

## §1. Synthesis of High Performance Nb<sub>3</sub>Sn Layers through the New Diffusion Process

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In the present JR (Jelly Roll) process Sn-based alloy sheet, 80  $\mu$  m in thickness, is laminated with Nb sheet and wound into a JR composite. The composite is encased in a Nb-based alloy sheath, and then fabricated into a wire. A swift fabrication is possible, and the leak out of Sn is prevented by the holding between Nb sheets at the extrusion and heat treatment. The volume fraction of residual bronze in the JR wire is much smaller than in the bronze and internal Sn processed wires resulting in the increase of non-Cu  $J_c$  and the reduction of strain effect.

At the initial stage of the heat treatment at 750°C Nb<sub>6</sub>Sn<sub>5</sub> layers are formed in the JR wire by the reaction between Nb and Sn sheets. The Nb<sub>6</sub>Sn<sub>5</sub> layers show a distinct columnar microstructure with an aspect ratio of  $\sim$ 10 owing to the solid / liquid reaction. Then Nb<sub>3</sub>Sn layers are formed by the solid state reaction between residual Nb and Nb<sub>6</sub>Sn<sub>5</sub>. The Nb<sub>3</sub>Sn layers show a homogeneous and almost equiaxed microstructure.

The JR processed Nb<sub>3</sub>Sn wires show an offset  $T_c$  of 18.1K with a transition width of less than 0.1K, which is appreciably higher and sharper transition than other Nb<sub>3</sub>Sn wires. Fig. 1 illustrates the relation between offset  $T_c$  and applied magnetic field of two JR wires. The  $B_{c2}$  increases almost linearly with decreasing temperature, which reaches  $\sim$ 13T at 12K. The result indicates that the JR wires are promising for refrigerator-cooled superconducting magnets capable of generating 10T at 12K. The offset  $B_{c2}$  (4.2K) of the JR wires is  $\sim$ 26.5T.

Fig. 2 is non-Cu  $J_c$  versus magnetic field curves of Sn-Ta based sheet JR wires with different diameters and Nb sheet thickness. The non-Cu  $J_c$  increases with reducing the wire diameter, e.g., 120A/mm<sup>2</sup> for 1.4mm  $\phi$  wire and 180A/mm<sup>2</sup> for 1.0mm  $\phi$  wire at 4.2K and 22T. The thinner wire diameter increases the areal fraction of Nb<sub>3</sub>Sn layer as shown in Fig. 2, which may be the main origin of the non-Cu  $J_c$  improvement.

The thickness of the Nb sheet produces a drastic change in both the structure and the performance of JR wires. In the wire starting from 100  $\mu$  m-thick Nb sheet, the Nb sheet is completely consumed and some Nb<sub>6</sub>Sn<sub>5</sub> phase is

present in the JR part after the heat treatment. The excess Sn diffuses into the Nb sheath and the central Nb core to form thick Nb<sub>3</sub>Sn layers. In the wire starting from 160  $\mu$  m-thick Nb sheet, the JR part is composed of equally spaced spiral Nb<sub>3</sub>Sn layers. The configuration of the JR part is much improved compared to the previous wires starting from 100  $\mu$  m-thick Nb sheet.

The non-Cu  $J_c$  of the 160  $\mu$  m-thick Nb sheet wire increases much more rapidly with decreasing field below 24T than the 100  $\mu$  m-thick Nb sheet wire as shown in Fig. 2. The non-Cu  $J_c$  at 4.2K and 21T of the 1.4mm  $\phi$  wires is  $\sim$ 160A/mm<sup>2</sup> and  $\sim$ 290A/mm<sup>2</sup> for 100  $\mu$  m-thick and 160  $\mu$  m-thick Nb sheet wires, respectively. The increase of starting Nb sheet thickness is found to yield much better JR configuration and larger non-Cu  $J_c$  in the wire.

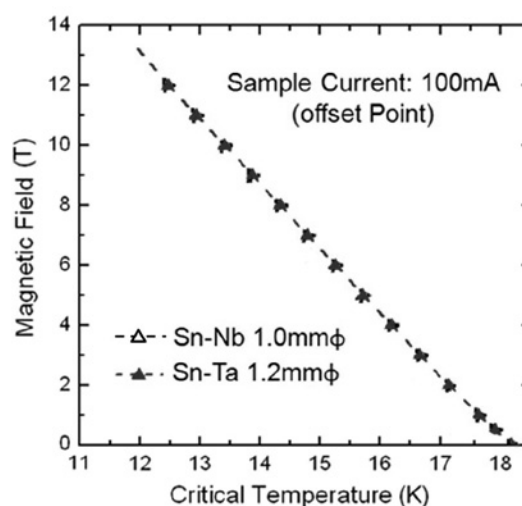


Fig.1 Relation between magnetic field and offset  $T_c$  for Sn-Ta and Sn-Nb based sheet JR wires.

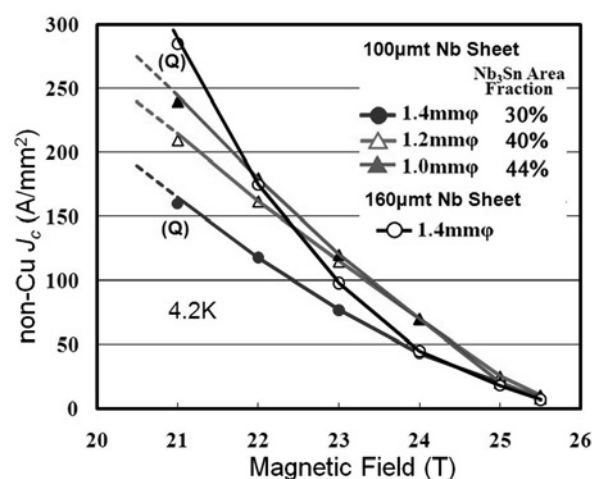


Fig.2 Non-Cu  $J_c$  versus magnetic field curves of wires with different diameter and Nb sheet thickness. The areal fraction of Nb<sub>3</sub>Sn layers in the 100  $\mu$  m-thick Nb sheet wires is inserted in the figure.