§51. Role of Energetic Electrons on Non-inductive Current Start-up and Formation of Inboard Poloidal Field Null Configuration in the Spherical Tokamak QUEST

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Non-inductive plasma current startup using ECH has been conducted on several tokamaks and spherical tokamaks (STs). Recently, it has been suggested that  $I_p$  is generated by trapped particles in STs. In QUEST, a new startup scenario has been tested using ECH at a high toroidal magnetic mirror ratio  $M_{OMFC}$  of ~2 (=decay index  $n^*\approx 0.5$ ) with high  $B_z$  and  $B_z/B_t$  of 10% at the fundamental resonance layer  $R_{res1}^{1,2)}$ . The increments of  $B_z$  and  $M_{OMFC}$  are favorable for trapped fast electron confinement. In these experiments, because of high  $\beta_p$  plasma, under high  $M_{OMFC}$  and ECW-driven current, the inboard poloidal null configurations are observed in steady state.

Next, we described the manner in which confined fast electrons contribute to electron pressure and plasma equilibrium through  $\beta_p$  using fast electrons effective temperature  $T_{HX}$  by measuring hard X-ray (HXR) energy spectrum. To study electrons in  $v_z > v_{l/s}$ , the measurement with vertical line of sight is conducted for different  $M_{OMFC}$  values. The relation between  $I_p$  and  $B_z$  is expressed as<sup>3</sup>,

$$B_{z} = \mu_{0} I_{p} (In(8R/a) + l_{i}/2 - 3/2 + \beta_{p})/4pR$$
(1)

where  $l_i$  is the internal inductance. The effects of  $M_{OMFC}$  on  $\beta_{\rm p} = (\beta_{\rm p}^{\rm bulk} + \beta_{\rm p}^{\rm hot})$  is investigated. Relations between  $I_p$  and  $B_z$ are plotted in Fig. 1(a) for different  $M_{OMFC}$  values. The solid lines stand for time traces in the  $B_z$ - $I_p$  space. The symbols correspond to steady state  $I_{ps}$  under constant  $B_{zs}$ . For  $M_{OMFC}$ =1.2,  $I_p$  increased linearly up to 5 kA as  $B_z$  increased to 1.7 mT, but it dropped to <1 kA at  $B_z$ >1.8 mT. On the other hand, at  $B_z$  ramp-up discharge,  $I_p$  reached 15 kA at 4 mT. These cases of  $B_z$  constant and ramp-up show the same proportional constant for  $M_{OMFC}$ =1.2. For  $M_{OMFC}$ =1.8, 2, and 2.7, a drop or saturation of  $I_p$  was not observed. In addition, the proportional constant of the  $I_p$ - $B_z$  relation was maintained. For  $M_{OMFC}=1.2$ , 1.4, and 2, the inverse of proportional constants  $B_z/I_p$  was 0.3, 0.5, and 1 mT/kA, respectively, in steady state  $I_p \approx 15$  kA. This suggests that stronger  $B_z$  is required to attain the same  $I_p$  as  $M_{OMFC}$  is increased up to 2. Fig. 1(b) to (e) shows the reconstructed closed flux surfaces for MOMFC=1.2, 1.4, 1.8, and 2 at  $I_p \approx 15 \pm 1$  kA taken at  $B_z = 4.79$ , 8.47, and 15 mT, respectively. The plasma shown in Fig. 1(b) is noncircular, which is characterized by an elongation  $\kappa$ =1.14 and triangularity  $\delta$ =0.24 with an aspect ratio  $A_p$ (= $R_0$ =<a>=0.69 m/0.45 m) of 1.63, and Shafranov shift  $\Delta$ =0.047. The value of  $\beta_{\rm p}$  is evaluated as 0.64 from Eq. (1) with an assumption of internal inductance  $l_i=1.2$  and the parabolic profile of  $I_p$ . A typical low-aspect-ratio ST was present. Fig. 1(c) shows the



Fig. 1(a)  $I_p$ - $B_z$  relations for  $M_{OMFC}$ =1.2, 1.3, 1.4, 1.8, 2, and 2.7. Poloidal flux contours for (b) $M_{OMFC}$ =1.2, (c) $M_{OMFC}$ =1.4, (d) $M_{OMFC}$ =1.8, and (e) $M_{OMFC}$ =2.



Fig. 2.  $\beta_p$  and  $T_{HX}$  as a function of  $M_{OMFC}$ .

oblate shape with parameters  $\langle a \rangle = 0.2 \text{ m}$ ,  $R_0 = 0.73 \text{ m}$ , A = 2.9,  $\kappa = 0.63$ ,  $\Delta = 0.106$ , and  $\beta_p$  evaluated as 3.7. For Figs. 1(d) and (e), the poloidal field null points appear, which is caused by increase of  $\beta_p$ . Fig. 2 shows  $\beta_p$  and  $T_{\text{HX}}$  at  $I_p \approx 15 \pm 1 \text{ kA}$  as a function of  $M_{OMFC}$ .  $T_{\text{HX}}$  is evaluated in the energy range of  $\langle 200 \text{ keV} \rangle$ , indicating that  $\beta_p$  and  $T_{\text{HX}}$  increased with  $M_{OMFC}$ . The high value of  $T_{\text{HX}}$  is probably due to the better confinement of fast electrons caused by the high  $M_{OMFC}$  and  $B_z$  configurations. Thereafter, the confined fast electrons contribute to the formation of high  $\beta_p$  plasma.

In summary, to study the effects of fast electron confinement on  $\beta_p$  equilibrium, experiments with different confinement regimes of trapped particles were performed. The experiments indicated that stronger  $B_z$  is required for equilibrium at the same  $I_p$  values. Furthermore, the natural poloidal field null configuration, which is caused by high  $\beta_p$ , appeared as  $M_{OMFC}$  increased. The peak position of the radial profile of  $T_{\rm HX}$  agrees with reconstructed magnetic flux surfaces.

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