Dynamics and Bifurcation of a Model for Hormonal Control of the Menstrual Cycle with Inhibin Delay *

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Abstract

A system of 13 ordinary differential equations with 42 parameters is presented to model hormonal regulation of the menstrual cycle. The model's fit to data is improved by including a time delay for the effect of inhibin on the synthesis of follicle stimulating hormone. Biological reasons for this improvement are discussed. Bifurcations with respect to changes in three important parameters are examined. One parameter represents the level of estradiol adequate for significant synthesis of luteinizing hormone. Bifurcation diagrams with respect to this parameter reveal an interval of parameter values for which a unique stable periodic solution exists and this solution represents a menstrual cycle during which ovulation occurs. The second parameter measures mass transfer between the first two stages of ovarian development and is indicative of healthy follicular growth. The third parameter is the time delay. Changes in the second parameter and the time delay affect the size of the uniqueness interval defined with respect to the first parameter. Saddle-node, transcritical and degenerate Hopf bifurcations are studied.

Key Words. Estradiol, follicle, Hill function, degenerate Hopf bifurcation

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1. Introduction.

Systems of ordinary and delayed differential equations have been used to model hormonal regulation of the human menstrual cycle, e.g., see Bogumil et al., 1972a, 1972b [3, 4], McIntosh and McIntosh, 1980 [15], Plouffe and Luxenberg, 1992 [20], Selgrade and Schlosser, 1999 [25], Schlosser and Selgrade, 2000 [22], Harris-Clark et al., 2003 [10], Reinecke and Deuflhard, 2007 [21], and Pasteur, 2008 [19]. Dual control of the menstrual cycle depends on hormones produced by the hypothalamus and the pituitary glands in the brain and by the ovaries. The pituitary secretes follicle stimulating hormone (FSH)and luteinizing hormone (LH) which control ovarian development and ovulation, see [11, 29, 30]. The ovaries produce estradiol (E_2) , progesterone (P_4) and inhibin (Inh)which affect the synthesis and release of FSH and LH, see [12, 14, 26]. Harris, Pasteur, Schlosser and Selgrade [10, 19, 22, 25] have derived a 13-dimensional system of delayed differential equations which captures these interacting mechanisms. Model parameters were identified using two different clinical data sets for normally cycling women (McLachlan et al., 1990 [16], and Welt et al., 1999 [27]). Model simulations with parameters from the McLachlan data [16] revealed two stable periodic solutions [10] — one fitting the McLachlan data for normally cycling women and the other being non-ovulatory because of no LH surge. The non-ovulatory cycle has similarities to an abnormal cycle of a woman with polycystic ovarian syndrome (PCOS) [28], the leading cause of female infertility. However, model simulations corresponding to the Welt parameters produced only one stable periodic solution and it fits the Welt data for normally cycling women. Selgrade et al., 2009 [24], explained this apparent inconsistency by showing that a change in only one sensitive parameter of the Welt system would result in the Welt model exhibiting bistability like the McLachlan model.

Abnormal cycling and non-ovulatory cycling have serious health and reproductive consequences. In fact, between 6% and 9% of adult women exhibit some symptoms of PCOS, see Azziz *et al.*, 2004 [2], and Alvarez-Blasco *et al.*, 2006 [1]. Since cycle irregularities are usually associated with abnormal hormone levels, mathematical models of hormonal regulation may provide information about parameter variations which result in abnormal cycling and may provide insights about possible hormonal therapies. In an effort to understand what parameter ranges result in normal and abnormal cycling, Selgrade [23] set the time-delays to zero in the Welt model and used the software XPPAUT [7] to study bifurcation diagrams with respect to two of the most sensitive parameters. Bifurcation diagrams for the resulting autonomous system could be drawn with the features of AUTO [5] in XPPAUT. This autonomous model gives an acceptable fit to the 28 day Welt data set [27] except some hormone peaks are lower than the data and the period for the normal cycle is only about 26 days, see Figure 1.



Figure 1: LH and E_2 simulations for 3 cycles of the Welt model with inhibin delay τ of 1.5 days (green curves) and no delay (red curves) with data points (84 black dots) corresponding to the 28 day data from Welt et al. [27] plotted 3 times. The vertical dashed line indicates day 29, where the second cycle begins and where both solution orbits are very close to one another.

The two key parameters for analysis in [23] are Km_{LH} and c_2 . Km_{LH} represents the level of E_2 sufficient for significant LH synthesis and the LH surge and c_2 indicates the ovarian mass transfer rate between the first two stages of ovarian development. The bifurcation diagram with respect to Km_{LH} reveals an interval of Km_{LH} parameter values for which a unique stable periodic solution exists and this solution represents a menstrual cycle with a LH surge adequate for ovulation. If Km_{LH} lies outside this cycle uniqueness interval then either no LH surge occurs or there are two stable cycles — one is ovulatory and the other is non-ovulatory. Changes in c_2 affect the size of this interval because of the positions of Hopf, saddle-node and transcritical bifurcations as discussed in [23].

In this study, we carry out a bifurcation analysis for the system of delayed differential equations using the DDEBIFTOOL [6], which is designed to handle the delay. The original model [10] had three discrete time-delays (one corresponding to each ovarian hormone) which represented the time interval between changes in ovarian hormone concentrations and subsequent changes they cause in synthesis rates of the pituitary hormones. Here we show that including only a delay of $\tau = 1.5$ days for inhibin is needed to improve the fit to the Welt data, see Figure 1. The other two delays were no more than a day and did not contribute significant additional improvement. Also, the system with the inhibin delay has larger uniqueness intervals than the model with no delay (see Table 2). In Section 3 we speculate about the biological reasons for this improvement in model behavior due to inhibin delay. We examine bifurcation diagrams with respect to Km_{LH} for the delayed system and show that the cycle uniqueness interval is usually determined by two saddlenode bifurcations. For the delay τ fixed at 1.5 days, we illustrate how this interval may be enlarged by varying c_2 due to the occurrence of two degenerate Hopf bifurcations. Then for fixed c_2 , we increase the delay parameter τ from 0 to 1.5 to unfold transcritical bifurcations and produce large cycle uniqueness intervals. Finally, we illustrate how loops in the Km_{LH} bifurcation diagrams may appear and disappear by varying the parameters τ and c_2 .

2. Biological Background and Model Equations.

The menstrual cycle for an adult female consists of the follicular phase, ovulation, the luteal phase and menstruation (e.g., see Odell [17] or Ojeda [18]). During the follicular phase, FSH produced by the pituitary gland promotes the development of 6 to 12 follicles. Typically one dominant follicle is selected to grow to maturity and to produce a large amount of E_2 which primes the pituitary for LH secretion. At mid-cycle, a surge of LHover a period of 4 or 5 days results in ovulation. After releasing its egg, the dominant follicle becomes the corpus luteum which produces hormones in preparation for pregnancy and produces P_4 and Inh which inhibit LH and FSH, respectively. If fertilization does not occur, the corpus luteum atrophies, menstruation follows and a rise in FSH marks the beginning of the next cycle.

Harris, Pasteur, Schlosser and Selgrade [9, 10, 19, 22, 25] developed a model for this endocrine control system based on 13 ordinary differential equations (S) with three auxiliary equations (A) and with discrete time-delays. Four of these differential equations (S1-S4) describe the synthesis, release and clearance of LH and FSH. The state variables RP_{LH} and RP_{FSH} represent the amounts of these hormones in the pituitary and LH and FSH represent the blood concentrations of these hormones. The biological literature (e.g., [12, 14, 29]) indicates that LH exhibits a biphasic response to E_2 . To capture this our model assumes that the effect of E_2 on LH synthesis is different than the effect on LH release, i.e., E_2 inhibits release (see the denominator of the second term in (S1)) but at high levels E_2 promotes synthesis (see the Hill function in the numerator of the first term of (S1)). On the other hand, P_4 inhibits LH synthesis but promotes release. The release term appears in (S1) as a decay term and in (S2) as a growth term, where it is divided by blood volume v. The equations (S3-S4) for FSH are similar except the synthesis term has Inh inhibition which is delayed by time τ . The parameters in (S1-S4) are named according to the traditional usage for chemical reactions, e.g., $V_{1,LH}$ denotes the velocity of the reaction (see Keener and Sneyd, 2009, [13]).

The state variables in (S5-S13) represent tissue masses of 9 distinct stages of the ovary during the follicular and luteal phases of the cycle. ReF, SeF and PrF denote the recruited follicles, the secondary follicles and the preovulatory or dominant follicle, respectively. Ov_1 and Ov_2 represent periovulatory stages and Lut_i , i = 1, ..., 4, denote four luteal stages. LH and FSH promote tissue growth within a stage and the transformation of tissue from one stage to the next. Since clearance from the blood of the ovarian hormones is on a fast time scale, we assume that blood levels of E_2 , P_4 , and Inh are at quasi-steady state [13] as did Bogumil *et al.* [3]. Hence, we take these concentrations to be proportional to the tissue masses during the appropriate stages of the cycle giving the three auxiliary equations (A1-A3) for the ovarian hormones. Here we study the 13-dimensional system (S) with (A), which has one time-delay τ . The forty-two parameters are listed in Table 1 and correspond to those which Selgrade [23] used to analyze bifurcation diagrams for the Welt system with time-delay set to zero.

Auxiliary Equations (A)

$$E_2 = e_0 + e_1 SeF + e_2 PrF + e_3 Lut_4$$
(A1)

$$P_4 = p_0 + p_1 L u t_3 + p_2 L u t_4 \tag{A2}$$

$$Inh = h_0 + h_1 PrF + h_2 Lut_2 + h_3 Lut_3.$$
 (A3)

System (S)

$$\frac{d}{dt}RP_{LH} = \frac{V_{0,LH} + \frac{V_{1,LH}E_2^8}{Km_{LH}^8 + E_2^8}}{1 + P_4/Ki_{LH,P}} - \frac{k_{LH}\left[1 + c_{LH,P}P_4\right]RP_{LH}}{1 + c_{LH,E}E_2}$$
(S1)

$$\frac{d}{dt}LH = \frac{1}{v}\frac{k_{LH}\left[1 + c_{LH,P}P_4\right]RP_{LH}}{1 + c_{LH,E}E_2} - a_{LH}LH$$
(S2)

$$\frac{d}{dt} RP_{FSH} = \frac{V_{FSH}}{1 + Inh(t - \tau)/Ki_{FSH,Inh}} - \frac{k_{FSH} \left[1 + c_{FSH,P} P_4\right] RP_{FSH}}{1 + c_{FSH,E} E_2^2}$$
(S3)

$$\frac{d}{dt}FSH = \frac{1}{v}\frac{k_{FSH}\left[1 + c_{FSH,P}P_4\right]RP_{FSH}}{1 + c_{FSH,E}E_2^2} - a_{FSH}FSH$$
(S4)

$$\frac{d}{dt}ReF = bFSH + [c_1FSH - c_2LH^{\alpha}]ReF$$
(S5)

$$\frac{d}{dt}SeF = c_2 LH^{\alpha} ReF + \left[c_3 LH^{\beta} - c_4 LH\right] SeF$$
(S6)

$$\frac{d}{dt}PrF = c_4 LH SeF - c_5 LH^{\gamma} PrF \tag{S7}$$

$$\frac{d}{dt}Ov_1 = c_5 LH^{\gamma} PrF - d_1 Ov_1 \tag{S8}$$

$$\frac{d}{dt}Ov_2 = d_1 Ov_1 - d_2 Ov_2 \tag{S9}$$

$$\frac{d}{dt}Lut_1 = d_2 Ov_2 - k_1 Lut_1 \tag{S10}$$

$$\frac{d}{dt}Lut_2 = k_1 Lut_1 - k_2 Lut_2 \tag{S11}$$

$$\frac{d}{dt}Lut_3 = k_2 Lut_2 - k_3 Lut_3 \tag{S12}$$

$$\frac{a}{dt}Lut_4 = k_3Lut_3 - k_4Lut_4.$$
(S13)

Eqs. $(S1-S4)$				
τ	1.5 days			
k_{LH}	$2.42 \mathrm{day}^{-1}$			
a_{LH}	$14.0 \mathrm{day}^{-1}$			
$V_{0,LH}$	500 IU/day			
$V_{1,LH}$	4500 IU/day			
Km_{LH}	200 pg/mL			
$Ki_{LH,P}$	12.2 ng/mL			
$c_{LH,E}$	0.004 mL/pg			
$C_{LH,P}$	0.26 mL/ng			
V_{FSH}	$375 \mathrm{IU/day}$			
a_{FSH}	$8.21 \mathrm{day}^{-1}$			
k_{FSH}	$1.90 day^{-1}$			
$c_{FSH,E}$	$0.0018 \ {\rm mL^2/pg^2}$			
Ki _{FSH,Inh}	3.5 ~ IU/mL			
$c_{FSH,P}$	12.0 mL/ng			
v	2.50 L			

Eqs. (S5-S13) $0.05 \text{ L} \ \mu\text{g}/(\text{IU day})$ b0.08 L/(IU day) c_1 $0.07 (L/IU)^{\alpha}/day$ c_2 $0.13 (L/IU)^{\beta}/day$ c_3 0.027 L/(IU day) c_4 $0.51 \ (L/IU)^{\gamma}/day$ c_5 0.50 day^{-1} d_1 0.56 day^{-1} d_2 0.55 day^{-1} k_1 0.69 day^{-1} k_2 0.85 day^{-1} k_3 $0.85 \, \rm day^{-1}$ k_4 0.79 α β 0.16 0.02 γ

Eq. (A)							
e_0	30 pg/mL	p_2	$0.048 \ \rm kL^{-1}$				
e_1	$0.11 \ \mathrm{L^{-1}}$	h_0	0.4 IU/mL				
e_2	$0.21 \ \mathrm{L^{-1}}$	h_1	$0.009 \text{ IU}/(\mu \text{g mL})$				
e_3	$0.45 \ \mathrm{L^{-1}}$	h_2	$0.029 \text{ IU}/(\mu \text{g mL})$				
p_0	0 ng/mL	h_3	$0.018 \text{ IU}/(\mu \text{g mL})$				
p_1	$0.048 \ \rm kL^{-1}$						

Table 1: Parameters and values for system (S) and auxiliary equations (A).

3. Effect of Inhibin Delay on Model Fit and Uniqueness Intervals.

System (A) and (S) with the inhibin delay is a better approximation to the 28 day Welt data [27] than the model with no delay [23]. The delay model has an asymptotically stable cycle of period 28 days instead of 26 days for the no-delay model. The LH data indicates a 14 day follicular phase and the position and height of the LH surge for the delay model is consistent with that (see Figure 1). Also, the delay E_2 follicular and luteal peaks are higher then the no-delay. To understand why the inhibin delay is responsible for these differences we examine hormone profiles and ovarian stages for both models over three carefully chosen consecutive cycles. MATLAB simulations of both models were run with the following initial conditions (rounded to two decimal places) given in the order of the 13 state variables in (S), {29.65, 6.86, 8.47, 6.15, 3.83, 11.51, 5.48, 19.27, 45.64, 100.73, 125.95, 135.84, 168.71}. The simulations were aligned so that both delay and no-delay periodic orbits are as close to one another as possible at the beginning of their second cycle, indicated by the vertical line at day 29 in Figures 1, 2 and 3. This was done so that the point of our comparison would be the second cycle in these figures and the preceding cycle would

also be plotted because hormone profiles during the luteal phase of the preceding cycle

influence behavior in our comparison cycle. The key feature to observe in Figure 2 is that the no-delay FSH (red curve) is higher than the delay FSH from day 19 until day 34, which includes the first six days of the follicular phase of our comparison cycle. Since FSH stimulates follicular development, the no-delay ovarian stages of the second cycle increase sooner than the delay ovarian stages and the no-delay cycle is advanced ahead of the delay cycle (see Figure 3). Nodelay FSH is higher because delay Inh has a greater inhibitory effect than no-delay Inhon FSH synthesis (see (S3)) during that period. Delay Inh (green curve) is greater than no-delay Inh from day 16 to day 24 where the curves cross and then both curves decrease in parallel until day 34. These Inh curves are so close to one another (see Figure 2) from day 24 to day 34 that the delay of 1.5 days results in delay Inh(t-1.5) being greater than no-delay Inh(t) for $24 \le t \le 34$. Effectively, from day 17.5 to day 34 the synthesis of delay FSH is suppressed more than the synthesis of no delay FSH. This causes the no-delay follicles to develop sooner then the delay follicles with the consequence via (A1) that no-delay E_2 rises sooner (see Figure 1). Since E_2 inhibits FSH release (see (S3-S4)), this earlier rise in E_2 tends to decrease no-delay FSH sooner than delay FSH with the result that the no-delay follicular stages develop to a lesser extent than the delay stages (Figure 3). Also, because E_2 promotes LH synthesis, the LH surge is earlier and smaller for the no-delay model (Figure 1). The cumulative effect of these profile differences is a shortening of the no-delay cycle length by 2 days and a reduction in no-delay hormone peaks.



Figure 2: FSH and Inh simulations for 3 consecutive cycles of the delay model (green curves) and the no delay model (red curves) with 84 data points (Welt et al. [27]). The vertical dashed line indicates day 29, the beginning of the second cycle. From day 17.5 to day 34 the synthesis of delay FSH is suppressed more than the synthesis of no delay FSH because of inhibin differences.



Figure 3: First 3 ovarian stages ReF, SeF and PrF for 3 consecutive cycles of the delay model (green curves) and the no delay model (red curves).

The parameters Km_{LH} and c_2 are two of the three most sensitive parameters when sensitivity is measured with respect to the E_2 follicular peak as system output [19, 24]. The parameter Km_{LH} is the half-saturation constant in the Hill function in (S1), $(V_{1,LH} E_2^8)/(Km_{LH}^8 + E_2^8)$. This sigmoidal shaped function (see Figure 4) acts like a threshold for the synthesis of LH in response to E_2 blood levels. Once E_2 concentration reaches the value Km_{LH} , half way up the sigmoid as indicated by the dashed line in Figure 4, then the pituitary is synthesizing LH in large amounts, which is necessary for ovulation. For larger values of Km_{LH} , E_2 must reach a higher level to produce the same LH synthesis rate. Because higher follicular E_2 levels may suggest a greater probability of abnormal cycling [23, 24], we construct bifurcation diagrams where LH is plotted against the parameter Km_{LH} to determine the number of stable cycles for a given Km_{LH} value and to determine LH surge height along each cycle. When similar bifurcation diagrams were drawn for the no-delay model [23], an interval of Km_{LH} values was observed for which a unique stable periodic solution existed and it represented an ovulatory cycle. The length of this cycle uniqueness interval varied as the parameter c_2 was changed [23]. The present study reveals that these uniqueness intervals are larger for the model with inhibin delay, (A) and (S), as indicated in Table 2.



Figure 4: Graphs of Hill functions, $(V_{1,LH} E_2^8)/(Km_{LH}^8 + E_2^8)$, for three values of Km_{LH} . The dashed line indicates the synthesis rate when $E_2 = Km_{LH}$, half-saturation.

Here, the software DDEBIFTOOL [6] is used to construct bifurcation diagrams where the maximal LH value along a periodic solution or at a steady-state solution is plotted against the parameter Km_{LH} . Figure 5 displays this bifurcation diagram where the remaining parameters are those in Table 1. Stable and unstable periodic orbits and equilibria are depicted. Saddle-node (SN) and Hopf (HB) bifurcations are labeled. The curve along the lower portion of Figure 5 tracks an equilibrium, which undergoes a destabilizing Hopf bifurcation as Km_{LH} increases through 64 and a stabilizing Hopf bifurcation at $Km_{LH} = 248$. The bifurcation at $Km_{LH} = 64$ results in a small amplitude, stable, periodic orbit which persists until $Km_{LH} = 73$. Stable and unstable cycles appear together at $Km_{LH} = 68$ via a saddle-node bifurcation of periodic orbits. The unstable orbit coalesces with the stable Hopf orbit at $Km_{LH} = 73$ and both disappear in another saddle-node. The stable cycle appearing at $Km_{LH} = 68$ grows in amplitude, continues across the top portion of the diagram and disappears in a saddle-node at $Km_{LH} = 282$. This branch of periodic solutions represents the ovulatory cycles of the model (S) with (A), where the * indicates the cycle corresponding to the Km_{LH} value of Table 1, 200 pg/ml. Analogous behavior occurs at the right side of the bifurcation diagram where the hysteresis character of the curve of periodic orbits is evident. Clearly, for Km_{LH} from 227 to 282 there is a stable, large amplitude ovulatory cycle and a stable, small amplitude non-ovulatory cycle or stable equilibrium. For Km_{LH} in the interval between the lower SN's in Figure 5 ($73 < Km_{LH} < 227$), there is only one stable cycle and it is ovulatory. Selgrade [23] referred to this Km_{LH} interval as the cycle uniqueness interval. In the context of this cycle regulation model, a woman's Km_{LH} parameter must fall within her cycle uniqueness interval for her to be assured of only a normal cycle. From Figure 5, we observe that decreasing Km_{LH} from 200 pg/mL keeps it within the interval and increases the height of the LH surge. However, increasing Km_{LH} to 227 moves Km_{LH} to a region of multiple stable cycles and possible non-ovulation. For $c_2 = 0.07$, the diameter of this cycle uniqueness interval is 154 for the delay model and only 114 for the no-delay model (see Table 2).



Figure 5: In this bifurcation diagram the maximal LH value along a periodic solution or at an equilibrium is plotted against Km_{LH} when $\tau = 1.5$ and $c_2 = 0.07$. HB and SN denote Hopf and saddle-node bifurcations. The * indicates the position of the cycle for the parameters of Table 1 and this cycle is the only stable solution at $Km_{LH} = 200 \text{ pg/mL}$. The cycle uniqueness interval is the interval between the lower saddle-nodes, i.e., 73 < $Km_{LH} < 227$.

Table 2: Size of cycle uniqueness interval for inhibin delay $\tau = 0$ (column 2) and $\tau = 1.5$ days (column 4) for increasing values of c_2 . $c_2 = 0.07$ and $Km_{LH} = 200$ pg/mL give the best fit to data.

c_2	size $(\tau = 0)$	Km_{LH} bounds $(\tau = 0)$	size $(\tau = 1.5)$	Km_{LH} bounds ($\tau = 1.5$)
0.03	126	$147 < Km_{LH} < 273$	271	$40 < Km_{LH} < 311$
0.04	50	$181 < Km_{LH} < 231$	226	$44 < Km_{LH} < 270$
0.05	81	$153 < Km_{LH} < 234$	173	$85 < Km_{LH} < 258$
0.06	118	$122 < Km_{LH} < 230$	167	$80 < Km_{LH} < 247$
0.07	114	$98 < Km_{LH} < 212$	154	$73 < Km_{LH} < 227$
0.08	102	$84 < Km_{LH} < 186$	141	$63 < Km_{LH} < 204$

For the no-delay model, Selgrade [23] investigated how variations in the ovarian transfer parameter c_2 changed the size of the cycle uniqueness interval. Increasing c_2 from $c_2 = 0.07$ causes an increased transfer of mass from the first follicular stage ReF to the second stage SeF which diminishes the development of not only ReF but of all subsequent ovarian stages. Effectively, ovarian hormone production is reduced and the cycle uniqueness interval is decreased for both delay and no-delay models. Table 2 lists the cycle uniqueness intervals for various values of c_2 which we compute for the delay model $(\tau = 1.5)$ and which were reported in [23] for the no-delay model $(\tau = 0)$. Decreasing c_2 from 0.07 in increments of 0.01 widens the cycle uniqueness interval for the delay model but shrinks it for the no-delay model until $c_2 = 0.03$. For the no-delay model as c_2 decreases, the hysteresis curves enlarge and the Hopf points move closer together resulting in a narrowing of the gap between the lower two saddle-nodes. Then, as described in [23], an unfolding of a transcritical bifurcation of periodic solutions as c_2 decreases through 0.0305 leads to the disappearance of the left hysteresis curve and a rapid expansion of the cycle uniqueness interval. In contrast, for the delay model, decreasing c_2 from 0.07 causes the hysteresis curves to enlarge only slightly (see Figure 6), the Hopf points to move apart and the left hysteresis curve to disappear due to two degenerate Hopf bifurcations described below. The uniqueness interval for $\tau = 1.5$ when $c_2 = 0.03$ (Figure 7) is over twice as large as that for the no-delay model.



Figure 6: Bifurcation diagram with respect to Km_{LH} when $\tau = 1.5$ and $c_2 = 0.05$. HB and SN denote Hopf and saddle-node bifurcations. The length of the cycle uniqueness interval is 173, i.e., $85 < Km_{LH} < 258$.



Figure 7: Bifurcation diagram with respect to Km_{LH} when $\tau = 1.5$ and $c_2 = 0.03$. HB and SN denote Hopf and saddle-node bifurcations. The left hysteresis curve has disappeared and the length of the cycle uniqueness interval is 271, i.e., $40 < Km_{LH} < 311$.

The broad expansion of the cycle uniqueness interval for c_2 less than 0.05 is due to two different unfoldings of degenerate Hopf bifurcations which occur for c_2 near 0.05. Each Hopf bifurcation is degenerate because the real part of the eigenvalue pair crossing the imaginary axis has a zero derivative with respect to the parameter at crossing. One of these degeneracy occurs when two Hopf points coalesce at $c_2 = 0.05147$ and $Km_{LH} = 69.8458$. At $c_2 = 0.05$, the left side of Figure 6 displays a branch of stable cycles lying just above a branch of unstable equilibria. Figure 8(a) blows these curves up at $c_2 = 0.0511$. They touch when $c_2 = 0.05147$ producing a degenerate Hopf point. Then as c_2 increases, the degenerate Hopf point separates into two nondegenerate, supercritical Hopf points with stable equilibria in between them pictured at $c_2 = 0.0516$ in Figure 8(b). As discussed in Golubitsky and Schaeffer [8], p. 375, the unfolding of this bifurcation may be described roughly by the equation

$$-x^{3} + (Km_{LH} - 69.8458)^{2} x + (0.05147 - c_{2}) x = 0$$
 (DegHB1)

where x represents the state variable LH and the line $\{x = 0\}$ represents the curve of equilibria. As c_2 continues to increase above 0.052, the two Hopf points on the left in Figure 9(a) coalesce in a second degenerate Hopf point at $c_2 = 0.05209$ and $Km_{LH} = 61.0174$ and that Hopf point disappears for $c_2 > 0.05209$. The unfolding of this bifurcation may be represented by the equation (see [8])

$$x^{3} + (Km_{LH} - 61.0174)^{2} x + (c_{2} - 0.05209) x = 0.$$
 (DegHB2)

As c_2 increases above 0.05209, the saddle-nodes which determine the cycle uniqueness interval move closer together causing the interval to shrink and move to the left, see Table 2. As c_2 increases from 0.04 to 0.045, the left hysteresis curve forms because of the appearance of a kink and two saddle-nodes along the left edge of the large loop of periodic solutions (Figure 10). This kink and the two unfoldings (DegHB1 and DegHB2) are embedded in the continuous display of bifurcation diagrams as c_2 increases from 0.04 to 0.055, see Figure 10.



Figure 8: (a): Blow up for $c_2 = 0.051$ (b): Blow up for $c_2 = 0.0516$ s indicates a stable and u, an unstable cycle or equilibrium. HB denotes a Hopf point.



Figure 9: (a): $c_2 = 0.052$ (b): Blow up of "two HB" for $c_2 = 0.052$ s indicates stable and u indicates unstable. HB denotes a Hopf point.



Figure 10: Bifurcation diagrams with respect to Km_{LH} for $\tau = 1.5$ as c_2 increases from 0.04 to 0.055 by increments of 0.5. A kink appears in the large loop at the left and bends to touch the curve of equilibria causing a degenerate Hopf bifurcation. s indicates stable and u indicates unstable. [CLICK HERE] for an animated display of bifurcation diagrams as c_2 increases from 0.04 to 0.055 in increments of 0.001.

4. Comparing Bifurcation Diagrams as Delay τ Varies.

For all c_2 values in Table 2, the delay model has a larger cycle uniqueness interval than the no-delay model. As discussed previously, when $Km_{LH} = 200$ and $c_2 = 0.07$ the delay in the effect of inhibin on FSH results in more vigorous growth of ovarian stages, a longer cycle and higher hormone peaks. Numerical simulations indicate that this is also true after reasonable variations in both model parameters, Km_{LH} and c_2 . It is conceivable that the more robust ovarian development of the delay model permits a broader range of half-saturation constants Km_{LH} for the successful surge response of LH to E_2 priming and, hence, a larger cycle uniqueness interval.

First we fix $c_2 = 0.07$, which is the parameter value fitting the data best (see Table 1). Then we draw bifurcation diagrams with respect to Km_{LH} to study how the cycle uniqueness interval opens up as the delay τ increase from 0 to 1.5. Figure 11 illustrates these diagrams for τ values increasing from $\tau = 0$ to $\tau = 1.5$ by increments of 0.5. As τ increases the Hopf points (HB) along the curve of equilibria spread apart as do the saddle-nodes (SN), which determine the cycle uniqueness interval. The qualitative features of these diagrams are similar. In particular, there are hysteresis curves on both the left and right edges of a large loop of periodic solutions. The hysteresis curves give rise to two regions of periodic bistability.

For other values of c_2 , these two hysteresis curves do not persist for all values of τ . For instance, if $c_2 = 0.04$ then the hysteresis curve on the left disappears as τ increases. The cycle uniqueness interval enlarges from 50 when $\tau = 0$ to 226 when $\tau = 1.5$. The primary reason for this drastic increase is a sequence of bifurcations that occur as τ increases from 0.7 to 1.2. A degenerate Hopf bifurcation similar to that described by (DegHB2) occurs at $\tau = 0.73$ resulting in a bump of stable cycles to the left of the large loop of periodic solutions as pictured in Figure 12. This Hopf bump of stable solutions is just below the branch of unstable cycles in the left hysteresis curve and, as τ increases, this bump grows and touches the curve of cycles above producing a transcritical bifurcation of periodic solutions in the parameter Km_{LH} when $\tau = 1.06$. The unfolding of this transcritical bifurcation is analogous to that discussed in [23] except here the second parameter is τ instead of c_2 . For τ values just above 1.06 the bump of stable cycles appears on the other side of the large loop of cycles (see Figure 12) and disappears via the following sequence of bifurcations. At $\tau = 1.11$ a degenerate Hopf bifurcation like (DegHB1) causes the Hopf bump to separate from the curve of equilibria producing a small closed loop of periodic solutions (Figure 13). Then this loop shrinks and disappears because the two saddle-nodes at each end of the loop coalesce and annihilate one another at $\tau = 1.173$. The unfolding of these bifurcations as τ increases from 0.7 to 1.2 is animated in Figure 12.



Figure 11: Bifurcation diagrams with respect to Km_{LH} for $c_2 = 0.07$ as τ increases from $\tau = 0$ to $\tau = 1.5$ by increments of 0.5. The cycle uniqueness interval enlarges from 114 to 154. HB and SN denote Hopf and saddle-node bifurcations. s indicates a stable and u, an unstable cycle or equilibrium.



Figure 12: Bifurcation diagrams with respect to Km_{LH} for $c_2 = 0.04$ as τ increases from $\tau = 0.9$ to $\tau = 1.2$ by increments of 0.1. HB and SN denote Hopf and saddle-node bifurcations. s indicates a stable and u, an unstable cycle or equilibrium. [CLICK HERE] for an animated display of transcritical and degenerate Hopf bifurcations as τ increasing from 0.7 to 1.2 in increments of 0.02.



Figure 13: Bifurcation diagram with respect to Km_{LH} when $\tau = 1.14$ and $c_2 = 0.04$. HB and SN denote Hopf and saddle-node bifurcations. The small closed loop of cycles has saddle-nodes at each end.

5. Summary and Conclusion.

The half-satuation parameter Km_{LH} in the Hill function in (S1) indicates the level of E_2 sufficient for significant LH synthesis. We study bifurcation diagrams where maximum LH along a periodic or equilibrium solution is graphed against Km_{LH} . We observe an interval of Km_{LH} values for which the model admits a unique stable periodic solution and this solution represents an ovulatory cycle. A large cycle uniqueness interval signifies a wide range of follicular E_2 levels which promote a LH surge sufficient for ovulation. This cycle uniqueness interval is usually determined by two saddle-nodes bifurcations which lie on hysteresis curves at the left and right sides of the bifurcation diagram.

The parameter τ is the time delay for the inhibition of FSH synthesis caused by inhibin. We explain why a delay up to 1.5 days (the value of τ fitting the data best) permits increased ovarian development and a larger interval of Km_{LH} values which result in a unique cycle. The ovarian growth parameter c_2 promotes mass transfer between the first two stages of ovarian development and is indicative of healthy follicular growth. For various values of c_2 , we illustrate how the cycle uniqueness interval grows as τ increases due to the occurrences of transcritical and degenerate Hopf bifurcations, e.g., see Figure 12. Also, for delay τ near 1.5 days, the cycle uniqueness interval increases as c_2 decreases because of additional growth of the first follicular stage as discussed in Section 3.

The transcritical bifurcation is a prominent feature of the left side of the bifurcation diagrams for smaller values of τ . When this bifurcation is present in the diagram, the cycle uniqueness interval has reduced length, e.g., only 125 for the first frame of Figure 14. The transcritical bifurcation persists for the parameter pairings in Figure 14 until bifurcation point coalesces with a saddle-node for $\tau \approx 1.45$. As τ increases, the cycle uniqueness interval grows although c_2 is also increasing (see the animation for Figure 14). Hence, a larger delay in the effect of inhibin may compensate for an apparent reduction in growth of the first follicular stage of a cycle. In fact, an increase in *FSH* inhibition during the luteal phase of the previous cycle due to the delay in inhibin results in greater early follicular development during the next cycle (see the middle cycle in Figure 3). Also, as the delay increases to 1.5 days, the model describes the data in the biological literature [27] better.



Figure 14: Bifurcation diagrams with respect to Km_{LH} as τ increases from 0.048 to 1.5 and c_2 increases from 0.031 to 0.044. HB, SN and T denote Hopf, saddle-node and transcritical bifurcations. s indicates a stable and u, an unstable cycle or equilibrium. [CLICK HERE] for an animated display of a sequence of transcritical bifurcations revealing the creation of a loop and the broadening of the cycle uniqueness interval as τ and c_2 increase.

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