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Optical resonances in microsphere photonic nanojets

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Abstract

We report microsphere photonic nanojets, using the Mie theory to analyse the behaviour of the Poynting vector fields for incident visible wavelengths. Optical field structures of optical whirlpools and rhombuses are observed when incident plane waves transmit through $5\ \mu\text{m}$ and $30\ \mu\text{m}$ dielectric microspheres of various refractive indices from $n = 1.4$ to 2.2 embedded in vacuum and water media. The waist of the photonic nanojets and the distance between the microsphere and the position of maximum electric intensity determines the potential for super-resolution imaging. These parameters vary with the refractive index and size of the dielectric microspheres. We have shown that the waists of the photonic nanojets are beyond the diffraction limit without significant optical diffraction.

Keywords: photonic nanojet, optical resonance, super-resolution, microsphere, diffraction limit

(Some figures may appear in colour only in the online journal)

1. Introduction

The spatial optical resolution of an imaging system is affected by the quality and configuration of the optics and more importantly by the diffraction limit. To resolve objects below the diffraction limit, metamaterial based super-lenses have been used to restore evanescent waves [1]. Microspheres have been successfully used for nano-fabrication [2, 3] and for super-resolution imaging [4]. Photonic nanojets can also approach sub-diffraction-limited detection enhanced by backscattering of light by dielectric micro-particles [5]. Backscattering enhancement induced by the photonic nanojets was experimentally demonstrated using barium titanate microspheres at visible wavelengths [6]. Dielectric microcylinders were also shown to be able to generate nano-scaled photonic jets and to enhance backscattering of light in the visible spectrum [7]. Such enhancements can overcome the optical diffraction limit without significant optical diffraction. The phenomenon of photonic nanojets was theoretically investigated in circular dielectric microcylinders [7, 8] and dielectric microspheres [9–12] with incident plane waves. The diffraction effect of a

dielectric microsphere array was demonstrated in far-field photonic nanojets [13, 14]. Photonic nanojets were also experimentally confirmed with a polystyrene microsphere on a glass substrate with a confocal microscope [15]. The applications of photonic nanojets have been investigated in various fields, including optical data storage [16], two-photon fluorescence depletion microscopy [17], plasmonic nanodisc lithography [18], and optical forces at resonance and off-resonance wavelengths [19]. In this paper, we report the interactions of optical resonances and photonic nanojets that are generated by dielectric microspheres with diameters of $d = 5$ and $30\ \mu\text{m}$ with refractive indices of $n_1 = 1.46$ (fused silica), 1.59 (polystyrene), 1.93 ($\text{TiO}_2\text{-BaO-SiO}_2$ glass), 2.20 ($\text{TiO}_2\text{-BaO-ZnO}$ glass) embedded in media with refractive indices of $n_2 = 1.0$ (vacuum) and 1.33 (water) for incident plane waves of wavelength $\lambda = 300\ \text{nm}$, $400\ \text{nm}$, $500\ \text{nm}$, $600\ \text{nm}$, $700\ \text{nm}$ and $800\ \text{nm}$, respectively. The distance between the microspheres and the position of maximum electric intensity (DMI) and full width at half maximum (FWHM) were determined. The potential for super-resolution imaging is explained by the analysis of the waist of the photonic nanojets and the optical field structures in the

Poynting vectors in the near and near-external fields of a microsphere.

2. Methods

Light scattering and the optical cross-section of light interacting with a microsphere is modelled using the Mie theory by considering the electric (\mathbf{E}), magnetic (\mathbf{H}) and Poynting vector ($\mathbf{S} = \mathbf{E} \times \mathbf{H}$) fields. Using analytical Mie theory with Poynting vectors in the electromagnetic field, we show that high index dielectric microspheres with diameters of not only $5 \mu\text{m}$ but also $30 \mu\text{m}$ can achieve the photonic nanojet effect in the visible spectrum when illuminated with incident plane waves. Previously, photonic nanojets were limited to small particles below $10 \mu\text{m}$ [6, 7, 10] because they could take advantage of producing a minimized FWHM and enhanced backscattering. The Poynting vector flow was used to analyse the field distribution from the direction of propagation [20, 21], because the image distribution of the Poynting vector can allow one to understand the local changes of the electromagnetic energy [22]. Thus, the capability for super-resolution imaging might be predicted with photonic nanojets in the near and near-external fields of a microsphere. In the simulations, when the incident plane waves propagate through a single dielectric microsphere along the z coordinate, the flux of electric and magnetic vectors is along the x and y coordinates, respectively. The time-averaged Poynting vector is applied to the power per unit area by a wave. The single microsphere embedded in a medium is considered in the two-dimensional xz plane. For the $30 \mu\text{m}$ diameter microsphere, the boundary dimensions equal a square that is $72 \mu\text{m}$ on each side, and the number of terms on each side is 2500 points. For the $5 \mu\text{m}$ microsphere, the boundary dimensions are $12 \mu\text{m}$, and the number of terms is 1000 points. The directional energy flux density is generated by streamlines and directional vector arrows. The distribution of the Poynting vectors was obtained, showing drastic transformations in the field distribution. This relationship includes several different parameters for the incident plane wavelengths, and the sizes and refractive indices of microspheres in different media, showing the optical resonances and complicated optical field structures of the near-field transformations [23, 24].

3. Results and discussion

The behaviour of the Poynting vectors in photonic nanojets for dielectric microspheres of diameter $d = 5 \mu\text{m}$ is shown in figure 1. The red coloured lines indicate high electric intensity and the blue coloured lines indicate low electric intensity. Dielectric microspheres with two different refractive indices of $n_1 = 1.46$ and $n_1 = 2.20$ are embedded in media with $n_2 = 1.0$ (vacuum) and $n_2 = 1.33$ (water) with incident plane waves of wavelength $\lambda_2 = 400 \text{ nm}$. The optical field structures of near-field transformations that contain optical whirlpools are observed in the area of low electric intensity in the near and near-external field of the microsphere. The near-field transformations can be generated

by interference between the evanescent and the propagating waves [20]. In figure 1(a), the dielectric microsphere with $d = 5 \mu\text{m}$ and refractive index $n_1 = 1.46$ in vacuum ($n_2 = 1.0$) can clearly generate the transformation of optical whirlpools without significant diffraction in the near and near-external fields of the microsphere. Furthermore, it was experimentally confirmed to have super-resolution imaging through microsphere optical nanoscopy [4]. The interactions between the evanescent and the propagating waves can encourage photonic nanojets. This interference may provide a theoretical explanation of super-resolution imaging by transforming the evanescent waves from a near-field to a far-field propagating wave. Moreover, the resonance effect can be improved in the near field of the microsphere by this interference relationship [25]. In figure 1(b) with water ($n_2 = 1.33$), the focal position is located far from the microsphere and optical whirlpools are not generated in the near field of the microsphere. The microsphere of $n_1 = 2.20$ refractive index in vacuum ($n_2 = 1.0$) has the photonic nanojet focal point inside the microsphere, as shown in figure 1(c). Optical interference around the microsphere is generated so that a second focal position is observed near the top boundary of the microsphere. The second focal position can divert the energy of optical resonances so that the transformations of optical whirlpools may not obviously be generated. Moreover, the resonances surrounding a microsphere can decrease the backscattering enhancement of the photonic nanojets [6]. Thus, super-resolution imaging might be interrupted by the interference of the second focal position. In figure 1(d) with water ($n_2 = 1.33$), the photonic nanojet is formed on the boundary of the microsphere. The optical whirlpools that are partially generated might not be sufficient for the near-field evanescent wave to propagate wave transformations.

The Poynting vector traces of photonic nanojets of dielectric microspheres of diameter $d = 30 \mu\text{m}$ are shown in figure 2. Optical whirlpools are clearly observed in figure 2(a) $n_1 = 1.59$, $n_2 = 1.0$ and (c) $n_1 = 2.20$, $n_2 = 1.33$. This could be the result of satisfactory interference between the evanescent and propagating waves, therefore providing super-resolution imaging to observe sub-diffraction-limit objects. In figure 2(b) with $n_1 = 1.59$ and $n_2 = 1.33$, the focal position is far from the microsphere and the optical whirlpools are observed in the near and near-external fields of the microsphere. This may cause insufficient near-field transformations. In figure 2(c) with $n_1 = 2.20$ and $n_2 = 1.0$, the vector flow near the top of the microsphere is disturbed, similar to the phenomenon of figure 1(c). This may cause a change in the location of the focal position inside the microsphere. Optical whirlpool and rhombus structures are generated in the field of Poynting vectors around a microsphere. Optical whirlpool structures can be found at the bottom part inside the microsphere. Two optical whirlpool structures are observed each side of the light axis inside the microsphere with diameter $d = 30 \mu\text{m}$ in figure 2(a) $n_1 = 1.59$, $n_2 = 1.0$ and (d) $n_1 = 2.20$, $n_2 = 1.33$. Optical rhombus structures are created near the microsphere boundary. The location of the optical rhombuses depends on the refractive indices of the microsphere and

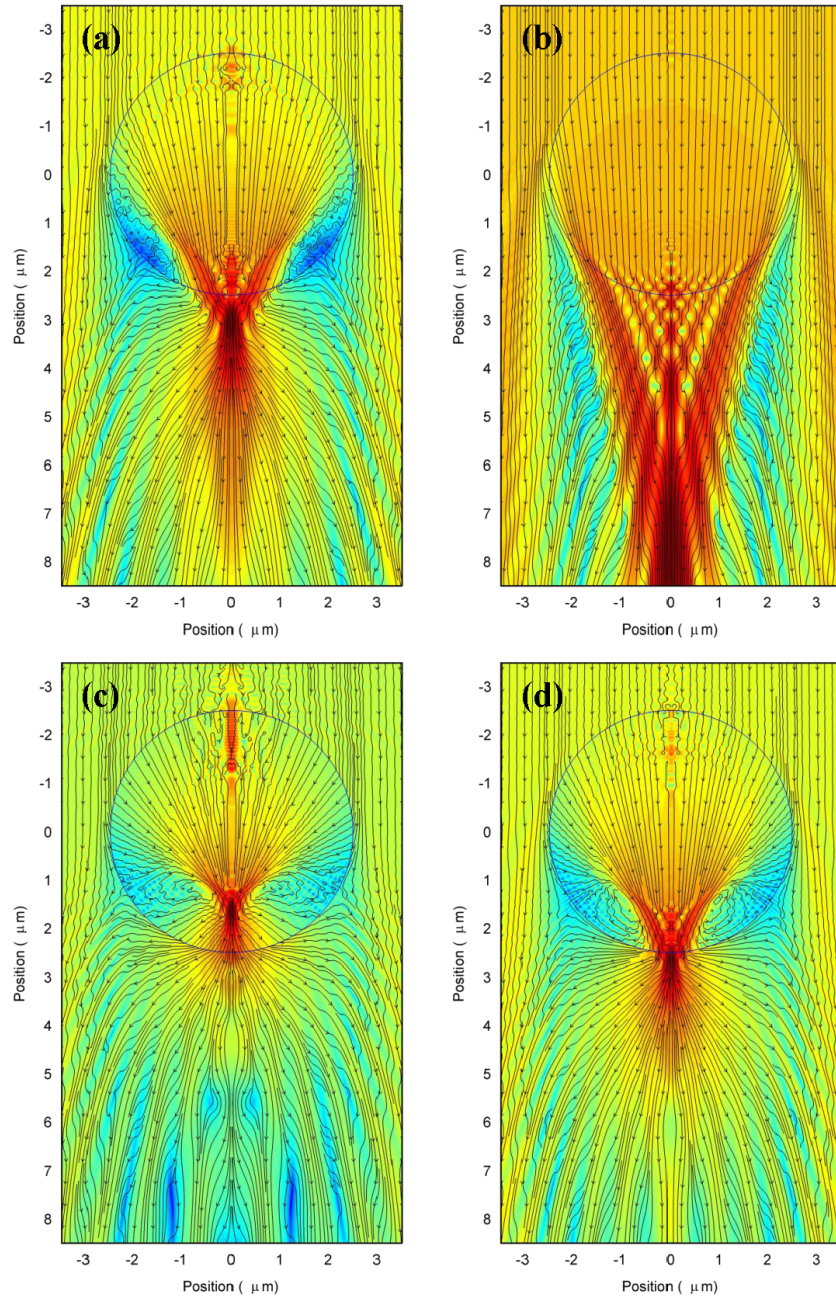


Figure 1. Poynting vectors of light transmitted through transparent dielectric microspheres of diameter $d = 5 \mu\text{m}$ and refractive index n_1 immersed in a medium of refractive index n_2 . Incident plane waves of wavelength $\lambda_2 = 400 \text{ nm}$ from the top. (a) $n_1 = 1.46, n_2 = 1.0$; (b) $n_1 = 1.46, n_2 = 1.33$; (c) $n_1 = 2.20, n_2 = 1.0$; (d) $n_1 = 2.20, n_2 = 1.33$.

the medium. Optical rhombuses are also observed outside the microspheres in figure 2(b) $n_1 = 1.59, n_2 = 1.33$, and on the boundary in figure 2(a) $n_1 = 1.59, n_2 = 1.0$, and both inside and outside in figure 2(d) $n_1 = 2.20, n_2 = 1.33$. Such optical field structures can be clearly observed in larger diameter microspheres because of the high energy field intensity. When the sub-diffraction-limited photonic nanojet is generated in the microspheres, the optical whirlpool and rhombus structures appear together with vector streamlines. Such optical field structures can be generated by clear optical interactions through the microsphere that prevent the second focal positions and optical interference around the

microsphere. Thus, optical rhombus and optical whirlpool structures may enhance optical interactions and contribute to the focusing of incident light to the microsphere. Furthermore, the energy flow can be disturbed by a second focal position if the focal position is inside the microsphere in figure 2(c) $n_1 = 2.20, n_2 = 1.0$. Such an interruption may break the optical field structures.

The FWHM and DMI of dielectric microspheres of diameter $d = 5 \mu\text{m}$ were determined with incident plane waves of wavelength $\lambda = 300, 400, 500, 600, 700$ and 800 nm , as shown in figure 3. The interactions with microspheres with refractive indices of $n_1 = 1.46, 1.59, 1.93$, and 2.20 are

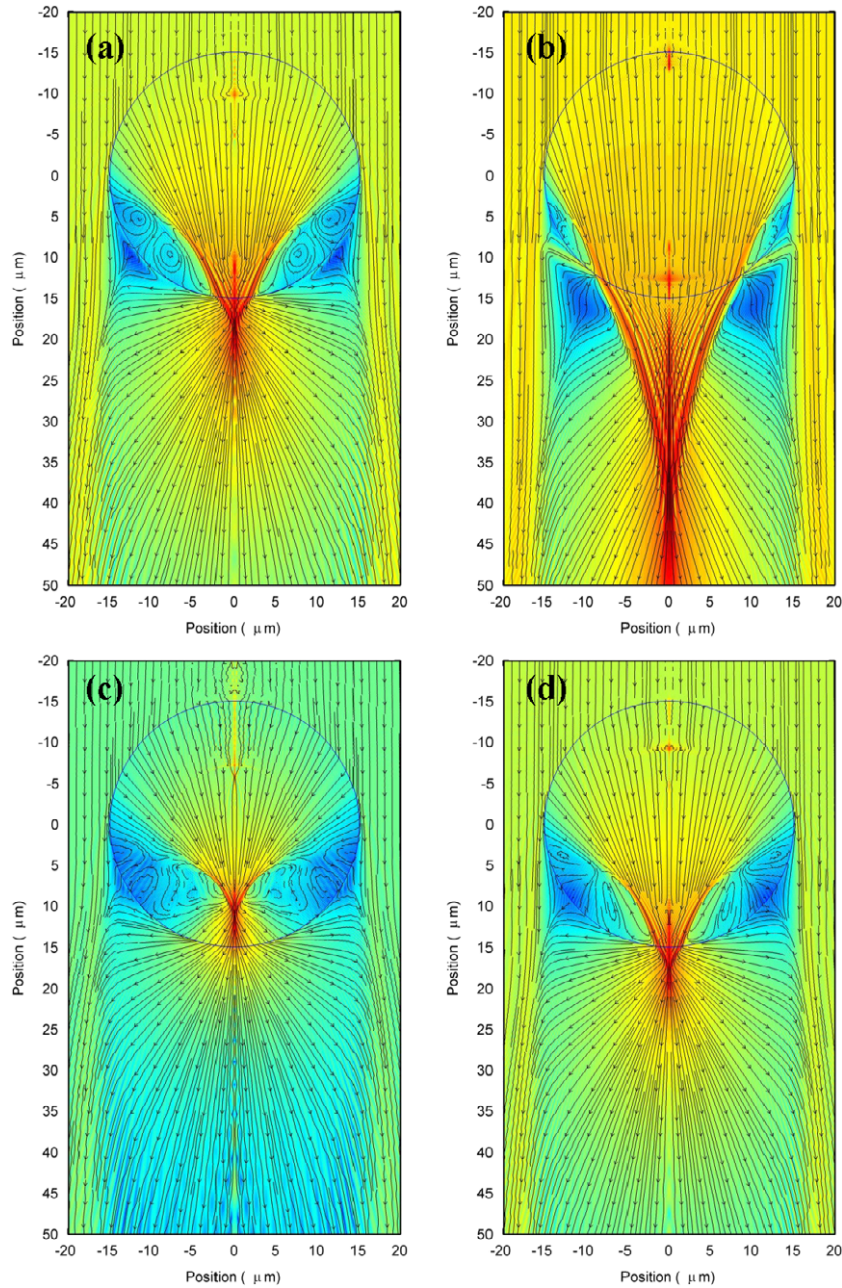


Figure 2. Poynting vector traces for plane wave light transmission through dielectric microspheres of diameter $d = 30 \mu\text{m}$ at a wavelength of $\lambda_2 = 400 \text{ nm}$. (a) $n_1 = 1.59, n_2 = 1.0$; (b) $n_1 = 1.59, n_2 = 1.33$; (c) $n_1 = 2.20, n_2 = 1.0$; (d) $n_1 = 2.20, n_2 = 1.33$.

compared for media with refractive indices $n_2 = 1.0$ (vacuum) and $n_2 = 1.33$ (water). The FWHM is increased when the light wavelength is increased. On the other hand, the DMI can be slightly decreased when the wavelength is increased. This effect can be observed clearly for water ($n_2 = 1.33$) because of optical aberrations. In vacuum ($n_2 = 1.0$), the focal positions of $n_1 = 1.93$ and 2.20 are located inside the microsphere. It can be expected that near-field transformations may not occur between evanescent and propagating waves in the near and near-external fields of the microsphere, therefore it may not support super-resolution imaging.

For the dielectric microspheres with diameter $d = 30 \mu\text{m}$, the FWHM and DMI have been computed with incident plane

waves in vacuum ($n_2 = 1.0$) and water ($n_2 = 1.33$), as shown in figure 4. The FWHM and DMI is increased as compared with the microspheres of diameter $d = 5 \mu\text{m}$. By increasing the size of the microspheres, the DMI can obviously emerge from the microspheres with refractive indices of $n_1 = 1.93$ and 2.20 and the activity of optical resonance can be stronger in the near and near-external field of the microspheres. As a result, clear optical whirlpools can be observed at the near horizontal plane of the microsphere.

The FWHMs divided by the incident wavelengths (λ) for dielectric microspheres of diameters $d = 5$ and $30 \mu\text{m}$ are shown in table 1. The refractive indices of microspheres and media are compared for wavelengths between 300 and

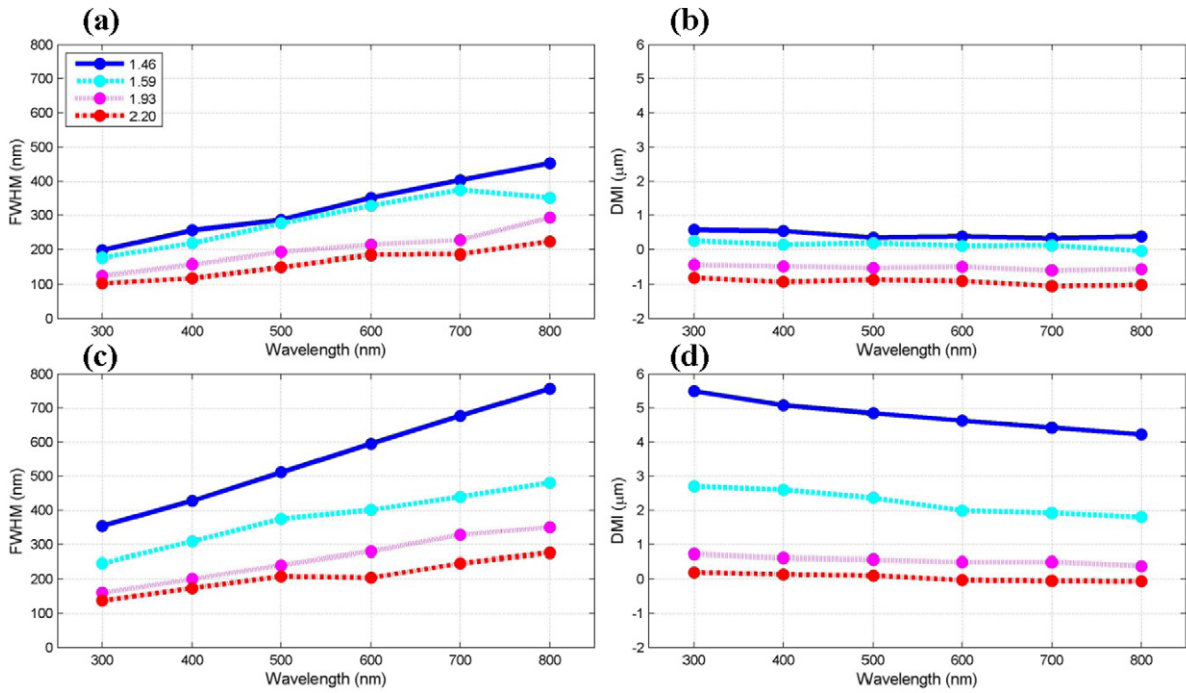


Figure 3. FWHMs and DMIs for dielectric spheres with a diameter $d = 5 \mu\text{m}$ with refractive indices of $n_1 = 1.46, 1.59, 1.93,$ and $2.20,$ respectively, in vacuum ($n_2 = 1.0$) and water ($n_2 = 1.33$) at wavelengths (λ) between 300 and 800 nm. (a) and (b) $n_2 = 1.0$; (c) and (d) $n_2 = 1.33$.

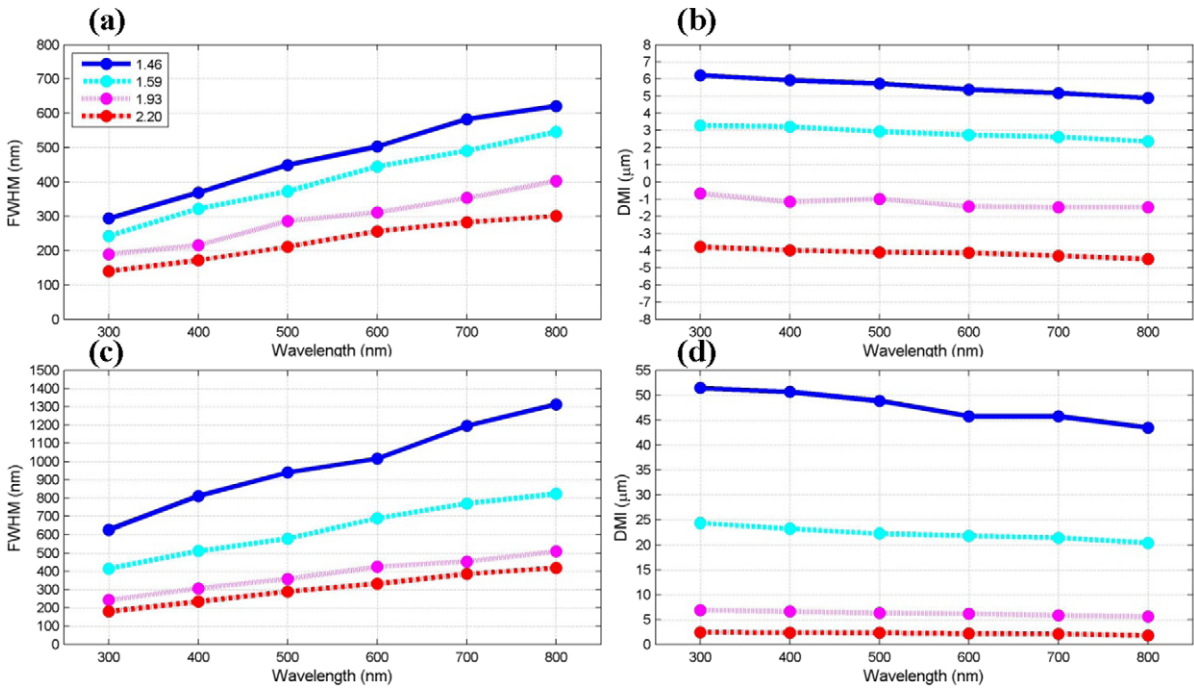


Figure 4. FWHMs and DMIs of dielectric microspheres with a diameter $d = 30 \mu\text{m}$ with four different refractive indices embedded in vacuum ($n_2 = 1.0$) and water ($n_2 = 1.33$) at wavelengths 300–800 nm. (a) and (b) $n_2 = 1.0$; (c) and (d) $n_2 = 1.33$.

800 nm. The value λ_{2-5} in the visible spectrum is calculated from the average FWHM of $\lambda_2 = 400$ to $\lambda_5 = 700$ nm. The FWHM divided by the incident wavelength (λ) seems to be independent of the wavelength, as the values are quite similar for wavelengths between 300 and 800 nm. The values can be

generally affected by the refractive indices of the microsphere and the media. The waist of the photonic nanojet is determined by the waist of the FWHM intensity [5]. In dielectric microspheres of diameter $d = 5 \mu\text{m}$, the waist of the photonic nanojet is similar to the diffraction limit for $n_1 = 1.46$ and

Table 1. FWHM divided by the incident wavelength (λ) for dielectric microspheres of diameter $d = 5$ and $30 \mu\text{m}$ in vacuum ($n_2 = 1.0$) and water ($n_2 = 1.33$). (a) $d = 5 \mu\text{m}$, $n_2 = 1.0$; (b) $d = 5 \mu\text{m}$, $n_2 = 1.33$; (c) $d = 30 \mu\text{m}$, $n_2 = 1.0$; (d) $d = 30 \mu\text{m}$, $n_2 = 1.33$.

n_1	$\lambda_1 =$ 300 nm	$\lambda_2 =$ 400 nm	$\lambda_3 =$ 500 nm	$\lambda_4 =$ 600 nm	$\lambda_5 =$ 700 nm	$\lambda_6 =$ 800 nm	$\lambda_{2-5} =$ visible
(a) $d = 5 \mu\text{m}$, $n_2 = 1.0$							
1.46	$0.66\lambda_1$	$0.64\lambda_2$	$0.57\lambda_3$	$0.59\lambda_4$	$0.58\lambda_5$	$0.57\lambda_6$	$0.59\lambda_{2-5}$
1.59	$0.59\lambda_1$	$0.55\lambda_2$	$0.55\lambda_3$	$0.55\lambda_4$	$0.53\lambda_5$	$0.44\lambda_6$	$0.55\lambda_{2-5}$
1.93	$0.42\lambda_1$	$0.40\lambda_2$	$0.39\lambda_3$	$0.36\lambda_4$	$0.32\lambda_5$	$0.37\lambda_6$	$0.37\lambda_{2-5}$
2.20	$0.34\lambda_1$	$0.30\lambda_2$	$0.30\lambda_3$	$0.31\lambda_4$	$0.27\lambda_5$	$0.28\lambda_6$	$0.29\lambda_{2-5}$
(b) $d = 5 \mu\text{m}$, $n_2 = 1.33$							
1.46	$1.18\lambda_1$	$1.07\lambda_2$	$1.02\lambda_3$	$0.99\lambda_4$	$0.97\lambda_5$	$0.95\lambda_6$	$1.01\lambda_{2-5}$
1.59	$0.82\lambda_1$	$0.77\lambda_2$	$0.75\lambda_3$	$0.67\lambda_4$	$0.63\lambda_5$	$0.60\lambda_6$	$0.70\lambda_{2-5}$
1.93	$0.54\lambda_1$	$0.50\lambda_2$	$0.48\lambda_3$	$0.47\lambda_4$	$0.47\lambda_5$	$0.44\lambda_6$	$0.48\lambda_{2-5}$
2.20	$0.46\lambda_1$	$0.43\lambda_2$	$0.42\lambda_3$	$0.34\lambda_4$	$0.35\lambda_5$	$0.35\lambda_6$	$0.39\lambda_{2-5}$
(c) $d = 30 \mu\text{m}$, $n_2 = 1.0$							
1.46	$0.98\lambda_1$	$0.92\lambda_2$	$0.90\lambda_3$	$0.84\lambda_4$	$0.83\lambda_5$	$0.78\lambda_6$	$0.87\lambda_{2-5}$
1.59	$0.81\lambda_1$	$0.81\lambda_2$	$0.75\lambda_3$	$0.74\lambda_4$	$0.70\lambda_5$	$0.68\lambda_6$	$0.75\lambda_{2-5}$
1.93	$0.63\lambda_1$	$0.54\lambda_2$	$0.57\lambda_3$	$0.52\lambda_4$	$0.51\lambda_5$	$0.50\lambda_6$	$0.53\lambda_{2-5}$
2.20	$0.47\lambda_1$	$0.43\lambda_2$	$0.42\lambda_3$	$0.43\lambda_4$	$0.40\lambda_5$	$0.38\lambda_6$	$0.42\lambda_{2-5}$
(d) $d = 30 \mu\text{m}$, $n_2 = 1.33$							
1.46	$2.09\lambda_1$	$2.03\lambda_2$	$1.88\lambda_3$	$1.69\lambda_4$	$1.71\lambda_5$	$1.64\lambda_6$	$1.83\lambda_{2-5}$
1.59	$1.38\lambda_1$	$1.28\lambda_2$	$1.16\lambda_3$	$1.15\lambda_4$	$1.10\lambda_5$	$1.03\lambda_6$	$1.17\lambda_{2-5}$
1.93	$0.81\lambda_1$	$0.77\lambda_2$	$0.72\lambda_3$	$0.71\lambda_4$	$0.65\lambda_5$	$0.64\lambda_6$	$0.71\lambda_{2-5}$
2.20	$0.60\lambda_1$	$0.58\lambda_2$	$0.58\lambda_3$	$0.55\lambda_4$	$0.55\lambda_5$	$0.52\lambda_6$	$0.57\lambda_{2-5}$

1.59 in vacuum ($n_2 = 1.0$). In dielectric microspheres of diameter $d = 30 \mu\text{m}$, the photonic nanojet can have a similar size to the diffraction limit for $n_1 = 2.20$ in water ($n_2 = 1.33$). These photonic nanojets may have the capability for super-resolution imaging due to enhanced backscattering of visible light by sub-diffraction-limit objects [7].

4. Conclusions

We have demonstrated photonic nanojets and Poynting vector interactions for dielectric microspheres embedded in vacuum and water illuminated by plane waves with wavelengths ranging from the near-ultraviolet to the near-infrared. The photonic nanojet divided by the incident wavelength can be affected by the size and the refractive indices of the dielectric microspheres and the surrounding medium rather than by the incident plane wavelengths. Optical resonances including optical whirlpools and rhombuses are observed in Poynting vector fields by analytical Mie theory. For the dielectric microspheres of diameter $30 \mu\text{m}$ with a refractive index of $n = 2.20$ in water, and of diameter $5 \mu\text{m}$ with refractive indices of $n = 1.46$ and 1.59 in vacuum, the waist of the photonic nanojets is similar to the optical diffraction limit without significant optical diffraction and the second focal position on the microspheres. For dielectric microspheres with refractive indices of $n = 2.20$ and 1.93 in vacuum, the waists of the photonic nanojets are smaller than the optical diffraction limit but the optical resonances can be disturbed because the second focal position can divert the energy of optical resonances. The phenomena of photonic nanojets from microspheres may provide potential applications for detecting sub-diffraction

features with visible light in the fields of materials, biology, chemistry, and medical sciences.

Appendix

Following the formulae [26] used in the Mie theory, the calculation of the electromagnetic fields outside and inside a sphere is expressed with incident plane wave propagation along the z coordinate, the electric vector (\mathbf{E}) field along the x coordinate, and the magnetic vector (\mathbf{H}) field along the y coordinate. The spherical coordinates $[r, \theta, \varphi]$ are indicated on the incident plane wave with the sphere positions for the electric and magnetic field,

$$E_r = e^{i k_m r \cos \theta} \sin \theta \cos \varphi, \tag{1}$$

$$H_r = \sqrt{\epsilon_m} e^{i k_m r \cos \theta} \sin \theta \sin \varphi, \tag{2}$$

$$E_\theta = e^{i k_m r \cos \theta} \cos \theta \cos \varphi, \tag{3}$$

$$H_\theta = \sqrt{\epsilon_m} e^{i k_m r \cos \theta} \cos \theta \sin \varphi, \tag{4}$$

$$E_\varphi = -e^{i k_m r \cos \theta} \sin \varphi, \tag{5}$$

$$H_\varphi = -\sqrt{\epsilon_m} e^{i k_m r \cos \theta} \sin \varphi. \tag{6}$$

The wavevectors for the embedding medium, sphere and vacuum are $k_m = 2\pi\sqrt{\epsilon_m}/\lambda$, $k_p = 2\pi\sqrt{\epsilon_p}/\lambda$ and $k_0 = 2\pi/\lambda$ respectively, where λ is the radiation wavelength. The dielectric permittivities of the embedding medium and the sphere are ϵ_m and ϵ_p respectively, and the magnetic permeability is μ . A dielectric sphere with radius a is identified by the complex refractive index $\sqrt{\epsilon_m} = n_m + ik_m$ for the medium and $\sqrt{\epsilon_p} = n_p + ik_p$ for the sphere.

A non-magnetic sphere ($\mu_p = 1$) in vacuum (superscript s for the scattered wave) gives

$$E_r^{(s)} = \frac{\cos \varphi}{(k_m r)^2} \sum_{n=1}^{\infty} n(n+1) {}^e B_n \zeta_n(k_m r) P_n^{(1)}(\cos \theta), \tag{7}$$

$$E_{\theta}^{(s)} = \frac{\cos \varphi}{k_m r} \sum_{n=1}^{\infty} \left[{}^e B_n \zeta_n'(k_m r) P_n^{(1)'}(\cos \theta) \sin \theta - i {}^m B_n \zeta_n(k_m r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} \right], \tag{8}$$

$$E_{\varphi}^{(s)} = \frac{\sin \varphi}{k_m r} \sum_{n=1}^{\infty} \left[{}^e B_n \zeta_n'(k_m r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} - i {}^m B_n \zeta_n(k_m r) P_n^{(1)'}(\cos \theta) \sin \theta \right], \tag{9}$$

$$H_r^{(s)} = \frac{\sqrt{\varepsilon_m} \sin \varphi}{(k_m r)^2} \sum_{n=1}^{\infty} n(n+1) {}^m B_n \zeta_n(k_m r) P_n^{(1)}(\cos \theta) \tag{10}$$

$$H_{\theta}^{(s)} = i \frac{\sin \varphi}{k_0 r} \sum_{n=1}^{\infty} \left[{}^e B_n \zeta_n(k_m r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} + i {}^m B_n \zeta_n'(k_m r) P_n^{(1)'}(\cos \theta) \sin \theta \right], \tag{11}$$

$$H_{\varphi}^{(s)} = i \frac{\cos \varphi}{k_0 r} \sum_{n=1}^{\infty} \left[{}^e B_n \zeta_n(k_m r) P_n^{(1)'}(\cos \theta) \sin \theta + i {}^m B_n \zeta_n'(k_m r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} \right]. \tag{12}$$

The inside of the non-magnetic sphere (superscript p for the scattered wave) is given by

$$E_r^{(p)} = \frac{\cos \varphi}{(k_p r)^2} \sum_{n=1}^{\infty} n(n+1) {}^e A_n \Psi_n(k_p r) P_n^{(1)}(\cos \theta), \tag{13}$$

$$E_{\theta}^{(p)} = \frac{\cos \varphi}{k_p r} \sum_{n=1}^{\infty} \left[{}^e A_n \Psi_n'(k_p r) P_n^{(1)'}(\cos \theta) \sin \theta - i {}^m A_n \Psi_n(k_p r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} \right], \tag{14}$$

$$E_{\varphi}^{(p)} = \frac{\sin \varphi}{k_p r} \sum_{n=1}^{\infty} \left[{}^e A_n \Psi_n'(k_p r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} - i {}^m A_n \Psi_n(k_p r) P_n^{(1)'}(\cos \theta) \sin \theta \right], \tag{15}$$

$$H_r^{(p)} = \frac{\sqrt{\varepsilon_m} \sin \varphi}{(k_p r)^2} \sum_{n=1}^{\infty} n(n+1) {}^m A_n \Psi_n(k_p r) P_n^{(1)}(\cos \theta), \tag{16}$$

$$H_{\theta}^{(p)} = i \frac{\sin \varphi}{k_0 r} \sum_{n=1}^{\infty} \left[{}^e A_n \Psi_n(k_p r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} + i {}^m A_n \Psi_n'(k_p r) P_n^{(1)'}(\cos \theta) \sin \theta \right], \tag{17}$$

$$H_{\varphi}^{(p)} = i \frac{\cos \varphi}{k_0 r} \sum_{n=1}^{\infty} \left[{}^e A_n \Psi_n(k_p r) P_n^{(1)'}(\cos \theta) \sin \theta + i {}^m A_n \Psi_n'(k_p r) \frac{P_n^{(1)}(\cos \theta)}{\sin \theta} \right], \tag{18}$$

with the coefficients

$${}^e B_n = i^{n+1} \frac{2n+1}{n(n+1)} a_n, \quad {}^m B_n = i^{n+1} \frac{2n+1}{n(n+1)} b_n, \tag{19}$$

$${}^e A_n = i^{n+1} \frac{2n+1}{n(n+1)} c_n, \quad {}^m A_n = i^{n+1} \frac{2n+1}{n(n+1)} d_n, \tag{20}$$

where

$$a_n = \frac{y \Psi_n'(x) \Psi_n(y) - x \Psi_n(x) \Psi_n'(y)}{y \zeta_n'(x) \Psi_n(y) - x \Psi_n'(y) \zeta_n(x)}, \tag{21}$$

$$b_n = \frac{y \Psi_n'(y) \Psi_n(x) - x \Psi_n(y) \Psi_n'(x)}{y \Psi_n'(y) \zeta_n(x) - x \Psi_n(y) \zeta_n'(x)}, \tag{22}$$

$$c_n = \frac{y \zeta_n(x) \Psi_n'(x) - y \zeta_n'(x) \Psi_n(x)}{y \zeta_n'(x) \Psi_n(y) - x \Psi_n'(y) \zeta_n(x)}, \tag{23}$$

$$d_n = \frac{y \zeta_n'(x) \Psi_n(x) - y \zeta_n(x) \Psi_n'(x)}{y \Psi_n'(y) \zeta_n(x) - x \Psi_n(y) \zeta_n'(x)}, \tag{24}$$

$$x = k_m a, \quad y = k_p a, \tag{25}$$

$$\zeta_n(\rho) = \rho h_n^{(1)}(\rho) = \sqrt{\frac{\pi \rho}{2}} H_{n+\frac{1}{2}}^{(1)}(\rho), \quad \zeta_n'(\rho) = \frac{\partial \zeta_n(\rho)}{\partial \rho}, \tag{26}$$

$$\Psi_n(\rho) = \rho j_n(\rho) = \sqrt{\frac{\pi \rho}{2}} J_{n+\frac{1}{2}}(\rho), \quad \Psi_n'(\rho) = \frac{\partial \Psi_n(\rho)}{\partial \rho}, \tag{27}$$

are the scattering coefficients and we make use of

(i) the associated Legendre function,

$$P_n^m(\cos \theta) \equiv \frac{(1 - (\cos \theta)^2)^{\frac{m}{2}}}{2^n n! d^{n+m}} ((\cos \theta)^2 - 1)^n, \quad (m \geq 0)$$

$$P_n^m(\cos \theta) \equiv (-1)^{-m} \frac{(n+m)!}{(n-m)!} P_n^{-m}(\cos \theta), \quad (m < 0) \quad (28)$$

(ii) the spherical Hankel function,

$$h_n^{(1)}(\rho) = \sqrt{\frac{\pi}{2\rho}} H_{n+\frac{1}{2}}^{(1)}(\rho), \quad (29)$$

(iii) the spherical Bessel function,

$$j_n(\rho) = \sqrt{\frac{\pi}{2\rho}} J_{n+\frac{1}{2}}(\rho). \quad (30)$$

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