_{0}\right)$. Hence $\left|\left(\log \zeta_{\epsilon} \circ F_{\epsilon, n}^{-1}\right)^{\prime}\right|=\left|\left(\log G_{\epsilon, n}\right)^{\prime}\right| \leq$ $K$. By the mean value theorem, for $x, y \in Y_{\epsilon}(n)$,

$$
\left|\log \zeta_{\epsilon}(x)-\log \zeta_{\epsilon}(y)\right|=\left|\left(\log \zeta_{\epsilon} \circ F_{\epsilon, n}^{-1}\right)\left(F_{\epsilon} x\right)-\left(\log \zeta_{\epsilon} \circ F_{\epsilon, n}^{-1}\right)\left(F_{\epsilon} y\right)\right| \leq K\left|F_{\epsilon} x-F_{\epsilon} y\right| .
$$

This proves condition (4) in Section 4.1, and so condition (i) in Definition 4.2 is satisfied.

Define $x_{1}=\frac{1}{2}$ and inductively $x_{n+1}<x_{n}$ (depending on $\epsilon$ ) with $T_{\epsilon} x_{n+1}=x_{n}$. Then $T_{\epsilon}\left(Y_{\epsilon}(n)\right)=\left[x_{n}, x_{n-1}\right]$ for $n \geq 2$ and it is standard that $x_{n}=O\left(1 / n^{\alpha}\right)$ as a function of $n$. By [25, Lemma 5.2], there is a constant $K$, depending only on $a_{-}$and $a_{+}$, such that $x_{n} \leq K n^{-1 / \alpha_{+}}$for all $n \geq 1, \epsilon \in\left[0, \epsilon_{0}\right)$. Hence

$$
m\left(\tau_{\epsilon}>n\right)=m\left(\left[1 / 2,\left(x_{n}+1\right) / 2\right]\right)=x_{n} / 2 \leq K n^{-1 / \alpha_{+}}
$$

Since $p<1 / \alpha_{+}$it follows that $\sup _{\epsilon \in\left[0, \epsilon_{0}\right)} \int_{Y}\left|\tau_{\epsilon}\right|^{p} d m<\infty$ so condition (ii) in Definition 4.2 is satisfied.

By [11, 25], $\nu_{\epsilon}$ is strongly statistically stable and the densities $\rho_{\epsilon}$ satisfy $R_{\epsilon}=$ $\int\left|\rho_{\epsilon}-\rho_{0}\right| d m=O\left(a_{\epsilon}-a_{0}\right)$. Hence, by Corollary 2.8, we obtain averaging in $L^{1}$ with respect to $\nu_{\epsilon}$ and also with respect to any absolutely continuous probability measure.

Finally, if $\epsilon \mapsto a_{\epsilon}$ is Lipschitz say, so that $R_{\epsilon}=O(\epsilon)$, then we obtain the rates described in Remark 4.4 with $p=\left(1 / a_{0}\right)-$.

Example 5.2 (Logistic family) We consider the family of quadratic maps $T$ : $[-1,1] \rightarrow[-1,1]$ given by $T(x)=1-a x^{2}, a \in[0,2]$, with $m$ taken to be Lebesgue measure.

Let $b, c>0$. The map $T$ satisfies the Collet-Eckmann condition [13] with constants $b, c$ if

$$
\begin{equation*}
\left|\left(T^{n}\right)^{\prime}(1)\right| \geq c e^{b n} \quad \text { for all } n \geq 0 \tag{5.2}
\end{equation*}
$$

In this case, we write $a \in Q_{b, c}$. The set of Collet-Eckmann parameters is given by $P_{1}=\bigcup_{b, c>0} Q_{b, c}$ and is a Cantor set of positive Lebesgue measure [19, 10]. When $a \in P_{1}$, the map $T$ has an invariant set $\Lambda$ consisting of a finite union of intervals with an ergodic absolutely continuous invariant probability measure $\nu_{a}$. The density for $\nu_{a}$ is bounded below on $\Lambda$ and lies in $L^{2-}$. The invariant set attracts Lebesgue almost every trajectory in $[-1,1]$.

There is also an open dense set of parameters $P_{0} \subset[0,2]$ for which $T$ possesses a periodic sink attracting Lebesgue almost every trajectory in [ $-1,1$ ]. By Lyubich [29], $P_{0} \cup P_{1}$ has full measure. For $a \in P_{0}$, we let $\nu_{a}$ denote the invariant probability measure supported on the periodic attractor, so we have a map $a \mapsto \nu_{a}$ defined on $P_{0} \cup P_{1}$.

It is clear that statistical stability holds on $P_{0}$, and that strong statistical stability fails everywhere in $P_{0} \cup P_{1}$. Moreover, Thunberg [33, Corollary 1] shows that on any full measure subset of $E \subset[0,2]$ the map $a \rightarrow \nu_{a}$ is not statistically stable at any
point of $P_{1} \cap E$. On the other hand, Freitas \& Todd [17] proved that strong statistical stability holds on $Q_{b, c}$ for all constants $b, c>0$. That is, the map $a \rightarrow \rho_{a}=d \nu_{a} / d m$ from $Q_{b, c} \rightarrow L^{1}$ is continuous. (See also [15, 16] for the same result restricted to the Benedicks-Carleson parameters [10].)

We consider families $\epsilon \rightarrow T_{\epsilon}$ where each $T_{\epsilon}$ is a quadratic map with parameter $a=a_{\epsilon}$ depending continuously on $\epsilon$. Fix $b, c>0$ such that $a_{0} \in Q_{b, c}$. We claim that

$$
\lim _{\substack{\epsilon \rightarrow 0 \\ a_{\epsilon} \in Q_{b, c}}} \int z_{\epsilon} d \nu_{a_{\epsilon}}=0
$$

Moreover, using Corollary 2.9 we obtain convergence in $L^{1}(\mu)$ for every absolutely continuous probability measure $\mu$. Given the above results on strong statistical stability, it suffices to verify that $T_{\epsilon}$ is a uniform family of nonuniformly expanding maps.

For the Benedicks-Carleson parameters, the method in [15, 16] is the approach of [2] and we can apply Remark 4.5. In the general case, a different method exploiting negative Schwarzian derivative and Koebe spaces [17, Proof of Theorem B in Section 6] shows that the conditions in [5] are satisfied. By Remark 4.5, this completes the proof of averaging with the possible exception of condition (3). However, a standard consequence of negative Schwarzian derivative and the Koebe distortion property (as discussed in [17, Lemma 4.1] and used in [17, Remark 3.2]) is that bounded distortion holds at intermediate steps and not just at the inducing time as in condition (4). Hence there is a uniform constant $\tilde{C}_{1}$ such that $\frac{\left|T_{\epsilon}^{j} x-T_{\epsilon}^{j} y\right|}{\operatorname{diam} T_{\epsilon}^{j} a} \leq \tilde{C}_{1} \frac{\left|F_{\epsilon} x-F_{\epsilon} y\right|}{\operatorname{diam} Y_{\epsilon}}$ for all partition elements $a$, all $x, y \in a$ and all $j \leq \tau_{\epsilon}(a)$. In particular, $\left|T_{\epsilon}^{j} x-T_{\epsilon}^{j} y\right| \leq\left(2 \tilde{C}_{1} / \operatorname{diam} Y_{\epsilon}\right)\left|F_{\epsilon} x-F_{\epsilon} y\right|$, yielding condition (3) uniformly in $\epsilon$.

Next, we discuss rates of convergence. By [17, Lemma 4.1], condition (ii) in Definition 4.2 is satisfied for any $p>1$. Hence $\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}=O\left(\epsilon^{\frac{1}{2}-}\right)$. If $\epsilon \mapsto a_{\epsilon}$ is $C^{1}$, then it follows from Baladi et al. 9] that $R_{\epsilon}=\int\left|\rho_{\epsilon}-\rho_{0}\right| d m=O\left(\epsilon^{\frac{1}{2}-}\right)$. By Remark 4.4, we obtain averaging with rate $O\left(\epsilon^{\frac{1}{2}-}\right)$ in $L^{q}\left(\nu_{\epsilon}\right)$ for all $q>0$ and in $L^{1}\left(\nu_{0}\right)$.

Example 5.3 (Multimodal maps) Freitas \& Todd [17] also consider families of multimodal maps where each critical point $c$ satisfies a Collet-Eckmann condition along the orbit of $T c$ with constants uniform in $\epsilon$. Hence the averaging result for the quadratic family in Example 5.2 extends immediately to multimodal maps.

Example 5.4 (Viana maps) Viana [34] introduced a $C^{3}$ open class of multidimensional nonuniformly expanding maps $T_{\epsilon}: M \rightarrow M$. For definiteness, we restrict attention to the case $M=S^{1} \times \mathbb{R}$. Let $S: M \rightarrow M$ be the map $S(\theta, y)=\left(16 \theta \bmod 1, a_{0}+a \sin 2 \pi \theta-y^{2}\right)$. Here $a_{0}$ is chosen so that 0 is a preperiodic point for the quadratic map $y \mapsto a_{0}-y^{2}$ and $a$ is fixed sufficiently small. Let
$T_{\epsilon}, 0 \leq \epsilon<\epsilon_{0}$ be a continuous family of $C^{3}$ maps sufficiently close to $S$. It follows from [1, 5] that there is an interval $I \subset(-2,2)$ such that, for each $\epsilon \in\left[0, \epsilon_{0}\right)$, there is a unique absolutely continuous $T_{\epsilon}$-invariant ergodic probability measure $\nu_{\epsilon}$ supported in the interior of $S^{1} \times I$. Moreover the invariant set $\Lambda_{\epsilon}=\operatorname{supp} \nu_{\epsilon}$ attracts almost every initial condition in $S^{1} \times I$.

By Alves \& Viana [5], $\nu_{0}$ is strongly statistically stable. Moreover, the inducing method of [4] and the arguments in [2] apply to this example, so $T_{\epsilon}$ is a uniform family of nonuniformly expanding maps by Remark 4.5. Also, Corollary 2.9 is applicable by Remark 2.10. Hence we obtain averaging in $L^{1}\left(\nu_{\epsilon}\right)$ and in $L^{1}(\mu)$ for all absolutely continuous $\mu$.

Finally we discuss rates. Although not stated explicitly in [5] , it follows from their estimates that $T_{\epsilon}$ is a uniform family of order $p$ for all $p$. Hence $\left|\delta_{\epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)}=O\left(\epsilon^{1 / 2-}\right)$ for all $q>0$ by Remark 4.4. To verify that $p$ is arbitrary, we mention the following steps in [5] (where their $q$ is our $p$ ). Note by Remark 4.5 that it remains to verify their condition (U2) for all $p$. On page 25 (calculation for (U2)), they give an estimate that works for all $p$ depending on two constants $C_{0}$ and $\gamma_{0}$ which are uniform in $\epsilon$ by [5, Remark 4.6].

## 6 Averaging for continuous time fast-slow systems

Let $\phi_{t}^{\epsilon}: M \rightarrow M, 0 \leq \epsilon<\epsilon_{0}$, be a family of semiflows defined on the metric space $\left(M, d_{M}\right)$. For each $\epsilon \geq 0$, let $\nu_{\epsilon}$ denote a $\phi_{t}^{\epsilon}$-invariant ergodic Borel probability measure. Let $a: \mathbb{R}^{d} \times M \times\left[0, \epsilon_{0}\right) \rightarrow \mathbb{R}^{d}$ be a family of vector fields on $\mathbb{R}^{d}$ satisfying conditions (2.3)- 2.5).

We consider the family of fast-slow systems

$$
\begin{aligned}
& \dot{x}^{(\epsilon)}=\epsilon a\left(x^{(\epsilon)}, y^{(\epsilon)}, \epsilon\right), \quad x^{(\epsilon)}(0)=x_{0}, \\
& y^{(\epsilon)}(t)=\phi_{t}^{\epsilon} y_{0},
\end{aligned}
$$

where the initial condition $x^{(\epsilon)}(0)=x_{0}$ is fixed throughout. The initial condition $y_{0} \in M$ is again chosen randomly with respect to various measures that are specified in the statements of the results.

Define $\hat{x}^{(\epsilon)}:[0,1] \rightarrow \mathbb{R}^{d}$ by setting $\hat{x}^{(\epsilon)}(t)=x^{(\epsilon)}(t / \epsilon)$. Let $X:[0,1] \rightarrow \mathbb{R}^{d}$ be the solution to the ODE $(2.2)$ and define

$$
z_{\epsilon}=\sup _{t \in[0,1]}\left|\hat{x}^{(\epsilon)}(t)-X(t)\right| .
$$

Recall that $E=\left\{x \in \mathbb{R}^{d}:\left|x-x_{0}\right| \leq L_{1}\right\}$. As in Section 2.1, define $\bar{a}(x, \epsilon)=$ $\int_{M} a(x, y, \epsilon) d \nu_{\epsilon}(y)$ and let $v_{\epsilon, x}(y)=a(x, y, \epsilon)-\bar{a}(x, \epsilon)$. We define the order function
$\delta_{\epsilon}=\delta_{1, \epsilon}+\delta_{2, \epsilon}: M \rightarrow \mathbb{R}$ where

$$
\begin{array}{lll}
\delta_{1, \epsilon}=\sup _{x \in E} \sup _{0 \leq t \leq 1 / \epsilon} \epsilon\left|v_{\epsilon, x, t}\right| & \text { where } & v_{\epsilon, x, t}=\int_{0}^{t} v_{\epsilon, x} \circ \phi_{s}^{\epsilon} d s \\
\delta_{2, \epsilon}=\sup _{x \in E} \sup _{0 \leq t \leq 1 / \epsilon} \epsilon\left|\mathbf{V}_{\epsilon, x, t}\right| & \text { where } & \mathbf{V}_{\epsilon, x, t}=\int_{0}^{t}\left(D v_{\epsilon, x}\right) \circ \phi_{s}^{\epsilon} d s .
\end{array}
$$

The next results is the continuous time analogue of Theorem 2.2 . The proof is entirely analogous, and hence is omitted.

Theorem 6.1 Let $S_{\epsilon}=\sup _{x \in E}\left|\int_{M} a(x, y, 0)\left(d \nu_{\epsilon}-d \nu_{0}\right)(y)\right|+\epsilon$. Assume conditions (2.3) -2.5. If $\delta_{\epsilon} \leq \frac{1}{2}$, then $z_{\epsilon} \leq 6 e^{2 L}\left(\delta_{\epsilon}+S_{\epsilon}\right)$.

As in the discrete time setting, we say that $\nu_{0}$ is statistically stable if $\nu_{\epsilon} \rightarrow_{w} \nu_{0}$. Proposition 2.6 goes through unchanged and statistical stability implies that $S_{\epsilon} \rightarrow 0$.

If the measures $\nu_{\epsilon}$ are absolutely continuous with respect to a reference measure $m$ on $M$, we define the densities $\rho_{\epsilon}=d \nu_{\epsilon} / d m$ and set $R_{\epsilon}=\int_{M}\left|\rho_{\epsilon}-\rho_{0}\right| d m$. Then $\nu_{0}$ is strongly statistically stable if $R_{\epsilon} \rightarrow 0$. Proposition 2.7 and Corollaries 2.8 and 2.9 go through unchanged from the discrete time setting.

Fix a Borel subset $M^{\prime} \subset M$ and a reference Borel measure $m^{\prime}$ on $M^{\prime}$. Let $h_{\epsilon}$ : $M^{\prime} \rightarrow \mathbb{R}^{+}$be a family of Lipschitz functions such that $\phi_{h_{\epsilon}(y)}^{\epsilon}(y) \in M^{\prime}$ for almost all $y \in M^{\prime}$. The map $T_{\epsilon}: M^{\prime} \rightarrow M^{\prime}, T_{\epsilon}(y)=\phi_{h_{\epsilon}(y)}^{\epsilon}(y)$, is then defined almost everywhere.

As usual, we suppose that there is a family $\nu_{\epsilon}^{\prime}$ of ergodic $T_{\epsilon}$-invariant probability measures on $M^{\prime}$. Define the suspension $M_{h_{\epsilon}}=\left\{(y, u) \in M^{\prime} \times \mathbb{R}: 0 \leq u \leq h_{\epsilon}(y)\right\} / \sim$ where $\left(y, h_{\epsilon}(y)\right) \sim\left(T_{\epsilon} y, 0\right)$. The suspension semiflow $f_{t}^{\epsilon}: M_{h_{\epsilon}} \rightarrow M_{h_{\epsilon}}$ is given by $f_{t}^{\epsilon}(y, u)=(y, u+t)$ computed modulo identifications. Let $\bar{h}_{\epsilon}=\int_{M^{\prime}} h_{\epsilon} d \nu_{\epsilon}^{\prime}$. Then $\nu_{\epsilon}^{\prime \prime}=\left(\nu_{\epsilon}^{\prime} \times\right.$ Lebesgue $) / \bar{h}_{\epsilon}$ is an ergodic absolutely continuous $f_{t}^{\epsilon}$-invariant probability measure on $M_{h_{\epsilon}}$. The projection $\pi_{\epsilon}: M_{h_{\epsilon}} \rightarrow M$ given by $\pi_{\epsilon}(y, u)=\phi_{u}^{\epsilon} y$ is a semiconjugacy between $f_{t}^{\epsilon}$ and $\phi_{t}^{\epsilon}$. Hence $\nu_{\epsilon}=\pi_{\epsilon *} \nu_{\epsilon}^{\prime \prime}$ is an ergodic $\phi_{t}^{\epsilon}$-invariant probability measure on $M$.

We suppose from now on that there are constants $K_{2} \geq K_{1} \geq 1$ such that for all $x, y \in M^{\prime}, \epsilon \in\left[0, \epsilon_{0}\right)$,

- $K_{1}^{-1} \leq h_{\epsilon} \leq K_{1}, \operatorname{Lip} h_{\epsilon} \leq K_{1},\left|h_{\epsilon}-h_{0}\right|_{\infty} \leq K_{1} \epsilon$.
- $d_{M}\left(\phi_{t}^{\epsilon} x, \phi_{t}^{\epsilon} y\right) \leq K_{2} d_{M}(x, y)$ and $d_{M}\left(\phi_{t}^{\epsilon} y, \phi_{t}^{0} y\right) \leq K_{2} \epsilon$ for all $t \leq K_{1}$.
(These assumptions are easily weakened; in particular changing the $\epsilon$ estimates to $\epsilon^{1 / 2}$ will not affect anything.)

Proposition 6.2 Let $v: M \rightarrow \mathbb{R}^{d}$ be Lipschitz. Define $\tilde{v}: M^{\prime} \rightarrow \mathbb{R}^{d}, \tilde{v}(y)=$ $\int_{0}^{h_{0}(y)} v\left(\phi_{u}^{0} y\right) d u$. Then

$$
\int_{M} v\left(d \nu_{\epsilon}-d \nu_{0}\right) \leq 3 K_{2}^{4}\|v\|_{\text {Lip }} \epsilon+K_{1}^{3}|v|_{\infty}\left|\int_{M^{\prime}} h_{0}\left(d \nu_{\epsilon}^{\prime}-d \nu_{0}^{\prime}\right)\right|+K_{1}\left|\int_{M^{\prime}} \tilde{v}\left(d \nu_{\epsilon}^{\prime}-d \nu_{0}^{\prime}\right)\right| .
$$

Proof We have

$$
\begin{aligned}
\int_{M} v\left(d \nu_{\epsilon}-d \nu_{0}\right) & =\int_{M_{h_{\epsilon}}} v \circ \pi_{\epsilon} d \nu_{\epsilon}^{\prime \prime}-\int_{M_{h_{0}}} v \circ \pi_{0} d \nu_{0}^{\prime \prime} \\
& =\left(1 / \bar{h}_{\epsilon}\right) \int_{M^{\prime}} \int_{0}^{h_{\epsilon}} v \circ \pi_{\epsilon} d u d \nu_{\epsilon}^{\prime}-\left(1 / \bar{h}_{0}\right) \int_{M^{\prime}} \int_{0}^{h_{0}} v \circ \pi_{0} d u d \nu_{0}^{\prime} \\
& =I_{1}+I_{2}+I_{3}+I_{4}
\end{aligned}
$$

where

$$
\begin{aligned}
& I_{1}=\left(1 / \bar{h}_{\epsilon}-1 / \bar{h}_{0}\right) \int_{M^{\prime}} \int_{0}^{h_{\epsilon}} v \circ \pi_{\epsilon} d u d \nu_{\epsilon}^{\prime}, I_{2}=\left(1 / \bar{h}_{0}\right) \int_{M^{\prime}} \int_{0}^{h_{\epsilon}}\left(v \circ \pi_{\epsilon}-v \circ \pi_{0}\right) d u d \nu_{\epsilon}^{\prime}, \\
& I_{3}=\left(1 / \bar{h}_{0}\right)\left(\int_{M^{\prime}} \int_{0}^{h_{\epsilon}} v \circ \pi_{0} d u d \nu_{\epsilon}^{\prime}-\int_{M^{\prime}} \int_{0}^{h_{0}} v \circ \pi_{0} d u d \nu_{\epsilon}^{\prime}\right), \\
& I_{4}=\left(1 / \bar{h}_{0}\right)\left(\int_{M^{\prime}} \int_{0}^{h_{0}} v \circ \pi_{0} d u d \nu_{\epsilon}^{\prime}-\int_{M^{\prime}} \int_{0}^{h_{0}} v \circ \pi_{0} d u d \nu_{0}^{\prime}\right) .
\end{aligned}
$$

Now

$$
\begin{aligned}
& \left|I_{1}\right| \leq K_{1}^{2}\left|\bar{h}_{\epsilon}-\bar{h}_{0}\right| K_{1}|v|_{\infty} \leq K_{1}^{3}|v|_{\infty}\left(K_{1} \epsilon+\left|\int_{M^{\prime}} h_{0}\left(d \nu_{\epsilon}^{\prime}-d \nu_{0}^{\prime}\right)\right|\right) \\
& \left|I_{2}\right| \leq K_{1}^{2} \sup _{y \in M^{\prime}} \sup _{0 \leq u \leq K_{1}} \operatorname{Lip} v d_{M}\left(\phi_{u}^{\epsilon} y, \phi_{u}^{0} y\right) \leq K_{1}^{2} K_{2} \operatorname{Lip} v \epsilon \\
& \left|I_{3}\right| \leq K_{1}|v|_{\infty}\left|h_{\epsilon}-h_{0}\right|_{\infty} \leq K_{1}^{2}|v|_{\infty} \epsilon, \quad\left|I_{4}\right| \leq K_{1}\left|\int_{M^{\prime}} \tilde{v}\left(d \nu_{\epsilon}^{\prime}-d \nu_{0}^{\prime}\right)\right| .
\end{aligned}
$$

The result follows from the combination of these estimates.
Corollary 6.3 Statistical stability of $\nu_{0}^{\prime}$ implies statistical stability of $\nu_{0}$.
Proof This follows from Proposition 6.2.

Corollary 6.4 Suppose that the measures $\nu_{\epsilon}^{\prime}$ are absolutely continuous with respect to $m^{\prime}$, with densities $\rho_{\epsilon}^{\prime}=d \nu_{\epsilon}^{\prime} / d m^{\prime}$. Then $S_{\epsilon} \leq 3 K_{2}^{4} L\left(\epsilon+\int_{M^{\prime}}\left|\rho_{\epsilon}^{\prime}-\rho_{0}^{\prime}\right| d m^{\prime}\right)$.

Proof This follows from Proposition 6.2 with $v(y)=a(x, y, 0)$ for each fixed $x$.
Next we show how the order function for the flows $\phi_{t}^{\epsilon}: M \rightarrow M$ is related to the order function for the maps $T_{\epsilon}: M^{\prime} \rightarrow M^{\prime}$. We restrict attention to $\delta_{1, \epsilon}$ since the corresponding statement for $\delta_{2, \epsilon}$ is identical.

Define the family of induced observables $w_{x, \epsilon}: M^{\prime} \rightarrow \mathbb{R}$,

$$
w_{\epsilon, x}(y)=\int_{0}^{h_{\epsilon}(y)} v_{\epsilon, x}\left(\phi_{u}^{\epsilon} y\right) d u
$$

Note that $\int_{M^{\prime}} w_{\epsilon, x} d \nu_{\epsilon}^{\prime}=0$ and $w_{\epsilon, x}$ is $d_{M^{\prime}}$-Lipschitz with $\left\|w_{\epsilon}\right\|_{\text {Lip }} \leq 2 K_{1} K_{2} L$. Let

$$
\Delta_{1, \epsilon}=\sup _{x \in E} \sup _{1 \leq n \leq 1+K_{1} / \epsilon} \epsilon\left|w_{\epsilon, x, n}\right| \quad \text { where } \quad w_{\epsilon, x, n}=\sum_{j=0}^{n-1} w_{\epsilon, x} \circ T_{\epsilon}^{j} .
$$

We can now state our main result for this section.
Lemma 6.5 Let $q \geq 1$. Then $\left|\delta_{1, \epsilon}\right|_{L^{q}\left(\nu_{\epsilon}\right)} \leq\left(\left|\Delta_{1, \epsilon}\right|_{L^{q}\left(\nu_{\epsilon}^{\prime}\right)}+4 K_{1} L \epsilon\right)$.
Proof Let $\hat{v}_{\epsilon, x}=v_{\epsilon, x} \circ \pi_{\epsilon}$ and define $\hat{v}_{\epsilon, x, t}=\int_{0}^{t} \hat{v}_{\epsilon, x} \circ f_{u}^{\epsilon} d u$. Let $N_{\epsilon, t}: M_{h_{\epsilon}} \rightarrow$ $\left\{0,1, \ldots, 1+\left[K_{1} t\right]\right\}$ be the number of laps by time $t$,

$$
N_{\epsilon, t}(y, u)=\#\left\{s \in(0, t]: f_{s}^{\epsilon}(y, u) \in M^{\prime} \times\{0\}\right\}
$$

Then

$$
\hat{v}_{\epsilon, x, t}(y, u)=w_{\epsilon, x, N_{\epsilon, t}(y, u)}(y)+H_{\epsilon, x} \circ f_{t}^{\epsilon}(y, u)-H_{\epsilon, x}(y, u),
$$

where $H_{\epsilon, x}(y, u)=\int_{0}^{u} \hat{v}_{\epsilon, x}\left(y, u^{\prime}\right) d u^{\prime}=\hat{v}_{\epsilon, x, u}(y, 0)$. Note that $\left|H_{\epsilon, x}\right|_{\infty} \leq 2 K_{1} L$. Hence

$$
\begin{aligned}
\sup _{s \leq t}\left|v_{\epsilon, x, s}\right| \circ \pi_{\epsilon}(y, u) & =\sup _{s \leq t}\left|\hat{v}_{\epsilon, x, s}(y, u)\right| \leq \sup _{s \leq t}\left|w_{\epsilon, x, N_{s}(y, u)}(y)\right|+4 K_{1} L \\
& \leq \max _{j \leq 1+K_{1} t}\left|w_{\epsilon, x, j}(y)\right|+4 K_{1} L .
\end{aligned}
$$

It follows that

$$
\epsilon \sup _{s \leq 1 / \epsilon}\left|v_{\epsilon, x, s}\right| \circ \pi_{\epsilon}(y, u) \leq \Delta_{1, \epsilon}(y)+4 K_{1} L \epsilon
$$

and so $\delta_{1, \epsilon} \circ \pi_{\epsilon}(y, u) \leq \Delta_{1, \epsilon}(y)+4 K_{1} L \epsilon$. The result follows.
As a consequence of Proposition 6.4 and Lemma 6.5, our results for maps go through immediately for semiflows. For example, suppose that the maps $T_{\epsilon}(y)=$ $\phi_{h_{\epsilon}(y)}^{\epsilon}(y)$ are a family of quadratic maps as in Example 5.2. Then for any $q>0$, we obtain averaging in $L^{q}\left(\nu_{\epsilon}\right)$ with rate $O\left(\epsilon^{\frac{1}{2}-}\right)$. If moreover, $\nu_{0}$ is strongly statistically stable, then we obtain averaging in $L^{1}\left(\nu_{0}\right)$ with rate $O\left(\epsilon^{\frac{1}{2}-}\right)$, and averaging in $L^{1}(\mu)$ for any absolutely continuous probability measure $\mu$ on $M$.

## 7 Counterexample for almost sure convergence

It is known [7] that almost sure convergence fails for fully-coupled fast-slow systems. Here we give an example to show that almost sure convergence fails also in the simpler context of families of skew products as considered in this paper.

We consider the family of maps $T_{\epsilon}:[0,1] \rightarrow[0,1]$ given by $T_{\epsilon} y=2 y+\epsilon \bmod 1$ with invariant measure $\nu_{\epsilon}$ taken to be Lebesgue for all $\epsilon \geq 0$. Let $a(x, y, \epsilon)=\cos 2 \pi y$. Since $a$ has mean zero, the averaged ODE is given by $\dot{X}=0$. We take $x_{0}=0$ so that $X(t) \equiv 0$. Nevertheless, we prove:

Proposition 7.1 For every $y_{0} \in[0,1]$, $\lim \sup _{\epsilon \rightarrow 0} \hat{x}^{(\epsilon)}(1)=1$.
Proof The idea of the proof is to show that every $y_{0} \in[0,1]$ is eventually fixed by $T_{\epsilon}$ for infinitely many arbitrarily small values of $\epsilon$. Namely, $y_{n}^{(\epsilon)}=T_{\epsilon}^{n} y_{0}=-\epsilon \bmod 1$ for all large enough $n$. Moreover, $y_{n}^{(\epsilon)}$ gets fixed sufficiently early to interfere with the averaging.

Let $y_{0} \in[0,1]$ and $\delta>0$. Let $N=\left[\delta^{-1 / 2}\right]$ and choose an integer $1 \leq k \leq 2^{N}$ such that

$$
y_{0} \in\left[-\delta+(k-1) 2^{-N},-\delta+k 2^{-N}\right] .
$$

Choose $\epsilon$ such that

$$
y_{0}=-\epsilon+(k-1) 2^{-N}
$$

Then

$$
\delta-2^{-N} \leq \epsilon \leq \delta
$$

If $\delta$ is small enough, then $\delta-2^{-N}=\delta-2^{-\left[\delta^{-1 / 2}\right]}>0$, and so $0<\epsilon \leq \delta$.
Now

$$
y_{n}^{(\epsilon)}=2^{n} y_{0}+\left(2^{n}-1\right) \epsilon \bmod 1=-\epsilon+(k-1) 2^{n-N} \bmod 1,
$$

for all $n \geq 0$. In particular, for $n \geq N$ we have $y_{n}^{(\epsilon)}=-\epsilon \bmod 1$, and $\cos 2 \pi y_{n}^{(\epsilon)} \geq 1-$ $\pi \epsilon$. Note that $N \leq \epsilon^{-1 / 2}$. Hence $\hat{x}^{(\epsilon)}(1)=\epsilon \sum_{n=0}^{\left[\epsilon^{-1}\right]-1} \cos \left(2 \pi y_{n}^{(\epsilon)}\right)=1+O\left(\epsilon^{1 / 2}\right)+O(\epsilon)$. Since $\epsilon \in(0, \delta]$ is arbitrarily small, the result follows.

Remark 7.2 A similar argument works for the family of maps $T_{\epsilon} y=2 y+\epsilon^{\beta} \bmod 1$ for any choice of $\beta>0$.

## A Proof of second order averaging

In this appendix, we prove Theorem 2.2. This is a quantitative version of a result due to [31] with a somewhat simplified proof. We work with discrete time rather than continuous time.

First, we consider the case where $T: M \rightarrow M$ is independent of $\epsilon$. Suppose that $a: \mathbb{R}^{d} \times M \rightarrow \mathbb{R}^{d}$ and $\bar{a}: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ are functions. Assume that $\|a\|_{\text {Lip }} \leq L$ and $\|D a\|_{\text {Lip }} \leq L$ where $D=\frac{d}{d x}$ and $L \geq 1$.

For $\epsilon>0$, consider the discrete fast-slow system

$$
x_{n+1}=x_{n}+\epsilon a\left(x_{n}, y_{n}\right), \quad y_{n+1}=T y_{n}
$$

with $x_{0} \in \mathbb{R}^{d}$ and $y_{0} \in M$ given.
Define $\hat{x}_{\epsilon}:[0,1] \rightarrow \mathbb{R}^{d}, \hat{x}_{\epsilon}(t)=x_{[t / \epsilon]}$, and let $X:[0,1] \rightarrow \mathbb{R}^{d}$ be the solution of the ODE $\dot{X}=\bar{a}(X)$ with initial condition $X(0)=x_{0}$. As in Section 2, we define
$\delta_{\epsilon}=\delta_{1, \epsilon}+\delta_{2, \epsilon}$ where

$$
\begin{aligned}
& \delta_{1, \epsilon}\left(y_{0}\right)=\sup _{x} \sup _{1 \leq n \leq 1 / \epsilon} \epsilon\left|\sum_{j=0}^{n-1}\left(a\left(x, y_{j}\right)-\bar{a}(x)\right)\right| \\
& \delta_{2, \epsilon}\left(y_{0}\right)=\sup _{x} \sup _{1 \leq n \leq 1 / \epsilon} \epsilon\left|\sum_{j=0}^{n-1}\left(D a\left(x, y_{j}\right)-D \bar{a}(x)\right)\right| .
\end{aligned}
$$

Theorem A. 1 Let $\epsilon>0, x_{0} \in \mathbb{R}^{d}$. For all $y_{0} \in M$ with $\delta_{\epsilon}\left(y_{0}\right) \leq \frac{1}{2}$ and $t \in[0,1]$,

$$
\left|\hat{x}_{\epsilon}(t)-X(t)\right| \leq 5 e^{2 L}\left(\delta_{\epsilon}\left(y_{0}\right)+\epsilon\right) .
$$

First we recall a discrete version of Gronwall's lemma.
Proposition A. 2 Suppose that $b_{n} \geq 0$ and that there exist constants $C, D \geq 0$ such that $b_{n} \leq C+D \sum_{m=0}^{n-1} b_{m}$ for all $n \geq 0$. Then $b_{n} \leq C(D+1)^{n}$.

Proof This follows by induction.
Define a function $u: \mathbb{R}^{d} \times\{0,1, \ldots,[1 / \epsilon]\}$ by setting $u(x, 0) \equiv 0$ and

$$
u(x, n)=\frac{\epsilon}{\delta_{\epsilon}} \sum_{j=0}^{n-1}\left(a\left(x, y_{j}\right)-\bar{a}(x)\right), \quad n \geq 1
$$

Proposition A. 3 For any $n \leq 1 / \epsilon$, we have $|u(\cdot, n)|_{\infty} \leq 1$ and $\operatorname{Lip} u(\cdot, n) \leq 1$.
Proof For all $x$,

$$
|D u(x, n)|=\frac{\epsilon}{\delta_{\epsilon}}\left|\sum_{j=0}^{n-1}\left(D a\left(x, y_{j}\right)-D \bar{a}(x)\right)\right| \leq \frac{\delta_{2, \epsilon}}{\delta_{\epsilon}} \leq 1 .
$$

Hence the second estimate follows from the mean value theorem, and the first estimate is easier.

Define a new sequence $w_{n}$ by setting $w_{0}=x_{0}$ and

$$
w_{n}=x_{n}-\delta_{\epsilon} u\left(w_{n-1}, n\right), \quad n \geq 1
$$

Lemma A. 4 For all $0 \leq n \leq 1 / \epsilon$ and for all $y \in M$ with $\delta_{\epsilon}(y) \leq \frac{1}{2}$,

$$
\left|w_{n}-w_{0}-\epsilon \sum_{k=0}^{n-1} \bar{a}\left(w_{k}\right)\right| \leq 4 L \delta_{\epsilon} .
$$

Proof By definition, for all $0 \leq k \leq 1 / \epsilon$,

$$
\delta_{\epsilon}\left[u\left(w_{k}, k+1\right)-u\left(w_{k}, k\right)\right]=\epsilon\left[a\left(w_{k}, y_{k}\right)-\bar{a}\left(w_{k}\right)\right] .
$$

Thus

$$
\begin{aligned}
& \delta_{\epsilon} u\left(w_{n-1}, n\right)=\delta_{\epsilon} \sum_{k=0}^{n-1}\left\{u\left(w_{k}, k+1\right)-u\left(w_{k-1}, k\right)\right\} \\
& \left.\quad=\delta_{\epsilon}\left(\sum_{k=0}^{n-1}\left\{u\left(w_{k}, k+1\right)-u\left(w_{k}, k\right)\right\}\right)+\sum_{k=0}^{n-1}\left\{u\left(w_{k}, k\right)-u\left(w_{k-1}, k\right)\right\}\right) \\
& \quad=\epsilon \sum_{k=0}^{n-1}\left\{a\left(w_{k}, y_{k}\right)-\bar{a}\left(w_{k}\right)\right\}+\mathrm{I}_{n},
\end{aligned}
$$

where

$$
\mathrm{I}_{n}=\delta_{\epsilon} \sum_{k=1}^{n-1}\left\{u\left(w_{k}, k\right)-u\left(w_{k-1}, k\right)\right\}
$$

This together with the definition of $x_{n}$ yields

$$
\begin{aligned}
w_{n} & =x_{0}+\epsilon \sum_{k=0}^{n-1} a\left(x_{k}, y_{k}\right)-\delta_{\epsilon} u\left(w_{n-1}, y, n\right) \\
& =w_{0}+\epsilon \sum_{k=0}^{n-1} a\left(x_{k}, y_{k}\right)-\epsilon \sum_{k=0}^{n-1}\left\{a\left(w_{k}, y_{k}\right)-\bar{a}\left(w_{k}\right)\right\}-\mathrm{I}_{n} \\
& =w_{0}+\epsilon \sum_{k=0}^{n-1} \bar{a}\left(w_{k}\right)-\mathrm{I}_{\mathrm{n}}+\mathrm{II}_{n},
\end{aligned}
$$

where

$$
\mathrm{II}_{n}=\epsilon \sum_{k=0}^{n-1}\left\{a\left(x_{k}, y_{k}\right)-a\left(w_{k}, y_{k}\right)\right\} .
$$

We claim that for all $0 \leq n \leq 1 / \epsilon$,

$$
\left|w_{n+1}-w_{n}-\epsilon \bar{a}\left(w_{n}\right)\right| \leq 4 L \epsilon \delta_{\epsilon} .
$$

The result follows by summing over $n$.
It remains to prove the claim. It is easy to check that $w_{1}-w_{0}-\epsilon \bar{a}\left(w_{0}\right)=0$. Inductively, suppose that $\left|w_{n}-w_{n-1}-\epsilon \bar{a}\left(w_{n-1}\right)\right| \leq 4 L \epsilon \delta_{\epsilon}$. Notice that

$$
\left|\mathrm{II}_{n+1}-\mathrm{II}_{n}\right| \leq \epsilon \operatorname{Lip} a\left|x_{n}-w_{n}\right| \leq \operatorname{Lip} a \epsilon \delta_{\epsilon}\left|u\left(w_{n-1}, n\right)\right| \leq L \epsilon \delta_{\epsilon} .
$$

Also,

$$
\left|\mathrm{I}_{n+1}-\mathrm{I}_{n}\right| \leq \delta_{\epsilon} \operatorname{Lip} u\left|w_{n}-w_{n-1}\right| \leq \delta_{\epsilon}\left(\epsilon\left|\bar{a}\left(w_{n-1}\right)\right|+4 L \epsilon \delta_{\epsilon}\right) \leq 3 L \epsilon \delta_{\epsilon}
$$

where the second inequality follows by the induction hypothesis and the third inequality uses $\delta_{\epsilon} \leq \frac{1}{2}$. Therefore

$$
\left|w_{n+1}-w_{n}-\epsilon \bar{a}\left(w_{n}\right)\right| \leq\left|\mathrm{I}_{n+1}-\mathrm{I}_{n}\right|+\left|\mathrm{II}_{n+1}-\mathrm{II}_{n}\right| \leq 4 L \epsilon \delta_{\epsilon},
$$

proving the claim.
Define the sequence

$$
z_{n+1}=z_{n}+\epsilon \bar{a}\left(z_{n}\right), \quad z_{0}=x_{0}
$$

Corollary A.5 $\left|x_{n}-z_{n}\right| \leq 5 \delta_{\epsilon} e^{2 L}$ for all $1 \leq n \leq 1 / \epsilon$, and for all $y \in M$ with $\delta_{\epsilon}(y) \leq \frac{1}{2}$,

Proof Write

$$
w_{n}-z_{n}=w_{n}-x_{0}-\epsilon \sum_{k=0}^{n-1} \bar{a}\left(z_{k}\right)=w_{n}-w_{0}-\epsilon \sum_{k=0}^{n-1} \bar{a}\left(w_{k}\right)+\epsilon \sum_{k=0}^{n-1}\left\{\bar{a}\left(w_{k}\right)-\bar{a}\left(z_{k}\right)\right\} .
$$

By Lemma A.4,

$$
\left|w_{n}-z_{n}\right| \leq 4 L \delta_{\epsilon}+\epsilon \sum_{k=0}^{n-1}\left|\bar{a}\left(w_{k}\right)-\bar{a}\left(z_{k}\right)\right| \leq 4 L \delta_{\epsilon}+L \epsilon \sum_{k=0}^{n-1}\left|w_{k}-z_{k}\right|
$$

By Proposition A.2, $\left|w_{n}-z_{n}\right| \leq 4 L \delta_{\epsilon} e^{L} \leq 4 \delta_{\epsilon} e^{2 L}$. Moreover, $\left|x_{n}-w_{n}\right| \leq \delta_{\epsilon}|u|_{\infty} \leq \delta_{\epsilon}$ and the result follows.

Lemma A. $6\left|X(n \epsilon)-z_{n}\right| \leq L^{2} e^{L} \epsilon$.
Proof Write

$$
X(n \epsilon)=x_{0}+\sum_{m=0}^{n-1} \int_{m \epsilon}^{(m+1) \epsilon}[\bar{a}(X(s))-\bar{a}(X(m \epsilon))] d s+\epsilon \sum_{m=0}^{n-1} \bar{a}(X(m \epsilon)) .
$$

Since $\left|\bar{a}\left(X\left(t_{1}\right)\right)-\bar{a}\left(X\left(t_{2}\right)\right)\right| \leq L\left|X\left(t_{1}\right)-X\left(t_{2}\right)\right| \leq L^{2}\left|t_{1}-t_{2}\right|$ for all $t_{1}, t_{2}$, we obtain that

$$
\left|X(n \epsilon)-x_{0}-\epsilon \sum_{m=0}^{n-1} \bar{a}(X(m \epsilon))\right| \leq L^{2} n \epsilon^{2} \leq L^{2} \epsilon
$$

Hence

$$
\left|X(n \epsilon)-z_{n}\right| \leq \epsilon \sum_{m=0}^{n-1}\left|\bar{a}(X(m \epsilon))-\bar{a}\left(z_{m}\right)\right|+L^{2} \epsilon \leq L \epsilon \sum_{m=0}^{n-1}\left|X(m \epsilon)-z_{m}\right|+L^{2} \epsilon
$$

The result follows from Proposition A.2.

Proof of Theorem A. 1 By Corollary A.5 and Lemma A.6,

$$
\begin{aligned}
\left|x_{[t / \epsilon]}-X(t)\right| & \leq|X(t)-X([t / \epsilon] \epsilon)|+\left|X([t / \epsilon] \epsilon)-z_{[t / \epsilon]}\right|+\left|z_{[t / \epsilon]}-x_{[t / \epsilon]}\right| \\
& \leq \epsilon L+L^{2} e^{L} \epsilon+5 e^{2 L} \delta_{\epsilon} \leq 5 e^{2 L}\left(\delta_{\epsilon}+\epsilon\right) .
\end{aligned}
$$

This completes the proof.
Proof of Theorem 2.2 Replacing $T, a(x, y)$ and $\bar{a}(x)$ in Theorem A.1 by $T_{\epsilon}$, $a(x, y, \epsilon)$ and $\bar{a}(x, \epsilon)$, we obtain that

$$
\left|\hat{x}_{\epsilon}(t)-X_{\epsilon}(t)\right| \leq 5 e^{2 L}\left(\delta_{\epsilon}+\epsilon\right)
$$

where

$$
\dot{X}_{\epsilon}=\bar{a}(X, \epsilon), \quad X_{\epsilon}(0)=x_{0} .
$$

Let $A_{\epsilon}=\sup _{x \in E}|\bar{a}(x, \epsilon)-\bar{a}(x, 0)|$. Then

$$
\begin{aligned}
\left|X_{\epsilon}(t)-X(t)\right| & \leq \int_{0}^{t}\left|\bar{a}\left(X_{\epsilon}(s), \epsilon\right)-\bar{a}(X(s), 0)\right| d s \\
& \leq t A_{\epsilon}+\int_{0}^{t}\left|\bar{a}\left(X_{\epsilon}(s), 0\right)-\bar{a}(X(s), 0)\right| d s \\
& \leq A_{\epsilon}+L \int_{0}^{t}\left|X_{\epsilon}(s)-X(s)\right| d s
\end{aligned}
$$

By Gronwall's lemma, $\left|X_{\epsilon}(t)-X(t)\right| \leq e^{L} A_{\epsilon}$ for all $t \leq 1$.
Next, $A_{\epsilon} \leq L \epsilon+\sup _{x \in E}\left|\int_{M} a(x, y, 0)\left(d \nu_{\epsilon}-d \nu_{0}\right)(y)\right|$. Combining these estimates we obtain that

$$
\left|\hat{x}_{\epsilon}(t)-X(t)\right| \leq 5 e^{2 L} \delta_{\epsilon}+6 e^{2 L} \epsilon+e^{L} \sup _{x \in E}\left|\int_{M} a(x, y, 0)\left(d \nu_{\epsilon}-d \nu_{0}\right)(y)\right|
$$

yielding the result.

Acknowledgements This research was supported in part by a European Advanced Grant StochExtHomog (ERC AdG 320977). We are grateful to Vitor Araújo, Jorge Freitas, Vilton Pinheiro, Mike Todd and Paulo Varandas for helpful discussions.

## References

[1] J. F. Alves. SRB measures for non-hyperbolic systems with multidimensional expansion. Ann. Sci. École Norm. Sup. (4) 33 (2000) 1-32.
[2] J. F. Alves. Strong statistical stability of non-uniformly expanding maps. Nonlinearity 17 (2004) 1193-1215.
[3] J. F. Alves, C. Bonatti and M. Viana. SRB measures for partially hyperbolic systems whose central direction is mostly expanding. Invent. Math. 140 (2000) 351-398.
[4] J. F. Alves, S. Luzzatto and V. Pinheiro. Markov structures and decay of correlations for non-uniformly expanding dynamical systems. Ann. Inst. H. Poincaré Anal. Non Linéaire 22 (2005) 817-839.
[5] J. F. Alves and M. Viana. Statistical stability for robust classes of maps with non-uniform expansion. Ergodic Theory Dynam. Systems 22 (2002) 1-32.
[6] D. V. Anosov. Averaging in systems of ordinary differential equations with rapidly oscillating solutions. Izv. Akad. Nauk SSSR Ser. Mat. 24 (1960) 721-742.
[7] V. Bakhtin and Y. Kifer. Nonconvergence examples in averaging. Geometric and probabilistic structures in dynamics. Contemp. Math. 469, Amer. Math. Soc., Providence, RI, 2008, pp. 1-17.
[8] V. Baladi. On the susceptibility function of piecewise expanding interval maps. Comm. Math. Phys. 275 (2007) 839-859.
[9] V. Baladi, M. Benedicks and D. Schnellmann. Whitney-Hölder continuity of the SRB measure for transversal families of smooth unimodal maps. Invent. Math. 201 (2015) 773-844.
[10] M. Benedicks and L. Carleson. On iterations of $1-a x^{2}$ on $(-1,1)$. Ann. of Math. 122 (1985) 1-25.
[11] V. Baladi and M. Todd. Linear response for intermittent maps. Comm. Math. Phys. 347 (2016) 857-874.
[12] D. L. Burkholder. Distribution function inequalities for martingales. Ann. Probability 1 (1973) 19-42.
[13] P. Collet and J.-P. Eckmann. Positive Liapunov exponents and absolute continuity for maps of the interval. Ergodic Theory Dynam. Systems 3 (1983) 13-46.
[14] D. Dolgopyat. Averaging and invariant measures. Mosc. Math. J. 5 (2005) 537576, 742.
[15] J. M. Freitas. Continuity of SRB measure and entropy for Benedicks-Carleson quadratic maps. Nonlinearity 18 (2005) 831-854.
[16] J. M. Freitas. Exponential decay of hyperbolic times for Benedicks-Carleson quadratic maps. Port. Math. 67 (2010) 525-540.
[17] J. M. Freitas and M. Todd. The statistical stability of equilibrium states for interval maps. Nonlinearity 22 (2009) 259-281.
[18] S. Gouzel. Decay of correlations for nonuniformly expanding systems. Bull. Soc. Math. France 134 (2006) 1-31.
[19] M. V. Jakobson. Absolutely continuous invariant measures for one-parameter families of one-dimensional maps. Comm. Math. Phys. 81 (1981) 39-88.
[20] G. Keller. Stochastic stability in some chaotic dynamical systems. Monatsh. Math. 94 (1982) 313-333.
[21] G. Keller and C. Liverani. Stability of the spectrum for transfer operators. Annali della Scuola Normale Superiore di Pisa, Classe di Scienze XXVIII (1999) 141152.
[22] Y. Kifer. Averaging in difference equations driven by dynamical systems. Astérisque (2003), no. 287, xviii, 103-123, Geometric methods in dynamics. II.
[23] Y. Kifer. Averaging principle for fully coupled dynamical systems and large deviations. Ergodic Theory Dynam. Systems 24 (2004) 847-871.
[24] Y. Kifer. Another proof of the averaging principle for fully coupled dynamical systems with hyperbolic fast motions. Discrete Contin. Dyn. Syst. 13 (2005) 11871201.
[25] A. Korepanov. Linear response for intermittent maps with summable and nonsummable decay of correlations. Nonlinearity 29 (2016) 1735-1754.
[26] A. Korepanov, Z. Kosloff and I. Melbourne. First (much longer) version of this paper, available at arXiv:1509.06772v2.
[27] A. Korepanov, Z. Kosloff and I. Melbourne. Explicit coupling argument for nonuniformly hyperbolic transformations. Preprint, September 2016. To appear in Proc. Roy. Soc. Edinburgh.
[28] C. Liverani, B. Saussol and S. Vaienti. A probabilistic approach to intermittency. Ergodic Theory Dynam. Systems 19 (1999) 671-685.
[29] M. Lyubich. Almost every real quadratic map is either regular or stochastic. Ann. of Math. 156 (2002) 1-78.
[30] I. Melbourne and R. Zweimüller. Weak convergence to stable Lévy processes for nonuniformly hyperbolic dynamical systems. Ann Inst. H. Poincaré (B) Probab. Statist. 51 (2015) 545-556.
[31] J. A. Sanders. On the fundamental theorem of averaging. SIAM J. Math. Anal. 14 (1983) 1-10.
[32] J. A. Sanders, F. Verhulst and J. Murdock. Averaging methods in nonlinear dynamical systems, second ed., Applied Mathematical Sciences 59, Springer, New York, 2007.
[33] H. Thunberg. Unfolding of chaotic unimodal maps and the parameter dependence of natural measures. Nonlinearity 14 (2001) 323-337.
[34] M. Viana. Multidimensional nonhyperbolic attractors. Inst. Hautes Études Sci. Publ. Math. (1997) 63-96.
[35] L.-S. Young. Recurrence times and rates of mixing. Israel J. Math. 110 (1999) 153-188.


[^0]:    ${ }^{1} q-$ denotes $q-a$ for all $a>0$

