

## §5. Orbit Effects on the ECRH Deposition Profile in Stellarator-devices<sup>[1]</sup>

Murakami, S., Nakajima, N., Okamoto, M.  
 U. Gasparino, , Maaßberg, H., Romé, M.  
 (Max-Planck-Institut für Plasmaphysik, Germany)  
 Marushchenko. N. (Institute of Plasma Physics,  
 NSC-KhPTI,Ukraine)

The particles trapped in the helical ripple tend to drift away from the starting magnetic surface in stellarators. Therefore, at low collisionalities, the ECRH-heated electrons can drift radially before being collisionally thermalized. And the absorbed power deposition profile will generally be broader than the peaked "birth profile" usually predicted by ray-tracing. Recent measurements suggesting the existence of a "broad component" in the ECRH deposition profile at W7-AS[2] have put a considerable interest in a quantitative analysis of this r. f.-driven transport.

In order to study the effects of radial drift motion of suprathermal electrons, a Monte Carlo simulation code is developed based on a technique similar to the adjoint equation for dynamic linearized problems. The linearized drift kinetic equation for the deviation from the Maxwellian background  $f_1(\underline{x}, \underline{v})$ ,

$$\vec{v} \cdot \nabla f_1 + \vec{a} \cdot \nabla_{\underline{v}} f_1 = C(f_1) + S_{ql}^0, \quad (1)$$

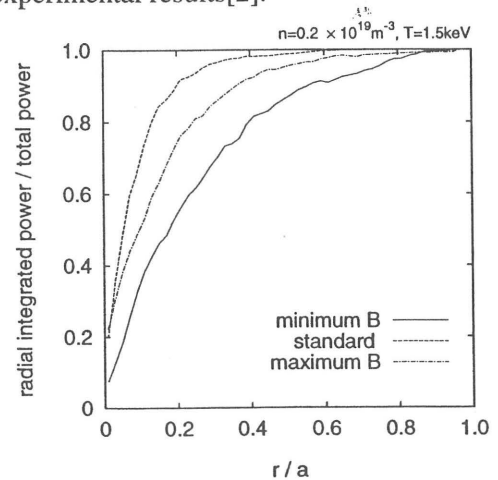
is evaluated as,

$$f_1(\underline{x}, \underline{v}) = \int_0^\infty dt \int d\underline{x}' \int d\underline{v}' S_{ql}^0(\underline{x}', \underline{v}') g(\underline{x}, \underline{v}, t | \underline{x}', \underline{v}'), \quad (2)$$

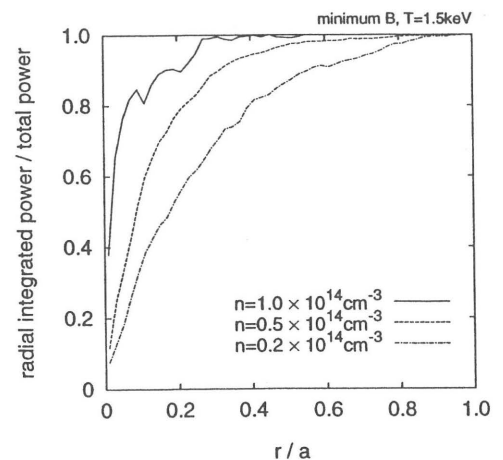
where  $C(f_1)$  is the linear Coulomb collision operator and  $S_{ql}^0$  is the wave induced flux in velocity space (quasi-linear diffusion term) which is assumed to be a given function. The time dependent Green function  $g(\underline{x}, \underline{v}, t | \underline{x}', \underline{v}')$  is the solution of the drift kinetic equation,  $\partial g / \partial t + \vec{v} \cdot \nabla g + \vec{a} \cdot \nabla_{\underline{v}} g = C(g)$ , with initial condition  $g(\underline{x}, \underline{v}, t = 0 | \underline{x}', \underline{v}') = \delta(\underline{x} - \underline{x}') \delta(\underline{v} - \underline{v}')$ . The function  $g$  is obtained using the Monte Carlo simulation in which the complex magnetic field configuration, the finite- $\beta$  effect, and the radial electric field can be included.

Figures 1 and 2 shows the normalized power deposition profiles with changing the configurations and

densities (X-mode, 2nd-harmonic,  $T = 1.5\text{keV}$ , and  $Z_{eff} = 2.0$ ). The quasi-linear diffusion term  $S_{ql}^0$  is estimated through an analytical model of ECRH heating. We can see a considerable broadening of the deposition profile due to the large orbit of suprathermal electrons, especially for "minimum B" configuration and for lower densities. These results are consistent with the experimental results[2].



**Fig. 1:** Normalized power deposition profiles for different configurations. [ $n_0 = 2.0 \times 10^{19} \text{m}^{-3}$ ,  $T = 1.5\text{keV}$ ].



**Fig. 2:** Same for different densities in the "minimum B" configuration.

## References

- 1) Murakami, S., et al, *Proc. 1996 Int. Conf. Plasma Phys.* (Nagoya, JSPSNFR) Vol. 1 (1997) p 1014.
- 2) Romé, M., et al, *Plasma Phys. Contr. Fusion* **39** (1997) 117.