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Review of pore network modelling of porous media: experimental characterisations, network constructions and applications to reactive transport

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7 8

9 Abstract

10 Pore network models have been applied widely for simulating a variety of different 11 physical and chemical processes, including phase exchange, non-Newtonian 12 displacement, non-Darcy flow, reactive transport and thermodynamically consistent oil 13 layers. The realism of such modelling, i.e. the credibility of their predictions, depends to 14 a large extent on the quality of the correspondence between the pore space of a given 15 medium and the pore network constructed as its representation. The main experimental 16 techniques for pore space characterisation, including direct imaging, mercury intrusion 17 porosimetry and gas adsorption, are firstly summarised. A review of the main pore 18 network construction techniques is then presented. Particular focus is given on how 19 such constructions are adapted to the data from experimentally characterised pore 20 systems. Current applications of pore network models are considered, with special 21 emphasis on the effects of adsorption, dissolution and precipitation, as well as biomass 22 growth, on transport coefficients. Pore network models are found to be a valuable tool 23 for understanding and predicting meso-scale phenomena, linking single pore processes, 24 where other techniques are more accurate, and the homogenised continuum porous 25 media, used by engineering community.

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27 Keywords: Porous media; Reactive transport; Pore network model

2 1 Introduction

3

4 Flow and transport phenomena in porous media play an important role in many diverse 5 fields of science and technology. For example, radioactive waste management is one of 6 the most pressing problems facing the world today, because of the longevity of 7 radionuclide and the possibility of their transport to the surface environment. For long-8 term performance assessment of nuclear repositories, knowledge concerning the 9 transport of radionuclide in the back-fill material is required. One porous medium, 10 bentonite, is presently considered as the best candidate for the high level waste disposal, 11 due to its large specific surface area, high ion-exchange capacity, and sorption affinity 12 for organic and inorganic ions. Sorption onto bentonite plays an important role in retarding the migration of radionuclide from a waste repository (Bourg et al., 2003; 13 14 Bourg et al., 2006; Bradbury and Baeyens, 2003). Another field where such knowledge 15 is highly demanded is the construction industry. Building materials such as bricks, 16 concrete and sandstone all are porous. These materials may interact with their 17 environment leading to degradation of the structures with time. A specific example is 18 that the constituent ions in the hydrated cement paste matrix may leach out from the 19 concrete and cause the concrete structure to become weak. Degradation processes like 20 this strongly relate to the transport phenomena in concrete. Transport mechanisms 21 involved in the degradation processes include: permeation of water or aqueous solutions, 22 diffusion of gaseous compounds and sorption of ionic contaminated water. In such a 23 case, transport properties in these porous media are usually considered as indicators for 24 evaluating the durability and ultimately predetermining the service life of these 25 structures (Zhang et al., 2013; Zhang et al., 2014, (Wang et al., 2015a; Wang et al., 26 2015b; Zhang et al., 2013). Another example is the contaminated underground water 27 transport, which is controlled by a variety of transport mechanisms including 28 permeation, diffusion and dispersion (McCarthy and Zachara, 1989).

29

There are several challenges in analysing transport in porous media. Firstly, the intricacy of the pore structure makes the transport processes in porous media very complex. Pores tend to have irregular surfaces and some of them even make dead ends.

1 These factors influence the flow and transport behaviour significantly. Another 2 difficulty of studying transport processes in porous media is the evolution of pore 3 structure during their service or operation, etc. These changes may result from chemical, 4 electrochemical or bacterial effects as well as from mechanical damage, such as micro 5 cracking. The evolution of the pore space may lead to changes in the macroscopic 6 transport properties, such as permeability and diffusivity. Furthermore, the transport 7 properties in porous media are a function of types of species in solution and the phase 8 (Appelo et al., 2010). Finally, many reactive pore network models (Li et al., 2006; 9 Raoof et al., 2012) employ a reaction rate or mass transfer coefficient measured in a 10 macroscopic experiment. For example, these authors determine reaction rate 11 coefficients from spinning disk, flow-through reactor, or continuous fluidized bed 12 reactor experiments. These results may not be applicable to reaction rates in individual 13 pores, as the smaller scale magnifies the effects of mass transfer resistance and 14 concentration gradients. Raoof et al. (2010) demonstrated the importance of 15 incorporating kinetics developed for the pore scale, rather than relying on macroscopic 16 averages.

17

18 Due to the complexity of the pore structure and the change of the environmental 19 conditions, a great deal of experimental, theoretical and numerical approaches have 20 been proposed and developed to study transport processes through porous media during 21 the past decades (Abichou et al., 2004; Aytas et al., 2009; Boult et al., 1998; Ghassemi 22 and Pak, 2011; Kohler et al., 1996; Li et al., 1992; Ren et al., 2009; Tsai and Chen, 23 1995; Wang et al., 2005). Field and reactive transport experiments as well as studies 24 involving X-ray, SEM, TEM etc. are usually expensive and time consuming. The 25 measurements are highly sensitive to the material composition, sample preparation, 26 methodology and testing environment. Analytical solutions are typically restricted to 27 problems with assumed homogeneous properties and specific boundary conditions, 28 some of which have limited practical relevance or are complicated to evaluate.

29

30 Pore-scale simulations have improved the understanding of large-scale natural processes 31 and informed large-scale geotechnical applications. Their importance comes from the 32 fact that they can produce rather cost-effective and accurate predictions for local

1 transport (diffusion/permeation), and at the same time allow for systematic variations of 2 the system's parameters (pore space geometries, fluid properties, and boundary 3 conditions) to assess their impact, which is much more difficult to achieve than with 4 experiments (Meakin and Tartakovsky, 2009). With pore-scale models one can make 5 improved assessments of macroscopic transport properties by varying the pore space 6 structure parameters. This offers a way to understand the scale dependence of 7 continuum transport parameters. Such scale dependence cannot be captured by an 8 effective medium Darcy approach. The pore-scale modelling is dominated by particle-9 based methods. These include the promising lattice Boltzmann method (Hao and Cheng, 10 2010; Zhang et al., 2014; Zhang et al., 2013) and smoothed particle hydrodynamics 11 (Tartakovsky and Meakin, 2006; Tartakovsky et al., 2007; Zhu and Fox, 2002).

12

13 The particle methods, while suitable for pore-scale analyses, become inefficient at the 14 meso-scale, e.g. when the system requiring analysis has tens or hundreds interconnected 15 pores in each direction (Tartakovsky et al., 2007). Further, these methods are time 16 consuming and only very limited pore volumes can be addressed. In order to increase 17 the pore volume with affordable computational resources and a reduced impact on the 18 reliability of the results, the pore network model (PNM) approach has been used to 19 study reactive transport phenomena. At the meso-scale the classical macroscopic 20 equations, such as the Darcy law, are yet not needed, i.e. fluid flow and solute transport 21 processes are simulated directly in the pores, but the precise particle dynamics, which 22 can be analysed by particle methods, is not accounted for. This is done by creating a 23 virtual representation of the porous medium consisting of pore bodies and pore throats 24 of different sizes (the "geometry" of the porous medium), connected to each other as 25 required (the "topology" of the porous medium). It is then possible to simulate the fluid 26 flow and solute transport process of interest at the meso-scale through this network, 27 with the relevant physics implemented on a pore-to-pore basis. The first network model 28 was constructed by Fatt (Fatt, 1956) who exploited the analogy between flow in porous 29 media and a random resistor network. Afterwards, the models have grown in 30 sophistication and now can deal with irregular lattices, wetting layer flow, arbitrary 31 wettability and any sequence of displacement in two- and three-phase flow, as well as a 32 variety of different physical processes, including phase exchange, non-Newtonian

displacement, non-Darcy flow, reactive transport and thermodynamically consistent oil
 layers, etc., (Balhoff and Wheeler, 2009; Blunt et al., 2002; Lopez et al., 2003;
 Ryazanov et al., 2009; Yiotis et al., 2006).

4

5 Compared to the pore-scale modelling methods, such as the lattice-Boltzmann and 6 particle methods, pore-network models take less time to compute and require less 7 computer capacity due to the inherent simplifications of the pore-space in the 8 construction of the model. Pore network simulations are much less computationally 9 demanding than direct methods. This allows researchers to incorporate more 10 heterogeneity in modelling larger rock volumes. As natural geological media can be 11 heterogeneous on all length scales, this is an important advantage of pore network 12 modelling. Pore network models have been used extensively to simulate multiphase and 13 single-phase fluid flow in porous media, and these models will continue to provide 14 important insight and information in the future. However, pore network models are 15 usually constructed based on simplified geometry due to the lack of experimental 16 resolution in pore space characterisation. This could make the predictions of transport 17 properties less credible. The greatest challenge in obtaining credible results is the 18 identification of features/phenomena relevant to the modelled process and neglecting 19 the rest to reduce the computational load. For example, to compute the single phase 20 flow field, it is necessary that the principal connected pore space is modelled; while the 21 very small pores below the image resolution can be neglected as they make little 22 contribution to the overall behaviour (Blunt et al., 2013).

23

24 The suitability of the pore-scale modelling techniques for a given application depends 25 on the governing equations, the underlying assumptions for the pore-scale flow and 26 transport equations, as well as the length-scales of the (computational) domain, etc. 27 While the lower scale limit of a pore-scale technique is determined by the scale of the 28 governing equations, the upper scale limit is set by the computational power. In this 29 review we focus on the pore network models with comprehensive account for the 30 methods for obtaining pore space information, constructing pore networks and the 31 application of pore network models to reactive transport in porous media.

2 Characterisation of pore space information

2

3 For pore network model construction, the geometry and topology of the pore space are 4 required. There are several ways to characterize the pore space. One is imaging 5 techniques such as producing 3D images by mapping the real interior structure of 6 original samples (the destructive approach of cutting and stacking serial 2D sections, 7 confocal laser scanning microscopy and non-destructive X-ray micro-tomography 8 (micro-CT)), and constructing synthetic 3D rock images from high resolution 2D thin 9 sections using statistical methods or geological process simulation. Other non-10 destructive approaches are mercury intrusion porosimerty and gas adsorption, which 11 produce pore size distribution and surface area, but not the full connectivity. The main 12 methods for space characterisations are reviewed in this section.

13

14 **2.1 Imaging techniques**

15

16 2.1.1 X-ray computed micro-tomography (micro-CT)

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Micro-CT is a non-destructive and non-invasive imaging technique used to characterize
cross-sectional and three-dimensional internal structures (Hazlett, 1995; Lindquist et al.,
1996; Schl üter et al., 2014).

21

22 To image geological porous materials at the micro-scale, three types of micro-CT 23 systems are in common use: medical CT, industrial X-ray generation tube and 24 synchrotron micro-tomography. They primarily differ in X-ray energy and source, 25 means of sample manipulation and detector geometry. Although the best image 26 resolution reported in the literature is from synchrotron micro-CT, the samples need to 27 be relatively small to achieve these resolutions. This may result in poor statistical 28 representation of the bulk material. Typical spatial resolution that medical CT systems 29 can achieve is between 200 to 500 µm, industrial systems range from 50 to 100 µm and 30 synchrotron based CT systems can reach from 1 µm to 50 µm (Wildenschild et al., 31 2002a). Presently, laboratory systems with genuine submicron capabilities exist, some 32 with voxel resolution of 20nm, providing spatial resolution of 50-60 nm. Three main

configurations are used in systems that seek submicron resolution; a good review can be
found in (Schlüter et al., 2014; Withers, 2007). Specific details of these imaging
techniques and their evolution can be found in previous reviews (Blunt et al., 2013;
Ketcham and Carlson, 2001).

5

It should be noted that X-ray CT may obscure significant features or misinterpret 6 7 attenuation values of a single material in different image sections resulting in 8 complicated quantitative image analysis (Ketcham and Carlson, 2001). Commonly 9 experienced problems include beam hardening, high-frequency noise, scattered X-rays, 10 defective detector pixels or poorly cantered samples (Ketcham and Carlson, 2001; 11 Wildenschild et al., 2002b). In addition, there might be errors and distortions arising 12 from the CT reconstruction (Ketcham and Carlson, 2001). These problems cannot be 13 completely avoided, but can be alleviated by careful detector calibration, metal filters 14 and sample centring. Application of wedge calibration and dual energy scans (Rebuffel 15 and Dinten, 2007) can also reduce beam hardening effects, but these techniques are 16 rarely implemented in standard scanning procedures (Ketcham and Carlson, 2001).

17

18 2.1.2 Focused ion beams and Scanning electron microscopy

19

20 Focused ion beams (FIB) and Scanning electron microscopy (SEM) are destructive 21 imaging techniques. SEM is a useful technique for extracting two-dimensional (2D) 22 images of the microstructures but does not provide the third spatial component of the 23 sample, which is important to find interconnected regions and pore volumes, shapes and 24 sizes. Depending on the instrument, the resolution achieved can be between 1 and 20 25 nm. The world's highest resolution conventional SEM can reach a point resolution of 26 0.4 nm by utilising a secondary electron detector (http://www.nanotech-27 now.com/news.cgi?story_id=42612). FIB is a well-established technique to acquire very high-resolution three-dimensional images, typically just a few micro-meters across 28 29 (Curtis et al., 2010; Lemmens et al., 2010; Tomutsa et al., 2007). This technique can 30 1 achieve less than nm imaging resolution 31 (http://www.fibics.com/fib/tutorials/introduction-focused-ion-beam-systems/4/).

32 Although this method provides a great potential for imaging hydrocarbon bearing rocks

at high resolution and produces images with better quality than the electron beam
imaging due to less charging of the surface, it is still significantly time consuming due
to the refocusing between milling and imaging as well as the sample repositioning
(Tomutsa et al., 2007). Due to these limitations, this technique exposes only small areas
of observation and cannot provide adequate sampling to characterize the sample, e.g.
from a shale reservoir.

7

8 The combination of FIB and SEM (FIB-SEM) is usually used to compute micro-9 structural properties of porous media. This combination allows for the observation of 10 fine macro-pores and meso-pores within the porous medium (Keller et al., 2011; 11 Tomutsa et al., 2007). This method can typically achieve voxel dimensions of tens of 12 nanometers, thus allowing for analysing volumes of around $(10 - 30)^3 \mu m^3$ within 13 practical measuring time (Michael D. Uchic, 2007).

14

Multi-scale imaging capabilities, such as the development and integration of a range of experimental and computational tools, can facilitate probing the structure of materials across various scales in an integrated fashion. Given their important function, multiscale imaging capabilities have been an area of focus of several research groups in the recent years.

20

21 The combination of FIB, SEM and TEM (transmission electron microscopy) is used to 22 analyse the structure of pore space (Wirth, 2009). For example, the backscattered 23 scanning electron microscopy (BSEM) and focused ion beam SEM (FIBSEM) have 24 been utilised to analyse rocks (Sok et al., 2010). These techniques are particularly 25 suitable for examining carbonates due to their multi-modal pore structure, which can 26 range from 10 nm to 10 cm. Sok et al (Sok et al., 2010) report on both plug to pore scale 27 registration in 3D, as well as pore-scale to submicron scale registration of features. The 28 plug to pore scale approach requires one or more subsets of the sample originally 29 imaged (at the scale of 4 cm and at 20 μ m/pixel) which then have to be reimaged with 30 Micro CT at a smaller sample volume (5 mm and at 2.5 µm/pixel). After that the data 1 should be carefully registered and integrated. The pore-scale to submicron scale 2 registration was done by cutting the subsamples and preparing thin sections from the 3 original field of view of the Micro CT images. The thin sections were then imaged with 4 BSEM and registered to the Micro CT image. Finally, in order to estimate reservoir 5 properties for complex materials, such as carbonate and mudstones, FIBSEM images 6 were also obtained at voxel resolutions of 50 nm. A longer term goal of Sok et al_group 7 is to undertake complete multi-scale registration from the whole core and/or plug scale 8 through the Micro CT length scale and down to the submicron (SEM or FIBSEM) scale.

9

10 2.1.3 Nuclear magnetic resonance

11

Recent NMR development has demonstrated the use of relaxation, cryoporometry, spectroscopy, diffusion, and imaging techniques to quantify pore structure (i.e. pore kized distribution, pore morphology, connectivity etc.), fluid properties, and rock heterogeneity. The NMR relaxometry poses the least demands on magnetic field quality, NMR imaging favours magnetic fields linear in space with a constant gradient, and NMR spectroscopy, places the highest demands on field homogeneity and field stability.

19 Magnetic resonance imaging allows the imaging of the interior of the rocks to obtain the 20 spatial distribution across a much larger scale (Blümich et al., 2009). Callaghan et al. 21 (Callaghan, 1993) discovered the diffusion diffraction phenomenon when the gradient 22 wavelength approaches the characteristic pore size. Since the water molecules in the 23 porous material will move randomly, they will probe the pore structure. This will 24 influence the transverse relaxation time. Therefore, NMR can provide information on 25 the pore-size distribution of a porous material by measuring this transverse relaxation 26 time. The major advantage of NMR in comparison with classical methods is the short 27 measurement time. This would allow the analysis of large quantities of samples as 28 required to characterize a field or catchment-scale hydraulic properties, which are 29 necessary for risk assessment (e.g., flood forecasting) and management (e.g., 30 fertilization and pest control). The determination of pore size distributions by NMR 31 relaxometry has its own drawbacks. First, the diffusion in induced magnetic field 32 gradients can shorten transverse relaxation times (T_2) . This must be checked and can be

1 minimized by choosing sufficiently small echo times or by measuring longitudinal 2 relaxation (T_1) . Second, diffusion in internal gradients may affect different modes of a 3 multimodal relaxation time distribution function in a different way. Third, one must be 4 aware that the derivation of PSD is always a scaling procedure, which requires 5 independent determination of the average specific surface area. The reason is that the 6 relaxation times are controlled by the pore sizes and surface relaxivity, but the average 7 S/V is controlled by the internal surface area of the porous media. Fourth, the 8 assumption of the homogeneous distribution of pores and paramagnetic centres and the 9 calculation of an average surface relaxivity parameter, especially for a multimodal 10 relaxation time distribution function, leads to an overestimation of the large pores. All 11 these issues, together with the necessary simplification of pore shape and geometry, can 12 deviate the calculated pore size distribution from the real one (Stingaciu et al., 2010).

13

14 Nuclear magnetic resonance (NMR) relaxometry uses the random motion of molecules, 15 whereas cryoporometry uses the melting-point depression of a confined liquid 16 determining pore size distributions. It is suitable for measuring pore diameters in the 17 range 2 nm-1 µm, depending on the absorbate. Whilst NMR cryoporometry is a 18 perturbative measurement, the results are independent of spin interactions at the pore 19 surface and so can offer direct measurements of pore volume as a function of pore 20 diameter (Mitchell et al., 2008). NMR cryoporometry (Strange et al., 1993) rely on the 21 Gibbs-Thomson equation concerning the relationship between the characteristic pore 22 length scale and the change in the freezing point of the liquid, or melting point of its 23 solid crystal, due to confinement within the porous matrix. However, the freezing and 24 melting behaviour of confined liquids/solids is often complex, with the thermodynamic 25 properties of the confined material being modified from those of the bulk (Christenson, 26 2001). NMR cryoporometry offers the advantage of a more direct measure of the open 27 pore volume.

28

The pore-size distribution obtained from cryoporometry is the derivative of the total liquid water content with respect to temperature. Because noise is always present in the NMR data, the derivative can have an irregular shape for pores having only a minor contribution to the total signal. However, no a priori shape of a distribution is imposed. 1 On the other hand, the pore-size distribution obtained from relaxometry is the result of a 2 numerical inverse Laplace transform. This is generally an ill-posed problem and hence 3 some sort of regularization method is used. The numerical code that we use always 4 yields a sum of log-normal Gaussian-shaped pore-size distributions. Therefore the shape 5 of the relaxometry pore-size distribution does not accurately reflect the actual pore-size 6 distribution. Only the position of the peaks and the total area of the distribution around a 7 specific peak are relevant parameters in this pore-size distribution (Valckenborg et al., 8 2002).

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2.2 Mercury Intrusion Porosimetry (MIP)

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12 Mercury intrusion porosimetry is by far the most popular method for characterizing 13 porous materials with pores in the range of 500 µm down to 3nm (Giesche, 2006; Le ón 14 y León, 1998; Rouquerol et al., 2012). Compared to alternative characterization 15 methods such as gas sorption, mercury porosimetry covers a much wider pore size 16 range, while it is based on simpler physicochemical principles and it is much faster in 17 operation. The most important limitations of the mercury intrusion when applied to 18 extract pore size distribution of a porous material is that it is based on the assumption 19 that the porous matrix can be represented by a bundle of cylindrical pores. In extracting 20 pore size distribution during mercury intrusion, all pores are assumed to be equally 21 accessible to the exterior mercury reservoir. This assumption can only be met if the pore 22 structure is represented in the form of a bundle of capillaries or if pore connectivity is 23 very high. In reality, however, pore network effects can be quite important resulting in 24 the so-called pore shadowing or ink-bottle phenomenon. In this case, a large pore has 25 smaller entrances connecting with the mercury reservoir and can only be filled at a 26 pressure higher than that required by its actual dimensions. Capillary network models of 27 varying complexity have been used over the past three decades and have been often 28 quite successful in proving structural characteristics of the porous materials examined 29 (Ioannidis and Chatzis, 1993a; Kikkinides and Politis, 2014). For obvious reasons this 30 method could be applied to study only the open porosity features, since the closed pores 31 are inaccessible to the mercury (Giesche, 2006).

1 **2.3 Gas Adsorption**

2

3 Adsorption behaviour of porous materials is a function of their microstructural 4 characteristics, such as the surface area, types of pores present in the material, topology 5 of the porous network and the available pore volume. Characterisation of porous 6 materials is, therefore, important in the development of adsorption applications. 7 Mercury porosimetry and gas adsorption are determined from surface tension, capillary 8 forces and pressure. The main difference is that with mercury porosimetry, large pores 9 at the intrusion phase are determined first, while with gas adsorption, the smallest pores 10 are measured first at the adsorption phase.

11

Physical gas adsorption has been used to study the pore characteristics of solid materials and the changes therein upon post-synthesis treatment. Gas adsorption is one of the most popular techniques used to characterize the pore space. It only allows determining the volume of open pores while closed porosity cannot be accessed. The advantage of this technique is that it allows assessing a wide range of pore sizes, covering essentially the completed micro-and meso-pore range in a timely and cost-effective fashion (Thommes, 2010).

19

20 Frequently used adsorptives are nitrogen (N₂), argon (A_r), and CO₂, depending on the 21 nature of the material (adsorbent) and the information required. Nitrogen at 77 K is 22 considered to be a standard adsorptive for surface area and pore size analysis, but it is 23 meanwhile generally accepted that nitrogen adsorption is not satisfactory with regard to 24 a quantitative assessment of the micro-porosity, especially for micro- pores with widths 25 smaller than 0.7 nm. Consequently, alternative probe molecules have been suggested, 26 e.g., argon and carbon dioxide. For many micro porous systems, argon adsorption at 27 87.3 K appears to be very useful as N₂ adsorption in microspores occurs at lower p/p_0 28 values than Ar (Groen et al., 2003; Ravikovitch et al., 2000; Serrano et al., 2009). Despite this advantage, the low pressures induced restricted diffusion prevents argon 29 30 molecules from entering the narrowest micropores, i.e., pores of widths < ca. 0.45 nm. 31 Furthermore, Ar adsorption at 77 K shows limited application for meso-pore size 32 determination, since the coolant temperature is below the bulk triple point. As a

1 consequence, pore condensation vanishes in case the pore diameter exceeds 2 approximately 12 nm (Thommes et al., 2002). When compared to nitrogen and carbon 3 dioxide, it exhibits weaker attractive fluid-pore wall attractions for most adsorbents, 4 which during adsorption does not give rise to specific interactions (like nitrogen and 5 carbon dioxide because of their quadrupole moments) with most of surface functional 6 groups and exposed ions. CO_2 is another usually preferred adsorptive, since these 7 adsorption measurements are mostly performed at temperatures near ambient, which 8 will enhance diffusion properties in the highly micro porous system compared to the 9 low temperatures used in N_2 and Ar adsorption (Rouquerol et al., 2013). A drawback of 10 CO₂ adsorption at ambient temperature is that in most commonly used equipments, 11 which predominantly operate in the pressure range of vacuum to 1 bar, only a limited 12 range of micropores can be measured, unless high-pressure CO₂ adsorption is used 13 (Cazorla-Amorós et al., 1998; Ravikovitch et al., 2000). While CO₂ adsorption at 273 K 14 is frequently used for the ultramicopore analysis of carbonaceous materials (Rodriguez-15 Reinoso, 2009), it is not a good choice for the pore size analysis of materials with polar 16 sites, mainly because of the very specific interactions that CO₂ can have with functional 17 groups on the surface.

18

19 After obtaining the physical adsorption isotherms, the interpretation of experimentally 20 measured isotherms is required in order to arrive the specific characteristics of the 21 material (Thommes, 2010). Classical theories include the Brunauer-Emmett-Teller 22 (BET) method (Brunauer et al., 1938), Barrett-Joiner-Halenda (BJH) (Barrett et al., 23 1951), Broekhoff-de Boer (BdB) method (Broekhoff and De Boer, 1967), Horvath-24 Kawazoe (HK) model, (Horvath and Kawazoe, 1983), and the Saito-Foley (SF) model 25 (Saito and Foley, 1991; Saito and Foley, 1995). Typically, these theories involve certain 26 assumptions, which are not universally applicable. For example, the notion of an 27 adsorbed monolayer required by the BET theory to assess the surface area is clearly not 28 applicable to materials with narrow pores and complex energy landscapes. The classical 29 BJH, based on the Kelvin equation and corrected for multi-layer adsorption, is most 30 widely used for calculations of the pore size distribution (PSD) over the mesopore and 31 part of the macropore range. However, the BJH and BdB theories of extracting PSD, 32 based on a description of a porous material as a collection of cylindrical or slit pores,

1 are not appropriate for zeolites, metal organic frameworks, polymers etc. The 2 conventional HK model for slit-shaped pores and SF model for cylindrical pore 3 geometry are mainly applied for micro-pore size calculations (Groen et al., 2003). Such 4 approaches allow for obtaining the pore size distribution in addition to the pore volume, 5 but rely on similar macroscopic and thermodynamic assumptions concerning the nature 6 of confined adsorbate. This leads to inaccurate determination of the pore size and 7 volume. New models based on non-local density functional theory (NLDFT) and 8 molecular simulations were developed (Lukens et al., 1999; Neimark and Ravikovitch, 9 2001; Neimark et al., 2000). It has been demonstrated that the application of these novel 10 theoretical and molecular simulation based methods leads to: (i) a much more accurate 11 pore size analysis (Neimark and Ravikovitch, 2001; Thommes et al., 2006), and (ii) 12 allows performing pore size analysis over the complete micro/mesopore size range 13 (Thommes et al., 2006). Currently there are methods for pore size quantification based 14 on NLDFT and molecular simulations, which are commercially available and applicable 15 to a range of adsorptive/adsorbent systems. They include hybrid methods that assume 16 various pore geometries for the micro- and meso-pore size range, as it can be found for 17 materials with hierarchical pore structures.

18

Previously published review on the use of gas adsorption for characterization of porous materials hardly comment on these phenomena (Sing, 2001). One drawback is that the process could be very time consuming, but the determinable pore diameter is from 0.3 to 300 nm, a range not completely covered by mercury porosimetry.

23

24 **3 Pore network model construction**

25

The different experimental techniques for pore space characterisation present different limitations to the way representative pore networks can be constructed. For example, if a 3D image is available (experimental or synthetically generated), pore networks can be constructed directly from this image, assuming what is considered to be a pore and what to be a throat between pores. The need for up-scaling from the typically small imaged volume to larger domains leads to the need for construction of statistically representative networks. This requires analysis of the image to extract size distributions of pores and throats and their connectivity. However, if the experimental data comes from non-imaging techniques such as mercury intrusion porosimetry and gas adsorption where not all pore space characteristics are readily available, regular pore network construction approach is usually applied with assumed connectivity (Jivkov and Olele, 2012).

6

7 To a large extent the success of pore network models depends on the way they represent 8 the real pore space in terms of its geometrical and topological characteristics for a given 9 application. Previous works have clearly demonstrated the importance of the geometric 10 properties of the porous media, such as the locations of pores and throats, the 11 distributions of sizes and shapes of pores and throats (Blunt et al., 1992; Knackstedt et 12 al., 1998; Oren, 1994)). Most early works on pore-scale modelling assumed that the 13 throats were cylinders with a circular cross-section or considered to be zero-volume 14 connections between pore chambers. Pores were either not modelled explicitly at all 15 (they simply represent throat junctions), or were spherical, cubical or cylindrical in 16 shape. These methods have been applied to study the convection, transient and steady-17 state diffusion, permeability, etc. (Bryntesson, 2002; Jivkov and Xiong, 2014; Laudone 18 et al., 2008; Meyers and Liapis, 1999; Meyers et al., 2001)). Recent advances have 19 allowed modelling a degree of irregularity in pore cross-sectional shape that was not 20 available in earlier PNM (Al-Gharbi and Blunt, 2005). For some applications, such as 21 multiphase flow, it may be necessary to pay specific attention to the local morphology 22 of the throats (Fenwick and Blunt, 1998; Gao et al., 2012; Hui and Blunt, 2000; Man 23 and Jing, 1999; Payatakes et al., 1973). For others, such as longer-scale permeability or 24 diffusivity predictions, the throats could be considered as straight channels with locally 25 averaged cross sections (De Josselin de Jong, 1958). In the latter case, the variability in 26 throat morphologies becomes of secondary importance and the transport is controlled 27 predominantly by the sizes and spatial positions of pores, and the connectivity of the 28 pore set via throats with different permittivity. In addition, the models have to reflect 29 basic topological properties, such as average pore coordination number and pore 30 coordination spectrum. The effects of the average pore coordination on transport 31 coefficients have been demonstrated in a number of works (Meyers and Liapis, 1999; 32 Raoof and Hassanizadeh, 2010a). The pore spectrum represents the relative number of

pores coordinated by different number of throats. Introducing the spectrum into the model poses a stronger constraint onto the construction than just utilising the average pore coordination number. This is due to the fact that totally different spectra could produce the same average coordination number and the effects of it have been demonstrated in a number of works (Jivkov et al., 2013; Meyers and Liapis, 1999).

6

7 Generally speaking there are three ways to construct a PNM representing a porous 8 medium. The first method is to create a statistically equivalent network using 9 distributions of basic morphologic parameters, while the second approach is to map a 10 network structure directly onto a specific porous medium void space. The fundamental 11 difference between the two methods is that the direct mapping provides a one-to one 12 spatial correspondence between the porous medium structure and the equivalent 13 network structure, whereas the other type of network is equivalent only in a statistical 14 sense to the modelled system. The last method is called the grain-based approach, which 15 is based on the diagenesis of porous media.

16

17 **3.1 Statistical reconstruction**

18

2D pore space images are routinely available at high resolution. 3D images can be 20 reconstructed using statistical methods with information obtained by analysing 2D thin 21 sections. Methods based on a truncated Gaussian random field are often used in 22 conjunction with the geometrical properties of the original pore space to reconstruct 3D 23 images (Adler and Thovert, 1998). These geometrical properties include porosity, also 24 called one-point correlation function, and two-point correlation function measuring the 25 probability of finding two points separated by a certain distance within the same phase.

26

However, the one and two-point correlations are insufficient to adequately replicate the
topology of the medium (Adler and Thovert, 1998; Ioannidis and Chatzis, 2000; Levitz,
1998; Roberts and Torquato, 1999; Yeong and Torquato, 1998a; Yeong and Torquato,
1998b). Yeong and Torquato (Yeong and Torquato, 1998a; Yeong and Torquato, 1998b)
used a combination of the two-point correlation function and the distribution of linear
path, which is the probability of finding a line segment with certain length in the void

1 space as a descriptor to characterize the pore geometry. In addition, Hilfer (Hilfer, 1991) 2 introduced local porosity distribution and local percolation probability to improve the 3 geometrical characterization. Other descriptors, such as pore chord length (the length in 4 the void between two solid voxels with a given direction) have proved useful in 5 characterising the structure and generating 3D images (Levitz, 1998; Roberts and 6 Torquato, 1999). The combination of one- and two-point correlation functions with 7 these geometrical descriptors improves the reconstruction of connectedness and 8 predictions of macroscopic properties such as permeability.

9

10 Despite this, these methods still fail to reproduce the long-range connectivity of the 11 original pore space. On the other hand, Yeong et al. (Yeong and Torquato, 1998a; 12 Yeong and Torquato, 1998b) developed a stochastic method based on simulated 13 annealing, which was later extended by Manwart et. al (Manwart et al., 2000). Their 14 approach is based on moving pore space voxels around to minimize the objective 15 function and they obtained the correct porosity. This method should be able to match 16 not only one- and two- point correlation functions but also multi- point correlation 17 functions.

18

Okabe and Blunt (Okabe and Blunt, 2004) developed a multi-point statistical method aiming to reconstruct the 3D volume from thin section images. The approach preserves the typical void space patterns observed in 2D and consequently preserves the longrange connectivity. These statistical methods discussed above produce 3D representation form 2D images of the porous media with similar morphological statistics.

25

26 3.2 Grain-based model

27

Bryant and co-workers pioneered the use of geologically realistic networks (Bryant and Blunt, 1992; Bryant et al., 1993a; Bryant et al., 1993b). Their models are based on random close packing of equally-sized spheres. They represented diagenesis by swelling the spheres uniformly and allowing them to overlap. Compaction was modelled by moving the centres of the spheres closer together in the vertical direction, again

1 allowing the spheres to overlap. Equivalent networks with a coordination number of 2 four or less were then constructed. Single and multiphase flow was simulated through 3 the pore space. They were able to predict the absolute and relative permeability, 4 capillary pressure, electrical and elastic properties of water-wet sand packs, sphere 5 packs and cemented quartz sandstone. They also reported a trend between permeability 6 and porosity for Fontainebleu sandstone. This represented a major triumph in pore-scale 7 modelling, since genuine predictions of transport and flow properties were made for the 8 first time. They showed that spatial correlations in the pore size distribution were 9 important for correct predictions: using the same pore size distribution, but assigning it 10 at random to the throats in the network gave erroneous predictions of permeability 11 (Bryant et al., 1993a). The main drawback with the work is its limited application—it 12 could only be applied to media that is predominantly composed of spherical grains of 13 the same size.

14

15 The next major advance came with the work of Oren, Bakke and co-workers at Statoil 16 (Bakke and Oren, 1997; Lerdahl et al., 2000; Oren and Bakke, 2002; Oren et al., 1998). 17 They developed a reconstruction method, where the packing of spheres of different size 18 was simulated. The grain size distribution was derived directly from analysis of thin 19 sections of the rock of interest. Compaction and diagenesis was modelled in a similar 20 manner to Bryant and co-workers (Bryant et al., 1993a). Clays were also included in the 21 model. Biswal et al. (Biswal et al., 1999) compared the pore space derived from this 22 geological reconstruction of Fontainebleau sandstone with an image obtained from 23 micro tomography. They also studied two stochastic models based on a correlation 24 function representation. It was shown that the stochastic models differed strongly from 25 the real sandstone in their connectivity properties. In contrast, the geological 26 reconstruction gave a good representation of the connectivity of the rock and as a 27 consequence could accurately predict transport properties (Oren and Bakke, 2002).

28

Oren and co-workers used the geological reconstructions to formulate topologically equivalent networks through which multiphase flow was simulated (Bakke and Oren, 1997; Oren et al., 1998). They predicted relative permeability for a variety of water-wet sandstones, showed promising results for a mixed-wet reservoir sample and matched

three-phase water-wet data (Lerdahl et al., 2000; Oren and Bakke, 2002). The works of
Pillotti, and Coehlo and co-workers have demonstrated a method to simulate the
deposition of grains of non-spherical shapes, in particular, enabling the reconstruction
technique to be applied more generally (Coelho et al., 1997; Pilotti, 2000).

5

6 There are three major concerns with the application of reconstruction methods to 7 reservoir samples. Firstly, the reconstruction algorithm is based on explicit simulation 8 of the geological processes by which the rock is formed. For many complex systems, 9 involving microporosity and clays as well as a variety of different sedimentary 10 processes, this may prove challenging. Furthermore, carbonate systems are not 11 modelled at all. Statistical reconstruction methods are more general, since they require 12 only a two-dimensional image of the system, and have been applied successfully to nonclassic rocks. However, their application so far, has been mainly to predict single-phase 13 14 flow simulated directly onto the pore-space reconstruction. The second problem, 15 regardless of the approach, is that characterisation of the pore space requires detailed 16 thin section analysis that might be unavailable or difficult to obtain. The final issue is 17 that the appropriate characterization of pore shape and wetability are not fully 18 understood.

19

20 **3.3 Direct mapping model**

21

Direct mapping from a real sample will yield an irregular lattice (Piri and Blunt, 2005). Such irregular networks allow for validating physical assumptions for flow simulations by comparison with 4D imaging of mass transport. 4D (3D + time) imaging methodology enables visualization and quantitative assessment of dynamic pore scale processes in real time over variable experimental durations. More details can be found at <u>http://www.solid-earth-discuss.net/se-2016-40</u>. Based on 3D images, the following approaches have been used to construct irregular pore network models.

29

30 3.2.1 Medial axis algorithm

1 The medial axis based methods transform the pore space images into a medial axis that 2 was the reduced representation of the pore space acting as a topological skeleton. The 3 topological skeleton is constructed roughly along the middle of the pore channels either 4 by a thinning algorithm (Baldwin et al., 1996; Liang et al., 1998) or a pore-space 5 burning algorithm (Lindquist et al., 1996). The algorithms developed by Lindquist et al. 6 (Lindquist et al., 2000) have been primarily applied to consolidated systems but it is 7 unknown whether they are applicable to unconsolidated porous media with higher 8 porosity. Al-Raoush et al. (Al-Raoush et al., 2003) developed algorithms which can be 9 applied to unconsolidated porous media and demonstrated their application on a variety 10 of unconsolidated porous media systems and the impacts of grain sizes, REV, and 11 image resolution (Al-Raoush and Willson, 2005). The pore throats are decided by local 12 minima along the branches whereas the pore bodies are determined at the nodes, which 13 verify pore space partitioning (Jiang et al., 2007; Lindquist et al., 1996; Prodanović et 14 al., 2006).

15

16 The medial axis mathematically preserves the topology of the pore space; it is difficult, 17 however, to identify pores unambiguously. Furthermore, pores normally encompass 18 more than one junction of the medial axis; therefore, various merging algorithms have 19 to be developed to trim the skeleton and fuse the junctions together while avoiding 20 unrealistically high coordination numbers (Jiang et al., 2007; Shin et al., 2005). The 21 choice of threshold value for throat quality can be problematic without being examined 22 by real flow simulations. In conclusion, medial axis algorithms readily capture the 23 interconnectivity of the pore spaces but encounter the problem in identifying pores.

24

25 3.2.2 Maximal ball algorithm

26

The maximal ball algorithm (Al-Kharusi and Blunt, 2007; Silin and Patzek, 2006) finds the largest inscribed spheres centred on each voxel of the image that just touch the grain or the boundary. Then those spheres, which are included in other spheres, are removed, while the rest are called maximal balls. The largest maximal balls usually identify pores, which are connected by smaller balls between them (these balls are called throats). Maximal balls were first used by Silin et al. (Silin and Patzek, 2006) to find the

1 dimensionless capillary pressure in drainage rather than to extract a pore network. The 2 method was then adopted and developed by Al-Kharusi and Blunt (Al-Kharusi and 3 Blunt, 2007). A more comprehensive set of criteria was set to determine the maximal 4 ball hierarchy including sphere clusters to handle equally sized balls. However, the 5 algorithms of Al-Kharusi and Blunt use a tremendous amount of computer memory and 6 time and consequently were limited to relatively small systems containing fewer than a 7 thousand pores. Moreover, their method tended to form pores with very high 8 coordination numbers. To overcome this issue, a two-step searching algorithm was 9 developed by Hu dong et al. (Dong and Blunt, 2009) based on Al-Kharusi and Blunt's 10 algorithm. The nearest solid was found to define a void bass instead of growing a ball 11 layer by layer. This algorithm invented a clustering process to define pores and throats 12 by affiliating the maximal balls into family trees according to their size and rank (Dong 13 and Blunt, 2009). This method clearly identified the larger pores, but tended to indentify 14 a cascade of smaller voids, which have sizes less than the image resolution.

15

16 One of the problems in network generation and modelling is the lack of a specific 17 definition for what constitutes pores and pore throats in real materials with complicated 18 pore geometries. Another problem is how the definition of pore morphology affects 19 network topology and discretization. Consequently, multiple network structures can be 20 created for the same material or 3D data set and there is no straightforward way to 21 define whether one network is more physically representative than another. Vogel 22 (Vogel, 2000) has demonstrated that different lattice-based topologies have similar 23 water retention curves. Inversely, Ams et al. (AMS et al., 2004) have demonstrated that 24 different network topology induce different relative permeability.

25

The effects of pore density (number of pores per unit volume) on transport in porous media have been studied by Bhattad et al (Bhattad et al., 2011). The large variations in pore density are the result of two factors: different methods for seeding the search for pores (i.e. defined as maximal inscribed spheres) and different settings in the merging criteria (that merges overlapping inscribed spheres to created single pores). Generally, the fast algorithms create low-density networks with little or no overlap in inscribed pores and therefore, can also generate longer pore–throats. More complete pore searches

are slow and result in higher pore densities and shorter pore-throats. The pore density
affects secondary parameters, such as pore-size distribution and pore interconnectivity.
However, the relationship is not always proportional or obvious. The most important
results have shown that single-phase permeability is relative insensitive to pore density,
while capillary pressure curves for quasi-static drainage are moderately sensitive to
network structure (Bhattad et al., 2011)

7

8 **3.4 Regular network model**

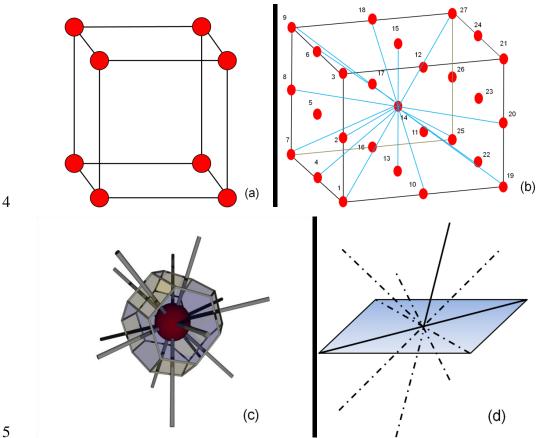
9

Irregular PNM are sample-specific and potentially not statistically representative. Direct construction of irregular network is impossible (especially for micro- and meso-porous materials) in the absence of data for throat locations and sizes, which cannot be observed by current imaging techniques. A regular PNM constructed within a larger volume allows for capturing statistics from a number of imaged samples, i.e. improved statistical correspondence, and for calculating transport at distances closer to the engineering length scale.

17

18 Early three-dimensional network models are usually based on cubic lattice with a 19 constant coordination number of six, Fig. 1(a) (Ioannidis and Chatzis, 1993b; Reeves 20 and Celia, 1996). However, in actual porous media, the connectivity is distributed and 21 can have coordination number greater than six (Dong and Blunt, 2009; Jivkov et al., 2013; Kwiecien et al., 1990). Therefore models with 26 pore coordination number have 22 23 been proposed, Fig. 1(b) (Raoof and Hassanizadeh, 2010a). This does not appear to be 24 physically realistic, however, because large numbers of throats intersect at points that 25 are not pores. A pore network construction based on the Kelvin solid has been recently 26 proposed and applied to permeability and diffusion analyses of porous media, Fig. 1(c) 27 (Jivkov et al., 2013; Xiong et al., 2014). This is based on truncated octahedral cells 28 providing variable pore coordination dependent on the allocation of pores and throats in 29 the cell, with maximum coordination of 14. Another application of a non-cubic lattice 30 support for pore-network construction uses rhombic dodecahedron as a unit cell, Fig. 31 1(d) (Vogel and Roth, 2001). This offers a maximum pore coordination of 12 with 32 throats of equal lengths, i.e. sufficiently large physically admissible coordination.

- 1 However, this cell is less representative for the volume around a pore than the truncated
- 2 octahedral cell.
- 3



5

6 Fig. 1 Unit cell shapes: (a) cubic lattice with a constant coordination number of six 7 (Ioannidis and Chatzis, 1993b; Reeves and Celia, 1996); (b) cubic lattice with 26 pore 8 coordination number (Raoof and Hassanizadeh, 2010a);(c) truncated octahedron (Jivkov 9 et al., 2013; Xiong et al., 2014); (d) rhombic dodecahedron (Vogel and Roth, 2001).

11 Historically, pores were related to the sites while pore throats were related to the bonds 12 of regular PNM (Dillard and Blunt, 2000; Meyers et al., 2001; Wilkinson and 13 Willemsen, 1983). If such correspondence is to be statistically representative of the 14 material modelled, sufficiently rich experimental information is required. This 15 information includes shape and size distribution of pores and throats, as well as the pore 16 coordination spectrum, i.e. percentages of pores coordinated by different numbers of 17 throats (Gao et al., 2012; Jivkov et al., 2013). These can be obtained in structures with 18 distinguishable pores and pore throats.

2 For porous media with macro-porosity, typically above 100 nm such as sandstone and 3 limestone, the throat and pore sizes distribution can be resolved as well as the 4 coordination spectra and the average coordination numbers can be calculated (Al-5 Raoush and Willson, 2005; Dong and Blunt, 2009). In such cases, topologically 6 representative networks can be constructed by a system of pores and throats connecting 7 all neighbouring pores, and subsequent elimination of throats to achieve the topological 8 constraints. This has been illustrated for average coordination number in simple bases, 9 such as cubic lattice (Raoof and Hassanizadeh, 2010a), as well as for full coordination 10 spectra in a bi-regular lattice (Jivkov et al., 2013).

11

12 However, for micro- and meso-porosity, e.g. below 100 nm, the pore size distribution 13 can usually be determined, but the resolution of the experimental techniques is not 14 sufficient to segment all the throats and calculate their sizes. For such cases, different 15 approaches to pore network construction are required. One possibility is to assume full 16 connectivity between neighbouring pores in a selected lattice, but control the diffusivity 17 of the throats through the sizes of the connected pores and the size of the solute 18 molecules (Xiong et al., 2014). This construction leads to a model length scale dictated 19 by the prescribed total porosity and the assumption that pores are located at each lattice 20 site. The approach was shown to provide insights into the effects of the structure on 21 diffusivity and the results correlated well with experimentally measured diffusion 22 coefficients (Xiong et al., 2014). This approach, however, is not sufficient to explain the 23 variability in reported experimental measurements of diffusion coefficients.

24

25 Clearly, the microstructure information used for the network construction is derived 26 from samples, which are different from the ones used for measurements of diffusion 27 coefficients. It is plausible, however, to assume that the statistical part of this 28 information, namely the distribution of pore sizes, is not vastly different between 29 samples of the same medium. This implies that there is an effect of an additional 30 geometrical parameter on the measured coefficients. It could be possible that a network 31 length scale that provides a constraint to the model rather than being determined from 32 total porosity and location of pores in all lattice sites (Jivkov and Xiong, 2014).

2 The two approaches that tackle incomplete pore space information, namely predefined 3 connectivity for calculable length scale (Xiong et al., 2014), and undetermined length 4 scale for realistic connectivity (Jivkov and Xiong, 2014), suffer from the lack of an 5 additional constraint. This cannot be found from the pore space information and 6 requires the consideration of the solid phase structure, e.g. the shape and size 7 distribution of mineral grains. A methodology for incorporating the structure of the 8 solid phase was then proposed which improves substantially the realism of the 9 constructed pore network, both in terms of geometry and topology (Xiong and Jivkov, 10 2015). The proposed methodology is conceptually different from previous works and it 11 has the added benefit of constructing a PNM which can be paired directly to lattice 12 models of the solid phase, developed for analysis of damage evolution via micro-13 cracking (Jivkov et al., 2012; Zhang and Jivkov, 2014). The approach is conceptually 14 similar to the dual models of pore spaces introduced by Glantz et al. (Glantz and Hilpert, 15 2007; Glantz and Hilpert, 2008). This facilitates mechanistic investigations of combined 16 mechanical-thermal-chemical-biological effects on diffusivity.

17

18 **3.5 Two-scale pore network models**

19

20 Despite the advances of pore scale modelling, the simulation of transport properties in 21 heterogeneous rocks still remains an open issue due to the very broad pore size 22 distributions in some porous media. Such examples could be of considerable scientific 23 and economic importance including many types of carbonates, clay-rich sandstones and 24 clays. For example, carbonate reservoirs or tight gas sandstones are very important for 25 hydrocarbon extraction and CO_2 sequestration. The failure of simulating most classical 26 empirical relations (e.g. the Brooks-Corey parameterization of relative permeability and 27 Archie's law for electrical behaviour) in these rocks due to the interaction of micro-and 28 macro-porosity in turn has promoted the development of modelling methods where 29 multiple pore scales are coupled (Mehmani and Prodanović, 2014). A short overview of 30 existing two-scale pore network models is presented here and illustrated in Fig. 2.

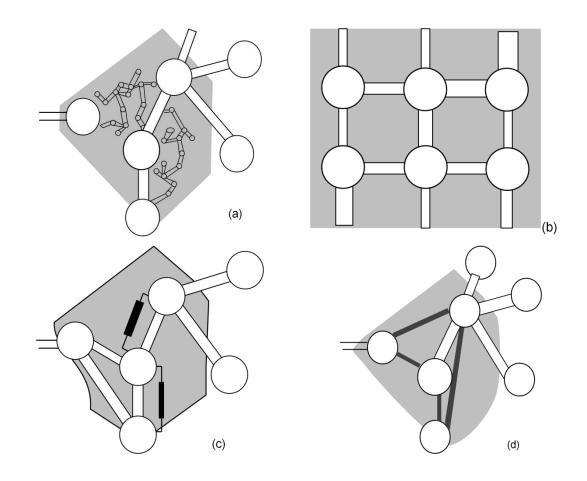


Fig. 2. Schematic depiction of existing DPNM methods. In all figures, light grey represents a zone of microporosity. (a) a model with individual micropores (Jiang et al., 2013; Prodanović et al., 2015). (b) a lattice-based model with upscaled microporosity properties (Bekri et al., 2005). (c) a model which takes upscaled microporosity in parallel with macro-throats (dark grey) into account (Bauer et al., 2012). (d) a model with connectivity added by the (upscaled) microporosity (Bultreys et al., 2015).

8

9 Jiang et al. (Jiang et al., 2013) developed a methodology to integrate networks extracted 10 from images at distinct length scales. The pore network model was generated at each 11 scale and then was integrated into a single two-scale network by characterizing the 12 cross-scale connection structure between the two networks. The shortcoming of this 13 method is that it is computationally costly due to the number of network elements 14 (Bultreys et al., 2015). Recognizing the computational problems when single micro-15 pores are taken into account, Mehmani and Prodanović (Mehmani and Prodanović, 16 2014) proposed a two-scale pore network by the use of packing algorithms. The macro-17 network is constructed by Delaunay tessellation of the grain centres. Micro-porous

1 networks are generated by downscaling existing networks extracted from macro-pores. 2 This approach was capable of investigating fundamental two-phase flow properties of 3 multi-scale porous media and micro-porosity was able to act in series (intergranular or 4 pore-filling micro-porosity) and in parallel (intra-granular or dissolution micro-porosity) 5 to the macro-pores. However, in the construction process, distorted pores were produced 6 when many small grains touched a large grain. A parameter, the ratio of the macro to 7 micro length scale, needs to be determined for micro-porous regions from the 8 measurements, which may be difficult to perform. Bultreys et al. (Bultreys et al., 2015) 9 developed a workflow to integrate networks of macro-pores and micro-porous regions 10 extracted from micro-CT images. This methodology allowed micro-porosity to act both 11 in parallel and in series with the macro-pore network. However, a representative 12 network for the micro-porosity is necessary. In addition, the pore networks from Jiang 13 et al. (Jiang et al., 2013) and Boultry et al. (Bultreys et al., 2015) were based on 14 experimental data from micro-CT images which did not take into account the micro-15 pores that cannot be resolved by micro-CT. As the truncated cone shape is used to 16 connect two neighbouring macro-pores, the tortuosity of the connection and geometric 17 details about the bulk of the micro-porous cluster are neglected, which can lead to 18 erroneous local conductivities.

19

20 **4.** Applications

Pore network models have been used to explore reactive transport processes, such as: adsorption (Acharya et al., 2005; Raoof et al., 2013), dissolution/precipitation (Kang et al., 2010; Knutson et al., 2001; Mehmani et al., 2012; Varloteaux et al., 2013; Zhou et al., 2000), and biomass growth (Dupin et al., 2001a; Gharasoo et al., 2012; Rosenzweig et al., 2013). In this section, we review some representative applications of the PNM in these areas.

27

28 4.1 Adsorption

29

30 PNM has been widely used to study adsorption in porous media which finds its 31 relevance in many areas of science and engineering including radioactive waste disposal (Xiong et al., 2015a; Xiong et al., 2014) and chromatographic separation processes
 (Câmara and Silva Neto, 2009; Meyers and Liapis, 1999).

3

There are two principal ways to simulate adsorption in PNM. The first one entails decoupling of adsorption and transport into two separate processes by the use of parameters associated with each process calculated separately. Firstly, the convectiondiffusion problem is solved to obtain the concentration and flow field. This is governed by the following equations:

9
$$q_{ij} = \frac{A_{ij}^2 G_{ij}}{\mu l_{ij}} \left(p_i - p_j \right)$$
(1)

10
$$V_{ij}\frac{dc_{ij}}{dt} = q_{ij}(c_i - c_j) + D_{ij}A_{ij}\frac{c_i - c_j}{l_{ij}}$$
(2)

11 where Eq. (1) is the Hagen- Poiseuille law, appropriate for describing flow in pores 12 (Jacob, 1972) and valid for laminar flow; and Eq. (2) is the classical convection-13 diffusion law. The equations are written for transport through a throat connecting pores *i* and *j*, with length l_{ii} [L], and cross section area A_{ii} [L²]. In Eq. (1) q_{ii} [L³T⁻¹] is the total 14 volumetric flow rate through the throat, p_i and p_i [ML⁻¹T⁻²] are pressures at pores *i* and *j*, 15 respectively, μ [ML⁻¹T⁻¹] is the dynamic viscosity, and G_{ii} is a cross-section shape 16 coefficient: for circular, equilateral triangle and square tube, G_{ii} is 0.5, 0.6 and 0.5623 17 respectively (Patzek and Silin, 2001). In Eq. (2) V_{ii} [L³] is the throat volume, c_i and c_i 18 [ML⁻³] are concentrations at pores *i* and *j*, respectively, and D_{ii} [L²T⁻¹] is the throat 19 20 diffusivity.

21

Then, a sorption isotherm, which is assumed to belong to some class of parameterized functions (Langmuir, Freundlich, Henry, etc.) with few tuning parameters, is used to describe the adsorption of species onto the walls. A coarse grained mathematical formula containing the sorption isotherm is used to reflect the effects of sorption on the change of pore sizes. For example, the adsorbate obstruction can be simulated by reducing the throats radii as a function of the concentration of the moving species (Meyers and Liapis, 1999; Xiong et al., 2015a; Xiong et al., 2014)

29
$$t^{2} - 2tR + 8\theta Rr_{M}/3 = 0$$
 (3)

1 where t is the thickness of the adsorbed layer, R is the throat radius, r_M is the radius of the diffusing species, and θ is the adsorption isotherm (e.g. $\theta_{ij} = k_{d,ij}c_{ij}$, $\theta_{ij} = ac_{ij}^{b}$). Thus, 2 the radius of a throat after adsorption at given concentration, C_A , and hence radius after 3 adsorption, becomes $R^* = R - t$. Thus, the sorption is accounted for by modifying the 4 5 pore geometries according to the changed pore sizes. This procedure is iterated until 6 reaching a certain criteria of interest. Using this method, various types of adsorption 7 reactions can be studied, including linear equilibrium, nonlinear equilibrium (Meyers 8 and Liapis, 1999; Xiong et al., 2014) and heterogeneous adsorption in which the 9 adsorption parameters are spatially varying (Xiong et al., 2015a).

10

11 Sorption isotherms in the published literature, however, are determined from bulk 12 measurements and may not be appropriate for occluded regions within the pore space. 13 Adsorption takes place at the water-solid interface and is thus controlled by solute 14 concentration in the fluid next to the solid grain in individual pores. The concentration 15 near the solid surface is different from the average concentration within the pore due to 16 the development of concentration gradients within the pore space. Simulations with 17 pore network models are usually based on the average concentration in individual pore 18 throats. An alternative way is to simulate the adsorption in each pore element and then 19 scale it up to the whole pore network. This has been demonstrated by Raoof and 20 Hassanizadeh (Raoof and Hassanizadeh, 2010b). They have performed pore-scale 21 simulations for a wide range of local-scale distribution coefficients and Peclet numbers. 22 These have provided relationships between the upscaled parameters and underlying 23 pore-scale parameters. Such relations are useful for upscaling by means of a pore-24 network model. Raoof and Hassanizadeh have shown also that even if there were 25 equilibrium adsorption at the pore wall (i.e. at the grain surface), one might need to 26 employ a kinetic description at larger scales. They have also shown this kinetic 27 behaviour through a volume averaging method, yielding very similar results for the 28 upscaled kinetic parameters. From upscaling, they have found a negligible dependency 29 of effective adsorption rate coefficients on pore water velocity. This finding is in 30 agreement with Zhang et al. (Zhang et al., 2008) who has observed that upscaled 31 sorption parameters are independent of pore-water velocity. The findings, however, 32 contradict some experimental results, where kinetic adsorption-desorption coefficients

have been shown to depended on velocity (Akratanakul et al., 1983; Brusseau, 1992;
 Maraqa et al., 1999; Pang et al., 2002).

3

4 The second principal method for simulating adsorption in PNM is by solving the 5 advection-diffusion-sorption equation directly. The adsorption reactions includes linear 6 equilibrium (Raoof et al., 2013) and nonlinear equilibrium (Acharya et al., 2005; Xiong 7 et al., 2015b). The equation can be expressed as

8
$$V_{ij}\frac{dc_{ij}}{dt} = V_{ij}u_{ij}\frac{c_i - c_j}{l_{ij}} + F_{ij} - S_{ij}\frac{ds_{ij}}{dt}$$
(4)

9 where $V_{ij}u_{ij}\frac{c_i-c_j}{l_{ij}}$ is the advection term which represents the effects of velocity u_{ij} on

10 particle transport, $F_{ij} = A_{ij}D_{ij}\frac{d^2c_{ij}}{dx^2}$ is the diffusive term, $s_{ij} = s_{ij}$ (c_{ij}) [ML⁻²] is mass 11 adsorbed per unit area of throat wall, S_{ij} [L²] is the throat surface area.

12

13 **4.2 Dissolution and precipitation**

14

15 Modelling of dissolution/precipitation is of high interest to applications such as CO₂ 16 geological storage, enhanced oil recovery and chemical degradation of the cement. For 17 example, CO_2 storage is one possible method to reduce atmospheric emissions of CO_2 , 18 and hence mitigates climate change. Some CO₂ may dissolve in the resident brine and 19 generates a weak acid after injection into the low-permeability caprock. The acidic CO₂-20 rich brine reacts with the host rock to produce solid carbonate which leads to 21 precipitation process (Blunt et al., 2013). In enhanced oil recovery, strong acid has been 22 injected near the well to dissolve the host rock and increase the reservoir permeability 23 (Algive et al., 2010). Accurately computing dissolution and precipitation processes 24 involve the development of advection, diffusion and reaction equations. In addition, for 25 better estimation of the abovementioned process the following also need to be 26 accounted for: the heterogeneities in the porous media (e.g. porosity, mineralogical and 27 chemical compositions), geometrical changes (i.e. pore and throat size changes), 28 topological changes (i.e. increase/decrease pore to pore connection) and the chemical 29 gradient filed inside pores. This is due to the fact that mineral dissolution and 1 precipitation reactions alter the geometry and topology of the porous media (Crandell et

2 al., 2012).

3

4 Li et al. (Li et al., 2006) used a pore-scale network model to simulate reaction of 5 kaolinite and anorthite during carbon sequestration. The processes of diffusion, 6 advection, aqueous reactions, dissolution and precipitation were simulated in individual 7 pores. The information was then summarised and averaged over the network to obtain 8 descriptions of processes at the continuum scale. As previously observed, they showed 9 qualitatively that in situ reaction rates were significantly different from values found in 10 batch lab reactors. Their network models, however, were 3D regular lattices and not 11 physically representative of the real media, which made quantitative analysis 12 challenging. Li et al. (Li et al., 2007)further investigated the influence of the spatial 13 distribution of mineral species using hypothetical networks, while Kim et al. (Kim et al., 14 2011a) used network models mapped from X-ray computed tomography(XCMT) of real 15 sandstones. They were able to determine the initial concentrations of minerals in the 16 porous medium from the XCMT images, which allowed for more predictive modelling. 17 Kim et al. (Kim et al., 2011b) studied the relationship of reaction rates with kinetic rate 18 and bulk flow rate based on previous work (Li et al., 2007; Li et al., 2006). The 19 reaction rates were found dependent on flow rates with an approximately power law 20 relationship. They noted that their models were limited by the size of the network model. 21 Notably, the effects of decreased throat conductivity as a result of precipitation were not 22 accounted for in their work. Thus the physical change in porosity is lacking in the 23 network flow model due to surface minerals dissolving or precipitating. These models, 24 therefore, do not capture the important topological changes that reactions produce and 25 the impact that such changes may induce on upscaled reaction rates. Algorithms for 26 incorporating physical change, due to dissolution and precipitation, need to be included 27 in the models. Based on these limitations, Kim et al. (Kim and Lindquist, 2013) further 28 extended the work to consider geometric changes (e.g. pore volume and reactive surface 29 area) and topological changes, which opened new connections between porosity 30 features. When geometrical changes are included, the steady state, core-scale reaction 31 rates could not be directly used as input into larger-scale simulations, as reactions are 32 time dependent. The Nernst-Planck (N-P) term in the governing transport equations,

which ensures the electrical neutrality when species of different charges diffuse at different rates, was also taken into account. The agreement between simulations with and without the N-P term improved as the flow rate decreased. In these studies, reactions far from equilibrium, such as anorthite reaction, and reactions close to equilibrium conditions, such as the kaolinite reaction, were considered.

6

7 Algive et al. (Algive et al., 2010) used a pore-scale model to simulate reactive transport 8 which caused precipitation in pores. The reactive transport and structural modifications 9 were treated separately and updated step by step. The reactive transport was solved in 10 each element of the pore network. The mean concentration in each pore was used to 11 determine the effective parameters of the macroscopic reactive transport equation. The 12 different pore geometries were also considered for reactive transport by using random 13 walk techniques. Then the reactive transport equation was solved analytically between 14 the nodes of the regular lattice network for the asymptotic regime and computed the 15 node concentration by a matrix inversion. Moreover, the system was based on the 16 assumption that it had nearly reached chemical equilibrium and mainly focused on 17 regular or periodic networks in one or two dimensions or systems, where pore-to-pore 18 heterogeneity was not taken into account.

19

20 Algive et al. (2012) developed a methodology to use a reactive pore network model to 21 extract upscaling factors to tie the pore-scale effects of reactive transport to the core-22 scale values of permeability and porosity. Their work simplified the geochemistry and 23 transport by incorporating most geochemical dynamics and transport into dimensionless 24 variables and in doing so prevented a full description of the different physical and 25 chemical parameters. Mehmani et al. (Mehmani et al., 2012) developed a novel 26 approach that coupled several pore-scale models using mortar coupling domain 27 decomposition to study the evolution of precipitation-induced cementation of calcite. 28 They were able to study large changes in permeability and porosity by coupling of 64 29 pore-scale models (1 mm \times 1 mm \times 1 mm each) but with a limited description of the 30 chemistry represented by two main parameters, the Damkohler number (D_a) and an 31 "alpha" parameter which described the deviation from equilibrium of the precipitation 32 reaction. Nogues et al. (Nogues et al., 2013) simulated carbonic acid-driven reactions in

1 a 3-D network of pores to predict the changes in permeability and porosity at continuum 2 scale. The pore network structure was based on a statistical characterization of a 3 synthetic microcomputed-tomography image of an oolithic dolostone. Only physical 4 heterogeneities were considered in comparison to the work done by Kim et al. (Kim et 5 al., 2013). To account for the changes caused by dissolution and precipitation, a 6 mathematical construct was developed to modify the pore-pore conductivities by 7 relating these changes to changes in pore volumes. It is important to note that while the 8 surface area fractions assigned to different minerals do change, the model does not 9 change the total surface area. Because the reactions are modelled at equilibrium, the 10 reaction rate is independent of surface area. Thus the simulation results are negligibly 11 affected in this situation. This model analysed the effects of different chemical, flow 12 and mixing conditions on permeability.

13

14 Raoof et al. (Raoof et al., 2013; Raoof et al., 2012) employed a coupled Complex Pore 15 Network Model (CPNM) and Biogeochemical Reaction Network Simulator (BRNS) to 16 solve reactive transport by pore networks. CPNM reproduced the microstructure of real 17 porous media and calculated solute concentration values for both pore bodies and pore 18 throats within the pore network. BRNS (Regnier et al., 2002) was responsible for 19 handling a comprehensive suite of multi-component complexion, mineral precipitation 20 and dissolution reactions, as well as reaction networks characterized by multiple kinetic 21 pathways. The evolution of porosity and permeability was discussed in this work as 22 well. In BRNS the processes of dissolution/precipitation were simulated by changes of 23 radius of pores and throats, while the effects of opening/clogging pores/throats were not 24 considered.

25

All modelling approaches mentioned above involve the assumption that mass transfer within the pore is not limited, and that the regular geometric solid-fluid interface and the aqueous concentrations are spatially uniform within a pore. Based on that, the pore network models do not resolve chemical gradients within pores and also for each single pore the average values for physical/chemical properties were used. The accuracy of this assumption in computing reactive flow requires further investigation.

1 **4.3 Biomass growth**

2

3 Biomass growth within porous media is of particular interest in a variety of industrial 4 scenarios such as microbial enhanced oil recovery (Ezeuko et al., 2011), bioremediation 5 of contaminated soil and water (Cunningham et al., 1991) and nuclear waste disposal. 6 As the mass of the microorganisms grows, hydrodynamic properties of porous media 7 change (Rolland du Roscoat et al., 2014; Stewart and Scott Fogler, 2002; Taylor and 8 Jaff é 1990). The relevance of PNM to biofilm growth modelling comes from the fact 9 that the biomass is not homogeneously dispersed within the porous media and hence 10 continuum models fail to adequately represent this pore-scale phenomena and its effect 11 on the macroscopic properties of the media.

12

13 The morphology of biomass growth in porous media depends strongly on the structure 14 of the porous medium, bacteria species and the prevailing hydrodynamic and nutritional 15 conditions. Wimpenny et al presented a good review of the factors affecting biofilm 16 formation (Wimpenny et al., 2000). Three principal morphologies have been reported in 17 the literature: flat and uniformly thick biofilms, mushroom-shaped structures and 18 streamers (Stoodley et al., 1999). Only the first two, however, have significant 19 relevance to porous media since streamers generally occur at turbulent flow conditions. 20 The most commonly modelled growth mode is a continuous layer on the surface of the 21 soil grains (Suchomel et al., 1998). Sometimes, however, the biomass does not cover 22 the entire surface of the soil grains, and thus microcolonies, i.e., a patchy biofilm, 23 develop (Molz et al., 1986). In other cases, biomass grows in pores in the form of 24 aggregates (Vandevivere and Baveye, 1992). When fungal growth occurs, mycelia can 25 develop throughout the entire pore volume and envelop several soil grains (Dupin and 26 McCarty, 2000). These forms of biomass growth can occur simultaneously (Dupin and 27 McCarty, 2000). Biofilm morphology is important because it results in different 28 clogging mechanisms. Rittmann (Rittmann, 1993) noted that the difference between flat 29 biofilms, and aggregate growth is crucial to model permeability reduction. It takes less 30 biomass to plug a pore if it forms aggregates as opposed to continuous biofilms. 31 Therefore, mushroom-shaped colonies have the potential to generate larger permeability 32 reductions in comparison to the uniform biofilm morphology (Thullner, 2010).

2 The development of robust methods to engineer biofilms in porous media requires 3 predictive mathematical models capable of determining the evolution of biofilms under 4 different flow conditions. However, reported studies of biofilm evolution in micro 5 models have revealed that the structural complexity of biofilms entrapped in a complex 6 pore space makes accurate modelling difficult (Dupin and McCarty, 2000; Kim and 7 Fogler, 2000). Several pore level models for biofilm growth under flow conditions have 8 been developed by adopting different approaches. They include lattice Boltzmann-based 9 models (Pintelon et al., 2009; Pintelon et al., 2012), bacterial cell level individual-based 10 models (Kreft et al., 2001; Picioreanu et al., 2004), and Pore Network Models (PNMs) 11 (Kim and Fogler, 2000; Stewart and Kim, 2004; Suchomel et al., 1998; Thullner and 12 Baveye, 2008). Ezeuko et al. modelled biofilm evolution in porous media using a new PNM that includes hydrodynamics and nutrient transport based on coupling of 13 14 advective transport with Fickian diffusion, and a reaction term to account for nutrient 15 consumption. Specifically, the PNM is used to examine the influence of biofilm 16 formation on the hydraulic properties of porous media. In PNM models, pore 17 geometries can range from a simple circular cylinder through to complicated voxel-by-18 voxel pore reconstructions from thin sections or micro-tomographical data. The PNM 19 approach facilitates simulations of important processes by adopting simplified but 20 realistic pore geometries while allowing implementations of other pore level physics 21 and kinetics relevant to the process of interest (Ezeuko et al., 2011).

22

23 Or and Tuller (Or and Tuller, 2000) used the classic capillary tube model to represent 24 the soil pores (i.e., the pores are simulated as a bundle of capillary tubes of different 25 diameters) and studied flow in unsaturated fractured porous media. The analysis showed 26 that the biofilm spatial distribution within the pore space has a significant effect on the 27 hydraulic properties of the soil. Although the capillary model is widely used and is easy 28 to obtain, it has a few shortcomings, primarily the over simplistic assumption regarding 29 the binary nature of the pores (i.e., each pore is either completely water-filled or 30 completely empty under a prescribed matric head). Moreover, the lack of pores 31 connectivity and the unrealistic cylindrical geometry of the capillaries further reduce the 32 model efficiency in determining soil hydraulic properties (Or and Tuller, 2000). Or and

1 Tuller (Or and Tuller, 2000) suggested to model the soil-pores by using channels having 2 a triangular cross-section. The channels can be either water-filled (for saturated flow) or 3 partially-filled (for unsaturated flow), where the water-filled fraction is related to the 4 matric head through the curvature of the water surface. The lack of pore connectivity 5 can be especially significant when dealing with biofilms. Thullner et al., (Thullner et al., 6 2004; Thullner et al., 2002) showed that under saturated conditions, modelling the pores 7 as one-dimensional capillary tubes could not reproduce connectivity-related effects. 8 These effects, such as flow bypasses generated from preferential biofilm growth, were 9 demonstrated in experiments and were successfully reproduced using pore-network 10 models. Rosenzweig et al. extended Or and Tuller's (Or and Tuller, 2000) concept and 11 propose a framework in which the soil pores are simulated as a network of channels 12 having a triangular cross-section. The representation of the soil pores by a network of 13 triangular channels is superior to the capillary model in two aspects. It eliminates the 14 binary (full/empty) behaviour and takes into account inter-pore connectivity. 15 Furthermore, it provides basis on simulating the combined water flow, solute transport 16 and biofilm growth under variably saturated conditions (Rosenzweig et al., 2013).

17

18 The coupled water flow, substrate transport, and biofilm growth under saturated 19 conditions has been studied by several researchers by using pore network models 20 (Dupin et al., 2001b; Kim and Fogler, 2000; Suchomel et al., 1998; Thullner and 21 Baveye, 2008; Thullner et al., 2002). To study the effects of biofilm on porous media, 22 the pore network models were built either of lattices of cylindrical tubes (Kim and 23 Fogler, 2000; Suchomel et al., 1998; Thullner and Baveye, 2008), or two-dimensional 24 channels (Dupin et al., 2001b). The biofilm was modelled as either a continuous layer 25 coating pore walls (Suchomel et al., 1998; Thullner and Baveye, 2008; Thullner et al., 26 2002) or as isolated microcolonies (Dupin et al., 2001b; Thullner et al., 2002). As 27 biomass growth depends on the environmental conditions and the type of bacteria, it is 28 problematic to represent this accumulation appropriately in terms of its effect on 29 porosity and permeability reductions. No universally applicable model has been 30 developed for saturated porous media systems. Moreover, no attempts have been made 31 to develop a model that is applicable to both saturated and unsaturated systems. This 32 situation could be particularly relevant to conditions that trigger gas generation as a

1 result of bacterial metabolic processes. Gas bubble formation and entrapment could also 2 severely block the pore space. It is usually characteristic of denitrifying and 3 methanogenic conditions where nitrogen (Soares et al., 1991) and methane gas 4 (Beckwith and Baird, 2001) production is observed but not exclusively as CO_2 and other gases are produced as part of the metabolic processes of different species. 5 6 Unsaturated conditions are crucial for understanding the flow and transport in 7 environments such as groundwater recharge basins, bioremediation sites, and effluent 8 irrigated fields. However, very few pore network models have been used to study the 9 biofilm morphology under unsaturated conditions (Rosenzweig et al., 2013). 10 Rosenzweig et al. proposed a pore network model of triangular channels to simulate 11 variably saturated flows in biofilm-affected porous media. The effects of biofilm on the 12 hydraulic properties of the network are demonstrated by examining three scenarios for 13 the biofilm pore-scale morphology including plug (or microcolony) morphology and a 14 uniform biofilm layer (Rosenzweig et al., 2013). The simulations demonstrated that the 15 effect of biofilms on the hydraulic properties of the network is a complicated and 16 nonlinear function that depends not only on the biofilm scenario but also on the 17 saturation.

18

19 A significant assumption in most existing biofilm simulation models is that biofilm-20 occupied regions are impermeable and that nutrient provision within the biofilm is 21 purely as a result of diffusion. However, experimental studies have demonstrated that 22 biofilm morphology is often extremely heterogeneous and can contain voids (e.g., 23 (Flemming et al., 2000)). As a result, biofilms are not fully impermeable and can 24 contain a significant amount of both static and dynamic water (Flemming et al., 2000). 25 In the absence of biofilm detachment, the effect of permeability is only relevant when 26 biofilm growth is nutrient mass transfer limited. Such conditions are prominent in 27 subsurface conditions (Kim and Fogler, 2000) and are very relevant to bio-barrier 28 scenarios where thick biofilms are desirable for permeability reduction. Thullner et al. studied the effects of permeability of bioflims on the hydraulic conductivity of porous 29 30 media by pore network models (Thullner and Baveye, 2008). The computational pore 31 network model consisted of cylindrical pores which were arranged in an orthogonal 32 array in two or three dimensions. The nodes connecting adjacent pores are assumed to 1 be volumeless. The biofilm is viewed as homogeneous with regard to its hydraulic and 2 biological characteristics in the model. In each pore, microbial activity is assumed to 3 take place exclusively within the biofilm. This model was able to simulate reductions of 4 the hydraulic conductivity of the pore networks by two to three orders of magnitudes. 5 These reductions are in good agreement with laboratory experiments with san columns 6 or field situations (Thullner and Baveye, 2008). The amount of biomass in any given 7 pore is controlled by biomass growth, decay processes inside the biofilm, and 8 detachment of biomass due to shear forces. However, few pore network models 9 simulate all the biomass growth or surface attachment mechanisms.

10

11 **5. Discussion and conclusions**

12

The paper reviewed three aspects of pore network modelling: main experimental 13 14 techniques for characterisation of pore spaces; main methods for construction of pore 15 networks reflecting available pore space characteristics; and application of pore network 16 models to analysis of phenomena, which have not been covered in previous reviews -17 sorption, dissolution and biofilm growth - but have increasing technological and 18 scientific importance. The works reviewed demonstrate that pore network modelling is 19 an efficient and useful approach to analysis of reactive transport at the meso-scale, less 20 computationally demanding than direct methods and able to incorporate heterogeneity 21 within larger rock volumes.

22

23 In many practical cases, the main source of uncertainty in pore space characterisation is 24 the finite resolution of the imaging techniques. These contain a hierarchy of pores from 25 nanometres to micrometres. The uncertainties could arise from hardware limitations or 26 from compromising resolution to capture heterogeneities. The errors may arise from the 27 relatively sparse numerical grid of the imaging techniques, whereby details of the 28 structure, e.g. very narrow passages in the porous space and surface roughness of fibres, may not be reproduced by the images. Furthermore, the results for each individual 29 30 sample may be affected by the methods used to minimise or remove noise and other 31 artefacts present in the original reconstructed images. These are the main factors that 32 currently limit the applicability of the imaging techniques especially for fine-structured 1 materials. The limitation is particularly relevant to materials with tight pore spaces, such 2 as clays and cement-based materials. For example in concrete, the so-called air porosity 3 can be resolved by existing methods, but appears as disconnected set of pores. The 4 transport pathways are formed by the gel porosity of the cement, as well as along the so-5 called interfacial transitions zones between cement and aggregates, the connectivity of 6 which is not resolvable with the current techniques. Similar is the situation with 7 compacted clays. In such cases, the connectivity of the transport pathways needs to be 8 assumed and tested against macroscopic transport experiments until suitable 9 connectivity is found (Xiong et al., 2015).

10

11 In addition to numerical uncertainties arising from discretisation and image quality 12 discussed above, this deviation includes statistical variation arising from relatively small 13 size and low number of samples used in computation. In principle, the accuracy can be 14 improved by using higher resolution images and better statistics. The former condition 15 may become feasible with the present rapid development of imaging techniques. The 16 latter can be obtained by larger size or larger number of samples. Clearly, these 17 improvements come with the cost of higher computational effort, which, however, does 18 not appear critical from the point of view of applicability of the method. Another 19 approach for overcoming the above mentioned problems is to use a combination of 20 different techniques such as FIB/SEM, nitrogen adsorption and FIB (Keller et al., 2013). 21

22 It must be noted that pore network modelling is conceptually scale indifferent, i.e. the 23 approach can be applied to any length interval, where the structure of the pore space has 24 been experimentally observed and analysed. For example if a particular experimental 25 technique allows for characterising pore features of sizes between 0.1µm and 10µm, 26 then the corresponding PNM will take these into account. The elements of PNM, sites 27 and bonds, are abstract and can be related to the measurable features in different ways 28 depending on the available information. The individual pore networks can be extracted 29 at each length scale and then integrated into a single multi-scale network by 30 characterizing the cross-scale connection structure between these networks (Bultreys et 31 al., 2015; Jiang et al., 2013).

1 Pore network models have been proven to be an effective research tool to upscale 2 reactive transport processes from the local scale to the continuum scale. The 3 microscopic dynamics can be understood best at the scale of individual pores. It is 4 challenging to directly correlate results obtained from pore scale simulations to improve 5 quantitative predictions based on large-scale simulations. However, the results obtained 6 from pore scale investigations contribute to the understanding of large-scale natural 7 processes and improve large scale geotechnical applications. It is well known that some 8 large scale reactive process may be difficult to predict or explain unless there is a well 9 defined experiment that has been used to study it precisely. Pore network models allow 10 for bridging the gap between laboratory simulations and larger scale transport. Although 11 pore scale is difficult to evaluate with traditional experimental studies, detailed 12 investigations allow for the identification of important pore-structure details that are 13 essential and control the reactive transport processes at the engineering/continuum scale. 14

15 It should be mentioned that, in spite of the attractiveness and successful applications of 16 pore network models, as any other numerical method they have certain limitations. First 17 of all, successful simulation of transport requires adequate representation of the real 18 porous media. Pore network models usually simplify the geometry of the pore space and 19 sometimes cannot capture all characteristics due to lack of resolution in most 20 simulations. Whether the pore networks could reasonably predict single and multi phase 21 properties depends on whether all relevant details are identified and included. In many 22 systems, the very small pores below the image resolution make a little contribution to 23 the overall behaviour. In these cases, transport properties are well predicted when the 24 principal connected pore space is modelled, while the smaller pores could be neglected. 25 Otherwise, assumptions are required. Secondly, the detailed microscopic features, such 26 as the exact molecular structures and intermolecular interactions cannot be represented 27 explicitly in pore network models. Therefore, in such situations, many approaches have 28 been developed to study upscaling reactive transport in porous media (Algive et al., 29 2010; Li et al., 2006; Varloteaux et al., 2013). Reactive transport processes are 30 simulated at the pore scale, which allows for incorporation of physical and chemical 31 effects' rate variations from pore to pore. The reactive transport is then upscaled from 32 the pore scale to a heterogeneous porous medium with the aid of a sufficiently large

1 three-dimensional pore network model. The PNM is used to determine the 2 phenomenological transport coefficients and the porosity/permeability relationship, 3 which can be used as inputs in reservoir scale simulation (Algive et al., 2010; Li et al., 4 2006; Varloteaux et al., 2013). The strength of the pore network models is that they can 5 be developed elegantly to study the effect of pore space changing mechanisms. In most 6 simulations of pore-scale mass transport the pore geometry is assumed to be constant, 7 apart from the effects of reaction. However, in many important geo-systems the 8 confining geometry changes slowly as a result of processes such as pressure solution 9 (Tada and Siever, 1989) or rapidly because of fracturing and changes in pore pressure or 10 matrix stress. It is important to include the coupling between diffusion and geo-11 mechanical processes for a number of important applications including oil recovery and 12 carbon dioxide sequestration.

13

14 Modelling and simulation of mineral precipitation and dissolution on the pore-scale 15 have not been subjects to extensive research in the geo-science community. In particular, 16 most simulations have employed very simple models for dissolution and/or precipitation 17 chemistry, and there has been little, if any, work on chemical processes coupled with 18 multiphase diffusion. It is believed that this situation is likely to change rapidly. There 19 is strong motivation for such simulations for both scientific discovery and practical 20 applications. Furthermore, the simulated results are usually compared with experimental 21 data, which will not give the same results realistically. For example, in the reactive 22 diffusion, a small change in the pore space geometry can lead to a large change in the 23 fluid-solid interface. Similarly, micro-scale roughness and trace impurities can also have 24 a strong impact on the outcome of experiments. Consequently, experimental validation 25 of computer models for reactive transport in porous media must be based primarily on 26 the comparison of statistical measures such as correlation functions, saturations, etc.

27

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