

## §10. Electrostatic Potential due to Induced Charge of Spherical Dust Particle on Plasma-Facing Wall in Non-uniform Electric Field

Tomita, Y.,  
Smirnov, R. (UCSD),  
Zhu, S. (ASIPP)

A spherical conducting dust particle with a radius  $R_d$  is attached on an infinitely extended conducting plane wall. The local electrostatic potential is composed of the sum of the external one  $\phi_{ex}$  and the one  $\phi_{in}$  due to induced charges on the conducting dust particle.

$$\phi(r, z) = \phi_{ex}(z) + \phi_{in}(r, z) \quad (1)$$

where  $(r, z)$  is the conventional cylindrical coordinate, where the origin is located at the contacting point of the spherical dust to the plane wall (Fig.1).

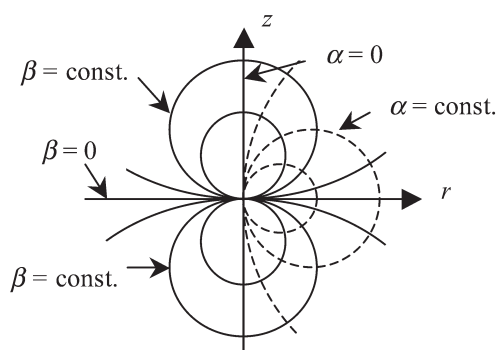


Fig.1 The cylindrical coordinate  $(r, z)$  and the bipolar coordinate  $(\alpha, \beta)$ . The infinitely extended conducting plane wall and the surface of the spherical dust are indicated by  $\beta = 0$  and  $1/2$ , respectively.

The potential  $\phi_{in}$  due to the induced charge satisfies the Laplace equation:

$$\frac{\partial}{\partial \alpha} \left( \frac{\alpha}{\alpha^2 + \beta^2} \frac{\partial \phi_{in}}{\partial \alpha} \right) + \frac{\partial}{\partial \beta} \left( \frac{\alpha}{\alpha^2 + \beta^2} \frac{\partial \phi_{in}}{\partial \beta} \right) = 0 \quad (2)$$

in the bipolar coordinate  $(\alpha, \beta)$  [5]. The relation between the cylindrical coordinate  $(r, z)$  and the bipolar coordinate  $(\alpha, \beta)$  is

$$z + ir = \frac{iR_d}{\alpha + i\beta}, \quad (3)$$

where  $i$  denotes the imaginary unit. The general solution of the Laplace equation (2), which satisfies the condition  $\phi_{in} = 0$  at  $\beta = 0$ , is given by

$$\phi_{in}(\alpha, \beta) = \sqrt{\alpha^2 + \beta^2} \sum_{n=1}^{\infty} c_n I_{\lambda n}(\alpha, \beta), \quad (4)$$

and

$$I_{\lambda n}(\alpha, \beta) \equiv \int_0^{\infty} d\lambda \lambda^n e^{-\beta\lambda} \frac{\sinh(\beta\lambda)}{\sinh(\beta_0\lambda)} J_0(\lambda\alpha) \quad (5)$$

where  $J_0$  is the first kind Bessel function of the 0-th order and the coefficients  $c_n$ 's are determined by the boundary condition at the surface of the spherical dust. In this study

the externally applied potential is approximated to be non-uniform toward the normal direction to the plane wall, i.e.  $z$ , which is expressed in  $(\alpha, \beta)$ :

$$\phi_{ex}(\alpha, \beta) = \sum_{k=0}^{k_{\max}} h_k z^k = \sum_{k=0}^{k_{\max}} h_k \left( \frac{R_d \beta}{\alpha^2 + \beta^2} \right)^k \quad (6)$$

where the coefficient  $h_0$  corresponds to the biased wall potential  $\phi_w$ . The total electrostatic potential  $\phi(\alpha, \beta)$  is obtained as a combination of the externally applied potential  $\phi_{ex}$  and the induced potential  $\phi_{in}$ . In case of  $k_{\max} = 4$ ,

$$\begin{aligned} \phi(\alpha, \beta) &= \phi_{ex}(\alpha, \beta) + \phi_{in}(\alpha, \beta) \\ &= \phi_w + \sqrt{\alpha^2 + \beta^2} \left[ (h_1 R_d + \frac{h_2 R_d^2}{3\beta} + \frac{h_3 R_d^3}{5\beta^2} \right. \\ &\quad \left. + \frac{h_4 R_d^4}{7\beta^3} \right] I_1(\alpha, \beta) + \left( \frac{h_2 R_d^2}{3} + \frac{h_3 R_d^3}{5\beta} + \frac{h_4 R_d^4}{7\beta^2} \right) I_2(\alpha, \beta) \\ &\quad + \left( \frac{h_3 R_d^3}{15} + \frac{2h_4 R_d^4}{35\beta} \right) I_3(\alpha, \beta) + \frac{h_4 R_d^4}{105} I_4(\alpha, \beta) \\ &\quad + \sqrt{\alpha^2 + \beta^2} \sum_{k=1}^{\infty} c_k I_{\lambda k}(\alpha, \beta) \end{aligned} \quad (7)$$

where the first and second terms of the RHS of eq.(7) are the external potential and the last term corresponds the induced potential. From the boundary condition at the surface of the spherical dust  $\beta = \beta_0 = 1/2$ , the coefficients  $c_k$ 's in Eq.(4) are expressed by the known quantities. Finally we obtain the total local potential consisting of the external and induced ones for the case  $k_{\max} = 4$ :

$$\begin{aligned} \phi(\alpha, \beta) &= \phi_{ex}(\alpha, \beta) + \phi_{in}(\alpha, \beta) \\ &= \phi_w + \sqrt{\alpha^2 + \beta^2} \left\{ h_1 R_d I_{s1} + \right. \\ &\quad \left. + \frac{h_2 R_d^2}{3} \left( \frac{\beta_0 - \beta}{\beta_0 \beta} I_1 + \frac{I_{s1}}{\beta_0} + I_{s2} \right) \right. \\ &\quad \left. + \frac{h_3 R_d^3}{5} \left( \frac{\beta_0^2 - \beta^2}{\beta_0^2 \beta^2} I_1 + \frac{I_{s1}}{\beta_0^2} + \frac{\beta_0 - \beta}{\beta_0 \beta} I_2 + \frac{I_{s2}}{\beta_0} + \frac{I_{s3}}{3} \right) \right. \\ &\quad \left. + \frac{2I_{s3}}{5\beta_0} + \frac{I_{s4}}{15} \right\} \end{aligned}$$

Here the function  $I_{sn}$  is defines as

$$I_{sn}(\alpha, \beta) \equiv \int_0^{\infty} d\lambda \lambda^n \frac{\sinh[(\beta_0 - \beta)\lambda]}{\sinh(\beta_0\lambda)} J_0(\lambda\alpha).$$

These expressions can be useful to obtain the induced charge on the dust surface.