

## §17. Comparison of Transfer Efficiency of High Energy Ion at Inward-Shifted and Standard Magnetic Axis

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In the minority ICRF heating the most of the RF power is absorbed by minority ions<sup>1)</sup>. The bulk plasma is heated by high energy ions via a slowing down process due to electrons in the following power balance equation about the tail and bulk plasmas;

$$\frac{dW_{tail}}{dt} = \eta_h P_A - \frac{W_{tail}}{\tau_E^{tail}},$$

$$\frac{dW_{bulk}}{dt} = \frac{W_{tail}}{\tau_s/2} - \frac{W_{bulk}}{\tau_E}.$$

The heating power source for the bulk plasma is a term of  $W_{tail}/\tau_s/2$ . A transfer efficiency  $\eta_{trms}$  is defined as what fraction of absorbed power is transferred to the bulk plasma using a following equation;

$$P_{abs} = \frac{W_{tail}}{\tau_E^{tail}}, \quad P_{trms} = \frac{W_{tail}}{\tau_s/2},$$

$$\eta_{trms} = \frac{P_{trms}}{P_{abs}} = \frac{\tau_E^{tail}}{\tau_s/2}, \quad \frac{1}{\tau_E^{tail}} = \frac{1}{\tau_s/2} + \frac{1}{\tau_E^{loss}}.$$

$\eta_{trms}$  is expressed using a ratio of  $\tau_E^{tail}/\tau_s/2$ . When the energy loss time  $\tau_E^{loss}$  is much longer than the electron slowing down time, the transfer efficiency becomes 100%. It is also expressed as the ratio of  $T_{tail}$  to  $T_{eff}$ :

$$T_{tail} = \frac{\tau_E^{tail}}{\tau_s/2} T_{eff},$$

$$\eta_{trms} = \frac{T_{tail}}{T_{eff}}.$$

In the Monte Carlo simulation  $\eta_{trms}$  is scaled as shown in the following equation<sup>2)</sup>:

$$\eta_{trms} = \frac{P_{trms}}{P_{abs}} = \frac{\tau_E^{tail}}{\tau_s/2} = \frac{T_{tail}}{T_{eff}} = \frac{1}{1 + C P_{abs} T_e^2 (n_H/n_e)^{-1} n_e^{-2}}.$$

Here C is a numerical factor to evaluate  $\eta_{trms}$ ; C depends on the magnetic configuration and the magnetic strength. It is determined using the ratio of  $T_{tail}$  to  $T_{eff}$  for 2 cases as described in the previous section and plotted in Fig.1, where the abscissa is  $P_{abs} T_e^2 (n_H/n_e)^{-1} n_e^{-2}$ . C is determined to be  $C_{exp}=0.032$  at  $R_{ax}=3.75m$  and  $0.005$  at  $R_{ax}=3.6m$ , respectively. C is also evaluated in the Monte Carlo simulation to be  $C_{MC}=0.043$  at  $R_{ax}=3.75m$  and  $0.004$  at  $R_{ax}=3.60m$  as shown in Fig.1. These experimental values fairly agree with the Monte Carlo simulation results.

The reduction of the transfer efficiency is caused by a decrease in the energy loss time of high energy ions  $\tau_E^{loss}$ . It is expressed using the transfer efficiency  $\eta_{trms}$  in the following equation;

$$\tau_E^{loss} = \frac{\tau_E^{tail} \tau_s / 2}{\tau_s / 2 - \tau_E^{tail}} = \frac{\eta_{trms}}{1 - \eta_{trms}} \tau_s / 2.$$

It is plotted in Fig.2, whose abscissa is  $n_e P_{abs}^{-1} T_e^{-0.5} (n_H/n_e)^{-1}$ . The energy loss time ranges in 0.1~0.2s in the plasma discharge at  $R_{ax}=3.75m$ ; however it is expected 10 times longer at  $R_{ax}=3.6m$  than that because C was 1/10.

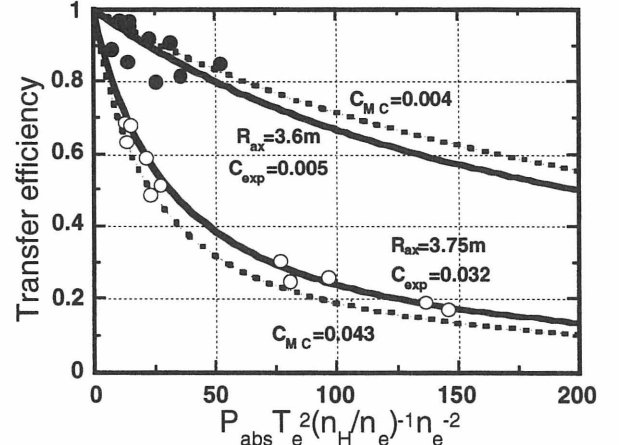


Fig.1 Dependence of transfer efficiency in the cases of  $R_{ax}=3.6m$  and  $3.75m$  with the calculation results from the Monte Carlo simulation.

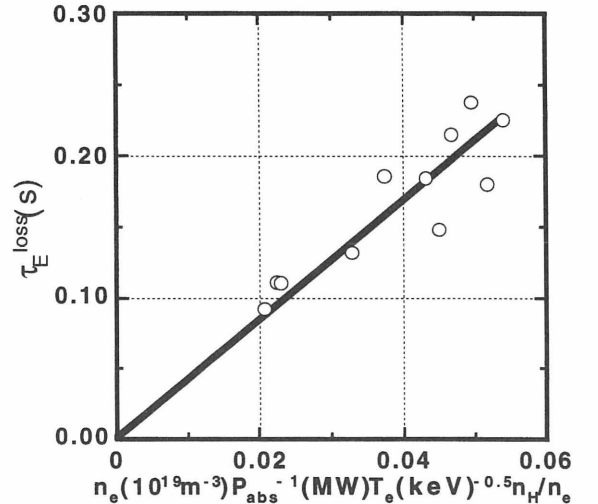


Fig.2 Dependence of energy loss time of high energy ions on  $n_e P_{abs}^{-1} T_e^{-0.5} (n_H/n_e)^{-1}$ .

### References

- 1) K.Saito et. al., Nucl. Fusion **41** (2001) 1021.
- 2) S.Murakami et. al., Nucl. Fusion **39** (1999) 1165.